

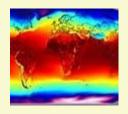
# Spectrally condensed turbulence in two dimensions

Michael Shats<sup>1</sup>, Hua Xia<sup>1</sup>, Gregory Falkovich<sup>2</sup>

<sup>1</sup> The Australian National University, Canberra, Australia

<sup>2</sup> Weizmann Institute of Science, Rehovot, Israel

Acknowledgements: H. Punzmann, D. Byrne



#### **Motivation**

Turbulence often coexists with coherent flow

2D turbulence is capable of generating such flows spectral condensation, crystallization

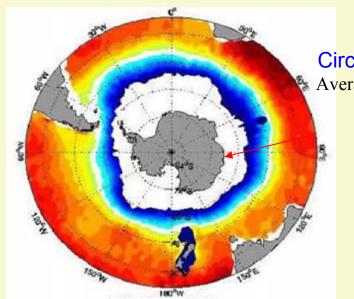
Turbulence-condensate interplay – dynamical steady-state energy transfer from turbulence to flows effects of shear flows on turbulence

#### Practical applications

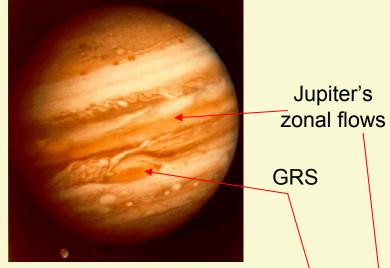
atmospheric and oceanic processes, magnetically confined plasma, etc.

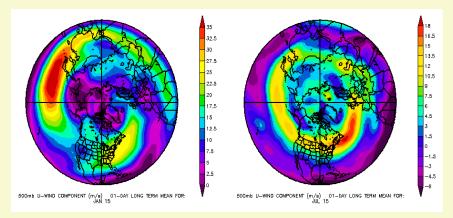


## Large coherent flows coexist with turbulence



Antarctic
Circumpolar Current
Average volume transport
~ 1.5×10<sup>8</sup> m<sup>3</sup>s<sup>-1</sup>



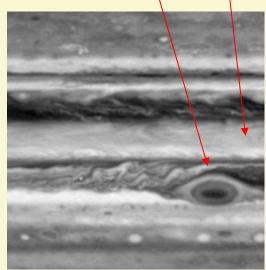


Earth: atmospheric zonal winds

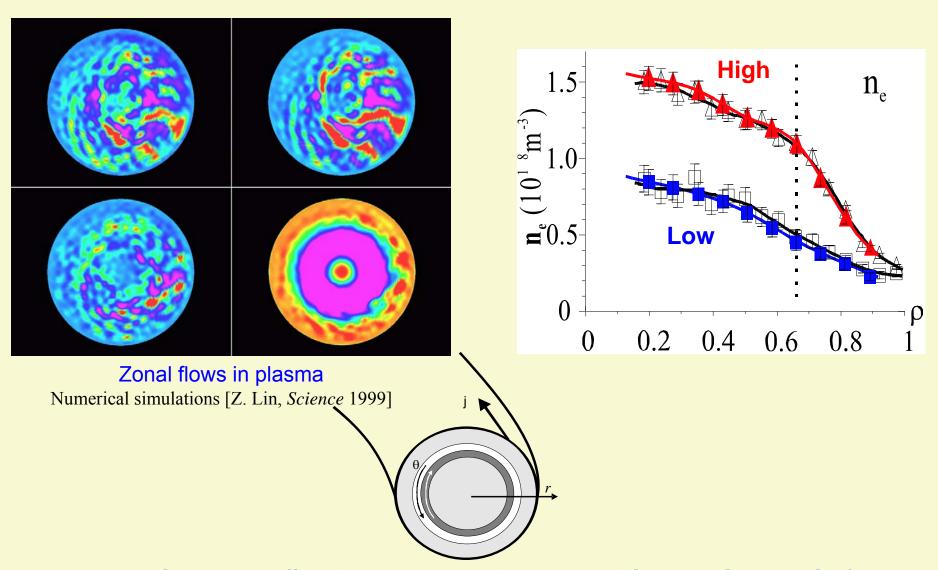
Planetary atmospheres are dominated by turbulent structures (cyclones, zonal winds, etc)

Cassini spacecraft

Courtesy NASA

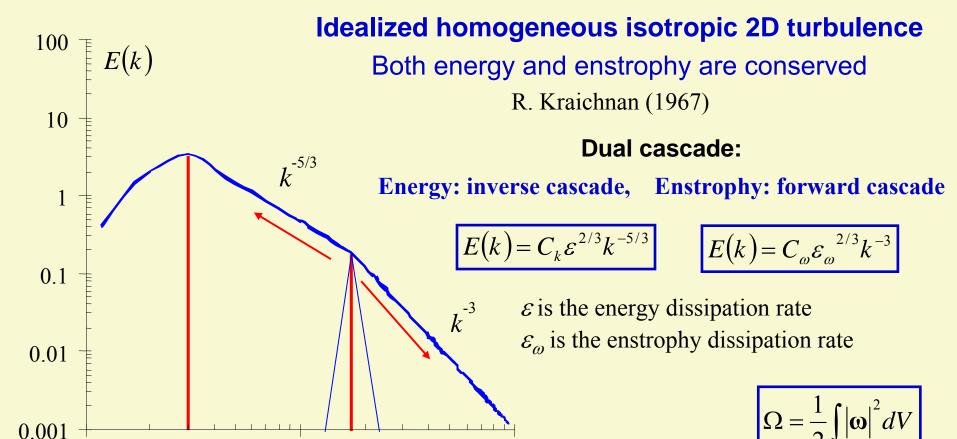


#### Turbulence-driven structures in fusion plasma



In magnetically confined plasma, turbulence-driven anisotropic flows develop, which inhibit radial transport of particles and energy

#### 2D turbulence



Opposite to 3D, energy flows from smaller to larger scales

enstrophy

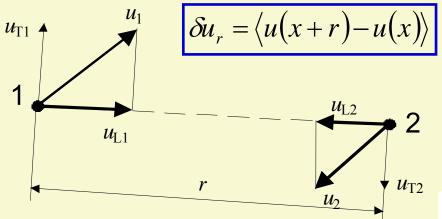
 $k_{\alpha}$ 

10

**Basis for self-organization** 

# Structure functions and Kolmogorov law

Label an 'eddy' by a velocity increment  $\delta u_l$  across a distance r:



Statistical moments of this increment are called *structure functions* of the n<sup>th</sup> order:

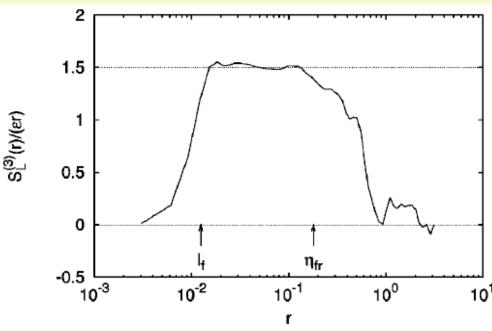
$$S_n(r) = \langle (\delta u_r)^n \rangle = \langle (u(x+r) - u(x))^n \rangle$$

#### Kolmogorov law

relates the third-order longitudinal structure function of turbulence to the mean energy dissipation per unit mass  $\varepsilon$ 

in 2D (e.g. [Lindborg 1999]):

$$S_{3L}(r) = \left\langle \delta V_L^3(r) \right\rangle = \frac{3}{2} \varepsilon r$$



[G. Boffetta, A. Celani, M. Vergassola, 2000]

# Spectral condensation of 2D turbulence

The maximum of the energy spectrum lies in the low-k range, at  $k_{\alpha}$ , in the absence of the energy dissipation at large scales  $k_{\alpha}$  can not be constant in time since it accumulates spectral energy

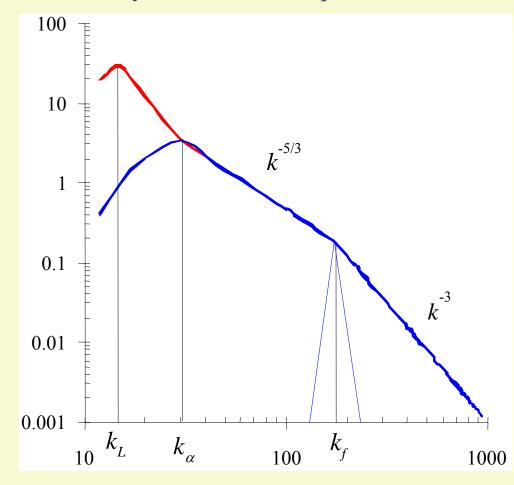
$$k_{\alpha} = f(\varepsilon, t)$$

Dissipation at large scales (bottom damping)  $\alpha$  stabilizes the maximum of the spectrum at the scale

$$k_{\alpha} \approx \left(\alpha^3/\varepsilon\right)^{1/2}$$

Kraichnan, 1967: predicted condensate

System size < dissipation scale



At low dissipation in a bounded system, at  $k_{\alpha} << k_L$  spectral energy accumulates in a box-size coherent structure

#### Spectral condensation of turbulence in thin layers

Experiments: Sommeria (1986), Paret & Tabeling (1998), Shats et al (2005, 2007)





Time evolution to condensed state [Shats *et al* (2005)]

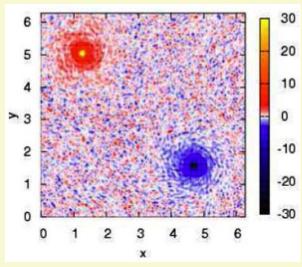
#### Numerical simulations of 2D turbulence:

Hossain (1983), Smith & Yakhot (1993)...

van Heijst, Clercx, Molenaar (2004-2006),

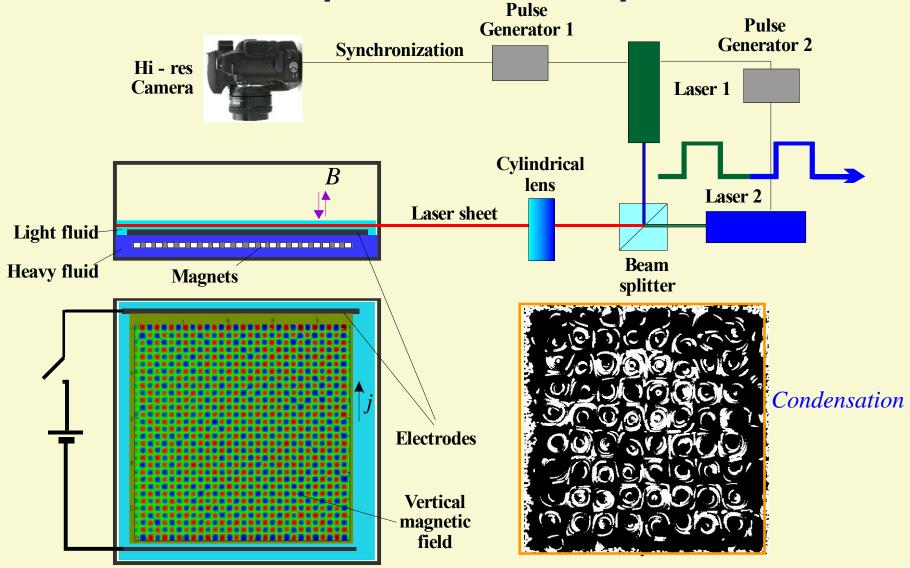
Chertkov et al. (2007)

Periodic boundary condition – dipole No-slip boundary – single vortex



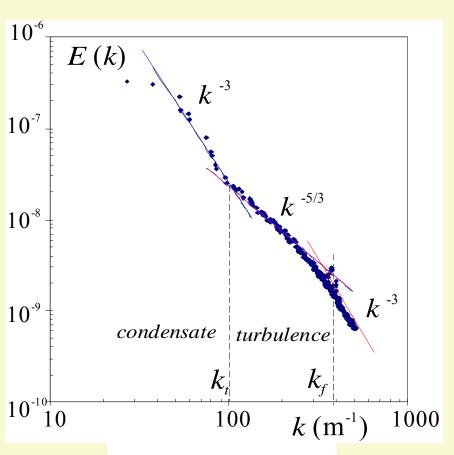
Vorticity of the condensate [M. Chertkov et al. (2007)]

**Experimental setup** 

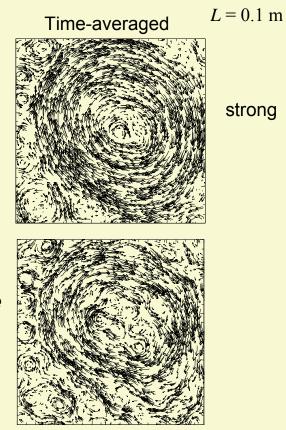


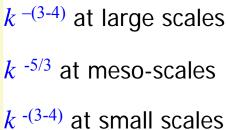
- □ Bottom layer: isolator Fluorinert FC-77 (resist. = 2x10<sup>15</sup> Ohm cm; SG = 1.78)
- □ Top layer: electrolyte NaCl solution (SG = 1.04)

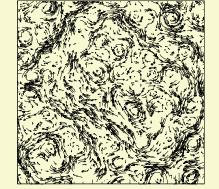
## Condensed turbulence spectrum is robust



-Bottom drag, -boundary size, -forcing affect condensate strength and topology





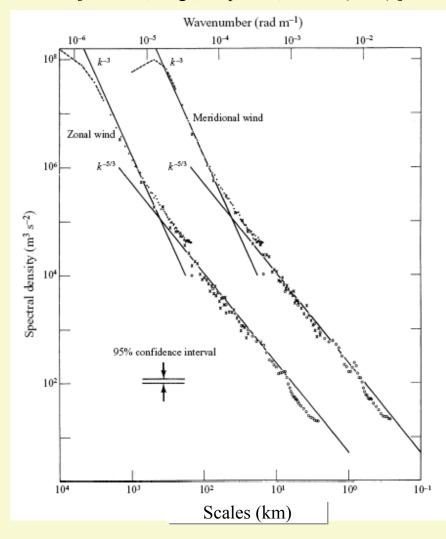


weak

# Nastrom-Gage spectrum of atmospheric winds

#### **Atmospheric spectrum**

[Nastrom, Gage, Jasperson, Nature (1984)]



 $k^{-3}$  and  $k^{-5/3}$  ranges are present but in the reversed order compared to the Kraichnan theory

$$E(k) = C_k \varepsilon^{2/3} k^{-5/3} \quad \text{at } k < k_f$$

$$E(k) = C_\omega \varepsilon_\omega^{2/3} k^{-3} \quad \text{at } k > k_f$$

What is the origin of

 $k^{-3}$  and  $k^{-5/3}$  ranges in atmosphere?

Meso-scale k<sup>-5/3</sup> range can be due to

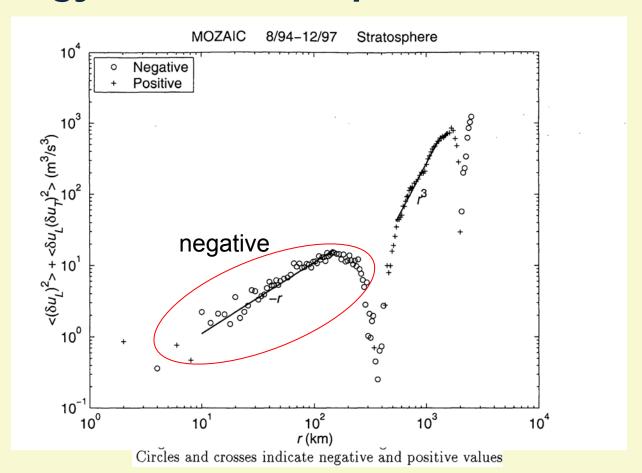
- •3D (downscale) direct energy cascade,
- •2D inverse (upscale) cascade

Large-scale k<sup>-3</sup> range can be due to

- direct enstrophy cascade (large-scale forcing)
- spectral condensation

Kinetic energy spectrum alone cannot resolve the question of the sources

## Energy flux in atmospheric turbulence



Third-order velocity moment gives the energy flux direction

$$S_3(r) = \frac{3}{2}\varepsilon r$$

Negative  $S_3$  at scales up to 500 km interpreted as evidence against inverse energy cascade in the mesoscale range

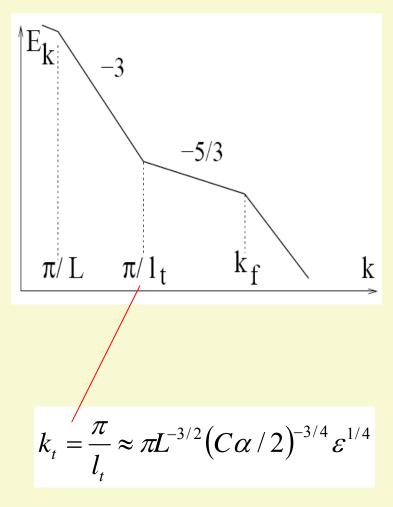
[Cho, Lindborg, J. Geophys. Res. (2001)]

# Need to understand spectral flux in the presence of large coherent flow, which may affect higher moments

Condensate – coherent flow – self-generated by turbulence

In the lab can control strength and spectral extent of condensate (?)

## Model of the spectrum



In the inverse cascade, the turnover time of the eddy of scale l is  $t_1 = l/\sqrt{S_2} \approx l^{2/3}C^{-1/2}\varepsilon^{-1/3}$ 

- 1. Assume that the condensate (vortex) appears when the system size L is such that  $t_L \alpha < 1$
- 2. Characterize the condensate amplitude by its mean velocity V.

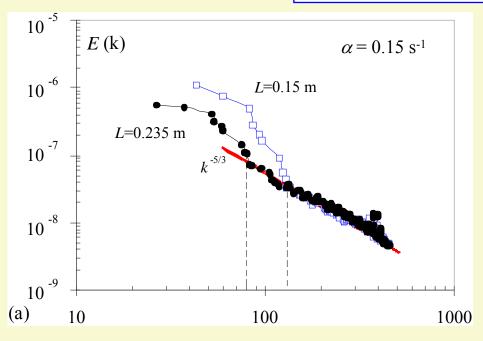
This velocity can be estimated from the energy balance,  $\alpha V^2 \cong 2\varepsilon$  which gives  $V \cong \sqrt{2\varepsilon/\alpha}$ 

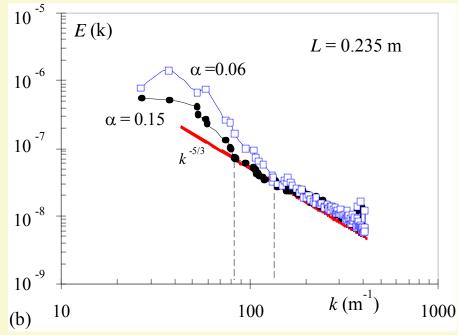
3. We estimate that the condensate related velocity fluctuation on the scale l as Vl/L. Then we expect the knee of the spectrum to be at the scale  $l_t$  defined by

$$Vl_t/L \cong C^{1/2} (\varepsilon l_t)^{1/3}$$
 This gives 
$$l_t \approx L^{3/2} (C\alpha/2)^{3/4} \varepsilon^{-1/4}$$

# Knee of the spectrum shifts with $\alpha$ and L

$$k_{t} = \frac{\pi}{l_{t}} \approx \pi L^{-3/2} (C\alpha/2)^{-3/4} \varepsilon^{1/4}$$





 $k_t$  increases with the decrease in the boundary size L

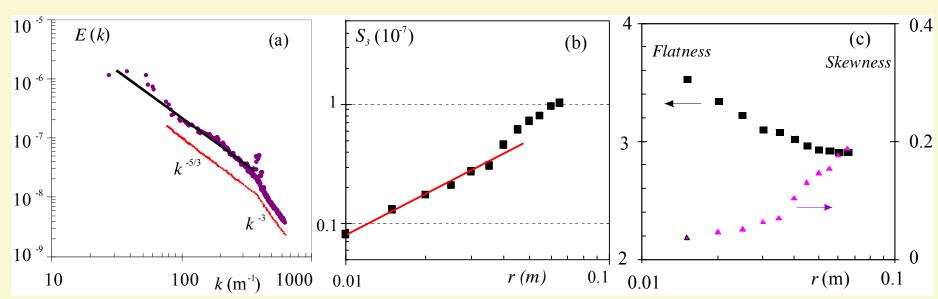
 $k_t$  increases with the decrease in the damping rate  $\alpha$ 

In the lab we can control strength, spectral extent of the condensate

#### Case of weak condensate

$$E(k) = C_k \varepsilon^{2/3} k^{-5/3} \qquad S_3(r) = \left( \left\langle \delta V_L^3 \right\rangle + \left\langle \delta V_L \delta V_T^2 \right\rangle \right) / 2 = \varepsilon r \qquad Sk = S_3 / (S_2)^{3/2}$$

$$F = S_4 / S_2^2$$



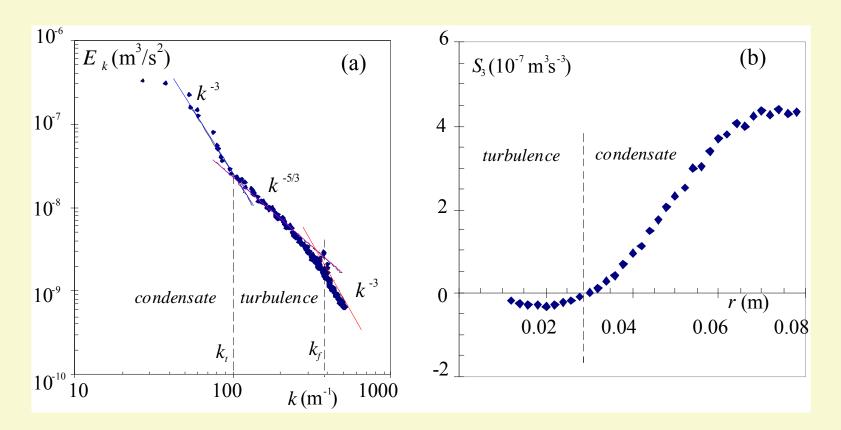
#### Weak condensate case shows small differences with isotropic 2D turbulence

 $\sim k^{-5/3}$  spectrum in the energy range

Kolmogorov law – linear  $S_3$  (r) dependence; Kolmogorov constant  $C \approx 5.6$ 

Skewness and flatness are close to their Gaussian values (Sk = 0, F = 3)

#### Case of stronger condensate



Mean shear flow (condensate)  $\delta V$  changes all velocity moments:

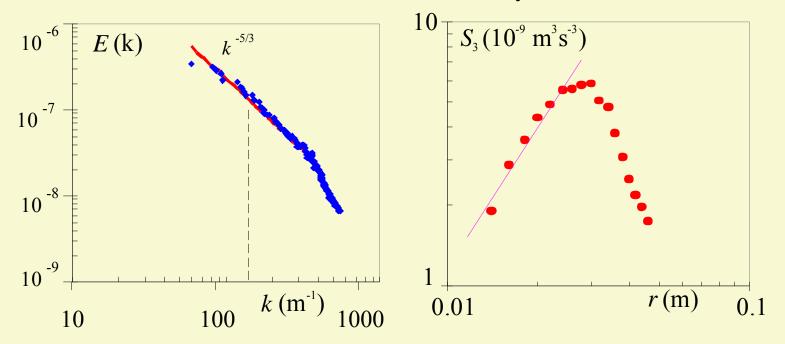
$$\delta V = \delta \overline{V} + \delta \widetilde{V}$$

$$\left\langle \delta V^2 \right\rangle = \left\langle \delta \overline{V}^2 + 2\delta \overline{V} \delta \widetilde{V} + \delta \widetilde{V}^2 \right\rangle$$

$$\left\langle \delta V^3 \right\rangle = \left\langle \delta \overline{V}^3 - 3\delta \overline{V}^2 \delta \widetilde{V} + 3\delta \overline{V} \delta \widetilde{V}^2 - \delta \widetilde{V}^3 \right\rangle$$

## Mean subtraction recovers isotropic turbulence

- 1. Compute time-average velocity field (N=400):  $\overline{V}(x,y) = 1/N \sum_{n=1}^{N} V(x,y,t_n)$
- 2. Subtract  $\overline{V}(x, y)$  from *N*=400 instantaneous velocity fields



Recover  $\sim k^{-5/3}$  spectrum in the energy range

S<sub>3</sub> (r) is positive – recovered inverse energy cascade

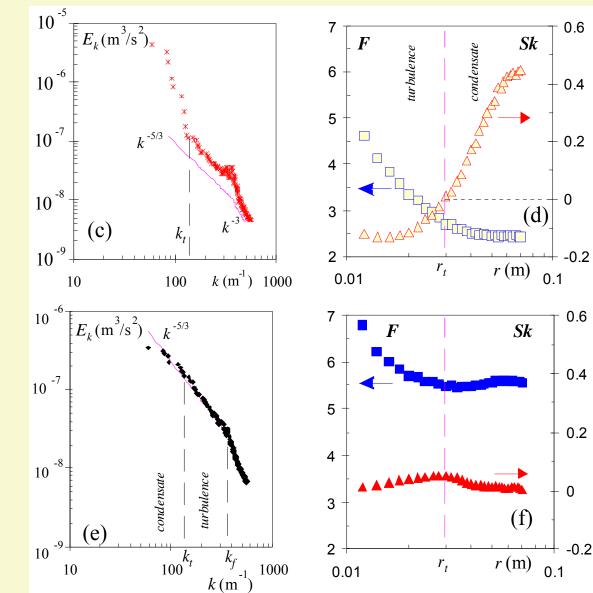
Kolmogorov law – linear  $S_3$  (r) dependence in the "turbulence range";

Kolmogorov constant C ≈ 7

#### Strong condensate: effect of mean subtraction

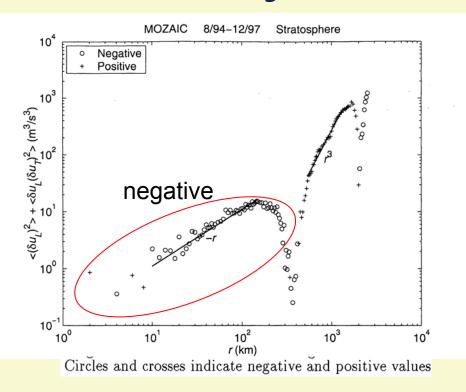
**Before:** 

After:



Normalized moments ~ scale-independent Flatness is higher than in isotropic turbulence

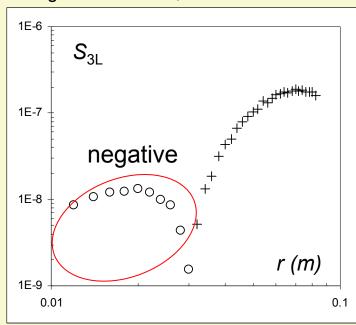
# Similarity with atmospheric turbulence



Mean shear flows present in the atmosphere affect velocity moments, similarly to laboratory experiments

[Cho, Lindborg, J. Geophys. Res. (2001)]

# Laboratory experiment Stronger condensate, no mean subtraction



# Different moments affected in different ranges

#### **Second moment**

Large-scale flow is spatially smooth:  $\delta V \approx sr$ ,  $\langle (\delta V)^2 \rangle \approx s^2 r^2$ 

Small-scale velocity fluctuations in turbulence  $\langle (\delta v)^2 \rangle \approx C(\varepsilon r)^{2/3}$ 

Small-scale fluctuations dominate at scales smaller than  $l < l_t \approx C^{3/4} s^{-3/2} \varepsilon^{1/2}$ 

#### **Third moment**

Large-scale flow:  $\langle \delta V^3 \rangle = \langle \delta V \delta v^2 \rangle \approx srC(\varepsilon r)^{2/3}$ 

Small-scale fluctuations:  $\langle (\delta v)^3 \rangle \approx \varepsilon r$ 

Large-scale flow dominates 3<sup>rd</sup> moment in a range to much smaller scales:

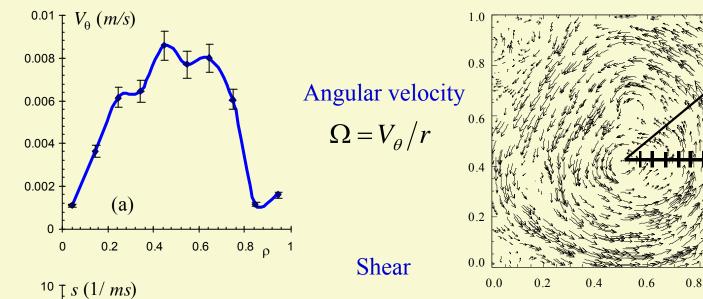
$$l_* \approx C^{-3/2} s^{-3/2} \varepsilon^{1/2}$$
 , since C > 1

Condensate imposes different scales on different moments

# The strongest condensate can suppress turbulence

- (a) Self-generated condensate
- (b) Externally imposed flow

#### Mean flow and its shear



Eddy lifetime

$$au_e pprox rac{l}{\sqrt{S_2(l)}}$$

$$s = \frac{d\Omega}{dx}$$

The condition for the turbulence suppression

$$\omega_s \tau_e > 1$$

Shearing rate

$$\omega_{s} = l \frac{d\Omega}{dr} = l \left[ \frac{1}{r} \frac{dV_{\theta}}{dr} - \frac{V_{\theta}}{r^{2}} \right] = sl$$



[M.G. Shats, H. Xia, H. Punzmann and G. Falkovich (2007)]

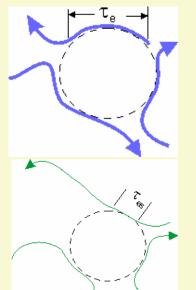
-10

-20

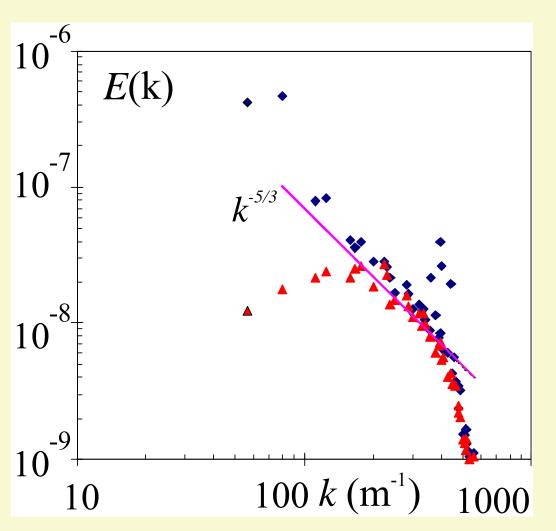
-30

0.2

0.6



# Condensate "shears" large scales first



Shearing acts more efficiently on large scales:

$$s = \omega_s \tau_e \sim l^{5/3}$$

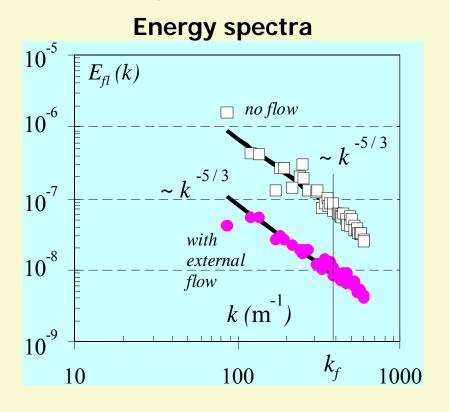
Suppression criterion s > 1 is satisfied for scales

$$l > 0.022 \text{ m}$$

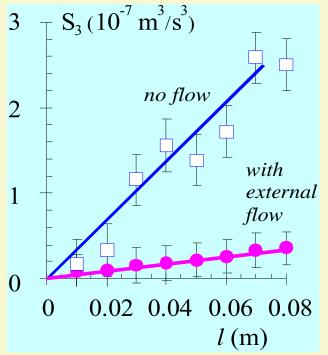
$$k < 145 \text{ m}^{-1} \text{ suppressed}$$

#### Imposed flow reduces turbulence

Externally imposed flow leads to turbulence reduction but retains k<sup>-5/3</sup> scaling



#### 3<sup>rd</sup> order structure function



The energy flux through inertial range  $\varepsilon = -2/3 S_3(l)/l$ 

 $\varepsilon$  is constant for all scales l

 $\varepsilon$  is reduced by a factor of ~10 in the presence of the strong flow

Mean flow reduces energy injected into turbulent cascade  $\varepsilon$ , leads to E(k) drop:

$$E(k) = C_k \varepsilon^{2/3} k^{-5/3}$$

#### Sweeping due to mean flow

Imposed flow affects scales down to the forcing scale

Mean flow sweeps forcing scale vortices relative to magnets

#### **Sweeping**

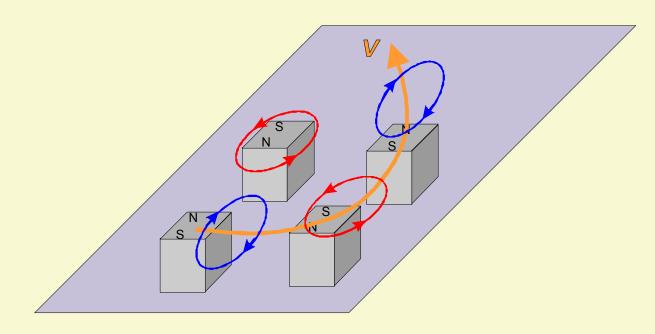
becomes important when

sweeping parameter

$$sw = \omega_{sw}\tau_e = \frac{V_\theta}{\sqrt{S_2}} > 1$$

Eddy life time

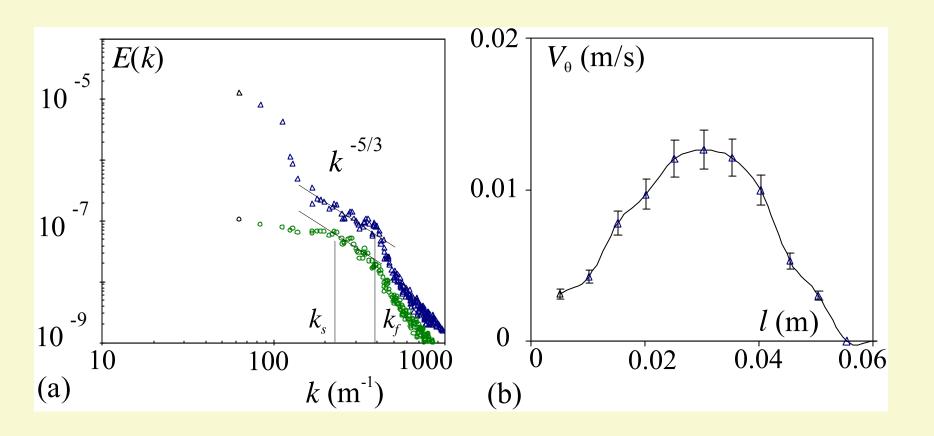
$$\tau_e = l / \sqrt{\left\langle \left| \delta u^2(l) \right| \right\rangle}$$



**Sweeping is more efficient on small scales:** 

$$sw = \omega_{sw}\tau_e = \frac{V_\theta}{S_1} \propto l^{-1/3}$$

# **Shearing and sweeping**



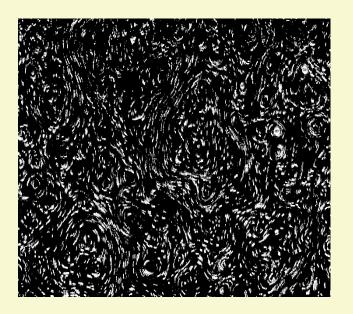
Strongest condensate causes both shearing and sweeping – reduces turbulence that feeds it

## **Summary**

- Spectral condensation leads to the generation of mean flow coherent across the system size
- Spectrum of condensed turbulence shows 3 power laws:
  - $\sim k^{-3}$  at large scales;  $k^{-5/3}$  in the meso-scales,  $k^{-(3-4)}$  at small scales
- Condensate modifies statistical moments of velocity fluctuations
- Different moments are affected in different ranges of scales
- Velocity moments of condensed turbulence similar to those in the atmosphere
- Coherent shear flow suppresses turbulence which generates it via shearing and sweeping

# Flow externally imposed on turbulence

Particle streak photo of turbulent flow



Turbulence with externally imposed mean flow

