

# Tracing the Dark Matter Sheet Through Phase Space

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- First Halos
- Phase Space Sheet
- Outlook

- mostly in collaboration with Oliver Hahn, Raul Angulo, Ralf Kähler, Devon Power, Sam Totorica, Arka Banerjee

Monthly Notices  
of the  
ROYAL ASTRONOMICAL SOCIETY

Mon. Not. R. Astron. Soc. **427**, 61–76 (2012)

doi:10.1111/j.1365-2966.2012.21754.x

**Tracing the dark matter sheet in phase space**

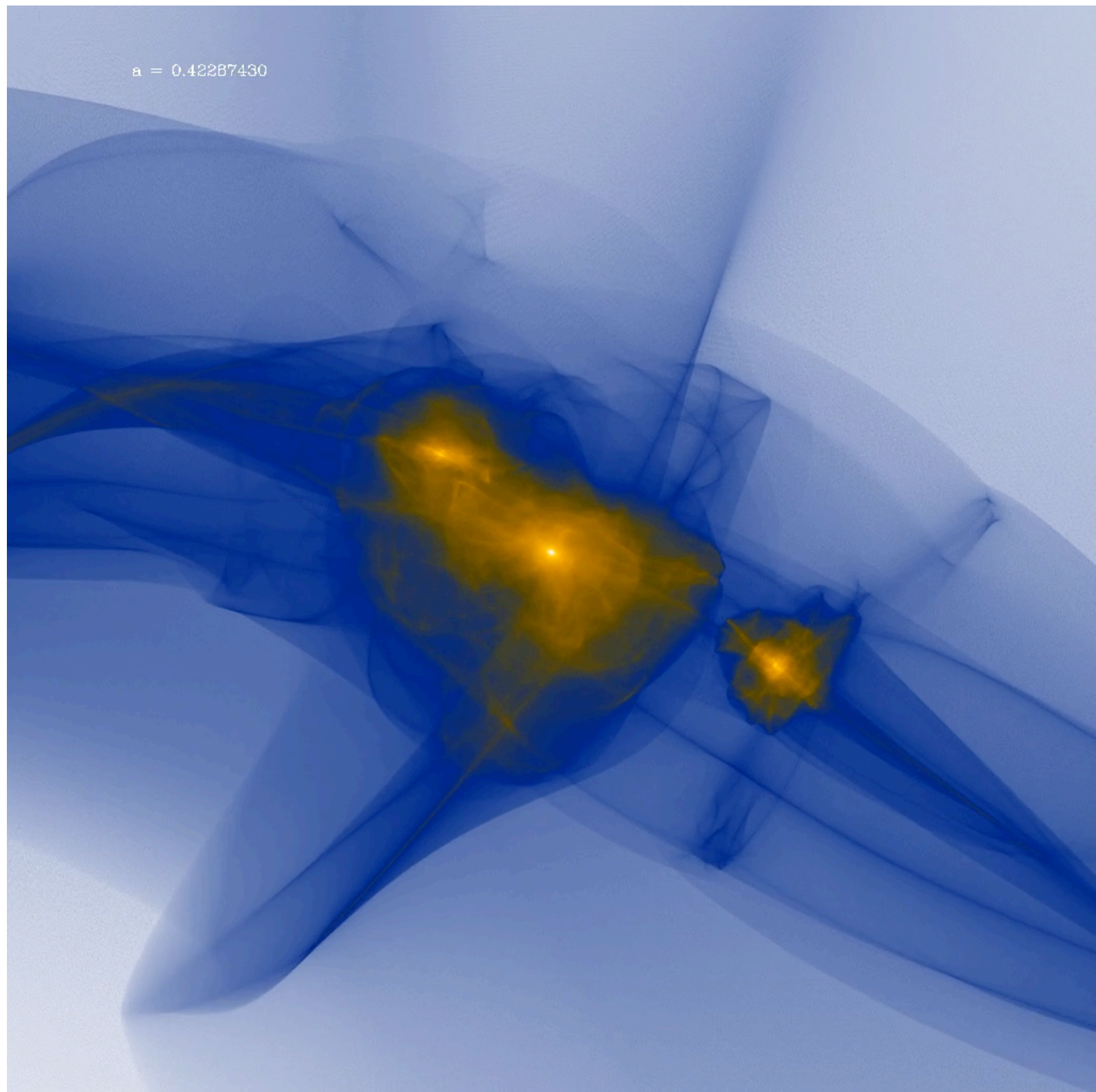
Tom Abel,<sup>★</sup> Oliver Hahn<sup>★</sup> and Ralf Kaehler<sup>★</sup>

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Warm dark  
matter halo with  
refinement and  
quadratic elements

Analogous to first  
cold dark matter halos

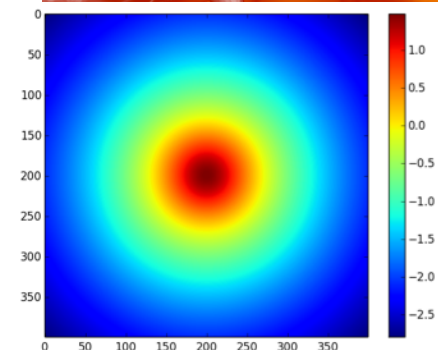
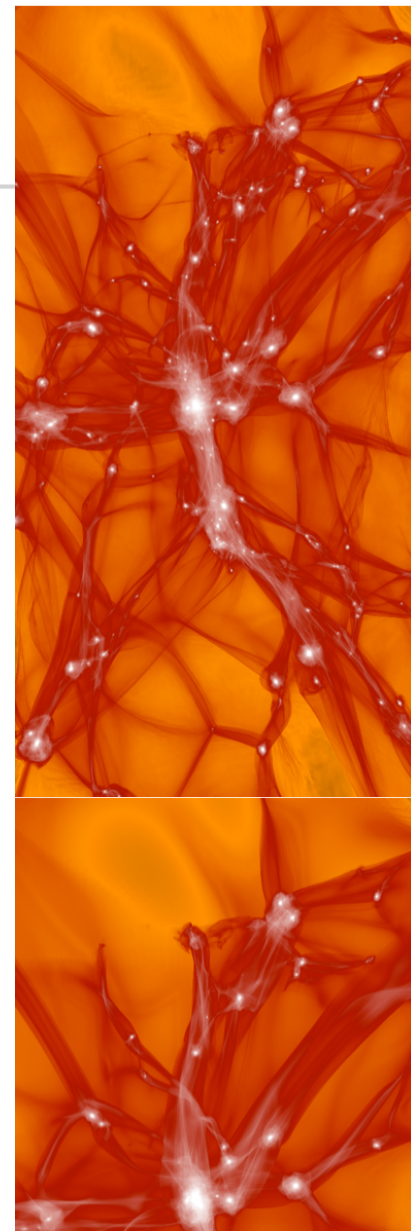


# Cosmological N-body simulations

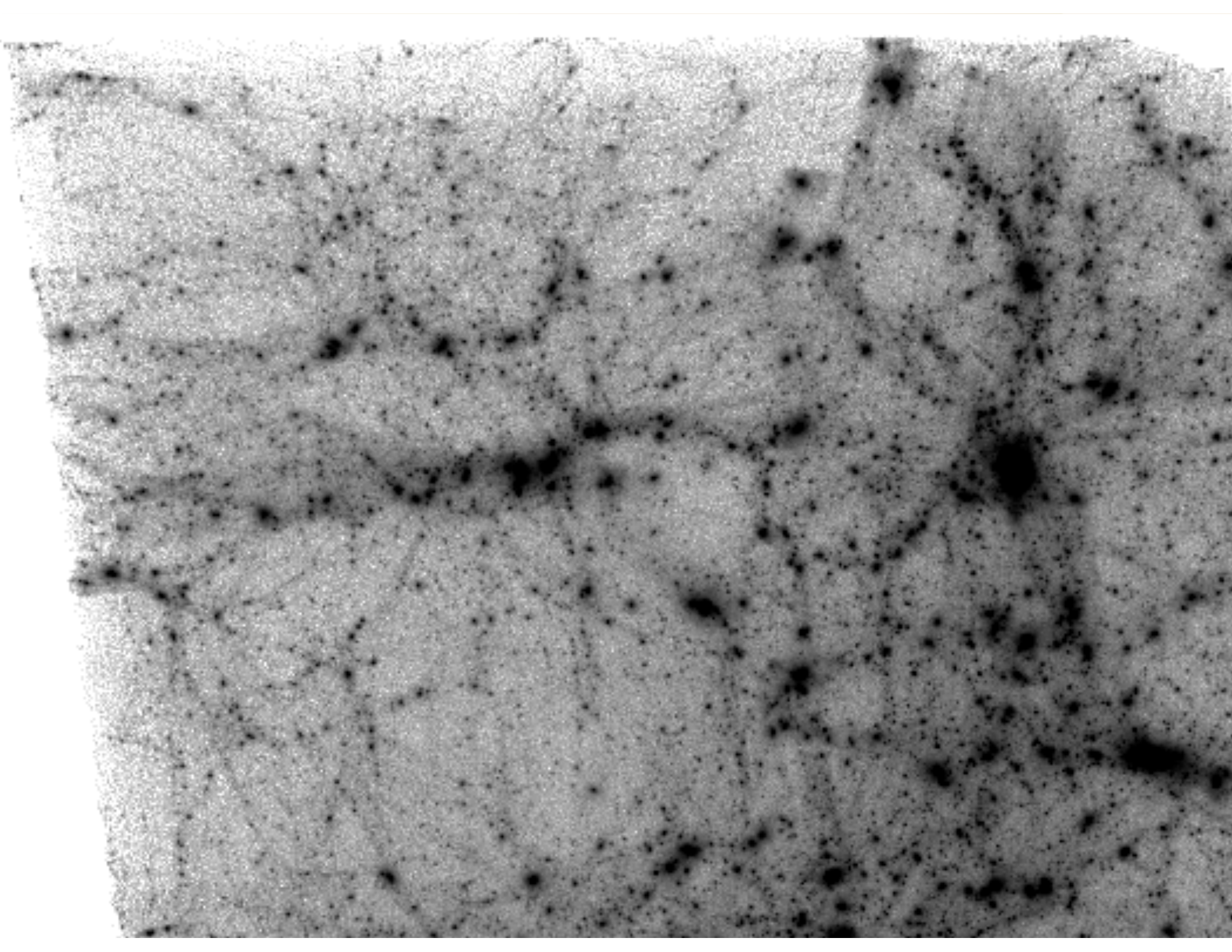
$$\dot{\mathbf{x}} = \mathbf{v}(t) \quad \dot{\mathbf{v}}_i = - \sum_{i \neq j}^N G m_i m_j \frac{(\mathbf{x}_j - \mathbf{x}_i)}{|\mathbf{x}_j - \mathbf{x}_i|^3}$$

- All modern cosmological simulation codes only differ in how they accelerate the computation of the sum over all particles to obtain the net force
- End result are simply the positions and velocities of all particles
- Softening of forces (add  $\epsilon^2$  in denominator) avoids singularities.
- Limit  $N$  goes to infinity must give correct answer, right?
- Plummer softening

$$\dot{\mathbf{v}}_i = - \sum_{i \neq j}^N G m_i m_j \frac{(\mathbf{x}_j - \mathbf{x}_i)}{(|\mathbf{x}_j - \mathbf{x}_i|^2 + \epsilon^2)^{3/2}}$$









# GRAVITY:

POISSON EQUATION :  $\nabla^2 \phi = 4\pi G \rho$

CONTINUUM DESCRIPTION

$$\vec{F}/m = -\nabla\phi$$

TOTAL MASS DENSITY

VLASOV EQUATION

$$\frac{\partial f}{\partial t} + v \nabla_x f - \nabla\phi \nabla_v f = 0$$

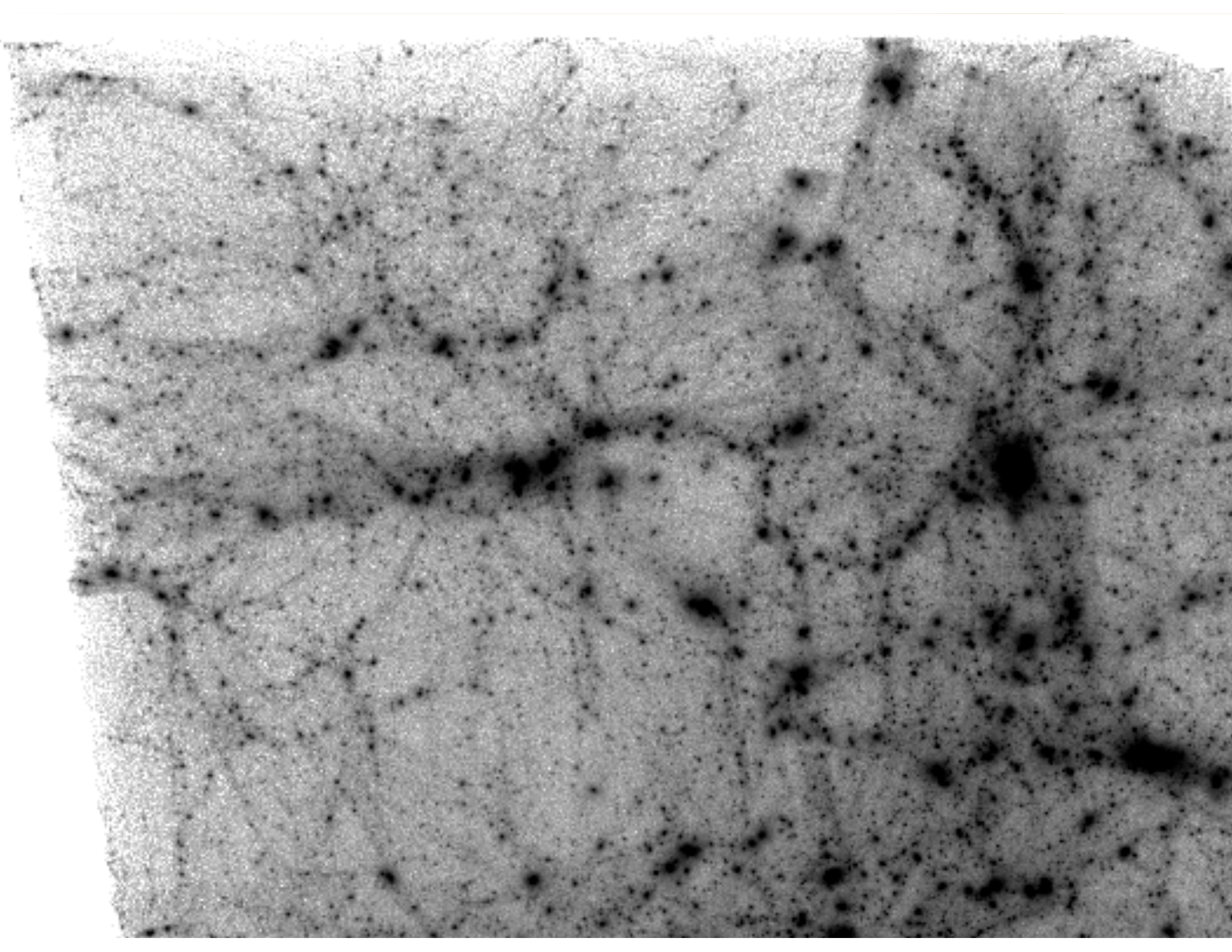
FOR  $N \rightarrow \infty$

IDENTICAL

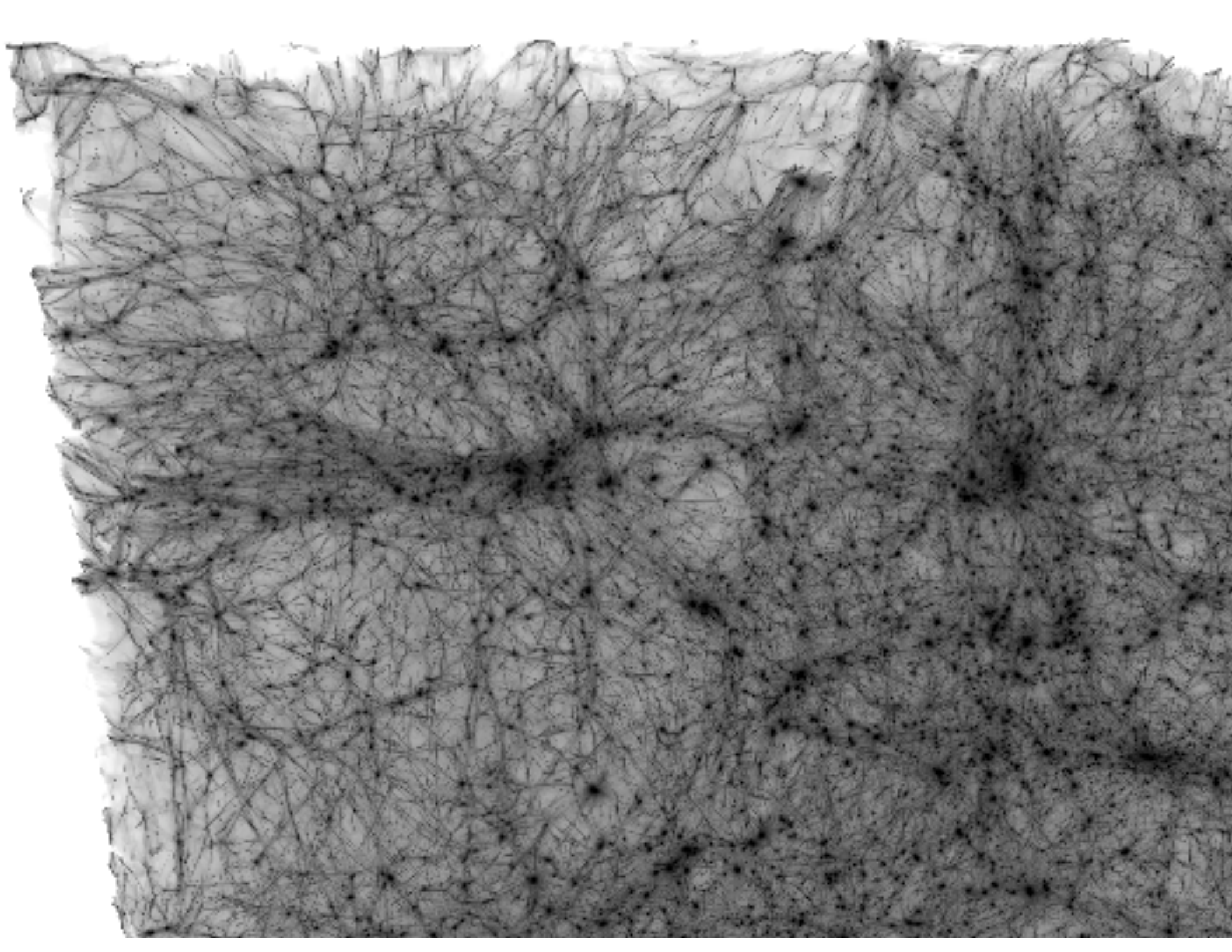
$N$  POINT MASSES :  $\vec{a}_j = - \sum_{i \neq j} \frac{G m_i}{|\vec{x}_j - \vec{x}_i|^2 + \epsilon^2} \frac{(\vec{x}_j - \vec{x}_i)}{|\vec{x}_j - \vec{x}_i|}$

PARTICLE PICTURE

Particle "advection":  $\frac{\partial \vec{x}_j}{\partial t} = \vec{v} ; \frac{\partial \vec{v}_j}{\partial t} = \vec{a}_j$







# The Dark Matter Sheet?

Quid

# OF DARK MATTER PARTICLES IN THE MILKY WAY :

$$N_{DM} \approx 10^{67} \left( \frac{100 \text{ GeV}}{m_{DM}} \right) \gg \# \text{ OF STARS IN THE UNIVERSE}$$

$\gg$  # OF PARTICLES THAT FIT ON A COMPUTER  
USING ALL THE COMPUTERS IN THE WORLD :  $\approx 10^{17}$  particles

SOLVE VLASOV-POISSON SYSTEM INSTEAD.

$$\frac{\partial f}{\partial t} + \vec{v} \cdot \nabla_x f + \vec{a} \cdot \nabla_v f = 0$$

$f$ : distribution function in PHASE SPACE

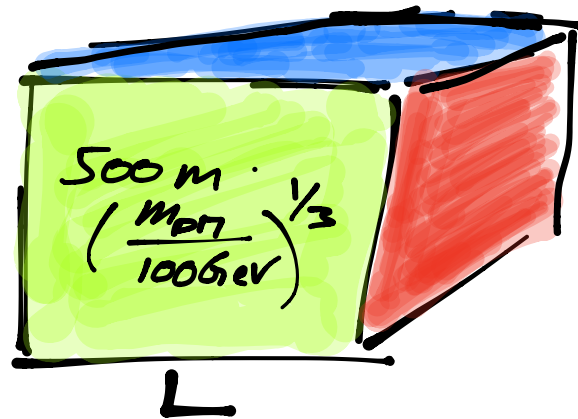
$\phi$ : potential

$$\vec{a} = -\nabla \phi$$

$$\nabla^2 \phi = 4\pi G \rho$$

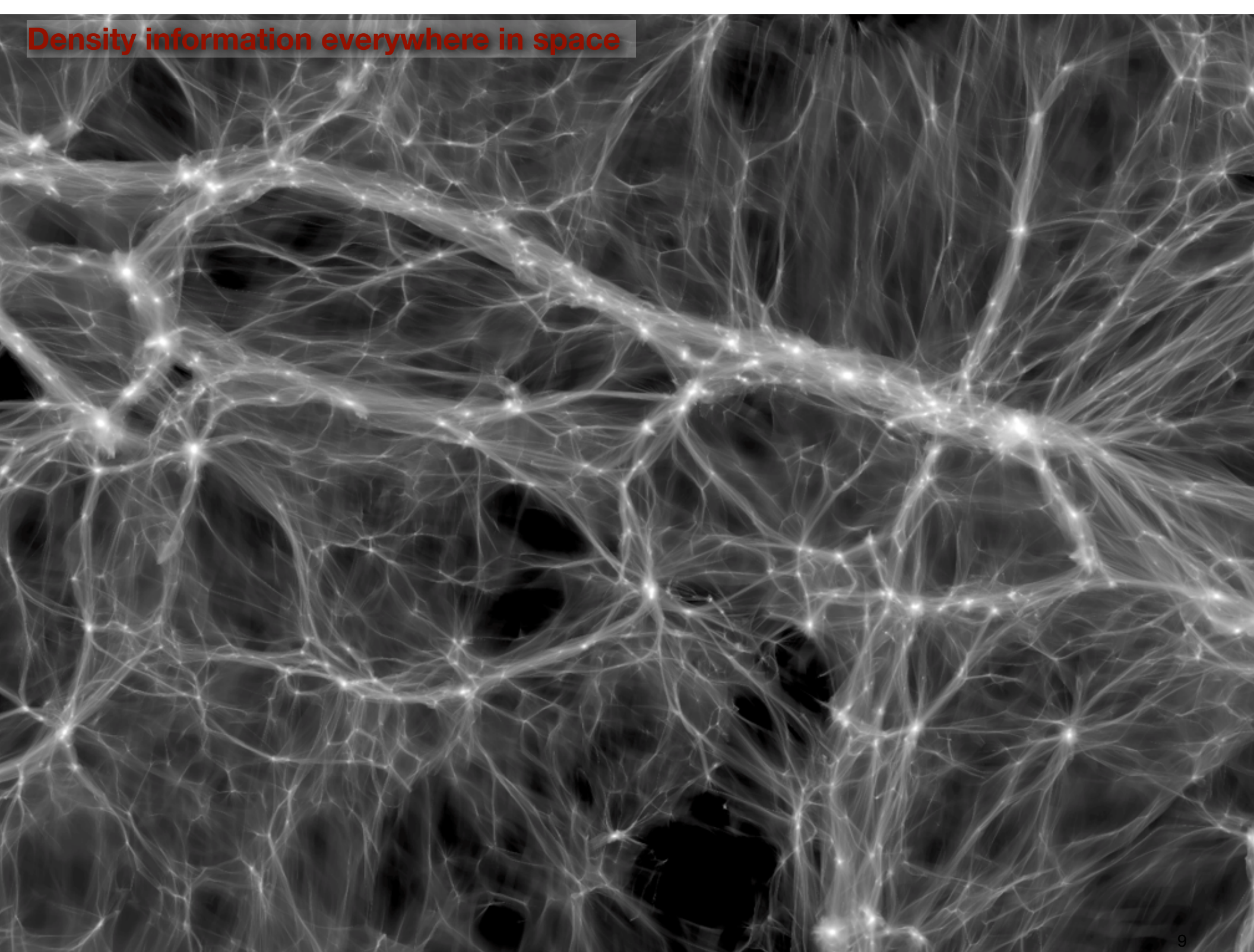
FOR PHASE SPACE ELEMENT TO CONTAIN  $10^6$  PARTICLES @ MEAN DENSITY IT HAS TO BE LARGER THAN

$$L \sim 500 \text{ m} \left( \frac{m_{DM}}{100 \text{ GeV}} \right)^{1/3}$$





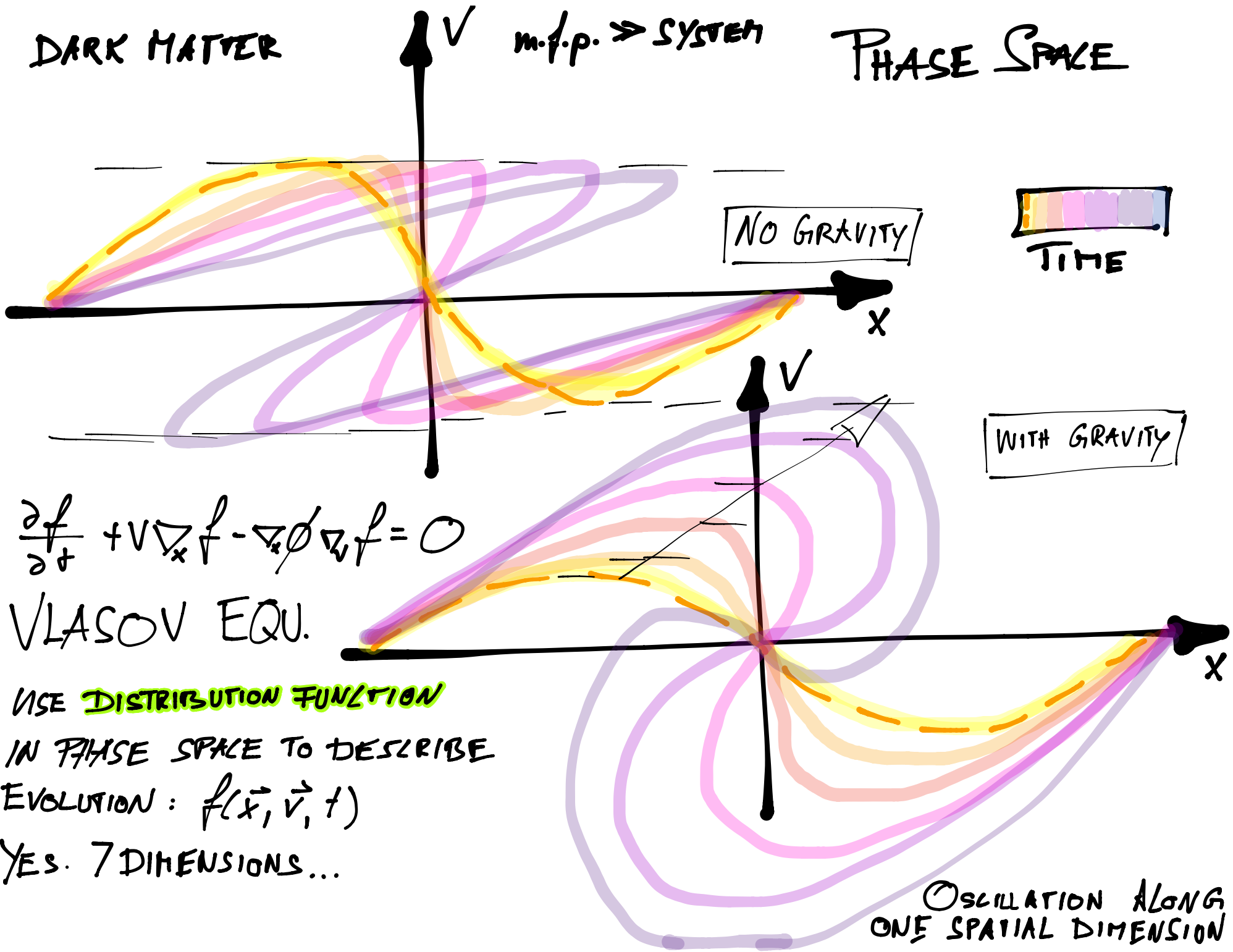
Density information everywhere in space



DARK MATTER

m.f.p.  $\gg$  SYSTEM

PHASE SPACE



$$\frac{\partial f}{\partial t} + v \nabla_x f - \nabla_x \phi \nabla_v f = 0$$

VLASOV EQU.

USE DISTRIBUTION FUNCTION

IN PHASE SPACE TO DESCRIBE

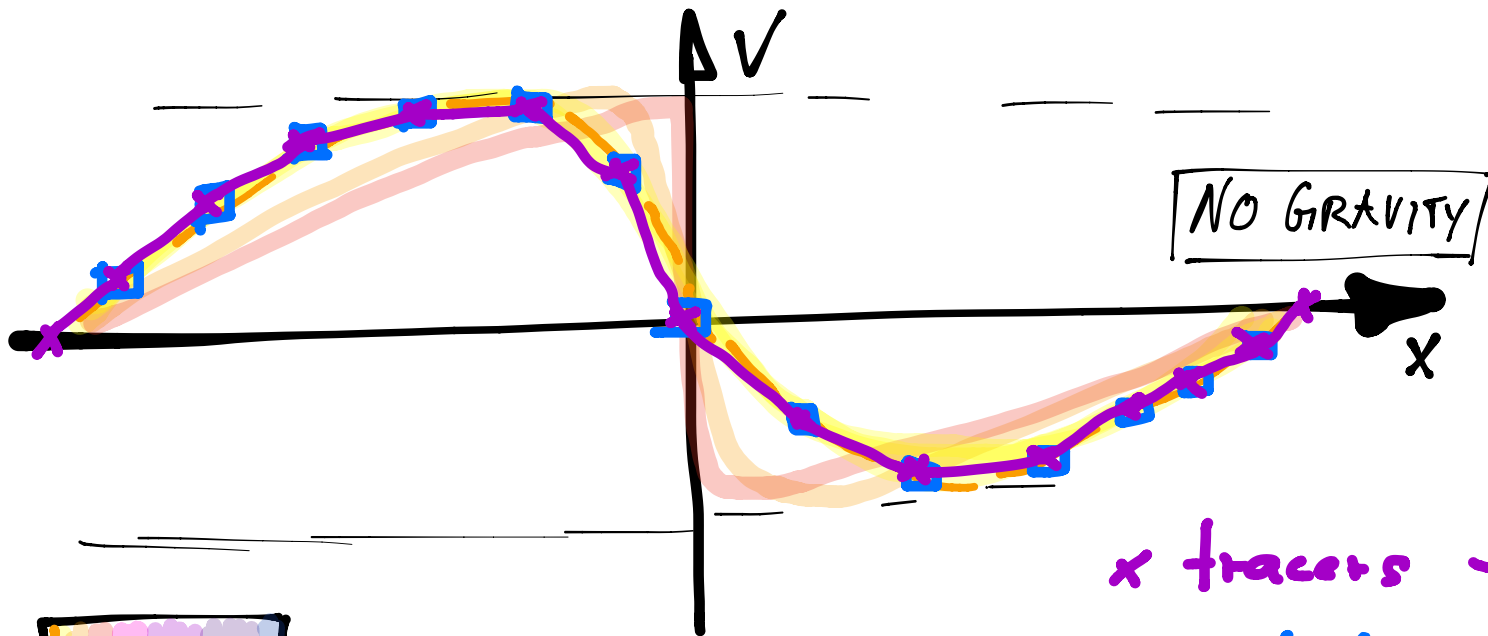
EVOLUTION:  $f(\vec{x}, \vec{v}, t)$

YES. 7 DIMENSIONS...

OSCILLATION ALONG ONE SPATIAL DIMENSION



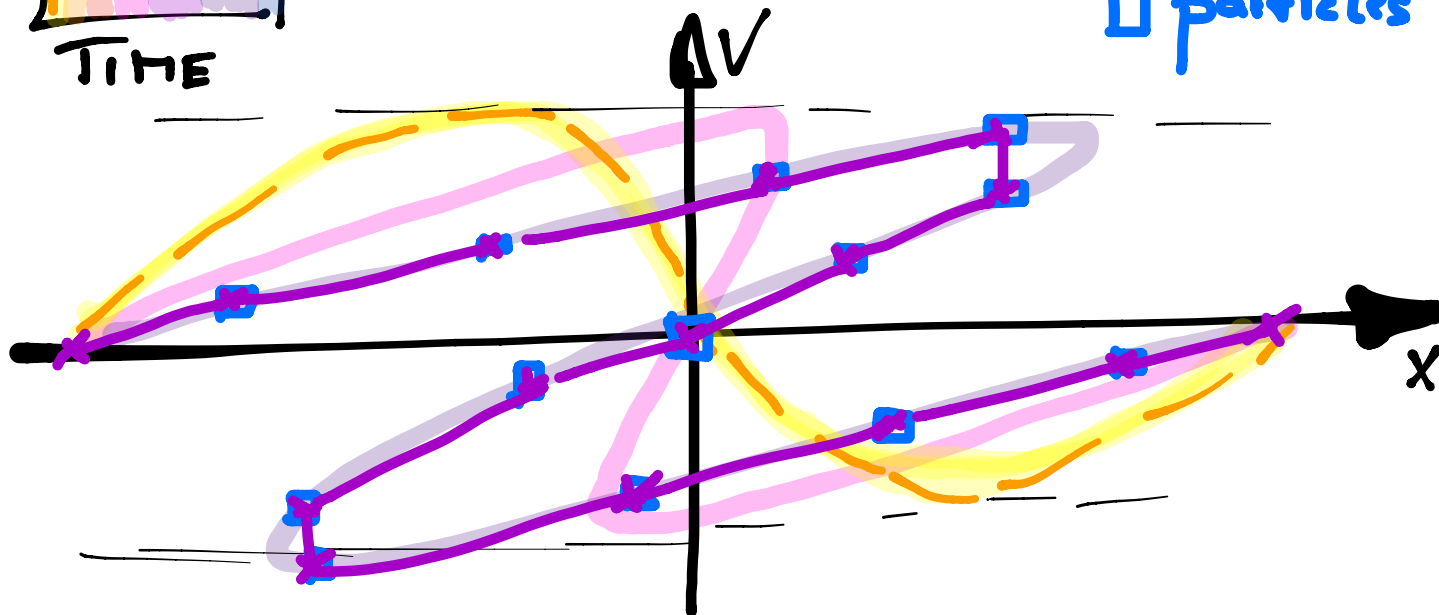
# DISCRETIZE DARK MATTER DISTRIBUTION: Mass or Volume?



COUNT PARTICLES  
IN RANDOM VOLUMES  
GIVES AVERAGE  
DENSITIES



x tracers — segments  
□ particles



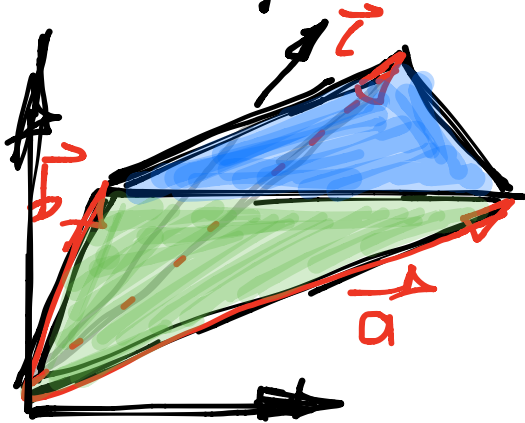
THINK MASS BETWEEN  
PARTICLES !

⇒ DENSITY KNOWN  
EVERYWHERE !

# TESSELATE 3D MANIFOLD & TRACK IN 6D PHASE SPACE

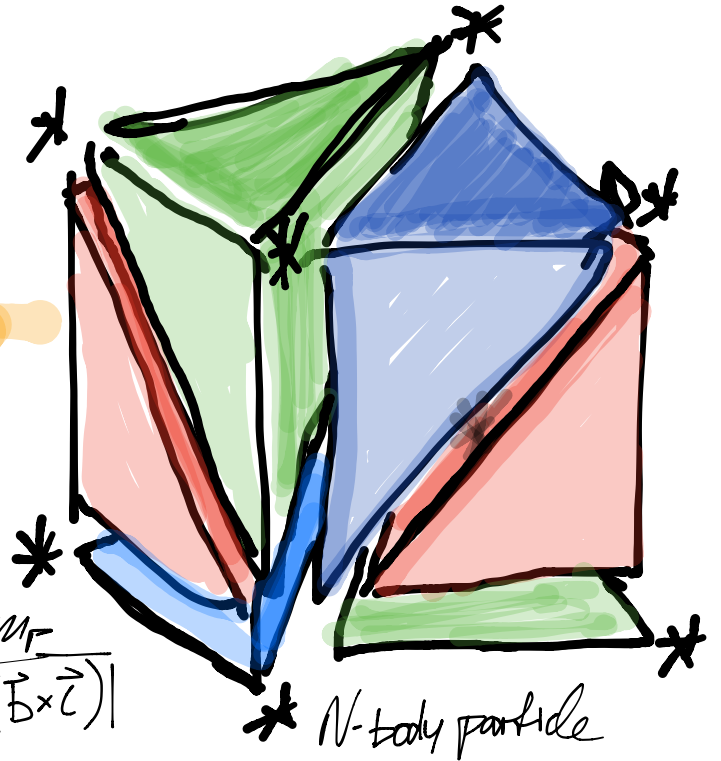
- NATURAL TESSELLATION SPLITS CUBE INTO 6 EQUAL SIZED TETRAHEDRA

- mass per tetrahedron =  $\frac{1}{6}$  of DM particle mass

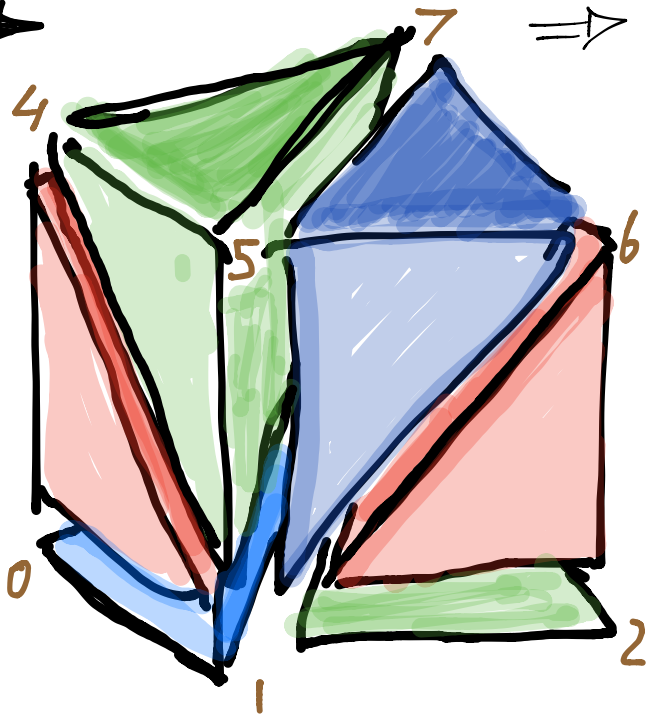


$$V = \frac{|\vec{a} \cdot (\vec{b} \times \vec{c})|}{6}$$

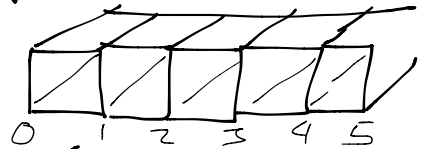
$$\Rightarrow \rho = \frac{M_P}{6V} = \frac{M_P}{|\vec{r} \cdot (\vec{b} \times \vec{c})|}$$



\* N-body particle



- Number the edges of the cube
- think of lattice
- Looping over

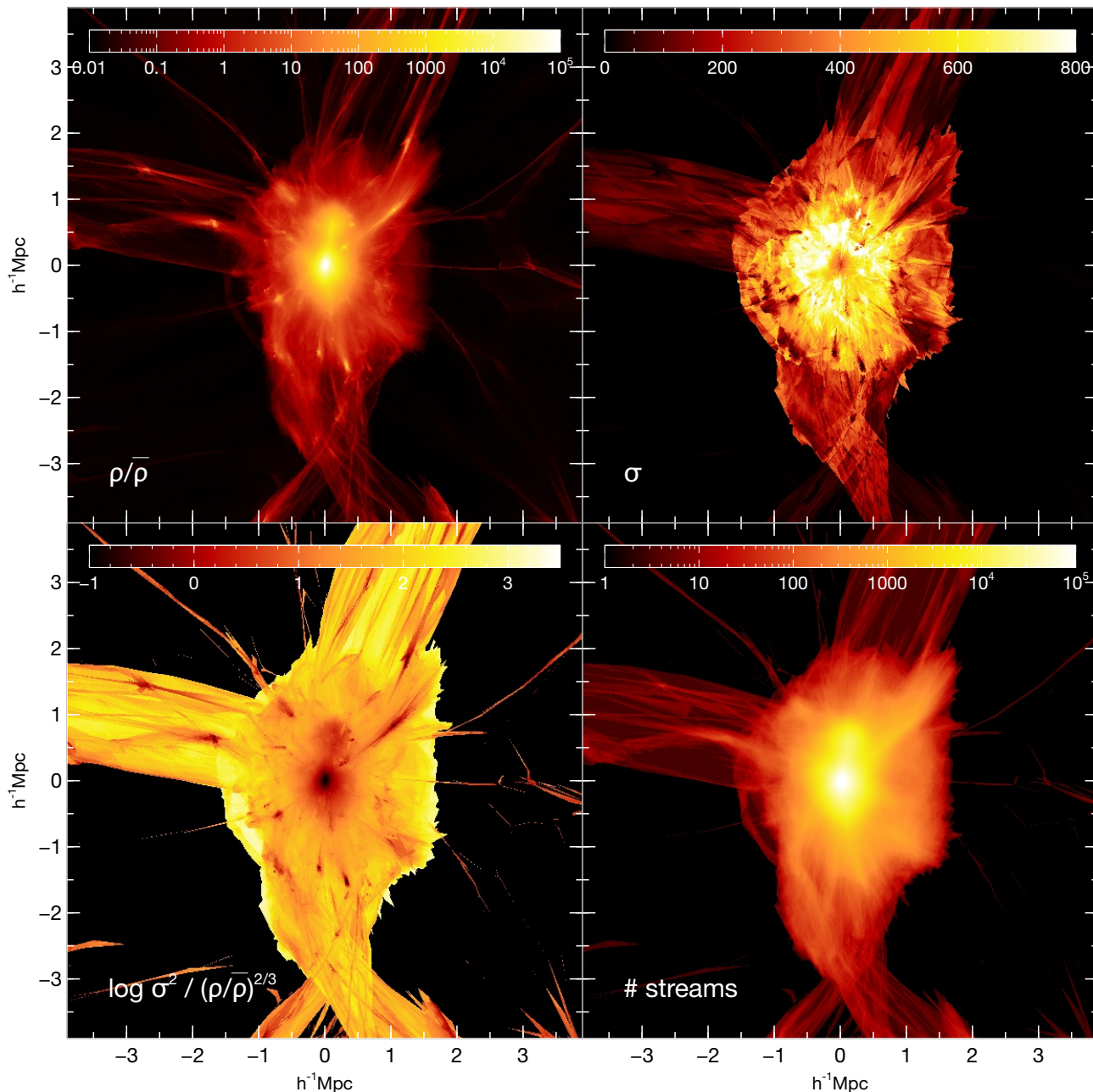
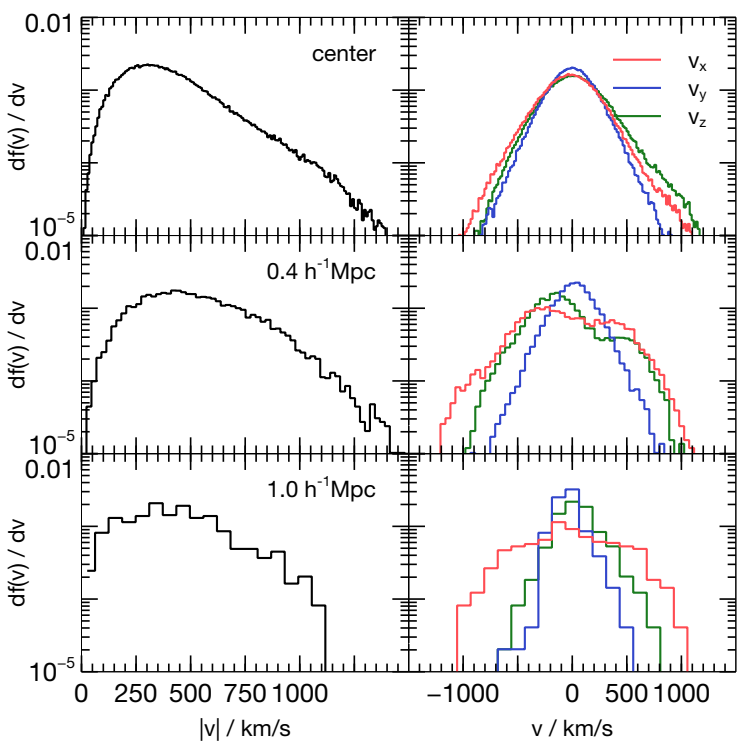


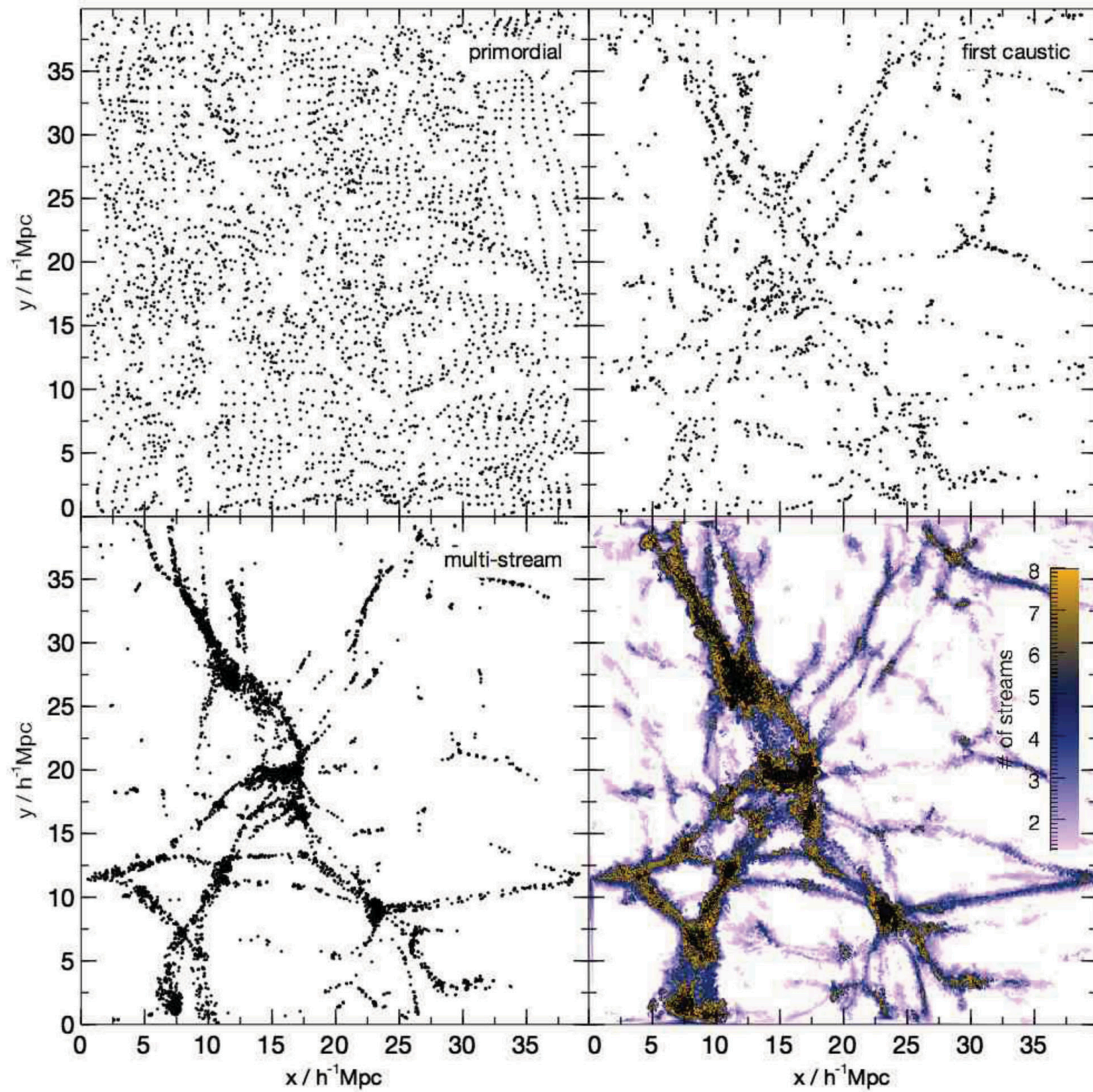
The initial cartesian (LAGRANGIAN) lattice generates the 6N tetrahedra.



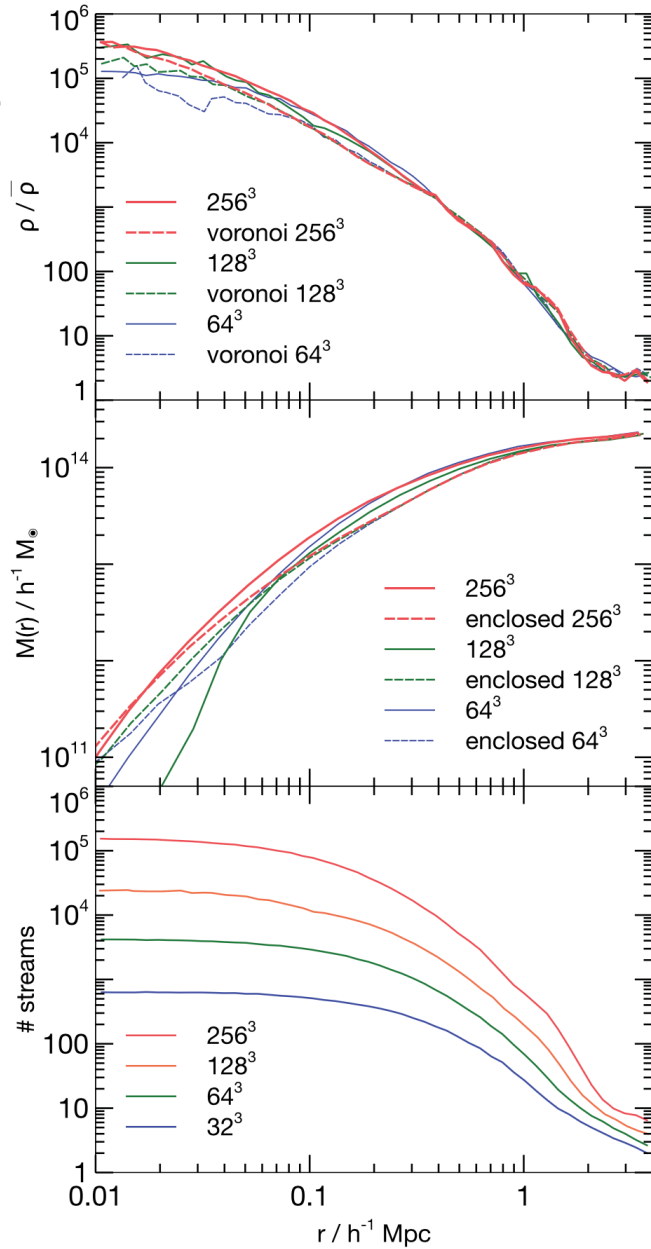
# All microphysical phase space information available

can probe  
fine-grained  
phase space  
structure.

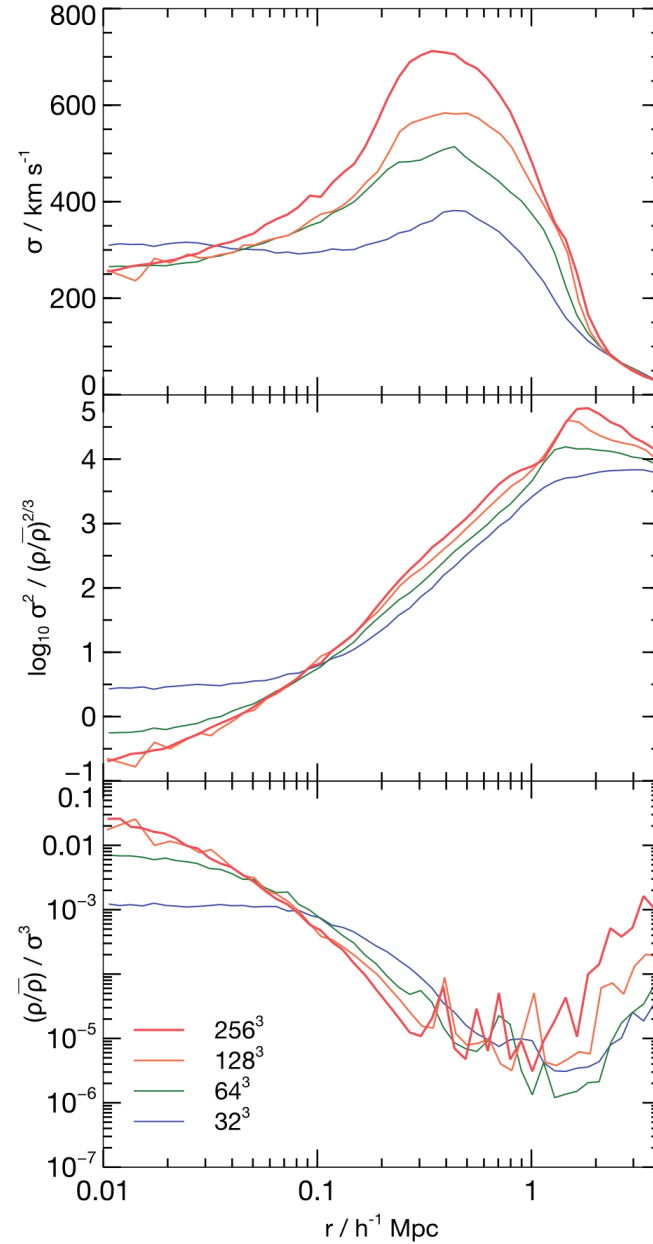




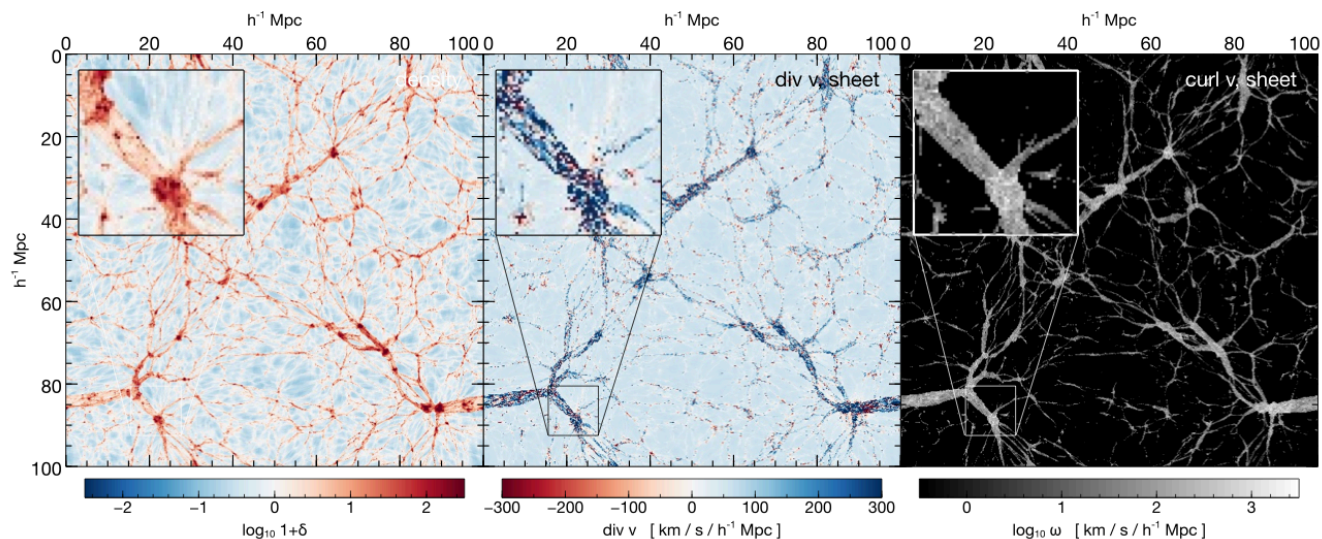
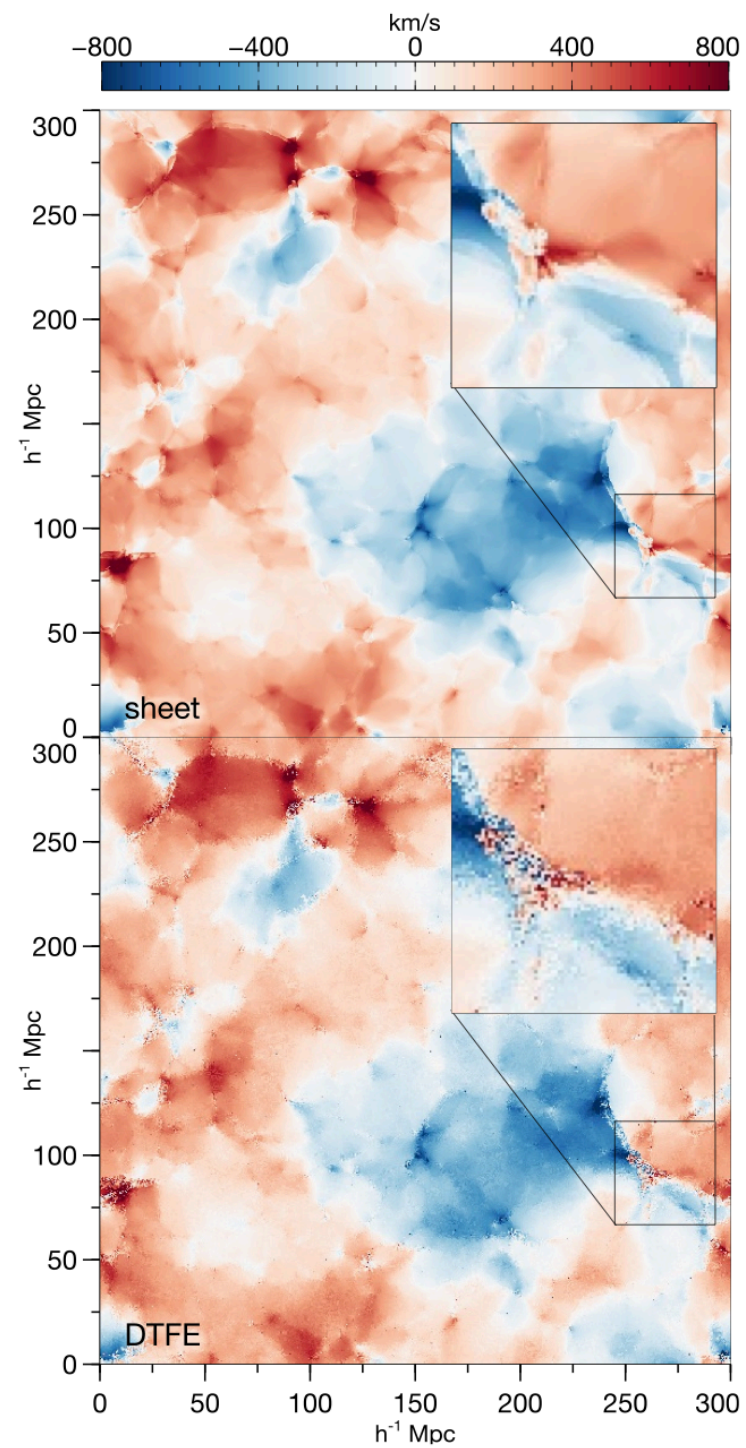




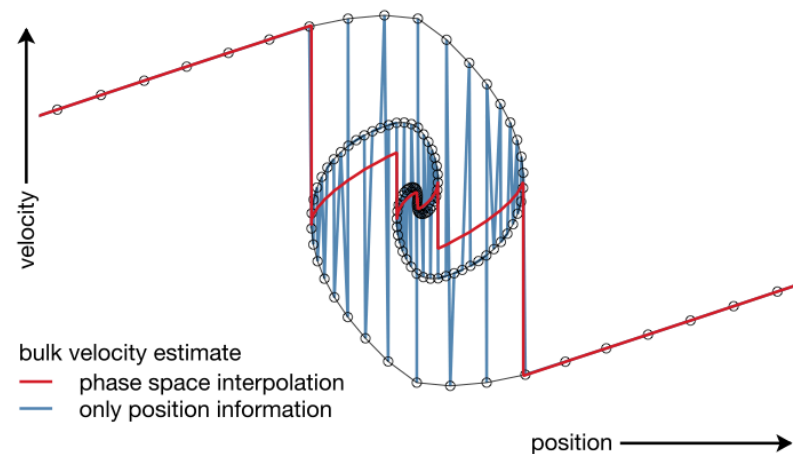
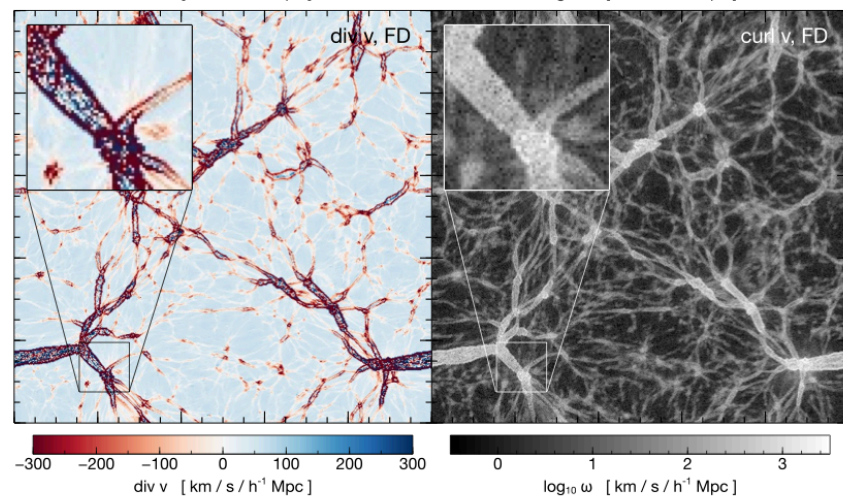
**Figure 14.** DM density profile in the most massive halo of  $2 \times 10^{14} M_{\odot}$  with  $R_{\text{vir}} \approx 1 \text{ Mpc}$  at redshift zero. The overdensity in the top, the enclosed DM (middle) as calculated from the density in the top panel and the number of streams (bottom panel) are shown for all the different resolution simulations studied in this paper. The density profile estimated from our method starts to differ from the conventional estimate at scales as large as one-third of the virial radius. As expected in CDM, the number of streams does not converge also in radial profiles as the resolution is increased.



**Figure 15.** Radial spherically averaged profiles of the velocity dispersion (top), the DM ‘entropy’  $\sigma^2 / (\bar{\rho})^{2/3}$  (middle) and the pseudo-phase-space density (bottom) for the same halo as in Fig. 14. The velocity dispersion is remarkably flat inside about one-tenth of the virial radius. The DM ‘entropy’ profile also shows signs of already converging at the modest resolutions employed here. Using the microscopic velocity dispersion of our approach which removes the bulk flows does not give the typical power-law behaviour in the pseudo-phase-space density found when using the total dispersion of particles at those radii.



# Cosmic Velocity Fields



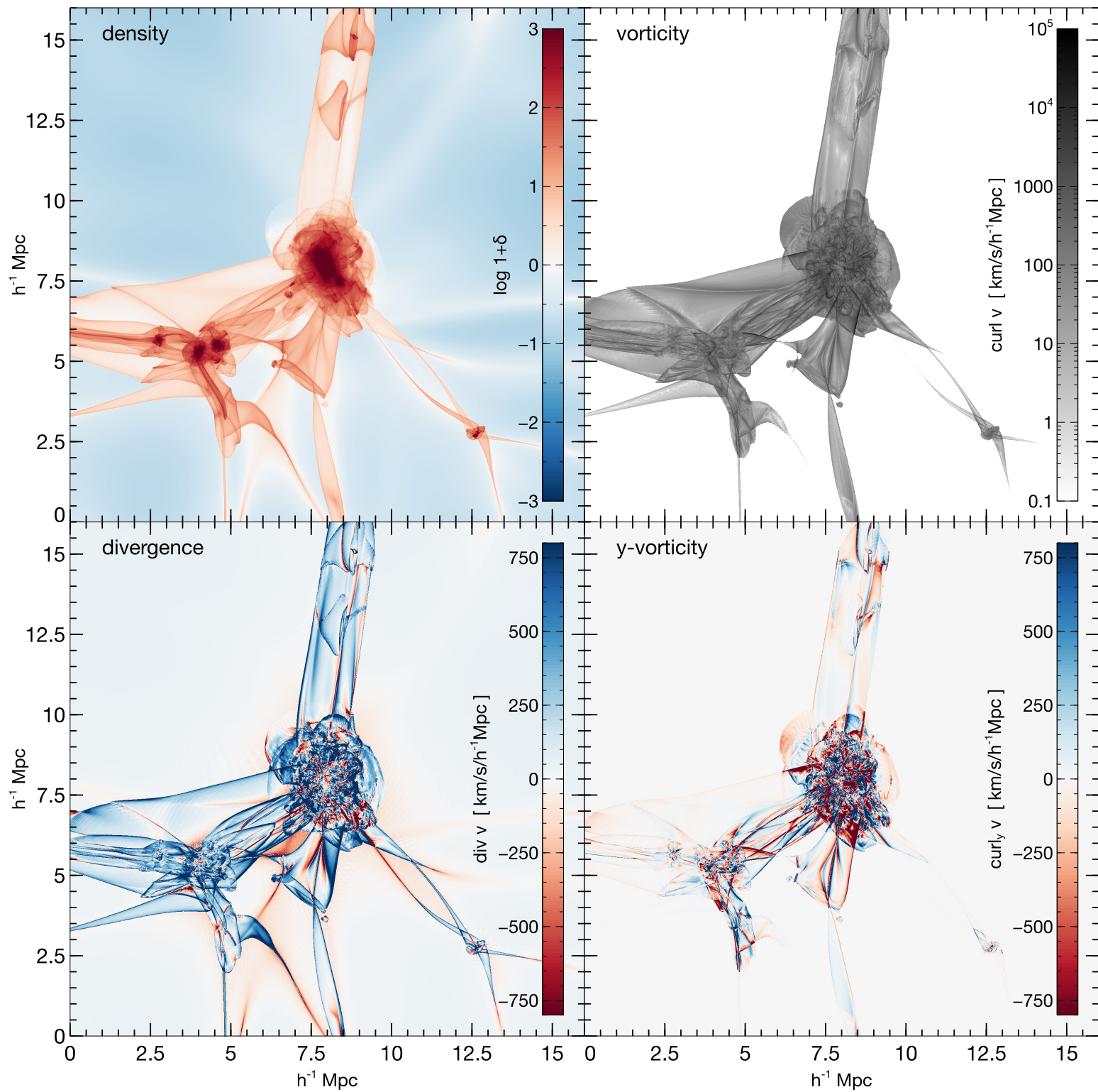
The properties of cosmic velocity fields

Oliver Hahn,<sup>1\*</sup> Raul E. Angulo<sup>2</sup> and Tom Abel<sup>3</sup>



# The properties of cosmic velocity fields

Oliver Hahn,<sup>1\*</sup> Raul E. Angulo<sup>2</sup> and Tom Abel<sup>3</sup>

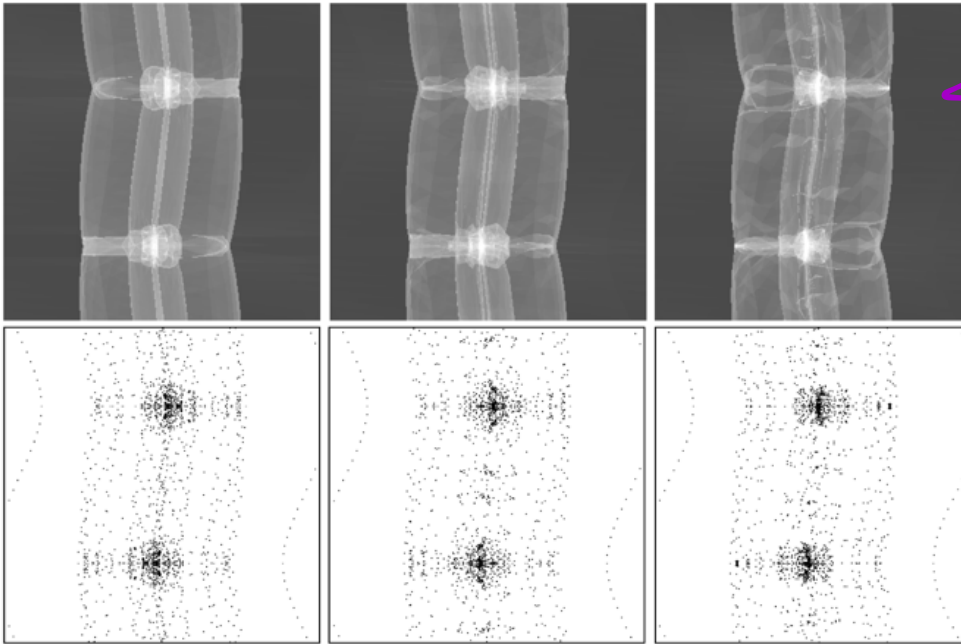


128<sup>3</sup>

256<sup>3</sup>

512<sup>3</sup>

std. PM

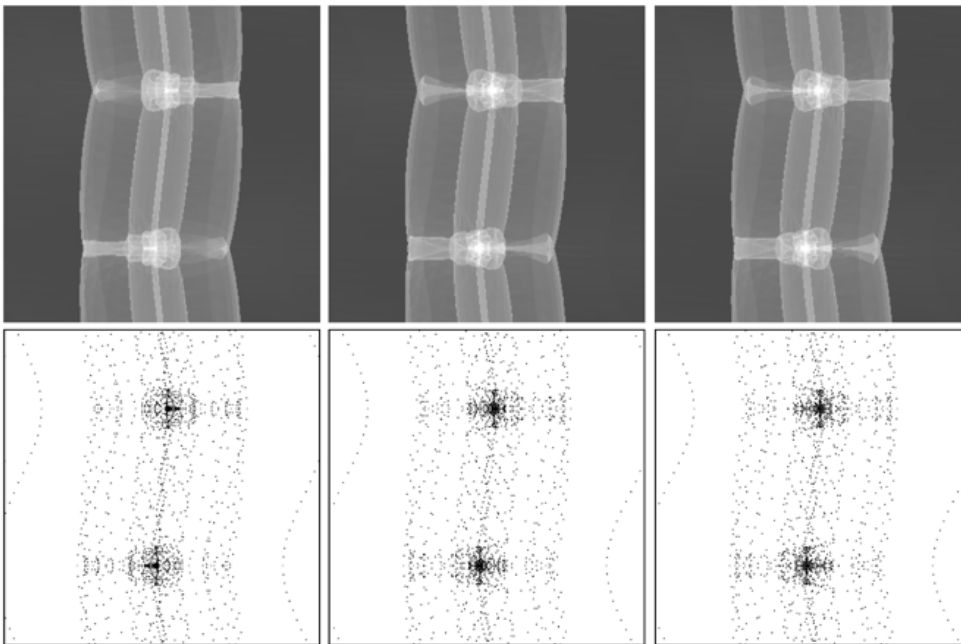


← two body scattering  
undesirable ...

Note how the new visualization  
technique helps in spotting  
errors in the N-body integration

Constant mass resolution  
Vary force resolution

TCM



← excellent convergence  
behavior of our new  
method.



# Solving the Vlasov equation in two spatial dimensions with the Schrödinger method

Michael Kopp,<sup>1,2,\*</sup> Kyriakos Vattis,<sup>1,3,†</sup> and Constantinos Skordis<sup>1,2,‡</sup>

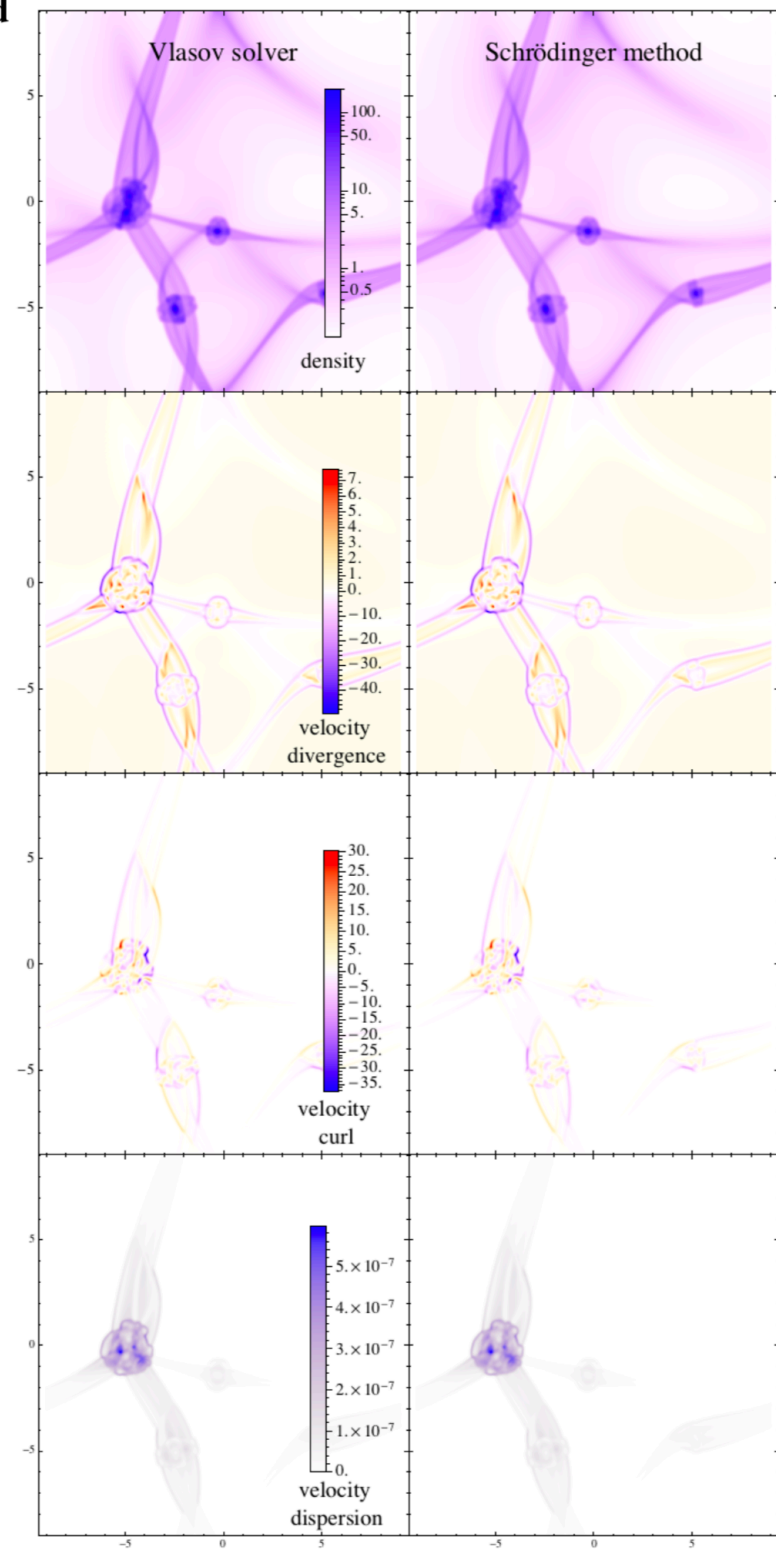
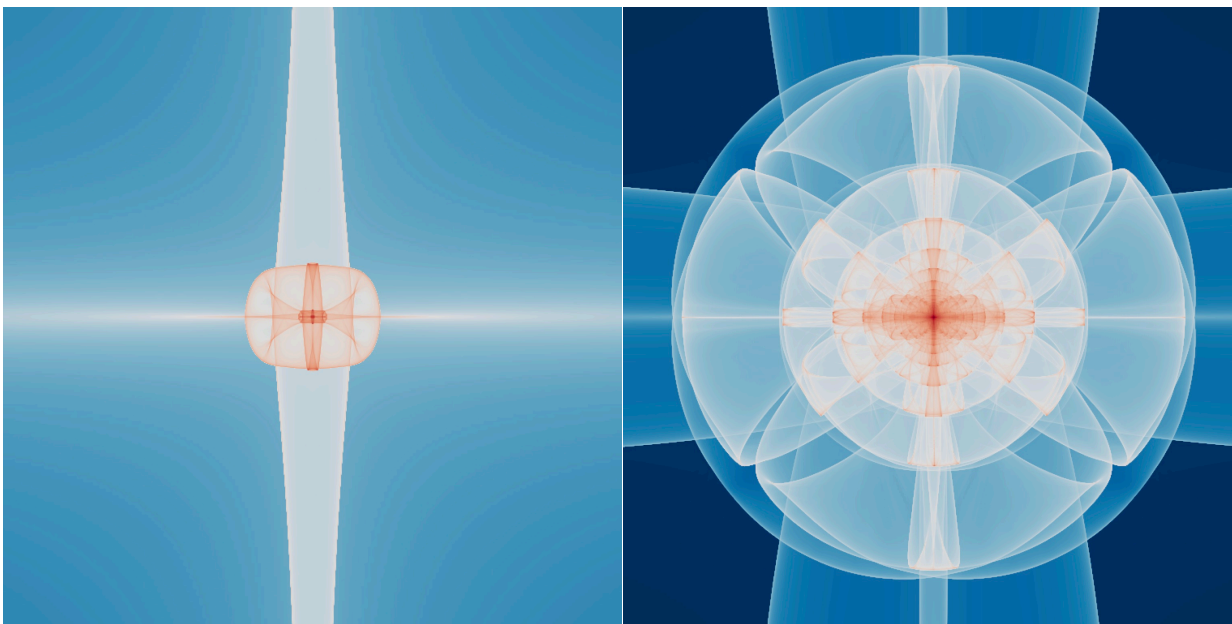
Physical Review D, Volume 96, December 2017

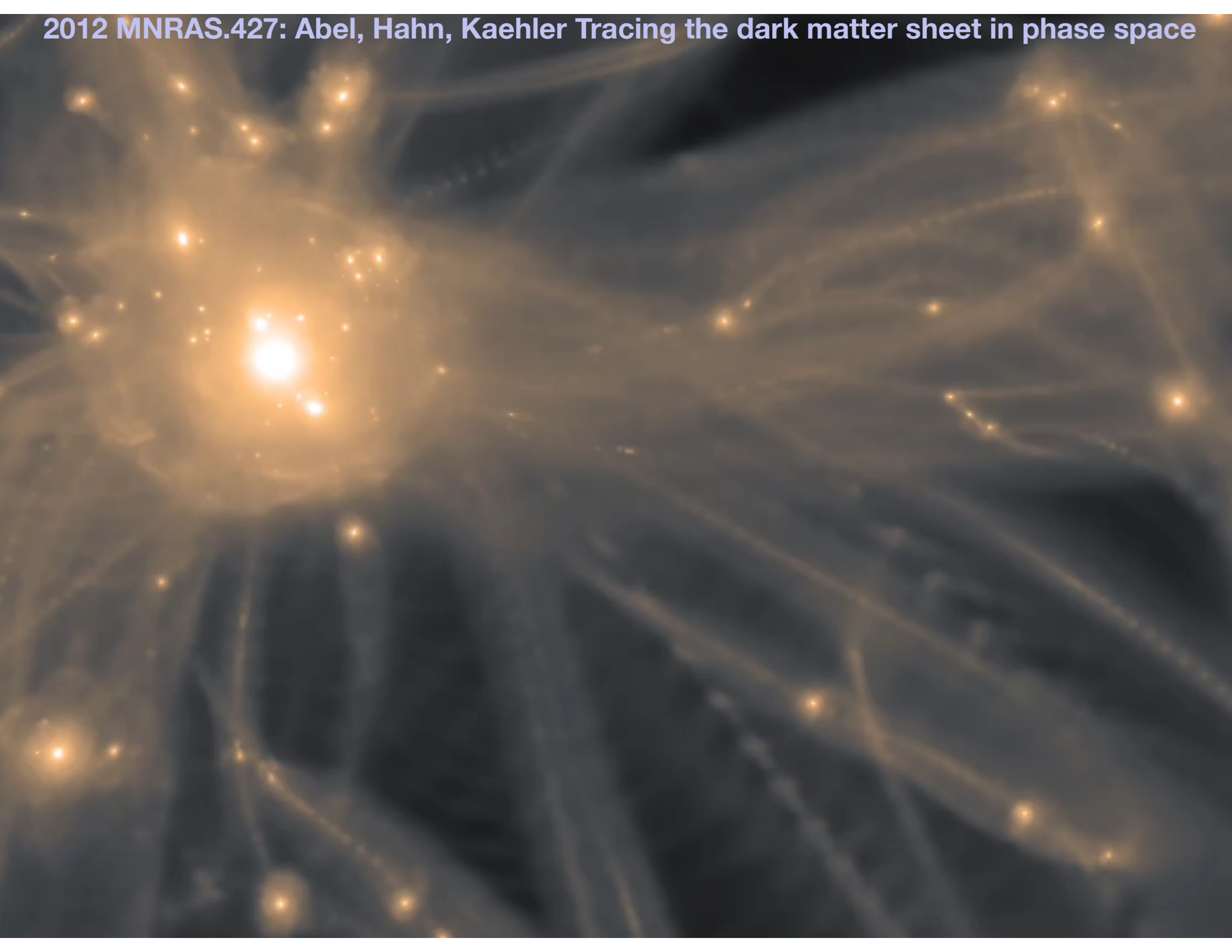
FIG. 1. Time snapshots of a two-dimensional cosmological simulation evolved to the present time  $a = 1$ . Both codes were started with the same single-stream initial conditions, set up using the Zel'dovich approximation at  $a_{\text{ini}} = 1/51$ . From top to bottom: density, velocity divergence, velocity curl and trace of the velocity dispersion tensor. The left column shows the smoothed results of the 2D-version of the Vlasov solver CoLDICE [13]. The result of the Schrödinger method is shown in the right column. The differences are barely visible by eye. A quantitative comparison is presented in Sec. IV.

used:

CoLDICE: a parallel Vlasov-Poisson solver using moving adaptive simplicial tessellation

Thierry Sousbie<sup>a,b,c,\*</sup>, Stéphane Colombi<sup>a</sup>



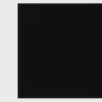
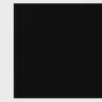




N-body

Quadrilaterals + refinement

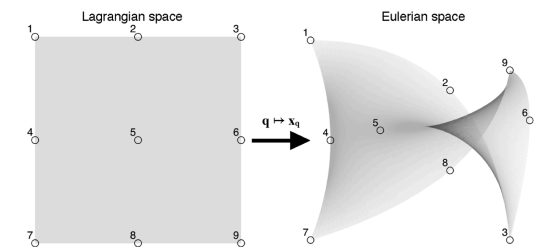
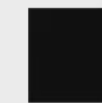
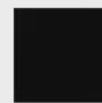
Quadrilaterals



125=5x5x5 points initially in cube  
N-body 2e6 particle reference

Linear tets + refinement

Linear tets



**Figure 1.** Mapping between Lagrangian space and Eulerian configuration space. The mapping can become multivalued in projection to configuration space and can have singular points and curves. Here, as an example, the mapping of nine points forming a square in Lagrangian space are shown under a biquadratic map to phase-space and then projected into configuration space.

**An adaptively refined phase-space element method for cosmological simulations and collisionless dynamics**

Oliver Hahn<sup>1\*</sup> and Raul E. Angulo<sup>2</sup>

<sup>1</sup>Department of Physics, ETH Zurich, CH-8093 Zurich, Switzerland

<sup>2</sup>Centro de Estudios de Física del Cosmos de Aragón, Plaza San Juan 1, Planta-2, E-44001 Teruel, Spain

Clear accuracy gains with higher order interpolation schemes.

Shown here in the test case of a cube evolving in a static potential.

Second order interpolation is clearly advantageous.

# Lagrangian Tessellation: What's it good for?

## Not complete list:

- Analyzing N-body sims, including web classification, velocity dispersion, profiles, resolution study (Abel, Hahn, Kaehler 2012)
- DM visualization (Kaehler, Hahn, Abel 2012)
- Better Numerical Methods (Hahn, Abel & Kaehler 2013, Hahn, Angulo & Abel 2014, Angulo and Hahn 2016, Sousbie & Colombi 2016)
- Finally reliable WDM mass functions below the cutoff scale Angulo, Hahn, Abel 2013
- Gravitational Lensing predictions, Angulo, Chen, Hilbert & Abel 2014
- Cosmic Velocity fields, Hahn, Angulo, Abel 2014
- The SIC method for Plasma simulations (Vlasov/Poisson) (Kates-Harbeck, Totorica, Zrake & Abel 2016, JCompPhys)
- Exact overlap integrals of Polyhedra (Powell & Abel 2015 JCompPhys)
- Adaptively refined phase-space, Hahn & Angulo 2016
- CoDICE: A parallel Vlasov-Poisson solver using moving adaptive simplicial tessellation, Sousby & Colombi 2016
- Void profiles, Wojtak, Powell, Abel 2016 MNRAS
- Stücker, Busch, & White 2017: Median density of the Universe
- East, Wojtak, Abel compare numerical GR and Newtonian cosmology 2017
- Powell & Abel 2018, Beam Tracing for radiation transport
- Banerjee, et al: Noiseless Cosmological Neutrino simulations, 2018
- Totorica, et. al. Weibel instabilities, shocks, particle acceleration in PIC simulations 2018
- “Nonthermal electron and ion acceleration by magnetic reconnection in large laser-driven plasmas” Totorica, Hoshino, Abel, Fiuza 2020
- Powell et al DM annihilation,
- Totorica & Abel: Vlasov-Maxwell (2023)