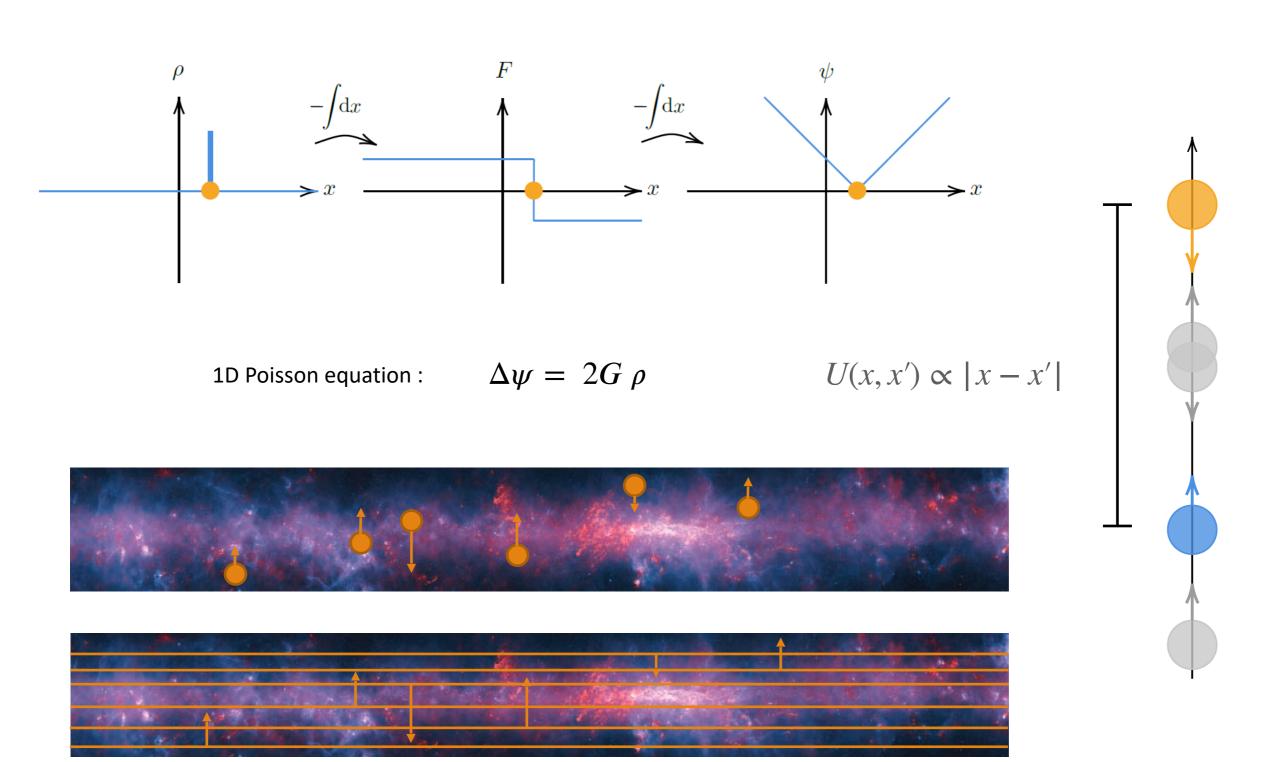
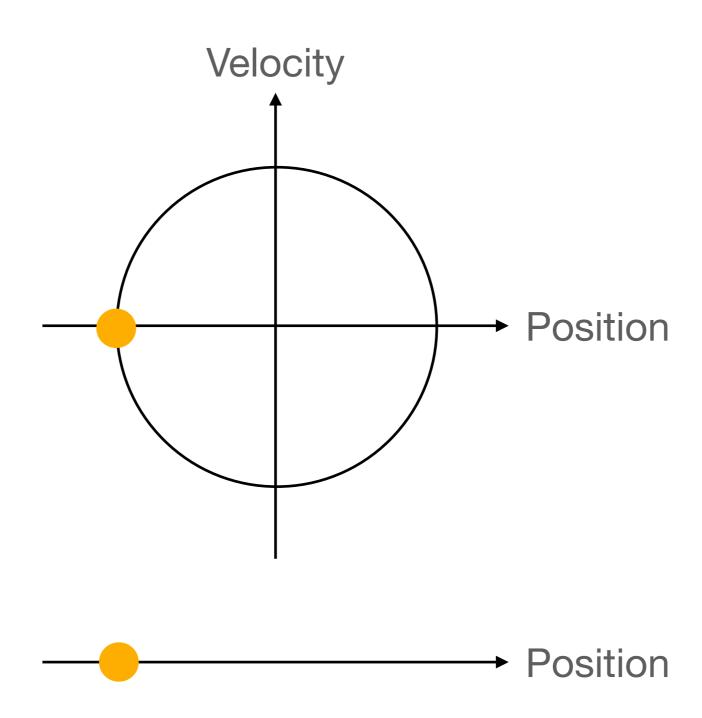
Secular relaxation of self-gravitating planes and disc

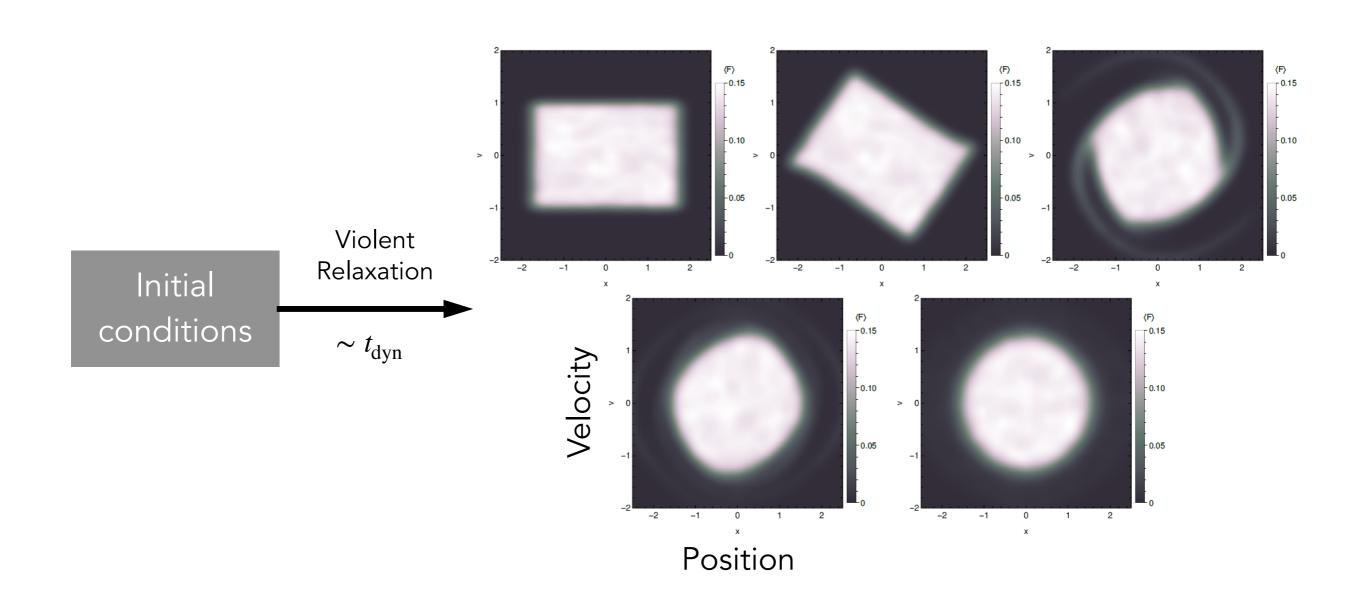
Mathieu Roule Institut d'Astrophysique de Paris roule@iap.fr

The one-dimensional self-gravitating system (planes)

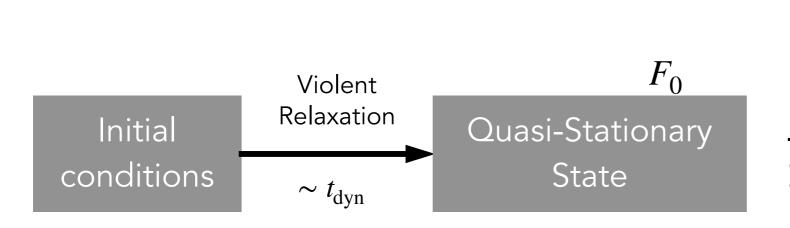


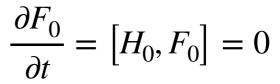
The one-dimensional self-gravitating system (planes)

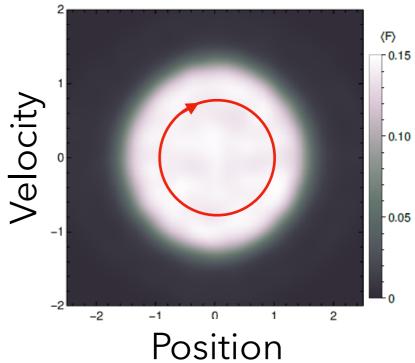


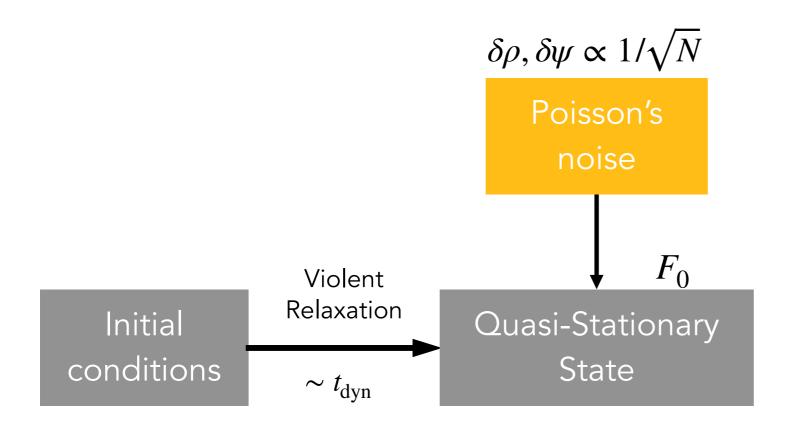


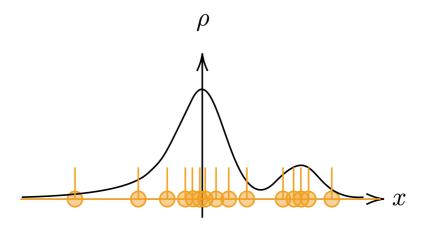


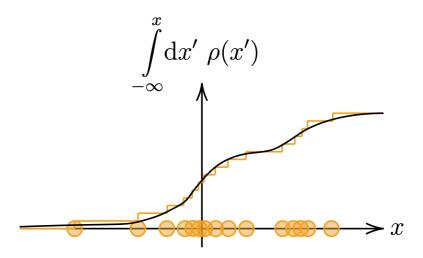


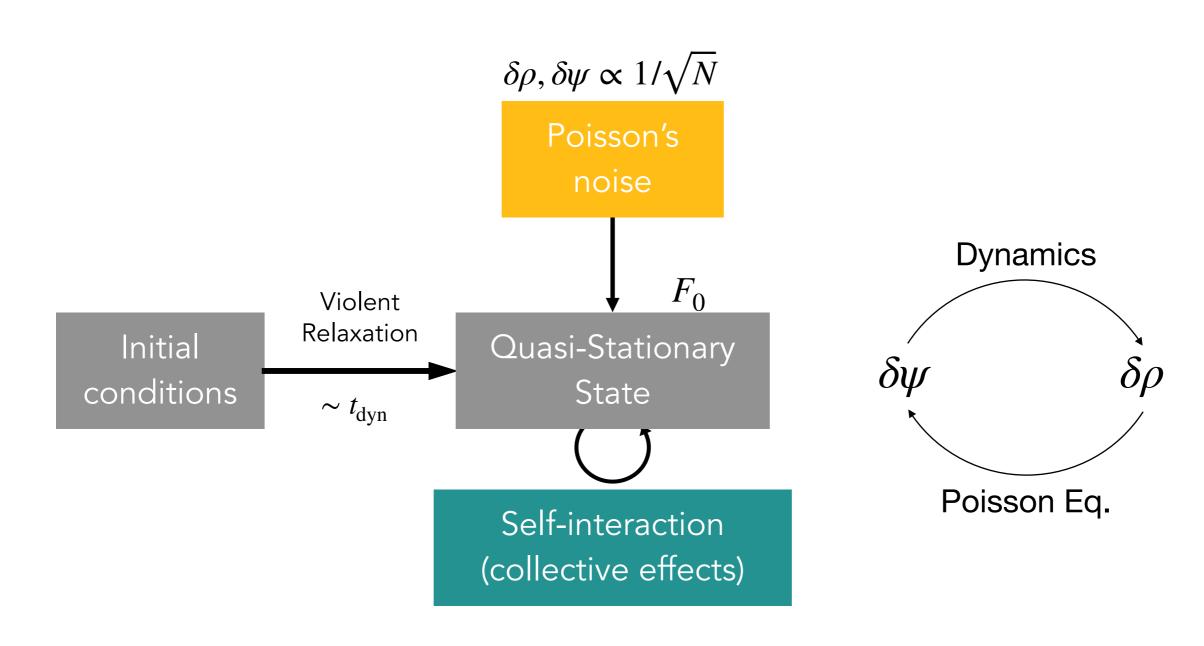


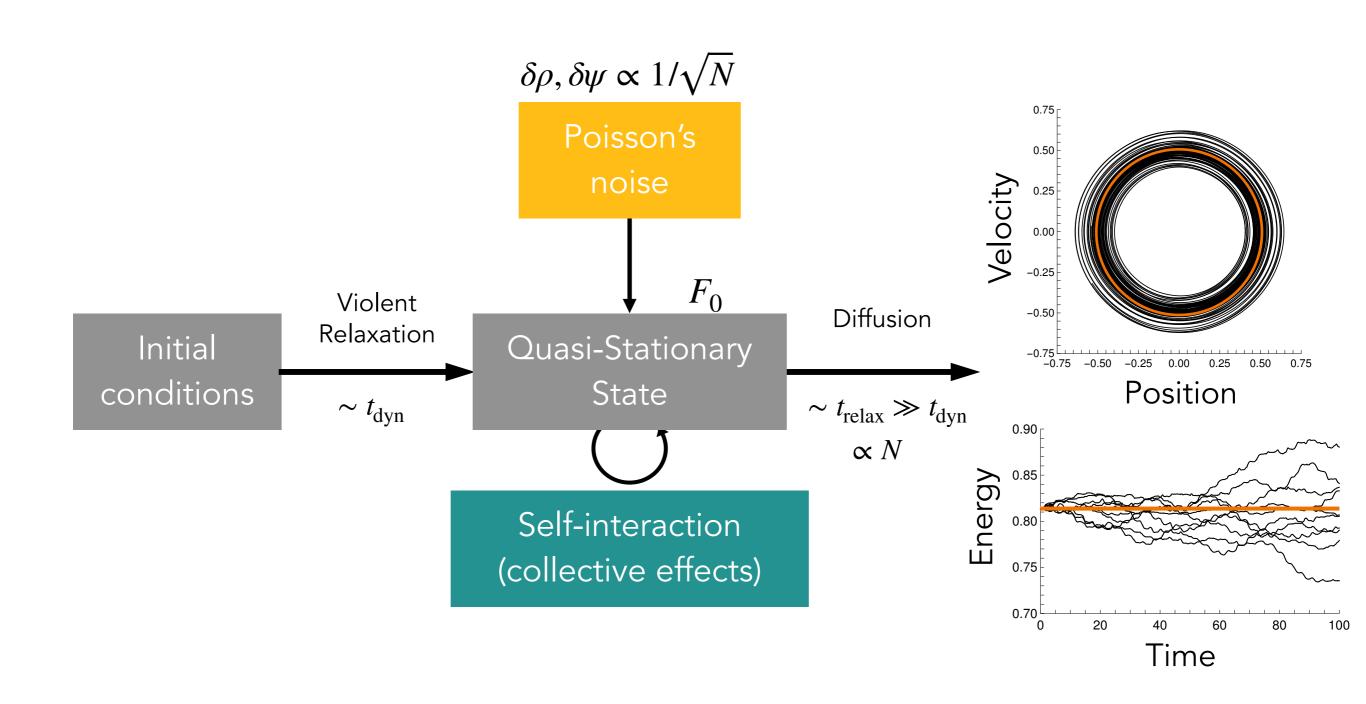


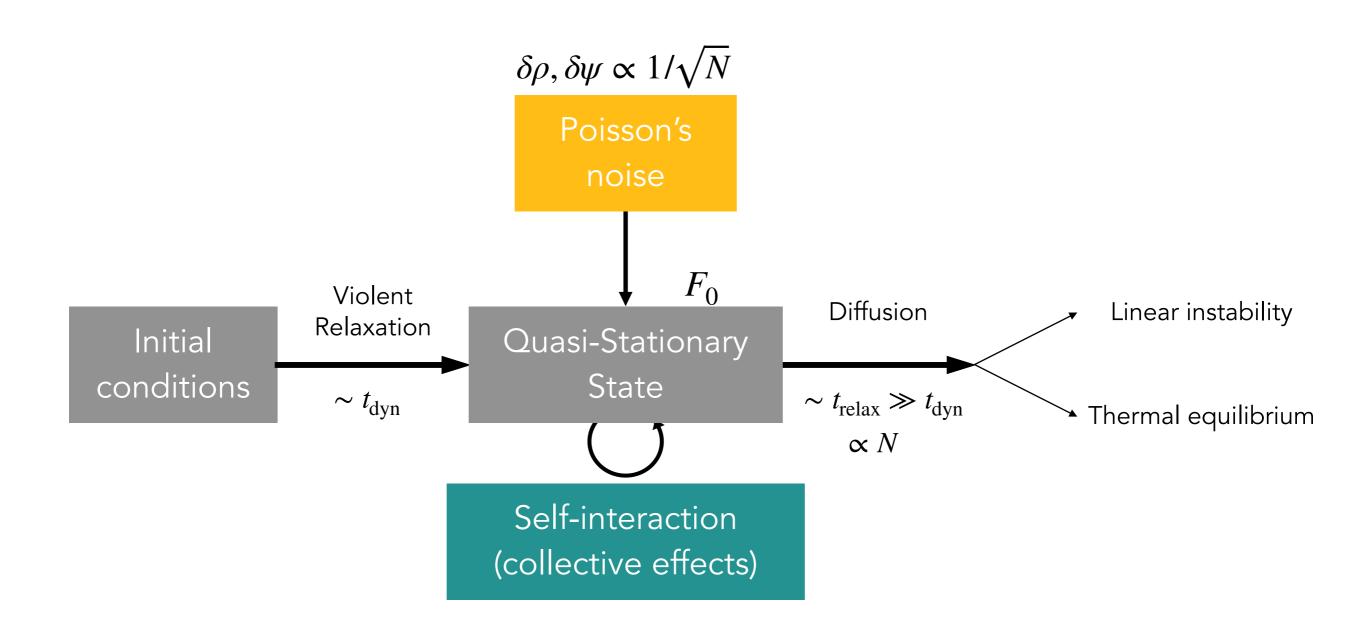






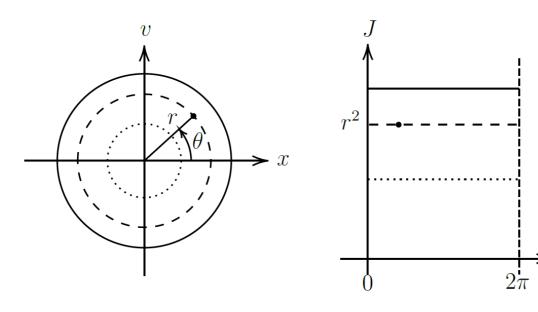






Key points:

► Angle-action variables



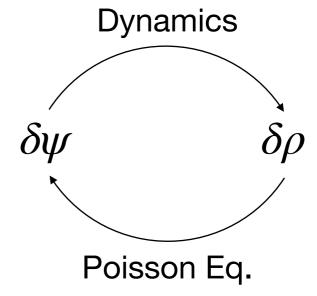
$$J = \text{cst}$$

$$\theta = \Omega(J) t + \theta_0$$

► Resonant coupling

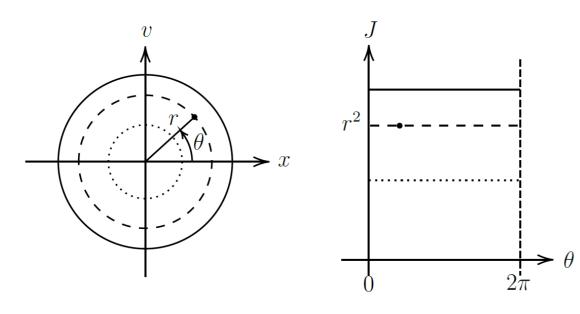
$$\delta_{\rm D}(n\Omega-n'\Omega')$$

► Collective effects



Key points:

► Angle-action variables



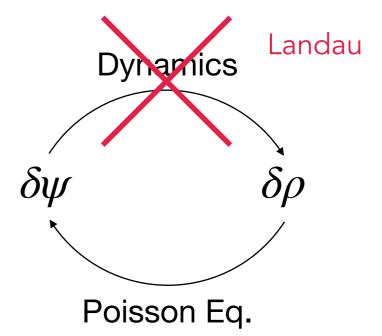
$$J = cst$$

$$\theta = \Omega(J) t + \theta_0$$

► Resonant coupling

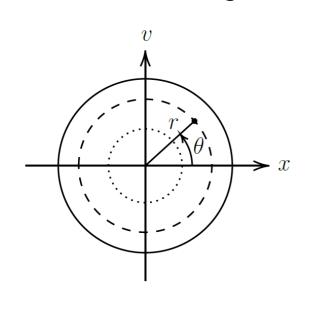
$$\delta_{\rm D}(n\Omega-n'\Omega')$$

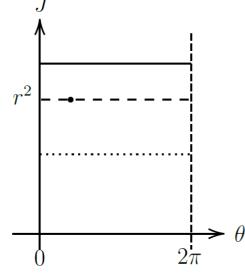
► Collective effects



Key points:

► Angle-action variables



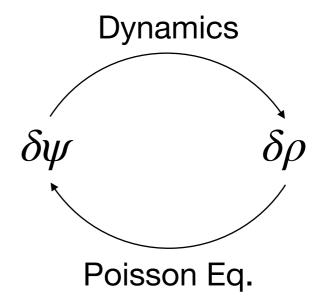


$$J = cst$$
$$\theta = \Omega(J) t + \theta_0$$

Resonant coupling

$$\delta_{\rm D}(n\Omega-n'\Omega')$$

► Collective effects



Hypotheses:

- $N \gg 1$
- ► Integrable system
- ► Resonances over encounters
- ► Time-scale separation

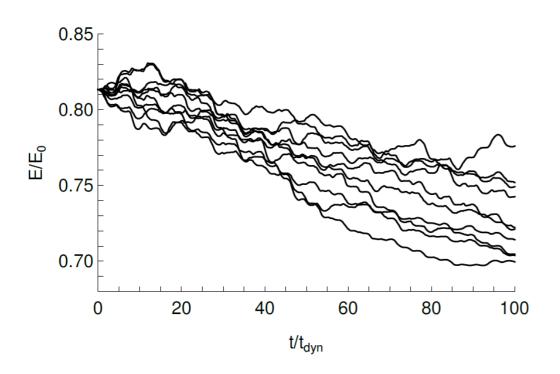
« Sufficiently » stable

► Uncorrelated Poisson noise

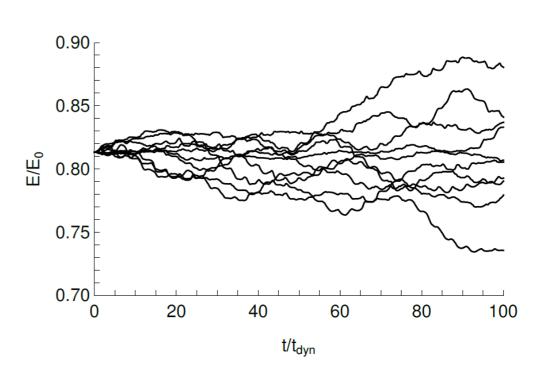
Fokker-Planck form

$$\frac{\partial F_0}{\partial t} = -\frac{\partial \mathcal{F}}{\partial J} = -\frac{\partial}{\partial J} \left[\mathbf{A}(J) F_0(J) - \frac{1}{2} \mathbf{D}(J) \frac{\partial F_0}{\partial J} \right]$$

Dynamical friction

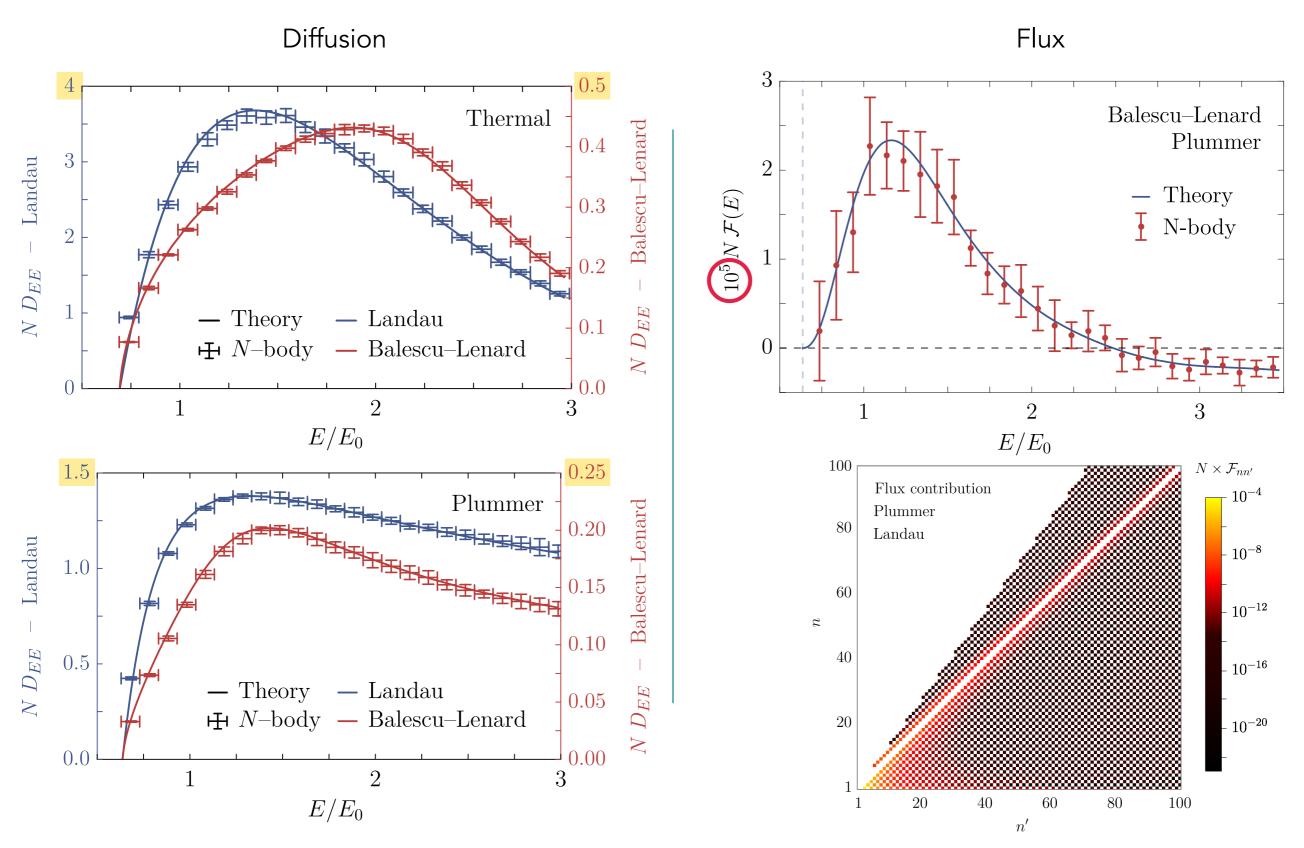


Diffusion



Prediction vs measurements

Roule+(2022)



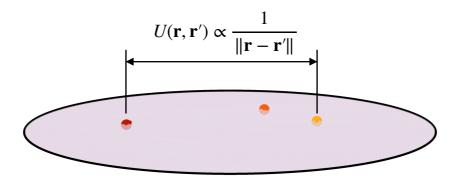
Conclusion on 1D self-gravitating system

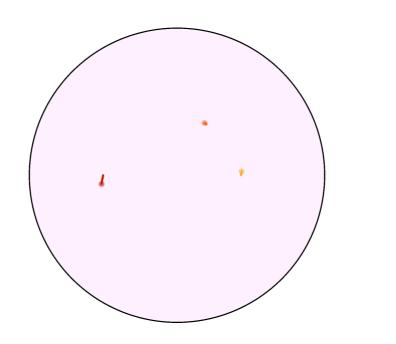
► Balescu-Lenard validation

- Quasi kinetic blocking
- Damping collective effects
- ► Resonances and collective effects matter

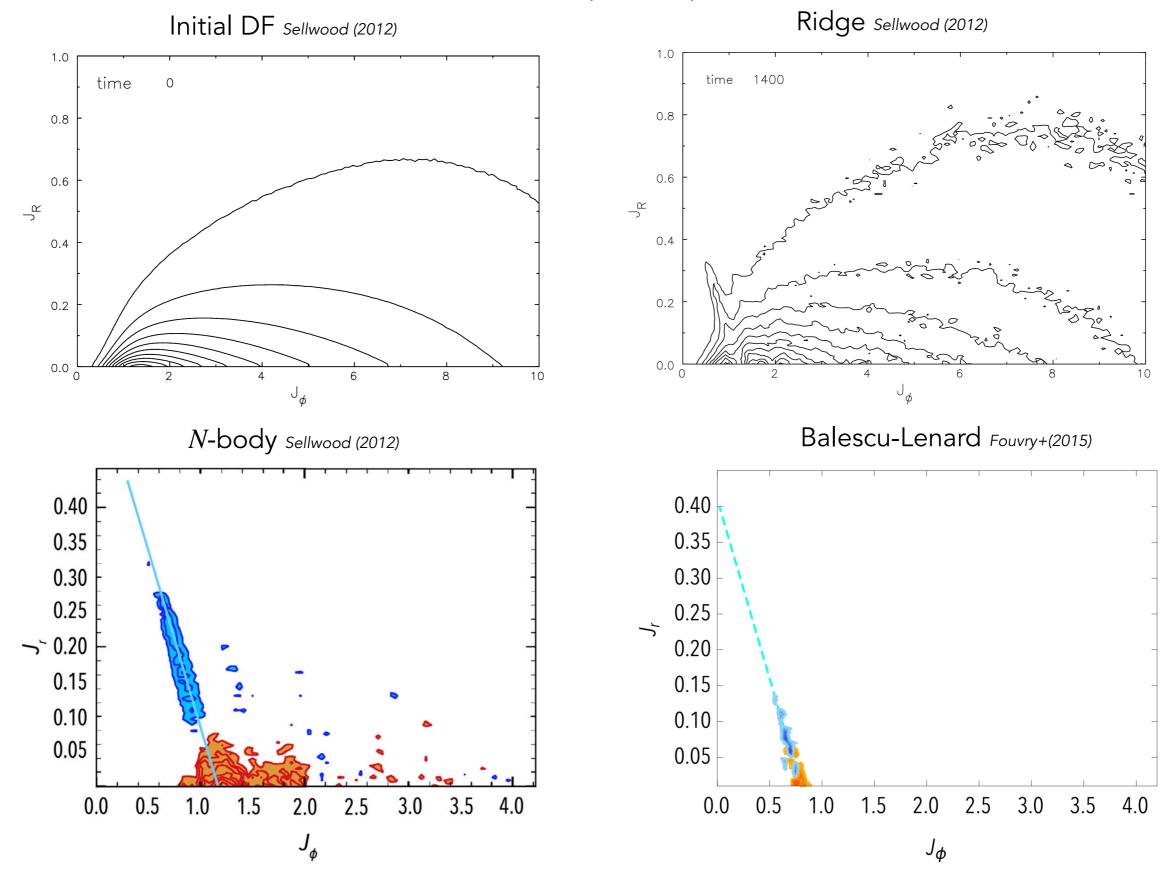
Razor-thin disc

Axisymmetric-potential motion

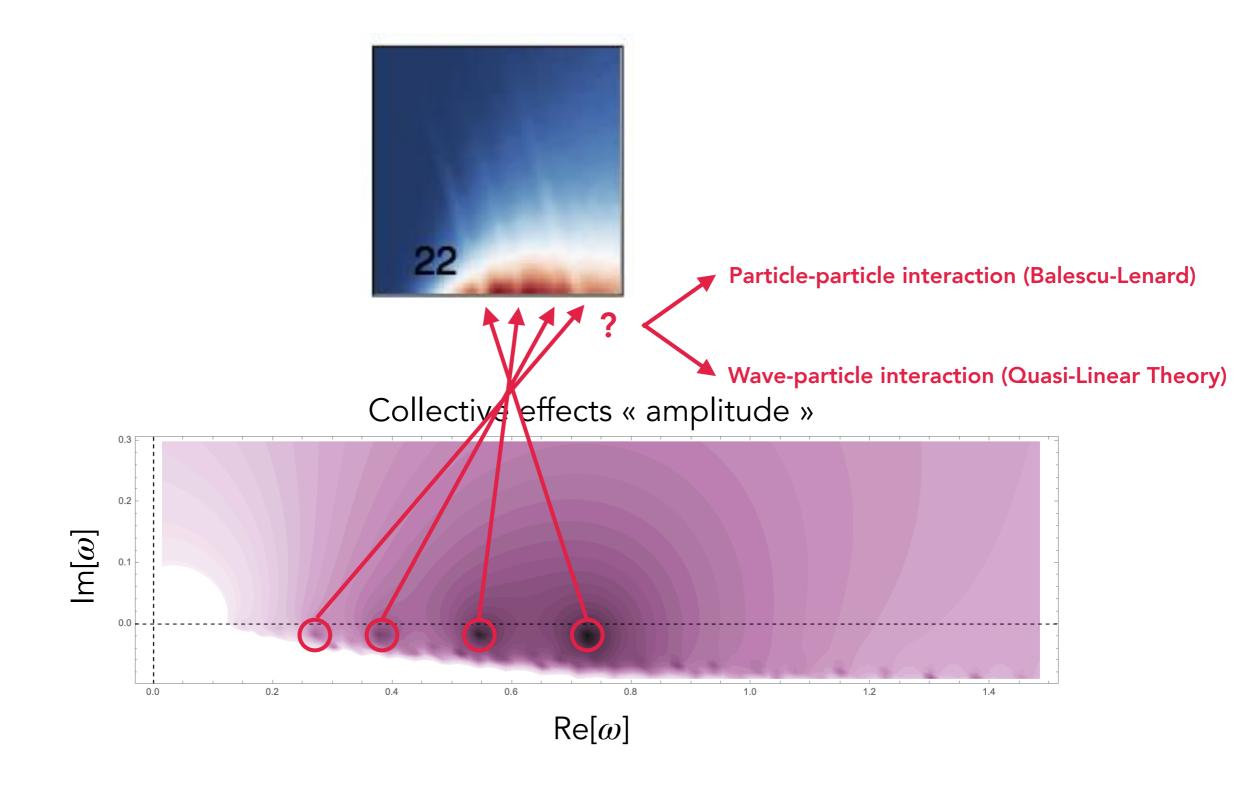




Sellwood (2012) disc

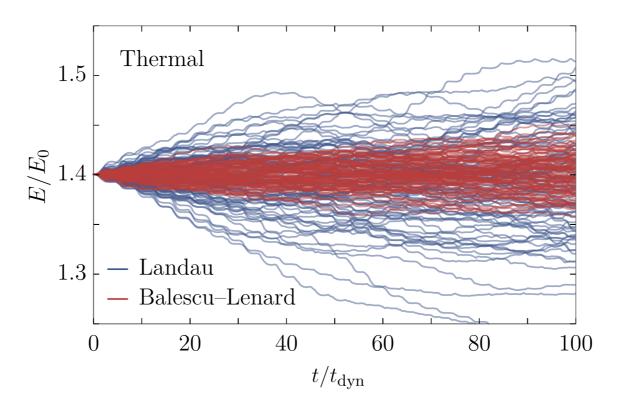


Sellwood (2012) disc

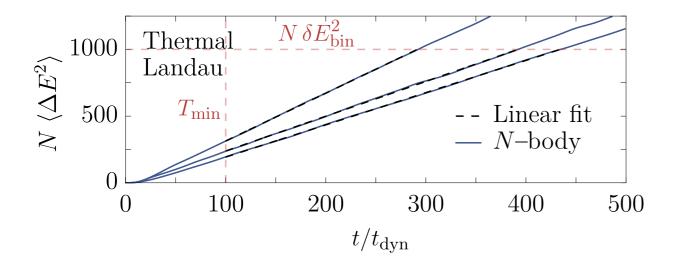


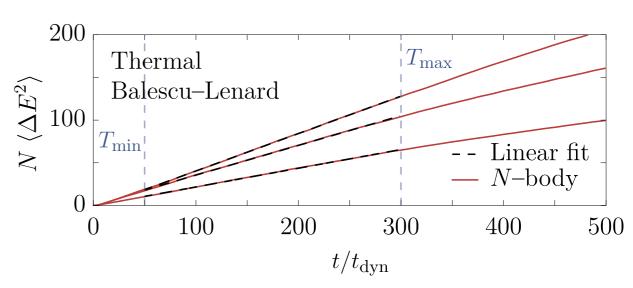
Back-up slides

N-body measurements



$$D(J) = \lim_{T \to +\infty} \frac{\left\langle \Delta J(T)^2 \right\rangle}{T}$$





$$\frac{\partial F_0}{\partial t} = 2\pi^2 m \frac{\partial}{\partial J} \sum_{n,n'} n \int dJ' \frac{\delta_D(n\Omega - n'\Omega')}{\delta_D(n\Omega - n'\Omega')} |\psi_{nn'}^d(J,J',n\Omega)|^2 \left(n \frac{\partial}{\partial J} - n' \frac{\partial}{\partial J'}\right) F_0(J) F_0(J')$$

 F_0 Mean-field orbit population in action space ${\it J}$

 $\sum_{n,n'}$ Scan over resonances $\int \! \mathrm{d}J'$ Scan over orbits

 $m \propto 1/N$ Individual mass of particles

 $\delta_{\rm D}(n\Omega-n'\Omega')$ Resonance condition

 $\frac{\partial}{\partial I}$ Divergence of a flux

 $|\psi_{nn'}^{\mathrm{d}}(J,J',n\Omega)|^2$ « Dressed » coupling

Coupling coefficients

$$\psi_{nn'}^{\mathbf{d}}(J,J',\omega) = -\sum_{p,q} \psi_{n}^{(p)}(J) \left[\mathbf{I} - \mathbf{M}(\omega) \right]_{pq}^{-1} \psi_{n'}^{(q)*}(J')$$

Bi-orthogonal basis

Fourier Transform in angle θ of bi-orthogonal basis elements $\left(\rho^{(p)}, \psi^{(p)}\right)$

$$\int dx \, \rho^{(p)}(x) \, \psi^{(q)*}(x) = -\delta_q^p$$

$$\psi^{(p)}(x) = \int dx' \, \rho^{(p)}(x') \, U(x, x')$$

Response Matrix

 $\mathbf{M}[F_0](\omega)$

$$\mathbf{M}_{pq}(\omega) = -2\pi \sum_{n} \int_{\mathcal{L}} dJ \, \frac{n \, \partial F_0 / \partial J}{n\Omega - \omega} \, \psi_n^{(p)*}(J) \, \psi_n^{(q)}(J)$$