Integrating out Astrophysical **Uncertainties in Direct** Detection Experiments Patrick Fox

‡ Fermilab

Based on work with G. Kribs, J Liu, T. Tait and N. Weiner

Dark Matter when Astrophysics Doesn't Matter

Patrick Fox

‡ Fermilab

Based on work with G. Kribs, J Liu, T. Tait and N. Weiner

- Direct detection and the standard assumptions
 Comparing multiple experiments, extracting g(v)
 CDMS recent results, some musings
- Conclusions

Direct Detection

$$\frac{dR}{dE_R} = \frac{N_T \rho_{\chi}}{m_{\chi}} \int_{v_{\min}}^{v_{\max}} d^3 \vec{v} f(\vec{v}(t)) \frac{d\sigma |\vec{v}|}{dE_R}$$

with,

$$\frac{d\sigma}{dE_R} = F_N^2(E_R)\frac{m_N}{\mu v^2}\bar{\sigma}(v, E_R)$$
$$\bar{\sigma}_i(v, E_R) = \begin{cases} \sigma_{i0} \\ \sigma_{i0}F_{\chi_i}^2(E_R) \\ \sigma_{i0}(v)F_{\chi_i}^2(E_R) \\ \sigma_{i0}(v, E_R) \end{cases}$$

Usual case (SI elastic WIMP)

$$\sigma_{N} = \frac{(Zf_{p} + (A - Z)f_{n})^{2}}{f_{p}^{2}} \frac{\mu_{N\chi}^{2}}{\mu_{n\chi}^{2}} \sigma_{p} \qquad v_{min} = \sqrt{\frac{m_{N}E_{R}}{2\mu_{N\chi}^{2}}}$$



Usual case (SI elastic WIMP)

$$\sigma_{N} = \frac{(Zf_{p} + (A - Z)f_{n})^{2}}{f_{p}^{2}} \frac{\mu_{N\chi}^{2}}{\mu_{n\chi}^{2}} \sigma_{p} \qquad v_{min} = \sqrt{\frac{m_{N}E_{R}}{2\mu_{N\chi}^{2}}}$$

Astrophysical inputs







Friday, 17 May 13



Direct Detection [see also, Drees and Shan, A. Peter, ...] Is a signal a measurement of particle physics or astrophysics?

The only way we have of probing our **local** DM distribution

 $f_1(v_{\min}(E_R)) = -\frac{4\mu^2 E_R^2}{m_N^2 E_R^2 - \mu^2 \delta^2} \frac{1}{\mathcal{N}\sigma_0(v_{\min}(E_R)) F_\chi^2(E_R)} \left(\frac{d\mathcal{R}}{dE_R} - \mathcal{R}\frac{1}{F_\chi^2(E_R)} \frac{dF_\chi^2(E_R)}{dE_R}\right)$

f-condition: $f(v) \ge 0$

(Deconvoluted) rate is a monotonically decreasing function, or there is non-standard particle physics e.g. inelastic or a increasing DM form factor **Direct Detection** [see also, Drees and Shan, A. Peter, ...] Is a signal a measurement of particle physics or astrophysics?

The only way we have of probing our **local** DM distribution

 $f_{1}(v_{\min}(E_{R})) = -\frac{4\mu^{2}E_{R}^{2}}{m_{N}^{2}E_{R}^{2} - \mu^{2}\delta^{2}} \frac{1}{\mathcal{N}\sigma_{0}(v_{\min}(E_{R}))F_{\chi}^{2}(E_{R})} \left(\frac{d\mathcal{R}}{dE_{R}} - \mathcal{R}\frac{1}{F_{\chi}^{2}(E_{R})}\frac{dF_{\chi}^{2}(E_{R})}{dE_{R}}\right)$ $f_{1}(v) = \int d\Omega f(\vec{v}).$

f-condition: $f(v) \ge 0$

(Deconvoluted) rate is a monotonically decreasing function, or there is non-standard particle physics e.g. inelastic or a increasing DM form factor **Direct Detection** [see also, Drees and Shan, A. Peter, ...] Is a signal a measurement of particle physics or astrophysics?

The only way we have of probing our **local** DM distribution

$$f_{1}(v_{\min}(E_{R})) = -\frac{4\mu^{2}E_{R}^{2}}{m_{N}^{2}E_{R}^{2} - \mu^{2}\delta^{2}} \frac{1}{\mathcal{N}\sigma_{0}(v_{\min}(E_{R}))F_{\chi}^{2}(E_{R})} \left(\frac{d\mathcal{R}}{dE_{R}} - \mathcal{R}\frac{1}{F_{\chi}^{2}(E_{R})} \frac{dF_{\chi}^{2}(E_{R})}{dE_{R}} \right)$$

$$f_{1}(v) = \int d\Omega f(\vec{v}). \qquad \mathcal{R} \equiv \frac{1}{F_{N}^{2}(E_{R})} \frac{dR}{dE_{R}}$$

f-condition: $f(v) \ge 0$

(Deconvoluted) rate is a monotonically decreasing function, or there is non-standard particle physics e.g. inelastic or a increasing DM form factor



Direct Detection without bias

$$\frac{dR}{dE_R} = \frac{N_T M_T \rho}{2m_\chi \mu^2} \int_{v_{min}}^{v_{max}} d^3 \vec{v} \frac{f(\vec{v}, \vec{v_E})}{v} \sigma(E_R)$$

$$\frac{g(v)}{v}$$

$$\frac{dR}{dE_R} = \frac{N_T M_T F_N^2(E_r)}{2\mu^2} \frac{\rho\sigma}{m_\chi} g(v)$$

$$v_{min} = \sqrt{\frac{M_T E_R}{2\mu^2}}$$

Recoil energy uniquely determines **minimum** DM velocity

Direct Detection without bias



$$v_{min} = \sqrt{\frac{M_T E_R}{2\mu^2}}$$

Recoil energy uniquely determines **minimum** DM velocity

Direct Detection without bias



$$v_{min} = \sqrt{\frac{M_T E_R}{2\mu^2}}$$

Recoil energy uniquely determines **minimum** DM velocity



Speed distribution is positive semidefinite $f(v) \geq 0$



Integral monotonically decreases



"Least" monotonic function is a step function $\Theta(v_1 - v_{\min})$







Two experiments allow us to test particle physics independent of astrophysics

- I) Make hypothesis about DM e.g. elastically scattering DM with mass 10 GeV and x-sec 10⁻⁴¹ cm²
- 2) Use experiment A to extract astrophysics i.e. rho x g(v)3) Use these extracted astrophysics properties to predict result at experiment B
- 4) Compare to B's measurement/bound
- 5) Rule in or out each particle physics hypothesis

Doesn't allow extraction of unique x-sec, mass Experiments must run over same part of year Other uncertainties (nuclear, atomic etc not addressed)

Comparing experiments

$$N_T = \kappa N_A m_p / M_T$$

Solve for g(v)

$$g(v_{min}) = \frac{2m_{\chi}\mu^2}{N_A\kappa m_p \rho \sigma(E_R)} \frac{dR_1}{dE_1}$$

$$\frac{dR_1}{dE_1} \iff g(v_{min}) \iff \frac{dR_2}{dE_2}$$

The master formula (SI):

$$C_T^{(i)} = \kappa^{(i)} \left(f_p \, Z^{(i)} + f_n \left(A^{(i)} - Z^{(i)} \right) \right)^2$$

$$\frac{dR_2}{dE_R} \left(E_2 \right) = \frac{C_T^{(2)}}{C_T^{(1)}} \frac{F_2^2(E_2)}{F_1^2 \left(\frac{\mu_1^2 M_T^{(2)}}{\mu_2^2 M_T^{(1)}} E_2 \right)} \frac{dR_1}{dE_R} \left(\frac{\mu_1^2 M_T^{(2)}}{\mu_2^2 M_T^{(1)}} E_2 \right)$$

Using vmin space



Using vmin space

Experiment I \longleftrightarrow Experiment 2 $[E_{low}^{(1)}, E_{low}^{(1)}] \iff [v_{min}^{low}, v_{min}^{high}] \iff [E_{low}^{(2)}, E_{high}^{(2)}]$

















Clearly some tension, but how much?

What do I have to believe in order to believe this result is DM?





Thanks to Peter Sorensen, see 1208.5046





3 events could be upward fluctuation, 90% lower limit: 1.7

CDMS expects 0.7 background events, so perhaps I event is background. Best case is the highest event Similarly, perhaps XENON is downward fluctuation from 5.3



3 events could be upward fluctuation, 90% lower limit: 1.7

CDMS expects 0.7 background events, so perhaps I event is background. Best case is the highest event Similarly, perhaps XENON is downward fluctuation from 5.3



Isospin dependent couplings

$$C_T^{(i)} = \kappa^{(i)} \left(f_p \, Z^{(i)} + f_n \left(A^{(i)} - Z^{(i)} \right) \right)^2$$

In going from Si to Xe typically get ~20 enhancement in With 3 events rate

$$\frac{f_n}{f_p} \approx -0.7$$

Suppression by ~170



Isospin dependent coupling see talk by Jason Kumar $\sigma^{(i)}$ (i) (i) $\sigma^{(i)}$ (i) $\sigma^{(i)}$ (j) $\gamma^{(i)}$

$$C_T^{(i)} = \kappa^{(i)} \left(f_p \, Z^{(i)} + f_n \left(A^{(i)} - Z^{(i)} \right) \right)^2$$

In going from Si to Xe typically get ~20 enhancement in With 3 events
rate

$$\frac{f_n}{f_p} \approx -0.7$$

Suppression by ~170



Isospin dependent couplings

Predicted rate at DAMA~amplitude of modulation DAMA ~100% modulated Consistent with event timing in CDMS

 $\frac{f_n}{f_n} \approx -0.7$

Rate too low to explain CoGeNT

Exothermic DM



Exothermic DM

$$v_{\min} = \left| \delta + \frac{m_{\mathrm{N}} E_{\mathrm{R}}}{\mu} \right| \frac{1}{\sqrt{2 E_{\mathrm{R}} m_{\mathrm{N}}}}$$



Exothermic DM



Astroindpendent rate at XENON100 integrating from 2keV

Astroindpendent rate at XENON100 integrating from 3keV



Friday, 17 May 13

Conclusions

•Should analyse data independent of astro uncertainties •With only one experiment results are limited, •Can extract f(v) by differentiating deconvoluted rate •With multiple experiments should compare g(v)•Under particle physics assumption can compare multiple experiments, test consistency •Ultimately find region of consistent parameter space Independent of all astrophysics inputs, CDMS-Si is at odds with XENON100, for simple elastic WIMP •XENON efficiency has to be a lot smaller, or nonstandard WIMP