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General 3D TFTs, BV and knot invariants

Chem-Schwinn theory (pert. theory)

$$Z = \int \mathcal{D}A e^{\frac{1}{\hbar} \int A \mathcal{D}A + \frac{2}{3} \int A^3}$$

AS $\rightsquigarrow Z = \int \mathcal{D}A e^{\frac{1}{\hbar} \int A^A \mathcal{D}A^3 \eta_{ABC} + \frac{2}{3} \int f_{ABC} A^A A^B A^C}$

$$\frac{k^\#}{\hbar} = \sum_n \frac{1}{\hbar} \sum_n^{\vee} V_n = \sum_n C_n \frac{1}{\hbar}$$

color inf. \swarrow
integrate over propagators \leftarrow

\nwarrow trivalent graphs

$$\Gamma = \bigcirc \quad C_n = f_{A_1 A_2 A_3} f_{B_1 B_2 B_3} \eta^{A_1 B_1} \eta^{A_2 B_2} \eta^{A_3 B_3}$$

$$b_n = \int \int$$

- graph complex — without ext legs
 — orient graphs
 — $(\Gamma, \partial) = (\Gamma, \omega)$

IHX relation

$$\int = \int - \int - \int$$

$$C_{\mathbb{Z}} = C_H - C_X$$

Jacobi ident.

$$\partial(\gamma)(\gamma) = \pm X$$

$$\sum_{\Gamma} c_{\Gamma} \Gamma = \text{graph cycle}$$

c_{Γ} weight system.

$$\sum_{\Gamma^*} b_{\Gamma} \Gamma^* \text{ graph cocycle}$$

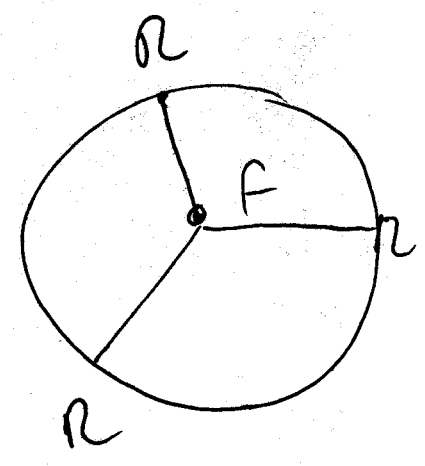
$$Z = \left\langle \sum_{\Gamma} c_{\Gamma} \Gamma, \sum_{\Gamma^*} b_{\Gamma} \Gamma^* \right\rangle$$

$$W_n(c) = \text{Tr}_R \text{Pexp}(SA)$$

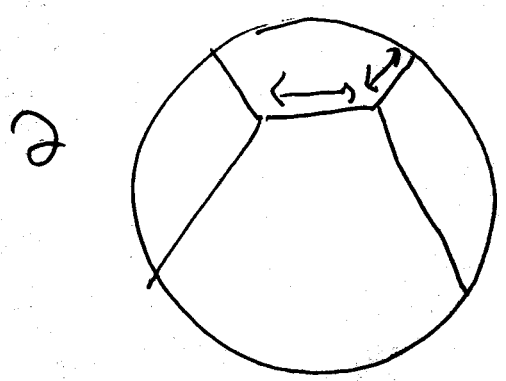
← Kontsevich integral

$$W_n(c) = \sum_{\Gamma} c_{\Gamma} b_{\Gamma}$$

↑
depend on \mathfrak{g}, R



$$c_{\Gamma} = R^{\alpha} A_{1\beta} R^{\beta} A_{2\gamma} R^{\gamma} A_{3\alpha} F A A \beta$$



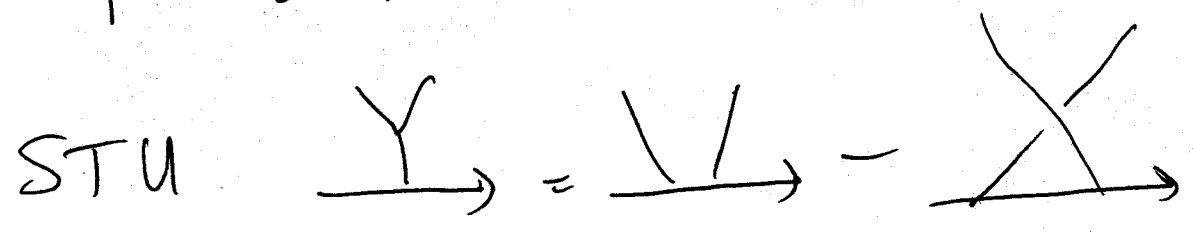
$$\sum_{\Gamma} c_{\Gamma} \Gamma \quad \text{graph cycle}$$

$$\sum_{\Gamma^*} b_{\Gamma} \Gamma^* \quad \text{graph cocycle}$$

$J_R(c)$

$$\langle W_R(c) \rangle = \sum h^n V_n \quad \leftarrow \begin{array}{l} \text{link invariant} \\ \text{(Vassiliev inv.)} \end{array}$$

$c_{\Gamma} = \text{weight system (for knots)}$



Physics: many 3D TFTs \rightsquigarrow L_{∞}

Math: all weight systems come from Lie alg.? (No)

gauge theories - BV formalism

$$\omega = \delta \Phi \wedge \delta \Phi^{\dagger} \quad \leftarrow -1$$

$$\{S, S\} = 0 \quad \text{classical master equation}$$

$$\int_{\mathcal{Z}} e^{-S}$$

AKSZ (195)

Lie algebra $\mathfrak{g}[1]$

$$\omega = \int da^A \wedge da^B \eta_{AB}$$

↑
Symplectic structure
1 degree 2

$$\mathbb{H} = f_{ABC} a^A a^B a^C \quad \text{deg } 3$$

$$\{ \mathbb{H}, \mathbb{H} \} = 0 \iff \text{Jacobi identity}$$

$$(\mathfrak{g}[1], \omega, \mathbb{H})$$

$$T[1] \mathcal{Z}_3 \rightarrow \mathfrak{g}[1]$$

$$\theta \in \mathfrak{g}$$

$$d\mathfrak{g}$$

$$A^A(\xi, \theta)$$

$$\omega = \int d^3\theta d^3\xi S A^A \wedge S A^B \eta_{AB} \quad \leftarrow -1$$

$$S = d^3\theta d^3\xi (A^A D A^B \eta_{AB} + \frac{2}{3} f_{ABC} A^A A^B A^C) \quad \leftarrow \mathbb{H}$$

$$D = \theta^{\alpha} \frac{\partial}{\partial \xi^{\alpha}}$$

$$\left(\mathcal{M}, \omega, \textcircled{H} \right) \quad \left. \begin{array}{l} \{ \textcircled{H}, \textcircled{H} \} = 0 \end{array} \right\} \text{Courant algebroid}$$

$$S = \int d^3 \theta d^3 \xi \left(\Phi^A D \Phi^B \omega_{AB} + \textcircled{H}(\Phi) \right)$$

$W_R(c)$

$$R_A R_B - R_B R_A = f_{AB}{}^C R_C \quad \text{Lie alg.}$$

$$Q = f_{AB}{}^C a^A a^B \frac{\partial}{\partial a^C} \quad Q^2 = 0$$

$$Q = f_{AB}{}^C a^A a^B \frac{\partial}{\partial a^C} + 2 \underbrace{R_{\alpha\beta}^{\gamma}}_{R_{\beta}^{\alpha}(\alpha)} a^{\alpha} \xi^{\beta} \frac{\partial}{\partial \xi^{\gamma}} \quad Q^2 = 0.$$

$(\mathcal{M}, \omega, \textcircled{H})$

$Q = \text{Hamiltonian vector field for } \textcircled{H}$

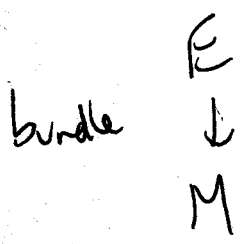
$$\boxed{Q \cdot R = -R R}$$

$T[M]$

$$Q = V^{\mu} \frac{\partial}{\partial x^{\mu}}$$

$$Q = V^{\mu} \frac{\partial}{\partial x^{\mu}} + \Gamma_{\mu}^{\alpha} (x) V^{\mu} \xi^{\alpha} \frac{\partial}{\partial \xi^{\alpha}}$$

↑
connection on E



$$Q^2 = 0 \Leftrightarrow d\Gamma + [\Gamma, \Gamma] = 0$$

M ← complex manifold

$T^{\mathbb{C}}[M]$

$$Q = V^{\bar{\alpha}} \frac{\partial}{\partial x^{\bar{\alpha}}}$$

⇒ representation: holomorphic vector bundles over M .

(M, ω, \mathbb{C})

$R \quad E \rightarrow M$

$$QR = -RQ$$

$$W_{\text{P}}(k) = \text{Tr} P \exp \int_C dt d\theta R$$



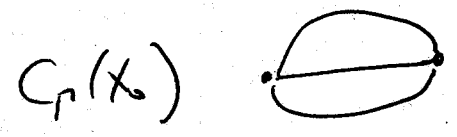
BV-invariant

$$Z = \int \mathcal{D}\Phi e^{\int \Phi^A D\Phi^B \omega_{AB} + (H)(\Phi)}$$

$\mathcal{X}_0 \text{ FM}$

$$\Rightarrow \partial_{A_1} \dots \partial_{A_k} (H)(\mathcal{X}_0)$$

$\sum_n C_n(\mathcal{X}_0) \hookrightarrow H^k(\mathcal{M})$
 function

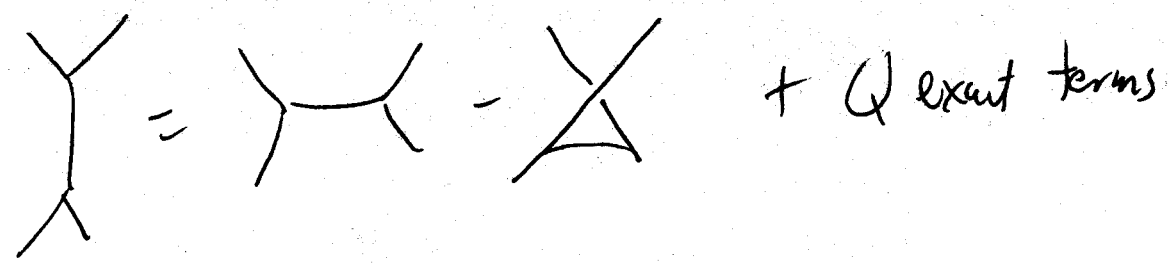


$$\partial_{A_1} \partial_{A_2} \partial_{A_3} (H)(\mathcal{X}_0) \partial_{B_1} \partial_{B_2} \partial_{B_3} \dots$$

$$= (H)(\mathcal{X}_0) \omega^{A_1 B_1} \omega^{A_2 B_2} \omega^{A_3 B_3} \dots$$

\hookrightarrow the same as in CS theory

\mathbb{IHX}

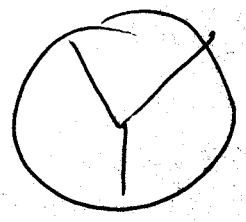
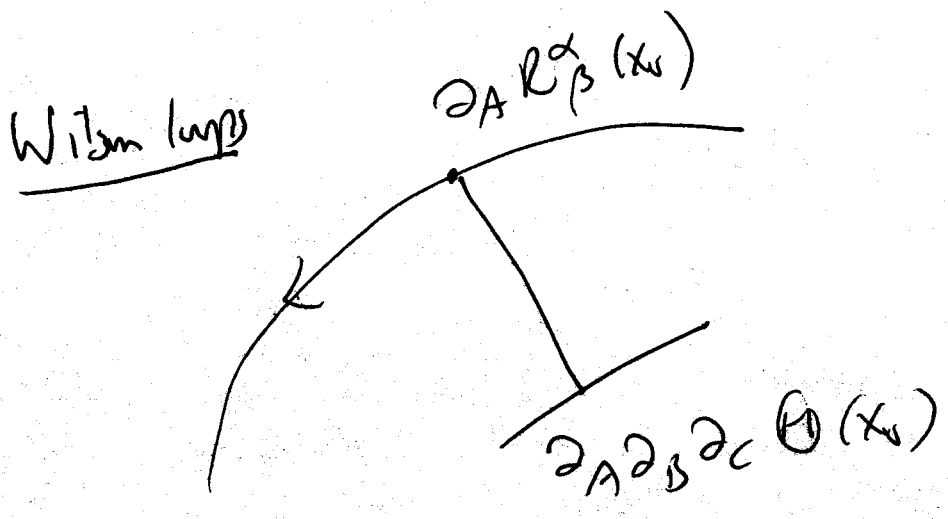


$$Q^A \partial_A = \omega^{AB} \partial_B (H)(\mathcal{X}_0)$$

$$\sum_n C_n(\mathcal{X}_0) \Gamma^n - \text{cycle in } H^k(\mathcal{M})$$

$$\int_2 \mathcal{D}\Phi e^{\int \mathcal{D}\Phi^B \omega_{AB} + (H)(\Phi)} \text{Tr Per} \int N$$

TUC



$$\sum_n C_p(x_0) b_\Gamma$$

↑
ω in Chem-Simons

$C_p(x_0)$ = new weight system with values in $H_Q(M)$

Let M be hyperkähler (or holosympl manifold)

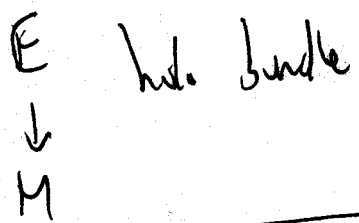
$$\omega = T^*[\omega] T^*[OM]$$

$$\omega = dx^{\mu} + dp_{\mu} + dg_{\mu} + dr^{\mu} + \Omega_{ij} dx^i dx^j$$

← degree 2

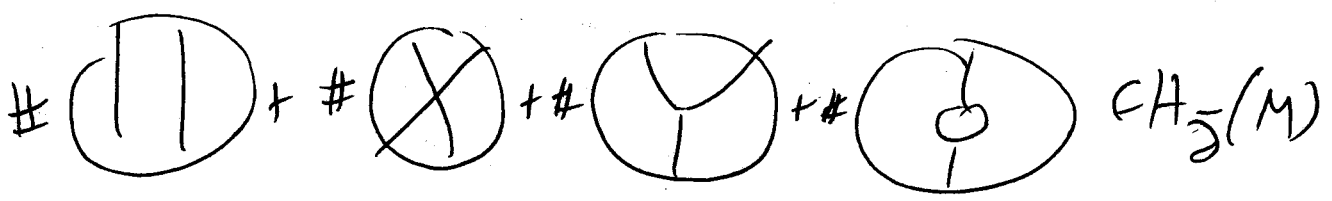
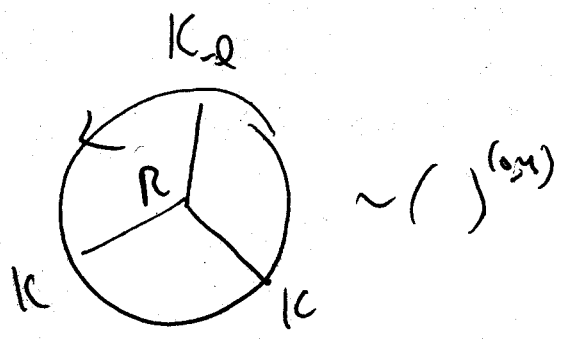
$$\textcircled{M} = P_{\bar{a}} v^{\bar{a}} \iff \bar{J}$$

$$S = \int d^3\theta d^3\xi \left(P_{\mu} D X^{\mu} + v^{\mu} D g_{\mu} + \Omega_{ij} X^i D X^j + P_{\bar{a}} v^{\bar{a}} \right)$$



$$R_{\bar{a} l_1 l_2 l_3} = R_{\bar{a} l_1 l_2 l_3} \Omega_{nlr}$$

$$\nabla_{\bar{a}} (K_{\bar{a}})^I_J \equiv K_{\bar{a} l}^I_J$$



Tr Bepn (g) K_{a} / H^+