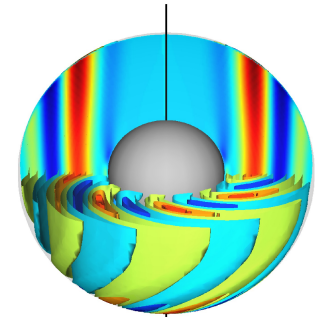


*Dynamo Theory*  
*KITP, Santa Barbara*  
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# ***Bistability and hysteresis of dipolar dynamos generated by chaotic convection in rotating spherical shells***

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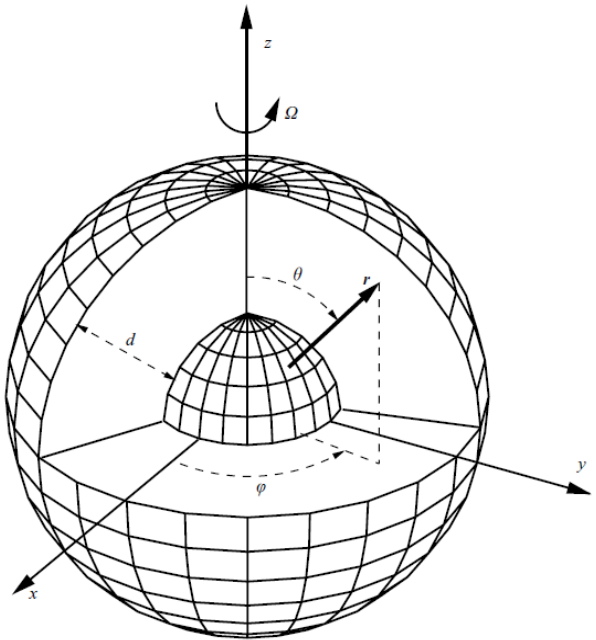
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of Glasgow**



# Convective spherical shell dynamos



## Model equations & parameters

*Boussinesq approximation*

$$\nabla \cdot \mathbf{u} = 0, \quad \nabla \cdot \mathbf{B} = 0,$$

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} =$$

$$-\nabla \pi - \tau \mathbf{k} \times \mathbf{u} + \Theta \mathbf{r} + \nabla^2 \mathbf{u} + \mathbf{B} \cdot \nabla \mathbf{B},$$

$$P (\partial_t \Theta + \mathbf{u} \cdot \nabla \Theta) = R \mathbf{r} \cdot \mathbf{u} + \nabla^2 \Theta,$$

$$P_m (\partial_t \mathbf{B} + \mathbf{u} \cdot \nabla \mathbf{B}) = P_m \mathbf{B} \cdot \nabla \mathbf{u} + \nabla^2 \mathbf{B}.$$

$$R = \frac{\alpha \gamma \beta d^6}{\nu \kappa}, \quad \tau = \frac{2\Omega d^2}{\nu}, \quad P = \frac{\nu}{\kappa}, \quad P_m = \frac{\nu}{\lambda}$$

## Basic state & scaling

$$T_S = T_0 - \beta d^2 r^2 / 2$$

$$\mathbf{g} = -d\gamma \mathbf{r}$$

Length scale:  $d$

Time scale:  $d^2 / \nu$

Temp. scale:  $\nu^2 / \gamma \alpha d^4$

Magn. flux density:  $\nu (\mu \rho)^{1/2} / d$

## Boundary Conditions

$$\mathbf{r} \cdot \mathbf{u} = \mathbf{r} \cdot \nabla \mathbf{r} \times \mathbf{u} / r^2 = 0,$$

$$\hat{\mathbf{e}}_r \cdot \mathbf{B}_{\text{int}} = \hat{\mathbf{e}}_r \cdot \mathbf{B}_{\text{ext}},$$

$$\hat{\mathbf{e}}_r \times \mathbf{B}_{\text{int}} = \hat{\mathbf{e}}_r \times \mathbf{B}_{\text{ext}},$$

$$\Theta = 0, \quad \text{at } r = r_i \equiv 2/3 \text{ and } r_o \equiv 5/3$$

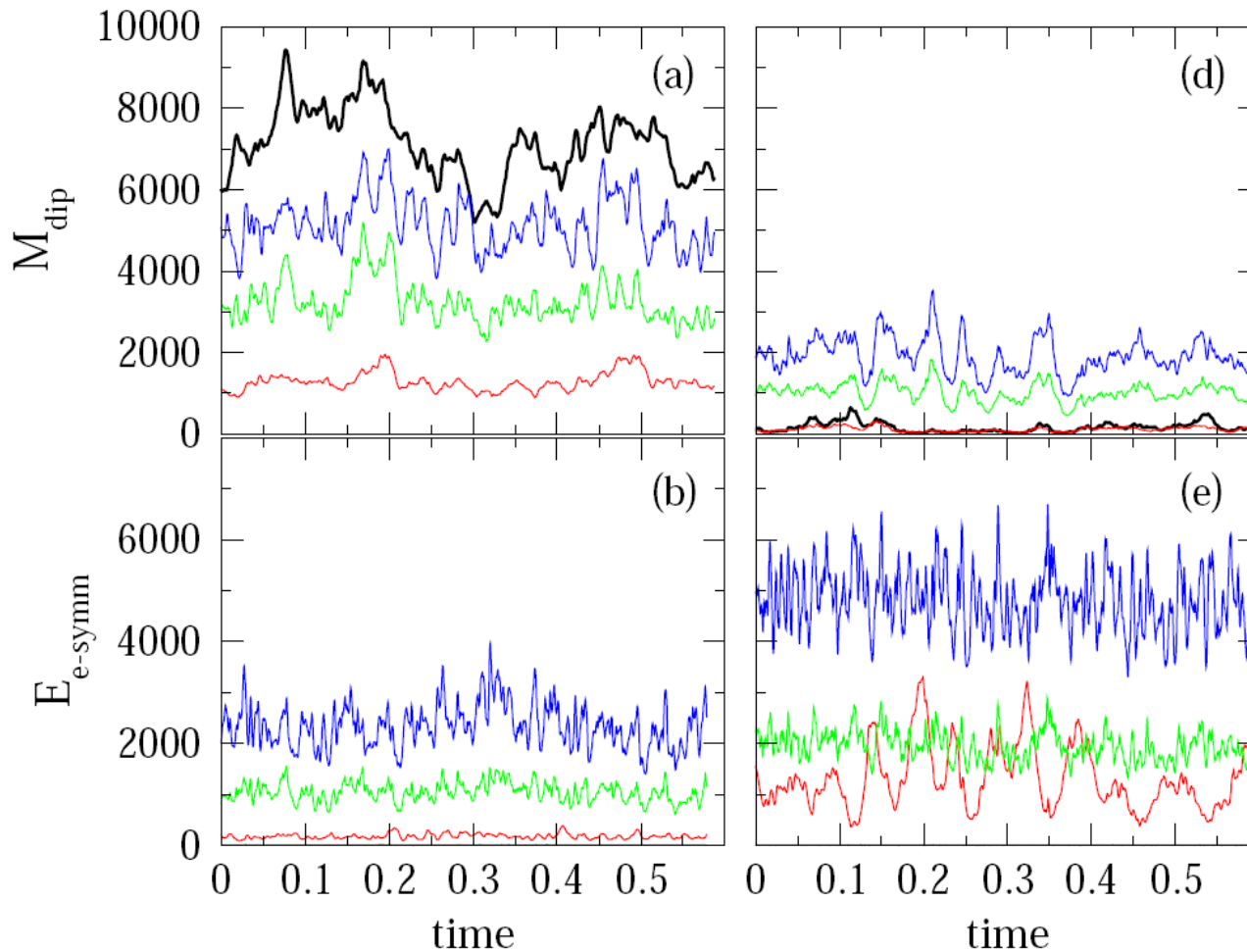
Simitev & Busse, JFM, 2004

# Two types of dipolar dynamos generated by chaotic convection

Energy densities

Mean Dipolar (MD)

Fluctuating Dipolar (FD)



- **Fully chaotic (large-scale turbulent) regime.**
- **Two chaotic attractors for the same parameter values.**
- **Essential qualitative difference: contribution of the mean poloidal dipolar energy**

	(ab)	(de)
Rm	133.6	196.5
Mdip/Mtot	0.803	0.527

black.....mean poloidal  
green.....fluctuating poloidal  
red.....mean toroidal  
blue.....fluctuating toroidal

$$R = 3.5 \cdot 10^6, \tau = 3 \cdot 10^4, P = 0.75 \text{ and } P_m = 1.5$$

# Regions and transition

Ratio of fluctuating to mean poloidal magn energy

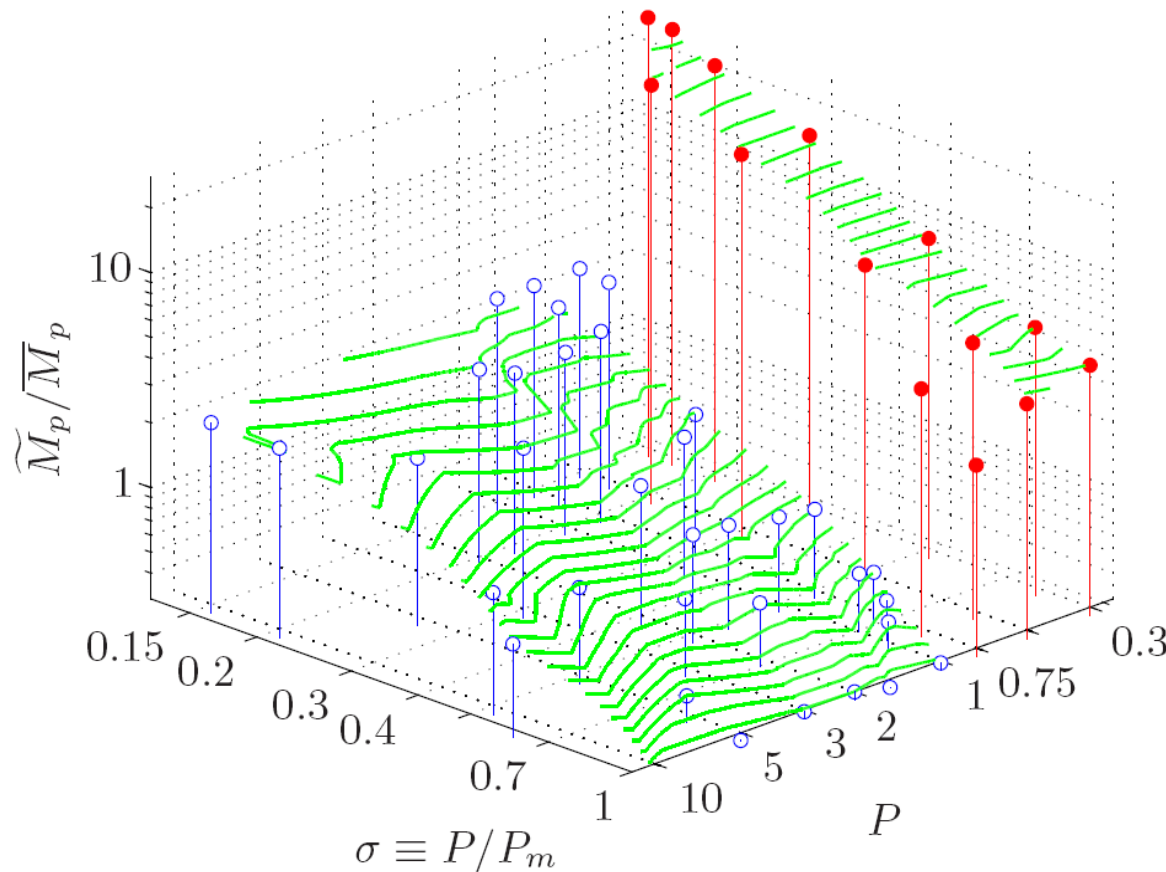
Two types of dipolar dynamos

⊕ **Mean Dipolar (MD)**

$$\widetilde{M}_p < \overline{M}_p$$

● **Fluctuating Dipolar (FD)**

$$\widetilde{M}_p > \overline{M}_p$$



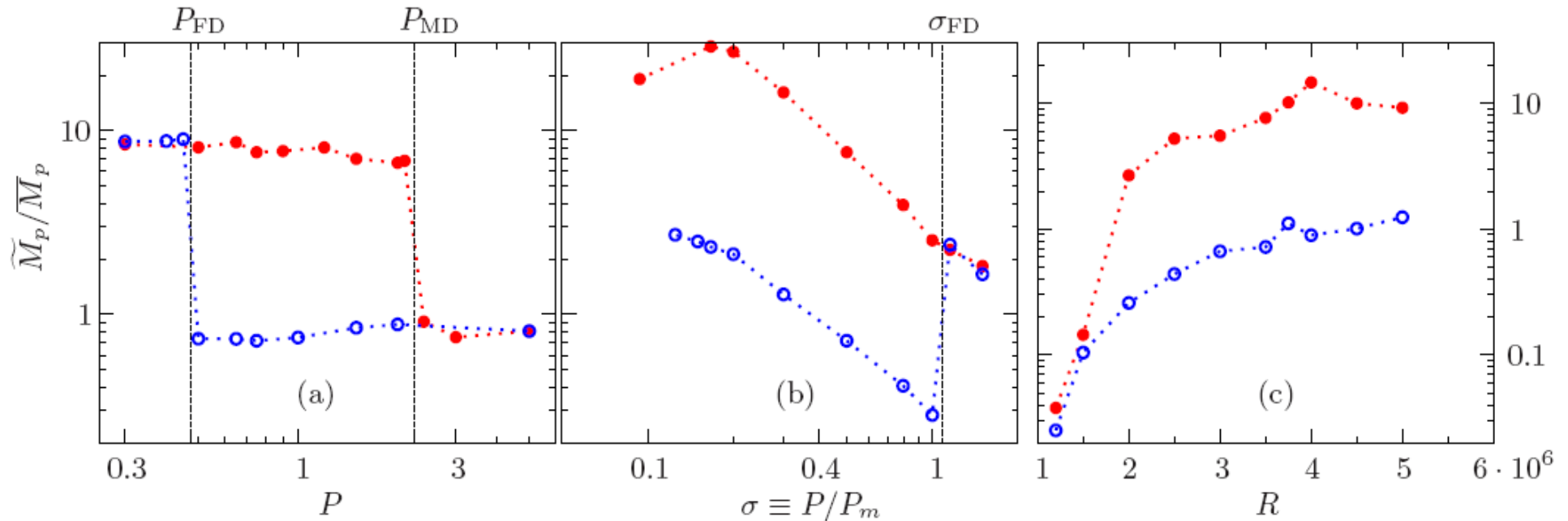
- MD and FD dynamos correspond to rather **different chaotic attractors** in a fully chaotic system
- The transition between them is not gradual but is an **abrupt jump** as a critical parameter value is surpassed.
- The nature of the transition is complicated.

$$R = 3.5 \cdot 10^6, \tau = 3 \cdot 10^4$$

	MD	FD
Mdip/Mtot	(0.62,1)	(0.41,56)

# Bistability and hysteresis in the MD $\Leftrightarrow$ FD transition

Bistability and hysteresis in the ratio of fluctuating poloidal to mean poloidal magn energy



(a)  $R = 3.5 \cdot 10^6$   $P/P_m = 0.5$

(b)  $R = 3.5 \cdot 10^6$ ,  $P = 0.75$

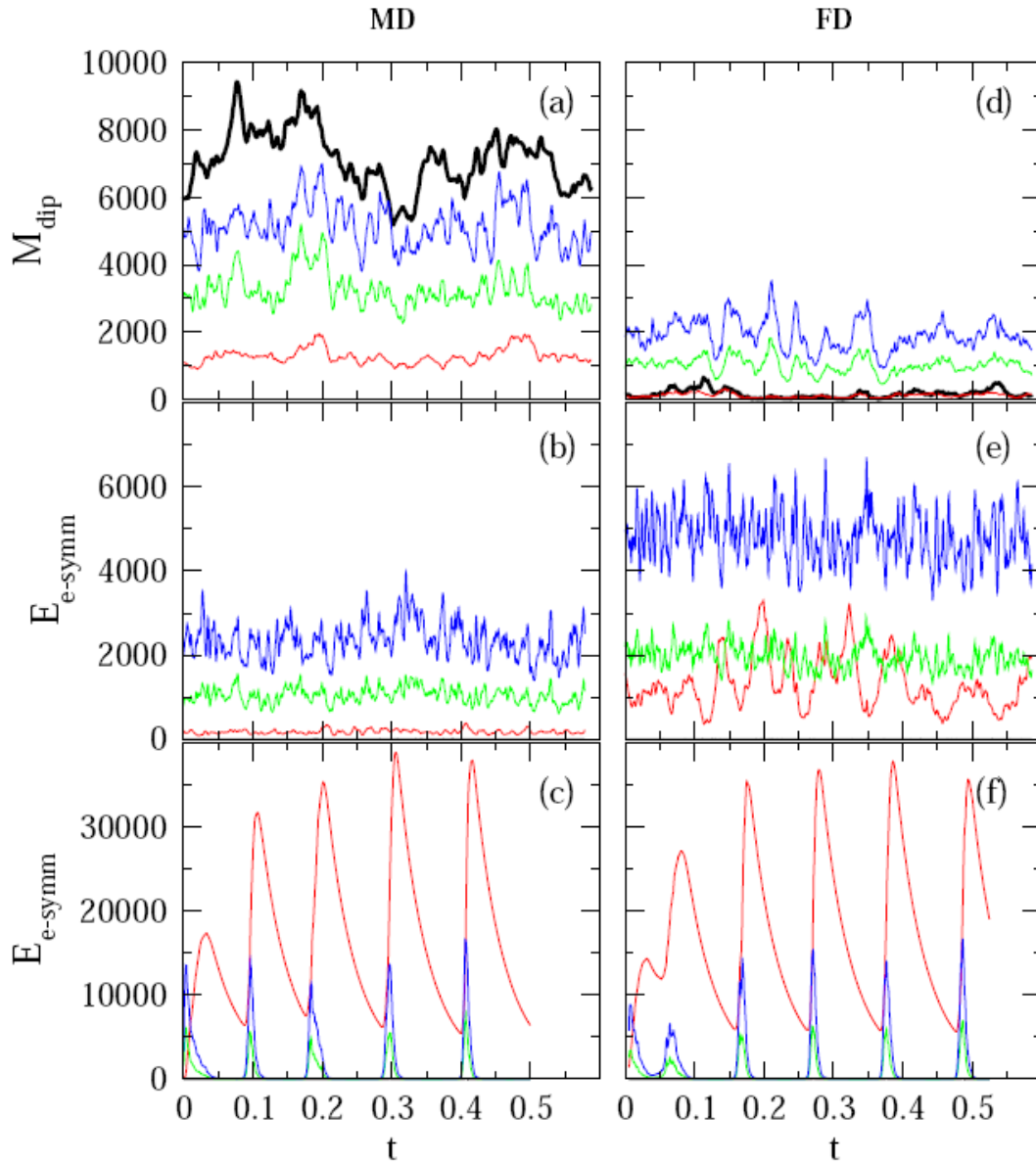
(c)  $P = 0.75$ ,  $P_m = 1.5$

in all cases:  $\tau = 3 \cdot 10^4$

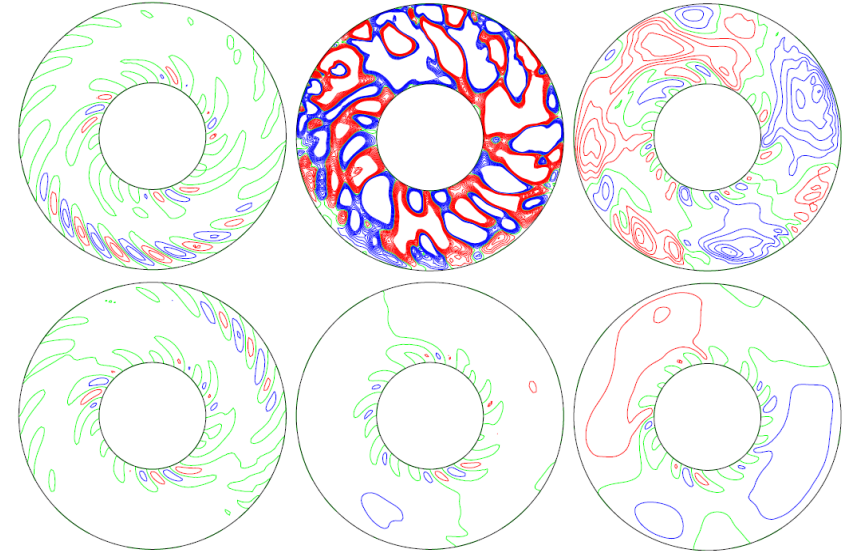
The coexistence is **not an isolated phenomenon** but can be traced with variation of the parameters.

$P_{MD} = 2.2$	$P_{FD} = 0.5$
$\sigma_{MD} = 0.07$	$\sigma_{FD} = 1$

# The hysteresis is a purely magnetic effect



After magnetic field is suppressed both MD and FD dynamos equilibrate to statistically **identical convective states** (period of relaxation oscillations, clockwise)

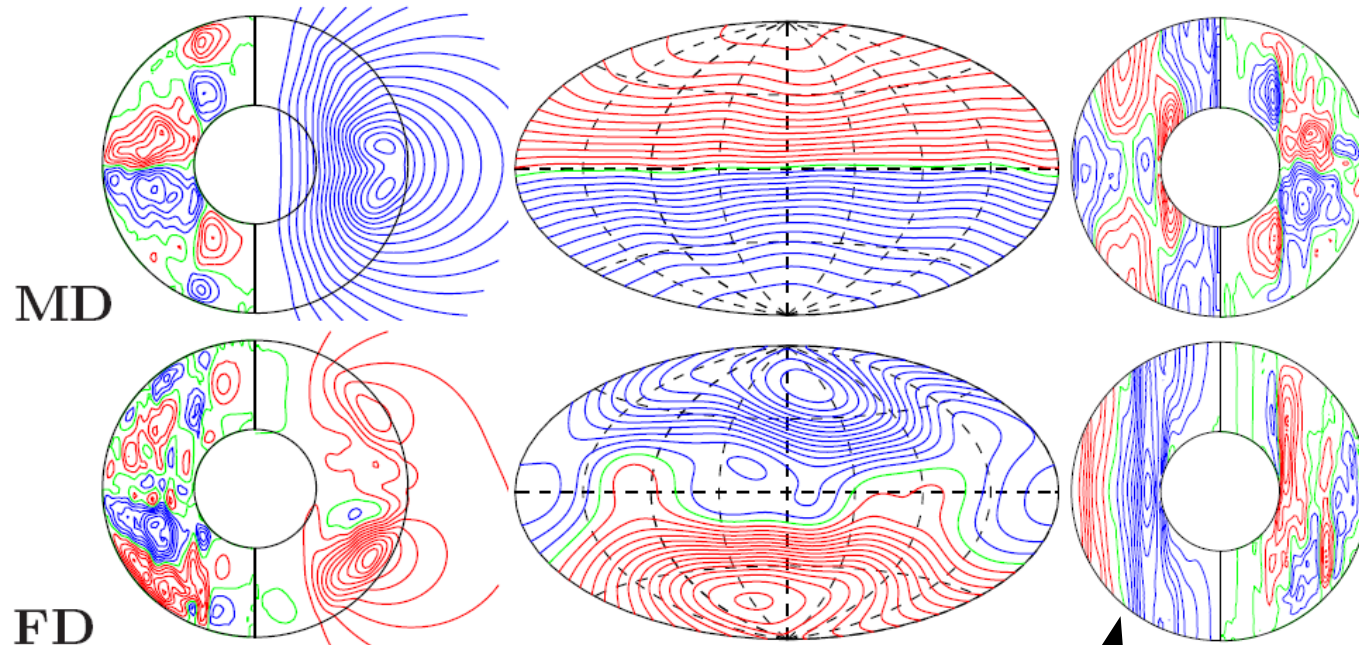


Magnetic field is artificially suppressed, i.e. **non-magnetic convection**

$$R = 3.5 \cdot 10^6, \tau = 3 \cdot 10^4 \quad P = 0.75 \text{ and } P_m = 1.5$$



# A property comparison of MD and FD dynamos (Spatial structures)



MD

FD

Snapshots of  
spatial structures

$$R = 3.5 \times 10^6, \tau = 3 \times 10^4$$

$$P = 0.75 \text{ and } P_m = 1.5$$

The stronger magnetic field of  
MD dynamos counteracts differential  
rotation (diff rotation, meridional streamlines)

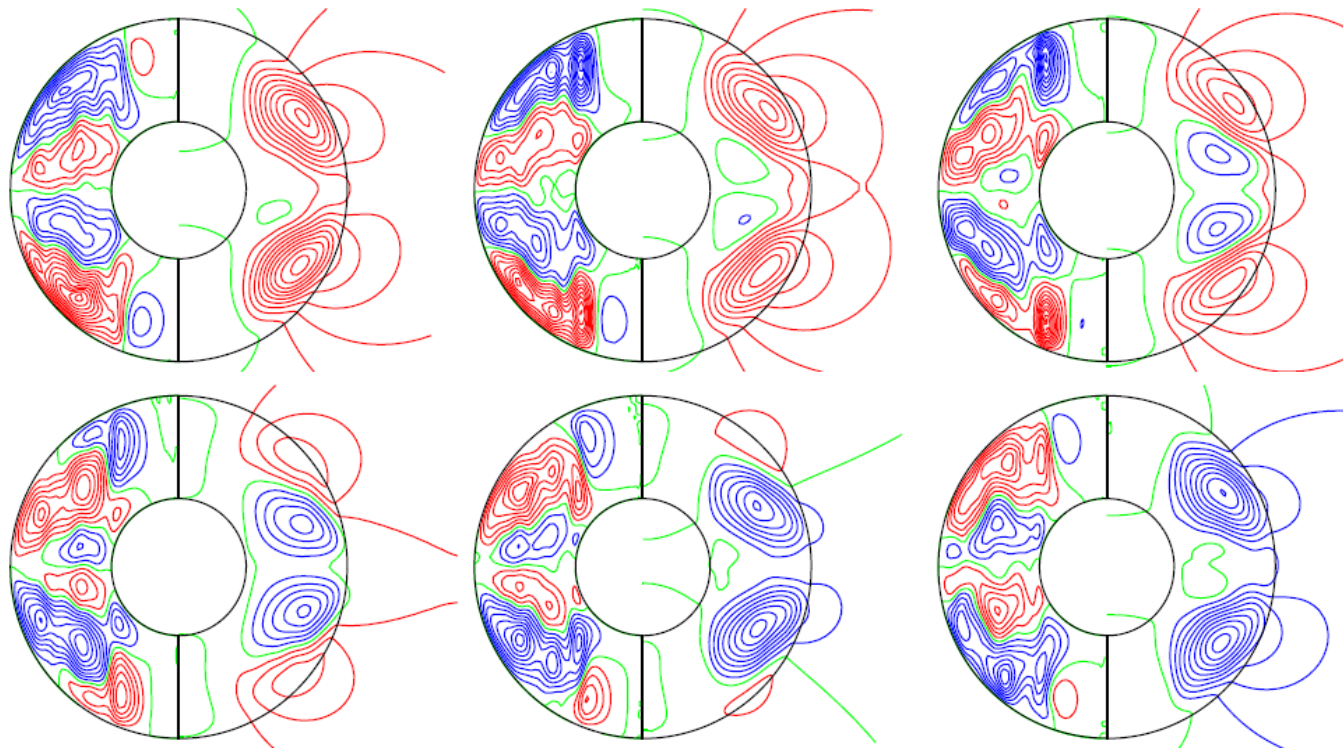
Both MD and FD dynamos appear dipolar from the outside (radial  
magn field at  $r=r_0+1.3$ , Earth's surface)

FD dynamos have a somewhat more irregular and small-scale internal structure  
( $B_{\phi}$  and meridional fieldlines)

# *A property comparison of MD and FD dynamos (Temporal variations)*

- **Mean Dipolar (MD) dynamos are non-oscillatory.**
- **Fluctuating Dipolar (FD) dynamos are oscillatory.**

*Half-period of oscillation in a FD dynamo (row-by-row)*



$$R = 3.5 \cdot 10^6, \tau = 3 \cdot 10^4, P = 0.75 \text{ and } P_m = 0.65$$



# Conclusion

- *Two types of dipolar dynamos can be distinguished:*
  - \* *Mean dipolar dynamos (MD)*
  - \* *Fluctuating dipolar dynamos (FD)*
- *MD and FD dynamos have rather different properties.*
- *FD dynamos are normally oscillatory. In some cases this may lead to reversals.*
- *The transition between MD and FD dynamos is hysteretic.*
- *Most geodynamo simulations have typical parameters values  **$P$  in  $[0.5, 2]$ ,  $Pm$  in  $[0.5, 10]$  and  $R < 10 R_c$  which are within the observed hysteresis region.***

*Thank you!*