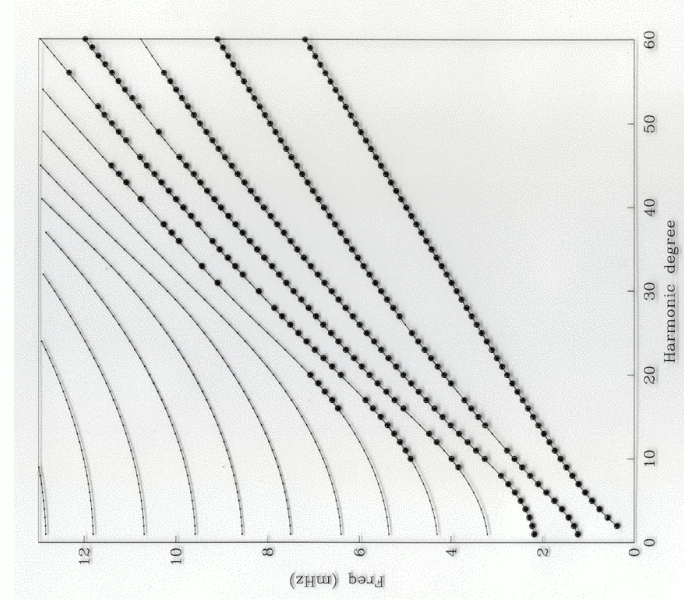
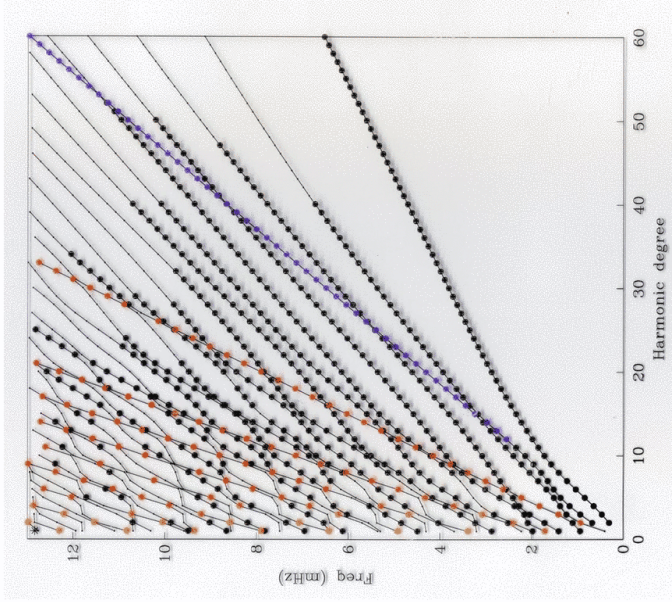


Surface waves

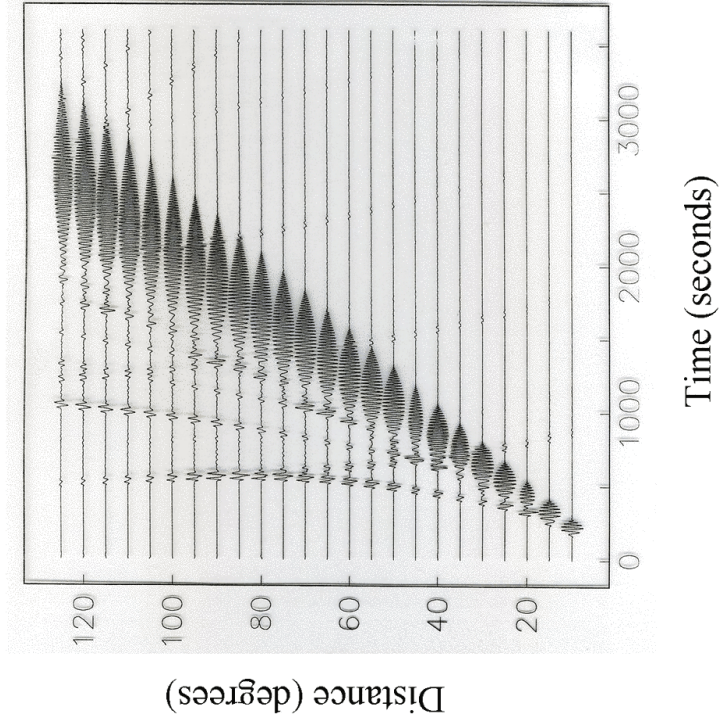
Guy Masters

Cider 2006





Synthetic seismograms



Standing waves and travelling waves

- Seismogram as a mode sum:

$$s(t) = \sum_k A_k \cos(\omega_k t + \phi_k) e^{-\alpha_k t}$$

In epicentral coordinates, A_k includes both the source excitation and the geometrical mode behavior as a function of epicentral distance, Δ . This latter term is proportional to a spherical harmonic Y_l^m . For surface waves, l is large and m is small, then

$$Y_l^m \simeq \frac{1}{\pi \sqrt{\sin \Delta}} \cos \left[\left(l + \frac{1}{2} \right) \Delta - \frac{\pi}{4} + \frac{m\pi}{2} \right] e^{im\phi}$$

Jeans formula is $ka = l + \frac{1}{2}$ so with arc distance given by $x = a\Delta$,

$$Y_l^m \simeq \frac{1}{\pi \sqrt{\sin \Delta}} \cos \left[kx - \frac{\pi}{4} + \frac{m\pi}{2} \right] e^{im\phi}$$

Seismogram of a single mode becomes proportional to a plane wave:

$$s(t) \propto e^{i(\omega t - kx)}$$

Phase and group velocity

In general $k = k(\omega)$ (dispersion). Mode sum becomes:

$$s(x, t) = \int B(\omega) e^{i(\omega t - k(\omega)x)} d\omega \quad c(\omega) = \frac{\omega}{k(\omega)}$$

$B(\omega)$ due to source – slowly varying. Consider phase, f ,

$$f = \omega t - k(\omega)x$$

Main contribution to integral when phase is stationary ($df/d\omega = 0$). If ω_s is the point where the phase is stationary occurs, we have:

$$t - \frac{dk}{d\omega}(\omega_s)x = 0$$

$$U = \frac{x}{t} = \frac{d\omega}{dk}$$

The energy associated with a particular frequency group centered on ω_s travels with the “group velocity”, $U(\omega_s)$.

Phase or group velocity?

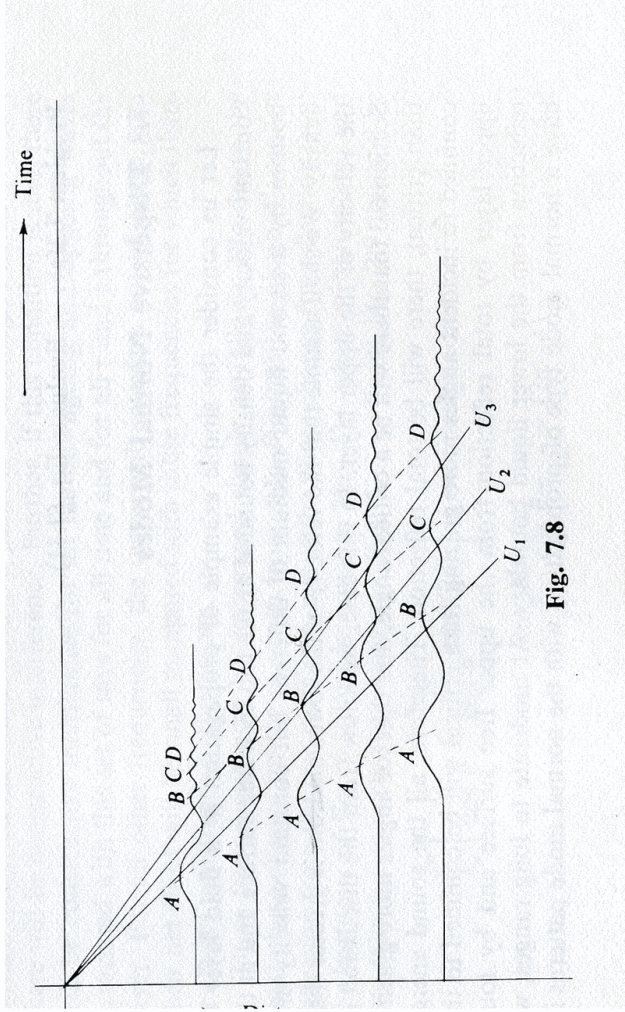


Fig. 7.8

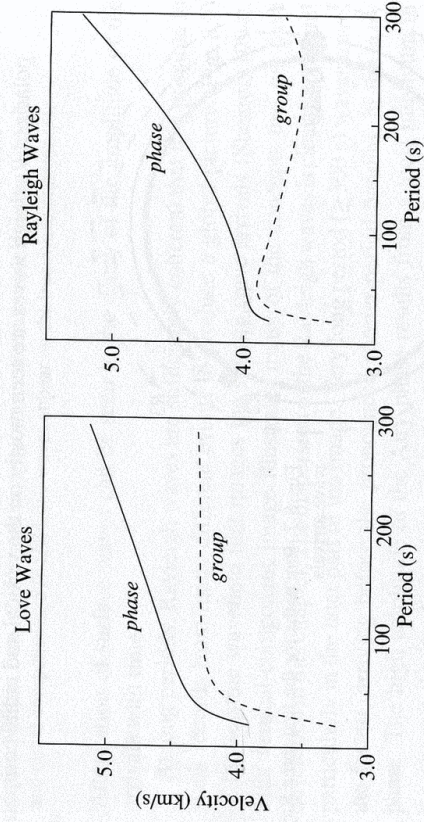
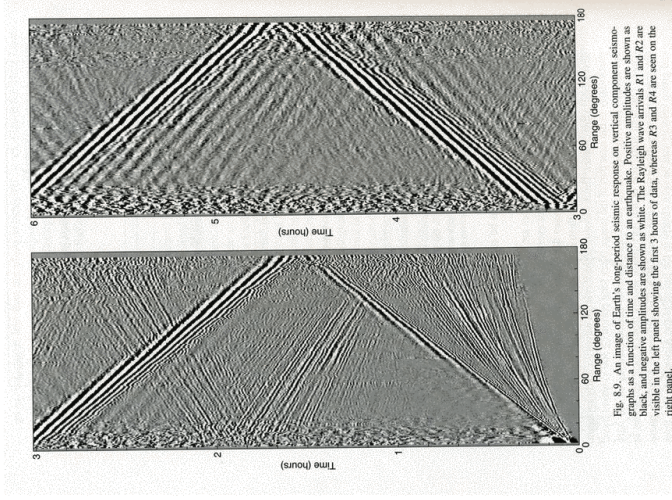


Fig. 8.6. Fundamental Love and Rayleigh dispersion curves computed from the isotropic PREM model (courtesy of Gabi Laske).



Measuring phase

Measuring phase

The phase of the Fourier transformed seismogram is just kx (with a source term) but we need to be careful how we measure this to avoid bias. Measure relative phase between data and synthetic or measure phase difference between orbits. Transfer function approach. Single spectral estimate:

$$S_{obs}(\omega) = T(\omega)S_{syn}(\omega)$$

$$T(\omega) = \frac{S_s^*(\omega)S_o(\omega)}{S_s^*(\omega)S_s(\omega)} = \frac{C_{12}(\omega)}{C_{11}(\omega)}$$

Better to use a multi-taper estimate. Minimize

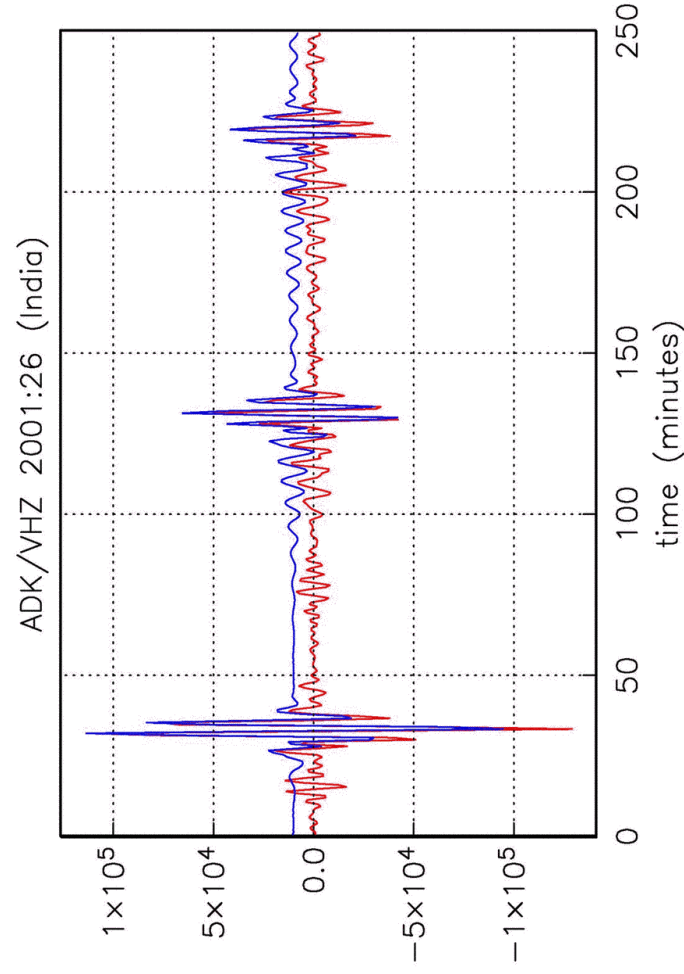
$$\|S_o(\omega) - T \cdot S_s(\omega)\|^2$$

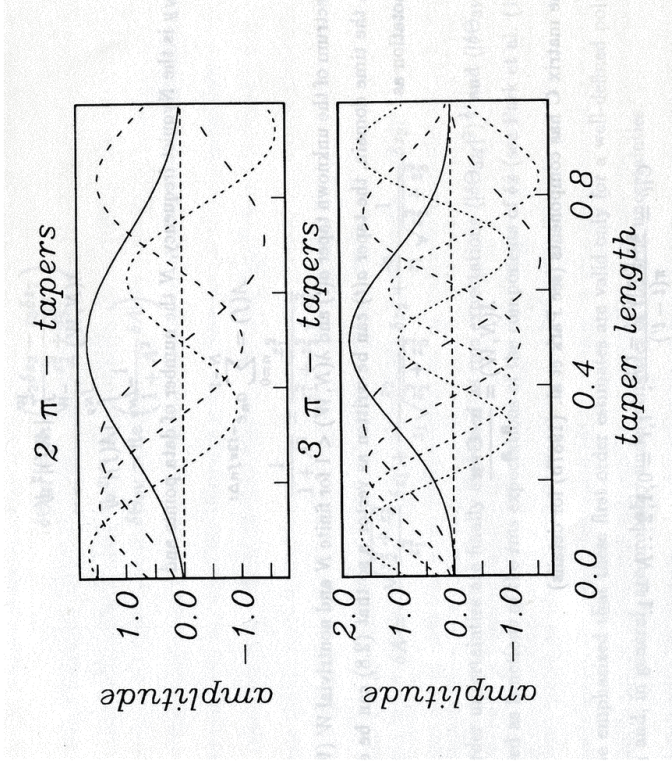
Solution is:

$$T(\omega) = \frac{S_o(\omega) \cdot S_s^*(\omega)}{S_s(\omega) \cdot S_s^*(\omega)}$$

T related to perturbation in (complex) wavenumber:

$$T(\omega) = e^{-\delta\bar{\gamma}(\omega)x} e^{-i\delta k(\omega)x}$$





Phase and group velocity maps

Phase and group velocity maps

- Ray theory (great circle propagation):

$$e^{ikx} \implies e^{i\vec{k}x} \quad \text{where} \quad \vec{k} = \frac{1}{x} \int_0^x k(x') dx'$$

where $k = k(\omega, \theta, \phi)$ or

$$c(\omega, \theta, \phi) = c_0(\omega) + \delta c(\omega, \theta, \phi)$$

In terms of measured phase:

$$\delta\Psi = \delta\vec{k}a\Delta = -\frac{\omega a}{c_0} \int_0^\Delta \frac{\delta c(\theta, \phi)}{c_0} d\gamma$$

- Ray theory for group velocity

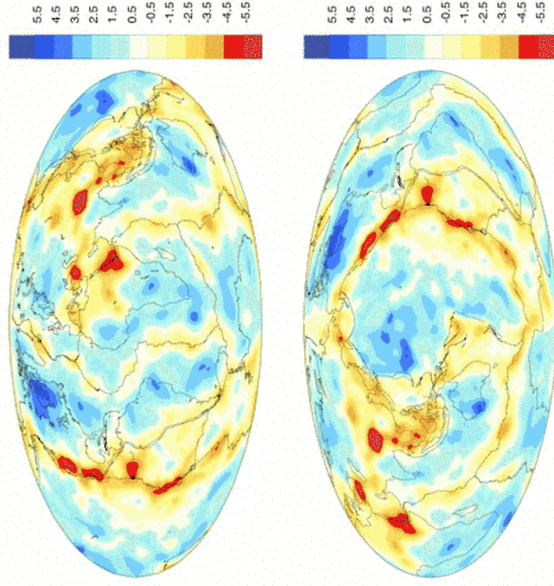
$$t_g = a \int_0^\Delta \frac{1}{U} d\gamma$$

Perturbation from reference group arrival time:

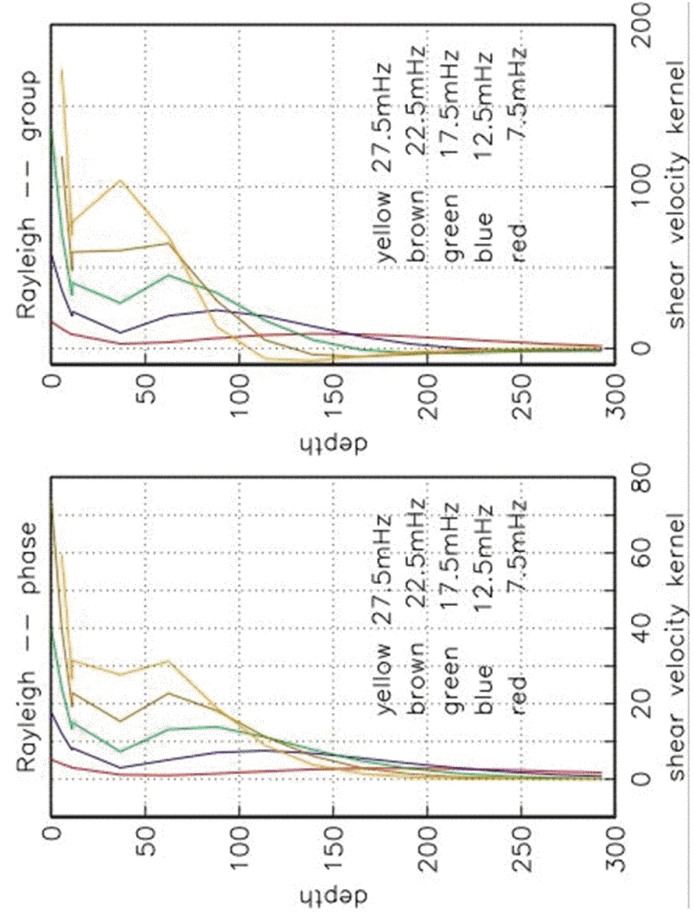
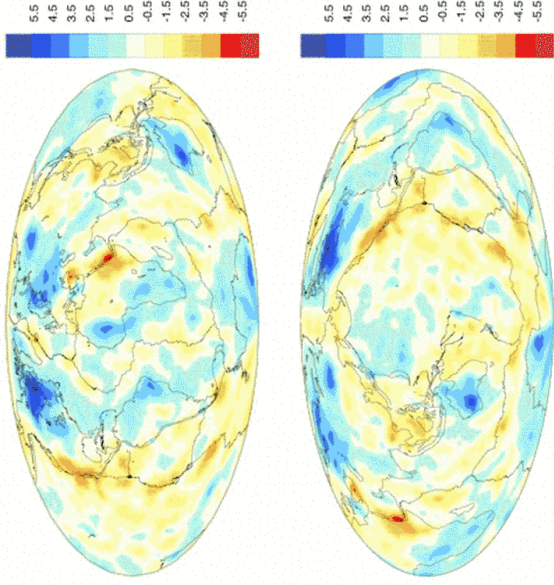
$$\delta t_g = -\frac{a}{U_0} \int_0^\Delta \frac{\delta U(\theta, \phi)}{U_0} d\gamma$$

(same form as phase velocity inversion)

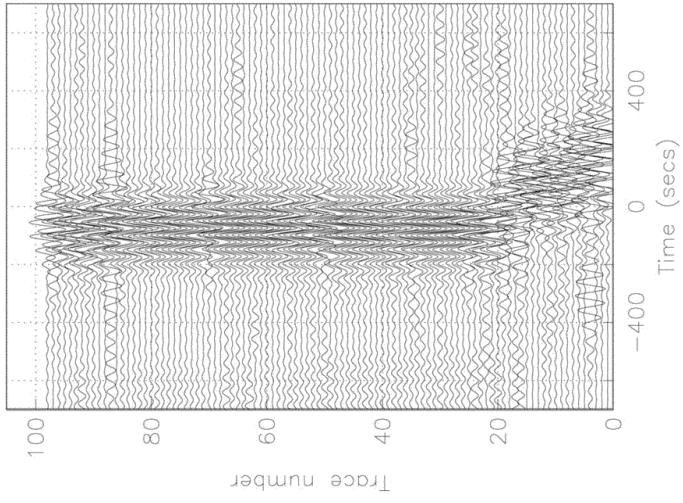
Rayleigh phase at 50 sec



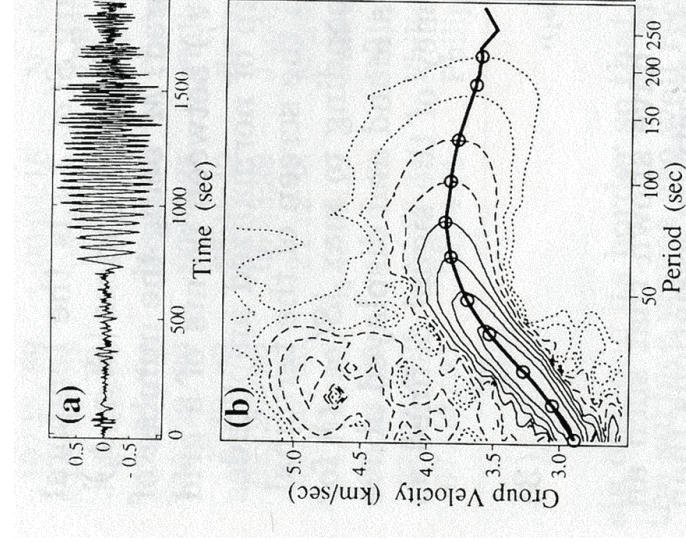
Rayleigh phase at 100 sec



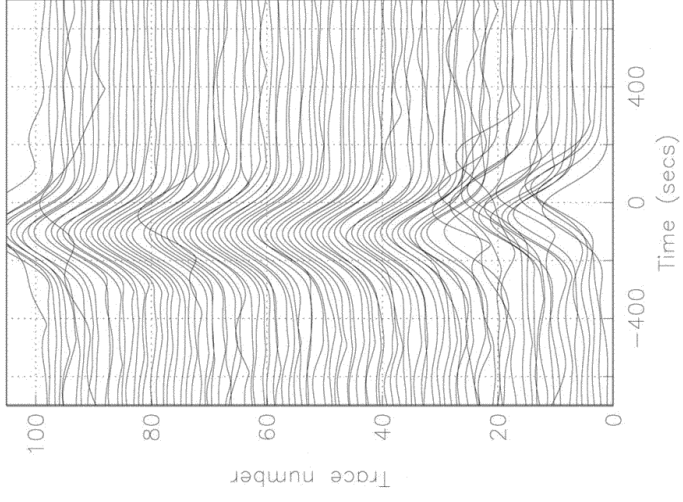
Undispersed 50 s Rayleigh waves



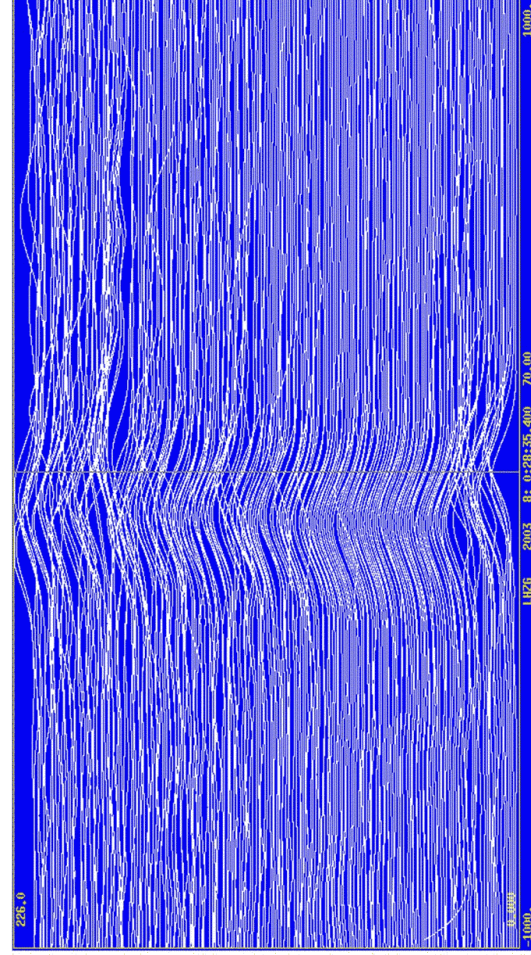
Note large number of cycle shifts possible



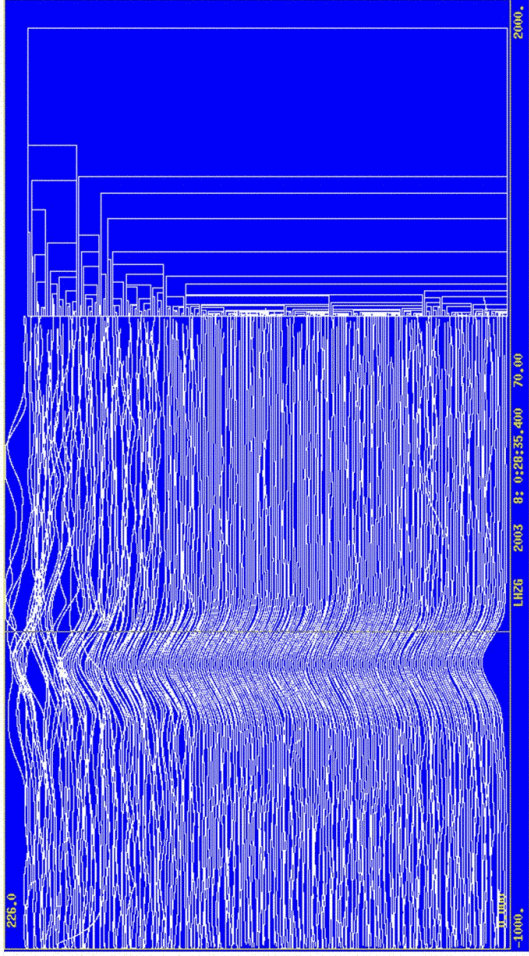
After computation of envelope functions



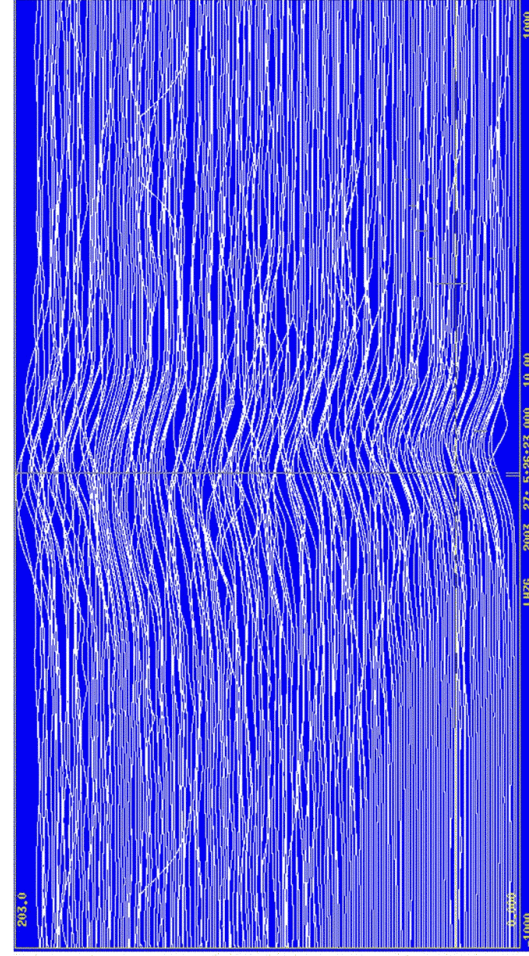
50 sec Rayleigh waves -- a typical example



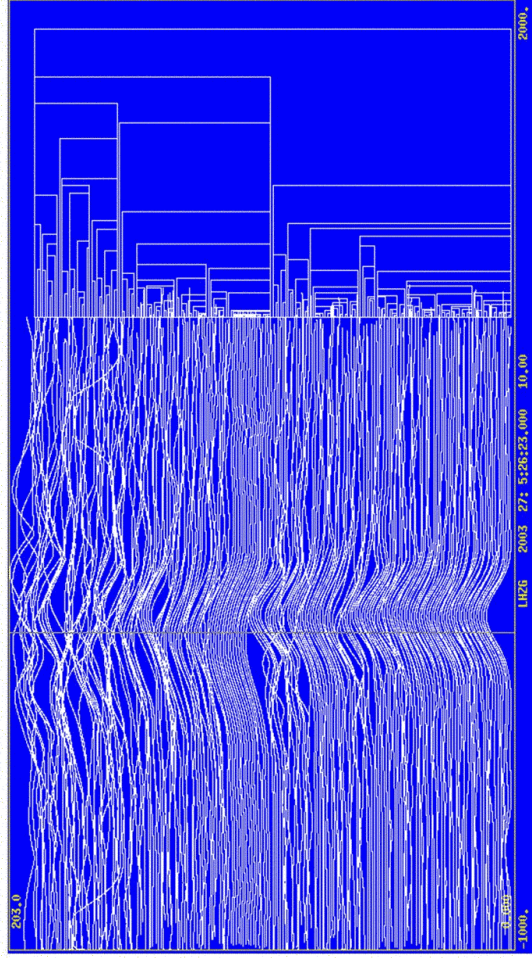
After clustering....



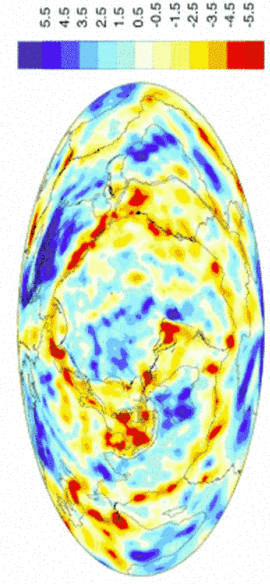
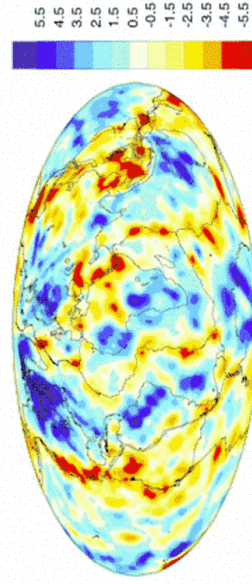
A more complicated example



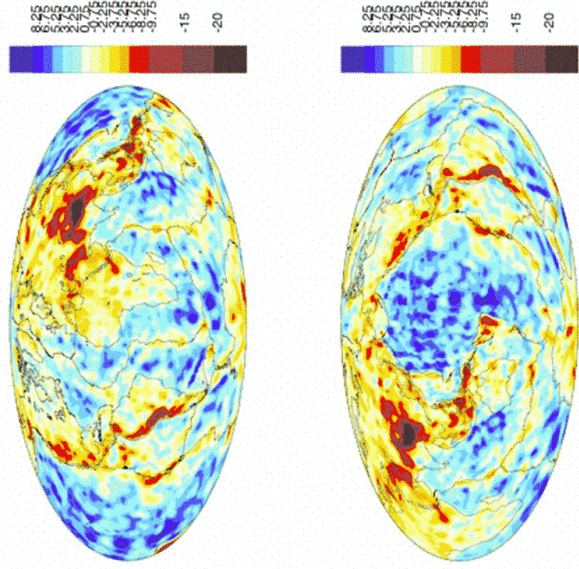
After clustering....



Rayleigh group at 100 seconds

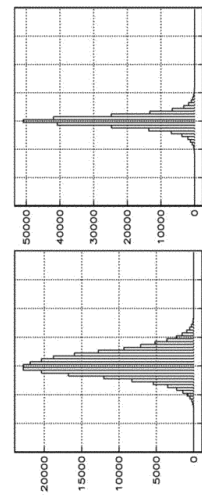


Rayleigh group at 50 seconds



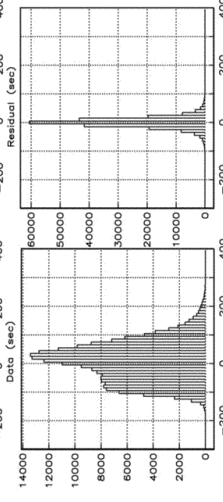
100s

71%, 239000



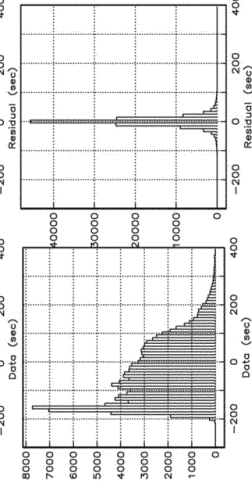
50s

93%, 215000



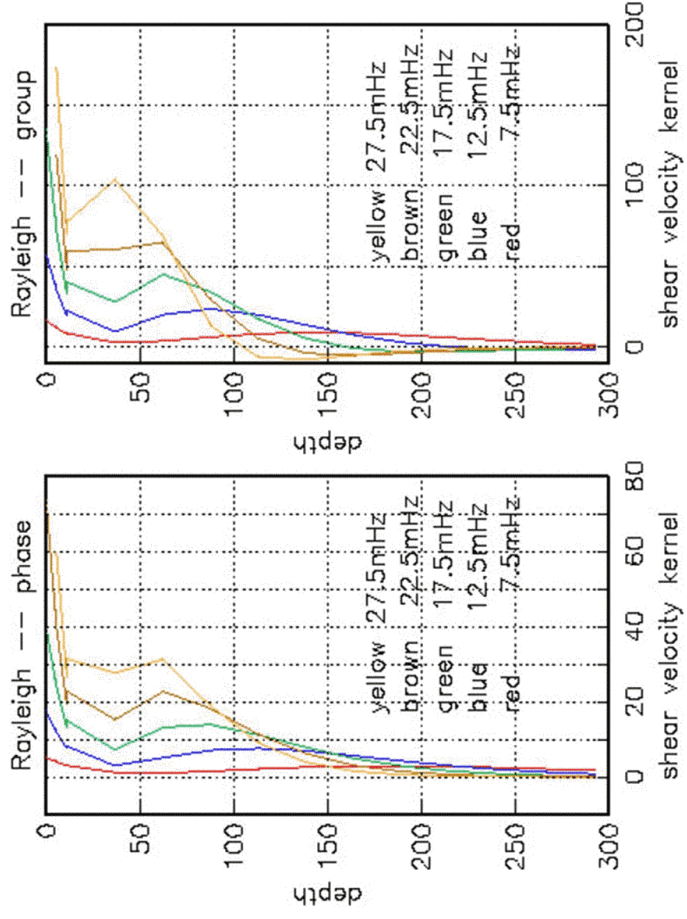
33s

98%, 125000

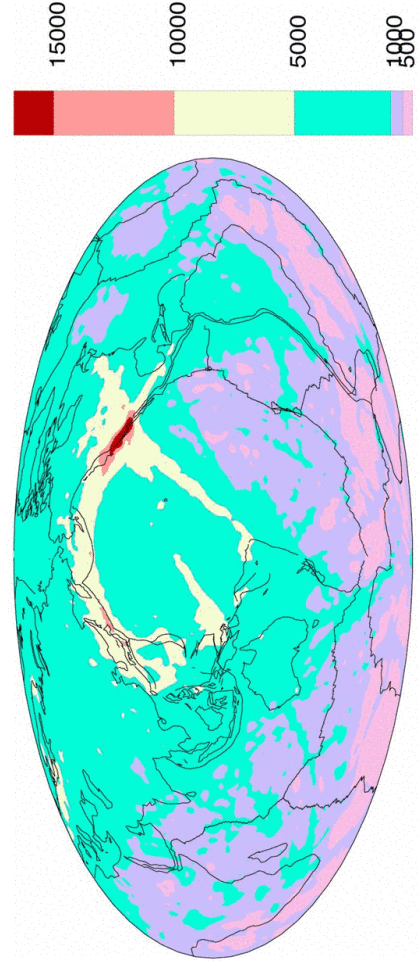


data

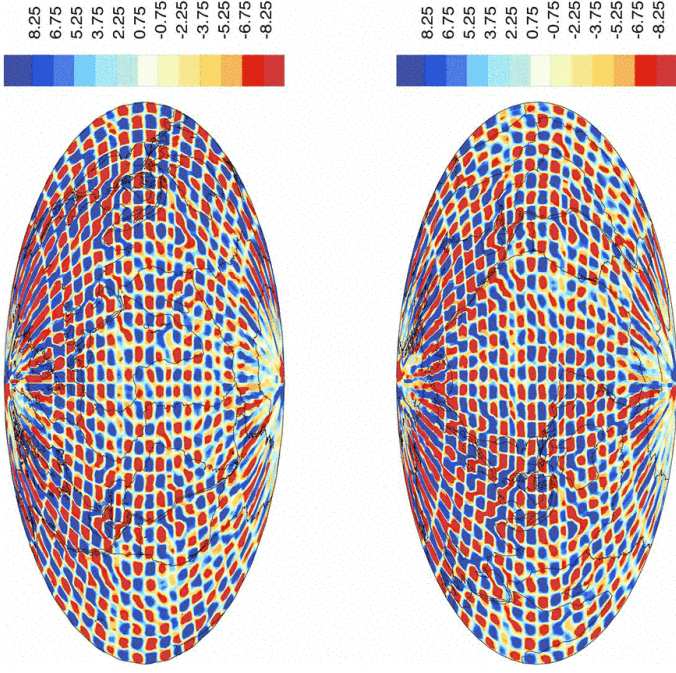
residual



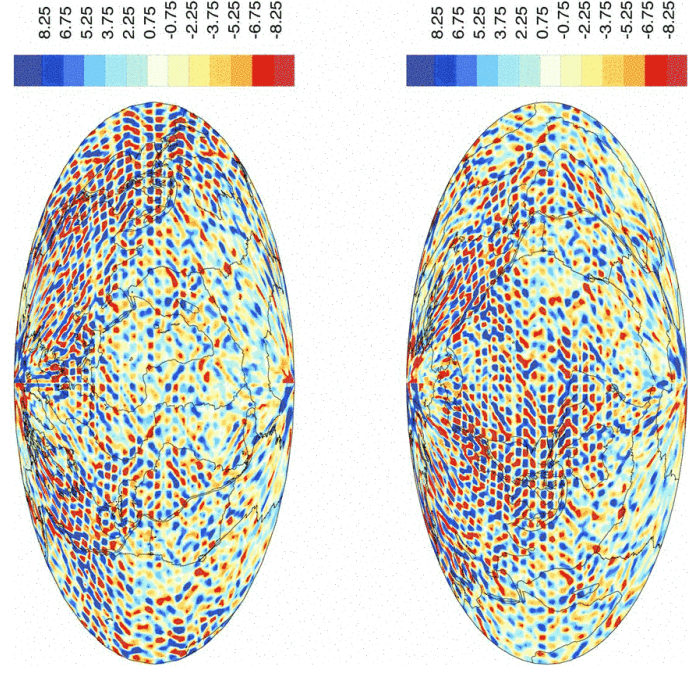
Hit count, 2 degree blocks



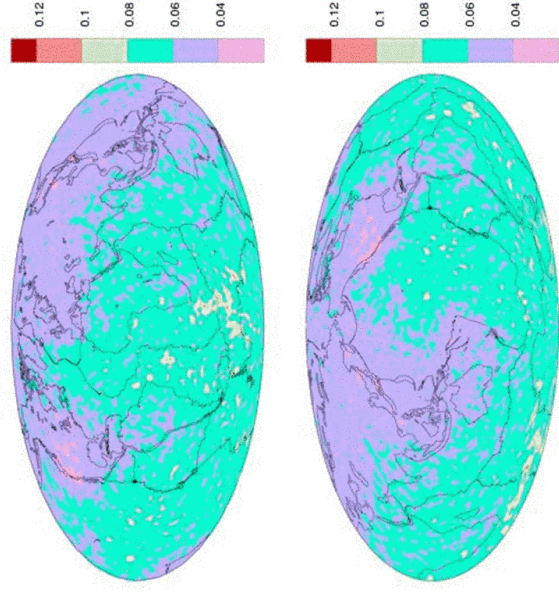
Resolution of 1000km structure is global



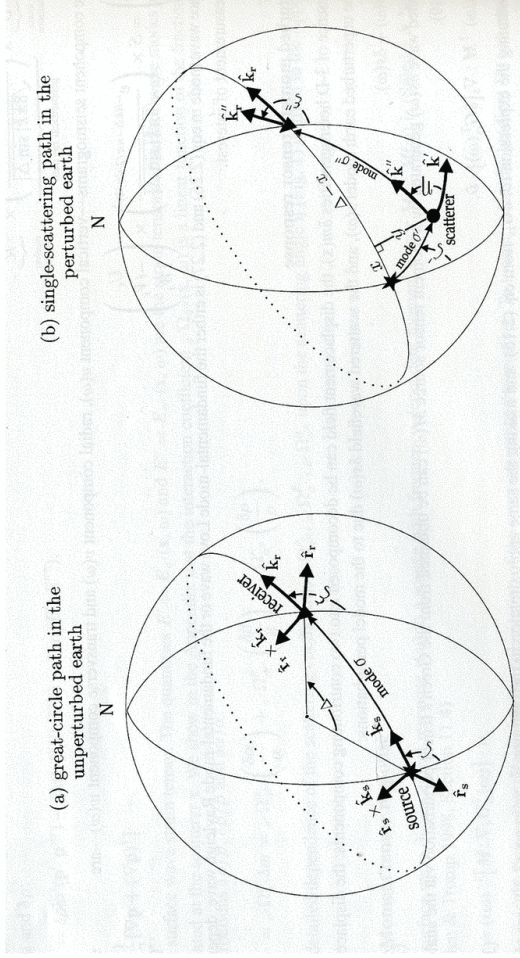
Resolution of finer structure is more variable



Error map for Rayleigh waves at 50 sec



Finite frequency kernels



Finite Frequency Kernels

- Born (single scattering) approximation

$$s(\omega) \rightarrow s(\omega) + \delta s(\omega)$$

$$\delta s(\omega) = \sum_{\sigma^i \sigma^{i'}} S^i \times \left(\frac{e^{-i(k^i \Delta^i - n^i \pi/2 + \pi/4)}}{\sqrt{8\pi k^i |\sin \Delta^i|}} \right) \times \sigma^{i'} \Omega_{\sigma^{i'}} \times \left(\frac{e^{-i(k^{i'} \Delta^{i'} - n^{i'} \pi/2 + \pi/4)}}{\sqrt{8\pi k^{i'} |\sin \Delta^{i'}|}} \right) \times \mathcal{R}^{i'} dV$$

Relation to measured transfer function:

$$s(\omega) + \delta s(\omega) = T(\omega) s(\omega)$$

3D kernel:

$$\delta \Psi(\omega) = \text{Im} \left(\frac{\delta s}{s} \right) = \int_V K_{\Psi} \delta m dV$$

K_{Ψ} is pretty ugly (see notes) but can be converted to a 2D kernel for a phase velocity perturbation with some approximations:

$$K_{\Psi} = \frac{2k^2 \sin [k(\Delta^i + \Delta^{i'} - \Delta) + \pi/4]}{\sqrt{8\pi k (|\sin \Delta^i| |\sin \Delta^{i'}| |\sin \Delta|)}}$$

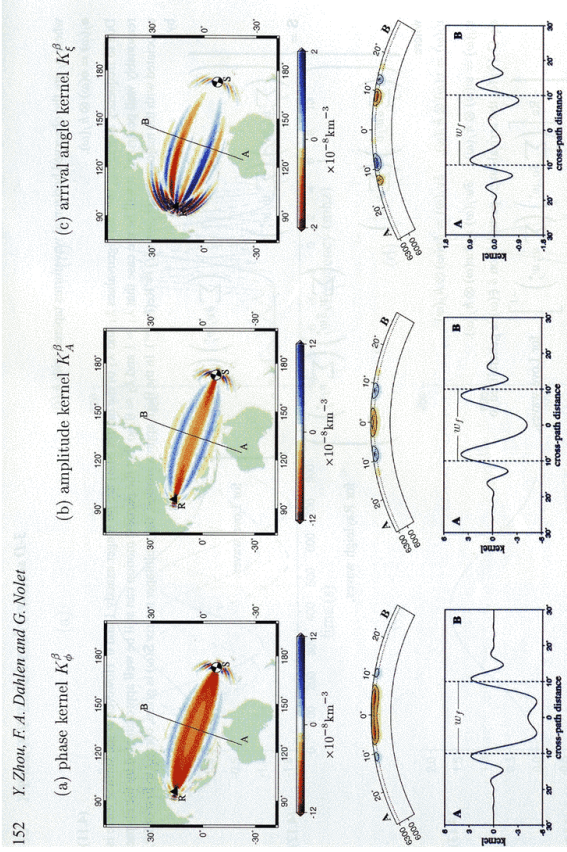
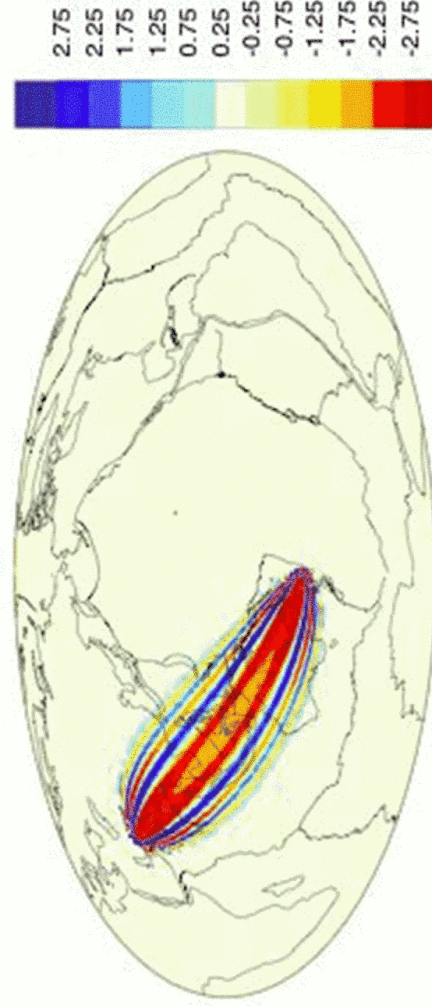
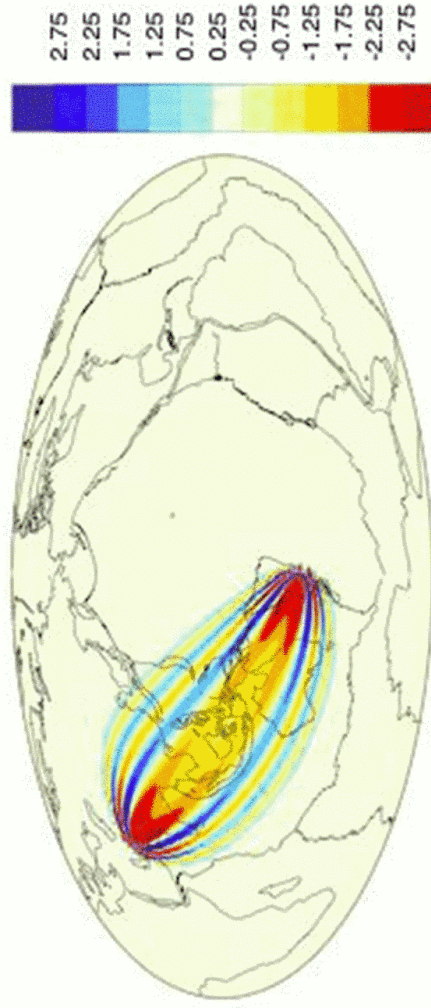


Figure 4. 3-D sensitivity kernels K_{ϕ}^{β} , K_A^{β} , K_{α}^{β} for a 10 mHz Love wave, excited by a 52 km deep strike-slip seismic source (S). Love wave radiation is maximum in the direction of the source-receiver geometrical ray (see beachball symbol). The epicentral distance to the receiver (R) is $\Delta = 80^\circ$. Sensitivity kernels are for 800 s cosine-taper measurements, with the taper centred at the group arrival time predicted by PREM. Top: Map view of kernels at the depth of approximately greatest sensitivity, 108 km. Middle: Slice view of cross-section AB, dotted lines are plotted at 108 km depth. Bottom: AB cross-section at a depth of 108 km; dashed lines indicate the width of the first Fresnel zone, w_f . Mode-coupling effects have been ignored, i.e. $\sigma_1 = \sigma_2 = \sigma$.

2D Phase kernel: 50s Rayleigh Waves



2D Phase kernel: 100s Rayleigh waves



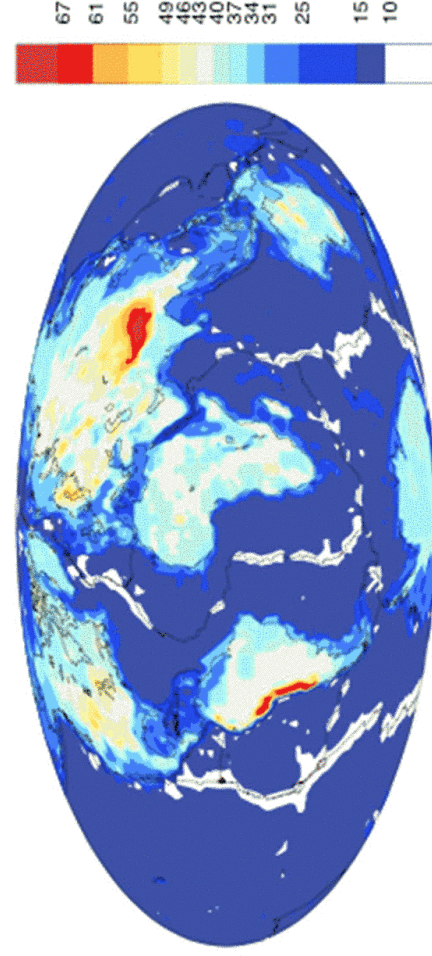
Finite frequency kernels

- Expect effect to be greatest at longest periods
- Difficult to make direct comparison with great circle result since different damping is implied
- Comparisons with analyses of SEM synthetics are ambiguous but 2D kernels do not seem effective -- need to use 3d kernels (Zhou et al, 2004)?
- What do you find?

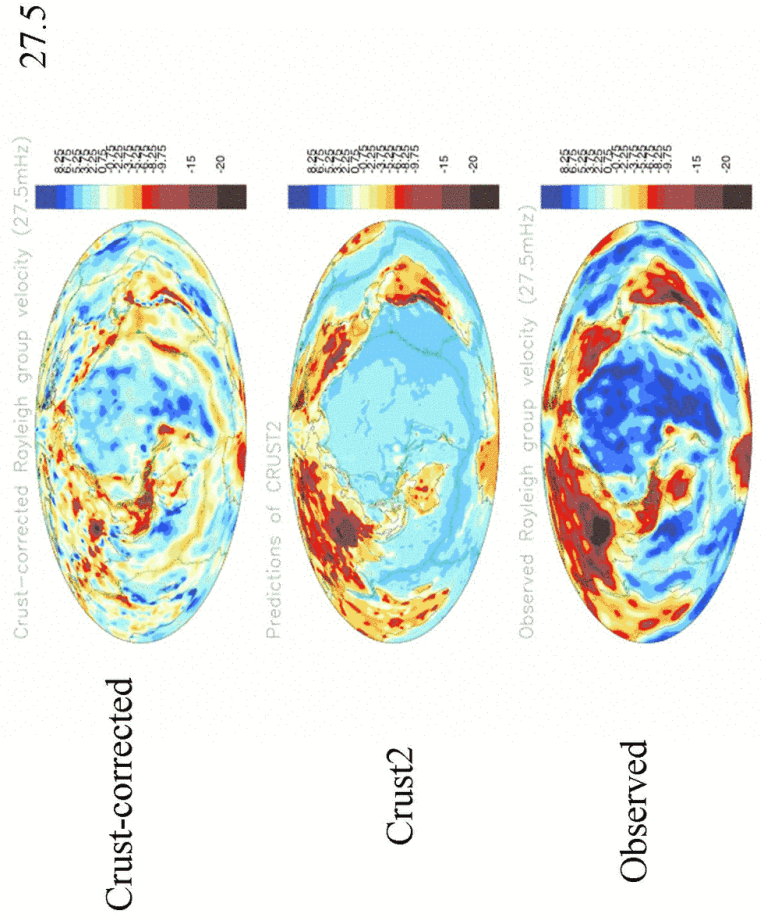
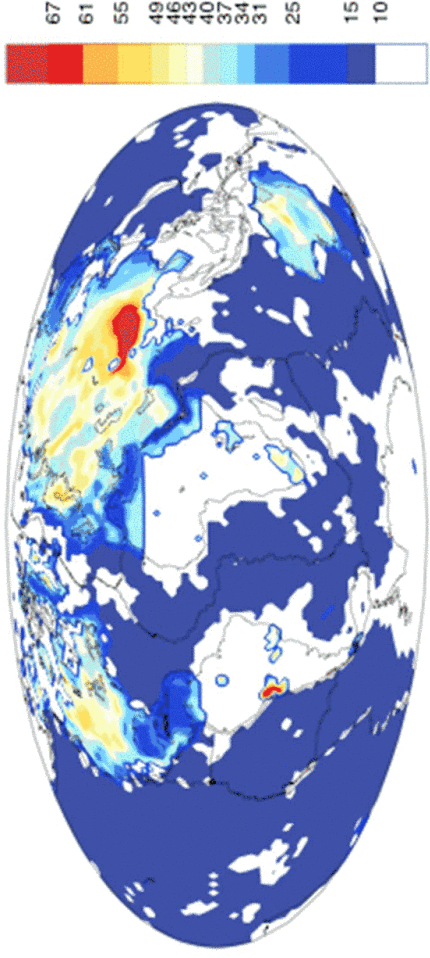
Inversion for crustal thickness

- Correct maps for the CRUST2 model
- Invert corrected maps with additional point constraints from receiver function studies for perturbation to CRUST2 moho depth
- Include only maps for periods shorter than 40 seconds to avoid mantle signal.

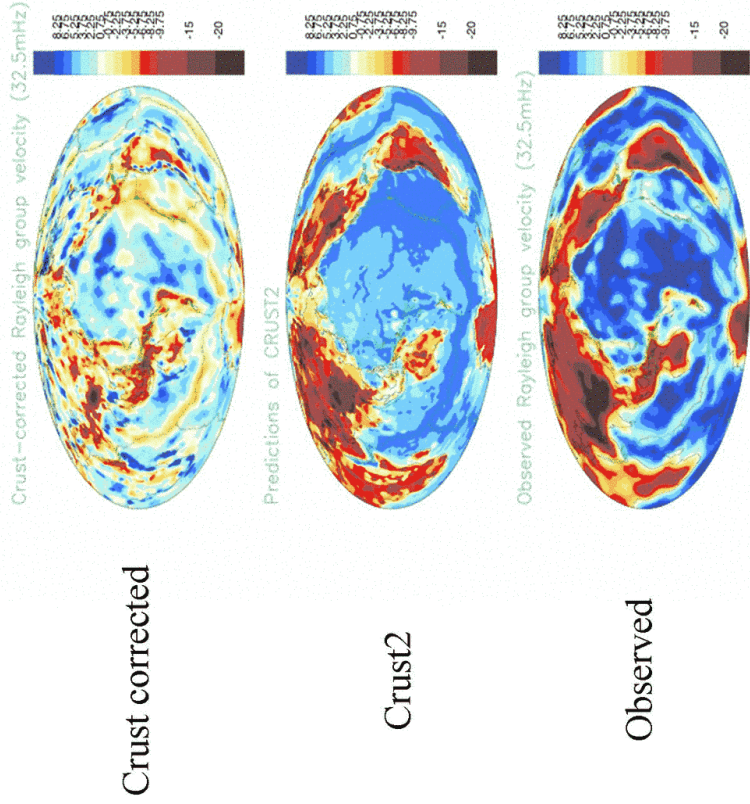
Depth to moho in CRUST2.0



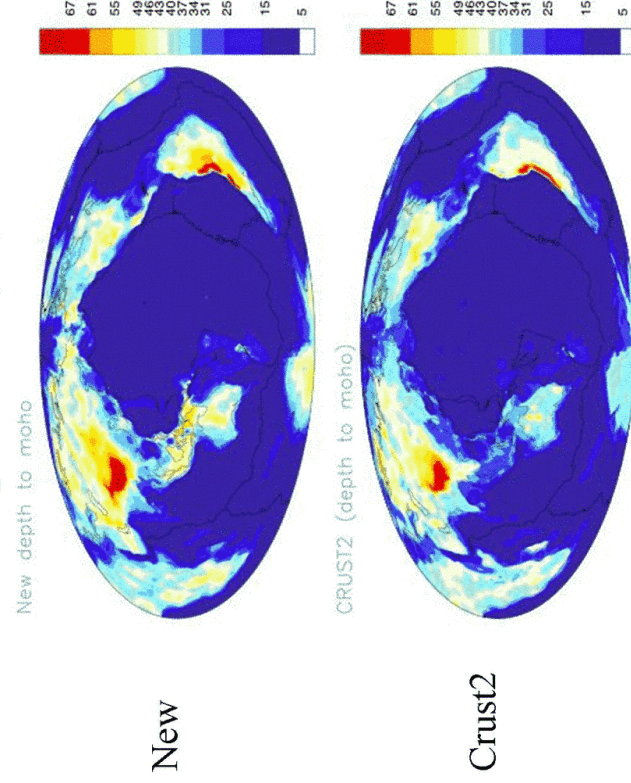
Observationally constrained locations

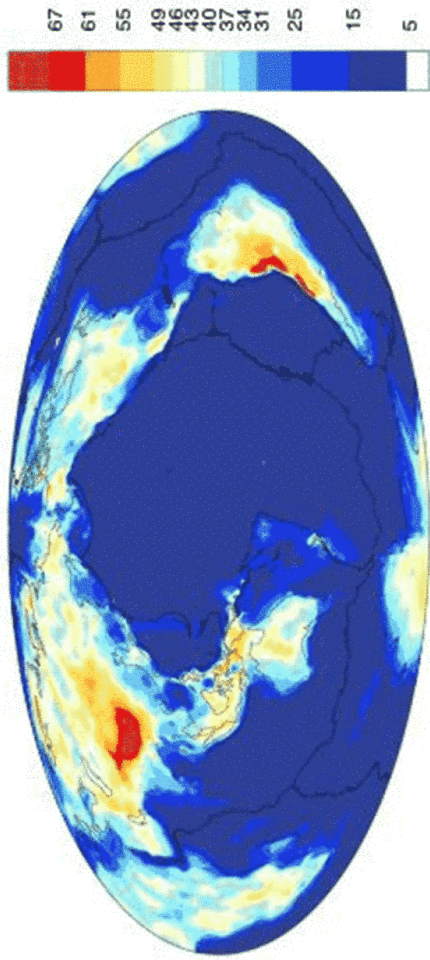


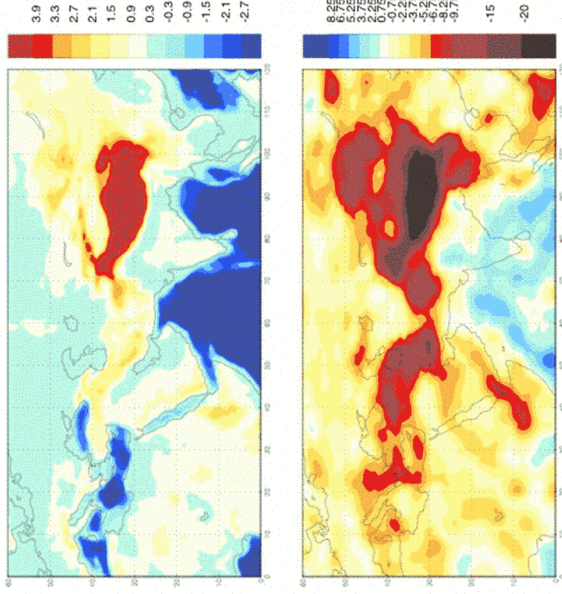
32.5



Depth to moho (km)

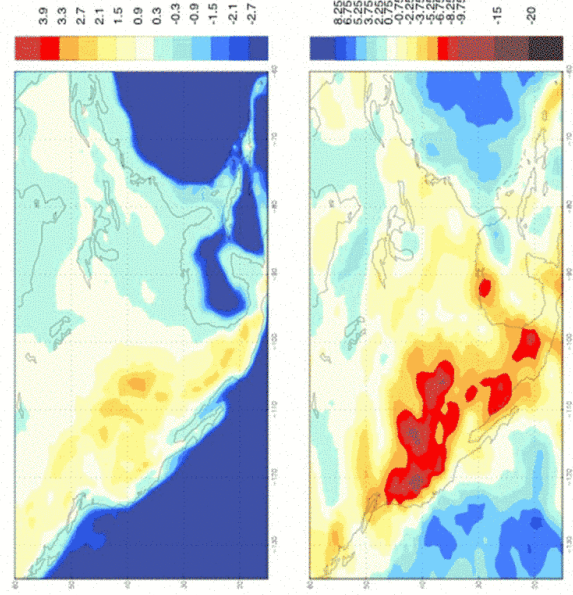






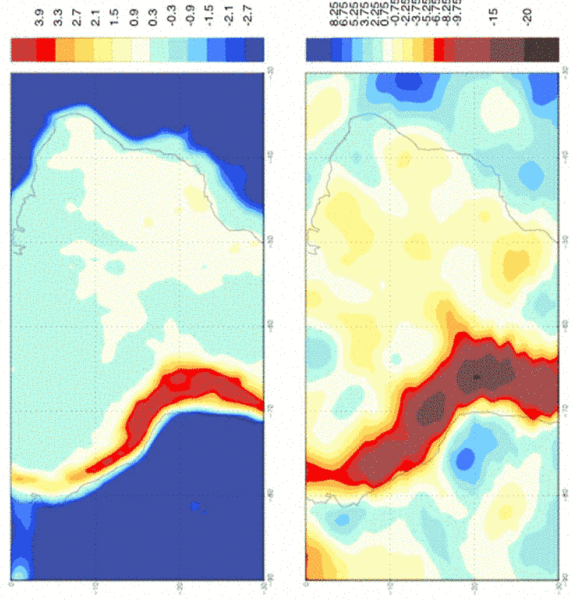
Topo

Group vel



Topo

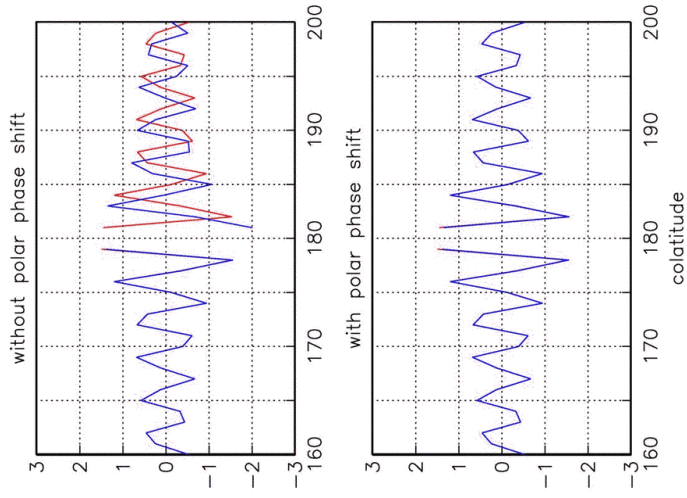
Group vel



Topo

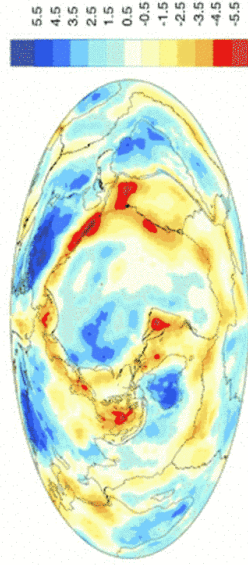
Group vel.

Where do those pole phase shifts come from?

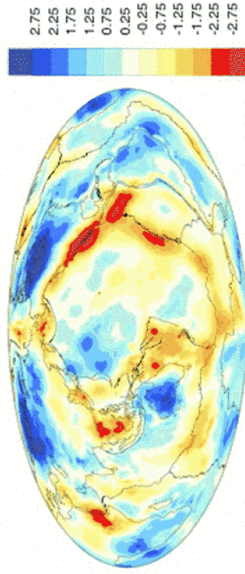


Comparison of X_{lm} with asymptotic form for $l=100$

Rayleigh, 100s, s20RTS+CRUST2

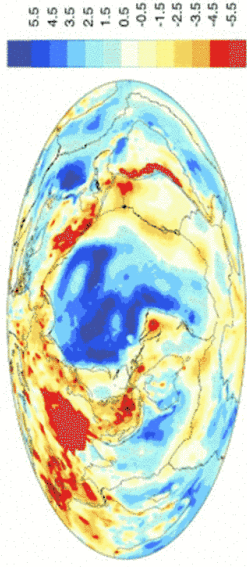


Group

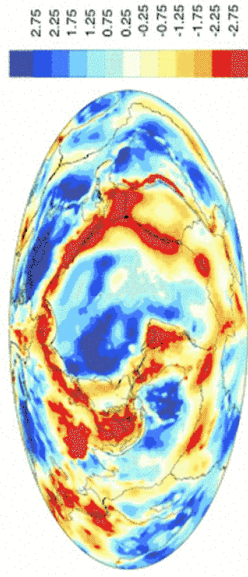


Phase

Rayleigh, 50s, S20RTS+CRUST2

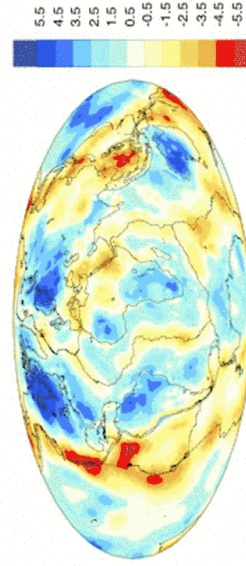


Group

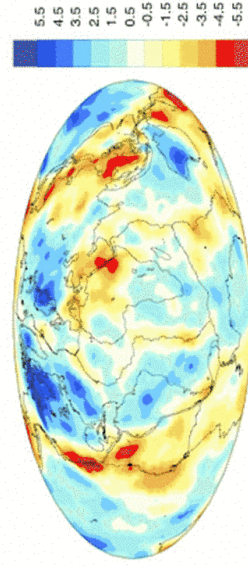


Phase

Rayleigh, group, 100s

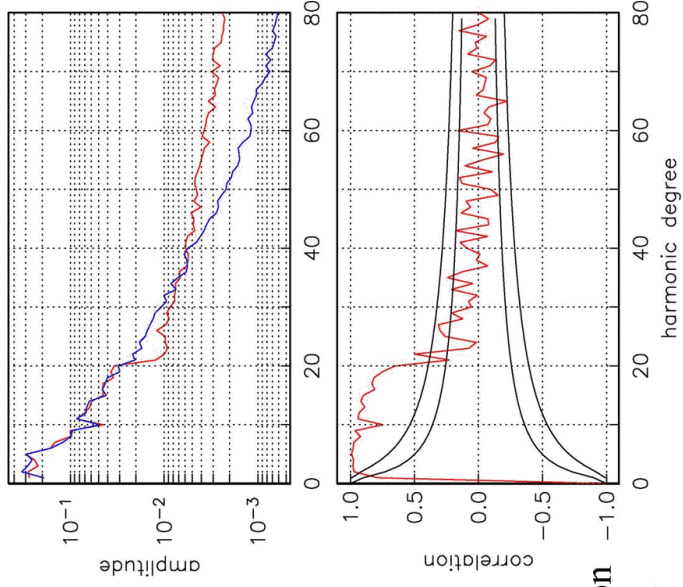


S20



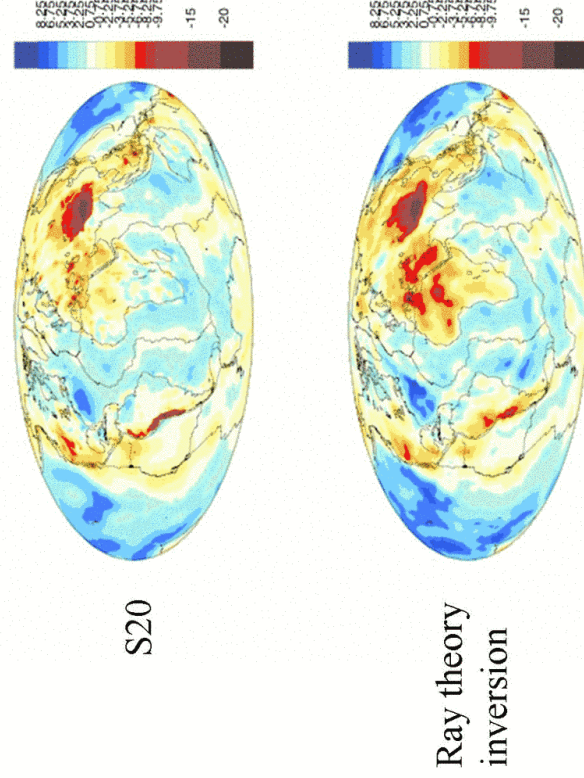
Ray theory
inversion

Synthetic Rayleigh, group, 100s

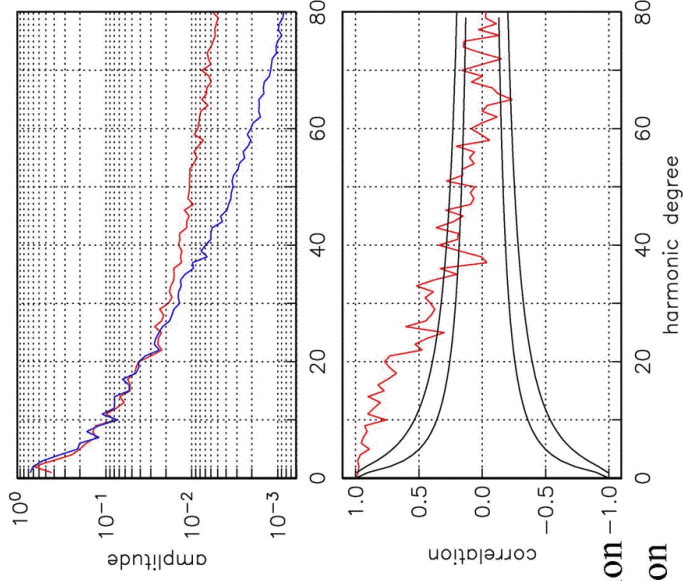


Red = prediction
Blue=inversion

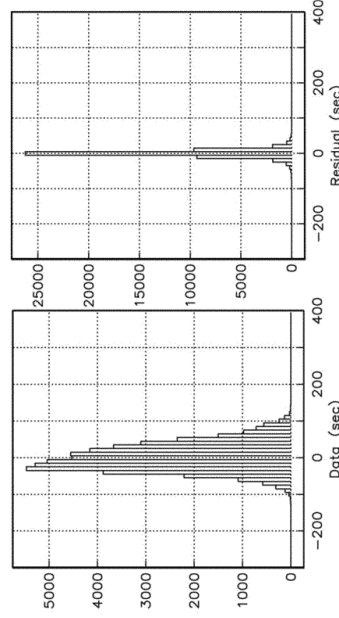
Rayleigh, group, 50s



Synthetic Rayleigh, group, 50s

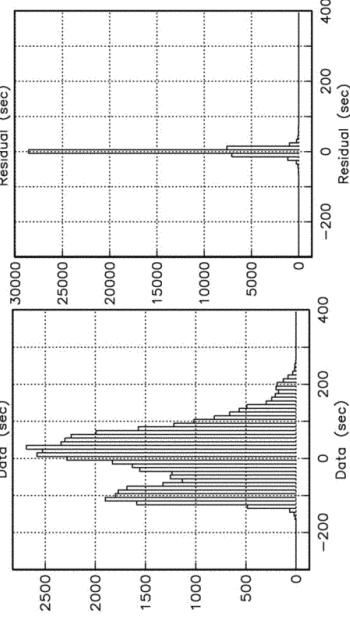


100s



93%, 50576

50s



99%, 46158

Inversion of SEM data using ray theory