

Regularities in Frequency/Period Spacings: How Asteroseismology Works

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Waves in Stars

Stars are gravitationally stratified \rightarrow two types of waves

Acoustic (Pressure) Waves



Dispersion Relation

$$\omega^2 = c^2(k_r^2 + k_h^2)$$

Sound Speed

$$c^2 = \frac{\gamma P}{\rho}$$

Gravity (Buoyancy) Waves



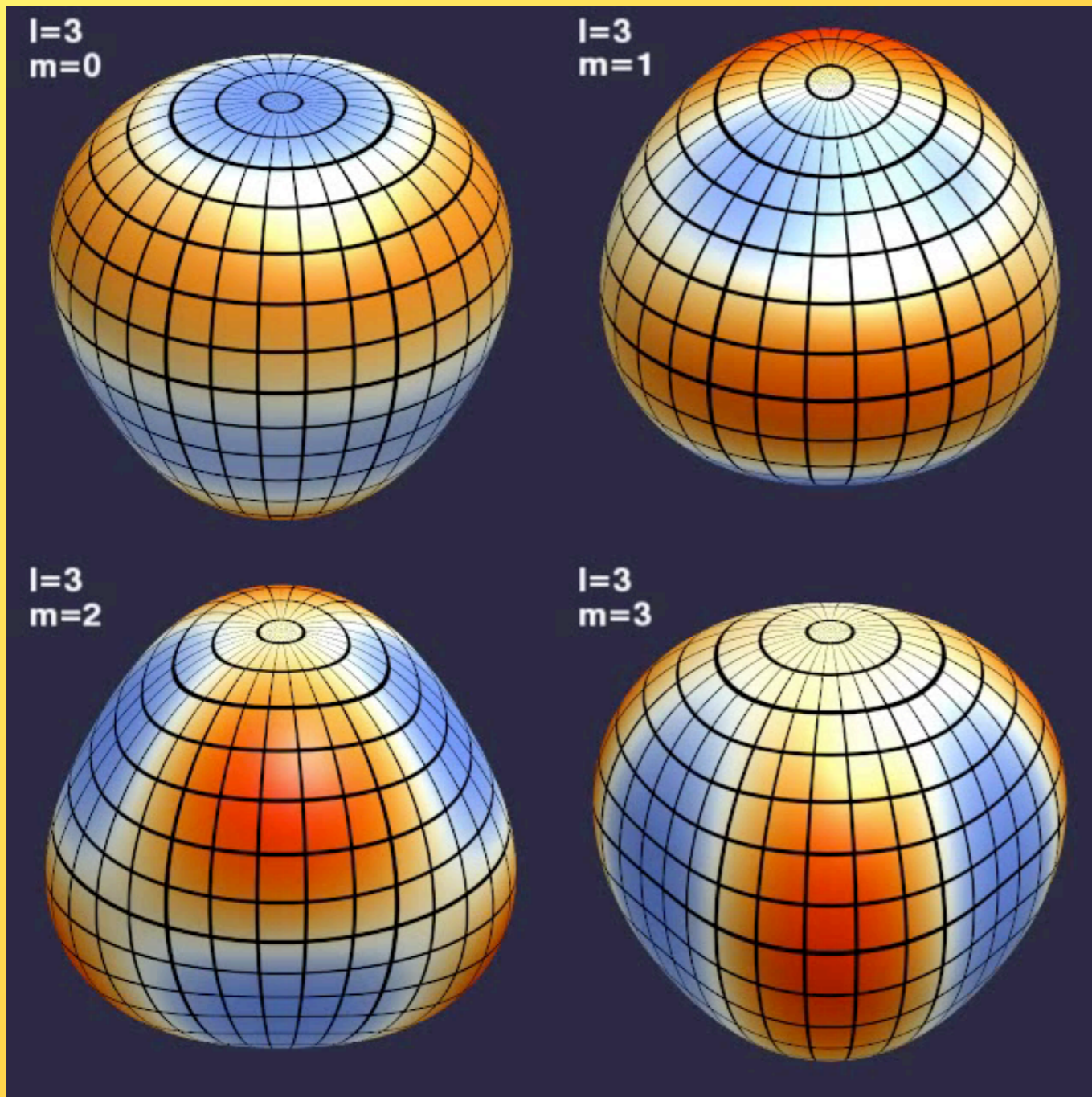
Dispersion Relation

$$\omega^2 = \frac{k_h^2}{k_r^2 + k_h^2} N^2$$

Brunt-Väisälä Frequency

$$N^2 = \frac{g}{r} \left(\frac{1}{\gamma} \frac{d \ln P}{d \ln r} - \frac{d \ln \rho}{d \ln r} \right)$$

Applying Periodic Boundary Conditions



Angular dependence of waves:
spherical harmonics

$$Y_{\ell}^m \propto P_{\ell}^m(\cos\theta) \exp(im\phi)$$

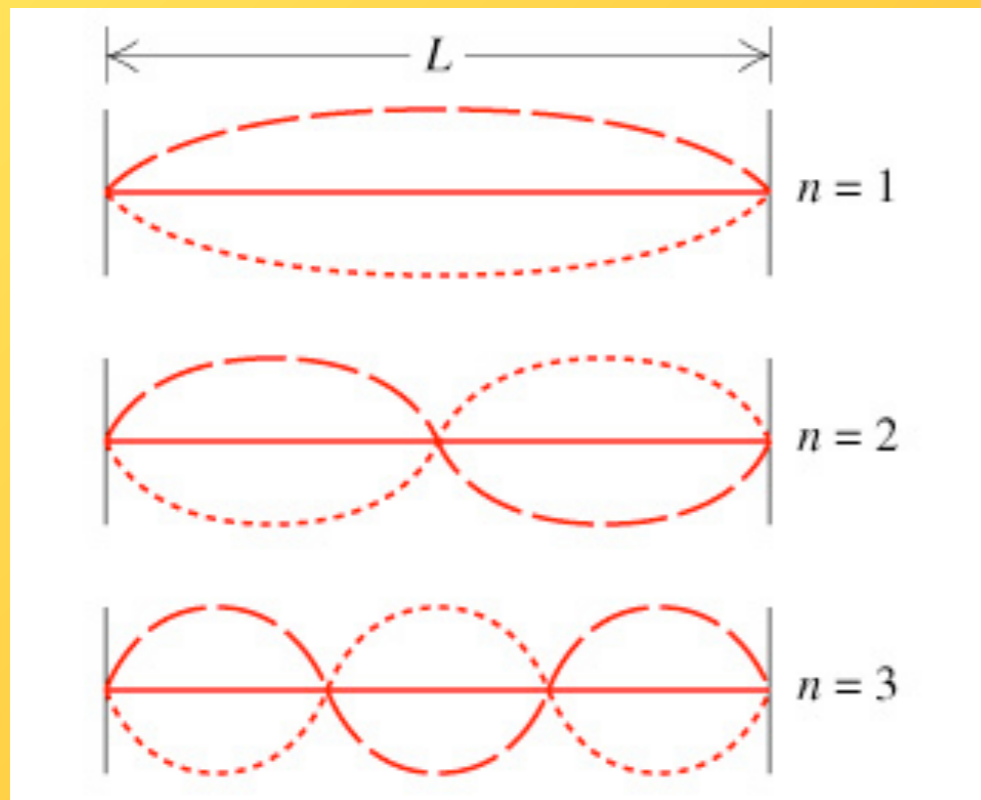
Spherical degree ℓ sets
horizontal wavenumber k_h

$$k_h^2 = \frac{\ell(\ell + 1)}{r^2}$$

Applying Radial Boundary Conditions

On a stretched string...

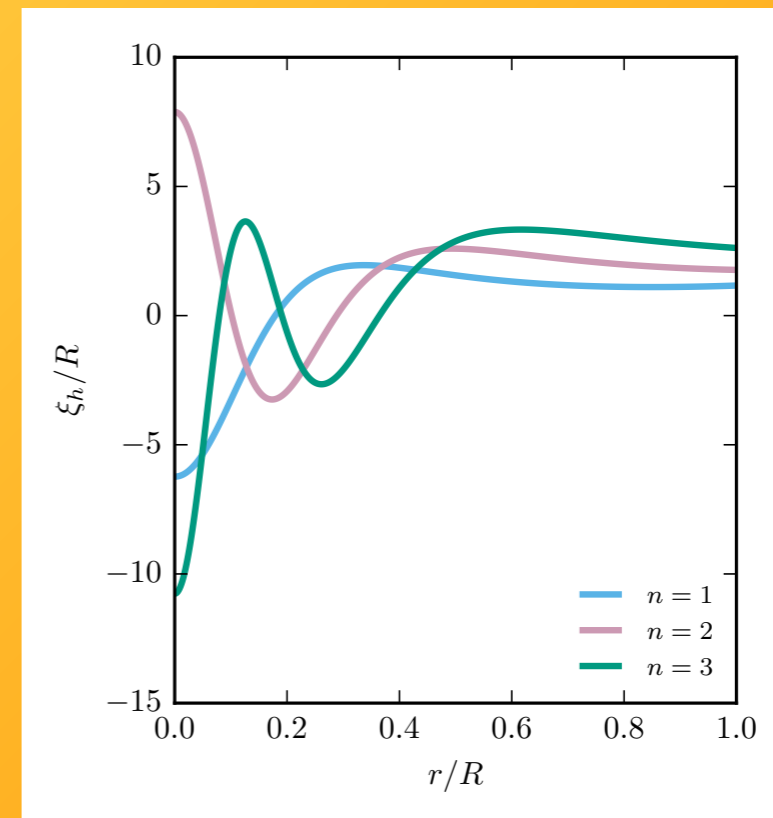
$$kL = n\pi$$



In a star...

$$\int_0^R k_r dr = (n + \alpha)\pi$$

boundary phase term
↓



Polytrope model

Asymptotic Frequencies of Pressure (p) Modes

Radial boundary condition

$$\int_0^R k_r dr = (n + \alpha)\pi$$

Asymptotic dispersion relation

$$\omega^2 = c^2(k_r^2 + k_h^2) \xrightarrow{k_r \gg k_h} k_r \approx \frac{\omega}{c}$$

Characteristic equation

$$\omega \int_0^R \frac{dr}{c} \approx (n + \alpha)\pi$$

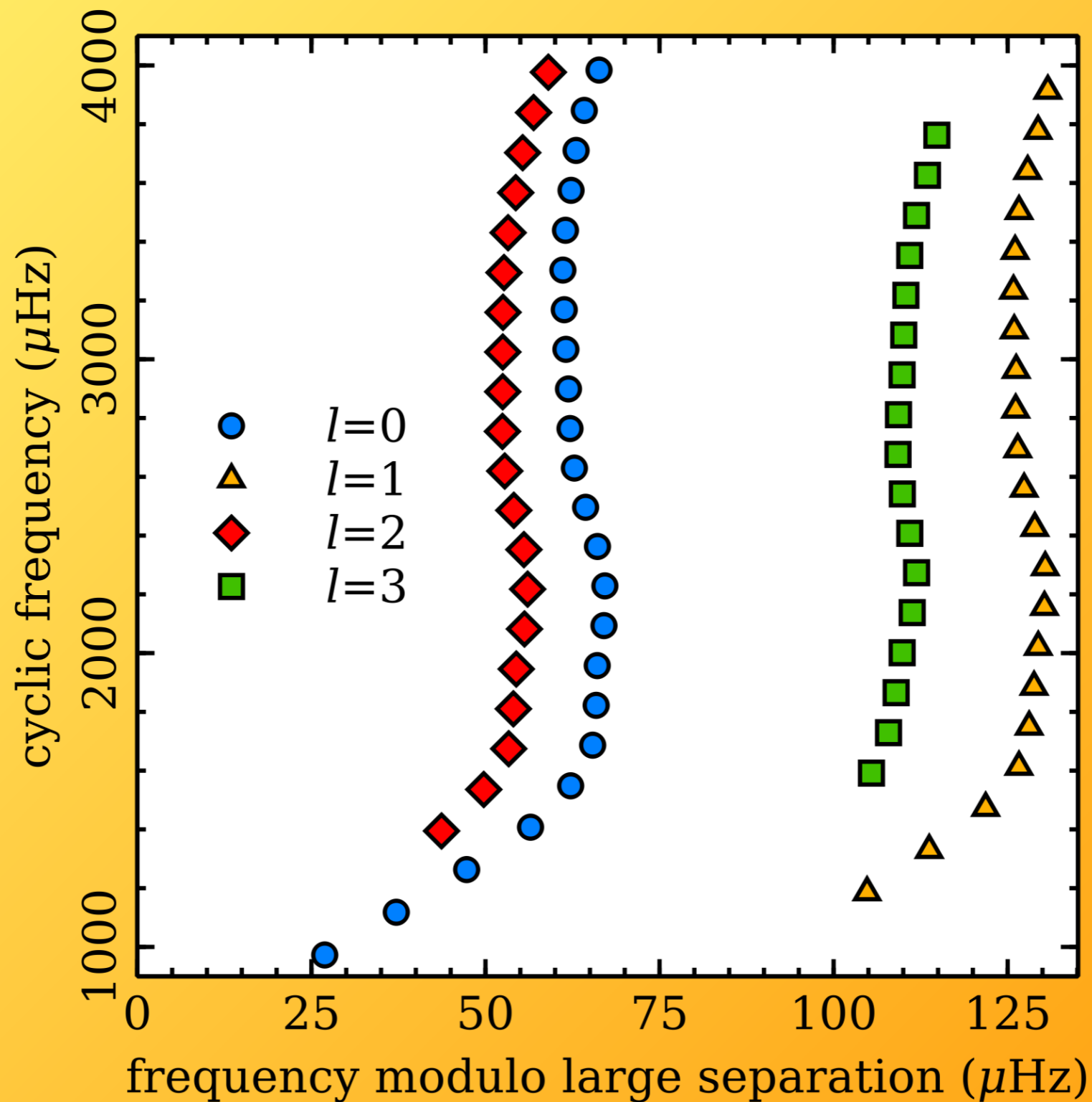
Solve for frequencies...

$$\nu_n \equiv \frac{\omega}{2\pi} \approx (n + \alpha) \underbrace{\left[2 \int \frac{dr}{c} \right]^{-1}}_{\Delta\nu}$$

...p-modes with consecutive n will show a uniform spacing in frequency $\Delta\nu$



Echelle Diagram for the Sun



Modes with consecutive n stack vertically in echelle diagram

$\Delta\nu$

Warrick Ball
(BiSON data)



p-Mode Frequencies Revisited

Asymptotic dispersion relation

$$\omega^2 = c^2 k^2 \xrightarrow{k_r \gg k_h} k_r \approx \frac{\omega}{c} \sqrt{1 - \frac{k_h^2 c^2}{\omega^2}}$$

Angular boundary condition

$$k_h^2 = \frac{\ell(\ell + 1)}{r^2}$$

Characteristic Equation

$$\omega \int_{r_1}^R \sqrt{1 - \frac{\ell(\ell + 1)c^2}{\omega^2 r^2}} \frac{dr}{c} \approx (n + \alpha)\pi$$

turning point
($k_r \rightarrow 0$)

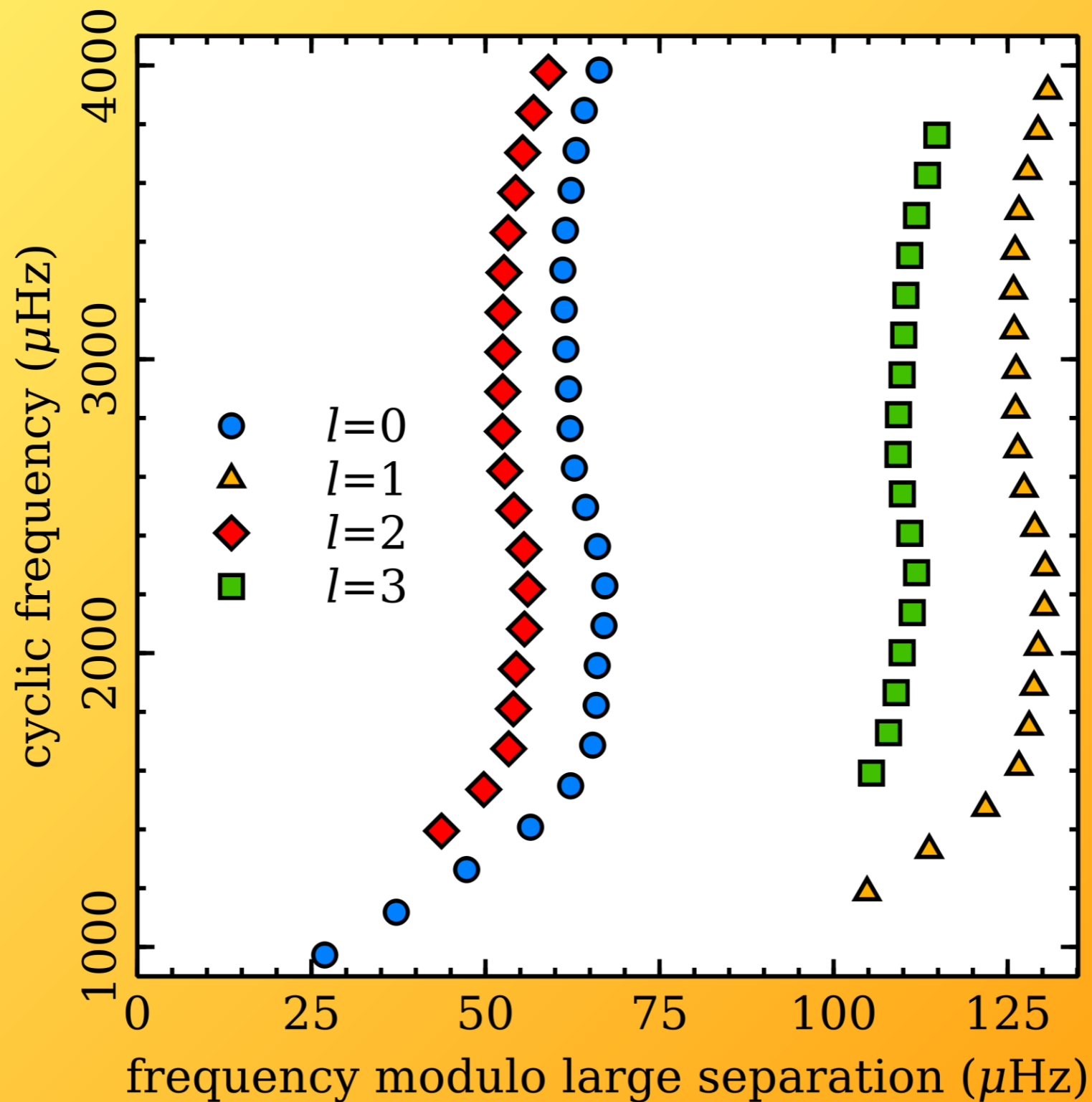
this step
involves math!

Solve for frequencies...

$$\nu_{nl} \equiv \frac{\omega}{2\pi} \approx \left(n + \frac{\ell}{2} + \frac{1}{4} + \alpha \right) \left[2 \int \frac{dr}{c} \right]^{-1}$$

...p-modes with $(n', \ell') = (n-1, \ell+2)$ will have (almost) the same frequencies

Echelle Diagram for the Sun



Odd- l modes are offset by $\Delta\nu/2$ from even- l modes

$\Delta\nu$

Asymptotic Periods of Gravity (g) Modes

Asymptotic dispersion relation $\omega^2 = \frac{k_h^2}{k^2} N^2 \xrightarrow{k_r \gg k_h} k_r \approx \frac{k_h N}{\omega}$

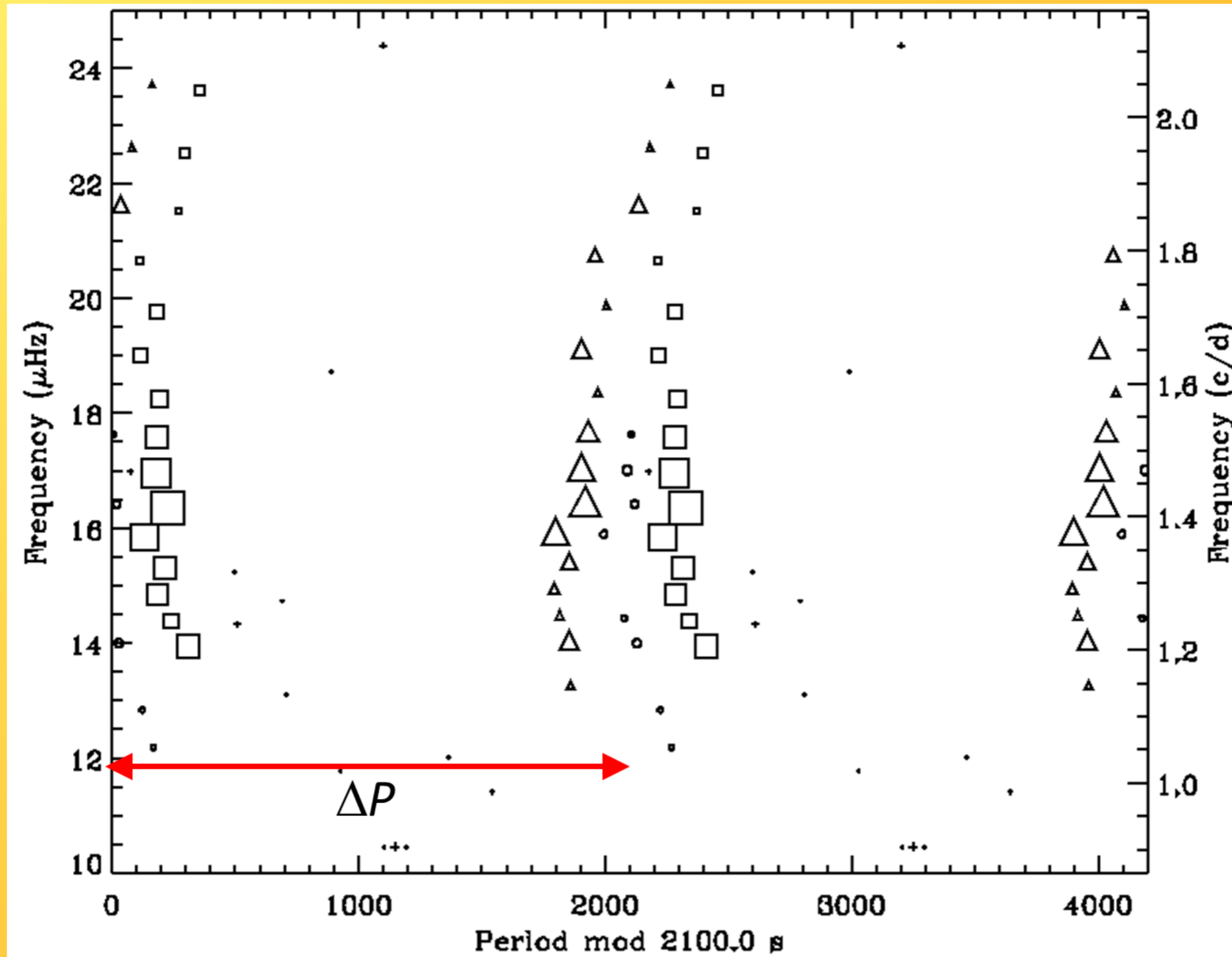
Angular boundary condition $k_h^2 = \frac{\ell(\ell + 1)}{r^2}$

Characteristic Equation $\frac{\sqrt{\ell(\ell + 1)}}{\omega} \int_{r_1}^{r_2} \frac{N}{r} dr \approx (n + \alpha)\pi$

Solve for periods... $P_{n\ell} \equiv \frac{2\pi}{\omega} \approx (n + \alpha) \overbrace{\left[\frac{\sqrt{\ell(\ell + 1)}}{2\pi^2} \int_{r_1}^{r_2} \frac{N}{r} dr \right]^{-1}}^{\Delta P}$

...g-modes with same ℓ and consecutive n will show a uniform period spacing ΔP

Period Echelle Diagram



Mixed Modes: Dispersion Relation

Gravito-acoustic dispersion relation

$$k_r^2 = \frac{1}{\omega^2 c^2} (\omega^2 - N^2) (\omega^2 - S_\ell^2)$$

$$S_\ell^2 = \frac{\ell(\ell + 1)c^2}{r^2}$$

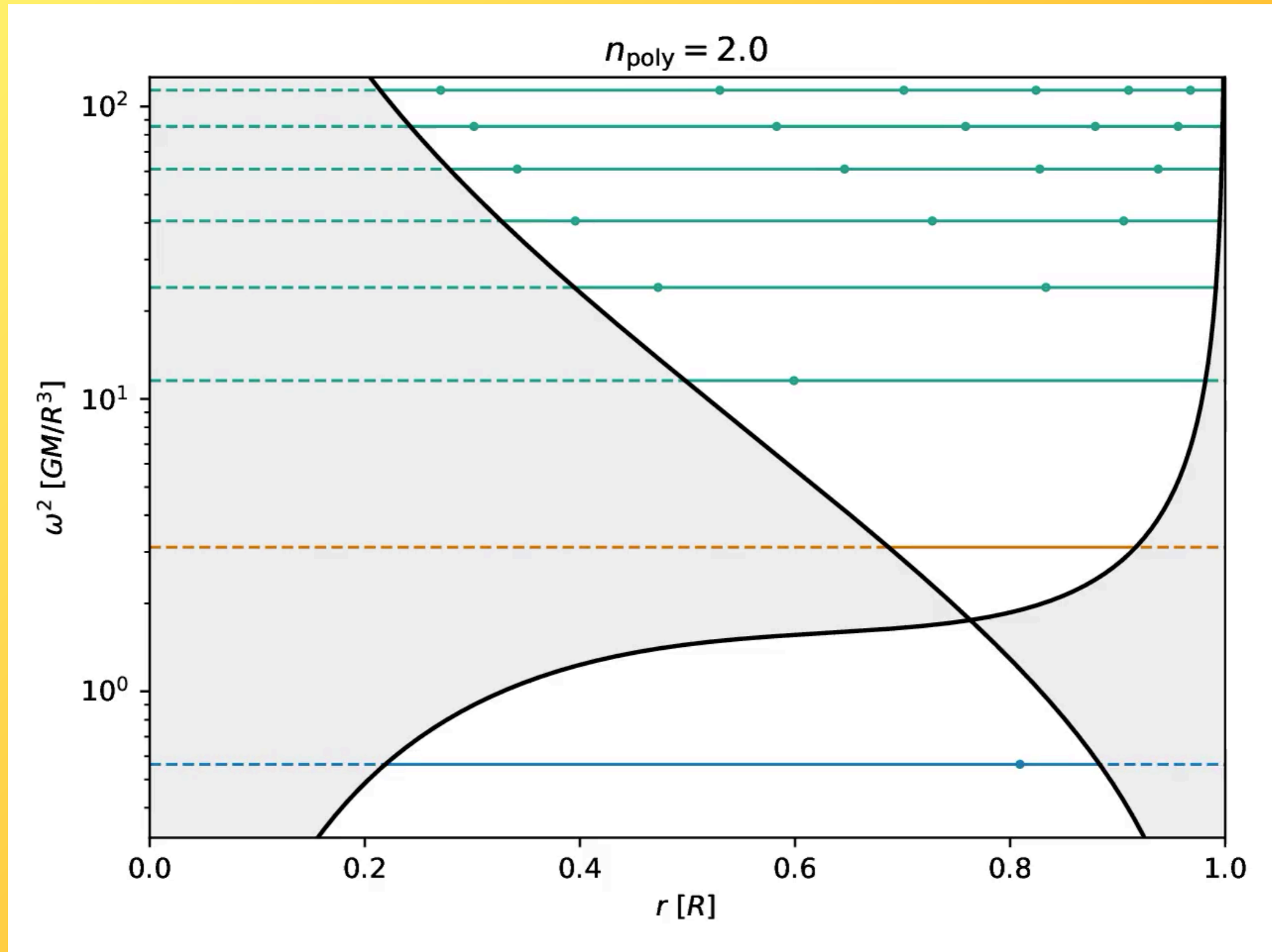
Lamb Frequency

$\omega^2 < N^2, S_\ell^2 \longrightarrow$ mode behaves locally as a gravity wave

$\omega^2 > N^2, S_\ell^2 \longrightarrow$ mode behaves locally as an acoustic wave

$\left. \begin{array}{l} N^2 < \omega^2 < S_\ell^2 \\ S_\ell^2 < \omega^2 < N^2 \end{array} \right\} \longrightarrow$ mode is locally evanescent

Mixed Modes: Propagation Diagram



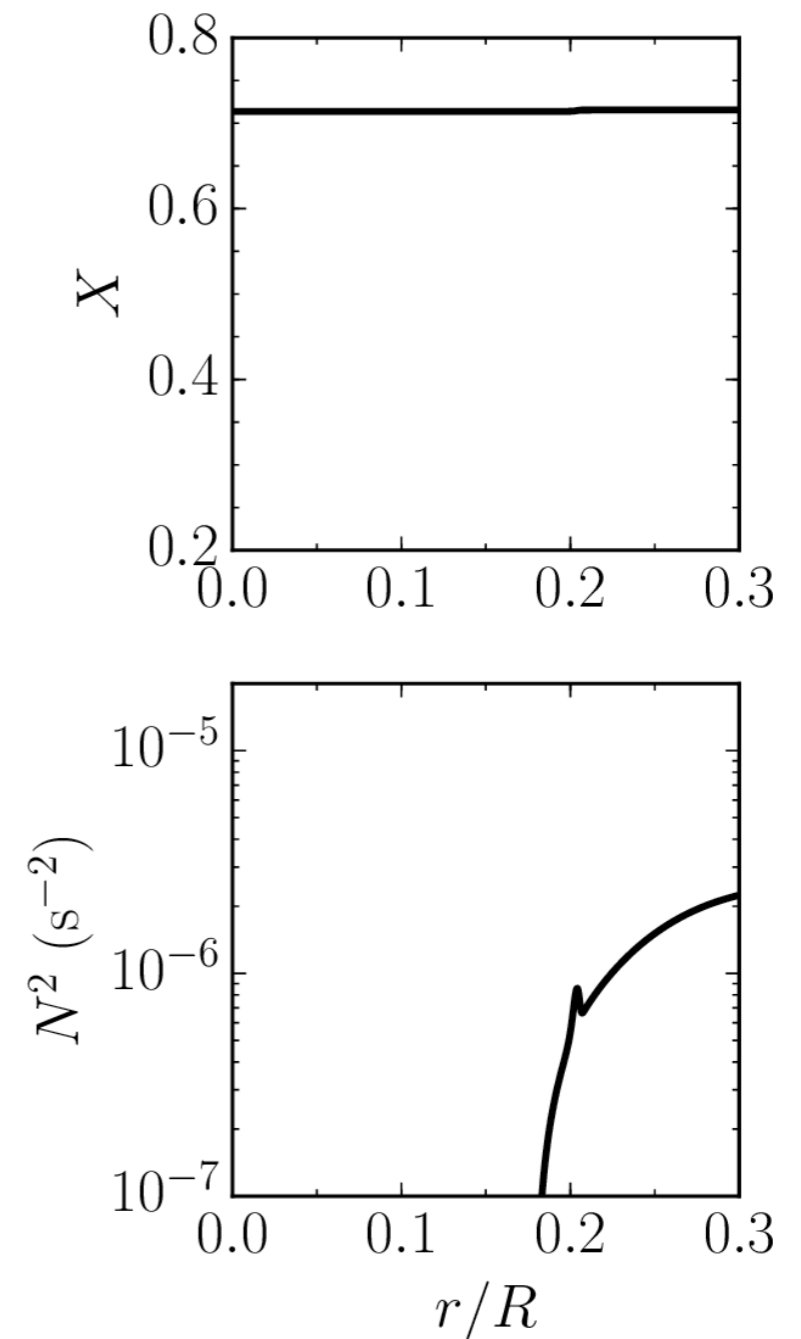
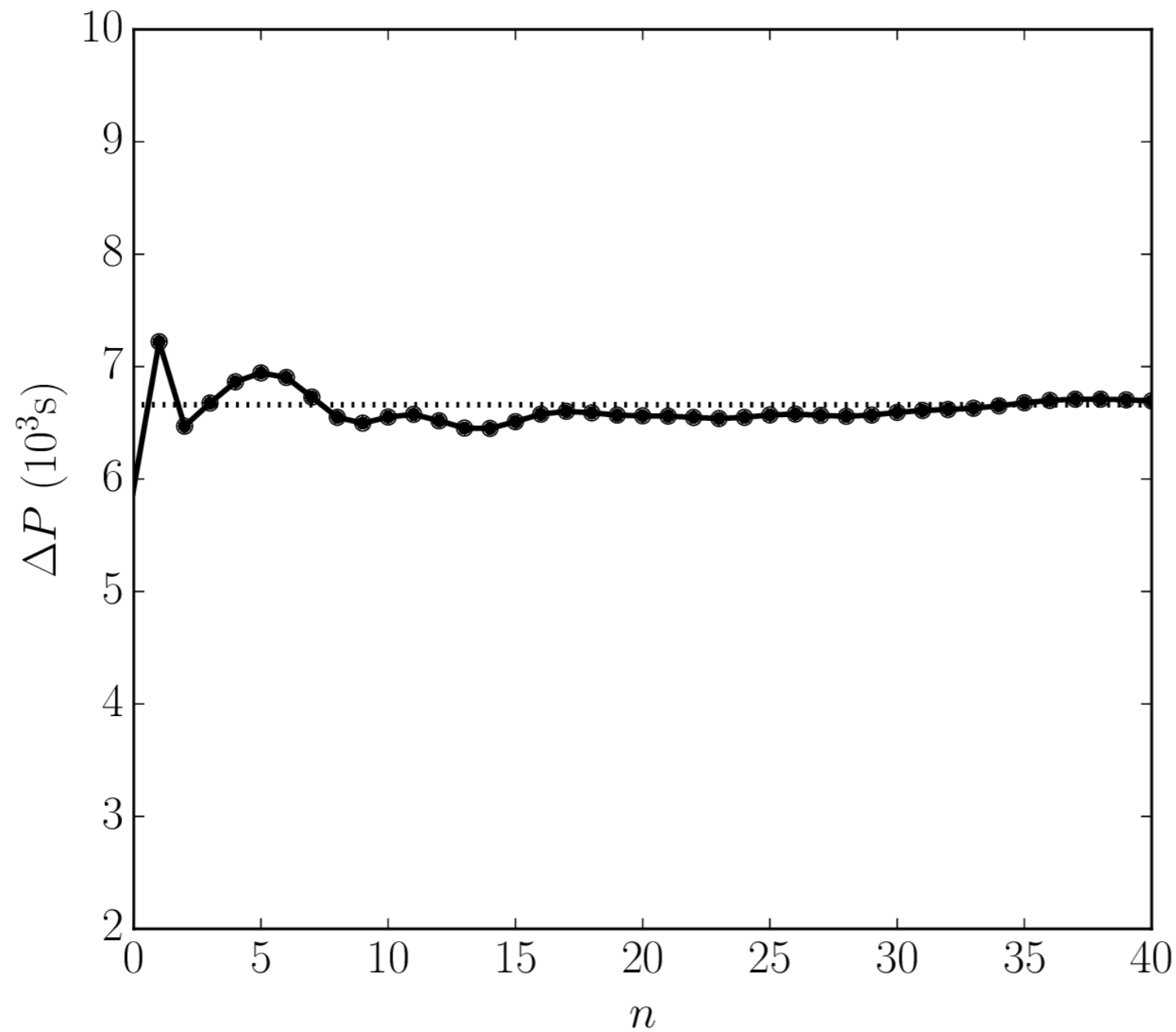
Calculations done with

GYRE

— what real stellar seismologists use!



Period Spacings in a $4.21 M_{\odot}$ SPB Model



...departures from uniform ΔP allow us to probe the core boundary – “real” asteroseismology

