

Electronic Structure of Fe-based superconductors: Magnetism and electron-phonon interaction

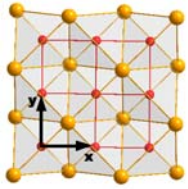
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M. Calandra, F. Mauri, *UPMC, Paris, France*

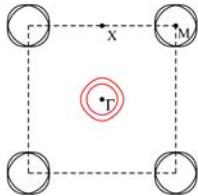
I.I. Mazin, *NRL Washington, USA*

DFT in Fe-based superconductors: What is special?



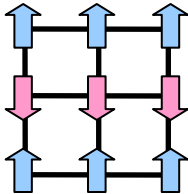
Chemistry:

Two details are crucial: presence and geometry of Fe(Pn/Ch) planes + band filling (d^6).



Non-magnetic (NM) Electronic Structure:

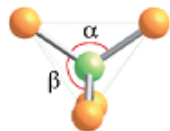
Presence of strongly nested hole-electron pockets, with nearly free-electron like masses.



Magnetism:

LSDA gives the correct AFM ground state; the value of the magnetic moment, and the doping dependence, is overestimated.

WHY?



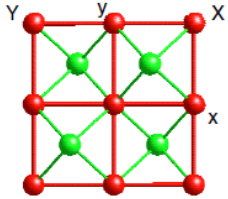
FeAs₄ tetrahedron

Lattice Properties:

The e-ph coupling estimated in the NM case is weak; BUT there is a strong sensitivity of the lattice properties to magnetism and vice-versa.

What are the consequences on the e-ph interaction?

Non-magnetic Electronic Structure

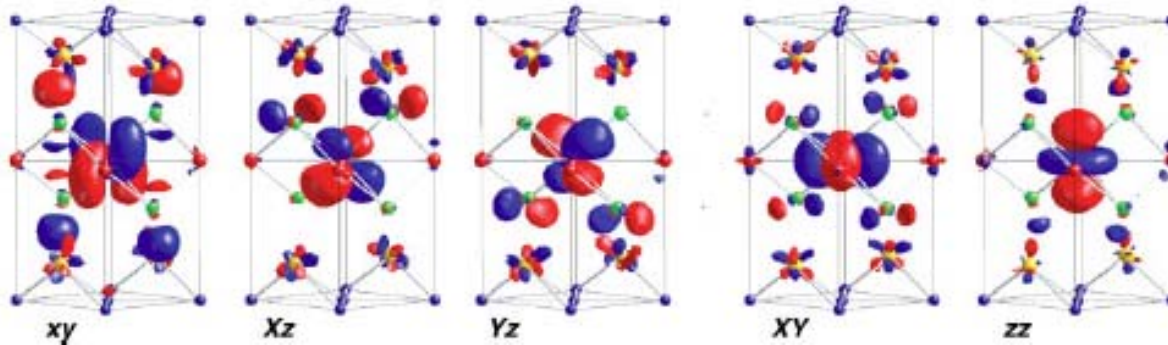


Geometry of the Fe-As planes:

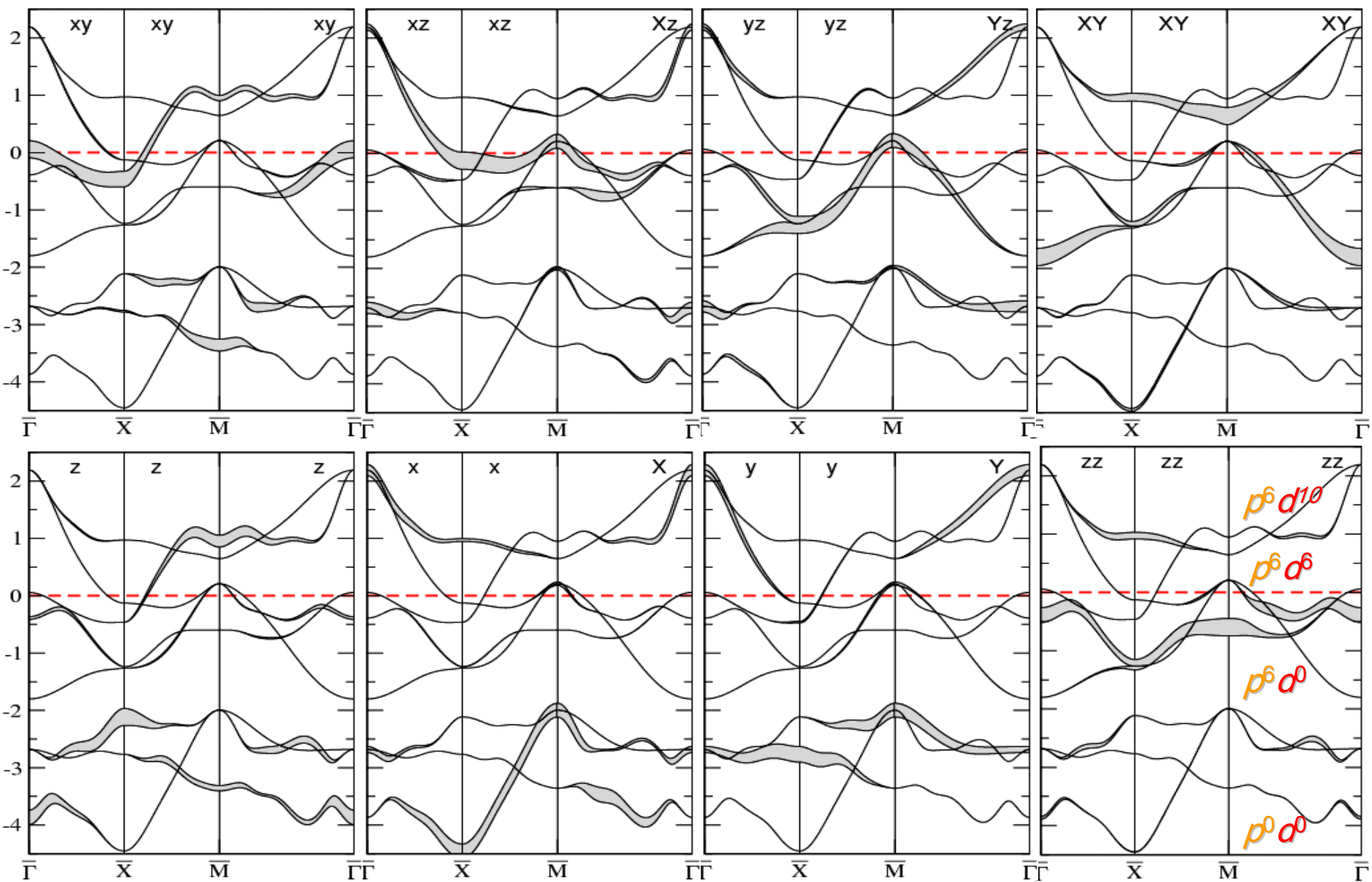
$d-p$ + direct $d-d$ hopping

$d-p$ hopping is large and sensitive to tetrahedral distortion

+ small (e-t) crystal field splitting (max crystal field splitting is (e-t) is 0.2 eV).

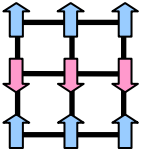


Long hopping ranges due to incomplete pairing of orbitals, i.e. the bonding is partly covalent and partly metallic.



Non-magnetic band structure of LaOFeAs, from ab-initio p-d TB model.

DFT, magnetism:

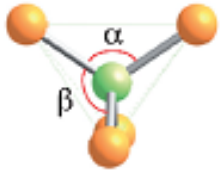


The ground state is a **striped AFM metal** ($m \sim 2.0 \mu_B$).

The **magnetic** transition is accompanied by a **structural** distortion.

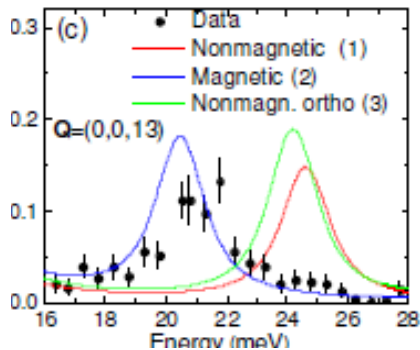
GOOD, BUT

Electrons: The tendency to magnetism is **overestimated** in DFT; the calculated magnetic moments are 2-5 times larger than those measured by neutrons, μ_{SR} ; magnetism survives at high dopings, where the samples are **paramagnetic**.



FeAs₄-tetrahedron

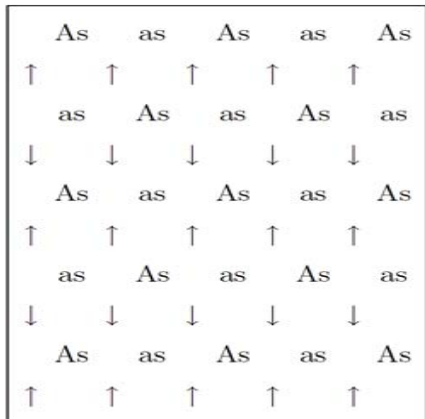
Structure: DFT calculations reproduce the “right” exp structure only if large magnetic moments are present; the FeAs₄ **tetrahedral angle** is very sensitive to the value of the magnetic moment.



Phonons: Phonon frequencies of both paramagnetic and SDW samples are well reproduced only from spin-polarized calculations with large magnetic moments...

Zbiri et al., PRB 2008; Yildirim, Physica C 2009, Reznik et al., PRB 2009

How can we understand the LSDFT magnetic response?



$$\mathbf{H}_{stripe} = \begin{pmatrix} \mathbf{H}(k_x, k_y) & \frac{1}{2}\Delta \\ \frac{1}{2}\Delta & \mathbf{H}(k_x, k_y + \pi) \end{pmatrix}$$

with $\Delta = \Delta \times$

SDW
Hamiltonian

Exchange
splitting

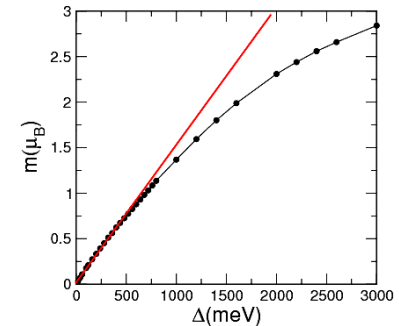
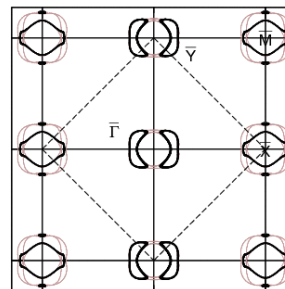
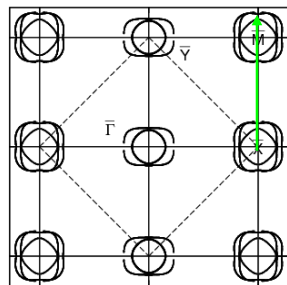
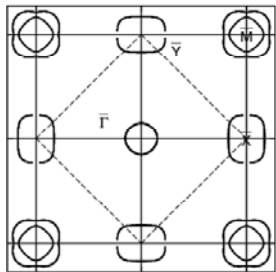
1	0	0	0	0	0	0	0
0	1	0	0	0	0	0	0
0	0	1	0	0	0	0	0
0	0	0	1	0	0	0	0
0	0	0	0	1	0	0	0
0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0

$$m(\Delta) = \sum_{\alpha=1}^5 m_{\alpha}(\Delta) = \sum_{\alpha=1}^5 \sum_{i\mathbf{k}} c_{\alpha, i\mathbf{k}} c_{\alpha, i\mathbf{k}+\mathbf{Q}}^* \quad \Delta = m(\Delta) I$$

I is the Stoner- or Hund's-rule exchange coupling constant for Fe

Magnetic moment

Self-consistency
condition



1) Start from PM bands

2) Fold in ...

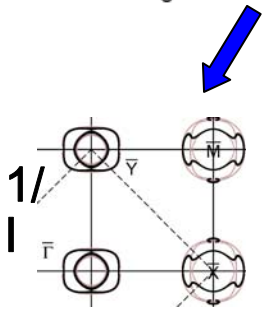
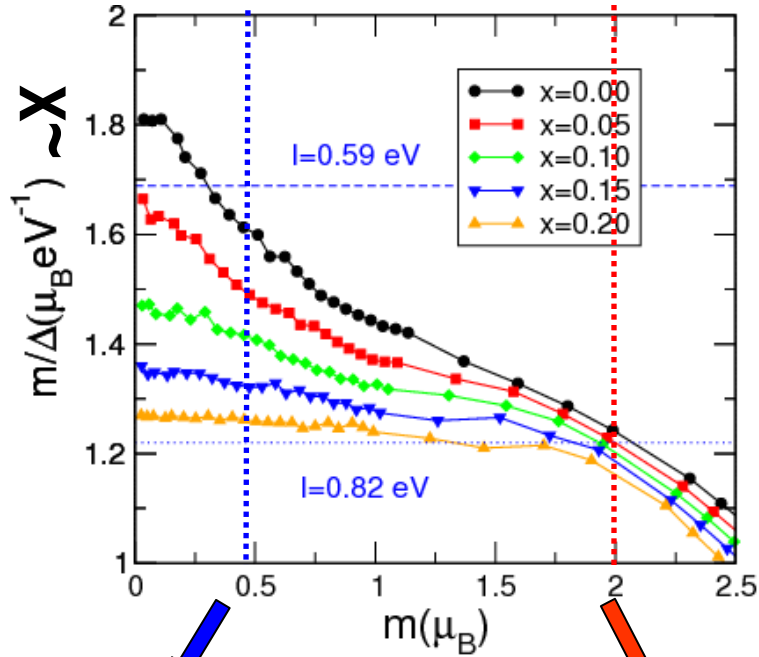
3) Couple with Δ ...

4) Compute $m(\Delta)$

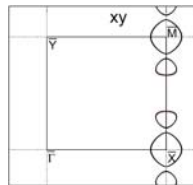
5) Solve $m(\Delta) = \Delta/I$

Itinerant Magnetism: linear response and strong-coupling regime

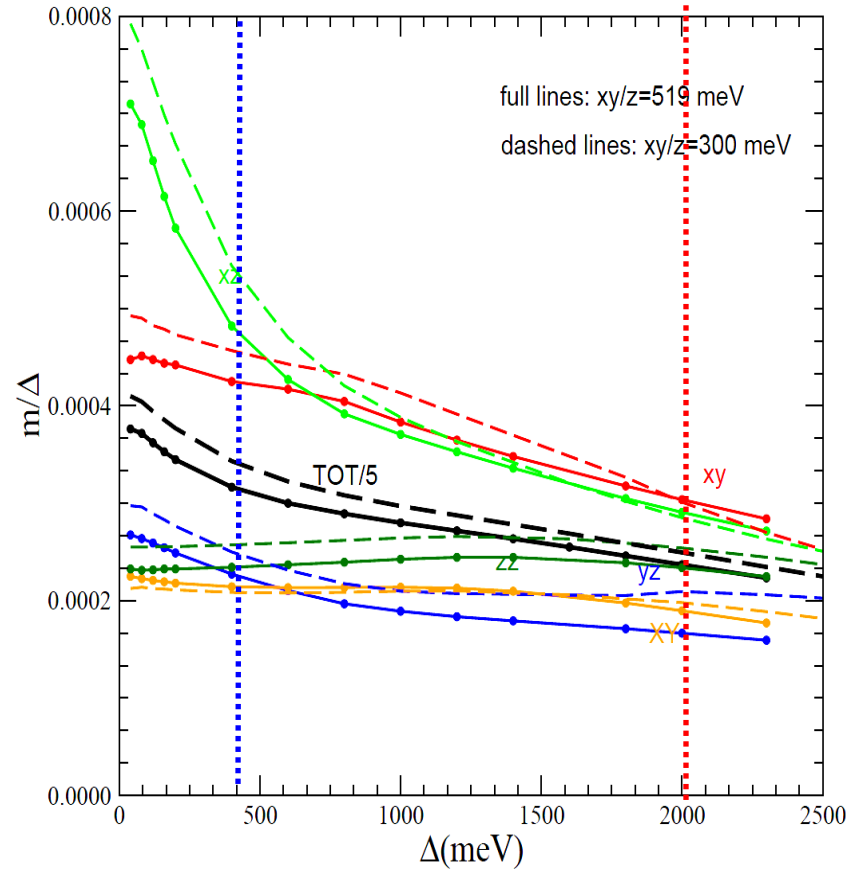
$$\Delta = mI$$



Linear response (nesting)



Intermediate coupling (LSDFT)



Orbital polarization is small and decreases with Δ !

In summary...

Non-Magnetic Electronic Structure: Long-range p-d interactions + Small crystal field splittings: hard to extract localized TB models of the band structure with less than 5 d orbitals; strong sensitivity of the band structure to tetrahedral angle.

Magnetism: LSDFT corresponds to an **intermediate coupling** regime – no nesting in strict sense; Orbital polarization is moderate -> hard to extract localized models.

What can we say about superconductivity?

Spin Fluctuations:

Mostly sensitive to intra-orbital χ of bands near E_F : Strong material dependence!

Mazin et al., PRL 2008, Kuroki et al., PRL 2008, Korshunov and Eremin, PRB 2008, Graser, Maier, Hirschfeld, Scalapino PRB 2008

Electron-Phonon Coupling:

Intrinsically **weak** due to non-bonding character of bands at E_F ;

Boeri et al., PRL 2008; Mazin et al., PRL 2008; Subedi et al., PRB 2008.

BUT

could be relevant for orbital fluctuation models, or to “switch” between almost degenerate gap symmetry solutions...

What can we say about the magnitude of the e-ph interaction in Fe pnictides?

Step 0 (2008): Estimate the magnitude of the e-ph coupling in the **non-magnetic state**; The coupling is low -> WHY?

BUT

The description of the lattice properties in non-spin polarized calculations is bad

The description of the lattice properties in spin-polarized calculations with large moments is good..



Local moment scenario

Step 1 (2010): Estimate the magnitude of the e-ph coupling in the **AFM state**; understand the relation between phonon softening, e-ph coupling, sensitivity of the magnetic moments to the crystal structure.



Step 2 (2010): Construct a model for the **Paramagnetic** state.

Step 0: Estimate the coupling in the NM state:

PRL 101, 026403 (2008)

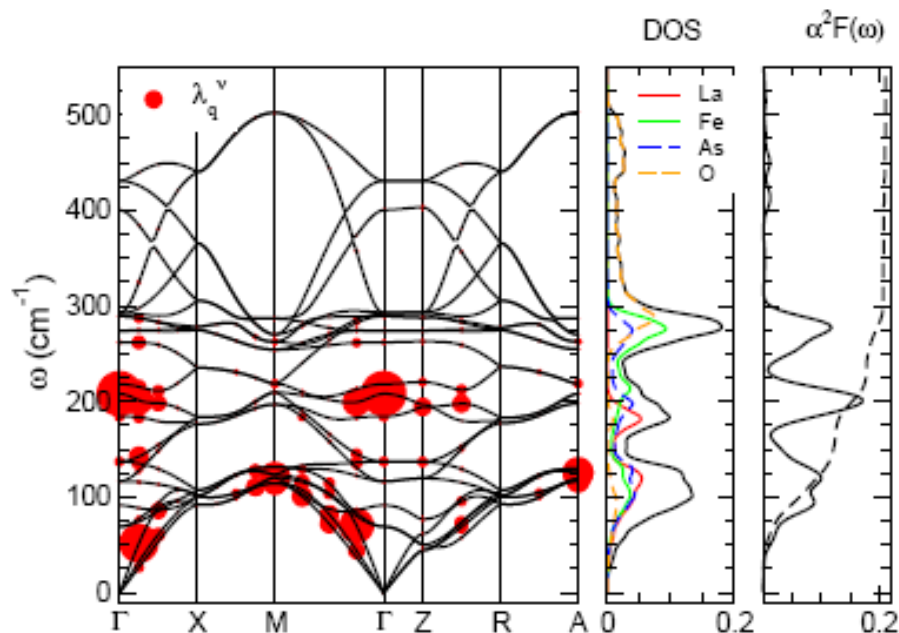
PHYSICAL REVIEW LETTERS

week ending
11 JULY 2008

Is $\text{LaFeAsO}_{1-x}\text{F}_x$ an Electron-Phonon Superconductor?

L. Boeri,¹ O. V. Dolgov,¹ and A. A. Golubov²

NO!



$$\lambda = 0.21$$

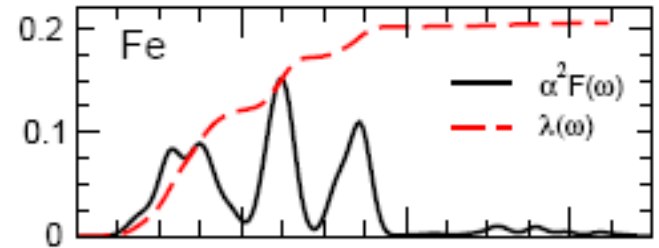
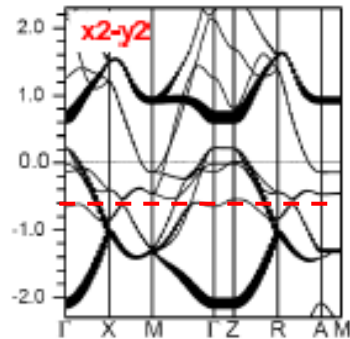
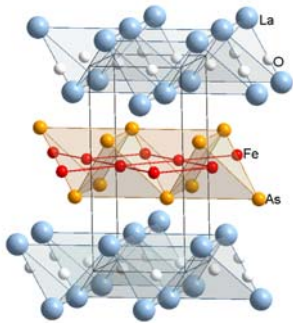
$$\omega_{\log} = 206\text{K}$$

$$T_c^{ME}(\mu^* = 0) = 0.5\text{K}$$

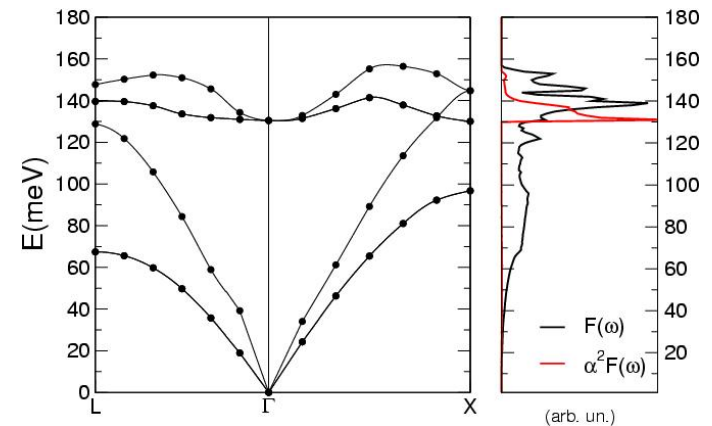
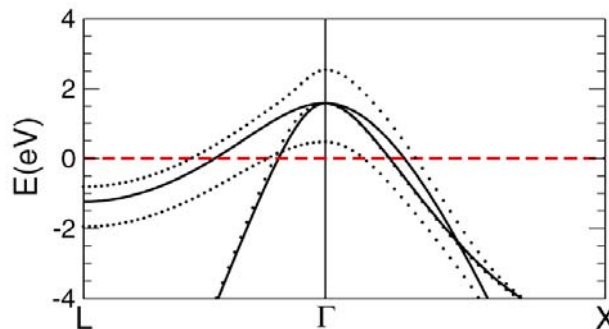
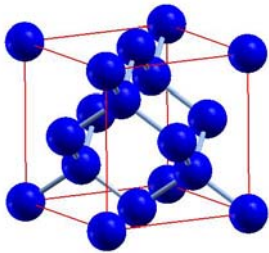
- The coupling is uniformly distributed over Fe-As modes.
- **E-ph interaction**: $\lambda \sim 0.2$ is **too small** to account for T_c
- $N(0)$ is large, λ is small because of **small e-ph matrix elements**.

... and what about phonons: why is the coupling so small ?

Because the only bands derived from **directed Fe-Fe bonds** sit far from E_F ...



In contrast to good e-ph superconductors, where strongly covalent bonds couple to bond-stretching phonons...



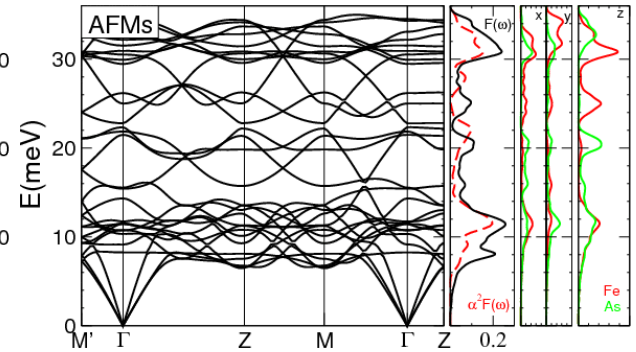
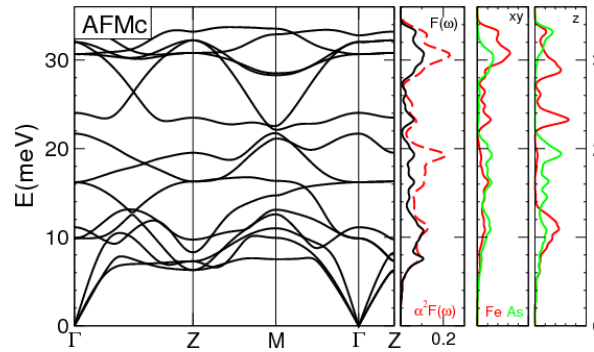
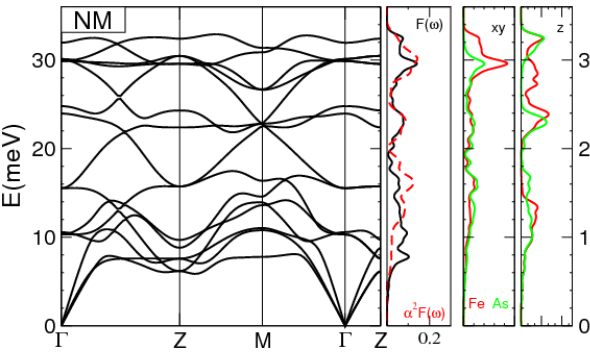
Hole-doped diamond is an extreme example ...

Step 1: Effect of Magnetism:

 **Non-magnetic (NM)**

 **Checkerboard (AFMc)**

 **Stripe (AFMs)**



$$\lambda = 0.18$$

$$N(0) = 1.18 \text{ev}^{-1}$$

$$m = 0$$

$$\lambda = 0.33$$

$$N(0) = 1.36 \text{ev}^{-1}$$

$$m = 2.4 \mu_B$$

$$\lambda = 0.18$$

$$N(0) = 0.68 \text{ev}^{-1}$$

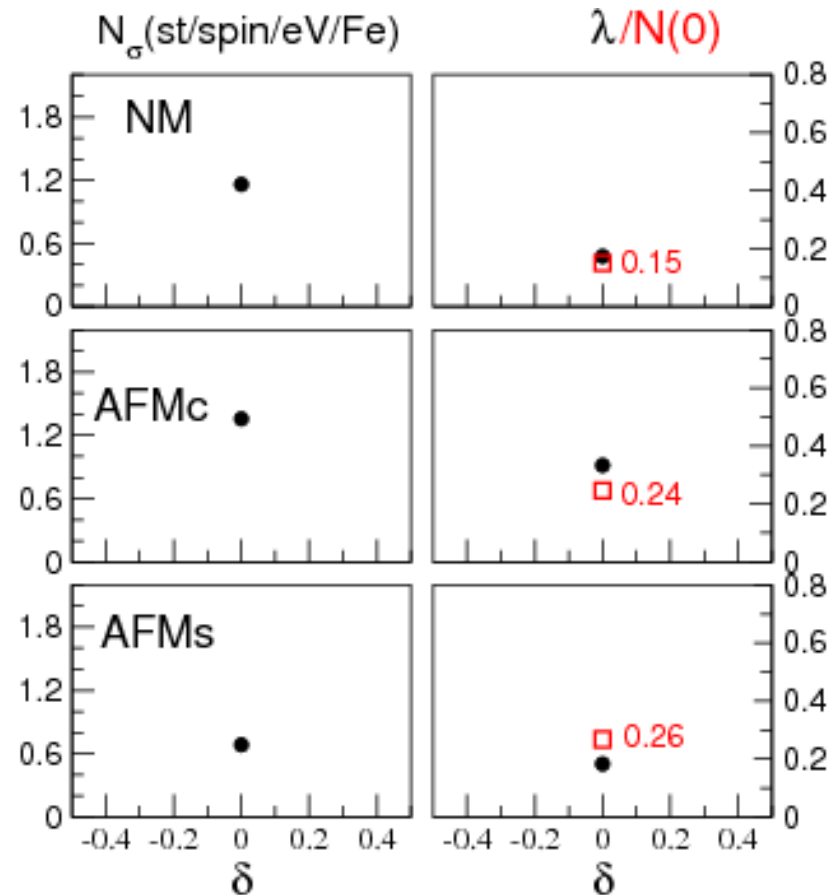
$$m = 2.6 \mu_B$$

Phonon Dispersions: AFM causes a strong shift of Fe and As modes (change of Fe-As bonding); in-plane polarized modes are sensitive to the direction of the AFM wave-vector.

Electron-phonon coupling: AFM and NM Eliashberg functions have the same three-peak structure; the values of λ are too small to explain T_c .

$$\lambda \approx \frac{N(0)I^2}{M\omega^2}$$

Electron-Phonon Coupling:



AFMc and AFMs have larger $\lambda/N(0)$ than NM...



Can we use this information to estimate the coupling in the paramagnetic state?

Paramagnetic State

Large local moments

+

Disorder

LDA: AFM phonon frequencies
(bonding)

NM Fermi surfaces
(wavefunctions)

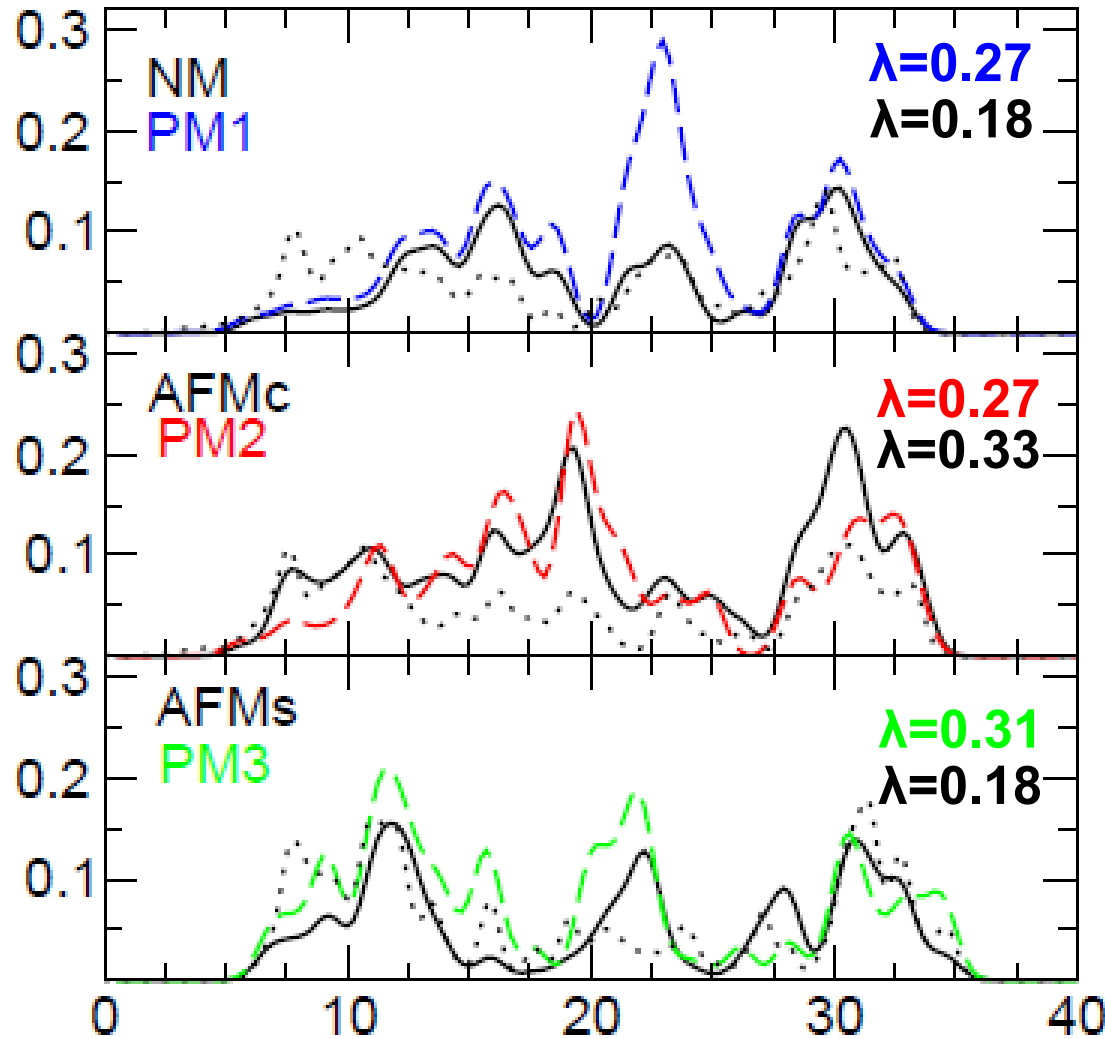


E-ph properties are given by **AFM potentials** and **frequencies** and **NM wavefunctions**

$$\alpha^2 F(\Omega) = \frac{1}{N(0)} \sum_{\mathbf{k}\mathbf{k}',\alpha} \delta(\varepsilon_{\mathbf{k}}) \delta(\varepsilon_{\mathbf{k}'}) |g_{\mathbf{k}\mathbf{k}'}^\alpha|^2 \delta(\Omega - \omega_{\mathbf{k}-\mathbf{k}'})$$

$$g_{\mathbf{k}+\mathbf{q}\mathbf{j}',\mathbf{k}\mathbf{j}}^{\mathbf{q}\nu} = \langle \mathbf{k} + \mathbf{q} \mathbf{j}' | \delta^{\mathbf{q}\nu} V_{sc} | \mathbf{k} \mathbf{j} \rangle$$

"Paramagnetic" model:

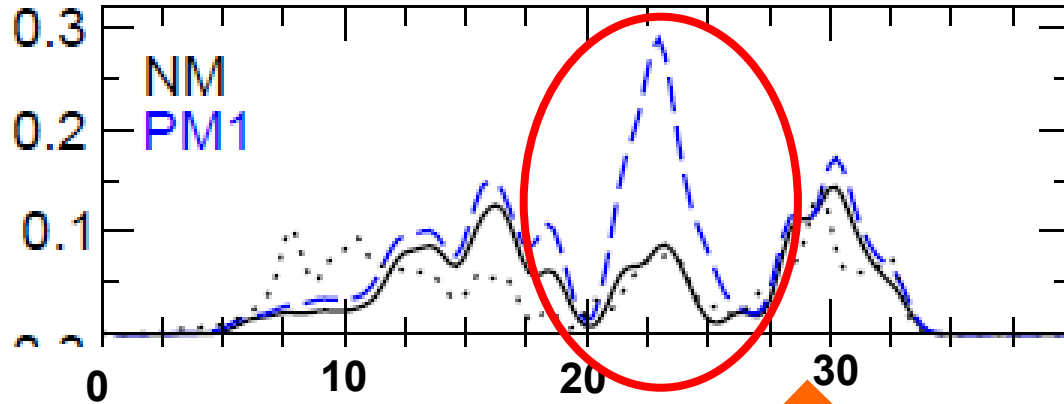


PM1:
NM wf + AFMc V + NM ω

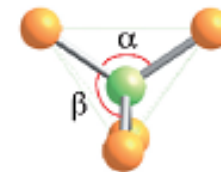
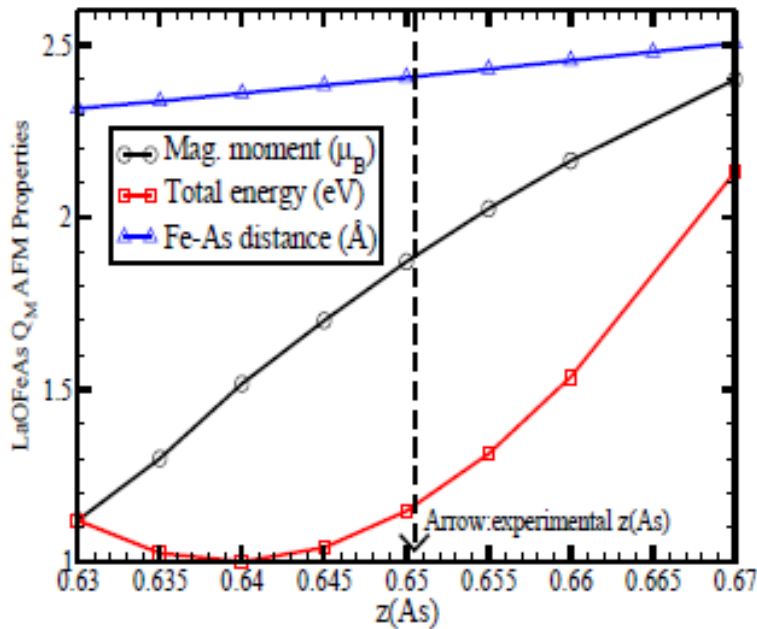
PM2:
NM wf + AFMc V + AFMc ω

PM3:
NM wf + AFMs V + AFMs ω

Where does the increase in λ come from?



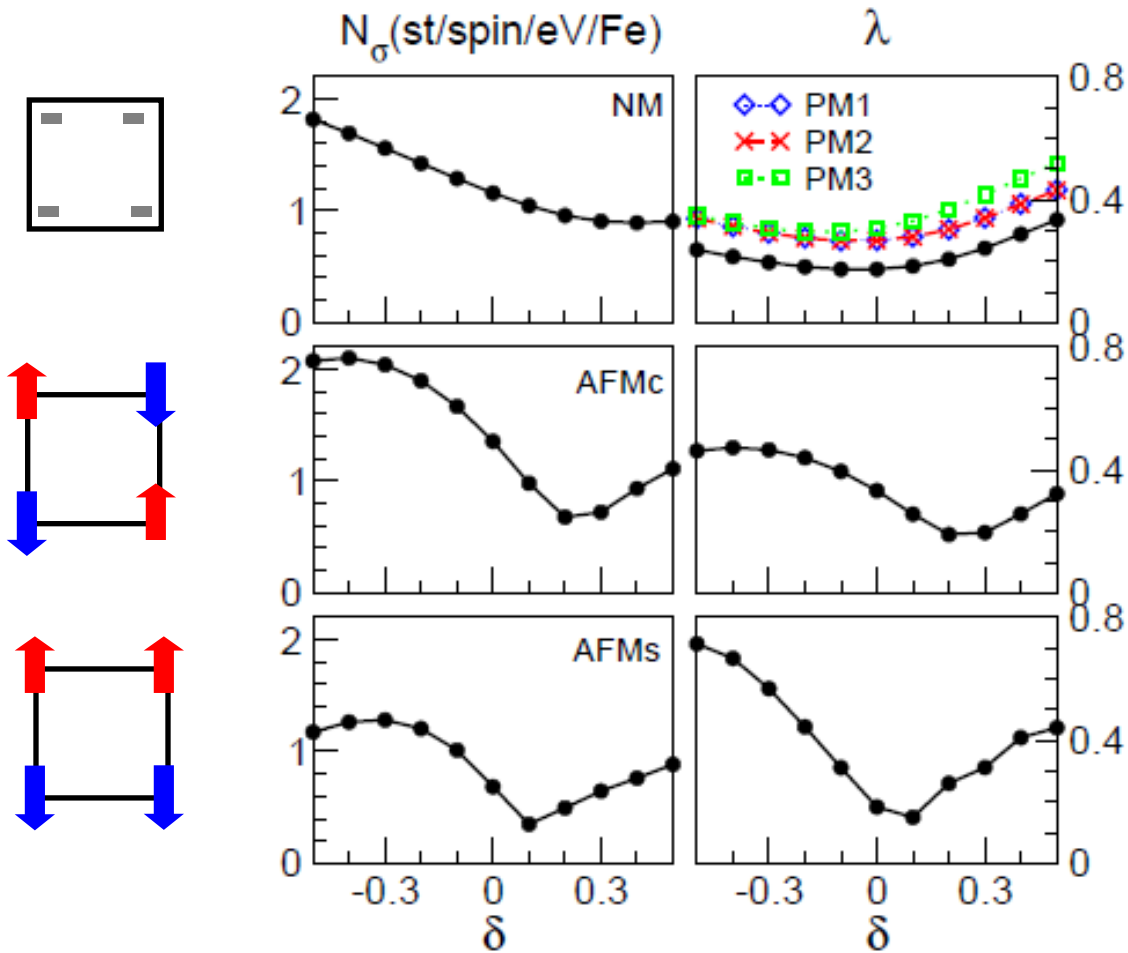
The modes at 20 meV modulate the FeAs₄ tetrahedra...



FeAs₄-tetrahedron

Which in turn modulate the magnetic moment!

Results:



Paramagnetic model: AFM potentials cause an almost uniform 50 % increase of the coupling.

Effect of doping: (Rigid-band) The coupling does not simply follow the shape of the DOS – matrix elements are important!

In the relevant range of dopings the e-ph coupling is always ≤ 0.35 !!!

... In Summary...

Early estimates of the **e-ph coupling** in Fe-based superconductors are **extremely small ($\lambda \sim 0.2$): non-magnetic (NM) calculations!**

BUT

Can we trust these estimates if NM calculations give a poor description of lattice dynamics?



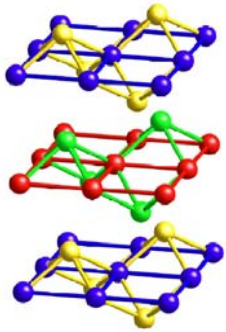
Linear response calculations of lattice dynamics and e-ph coupling in Ba-122

? What is the **effect of magnetism** on the phonon dispersion and e-ph coupling constant? *Strong **renormalization of Fe-As modes** (change of bonding); ~ 50 % increase of the e-ph coupling matrix element.*

? Can we model the **paramagnetic** state? *Combine AFM potentials and NM wavefunctions.*

? Can we estimate an upper bound for λ in Fe-based SC?

For realistic values of doping $\lambda \leq 0.35$!



What next?

More than 300 **RETM_xP_nO** compounds exist ...

Electronic and magnetic properties of Mn, Fe, Co, Ni and Zn in LaTPnO

<i>TM</i>	Mn		Fe		Co		Ni		(Cu)	Zn	
<i>Pn</i>	P	As	P	As	P	As	P	As		P	As
Elect. Prop.	Mott Insulator		Superconductor				Metal		—	Semiconductor	
Magnetism	AFM						FM		—	Non-magnetic	
E_g	-1 eV		—				—		—	-1.5 eV	
T_c (SC)	> 400 K		Undoped: 4 K	F-doped: 26 K	43 K	66 K	Undoped: 3 K	Undoped: 2.4 K		—	
T_{Nc} (Mag.)											
Ref.	Yanagi et al. JAP (submitted)		Kamihara et al. JACS(2006), Kamihara et al. JACS (2008)		Yanagi et al. PRB (2008)		Watanabe et al. IC (2007), Watanabe et al. JSSC (2007)		—	Kayanuma et al. PRB (2007), Kayanuma et al. TSG (2008)	

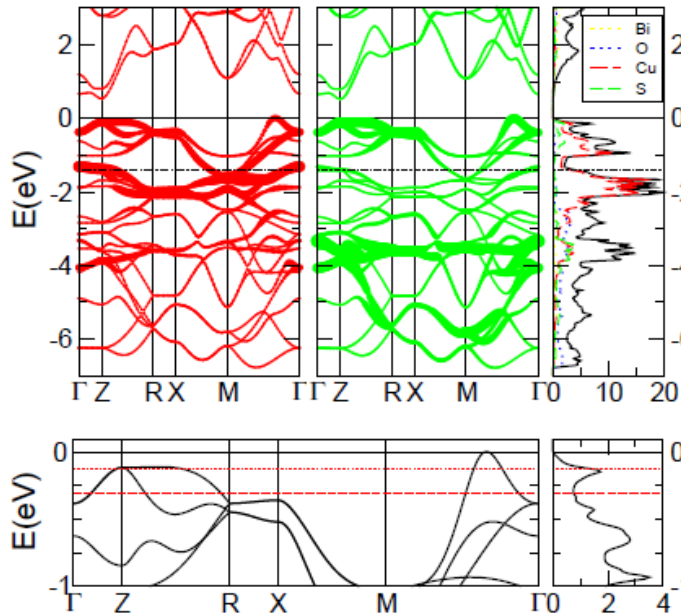
Fig. 10. Electromagnetic and optical properties elucidated in LaT_xPnO.

Due to the rich electronic structure the magnetic and SC properties can be tuned by doping/pressure...

TM pnictides represent an ideal playground to study the interplay between magnetism, SC and e-ph interaction...

A related example: BiOCu_{1-x}S, cond-mat/0911.5305

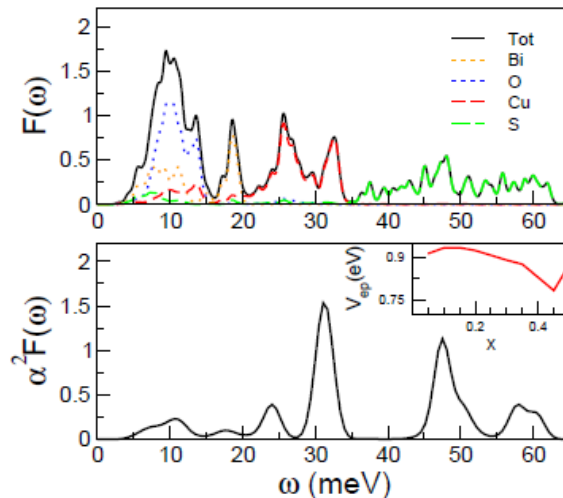
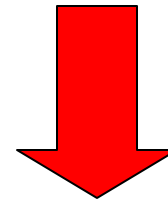
A related example: $\text{BiOCu}_{1-x}\text{S}$, cond-mat/1012.5761



Same crystal structure as LaOFeAs.
 Cu-S layers play the role of FeAs layers.
 Superconductivity is observed at $x=0.5$
 Electron count is d^{10-x} .
 The Fermi level sits in the S antibonding region.

LSDFT gives:

-> Instability to magnetism (FM) (Cu)
 Large e-ph interaction (S/Cu hybridization)



Competition of SF (FM) and e-ph interaction?

References:

- ➔ **On the e-ph interaction in Fe and Ni-based pnictides:**
L. Boeri, O.V. Dolgov and A. Golubov, Phys. Rev. Lett 2008 and Physica C 2009.
- ➔ **On problems connected with DFT in FeBSC:**
I.I. Mazin, M. D. Johannes, L. Boeri, K. Koepernik, and D. J. Singh,
Phys. Rev. B 78, 085104 (2008).
- ➔ **On the AFM e-ph coupling in Fe-BSC:**
L. Boeri, M. Calandra, I. I. Mazin, O. V. Dolgov, and F. Mauri,
Phys. Rev. B 82, 020506R (2010)
- ➔ **On the electronic structure and magnetism in Fe-BSC**
O.K. Andersen and L. Boeri, cond-mat/1011.1658.
- ➔ **On the superconductivity of BiOCuS,**
L. Ortenzi et al., arXiv:1012.4761

THANK YOU!

Experimental estimates of the e-ph interaction

$$T_c = \frac{\omega_{\log}}{1.2} \exp\left[-\frac{1.04(1+\lambda)}{\lambda - \mu^*(1+0.62\lambda)}\right]$$

$$T_c = 5 - 55 K$$



$$\lambda \approx 0.5 - 1.5,$$

$$\omega_{\log} \approx 200 K$$

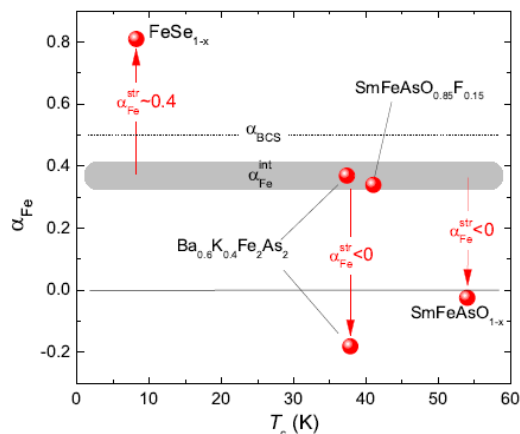
The *e-ph* coupling has been estimated by various **exp techniques**:

- Pump-probe ARPES (*cond-mat/1008.1561*), Pump-probe optics (*PRB* **82**, 024513 (2010)), Raman (*PRB* **78**, 212503 (2008)), ...

$$\lambda \approx 0.2 - 0.4$$

- ARPES (*cond-mat/1002.3149*), Raman (*PRB* **80**, 064509 (2009))

$$\lambda \approx 1.0 - 1.5$$



The isotope effect is also controversial...

Khasanov et al., cond-mat/1008.4540