

# Binding of heavy fermions by a single light atom in one dimension

Andrea Tononi, Jules Givois, and Dmitry Petrov

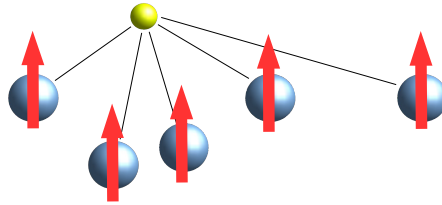
Laboratoire Physique Théorique et Modèles Statistiques (Orsay)

arXiv:2205.01018



# (N+1)-body problem

How many heavy fermions can be bound by a single light atom?



Kinetic energy of the heavy atoms  $\sim 1/M$

competes with

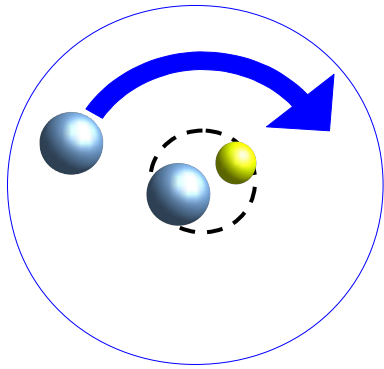
Attractive exchange potential of the light atom  $\sim 1/m$

Parameters of the **free-space zero-range** N+1-body problem:

- space dimension  $D$
- number of heavy atoms  $N$ 
  - mass ratio  $M/m$
- dimer size  $a$  (can be used as the length unit)

# 3D 2+1-trimer

Emergence of a trimer state for  $M/m > 8.2$  [Kartavtsev & Malykh'2006]

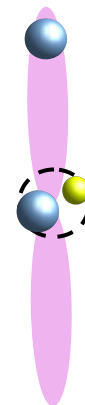
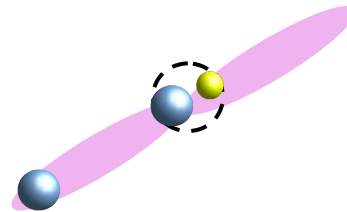
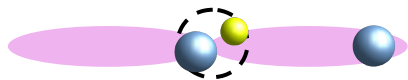


$M/m < 8.2$  p-wave atom-dimer scattering resonance

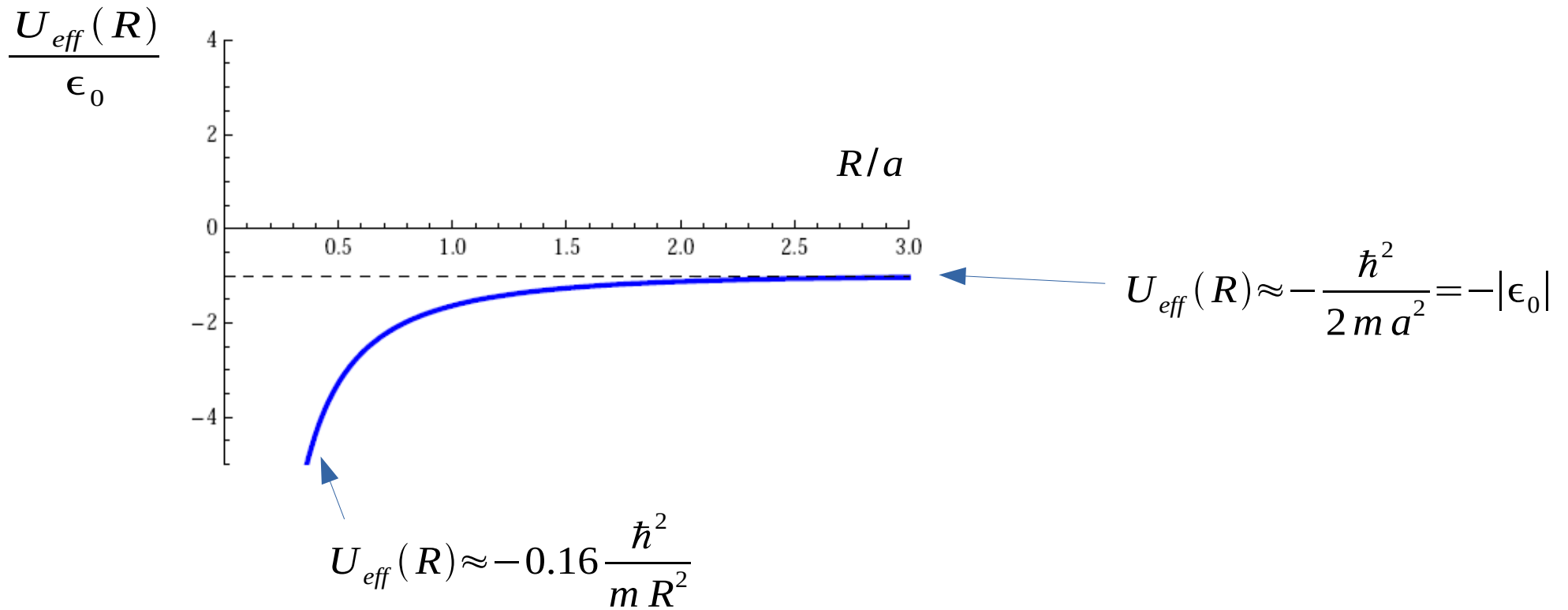
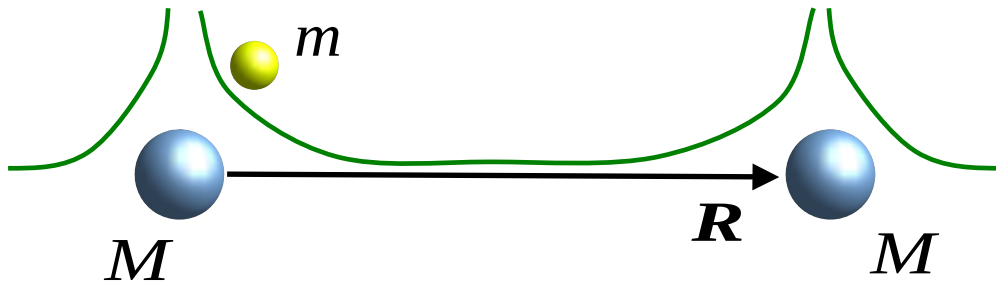
$M/m > 8.2$  trimer state with  $L=1$

**Bound trimer state, NOT EFIMOV**

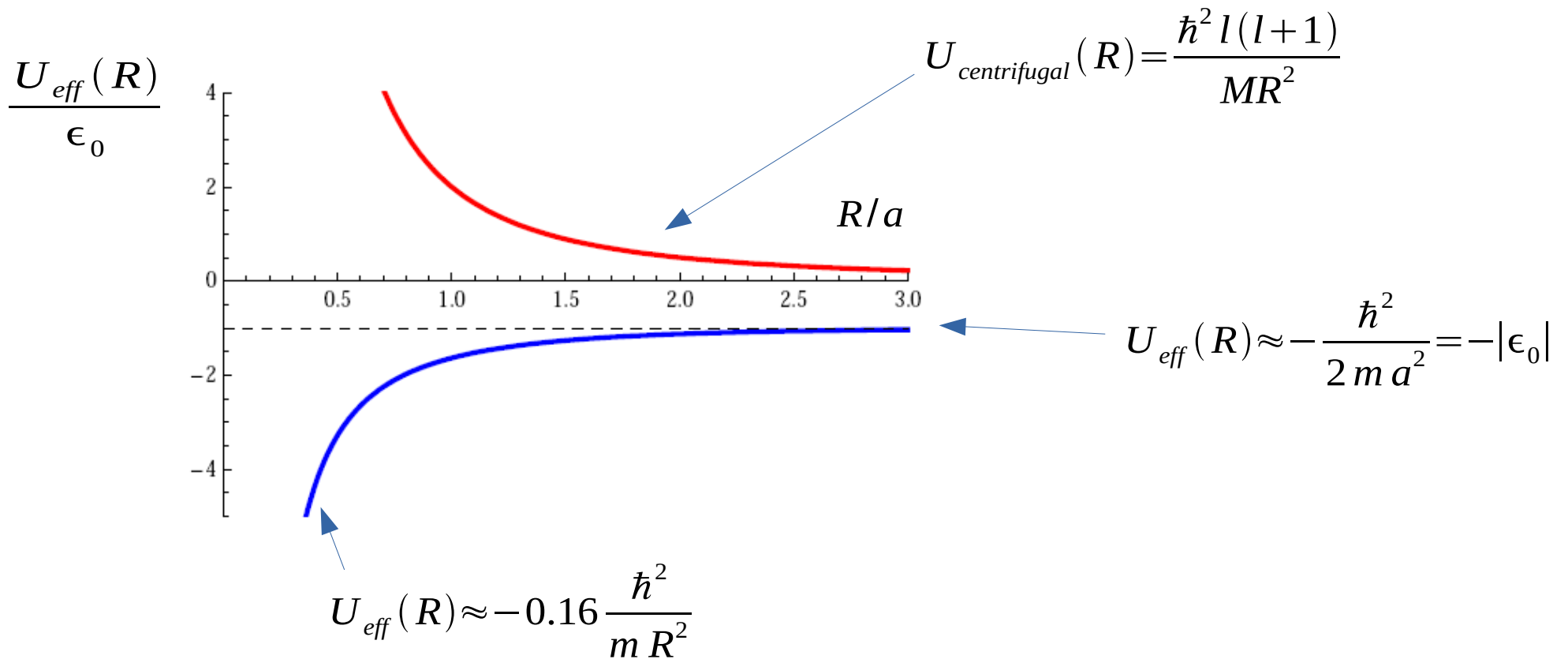
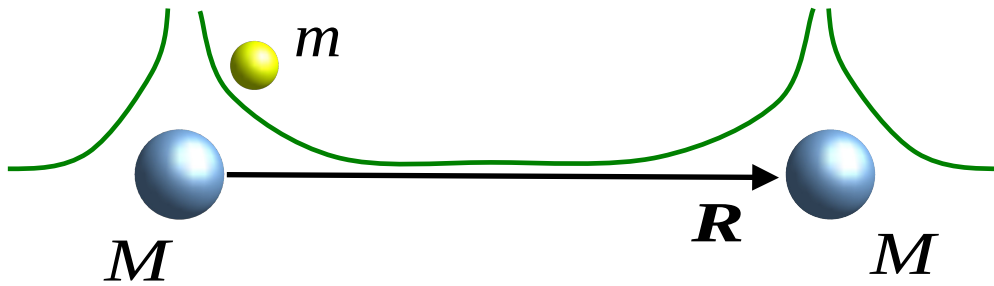
$p_x$ ,  $p_y$ , and  $p_z$  orbitals:



# Born-Oppenheimer picture



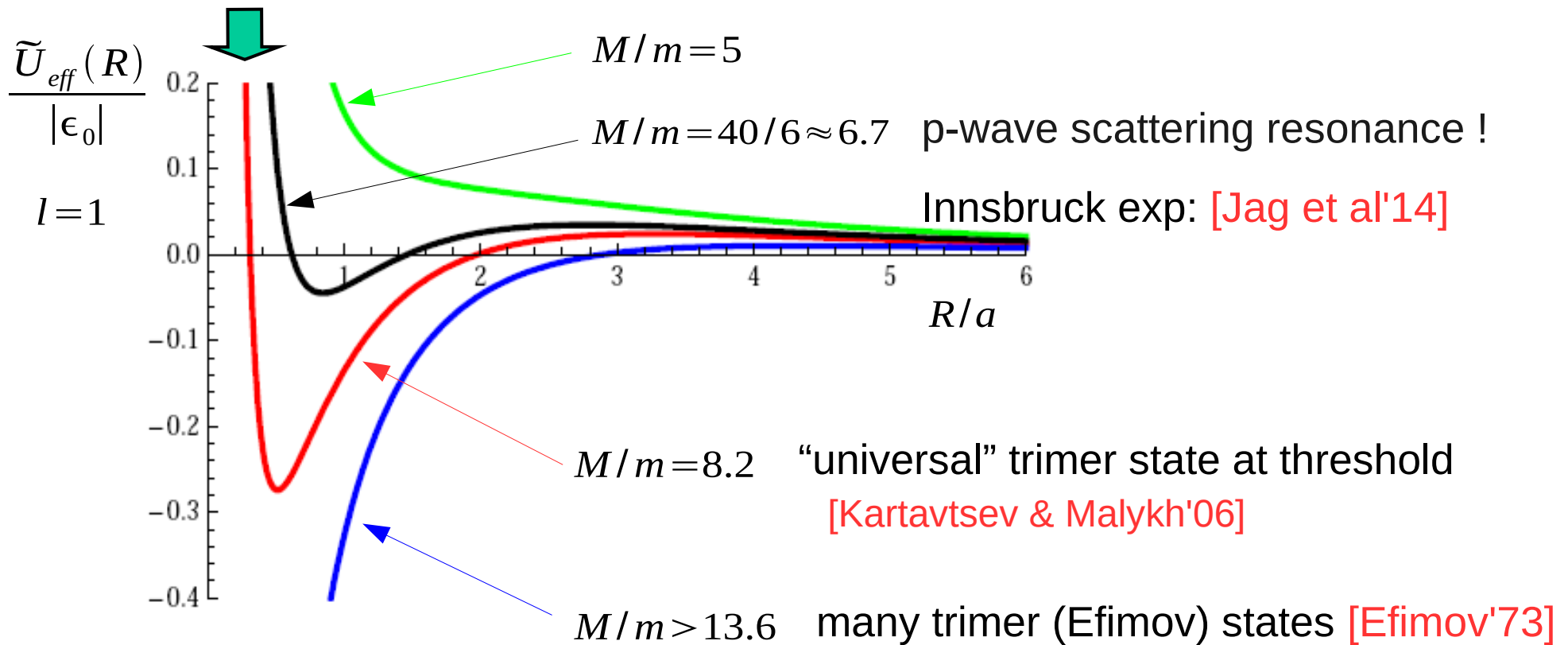
# Born-Oppenheimer picture



$$\left[ -\frac{\hbar^2}{M} \frac{\partial^2}{\partial R^2} + \tilde{U}_{eff}(R) \right] \chi(R) = E \chi(R)$$

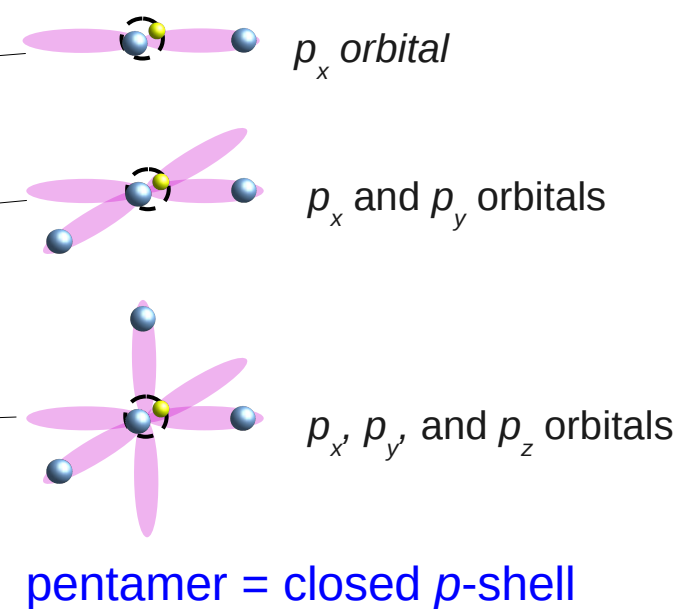
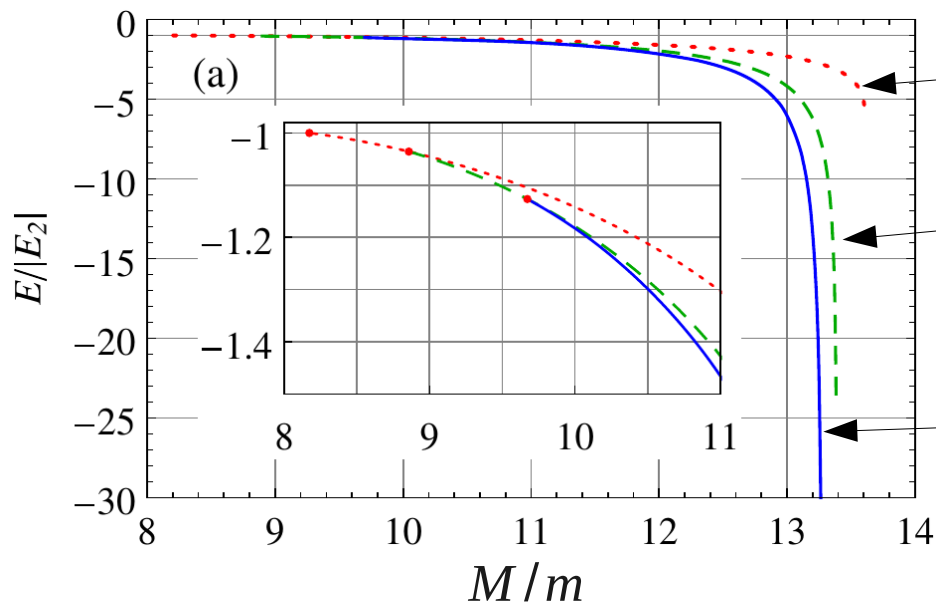
$$\frac{\hbar^2}{MR^2} \left( l(l+1) - 0.16 \frac{M}{m} \right)$$

$$\tilde{U}_{eff}(R) = U_{eff}(R) + |\epsilon_0| + \frac{\hbar^2 l(l+1)}{MR^2}$$



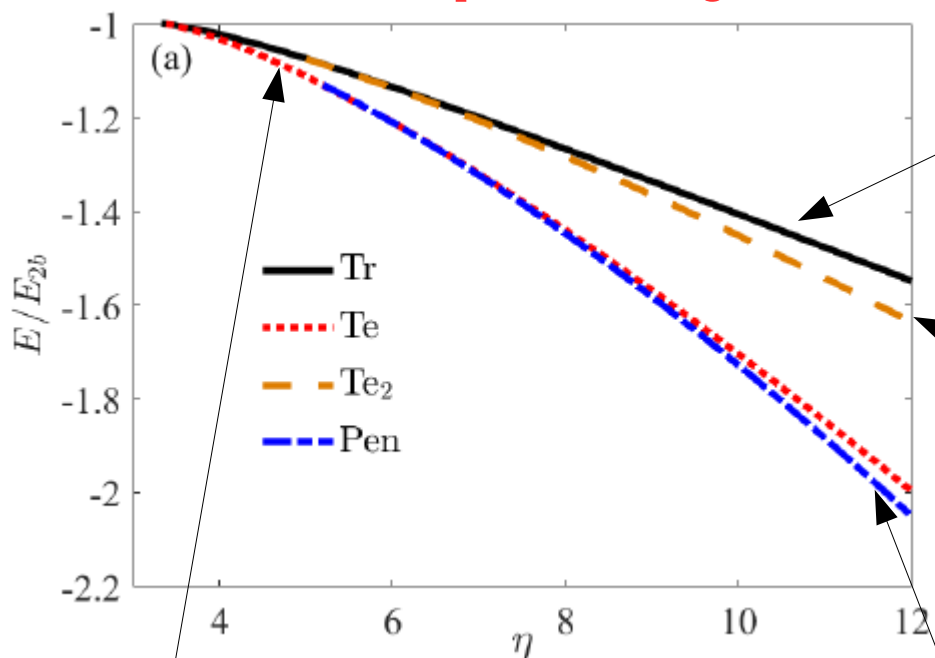
# 3D trimer, tetramer, pentamer,...

	Symmetry $L^\pi$	appear at $M/m >$	Efimovian for $M/m >$
2+1 trimer	$1^-$	8.173 <small>Kartavtsev&amp;Malykh'06</small>	13.607 <small>Efimov'73</small>
3+1 tetramer	$1^+$	8.862(1) <small>Blume'12, Bazak&amp;DSP'17</small>	13.384 <small>Castin,Mora&amp;Pricoupenko'10</small>
4+1 pentamer	$0^-$	9.672(6) <small>Bazak&amp;DSP'17</small>	13.279(2) <small>Bazak&amp;DSP'17</small>
N+1-mer	?	?	?



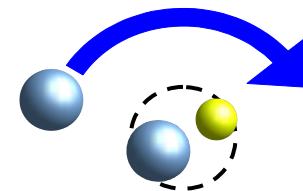
# 2D trimer, tetramer, pentamer...

[Liu & Peng & Cui'22]



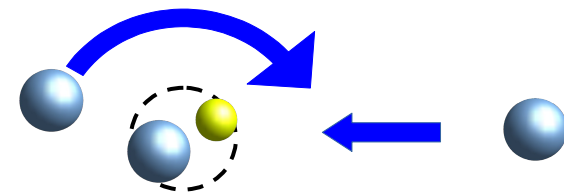
L=1 trimer  $(M/m)_c = 3.33$

[Pricoupenko & Pedri'10]



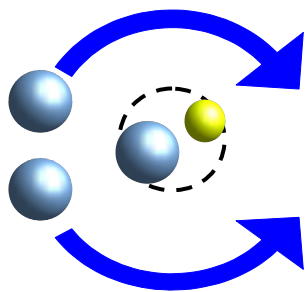
L=1 tetramer  $(M/m)_c^{2D} = 5.0$

[Levinsen & Parish'13]



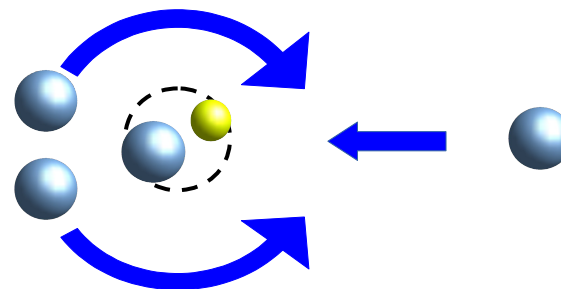
L=0 tetramer  $(M/m)_c = 3.38$

[Liu & Peng & Cui'22]



L=0 pentamer  $(M/m)_c = 5.14$

[Liu & Peng & Cui'22]

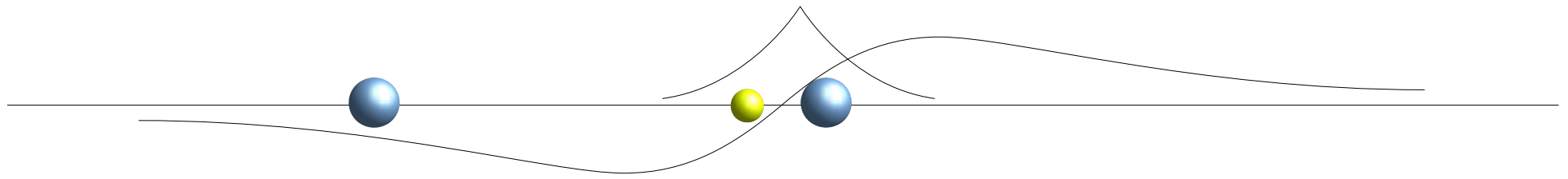




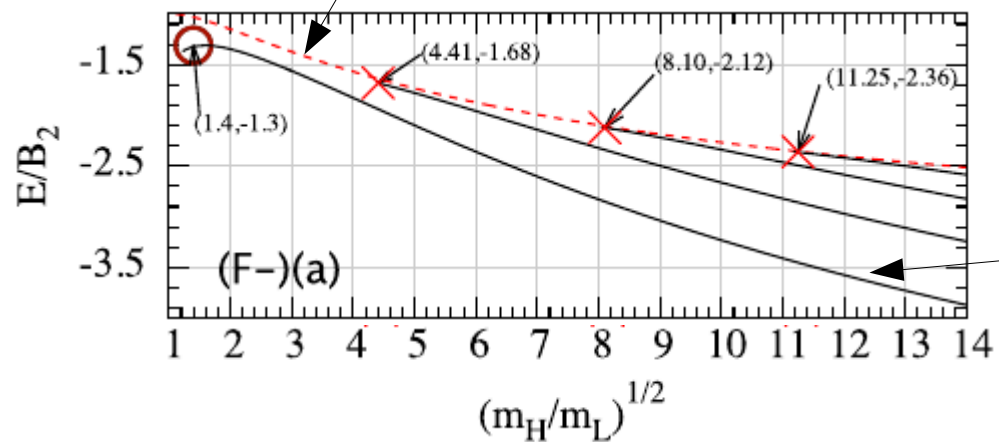
# 1D trimer, tetramer...

1D N+1 problem: no shell effects, no Efimov effect, “easy” calculations...

Ideal testbed for various few-body and many-body approaches!

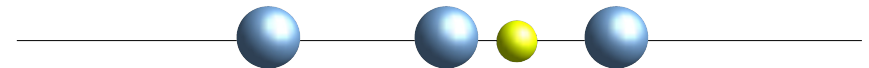


Odd-parity trimer for  $M/m > 1$  [Kartavtsev & Malykh & Sofianos'09]



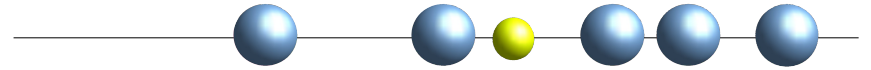
[Mehta'14]

Born-Oppenheimer treatment of the tetramer

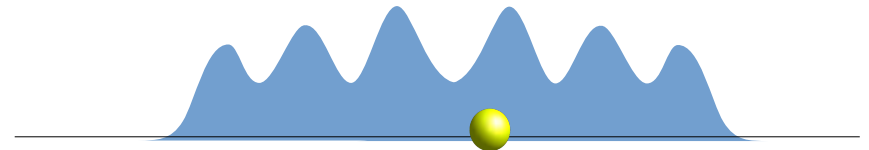
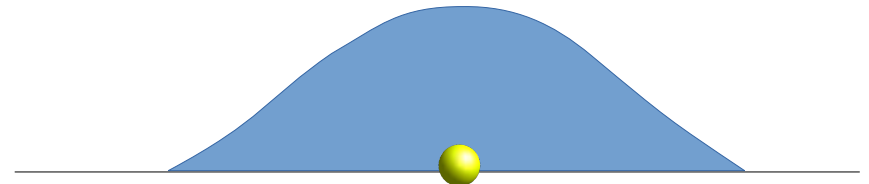


# Outline

- Exact Skorniakov & Ter-Martirosian (STM) calculations for  $N \leq 5$



- Mean-field theory for  $N \rightarrow \infty$ 
  - Thomas-Fermi for the heavies
  - Hartree-Fock for the heavies



- Applying MF for  $N \leq 5$ 
  - trying to understand...
  - trying to improve...

- Summary/Outlook

# STM equation

$$\left[ -\frac{1}{2m} \frac{\partial^2}{\partial x^2} - \sum_{i=1}^N \frac{1}{2M} \frac{\partial^2}{\partial x_i^2} + g \sum_{i=1}^N \delta(x_i - x) - E \right] \Psi(x_1, \dots, x_{N-1}, x_N, x) = 0$$

$$g = -1/(m_r a) < 0 \quad E_{1+1} = -1/(2m_r a^2) \quad m_r = Mm/(M+m)$$

$$\Psi(x_1, \dots, x_{N-1}, x_N, x = x_N) \xrightarrow{\text{Fourier}} F(q_1, \dots, q_{N-1}, q_N) \xrightarrow{q_N = -\sum_{i=1}^{N-1} q_i} F(q_1, \dots, q_{N-1})$$

Skorniakov and Ter-Martirosian equation (STM) [L.Pricoupenko'11, A.Pricoupenko&DSP'19]:

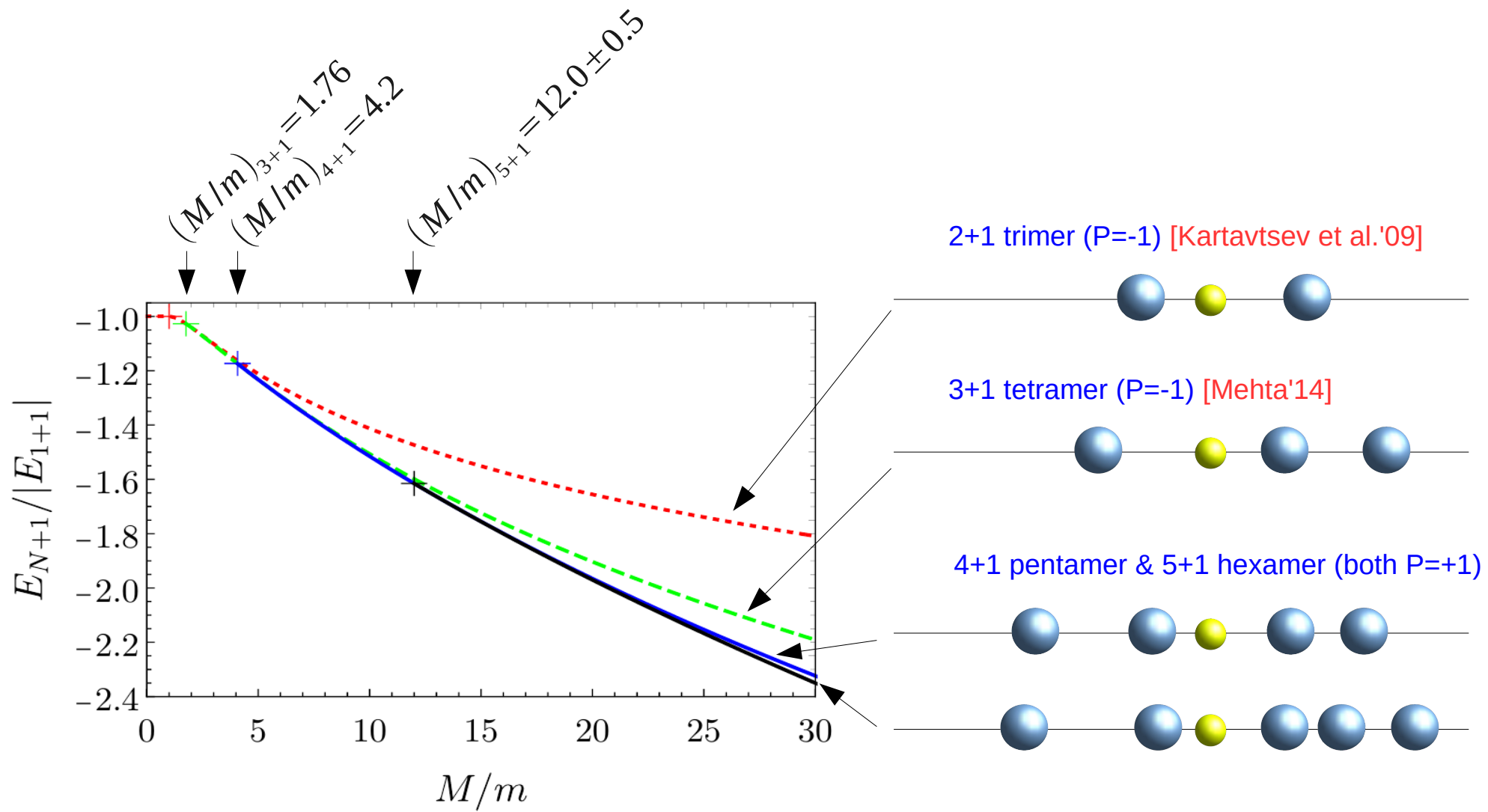
$$\left[ \frac{a}{2} - \frac{1}{2\kappa(q_1, \dots, q_{N-1})} \right] F(q_1, \dots, q_{N-1}) = - \int \frac{dp}{2\pi} \frac{\sum_{j=1}^{N-1} F(q_1, \dots, q_{j-1}, p, q_{j+1}, \dots, q_{N-1})}{\kappa^2(q_1, \dots, q_{N-1}) + (p + \frac{m_r}{m} \sum_{i=1}^{N-1} q_i)^2}$$

$$\kappa(q_1, \dots, q_{N-1}) = \sqrt{-2m_r E + \frac{m_r}{M+m} (\sum_{i=1}^{N-1} q_i)^2 + \frac{m_r}{M} \sum_{i=1}^{N-1} q_i^2}$$

STM:

- natural account of zero-range interactions
- removes one relative coordinate (1+1 dimer problem is analytic)
- ...

# STM results



# Mean field #1 Thomas-Fermi

$$\Omega = \int [|\phi'(x)|^2/2m + gn(x)|\phi(x)|^2 + \pi^2 n^3(x)/6M - \epsilon|\phi(x)|^2 - \mu n(x)] dx$$

Mean field

Kinetic energy in the TF approximation

Lagrange multipliers

$$\int |\phi(x)|^2 dx = 1$$

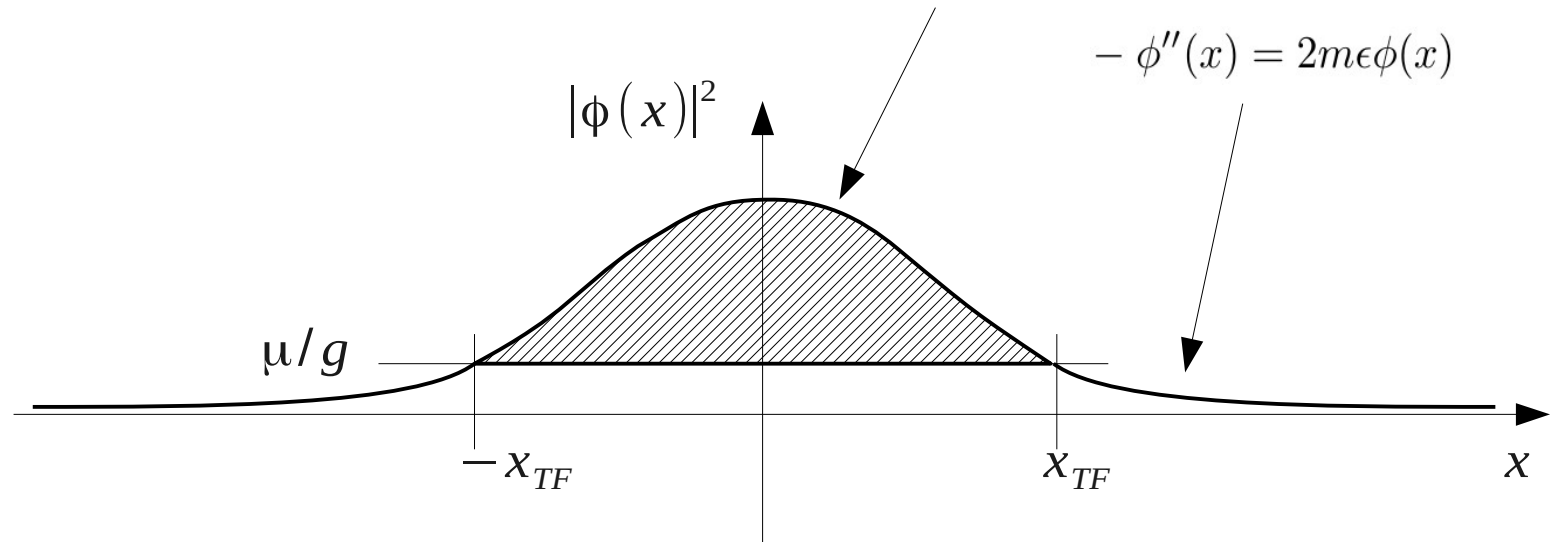
$$\int n(x) dx = N$$

Minimization with respect to  $n(x)$   $\longrightarrow$

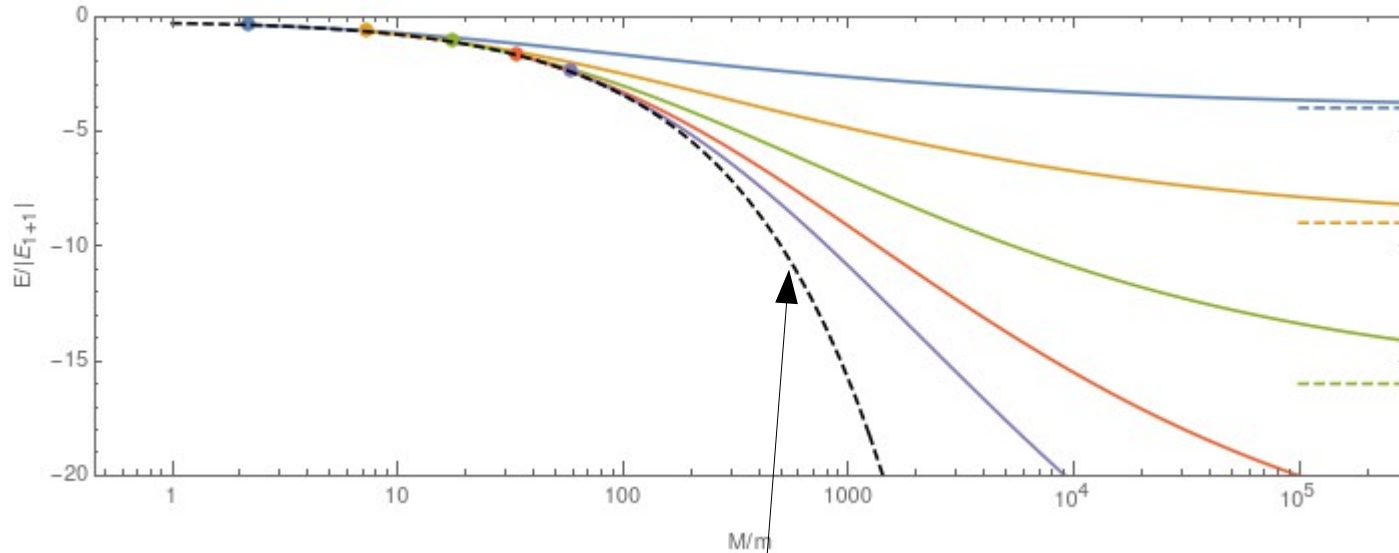
$$n(x) = \begin{cases} \sqrt{-2Mg[|\phi(x)|^2 - \mu/g]}/\pi, & |\phi(x)|^2 > \mu/g \\ 0, & |\phi(x)|^2 \leq \mu/g. \end{cases}$$

Minimization with respect to  $\phi(x)$   $\longrightarrow$

$$-\phi''(x) - \frac{\sqrt{-8Mm^2g^3}}{\pi} \sqrt{|\phi(x)|^2 - \mu/g} \phi(x) = 2m\epsilon\phi(x)$$



# Mean field #1 Thomas-Fermi



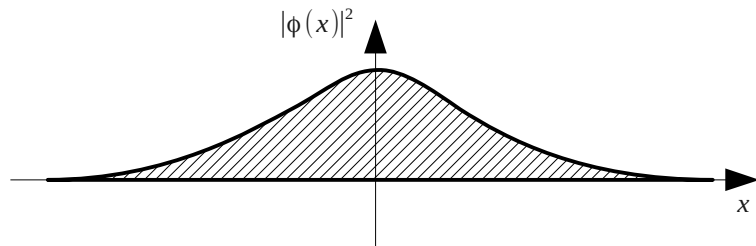
$E_{N+1}/|E_{1+1}| \approx -N^2$   
 Light atom bound in the potential  $N g \delta(x)$

N+1 cluster binds at

$$(M/m)/N^3 = \pi^2/36$$

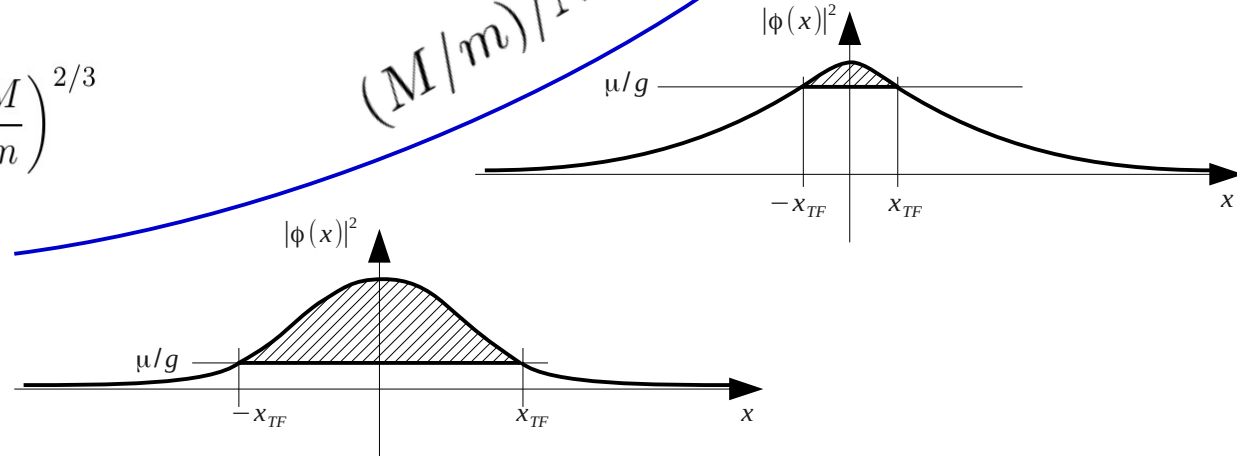
$$E = \epsilon + \int \frac{\pi^2 n^3(x)}{6M} dx = -\frac{mg^2}{30} N^2 = -\frac{mg^2}{30} \left( \frac{36 M}{\pi^2 m} \right)^{2/3}$$

$$\phi(x) = \frac{-3\pi\epsilon}{\sqrt{-8Mg^3} \cosh^2(\sqrt{-m\epsilon/2x})}$$



$\mu/g = 0$       $x_{TF} = \infty$

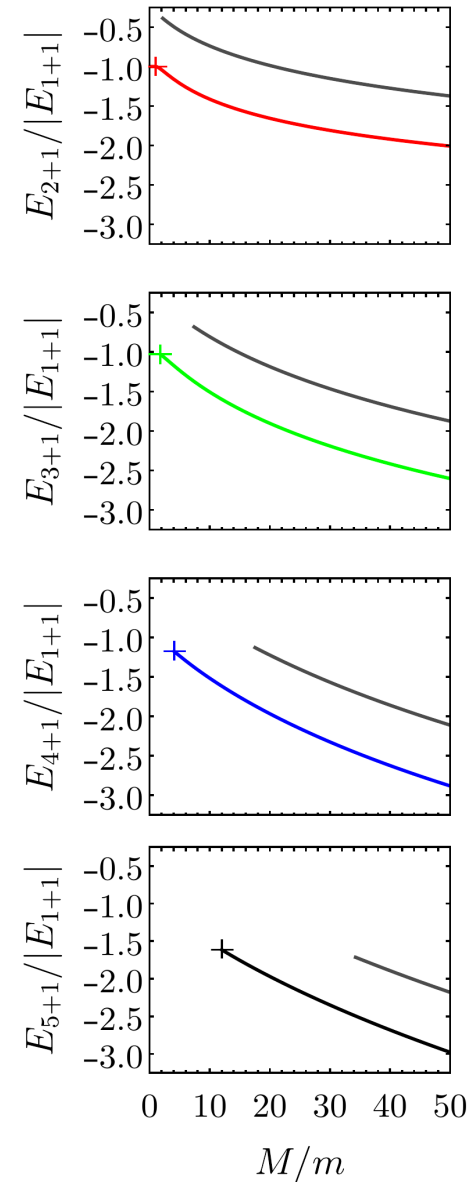
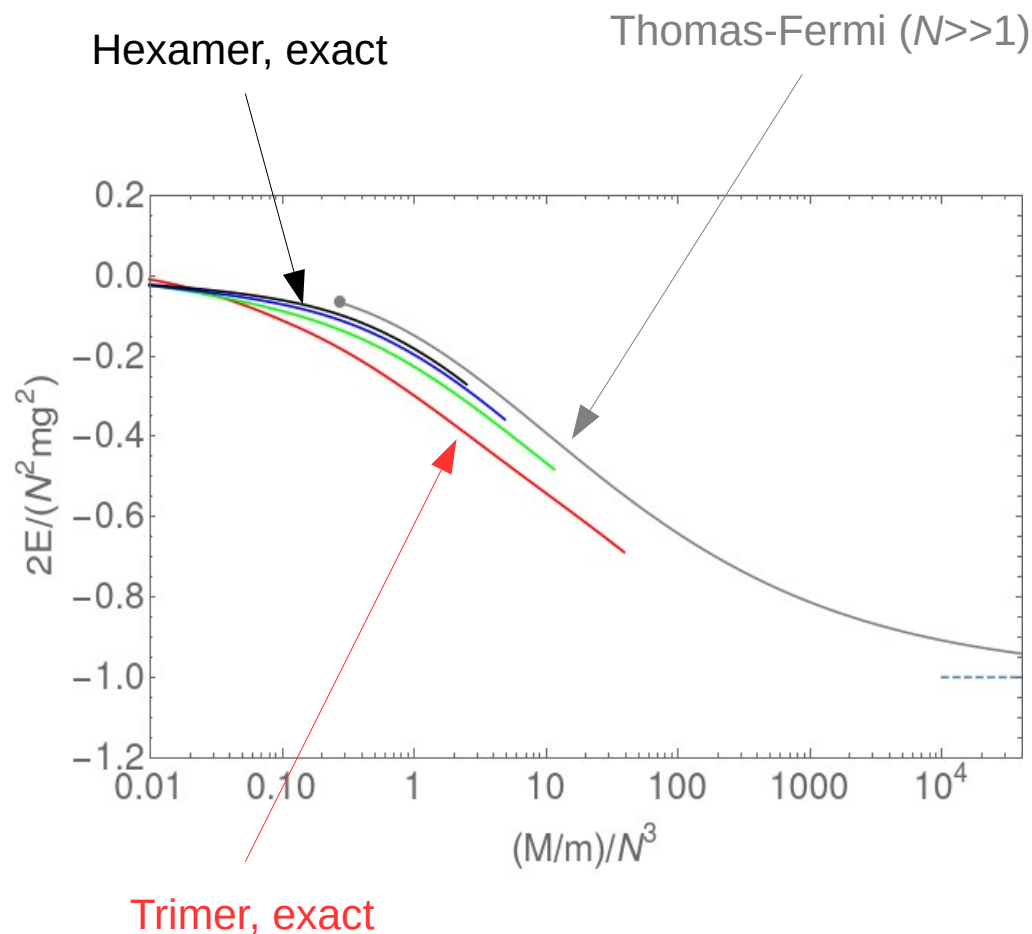
$(M/m)/N^3 \rightarrow \infty$



$(M/m)/N^3$  controls everything:  $E_{N+1}/mg^2 = N^2 f[(M/m)/N^3]$

Typical lengthscales  $\sim a/N$   $\longrightarrow$  validity condition:  $N \gg 1$

# Thomas-Fermi VS exact



Who is (mostly) responsible for the discrepancy: the MF treatment of the interaction or the local TF approximation for the kinetic energy?

# Mean field #2 Hartree-Fock

$$H = \int \left( -\frac{\hat{\Psi}_x^\dagger \partial_x^2 \hat{\Psi}_x}{2M} - \frac{\hat{\phi}_x^\dagger \partial_x^2 \hat{\phi}_x}{2m} + g \hat{\Psi}_x^\dagger \hat{\phi}_x^\dagger \hat{\Psi}_x \hat{\phi}_x \right) dx$$

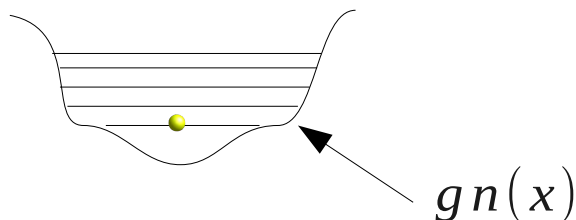
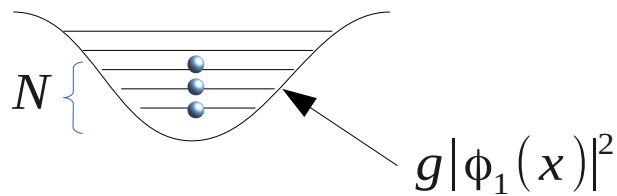
Variational ansatz

$$|v\rangle = \int dx \phi_1(x) \hat{\phi}_x^\dagger \int dx_1 \dots dx_N \frac{\det[\Psi_\nu(x_\eta)]}{\sqrt{N!}} \prod_{\eta=1}^N \hat{\Psi}_{x_\eta}^\dagger |0\rangle$$

Minimization with respect to the orbitals

$$\begin{aligned} -\partial_x^2 \phi_1(x)/2m + gn(x)\phi_1(x) &= \epsilon_1 \phi_1(x), \\ -\partial_x^2 \Psi_\nu(x)/2M + g|\phi_1(x)|^2 \Psi_\nu(x) &= E_\nu \Psi_\nu(x) \end{aligned}$$

$$n(x) = \sum_{\nu=1}^N |\Psi_\nu(x)|^2$$



$$E_{N+1} = \langle v | H | v \rangle = \epsilon_1 + \sum_{\nu=1}^N E_\nu - g \int |\phi_1(x)|^2 n(x) dx$$

Very similar energies as in the mean-field Thomas-Fermi case!

TF is good enough?



# “Improved” Hartree-Fock

$$-\partial_x^2 \phi_1(x)/2m + gn(x)\phi_1(x) = \epsilon_1 \phi_1(x),$$

$$-\partial_x^2 \Psi_\nu(x)/2M + g|\phi_1(x)|^2 \Psi_\nu(x) = E_\nu \Psi_\nu(x)$$



Orthonormal single-particle basis

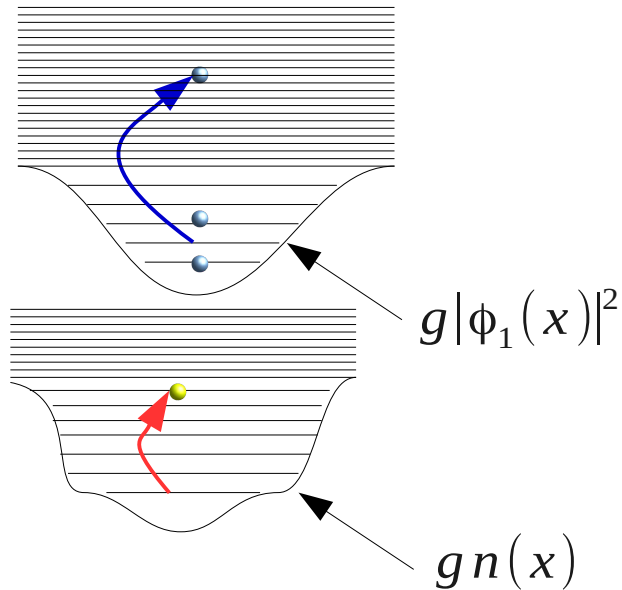


Hamiltonian couples  $|v\rangle$  with states  $|v'\rangle$  where the light atom and one of the heavy atoms are excited

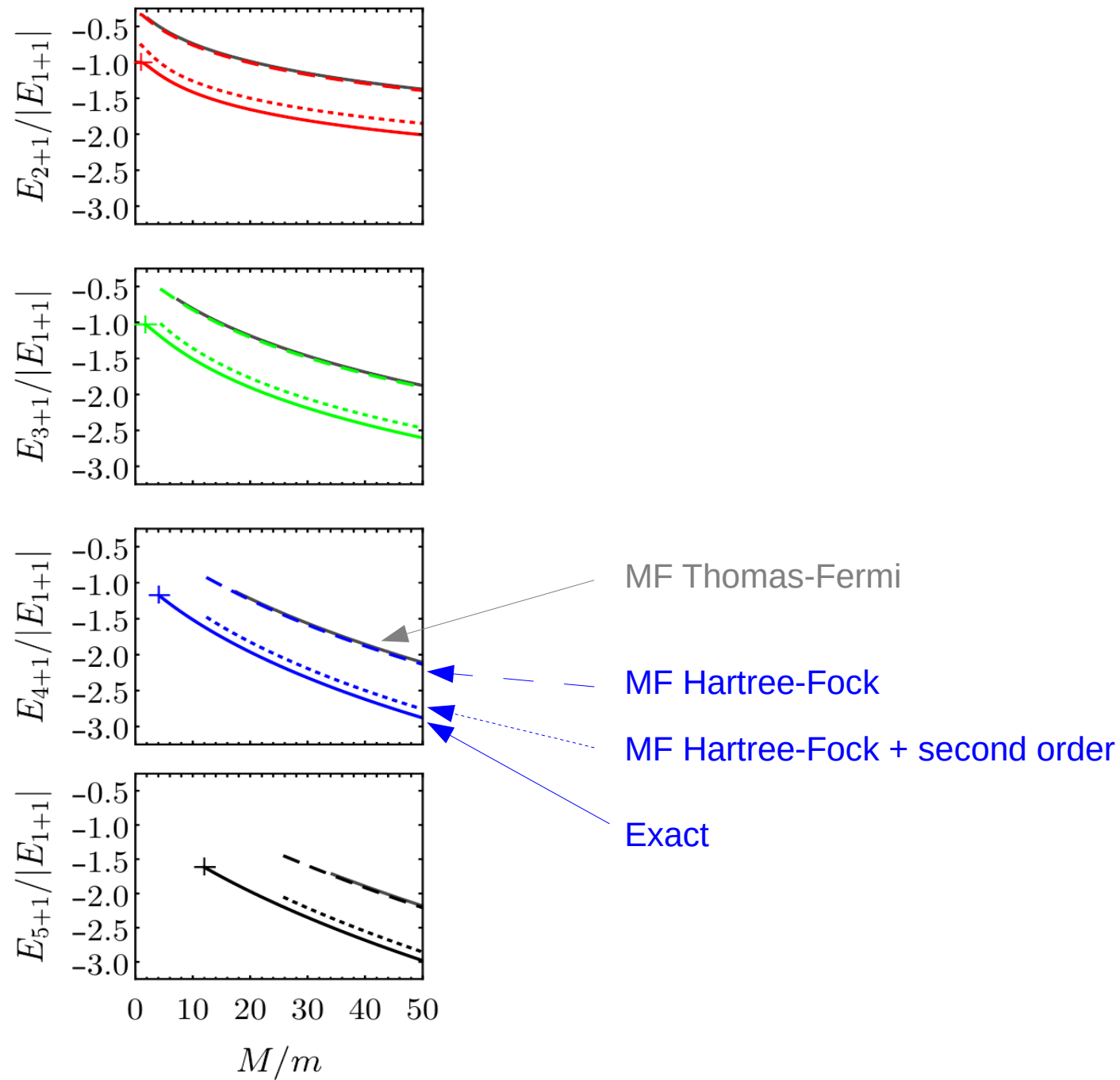


$$E_{N+1} = \langle v | H | v \rangle + \sum_{v' \neq v} \frac{|\langle v' | H | v \rangle|^2}{E_v^{(0)} - E_{v'}^{(0)}}$$

Computationally cheap, further corrections are more difficult

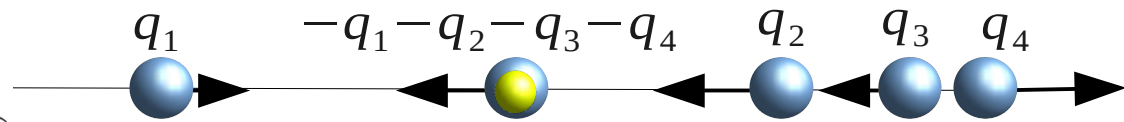
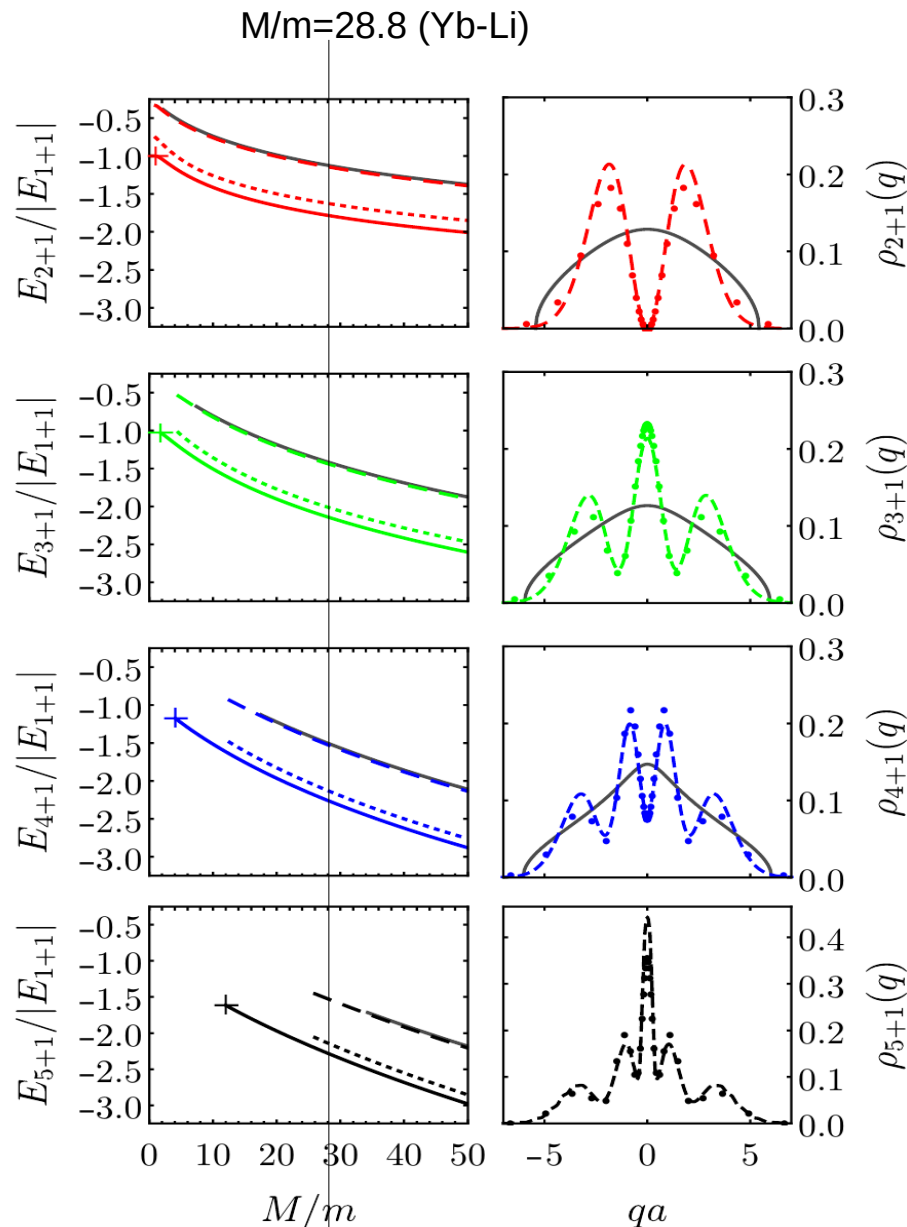


# Results



# Momentum correlations

$$\Psi(x_1, \dots, x_{N-1}, x_N, x = x_N) \xrightarrow{\text{Fourier}} F(q_1, \dots, q_{N-1}, q_N) \xrightarrow{q_N = -\sum_{i=1}^{N-1} q_i} F(q_1, \dots, q_{N-1})$$



Exact (STM)

Hartree-Fock  $|v\rangle$

$$\Psi(x_1, \dots, x_{N-1}, x_N, x)$$

$$F(q_1, \dots, q_{N-1})$$

Thomas-Fermi

$$\rho_{N+1}(q) = \int |F(q, q_2, \dots, q_{N-1})|^2 dq_2 \dots dq_{N-1}$$

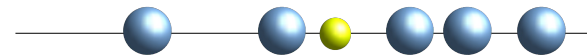
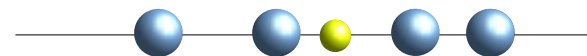
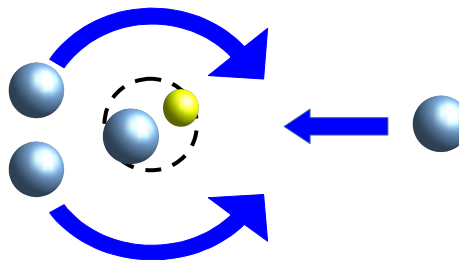
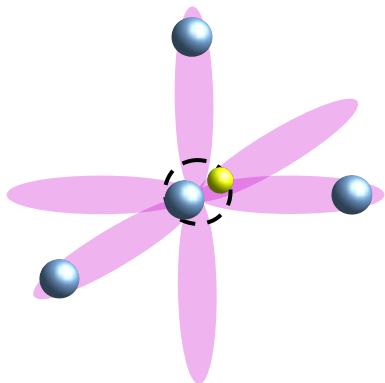
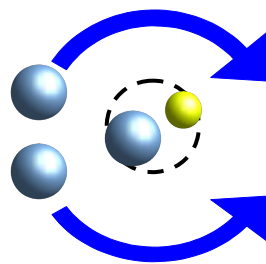
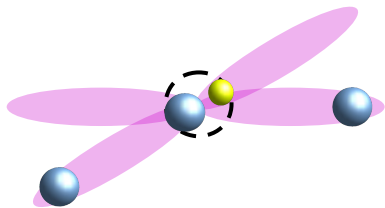
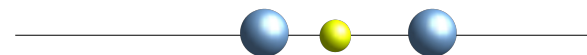
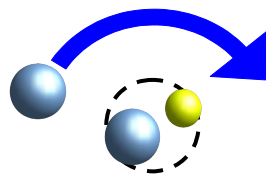
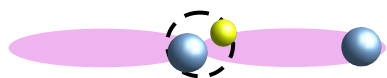
Momentum distribution of N-1 heavy fermions once the heavy-light pair is removed from the system (photoassociation)

# Summary and outlook

3D

2D

1D



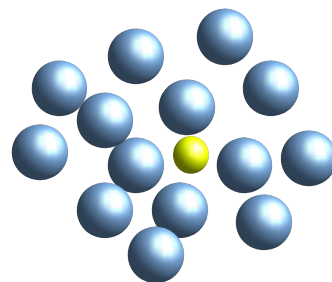
?

⋮

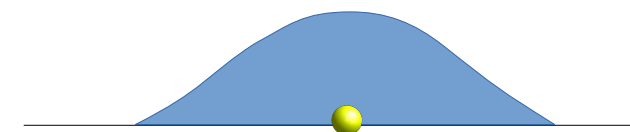
⋮

 $M/m \approx 13$  Efimov physics

?

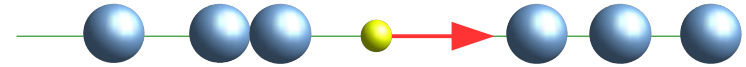


?


 Binding for  $\frac{M}{m} > \frac{\pi^2}{36} N^3$ 
 $M/m$  or  $N$

# Related (1D) research

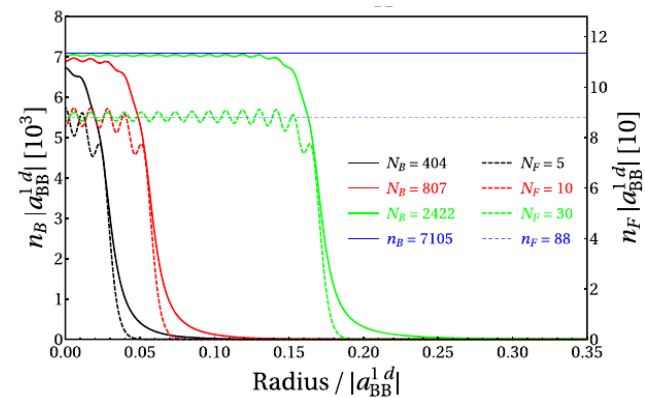
1D polaron or impurity (Fermi, Bose, solvable or not) [lots of literature]



Self-bound Bose-Fermi liquids in lower dimensions

Debraj Rakshit<sup>1,2</sup>, Tomasz Karpiuk<sup>3</sup>, Paweł Zin<sup>4</sup>, Mirosław Brewczyk<sup>3</sup>, Maciej Lewenstein<sup>2,5</sup> and Mariusz Gajda<sup>1</sup>

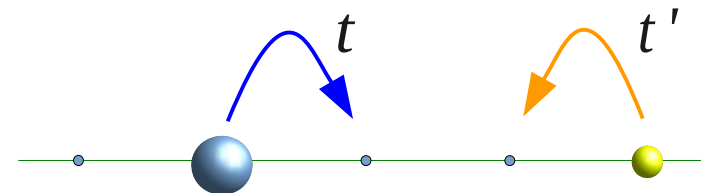
Mean field, Hartree-Fock...



1D lattice mixtures with asymmetric tunnelling (many-body phases)

DMRG

[Burovski-Orso-Jolicoeur-Roux'09-11,  
Dalmonte-Roschilde et al.'12]



Thank you!