

Gapped and gapless spin liquid phases on the Kagome lattice from chiral three-spin interactions

Bela Bauer (Station Q)

Andreas Ludwig & Brendan Keller (UCSB), Simon Trebst (U Cologne),
Michele Dolfi (ETH Zurich)



Today's plan

Intuitive picture for emergent phases in a network of edge states

Numerical evidence for a gapped and a gapless spin liquid in a local $SU(2)$ -invariant spin-1/2 system on the Kagome lattice

A very short overview

Gapped topological SL

- Two flavors: ***broken*** or ***unbroken*** time reversal (chiral - non-chiral)
- Long history for chiral spin liquids
 - Kalmeyer & Laughlin '89: triangular lattice Heisenberg AFM = Laughlin state?
 - Exact, but complicated parent Hamiltonians (Schroter, Thomale, Greiter/Yao)
 - Topological flat band models
- More recently interest in non-chiral topological spin liquids (Z_2 spin liquid)

A very short overview

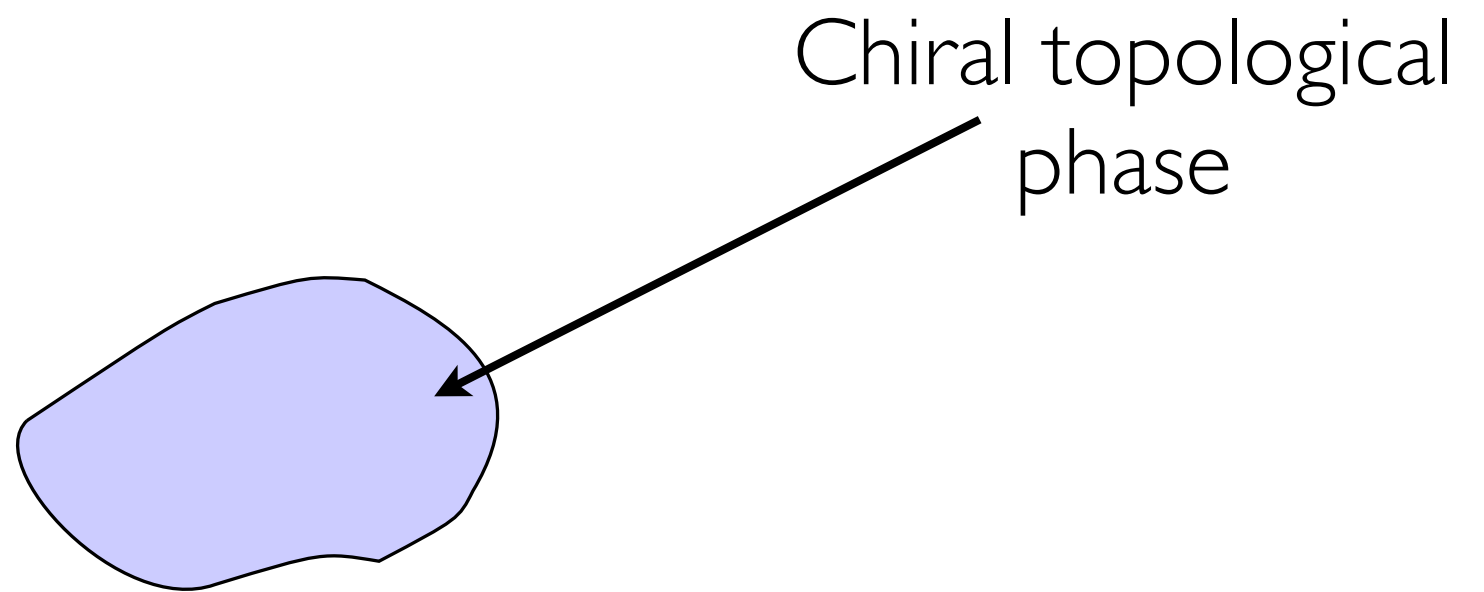
Gapped topological SL

- Two flavors: **broken** or **unbroken** time reversal (chiral - non-chiral)
- Long history for chiral spin liquids
 - Kalmeyer & Laughlin '89: triangular lattice Heisenberg AFM = Laughlin state?
 - Exact, but complicated parent Hamiltonians (Schroter, Thomale, Greiter/Yao)
 - Topological flat band models
- More recently interest in non-chiral topological spin liquids (Z_2 spin liquid)

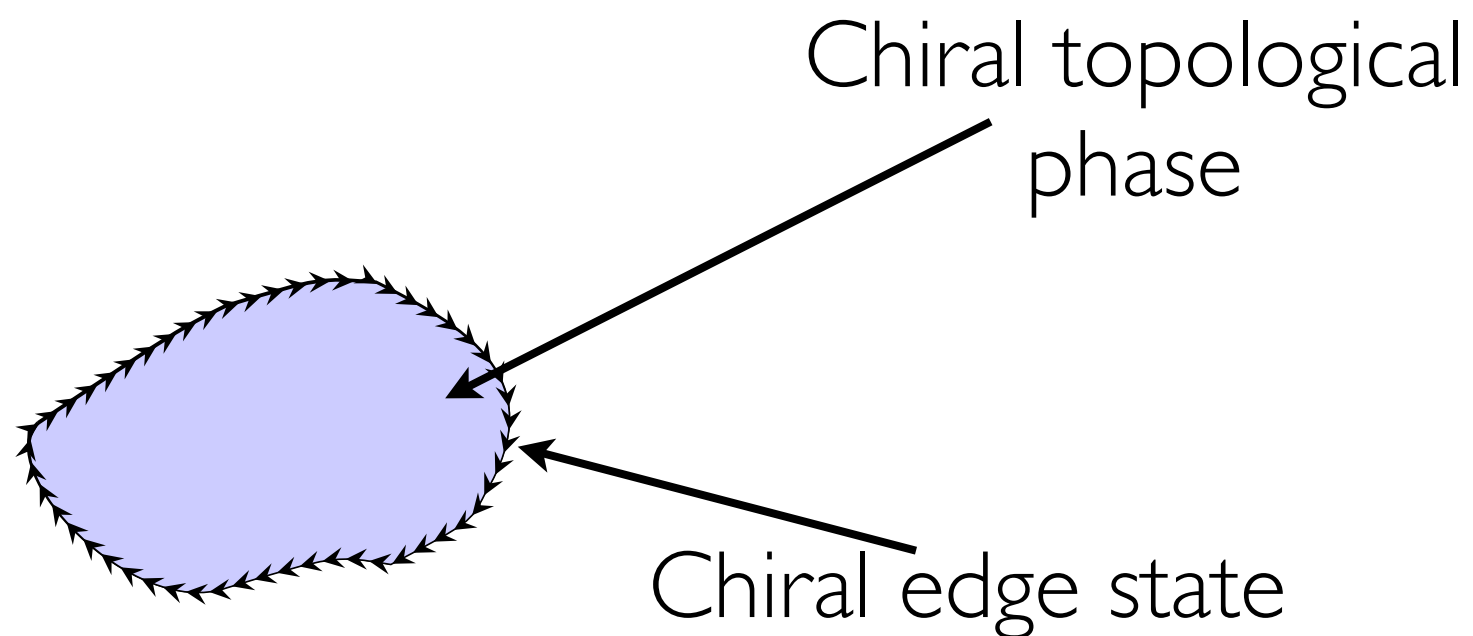
Gapless SL

- Two flavors: gapless excitations at points in momentum space, or “Fermi surfaces”
 - Points in momentum space: *algebraic* spin liquid - Kitaev's honeycomb model
 - “Fermi surfaces”: d-wave correlated Bose liquid, spinon Fermi surface

Network of edge states

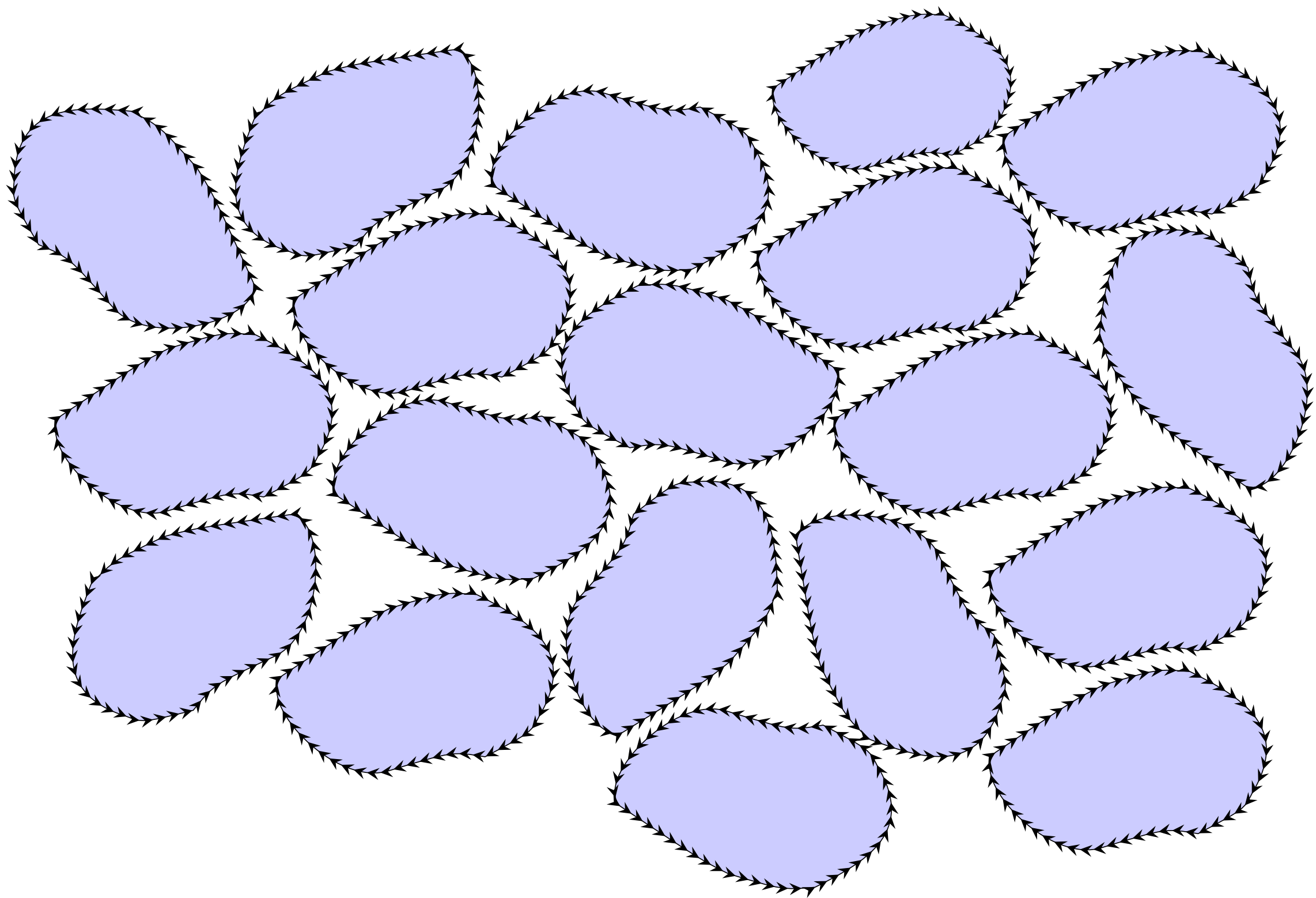


Network of edge states

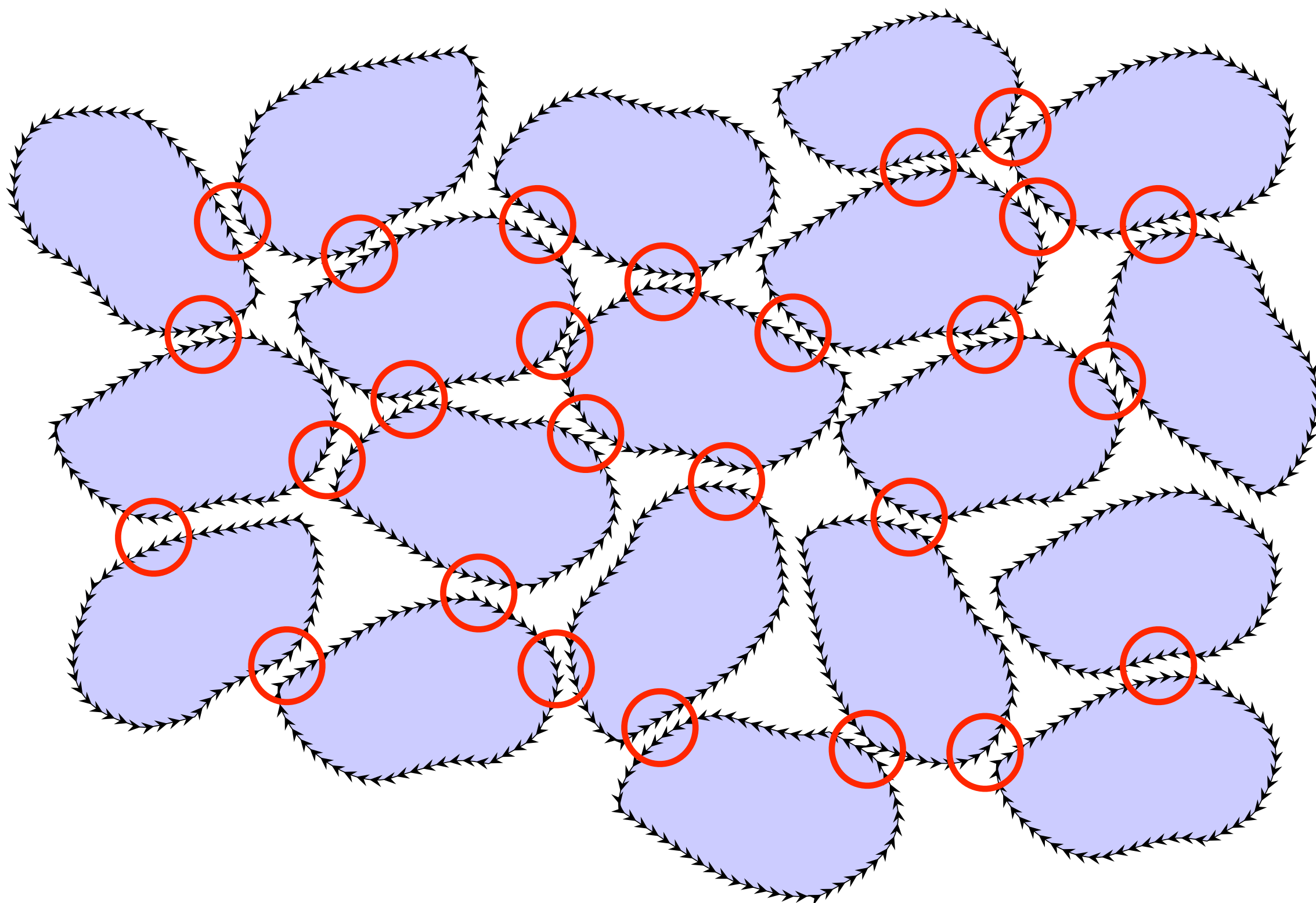


Network of edge states

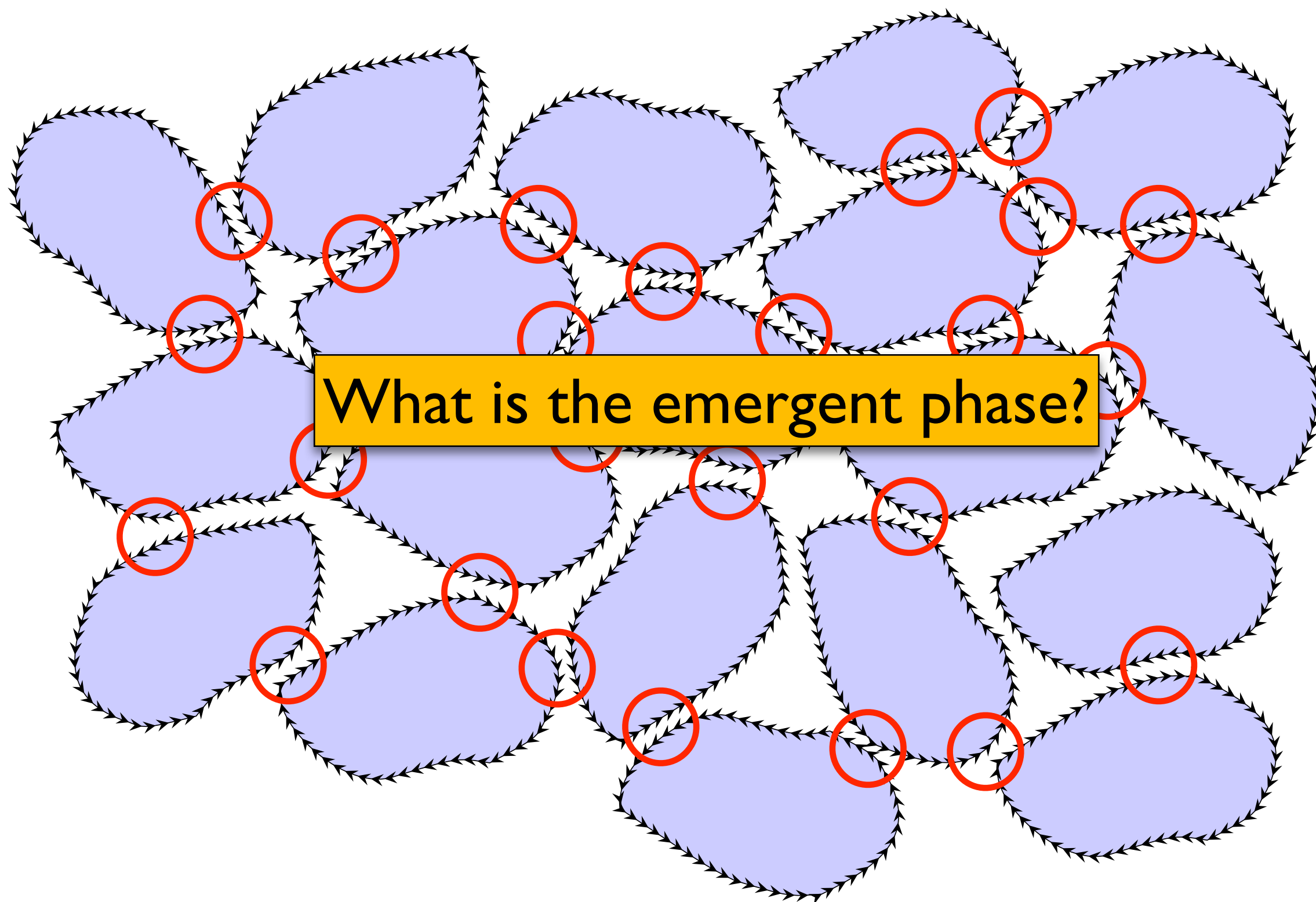
Network of edge states



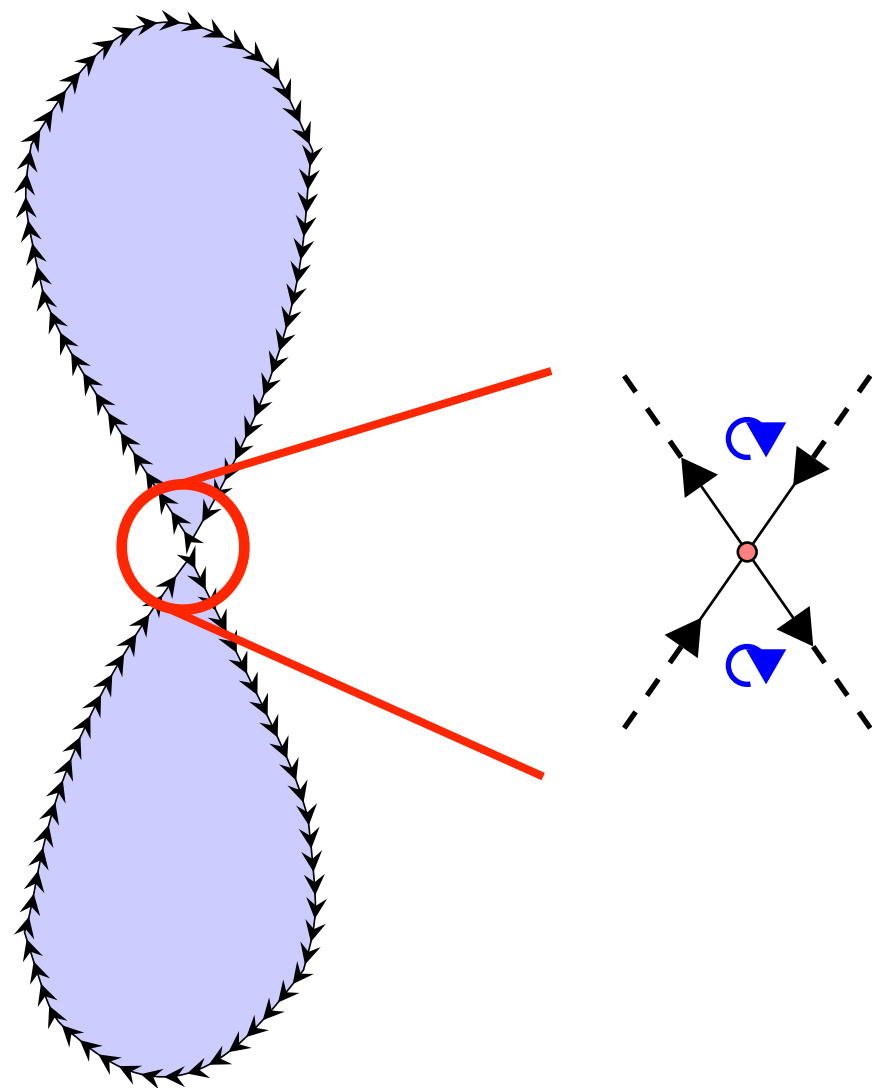
Network of edge states



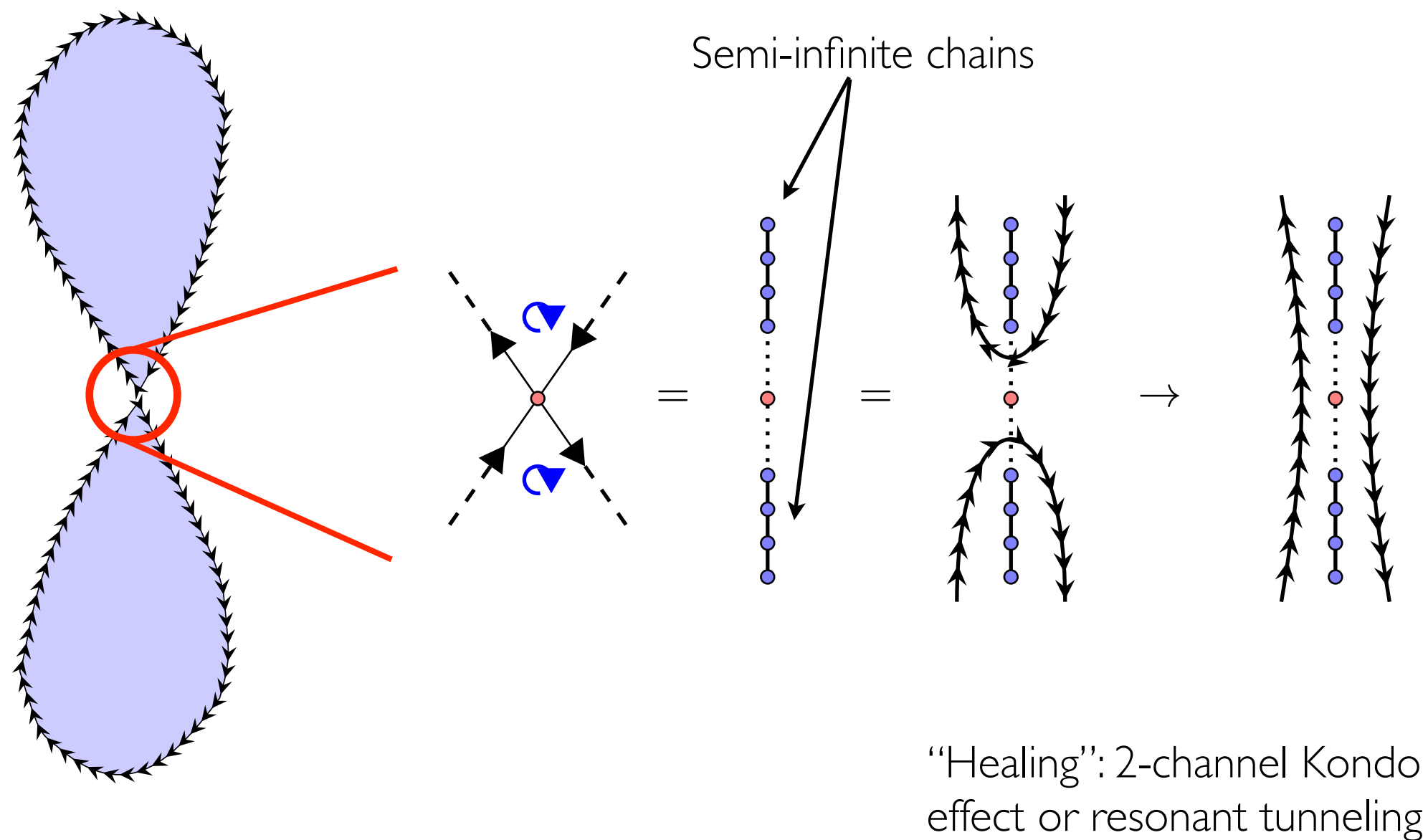
Network of edge states



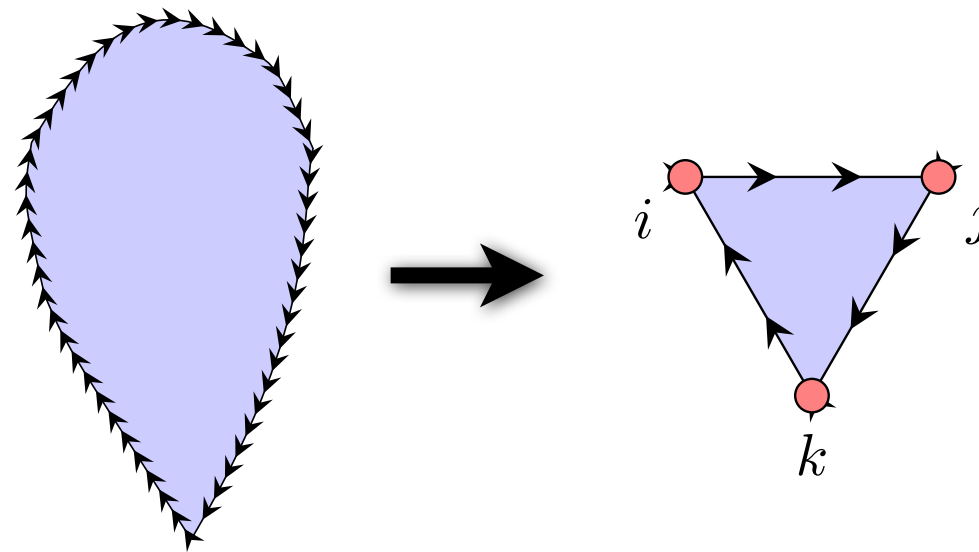
Building block: two puddles



Building block: two puddles



Puddles to triangles

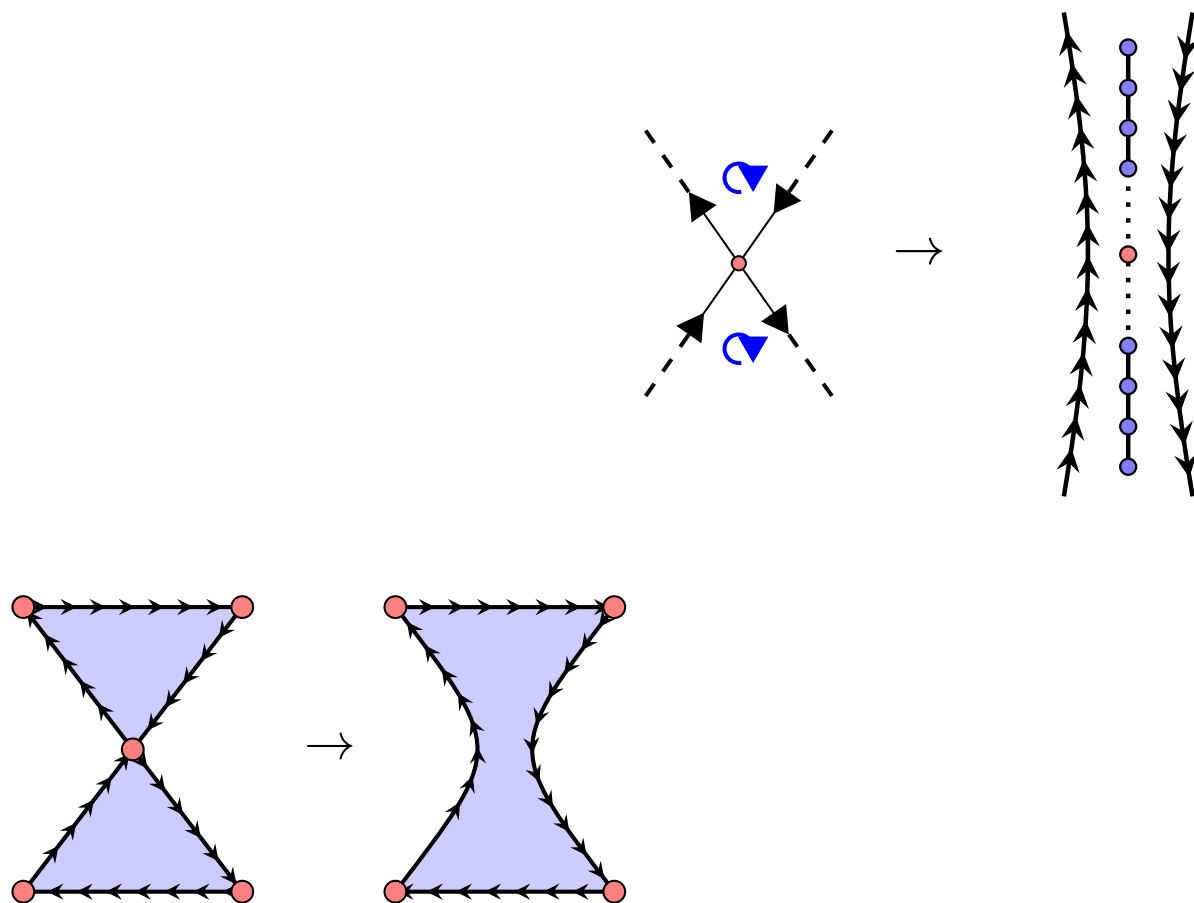


Break time reversal
on each triangle!

Spin 1/2: $\chi_{ijk} = \frac{i}{2} \vec{S}_i \cdot (\vec{S}_j \times \vec{S}_k)$

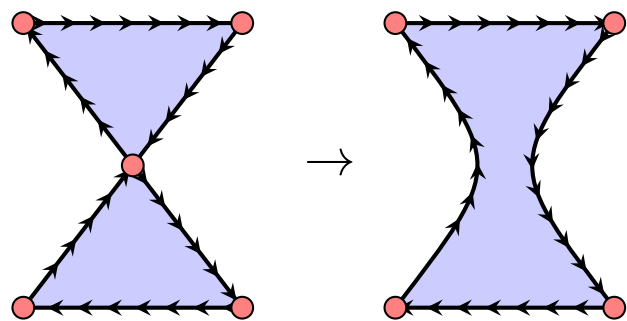
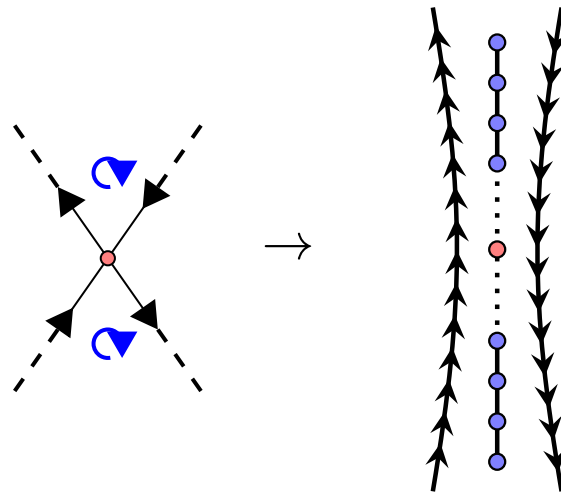
Majorana fermions: $\tilde{\chi}_{ijk} = i(\gamma_i \gamma_j + \gamma_j \gamma_k + \gamma_k \gamma_i)$

Two triangles

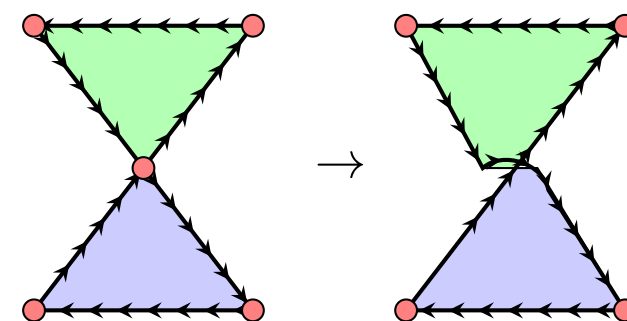


Equal chirality

Two triangles

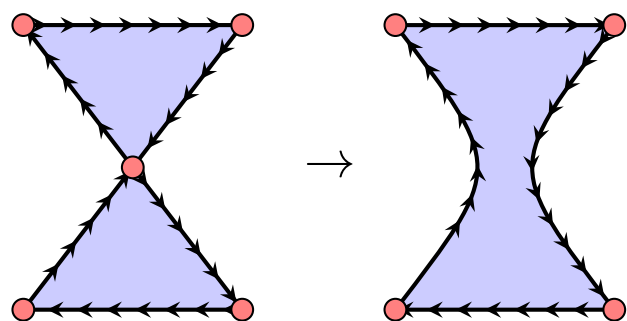
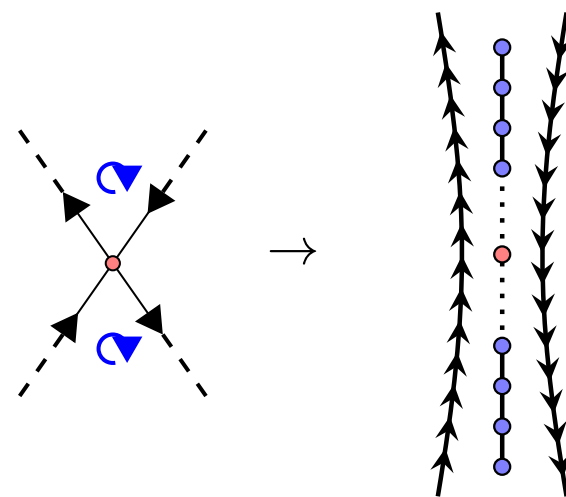


Equal chirality

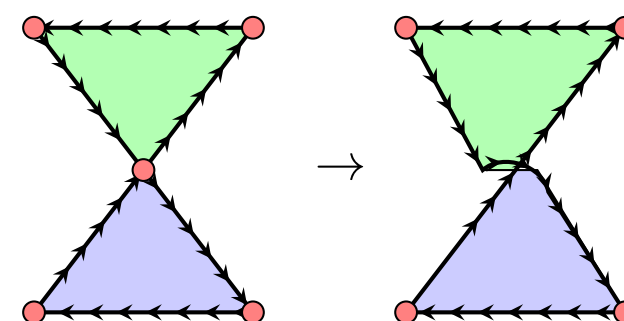


Different chirality

Two triangles



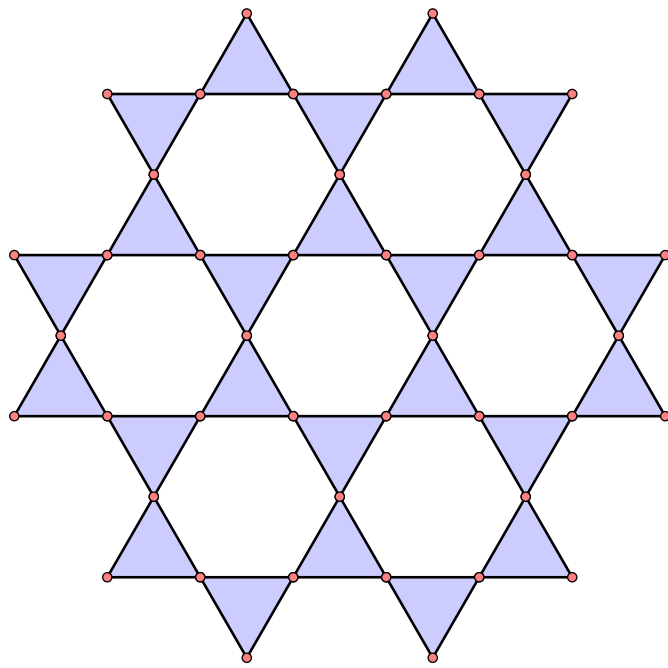
Equal chirality



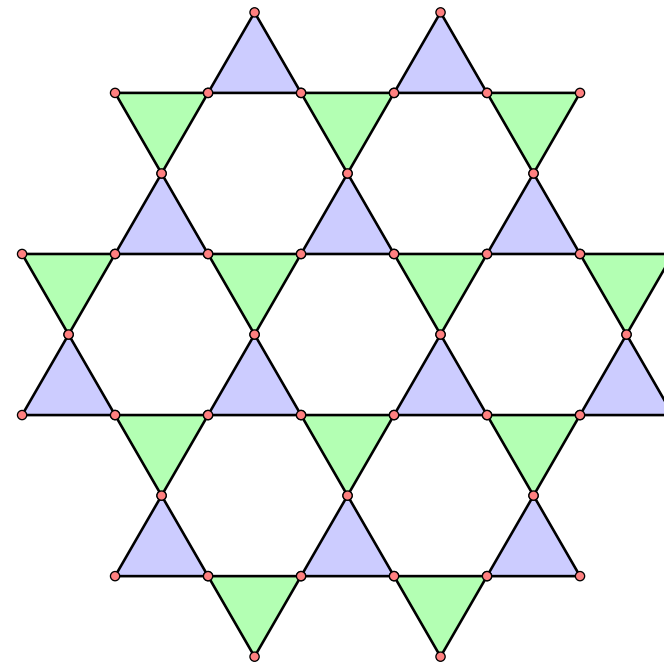
Different chirality

Both cases: **one** edge state remains

Kagome lattice: two cases

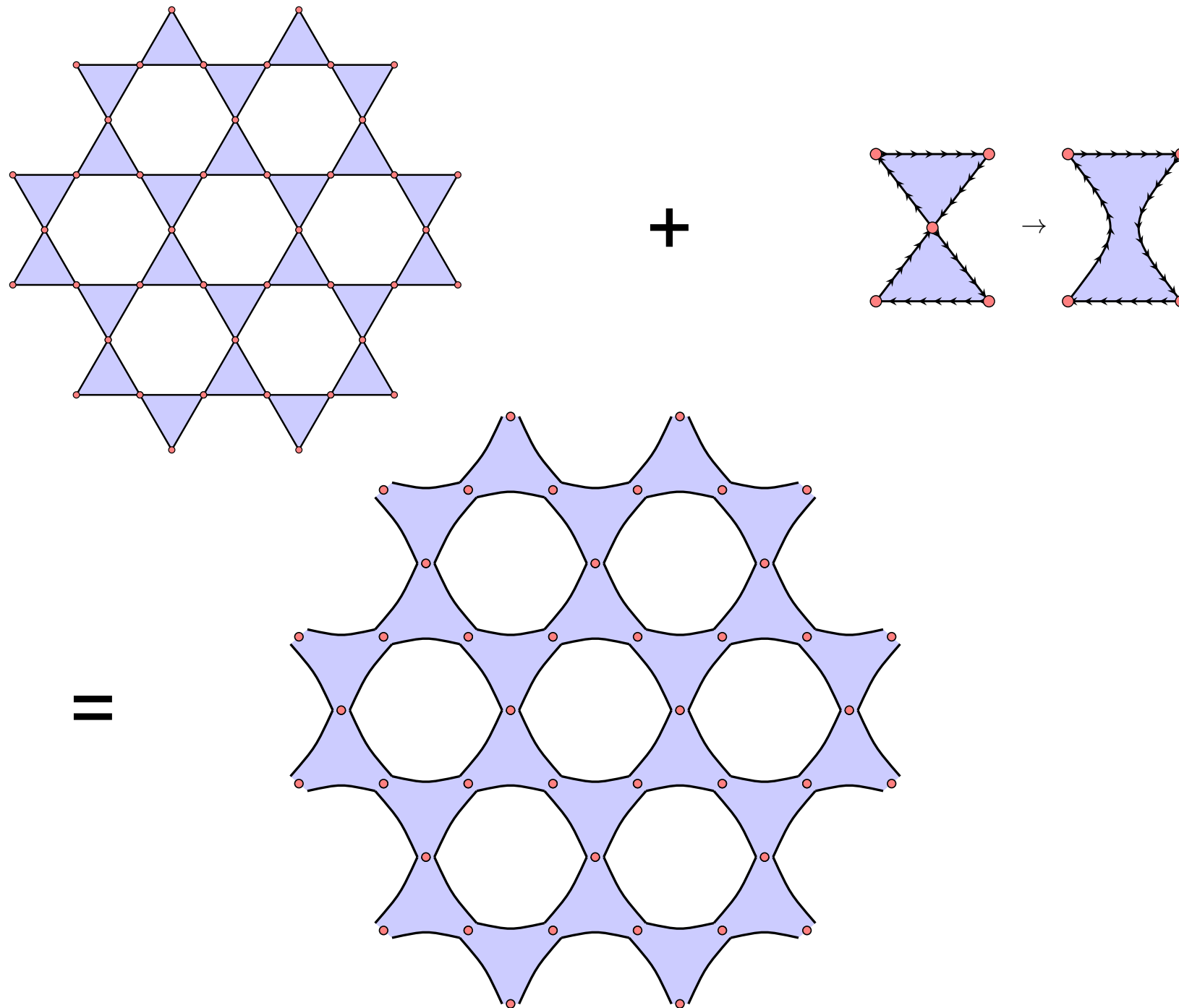


Homogeneous

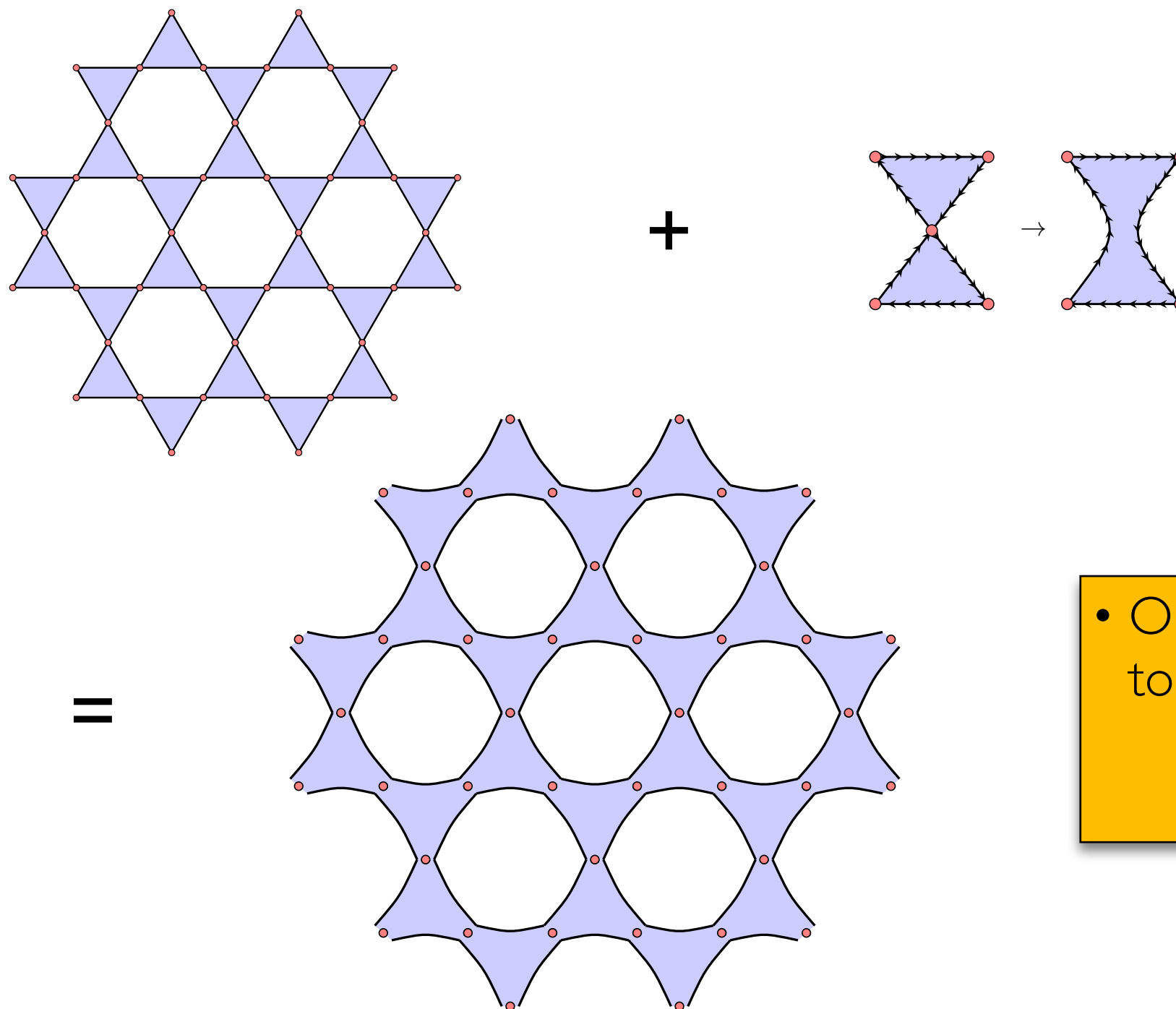


Staggered

Homogeneous phase

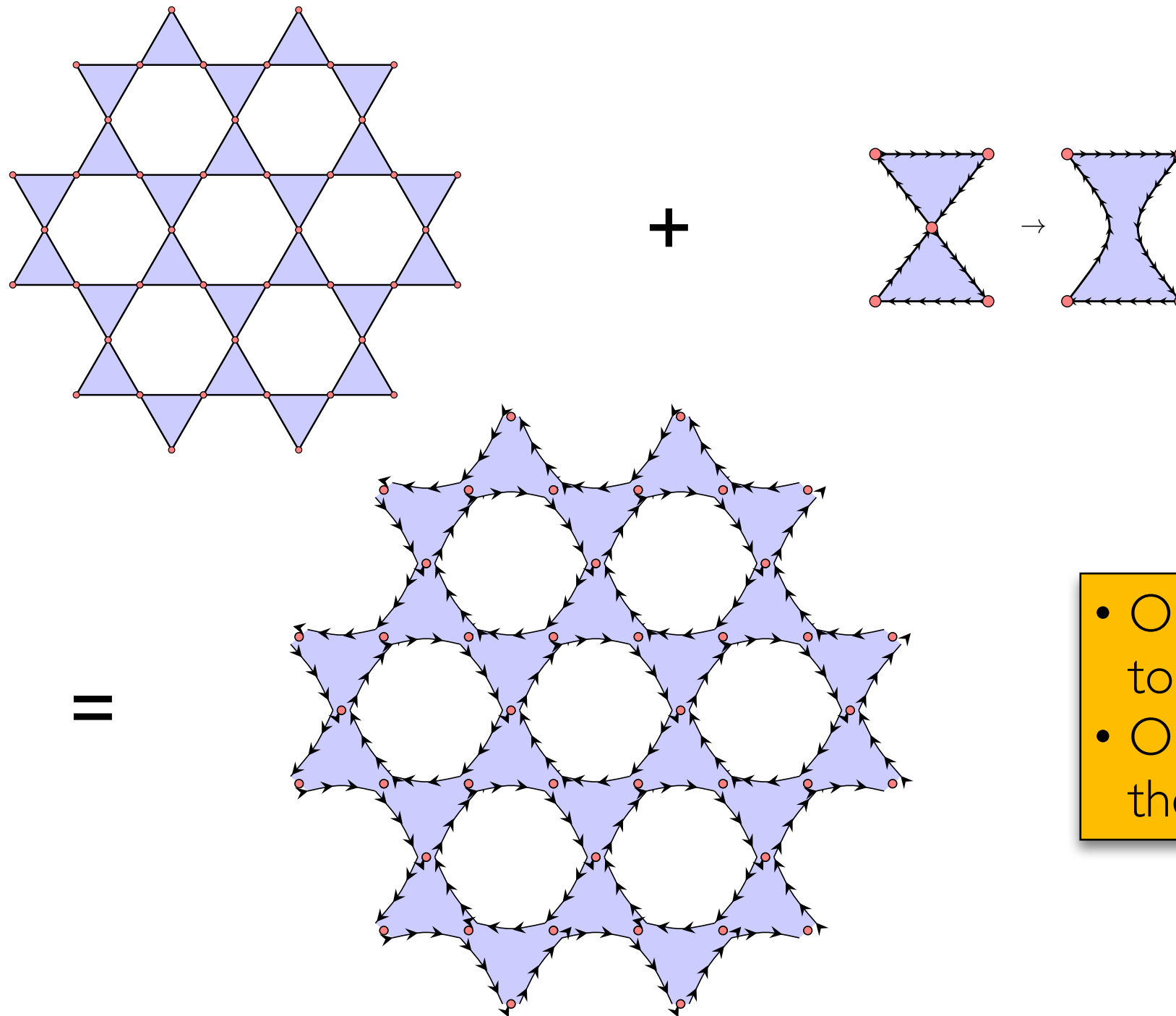


Homogeneous phase



- One extended region of the topological phase

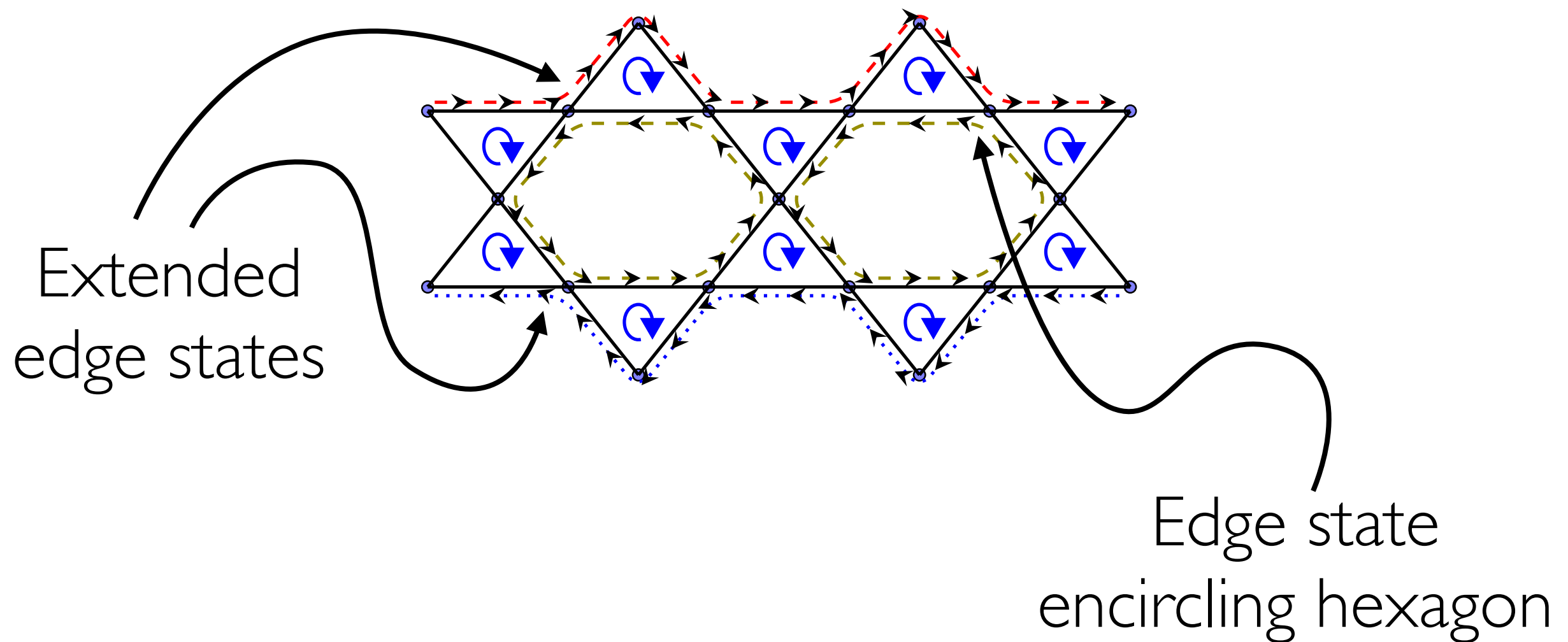
Homogeneous phase



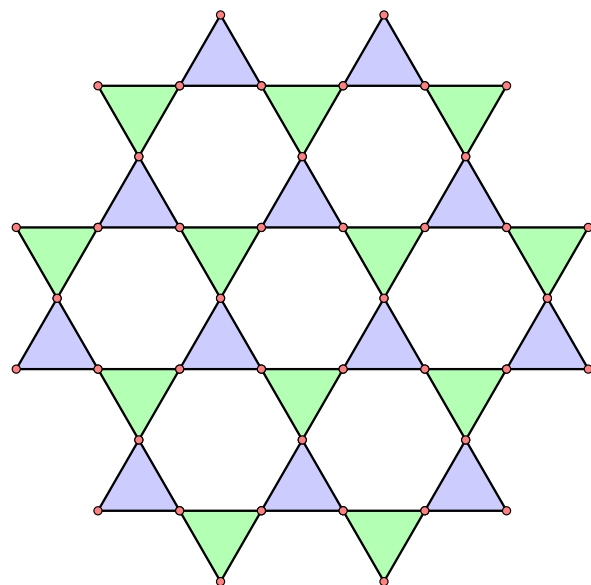
- One extended region of the topological phase
- One edge state encircling the whole system

Homogeneous phase

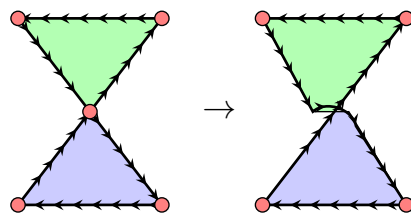
$2 \times L$ cluster



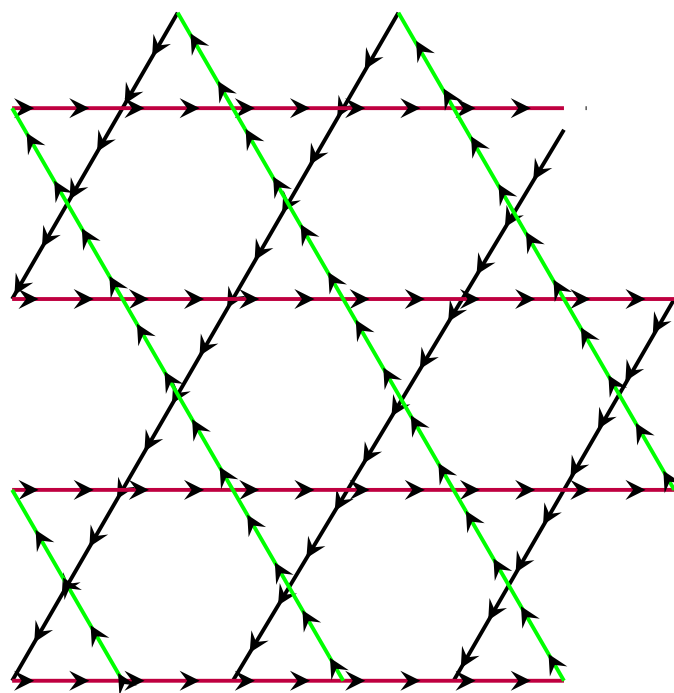
Staggered phase



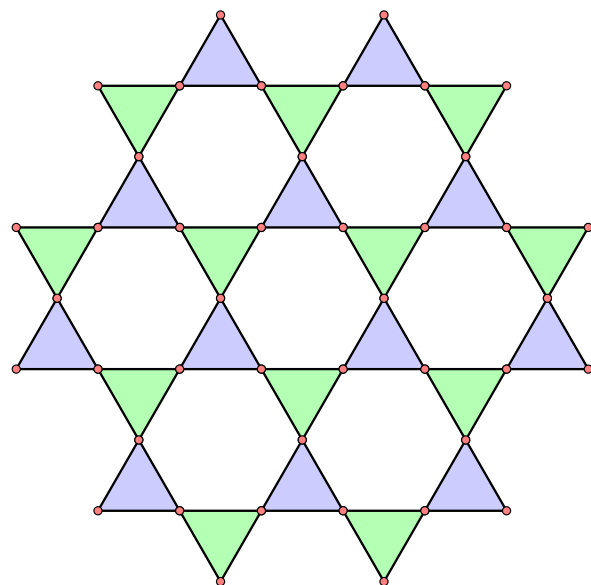
+



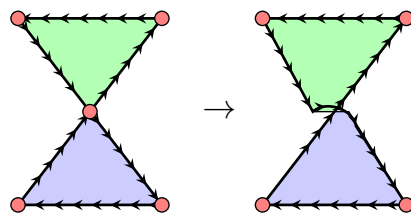
=



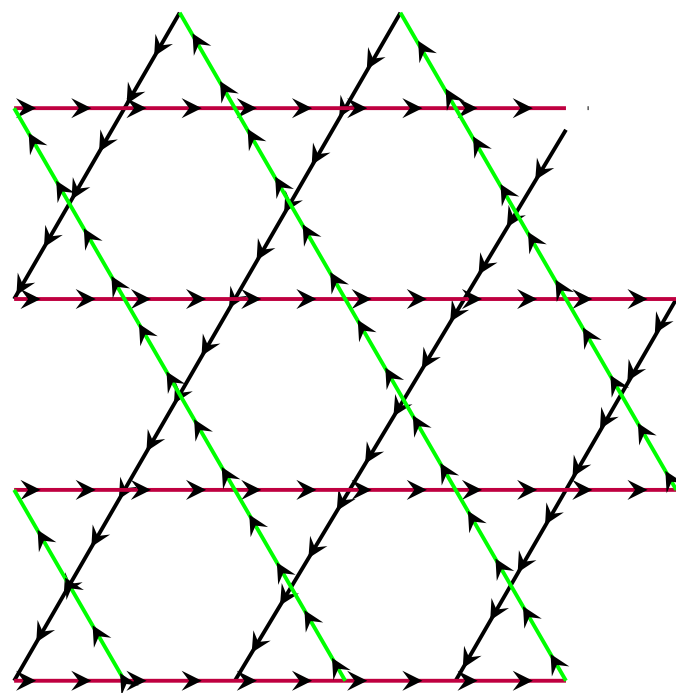
Staggered phase



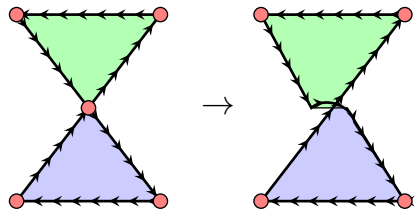
+



=

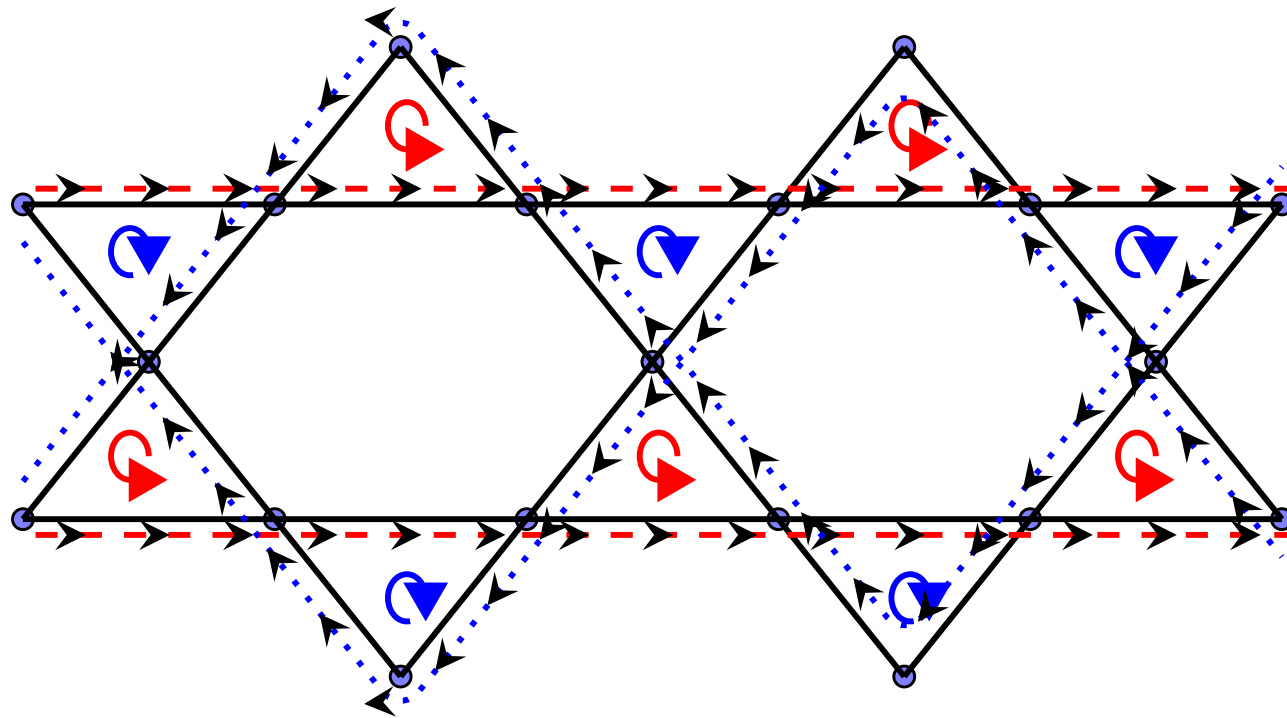


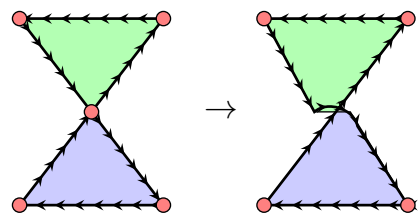
Uncoupled gapless edge states on the chains give rise to bulk gapless phase!



Staggered phase

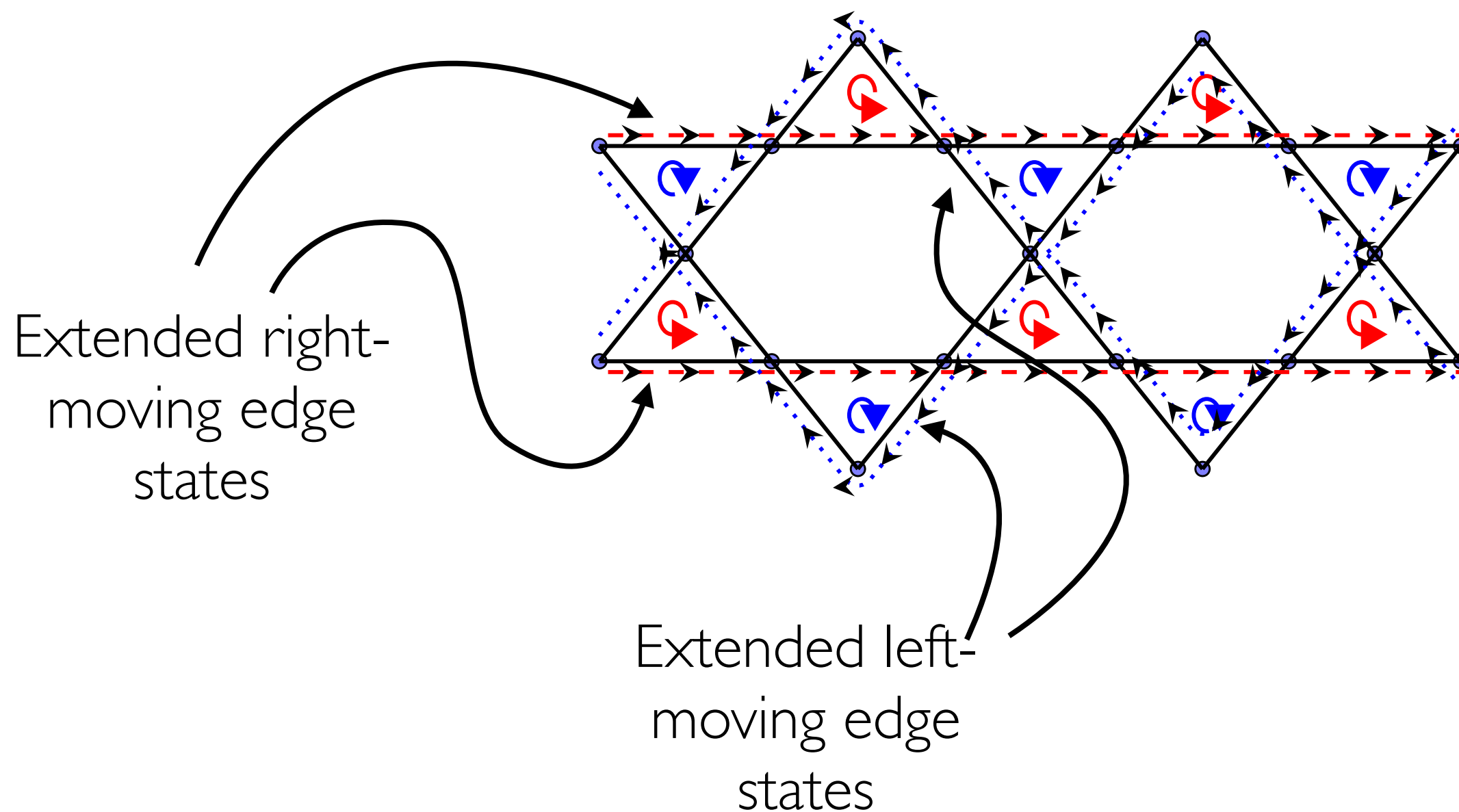
$2 \times L$ cluster





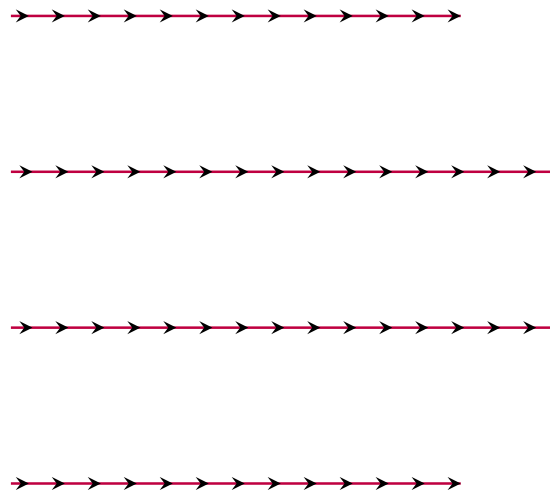
Staggered phase

$2 \times L$ cluster

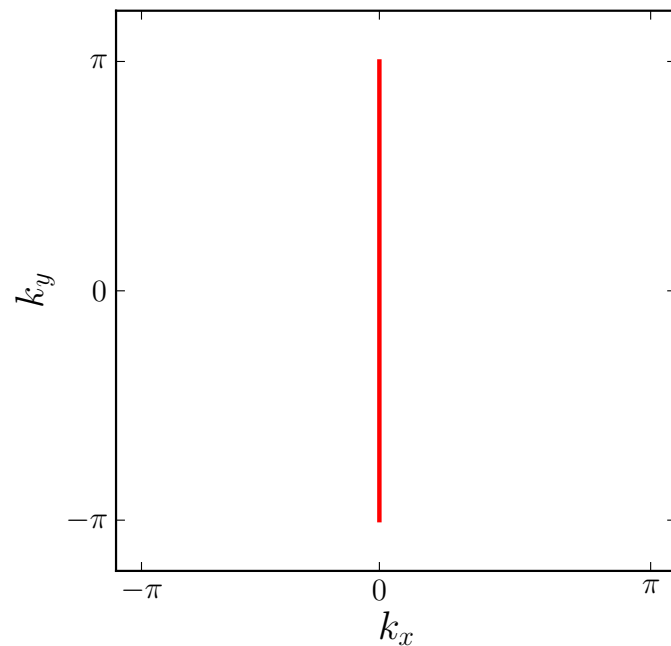


Fermi surface

Real space

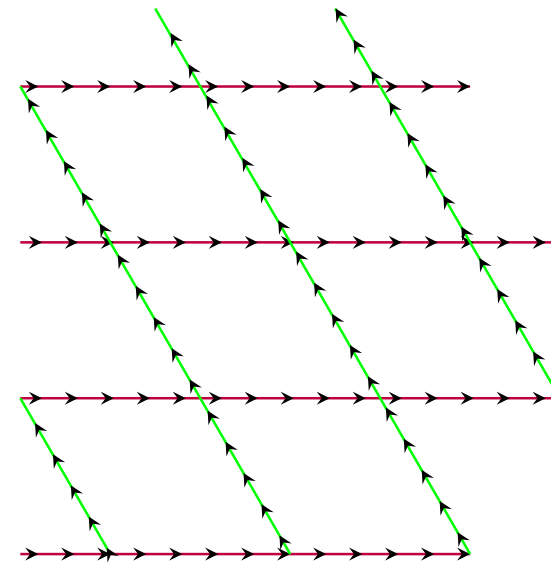
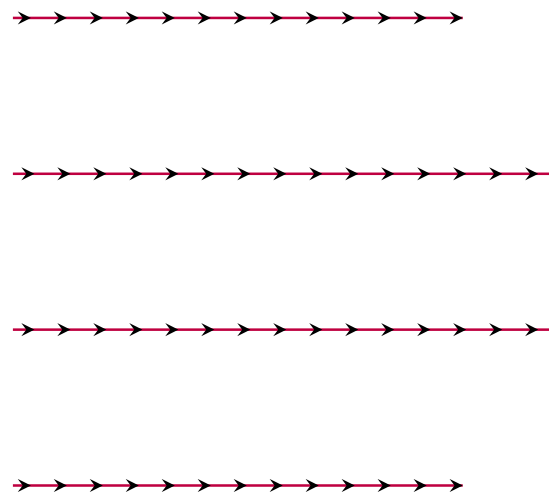


Momentum space

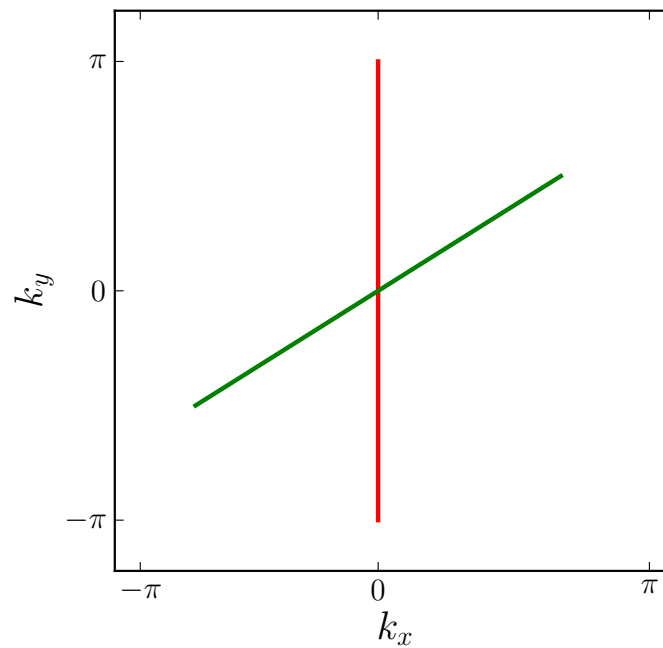
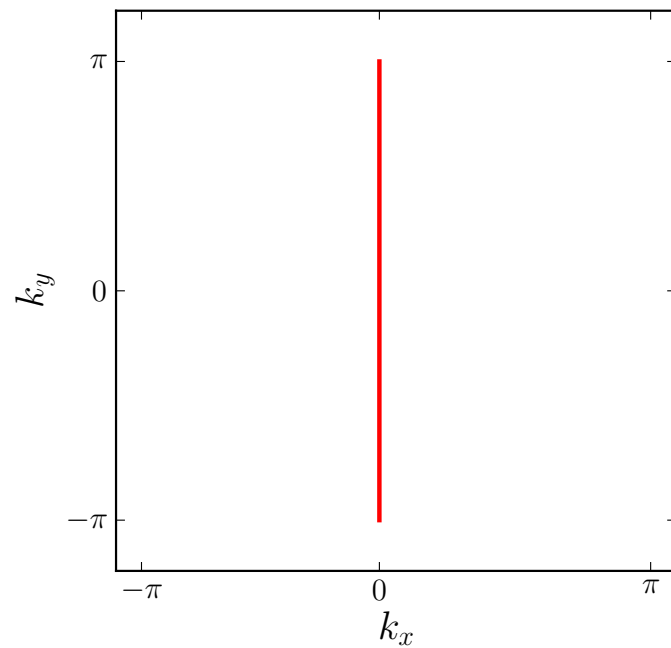


Fermi surface

Real space

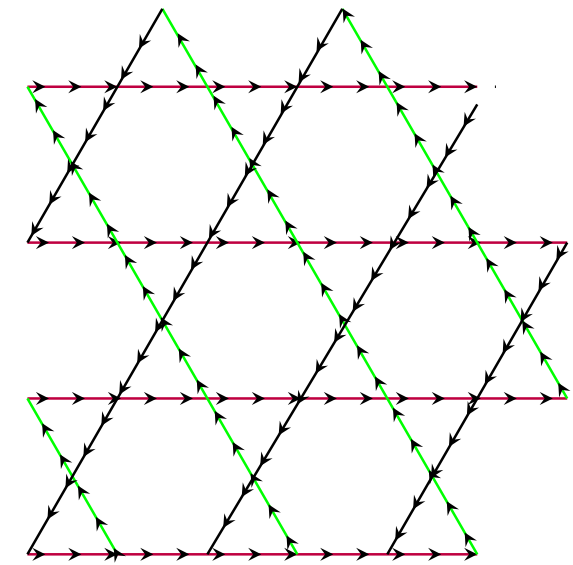
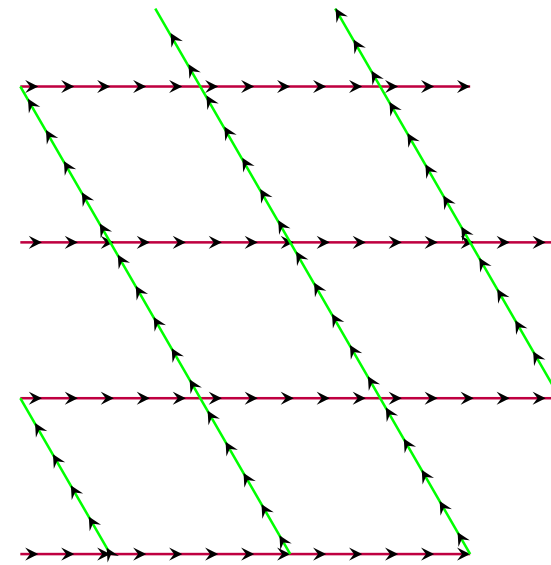
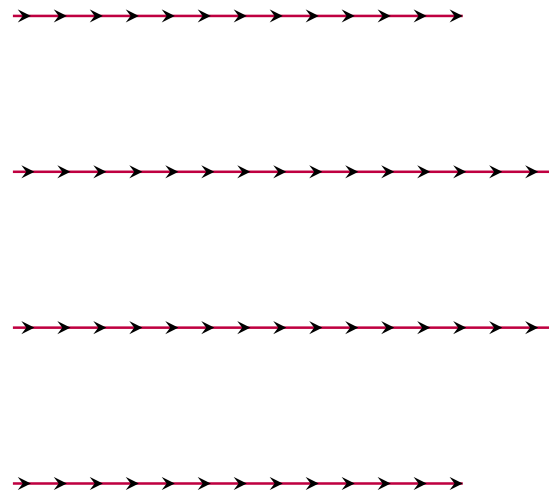


Momentum space

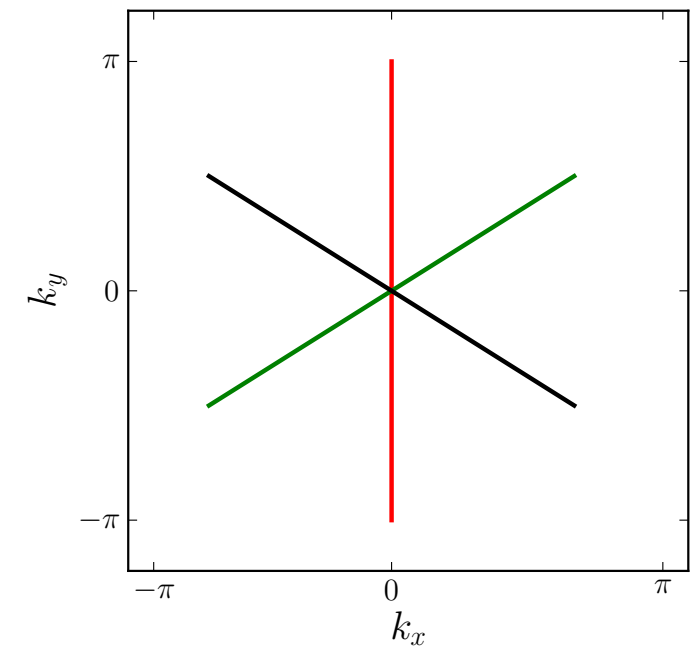
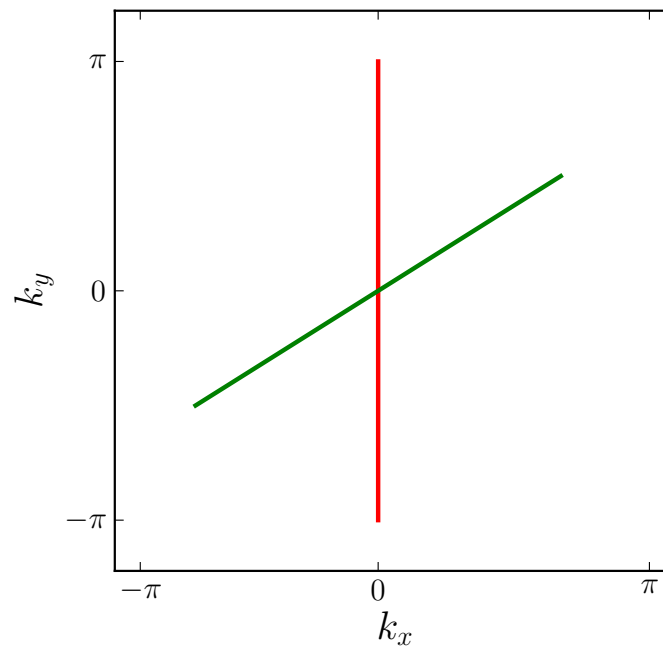
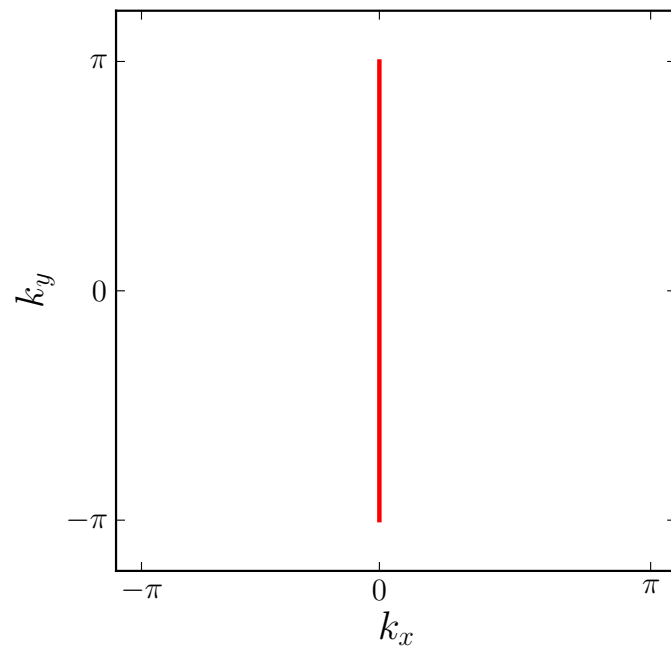


Fermi surface

Real space

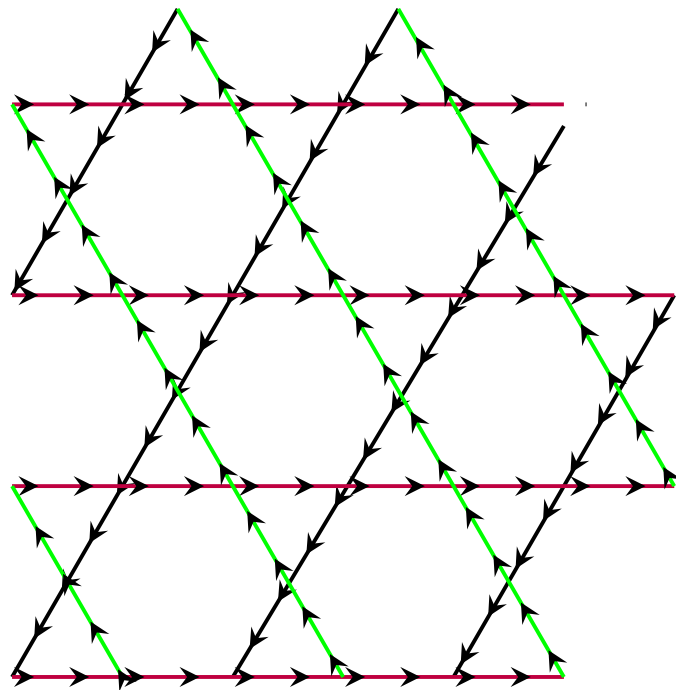


Momentum space

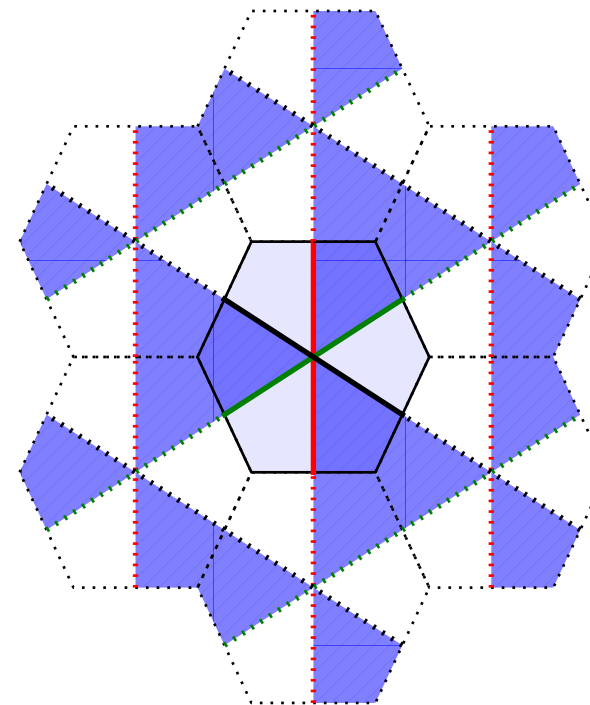


Fermi surface

Real space



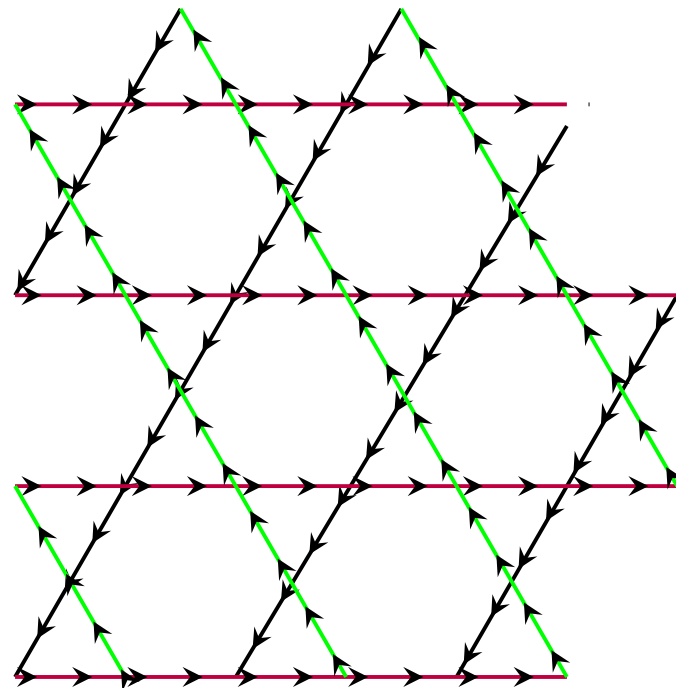
Momentum space



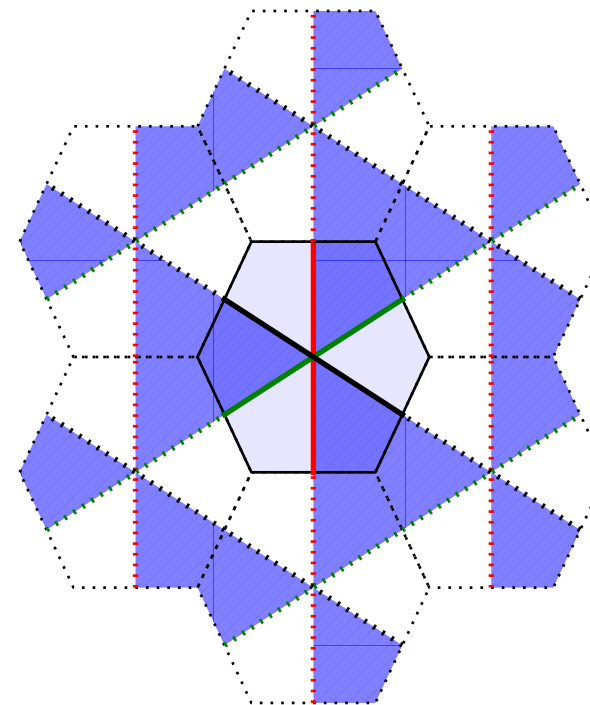
- Light blue: first BZ
- Dark blue: filled regions

Fermi surface

Real space



Momentum space

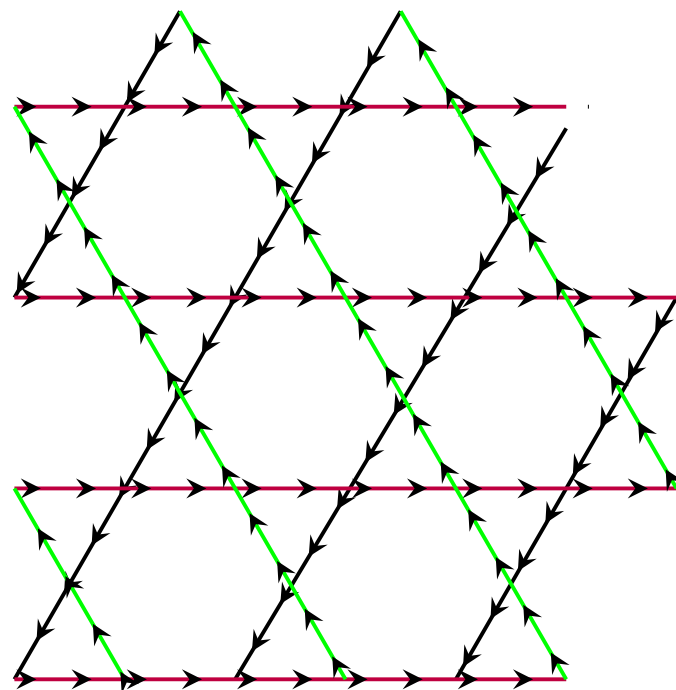


- Light blue: first BZ
- Dark blue: filled regions

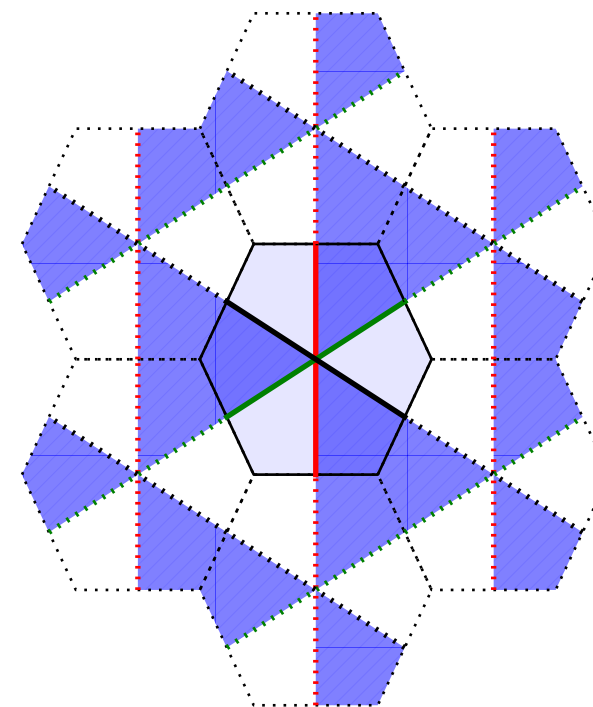
- Each direction of edge states leads to a gapless line in the BZ
- Lines enclose triangles of filled states

Symmetries

Real space



Momentum space



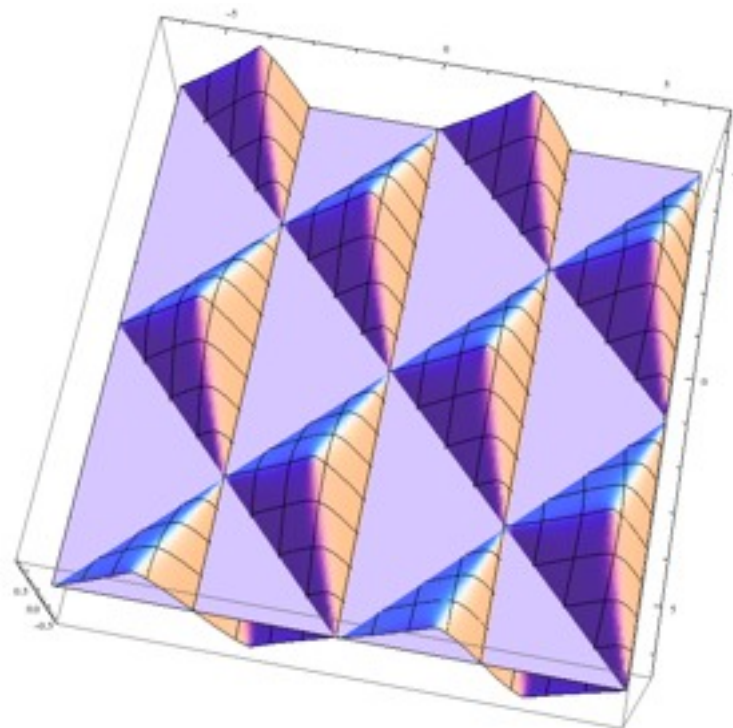
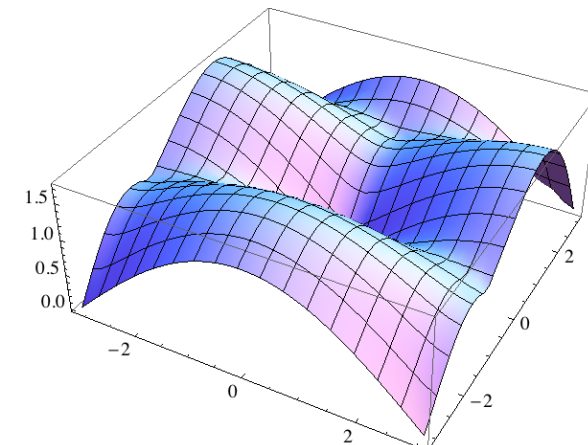
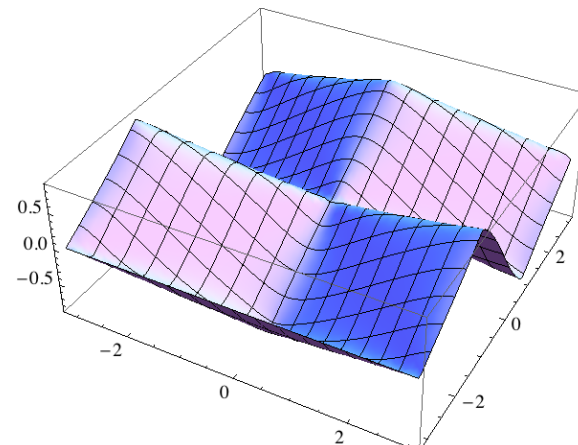
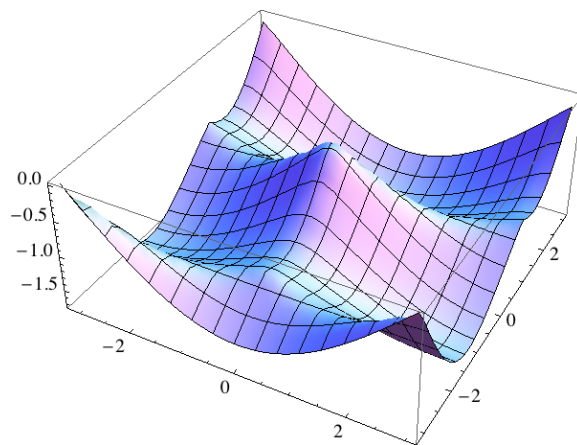
—

- Rotational symmetry: 6-fold to 3-fold
- Inversion symmetry $x \leftrightarrow -x$: broken
- Inversion symmetry $y \leftrightarrow -y$: not broken
- (Inversion $x \leftrightarrow -x$) \times (time reversal): additional symmetry

Majorana fermions: staggered

$$\tilde{\chi}_{ijk} = i(\gamma_i \gamma_j + \gamma_j \gamma_k + \gamma_k \gamma_i)$$

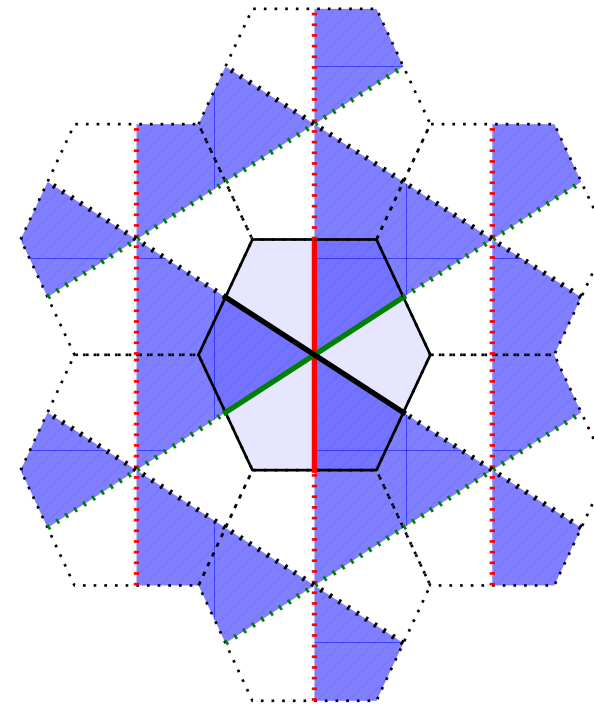
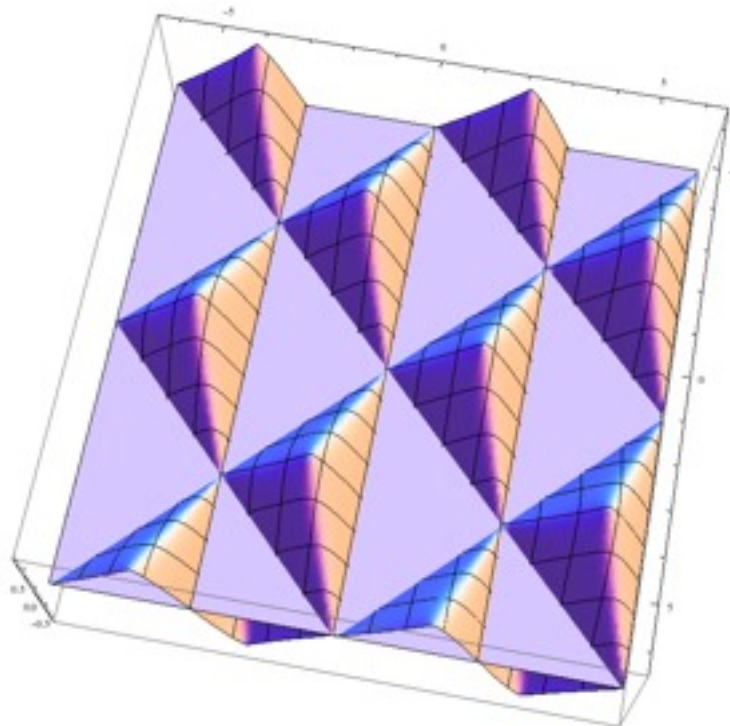
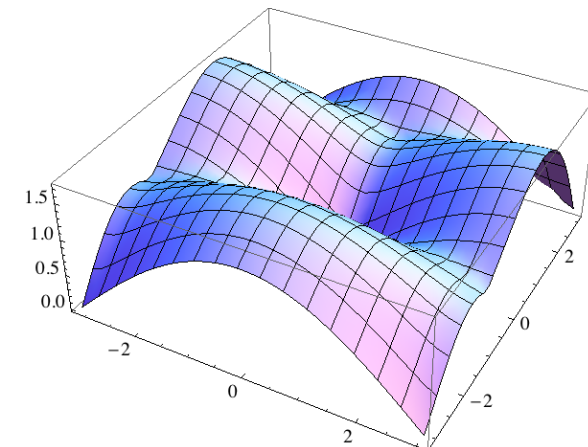
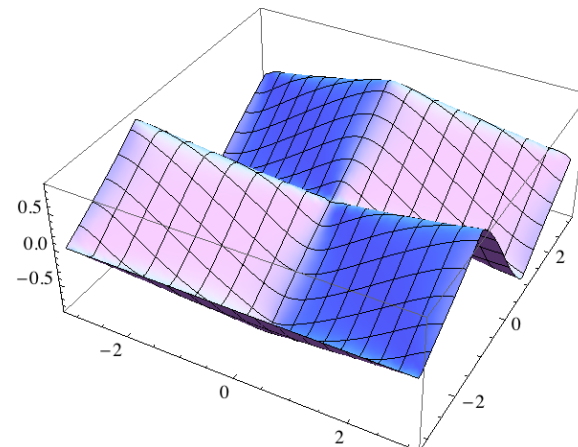
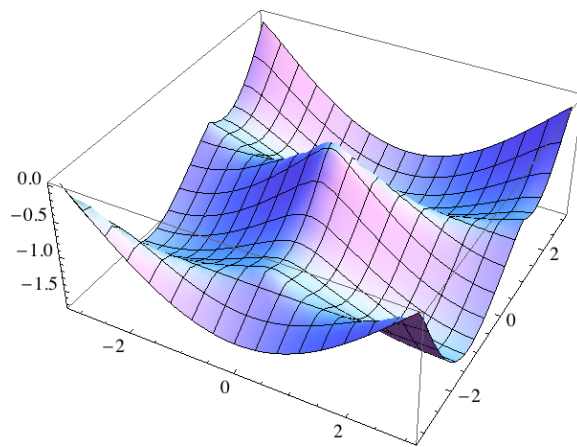
$$H = \sum_{\triangle} \tilde{\chi}_{ijk} - \sum_{\nabla} \tilde{\chi}_{ijk}$$



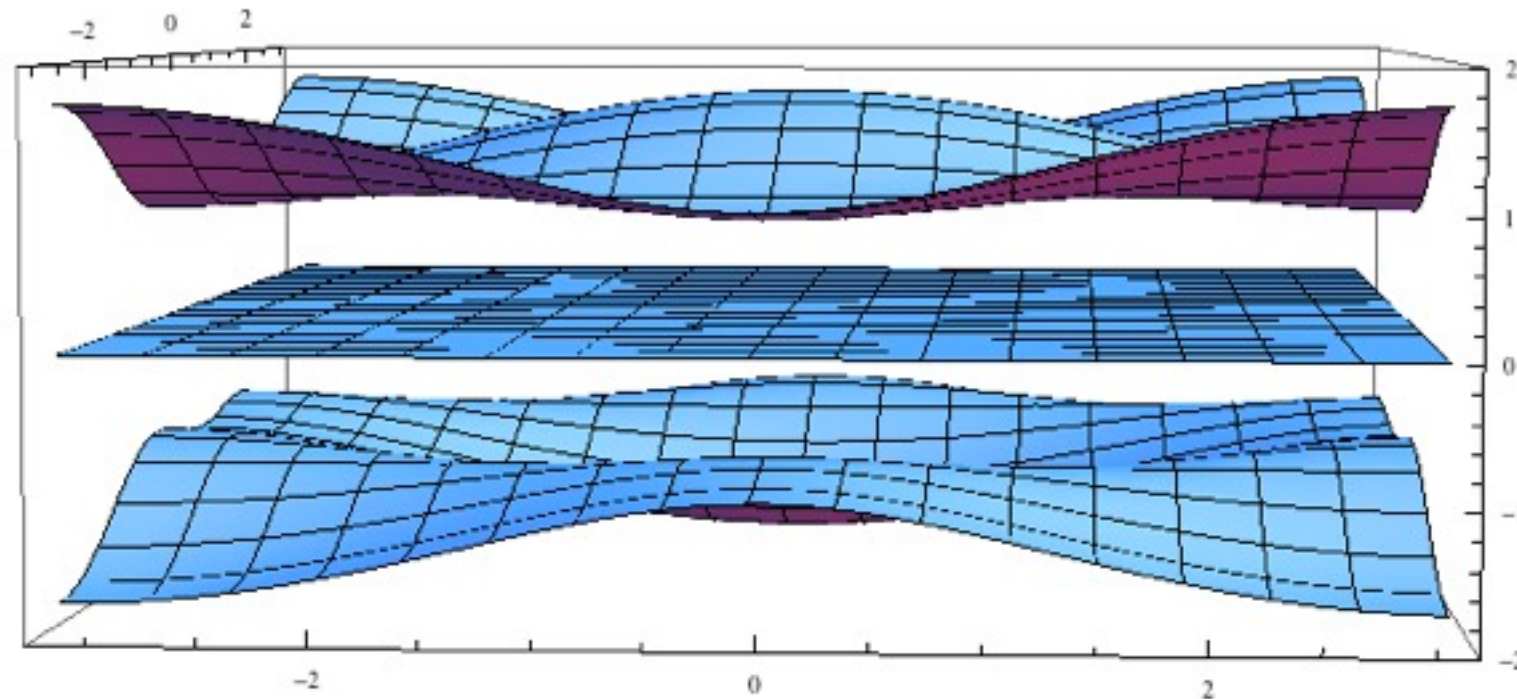
Majorana fermions: staggered

$$\tilde{\chi}_{ijk} = i(\gamma_i \gamma_j + \gamma_j \gamma_k + \gamma_k \gamma_i)$$

$$H = \sum_{\triangle} \tilde{\chi}_{ijk} - \sum_{\nabla} \tilde{\chi}_{ijk}$$



Majorana fermions: homogeneous



- Gapped spectrum
- Chern number of the top and bottom bands: $C = \pm 1$
- Dispersionless band: localized zero-energy states on hexagons

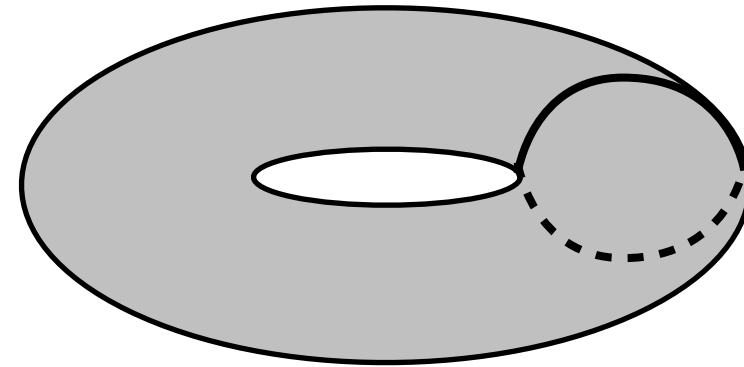
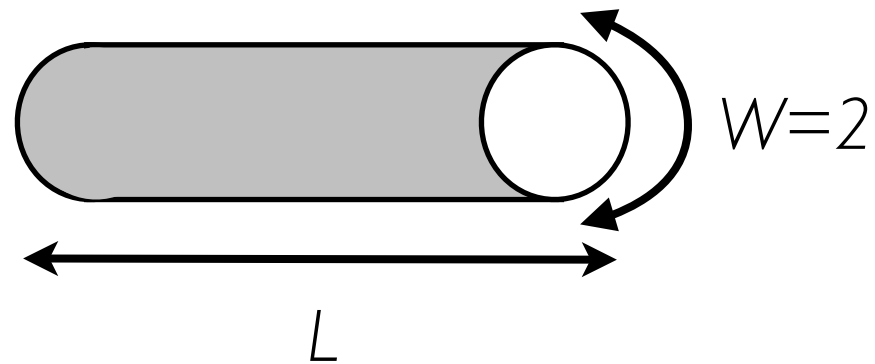
The same with spins?

$$\chi_{ijk} = \frac{i}{2} \vec{S}_i \cdot (\vec{S}_j \times \vec{S}_k)$$

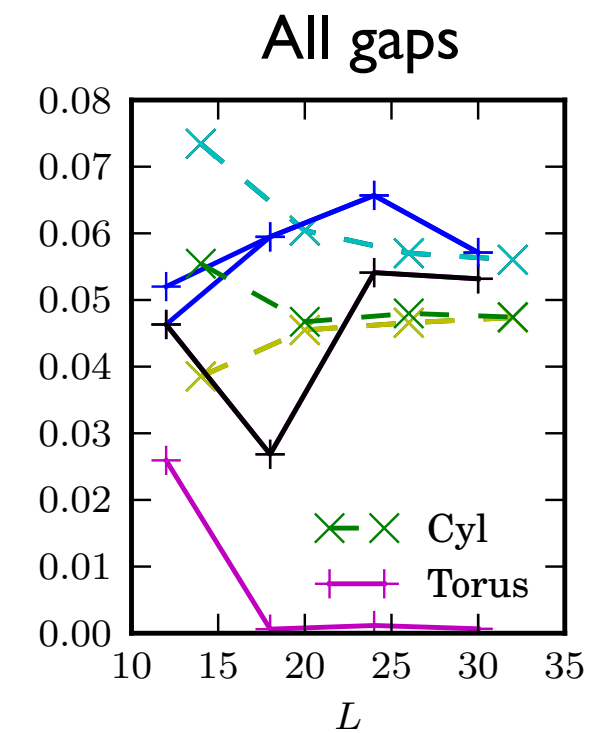
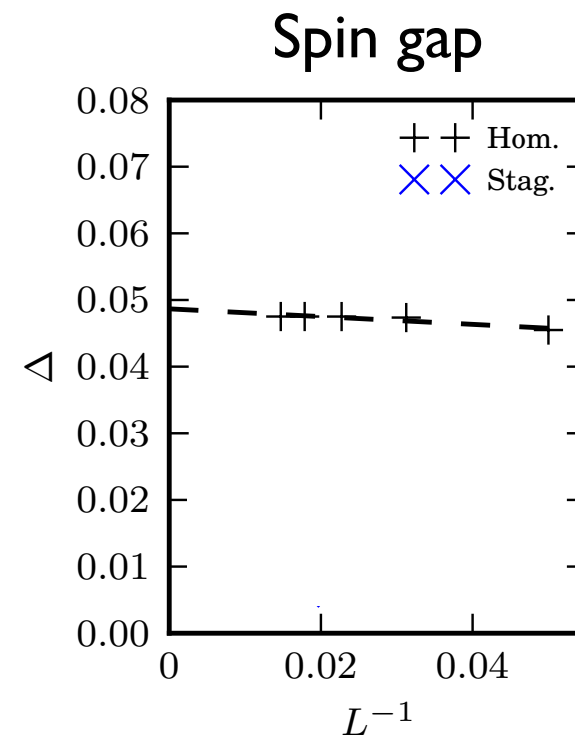
$$H = \sum_{\triangle} \chi_{ijk} \pm \sum_{\nabla} \chi_{ijk}$$

- Can only be solved numerically: DMRG and ED
 - quasi-1d systems with DMRG: “thin torus” or “thin cylinder” limit
- Our numerical agenda:
 - *Topological phase*
 - Spin gap
 - Degeneracy for torus vs cylinder
 - Edge state on strip vs cylinder
 - *Gapless phase*
 - Spin gap
 - Central charge

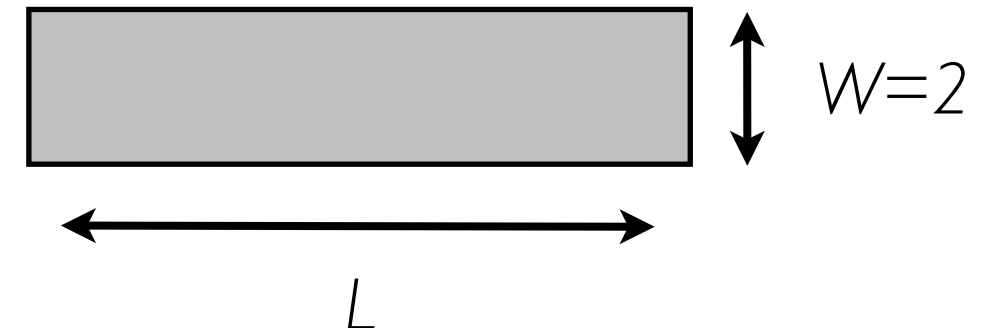
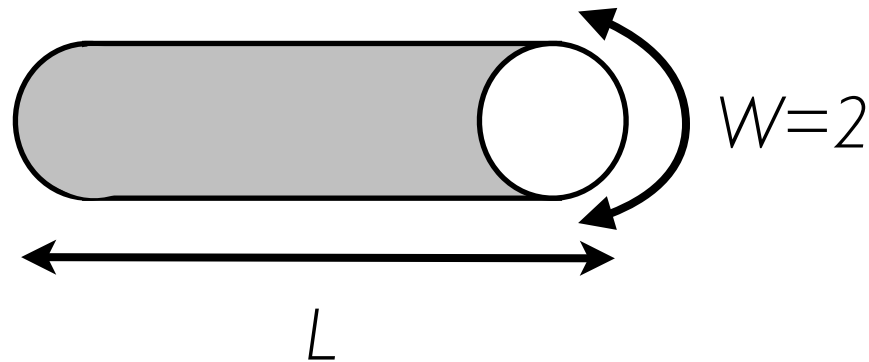
Spin gap



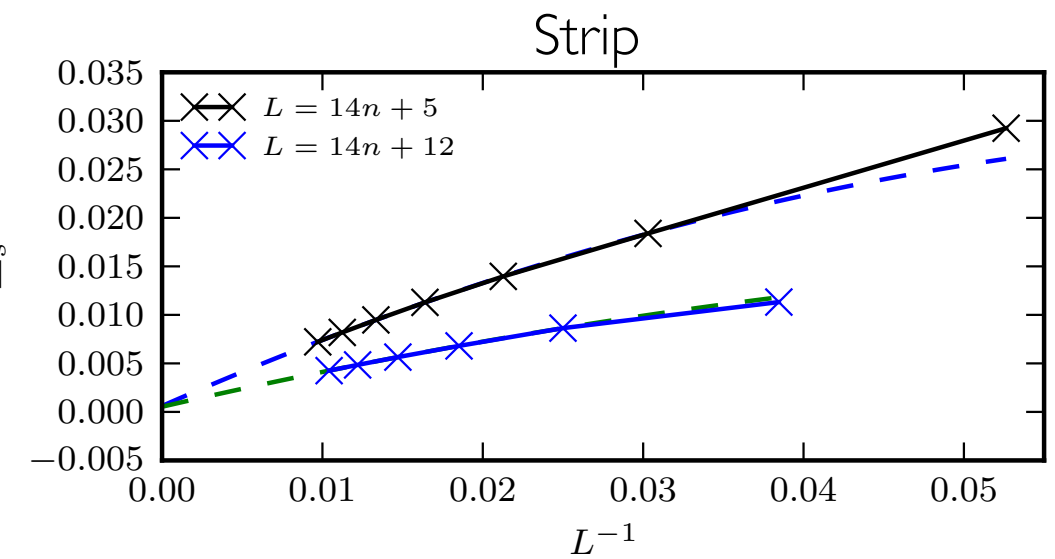
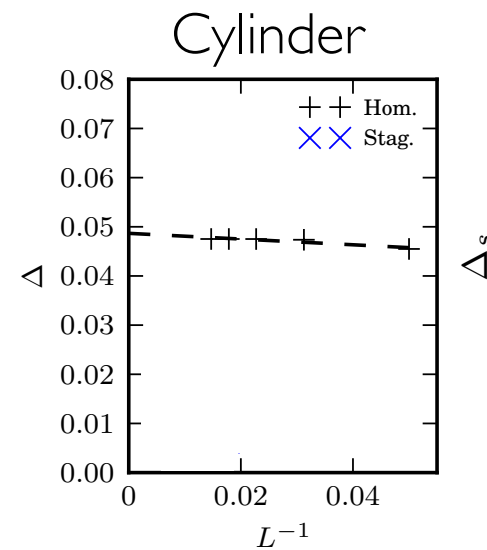
- $2 \times L$ cylinders (DMRG & ED) and tori (ED)
- DMRG with up to 2400 states
- ED up to 32 sites ($2 \times 5 \times 3 + 2$)
- *Finite gap ($\Delta \sim 0.05$)*
- *Additional low-lying state on the torus*



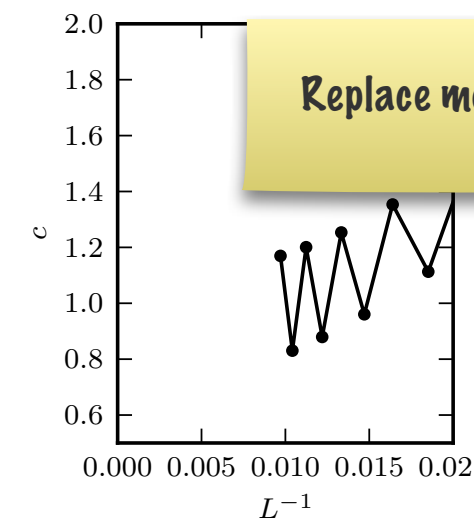
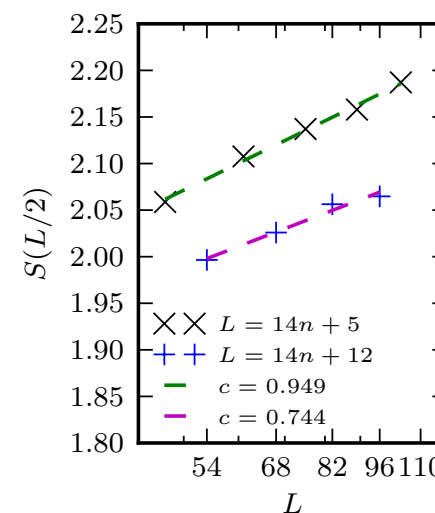
Edge state






- Spin gap for a fully open system:
 $\Delta \sim L^{-1}$



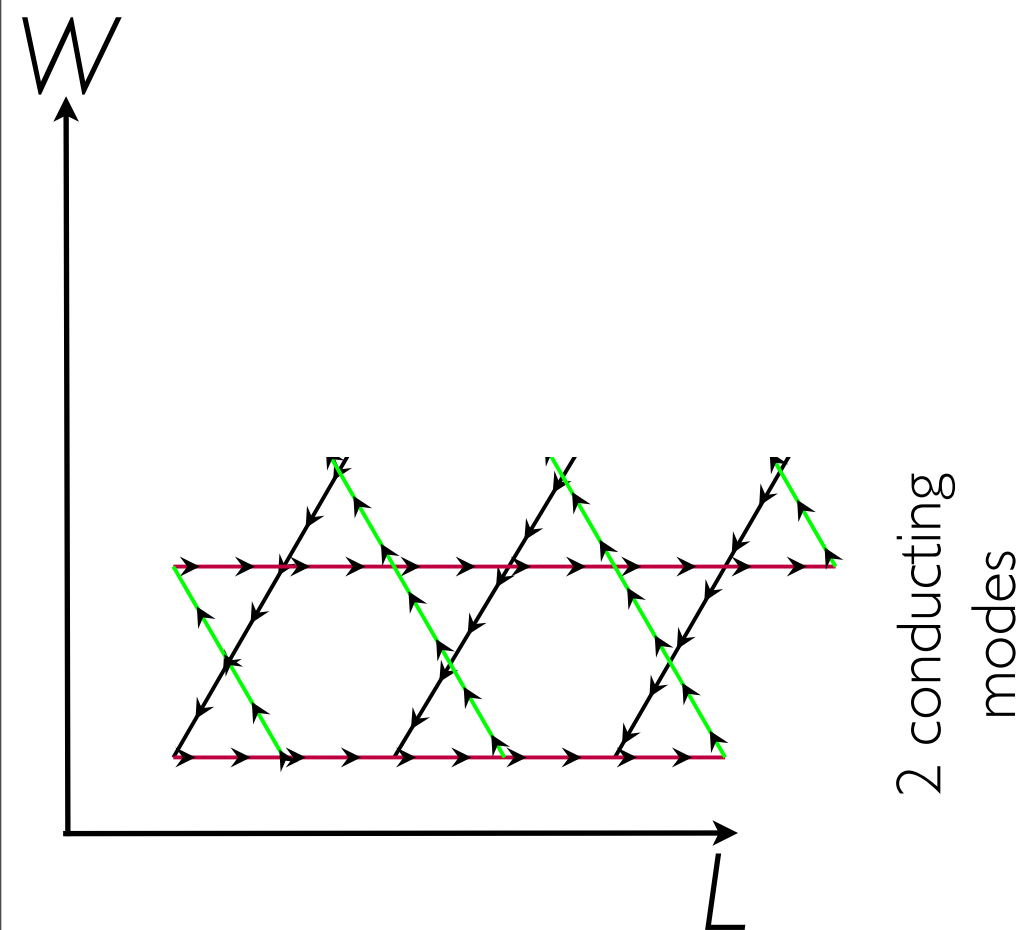
- Entanglement entropy:
 $S \sim \log(L)$
- Fit consistent with central charge $c=1$



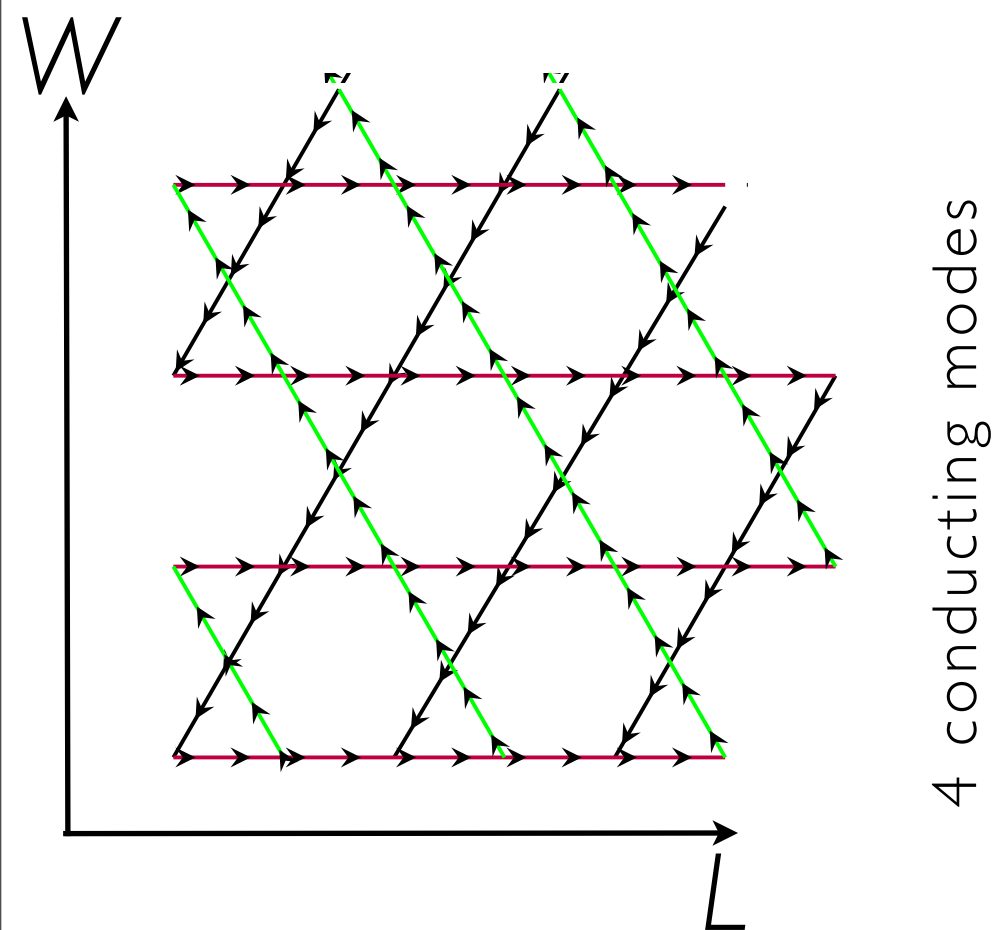
The same with spins?

- Our numerical agenda:
 - *Topological phase: consistent with $\nu=1/2$ bosonic Laughlin state*
 -  Spin gap
 -  Degeneracy for torus vs cylinder
 -  Edge state on strip vs cylinder
 - *Gapless phase*
 - Spin gap
 - Central charge

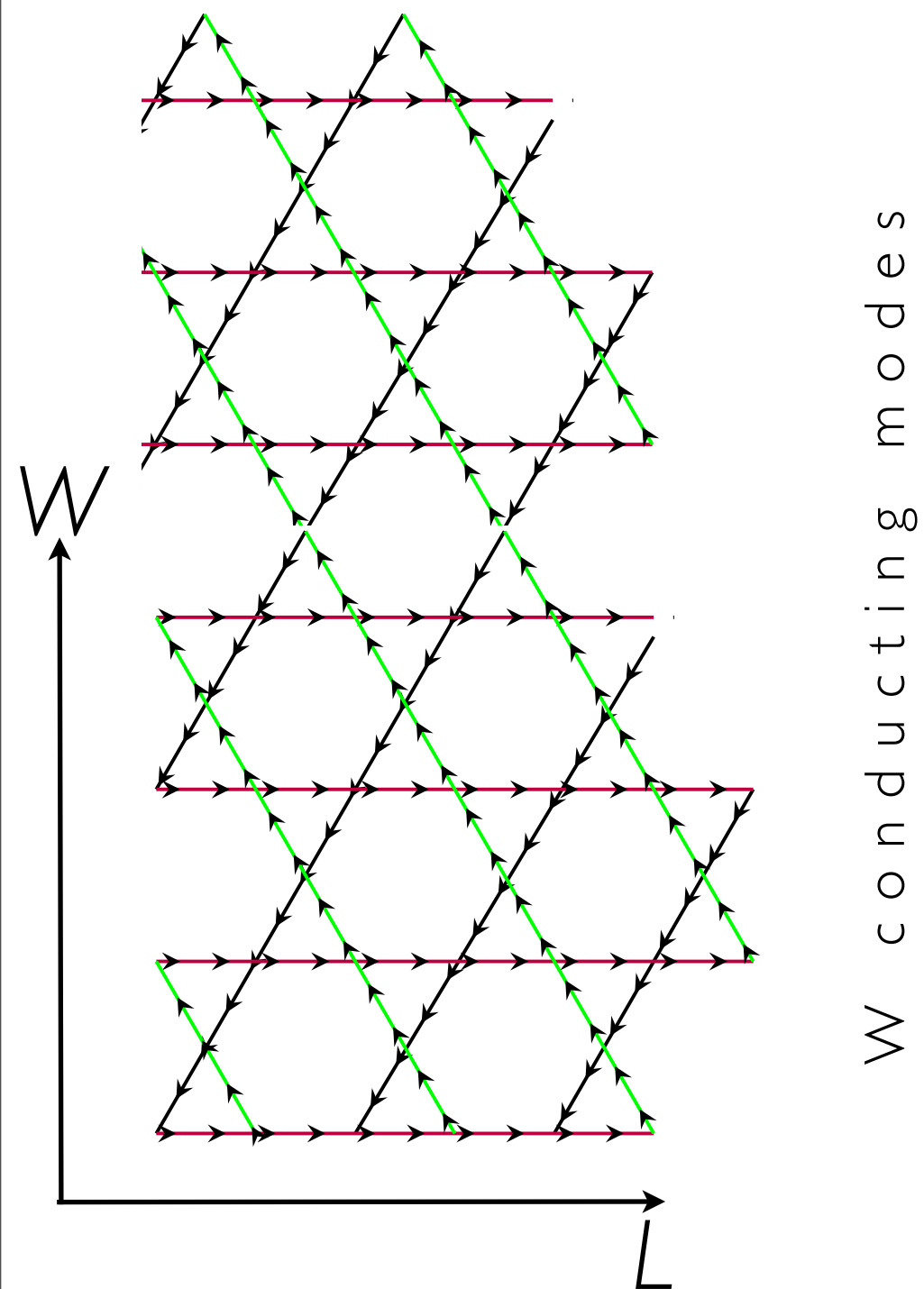
Quasi-1d predecessors



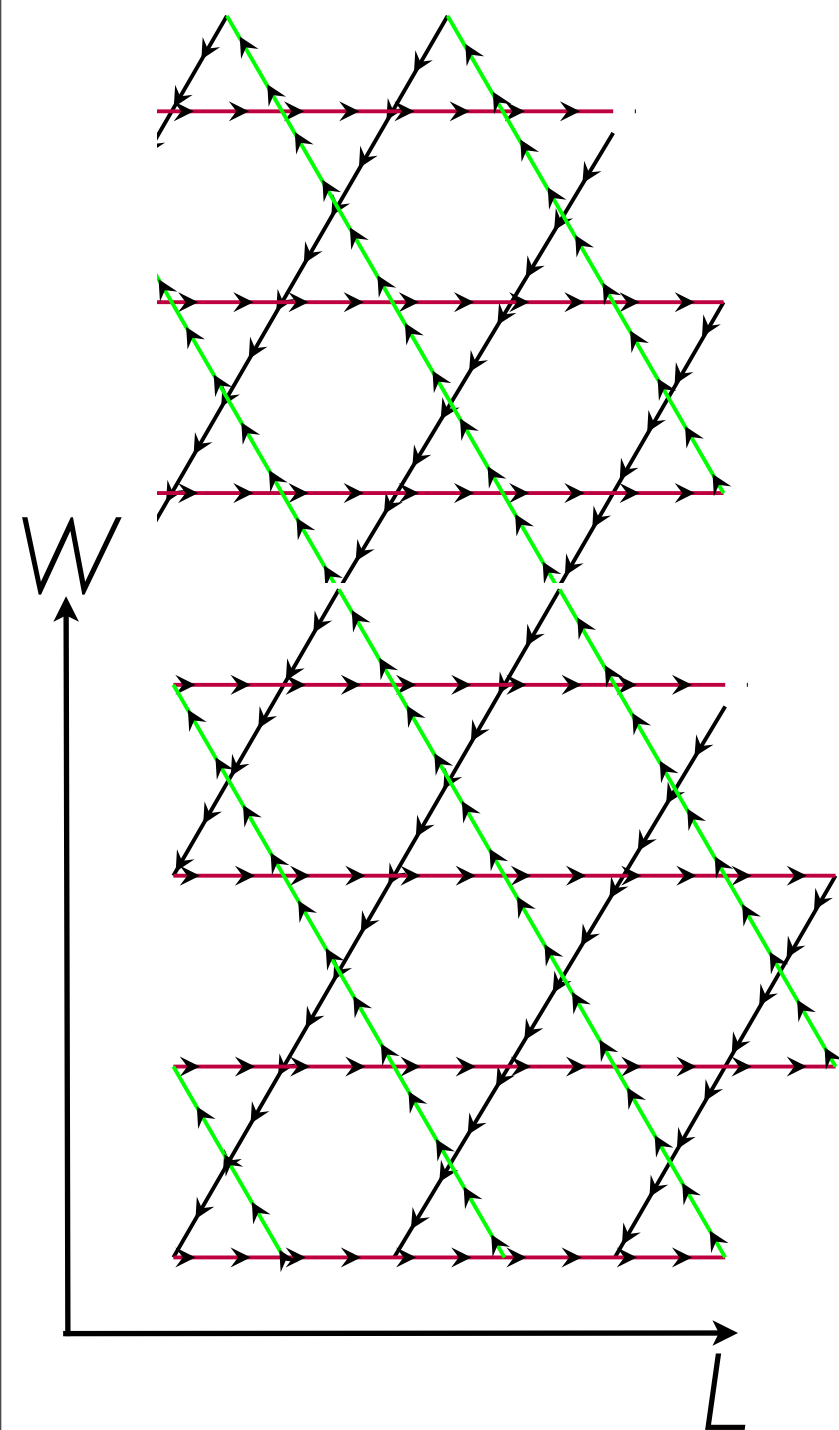
Quasi-1d predecessors



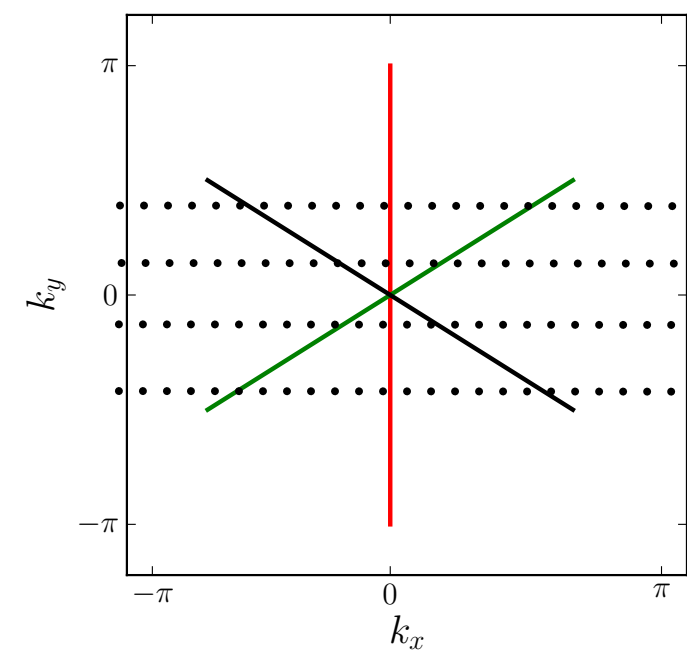
Quasi-1d predecessors



Quasi-1d predecessors



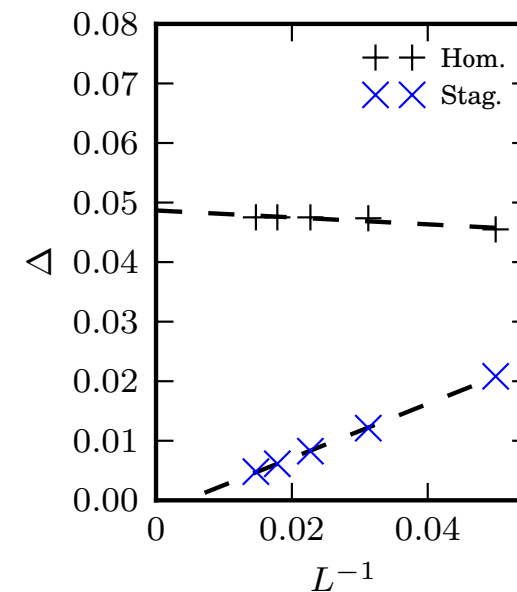
W conducting modes



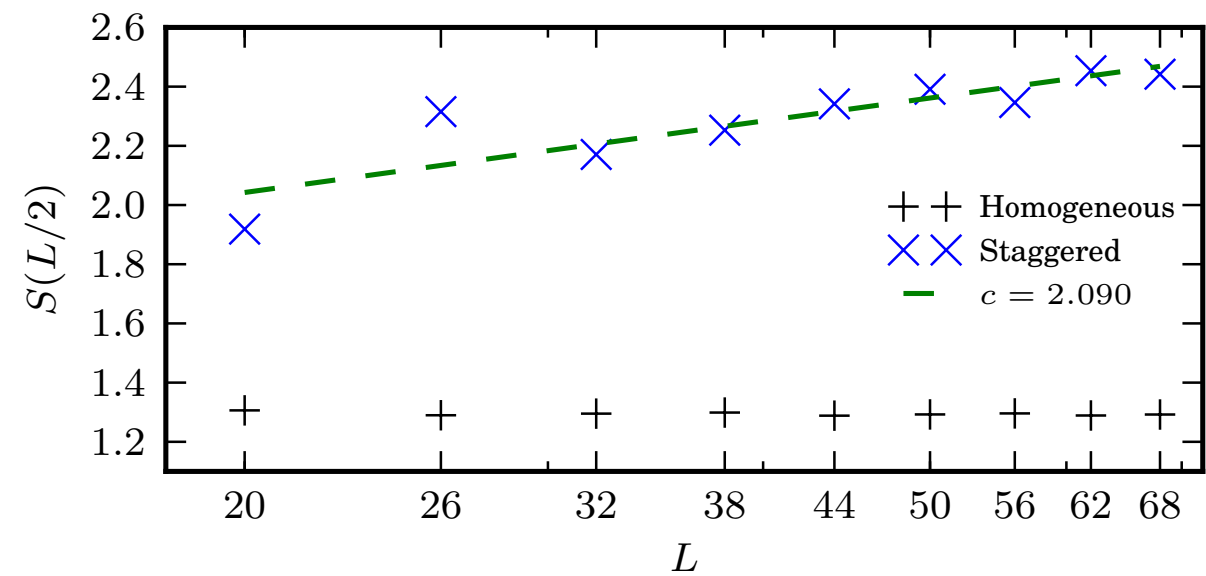
W
 \sim # of k_y points
 \sim number of gapless modes

Staggered phase






- Gap vanishes as $1/L$



- Entanglement entropy at the center of the system
 - Gapped: $S \sim \text{const}$
 - Gapless: $S \sim \log(L)$
- Fit yields expected $c=2$*



The same with spins?

- Our numerical agenda:
 - *Topological phase: consistent with $\nu=1/2$ bosonic Laughlin state*
 -  Spin gap
 -  Degeneracy for torus vs cylinder
 -  Edge state on strip vs cylinder
 - *Gapless phase*
 -  Spin gap
 -  Central charge for $W=2$

Conclusions & Outlook

- We can construct a topological spin liquid and a gapless phase with an emergent surface of excitations from coupling the edge states encircling seeds of a topological phase.
- We can predict the shape of the Fermi surface from the edge state picture.
- Ongoing work:
 - Understand the effect of interactions on the gapless lines in the $SU(2)$ spin liquid (stability of the gapless phase)
- Future work:
 - Other numerical approaches for gapless phase? VMC?
 - Possible relations to other models?

Thank you for your attention!