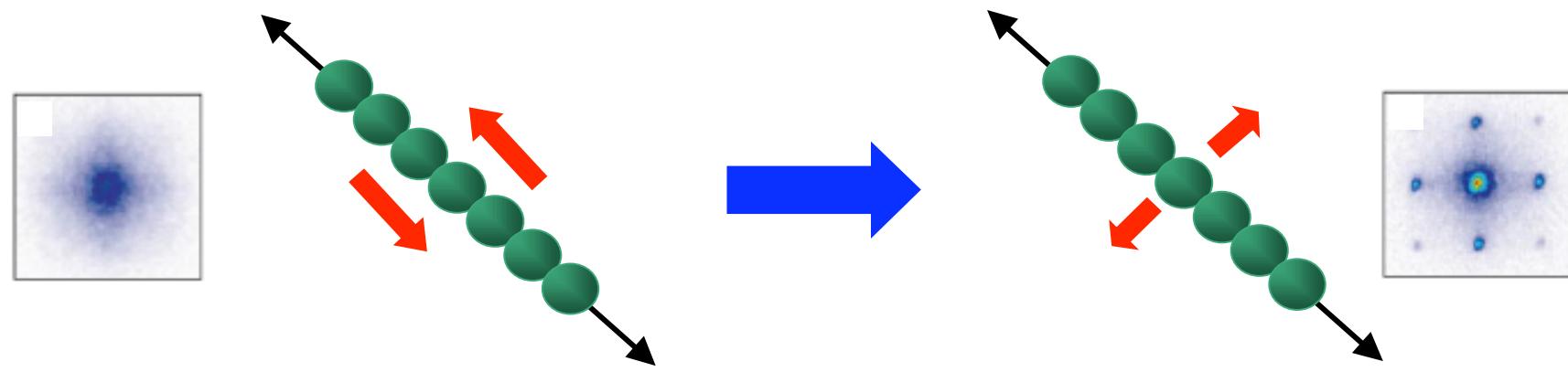


How cold atoms live in (and escape from) 1D



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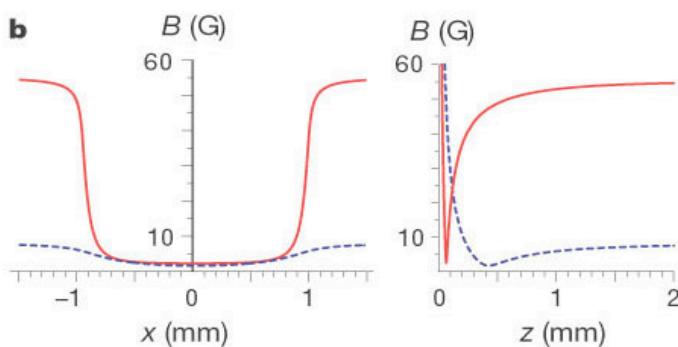
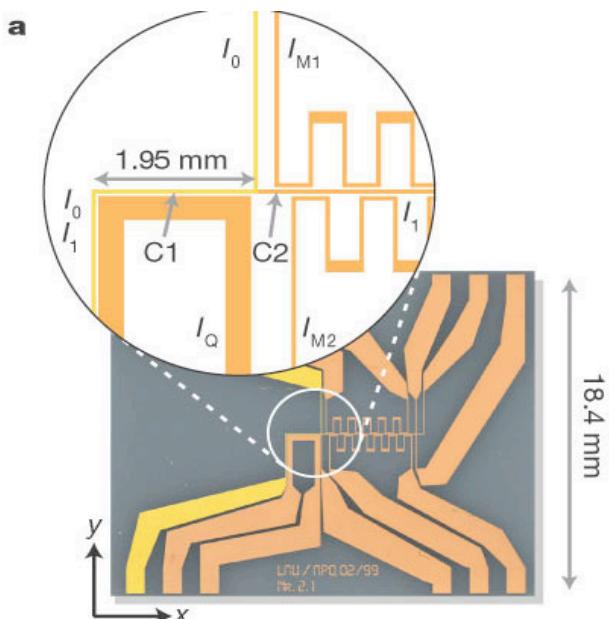
²⁾ Donostia Int'l Physics Center (DIPC), San Sebastian (Spain)

³⁾ École de Physique, University of Geneva (Switzerland)

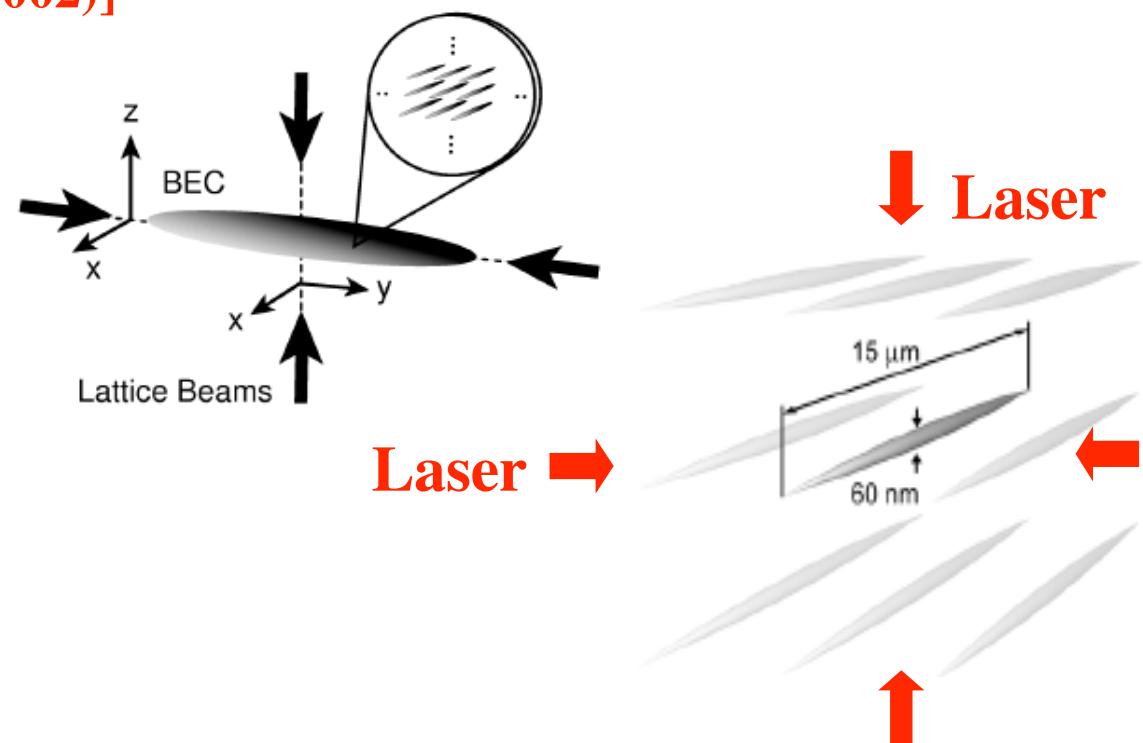
The 1D Cold Atom Zoo

- Cold atoms on a chip:

[W. Haensel *et al.* *Nature* **413** (2002)]



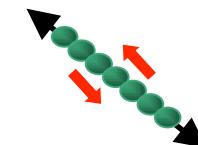
- 2D optical lattices:



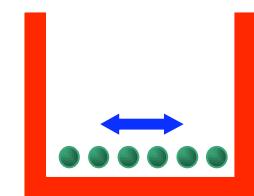
[M. Greiner *et al.* *PRL* **87** (2001),
H. Moritz, *PRL* **91** (2003),
T. Stoferle *et al.* *PRL* **92** (2004)]

Outline:

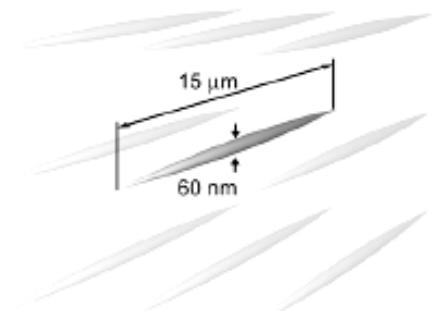
- Life in one dimension: **Luttinger liquids**



- Life in a trap: **bosonic atoms in a 1D box**

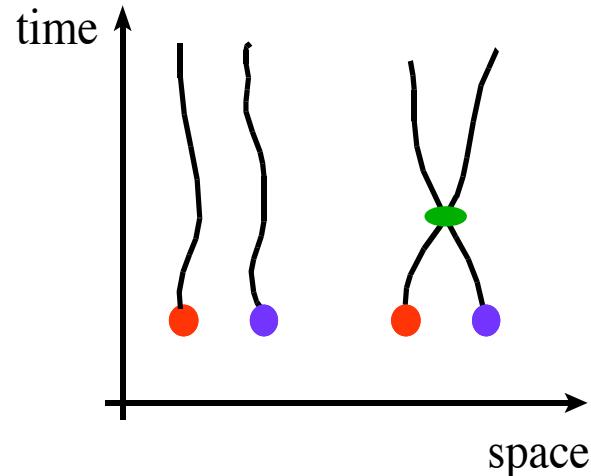


- Experiments in 3D and 2D **optical lattices**



- Escaping from 1D: **Phases of a 2D optical lattice**

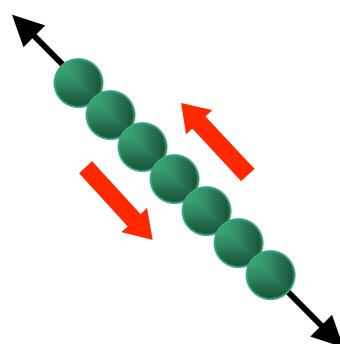
A crash course in Luttinger liquids (1)



- What is it made of ? Bosons or Fermions?

“In 1D [...] the symmetry of the wave function cannot be tested by a continuous change of coordinates that exchanges particles without close approach (collision). Thus interaction and statistics effects cannot be separated.”

[FDM Haldane, PRL 47 (1981)]



- Collective modes exhaust the low-energy spectrum:

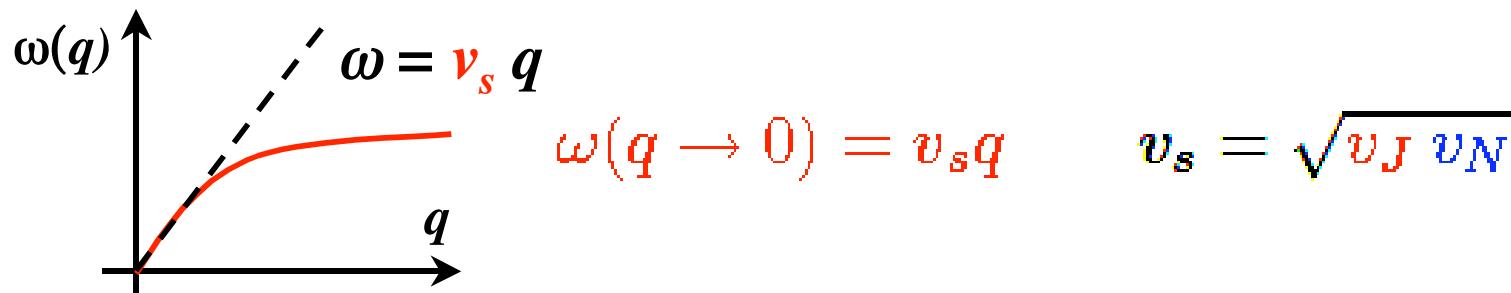
$$H = \frac{\hbar}{2\pi} \int dx \left[v_J (\partial_x \phi)^2 + v_N (\partial_x \theta)^2 \right]$$

phase stiffness

density stiffness

A crash course in Luttinger liquids (2)

- Collective modes have linear dispersion:



$$K = \sqrt{\frac{v_J}{v_N}}$$

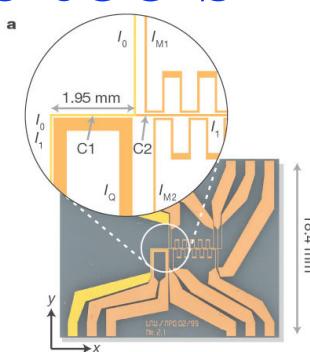
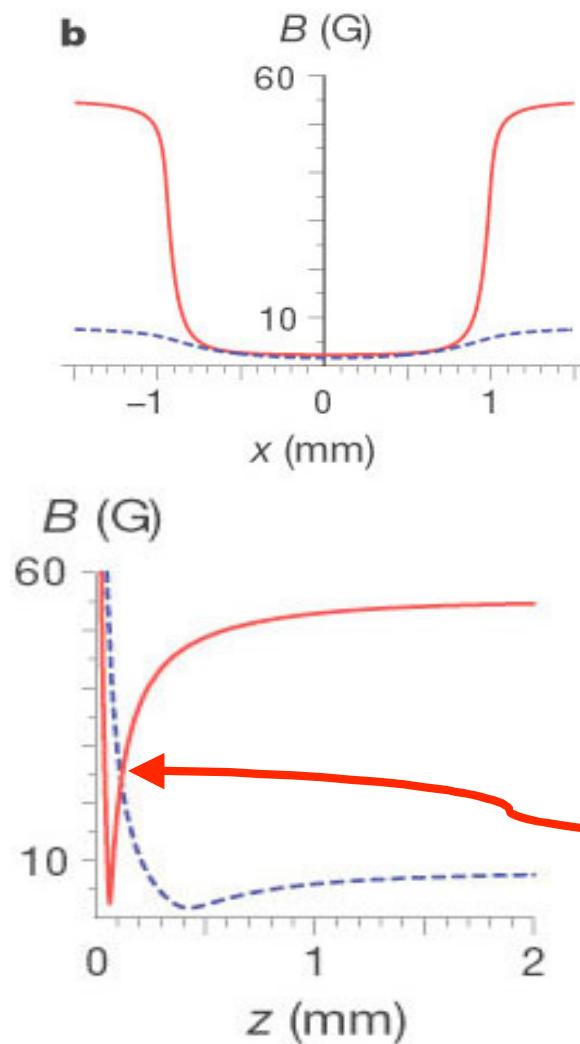
- The road map: range of K for the Bose-Hubbard model



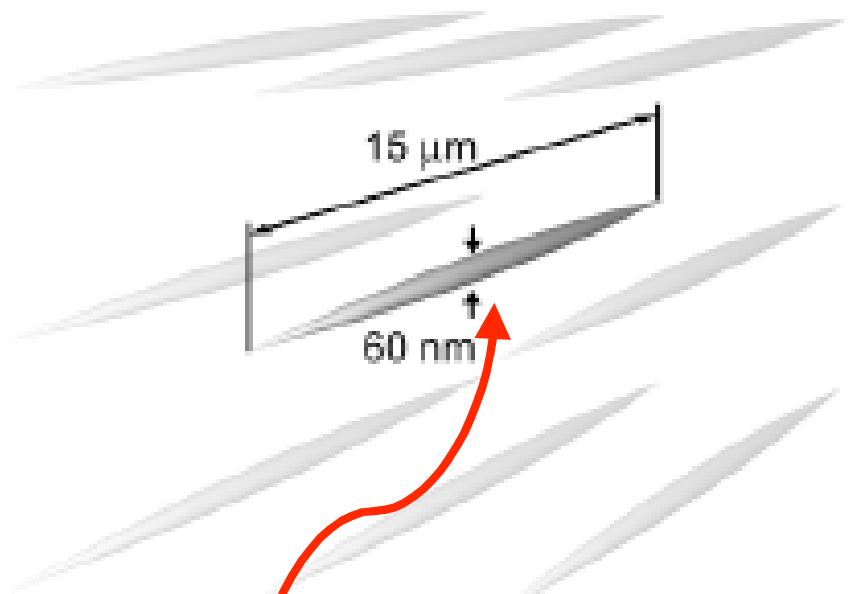
[T. Giamarchi, “Quantum Physics in One dimension”, Oxford University Press (2003)]

When your cage is too small: mesoscopic LL's

- Cold atoms on a chip:



- 2D Optical lattices:



$N_0 \sim 10$ to 10^3 atoms

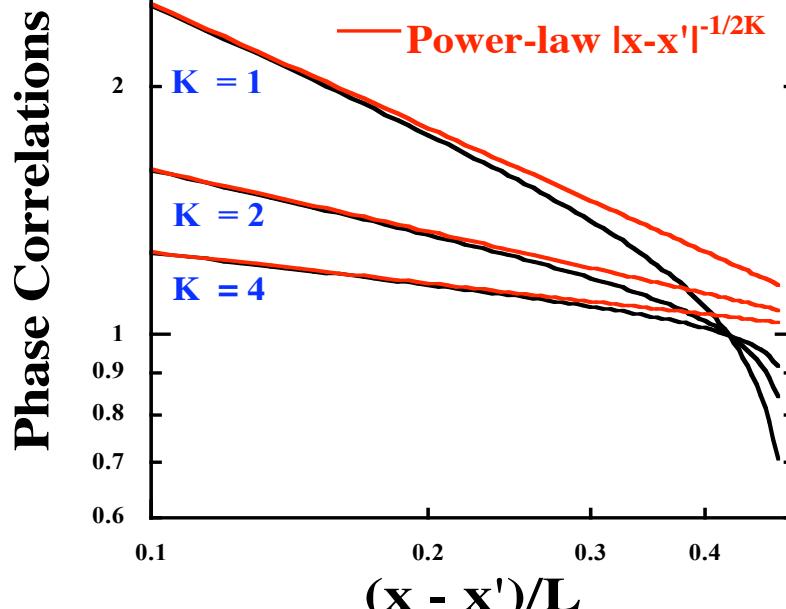
A toy model: bosonic atoms in a 1D box

- Phase correlations: $\langle \Psi^\dagger(x)\Psi(x') \rangle \simeq \rho_0 \langle e^{-i\phi(x)} e^{i\phi(x')} \rangle$

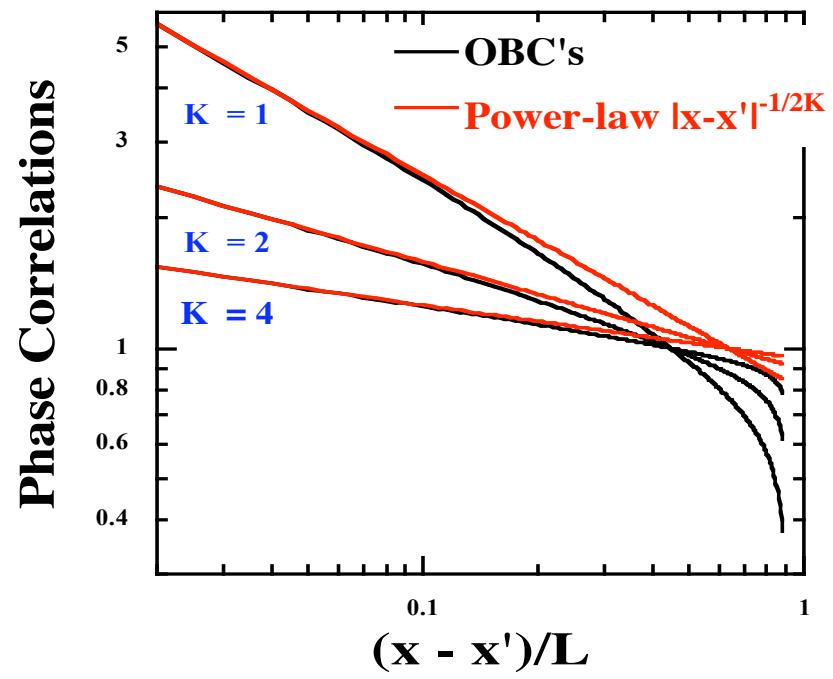
[Thermodynamic limit ($T = 0$)

$$g_1(x) = \langle \Psi^\dagger(x)\Psi(0) \rangle = \frac{A}{x^{1/2K}} \quad]$$

$x'/L = 0.5$



$x'/L = 0.1$

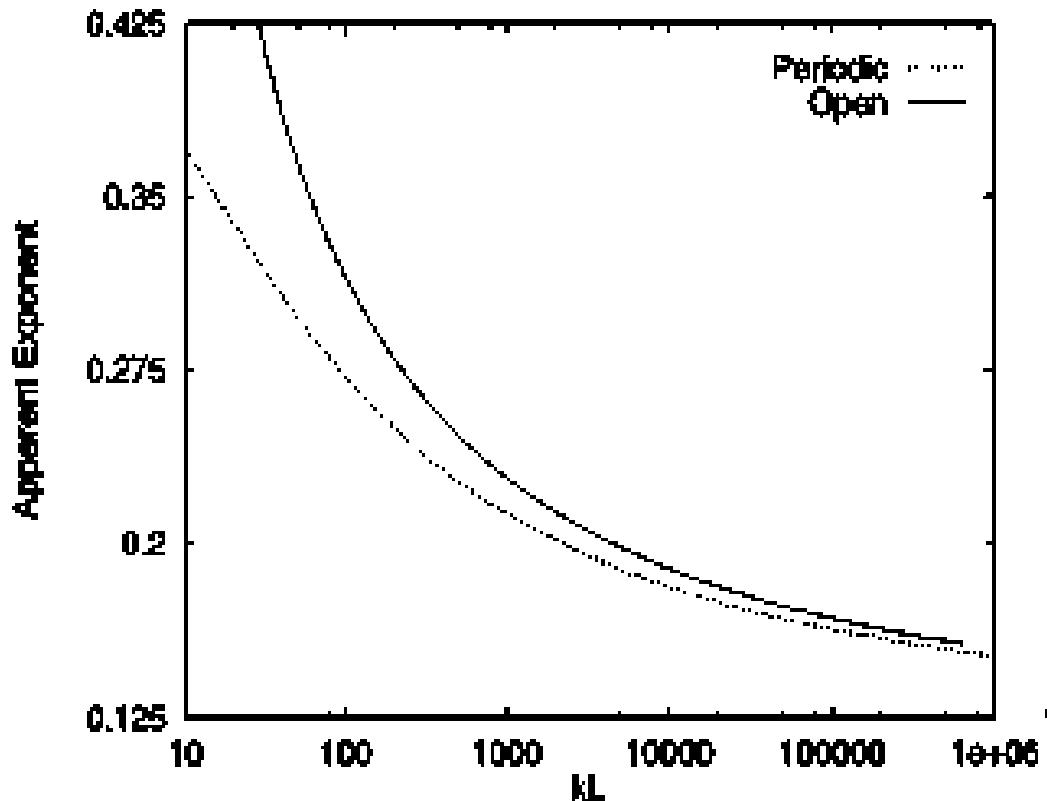


$\nabla = x$

$\bullet = x'$

[MAC, *Europhys. Lett.* **59** (2002) & *J. Phys. B* **37** (2004)]

Momentum distribution in a 1D finite LL



- Fermi Hubbard model $U = +\infty$
(away from half-filling)

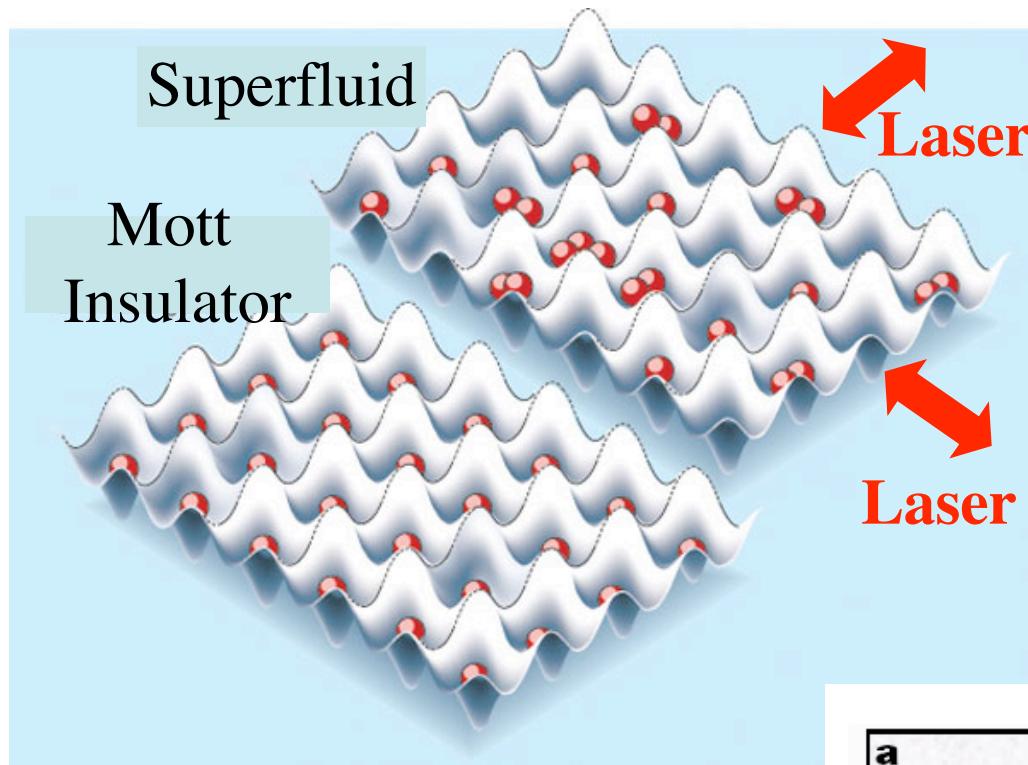
$$\alpha_{\text{app}}(kL) = \frac{dn(k, L)}{d \ln kL}$$
$$n(k, L) \sim |k|^{\alpha_{\text{app}}}$$

$$\alpha(L \rightarrow \infty) = 0.125$$

[S. Eggert *et al.* PRL **26** (1996)]

- Momentum distribution: $n(p, L) = (\rho_0 L)^{1 - \frac{1}{2K}} I(pL)$

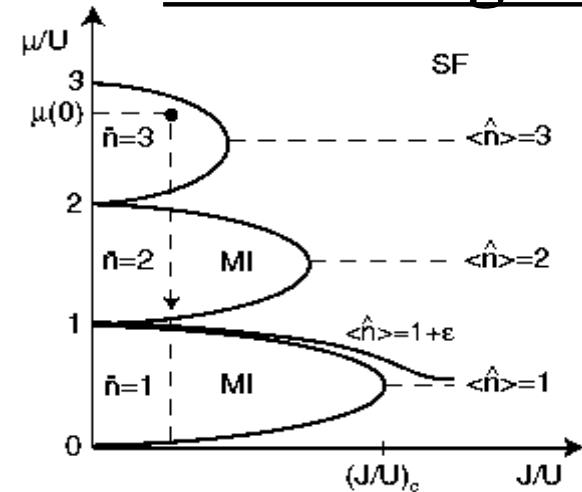
Experiments: 3D optical lattices



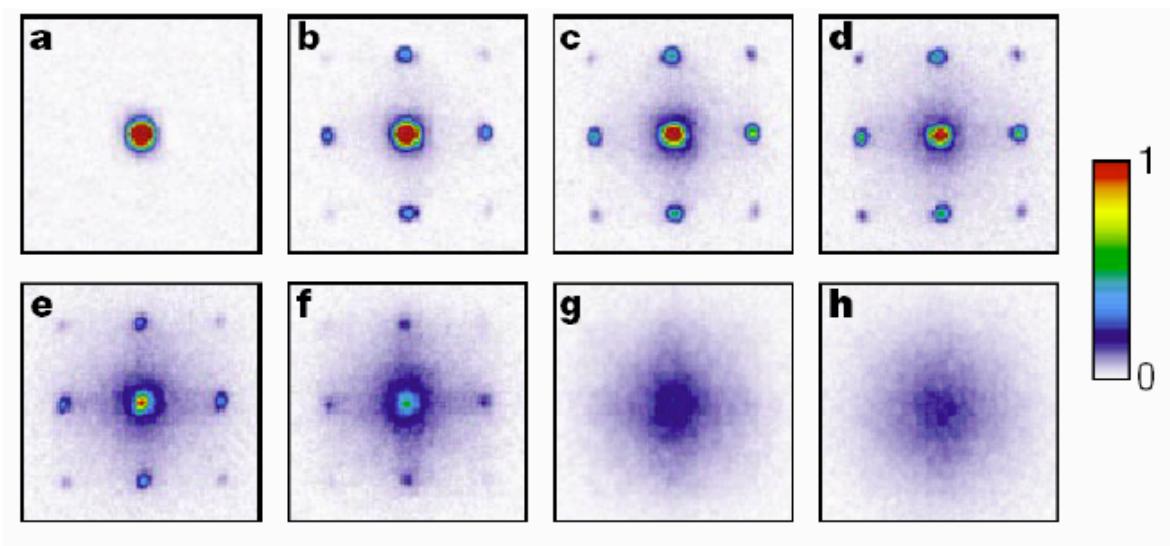
Superfluid to Mott insulator
transition in a 3D optical lattice

[D. Jaksch *et al.* PRL 81 (1998)]
[M Greiner *et al.* Nature, 415 (2002)]

- Phase diagram:



[MPA Fisher *et al.* PRB 40 (1989)]



Excitation spectrum: Bragg spectroscopy (3D)

- Bose-Hubbard model:

$$\begin{aligned}
 H_{\text{BH}} &= \sum_{\mathbf{R}, m} \left[-\frac{J_x}{2} \left(b_{m+1}^\dagger(\mathbf{R}) b_m(\mathbf{R}) + b_m^\dagger(\mathbf{R}) b_{m+1}(\mathbf{R}) \right) + \epsilon_m(\mathbf{R}) b_m^\dagger(\mathbf{R}) b_m(\mathbf{R}) \right] \\
 &\quad - J \sum_{\langle \mathbf{R}, \mathbf{R}' \rangle, m} b_m^\dagger(\mathbf{R}) b_m(\mathbf{R}') + U \sum_{\mathbf{R}, m} b_m^\dagger(\mathbf{R}) b_m^\dagger(\mathbf{R}) b_m(\mathbf{R}) b_m(\mathbf{R})
 \end{aligned}$$

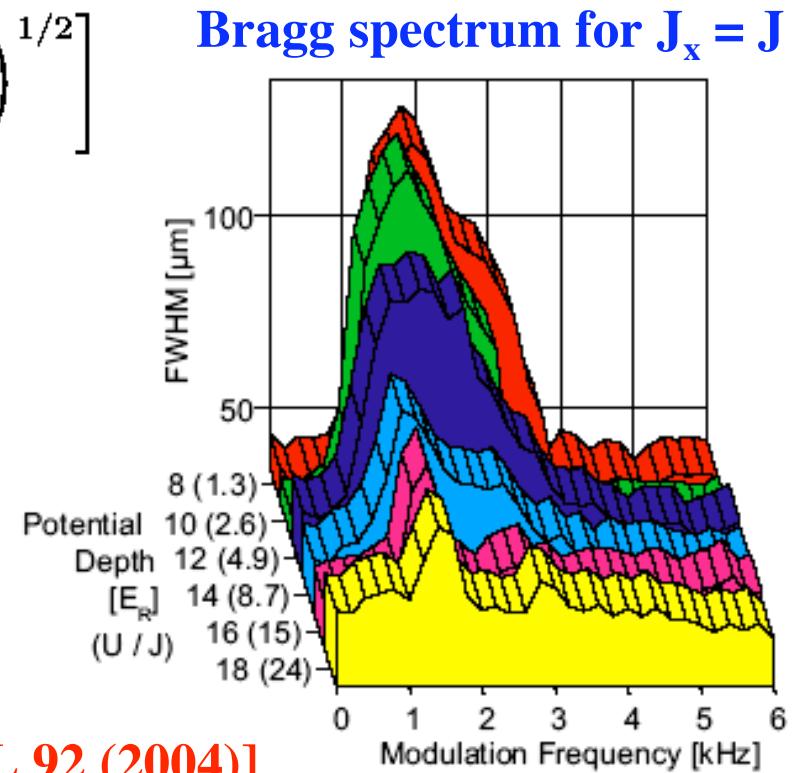
$$J_\alpha \left(\frac{V_{0\alpha}}{E_R} \gg 1 \right) = \frac{4E_R}{\sqrt{\pi}} \left(\frac{V_{0\alpha}}{E_R} \right)^{1/4} \exp \left[-2 \left(\frac{V_{0\alpha}}{E_R} \right)^{1/2} \right]$$

$(\alpha = x, y, z)$

- 2-photon Bragg spectroscopy:

$$V_{0x} \rightarrow V_{0x}(t) = [V_{0x} + A_{\text{mod}} \sin(2\pi\nu_{\text{mod}} t)]$$

(i.e. modulate the axial optical potential)

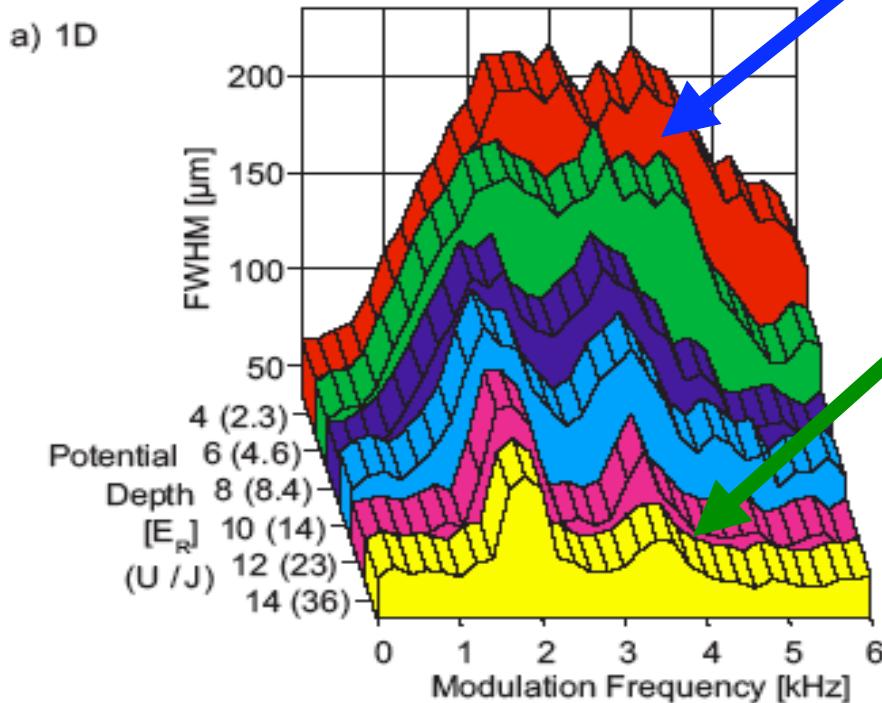


[T. Stoferle *et al.* PRL 92 (2004)]

Exc. spectrum: Bragg spectroscopy ($J_x \gg J$)

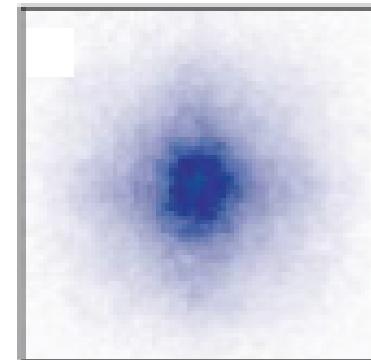
- $V_{0x,y} = 30 E_R$:

Bragg spectrum for $J_x \gg J$



Broad Spectrum: 1D SF (LL)

Discrete features: 1D MI



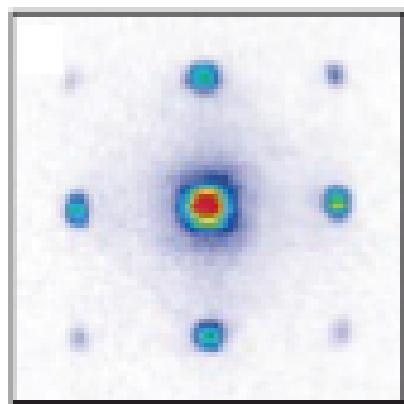
$$\left(\frac{J_x}{J}\right)_{\max} \simeq 25$$

[T. Stoferle *et al.* PRL 92 (2004)]

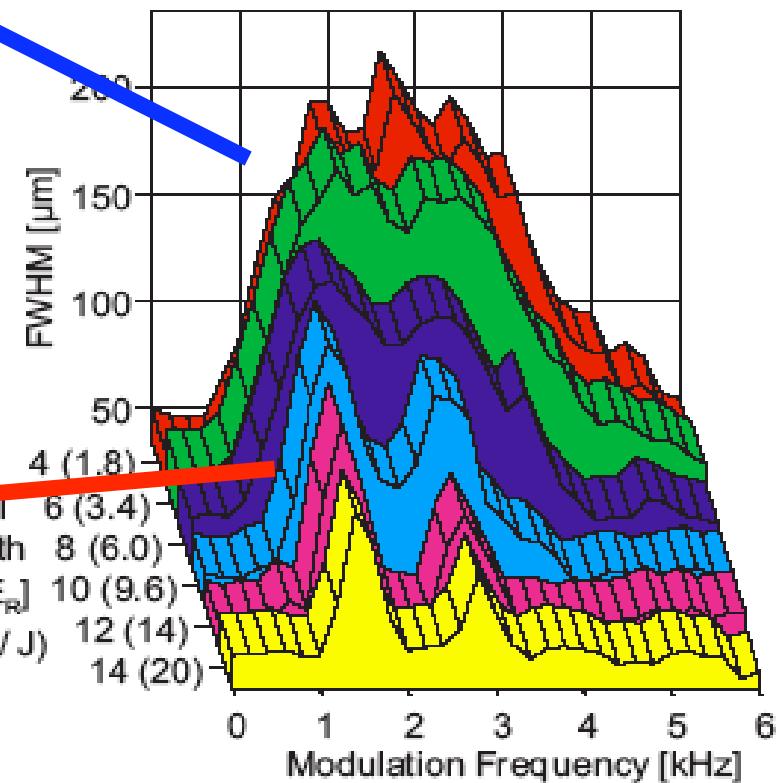
Excitation spectrum: deconfinement!

[T. Stoferle *et al.* PRL 92 (2004)]

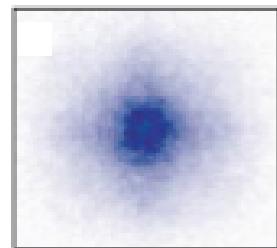
- 3D Superfluid



- $V_0 = 20 E_R$:
3D SF to 1D MI



- 1D Mott Insulator

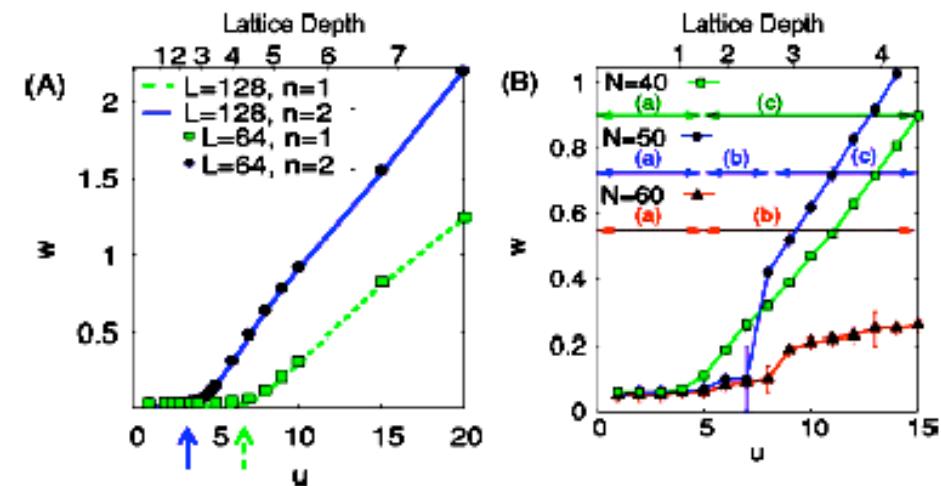
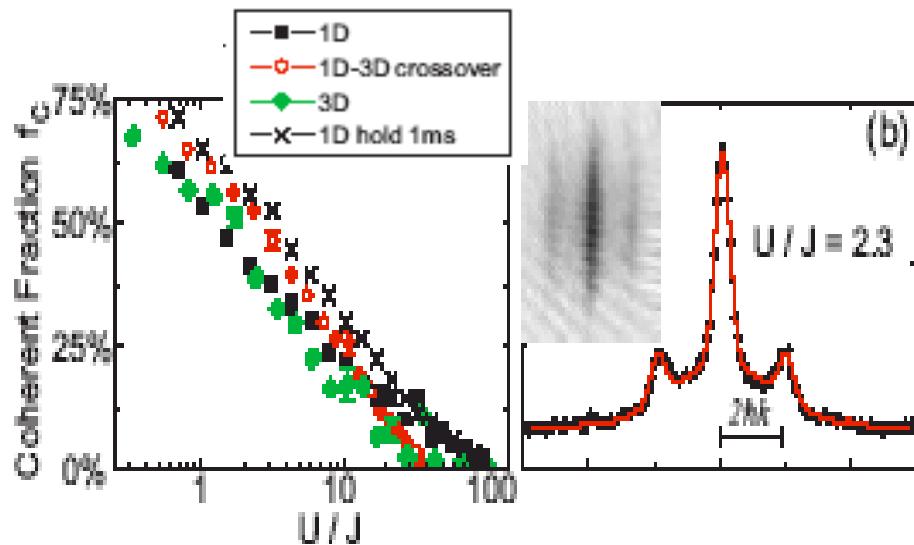


$$\frac{J_x}{J} \simeq 10$$



By reducing the axial hopping intertube coherence is destroyed!!

Where does the phase transition takes place?



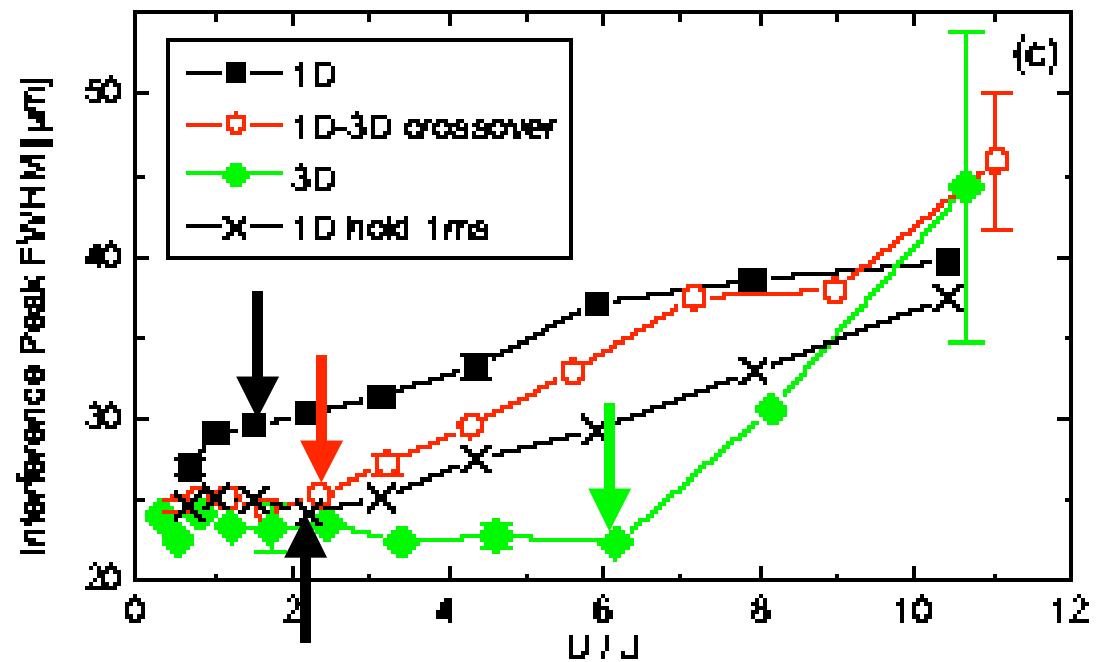
[C. Kollath *et al.* PRA **69** (2004)]

Large quantum depletion!!

$$\left(\frac{U}{zJ}\right)_{1D} \approx 1.9$$

$$\left(\frac{U}{zJ}\right)_{3D} \approx 5.8$$

[T. Stoferle *et al.* PRL **92** (2004)]

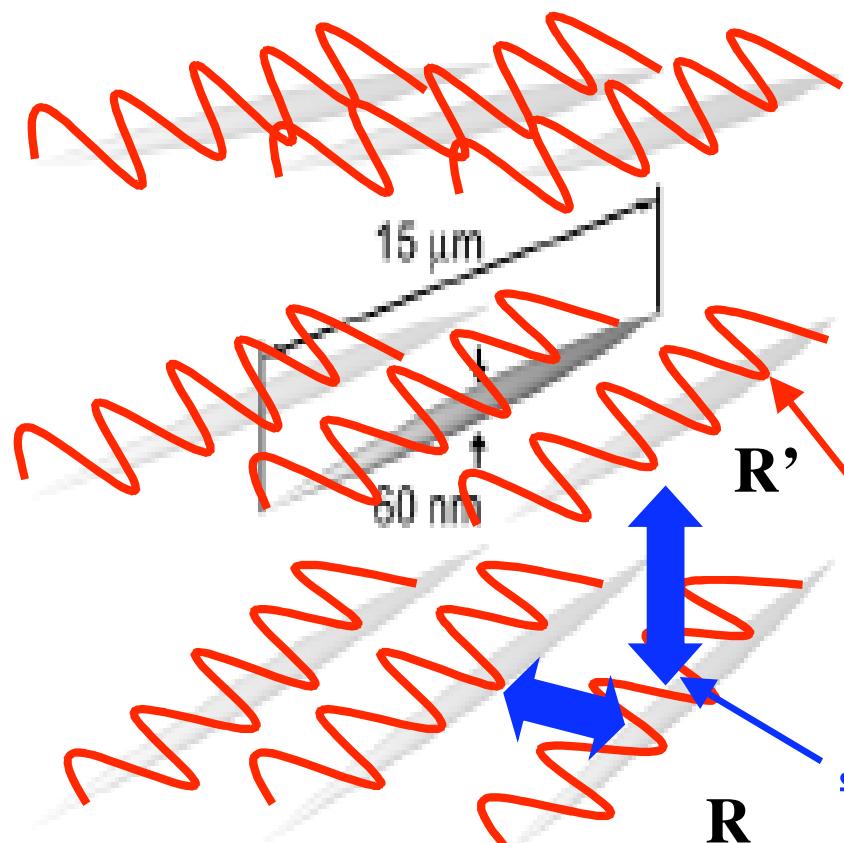


2D optical lattices: effective low-energy theory

- Through “bosonization”:

$$H_{\text{eff}} = \frac{\hbar v_s}{2\pi} \sum_{\mathbf{R}} \int_0^L dx \left[\frac{1}{K} (\partial_x \theta_{\mathbf{R}}(x))^2 + K (\partial_x \phi_{\mathbf{R}}(x))^2 \right]$$

$$+ \frac{\hbar v_s g_u}{2\pi a^2} \sum_{\mathbf{R}} \int_0^L dx \cos(2\theta_{\mathbf{R}}(x) + \delta\pi x)$$
$$- \frac{\hbar v_s g_J}{2\pi a^2} \sum_{\langle \mathbf{R}, \mathbf{R}' \rangle} \int_0^L dx \cos(\phi_{\mathbf{R}}(x) - \phi_{\mathbf{R}'}(x))$$



“Mott” potential: localizes atoms

Josephson coupling: delocalizes atoms

[AFH, MAC & T Giamarchi, PRL 92 (2004)]

2D optical lattices: phase diagram at T = 0

- Renormalization-group flow:

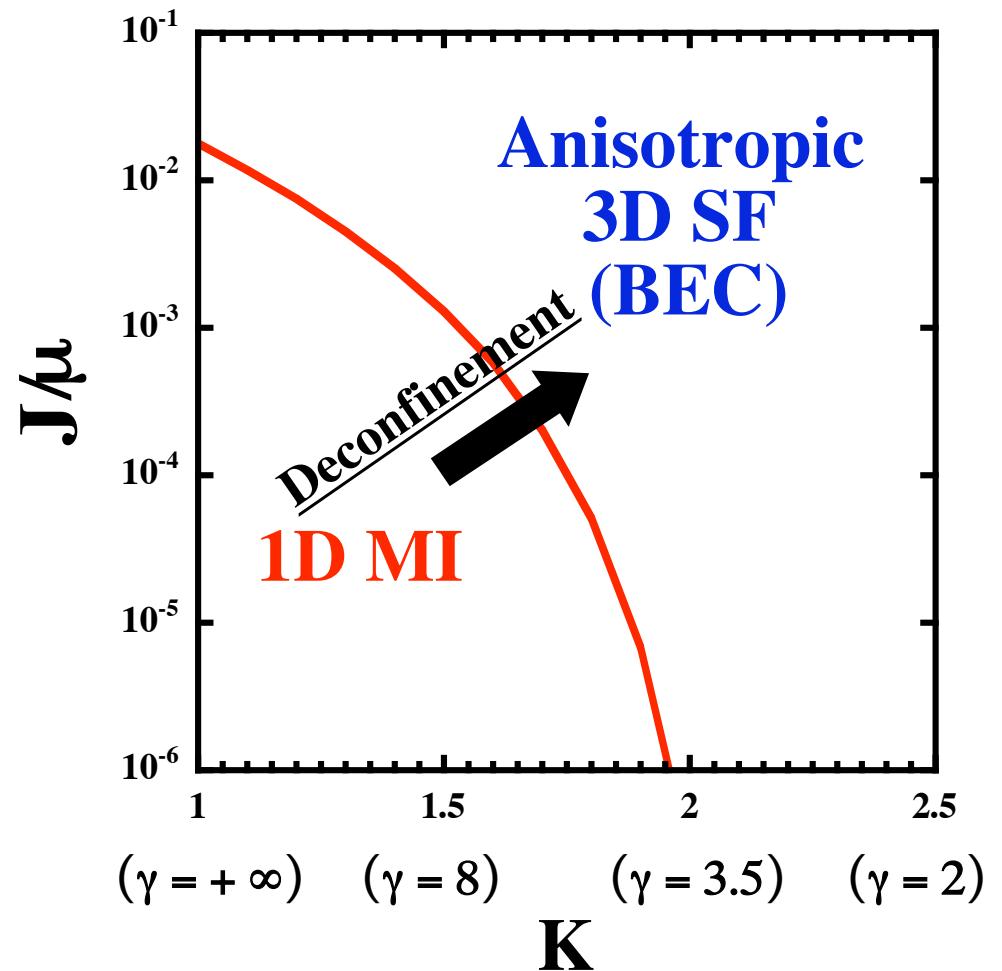
$$\frac{dg_F}{d\ell} = \frac{g_J^2}{K},$$

$$\frac{dg_J}{d\ell} = \left(2 - \frac{1}{2K}\right) g_J + \frac{g_J g_F}{2K},$$

$$\frac{dg_u}{d\ell} = (2 - K) g_u,$$

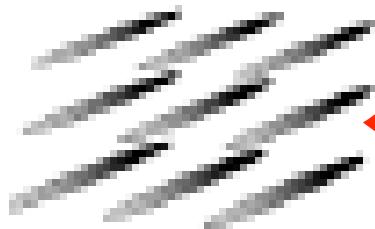
$$\frac{dK}{d\ell} = 4g_J^2 - g_u^2 K^2,$$

$$\ell \approx \ln \mu/T$$



[AFH, MAC & T Giamarchi, PRL 92 (2004)]

2D optical lattice of finite tubes: phase diagram



Array of atomic
'quantum dots'

- Quantum phase Hamiltonian ($\mathbf{g}_u = 0$) :

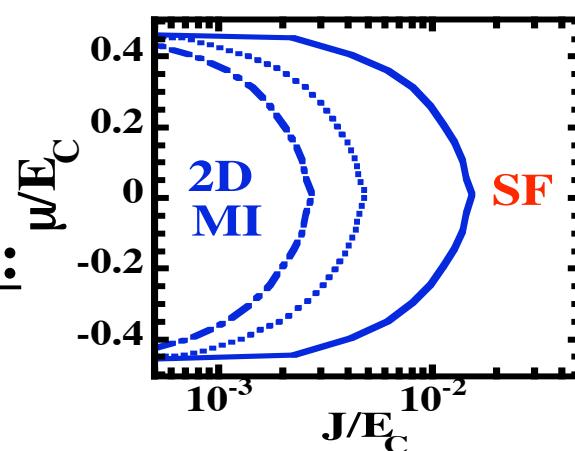
$$H_{QP} = -E_J \sum_{\langle \mathbf{R}, \mathbf{R}' \rangle} \cos(\phi_{0\mathbf{R}} - \phi_{0\mathbf{R}'}) + \frac{E_C}{2} \sum_{\mathbf{R}} (N_{\mathbf{R}} - N_0)^2 - \mu \sum_{\mathbf{R}} N_{\mathbf{R}}$$

$$E_J \approx J N_0^{1 - \frac{1}{2K}}$$

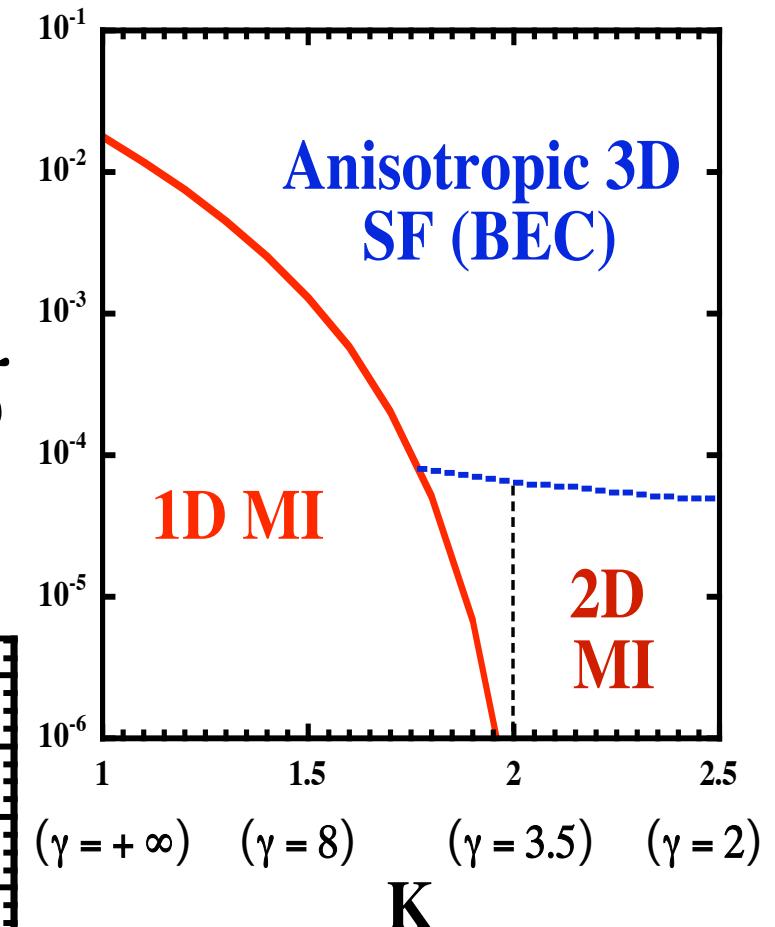
$$E_C = \frac{\hbar \pi v_s}{KL}$$

- Incommensurate fillings:

[AFH, MAC & T Giamarchi,
PRL 92 (2004)]



- 2D optical lattice: phase diagram



2D optical lattices: 3D Superfluid (BEC) phase

- Mean-field theory: condensate fraction $\psi_0^2(T = 0) \sim \rho_0 \left(\frac{J}{\mu}\right)^{1/(4K-1)}$
- Variational approach: momentum distribution at $T = 0$

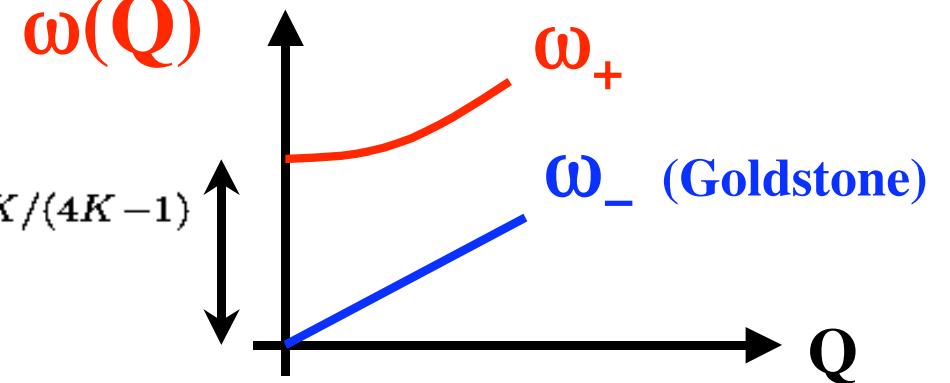
$$\frac{n(\mathbf{Q}, q)}{|w(\mathbf{Q})|^2} \simeq \psi_0^2 \delta(\mathbf{Q}) \delta(q) + \frac{\pi b^2 \psi_0^2 / 2K}{\left[q^2 + (v_{\perp} \mathbf{Q} / v_s)^2\right]^{1/2}},$$

transverse velocity:
 $v_{\perp} \sim \mu b (J/\mu)^{2K/(4K-1)} / \hbar$

- RPA : condensation temperature and excitation spectrum

$$\left(\frac{2\pi T_c}{\hbar v_s \rho_0}\right)^{2-1/2K} = f(K) \frac{4J}{\hbar v_s \rho_0}$$

[AFH, MAC & T Giamarchi, $\Delta_+ \sim \mu \left(\frac{J}{\mu}\right)^{2K/(4K-1)}$
 PRL 92 (2004)]



- 2D optical lattice: phase diagram

