

Lagrangian Statistics of Particles in Geophysical Turbulence

Herman Clercx Fluid Dynamics Laboratory Physics Department Eindhoven University of Technology The Netherlands

J.M. Burgerscentrum

e Technische Universiteit Eindhoven University of Technology

TU

Where innovation starts



- Motivation
- Inertial particles in stratified turbulence (ST)
- Summary ST
- Lagrangian statistics in rotating turbulence (RT)
- Summary RT





#### **Table-Top Geophysical Turbulence**

• Dispersion in Geophysical Turbulence





# **Table-Top Geophysical Turbulence**

- Dispersion in Geophysical Turbulence
- Dispersion in Forced Stratified Turbulence
  - Rotating Turbulence
  - Dispersion in Rotating Turbulence
    - Convection and Rotation
      - Shallow (Q2D) Flows

TU

2013





### **Table-Top Geophysical Turbulence:**

Staff "Vortex Dynamics and Turbulence" group at TU/e working on transport in (geophysical) turbulence:

Herman Clercx, GertJan van Heijst & Federico Toschi Rudie Kunnen & Bas van de Wiel

PhD students (who contributed to this talk):

Marleen van Aartrijk, Laurens van Bokhoven & Lorenzo Del Castello



#### May 9, 1998



Phytoplankton around the Galapagos Islands in May 1998 showing the rapid demise of El Niño.

• El Niño waters are blamed for choking off essential nutrients



May 9, 1998

May 18, 1998



Phytoplankton around the Galapagos Islands in May 1998 showing the rapid demise of El Niño.

- El Niño waters are blamed for choking off essential nutrients
- Upwelling during replacement of the warm El Niño waters



May 9, 1998

May 18, 1998



May 22, 1998



Phytoplankton around the Galapagos Islands in May 1998 showing the rapid demise of El Niño.

- El Niño waters are blamed for choking off essential nutrients
- Upwelling during replacement of the warm El Niño waters
- Cold nutrient-rich waters initiate phytoplankton bloom



SeaWiFS: Sea-viewing Wide Field-of-view Sensor.





Chlorophyll a Concentration mg/m3 Upwelling off Cape Town, South Africa (SeaWiFS Project, NASA/Goddard Space Flight Center).





Phytoplankton off the coast of California.

Image courtesy the SeaWiFS Project, NASA/Goddard Space Flight Center, and ORBIMAGE.

Strong currents of the coast of California are pulling cold, deep currents up from the sea floor to the surface.

These nutrient-rich waters support populations of microscopic marine plant life.

Rivers and urban areas provide additional nutrients.



- transition zones at land-sea interface
- influenced by riverine input
- nutrient-rich (land)
- strong (vertical) turbulent mixing
- particle-rich and high turbidity
- strong pelagic-benthic interaction
- tide- and wind-induced flows
- interaction with the ocean

J.E. Cloern, *Phytoplankton bloom dynamics in coastal ecosystems*, Reviews of Geophysics, 34, 127-168 (1996). KITP GEOFLOWS13 October 2013



Phytoplankton blooms in Shallow Coastal Ecosystems (Chesapeake Bay, USA).



What is the relevant range of scales?

1-100 μm: smallest organisms;
0.1-10 cm: Kolmogorov scale;
0.01-10 m: stratification;
10-1000 km: rotation and stratification.

Processes that regulate the total biomass, species composition, and spatial distribution of phytoplankton include:

- turbulent vertical mixing;
- (turbulent) horizontal advection;
- in situ processes of population change.





# Algal blooms

#### Aerosol dispersion



How is particle dispersion affected by the particle's inertial properties? Does preferential concentration persist for small particle-to-fluid density ratios? How is preferential concentration affected by stratification?



# **Direct Numerical Simulation**

Two steps:

- 1. Eulerian approach for velocity field
- 2. Lagrangian approach for particle tracking



# **Direct Numerical Simulation**



- Boussinesq approx.
- Periodic domain
- 128<sup>3</sup> (256<sup>3</sup>)
- Forced DNS
- Parallel



# Stable stratification



Buoyancy frequency N:





# **Density profiles**





N~0.1 (s<sup>-1</sup>)

 $N \sim 1$  (s<sup>-1</sup>) TU

e Technische Universiteit Eindhoven University of Technology

October 2.

2013

# Isovorticity





N~0.1 (s<sup>-1</sup>)

N~1 (s<sup>-1</sup>) TU/e Technische Universiteit Eindhoven University of Technology October 2, 2013

# Fluid particle tracking

$$\frac{d \vec{x}_p}{dt} = \vec{u}_p$$

$$\vec{u}_p = \vec{u}_f(\vec{x}_p)$$

# Cubic spline interpolation











# Horizontal dispersion



M. van Aartrijk, H.J.H. Clercx & K.B. Winters, *Phys. Fluids* **20**, 025104 (2008).





# Vertical dispersion



M. van Aartrijk, H.J.H. Clercx & K.B. Winters, Phys. Fluids 20, 025104 (2008).





# Inertia effect on dispersion

$$\overline{x^2} = 2\overline{u'_p}^2 \int (t - \tau) R(\tau) d\tau \text{ Taylor (1921)}$$

- Increasing inertia  $\rightarrow u'_p^2 \downarrow \rightarrow$ decreasing dispersion
- Increasing inertia  $\rightarrow$  memory,  $R(\tau) \uparrow \rightarrow$  increasing dispersion
- Dispersion optimum around  $\tau_p = \tau_K$  (iso)

2013

# Inertial particle tracking

$$m_p \frac{d\vec{u}_p}{dt} = 6\pi a\mu \left(\vec{u} - \vec{u_p} + \frac{1}{6}a^2\nabla^2\vec{u}\right) + m_f \frac{D\vec{u}}{Dt}$$
$$- (m_p - m_f) g\vec{\hat{z}} + \frac{1}{2}m_f \left(\frac{D\vec{u}}{Dt} - \frac{d\vec{u}_p}{dt} + \frac{1}{10}a^2\frac{d}{dt}\nabla^2\vec{u}\right)$$
$$ct = d\vec{z}/dz = d\vec{z}/dz + \frac{1}{2}a^2d\nabla^2\vec{z}/dz$$

$$+6\pi a^{2}\mu \int_{0}^{t} d\tau \frac{d\vec{u}/d\tau - d\vec{u}_{p}/d\tau + \frac{1}{6}a^{2}d\nabla^{2}\vec{u}/d\tau}{\left[\pi\nu(t-\tau)\right]^{1/2}}$$

M.R. Maxey and J.J. Riley, *Phys. Fluids* **26**, 883 (1983).

University of Technology

October 2.

2013



# Forces on the particles (N~0.3)



2013

# Inertial particle tracking

Heavy inertial particles:

$$\frac{\rho_p}{\rho_f} \gg 1$$

October 2.

2013

$$\begin{aligned} \frac{du_p}{dt} &= \frac{1}{\tau_p} \left( u - u_p \right) \\ \tau_p &= \frac{\left( \rho_p / \rho_f \right) d_p^2}{18\nu} \end{aligned} \qquad St = \frac{\tau_p}{\tau_k} \\ \text{TU(e) Extracted Unversited} \end{aligned}$$

# Horizontal dispersion (N~1)



#### Vertical dispersion (N~1) 10<sup>3</sup> St=0.30 St=0.55 St=0.80 10<sup>2</sup> St=1.38 St=1.81 St=6.75

 $\overline{z(0))^2}N^2/\overline{w'^2}$  $10^{1}$ fluid 10<sup>0</sup>  $t^2$ 10<sup>-1</sup> 8 M. van Aartrijk & H.J.H. Clercx, Phys. Fluids 22, 013301 (2010). 10<sup>-2</sup> 10<sup>-1</sup>  $10^{0}$  $10^{1}$ 

### Vertical dispersion (N~0.3)



- P = no Basset
- Q = no Faxèn corrections
- R = only Stokes drag, pressure gradient, added mass,

St = 1.4 $\rho_p / \rho_f = 25$ 

Technische Universiteit Eindhoven University of Technology

October 2.

2013

# **Preferential concentration**

![](_page_31_Figure_2.jpeg)

![](_page_31_Picture_3.jpeg)

# High strain, low vorticity

![](_page_31_Picture_5.jpeg)

# Preferential concentration isotropic turbulence

![](_page_32_Figure_2.jpeg)

2013

# Preferential concentration stratified turbulence

![](_page_33_Picture_2.jpeg)

![](_page_33_Picture_3.jpeg)

# Horizontal

M. van Aartrijk & H.J.H. Clercx, *Phys. Rev. Lett.* **100**, 254501 (2008).

 $N \sim 1 (s^{-1})$ 

Technische Universiteit Eindhoven University of Technology

KITP GEOFLOWS13 October 2013

October 2, 2013

## Pref. conc. isotropic turbulence

![](_page_34_Figure_2.jpeg)

Minimum distance to closest neighbor

 $\tau_d = (d^2/\mathcal{E})^{1/3}$ 

# Preferential concentration stratified turbulence

![](_page_35_Figure_2.jpeg)

# Intermediate Summary

- Stratification enhances horizontal dispersion and reduces vertical dispersion
- Inertia has negligible influence on horizontal dispersion in stratified turbulence
- With increasing inertia (  $St\uparrow$  ), the long-time vertical dispersion in stratified turbulence increases
- Stratification enhances mixing; only within the horizontal plane
- Basset force needs to be taken into account for vertical light
   particle dispersion in stratified turbulence

![](_page_37_Picture_1.jpeg)

![](_page_37_Picture_2.jpeg)

![](_page_37_Picture_3.jpeg)

**Engineering applications** 

Astrophysical and geophysical applications

- How does rotation affect velocity and acceleration PDFs?
- How does rotation affect Lagrangian autocorrelations?

- How do they compare with Eulerian autocorrelations? KITP GEOFLOWS13 October 2013

![](_page_37_Picture_9.jpeg)

Navier-Stokes equations in rotating frame of reference:  $\rho(\partial \mathbf{u}/\partial \mathbf{t} + \mathbf{u} \cdot \nabla \mathbf{u}) + 2\rho \Omega \times \mathbf{u} = -\nabla \rho + \rho \nabla^2 \mathbf{u}$  $\nabla \cdot \mathbf{u} = \mathbf{0}$ inertial forces u·∇u IJ Rossby number: Coriolis force  $|2\Omega \times \mathbf{u}|$  $2\Omega L$ Ro Technische Universiteit University of Technology October 2. KITP GEOFLOWS13 October 2013 2013

Navier-Stokes equations in rotating frame of reference:  $\rho(\partial \mathbf{u}/\partial \mathbf{t} + \mathbf{u} \cdot \nabla \mathbf{u}) + 2\rho \Omega \times \mathbf{u} = -\nabla \rho + \rho \nabla^2 \mathbf{u}$  $\nabla \cdot \mathbf{u} = 0$ viscous forces Ekman number: Coriolis force  $2\Omega L^2$  $|2\Omega \times \mathbf{u}|$ Ek University of Technology October 2. KITP GEOFLOWS13 October 2013 2013

Navier-Stokes equations in rotating frame of reference:

$$\rho(\partial \mathbf{u}/\partial t + \mathbf{u} \cdot \nabla \mathbf{u}) + 2\rho \Omega \times \mathbf{u} = -\nabla p + \nu \rho \nabla^2 \mathbf{u}$$

 $\nabla \cdot \mathbf{u} = 0$ 

steady, Ro << 1, viscous effects negligible (Ek << 1):  $2\Omega \times \mathbf{u} = -(1/\rho)\nabla p$ 

Taylor-Proudman theorem: under strong rotation

 $\frac{\partial \mathbf{v}}{\partial z} = 0$ 

 $\rightarrow$  columnar flow structuring

![](_page_40_Picture_8.jpeg)

# Taylor–Proudman theorem ↔ no vertical variation of velocity under geostrophic conditions

![](_page_41_Figure_2.jpeg)

"Taylor column" above object dragged through a rotating fluid

![](_page_41_Picture_4.jpeg)

#### What is the effect of rotation on turbulence?

![](_page_42_Figure_2.jpeg)

Experiment by Hopfinger, Browand and Gagne, JFM 125 (1982).

![](_page_42_Picture_4.jpeg)

![](_page_43_Figure_1.jpeg)

#### **Set-up for Rotating Turbulence Studies**

![](_page_43_Picture_3.jpeg)

![](_page_44_Picture_1.jpeg)

![](_page_44_Figure_2.jpeg)

![](_page_44_Picture_3.jpeg)

![](_page_45_Picture_1.jpeg)

![](_page_45_Picture_2.jpeg)

![](_page_46_Picture_1.jpeg)

![](_page_46_Picture_2.jpeg)

KITP GEOFLOWS13 October 2013

October 2, 2013

![](_page_47_Figure_1.jpeg)

![](_page_48_Figure_1.jpeg)

 $Re_{\lambda}$ ~150-250 , H=5 cm

Ω=0: stationary, reproducible, and  $(u')^2 \sim (v')^2 \sim (w')^2$ .

Characterization of rotating turbulence at several heights in the rotating fluid;  $\Omega$ =0.1, 0.2, 0.5, 1, 2, 5, 10 rad/s.

![](_page_48_Picture_5.jpeg)

![](_page_49_Figure_1.jpeg)

![](_page_50_Picture_1.jpeg)

![](_page_50_Picture_2.jpeg)

Ехр		1	2	3	4	5	6
Ω	rad/sec	0	0.2	0.5	1.0	2.0	5.0
U <sub>rms</sub>	mm/sec	12.4	11.0	13.5	16.3	22.9	16.9
Ro			0.86	0.30	0.11	0.07	0.04
T <sub>eddy</sub>	sec	5.66	6.37	5.20	4.31	3.05	4.15
ΔT	sec	160	160	160	160	320	320
$\Delta T/T_{eddy}$		28.3	25.1	30.8	37.1	104.9	77.1

 $0.6 < \eta < 0.8 \mbox{ mm}$  ,  $0.2 < \tau_\eta < 0.5 \mbox{ sec}$ 

 $Ro = U_{rms}/(2\Omega L)$ 

October 2.

2013

TU

 $50 < \text{Re}_{\lambda} < 150$ 

 $\Delta t_{PTV} = 16.7 \ \mu m$  for  $\Omega$  up to 1 rad/sec and  $\Delta t_{PTV} = 33.3 \ \mu m$  for  $\Omega = 2$  and 5 rad/sec

![](_page_52_Figure_1.jpeg)

- Horizontal homogeneous and isotropic
- Vertical decay of energy (but reduced for rotating turbulence)

For more details, see Van Bokhoven et al., PoF 21, 096601 (2009) TU/e Technische Universiteit Eindhoven University of Technology Notober 2013

#### **Comparison Eulerian and Lagrangian spatial correlations**

![](_page_53_Figure_2.jpeg)

#### **Horizontal and vertical velocity PDFs**

![](_page_54_Figure_2.jpeg)

#### **Horizontal and vertical acceleration PDFs**

![](_page_55_Figure_2.jpeg)

Agreement with literature for small  $\boldsymbol{\Omega}$ 

Note: PDFs converged with  $32\Delta t_{PTV} \rightarrow \Delta t_{PTV}$ 

echnische Universiteit

#### Lagrangian longitudinal auto-correlation coefficients of Cartesian <u>velocity</u> components

![](_page_56_Figure_2.jpeg)

2013

# Lagrangian longitudinal auto-correlation coefficients of<br/>Cartesian acceleration components $\Omega=0$

![](_page_57_Figure_2.jpeg)

#### Lagrangian longitudinal auto-correlation coefficients of Cartesian <u>acceleration</u> components all Ω>0

![](_page_58_Figure_2.jpeg)

# Decomposition of the Lagrangian acceleration vector

![](_page_59_Figure_2.jpeg)

#### Lagrangian longitudinal auto-correlation coefficients of $a_l$ , $a_{tv}$ , $a_{th}$ acceleration components $\Omega=0$

![](_page_60_Figure_2.jpeg)

#### Lagrangian longitudinal auto-correlation coefficients of $a_l$ , $a_{tv}$ , $a_{th}$ acceleration components all $\Omega > 0$

![](_page_61_Figure_2.jpeg)

# Lagrangian statistics in rotating turbulence

- Horizontal and vertically velocity PDFs remain Gaussian except for highest rotation rate (two-dimensionalization). Horizontal acceleration PDFs show enhancement of tails with increasing rotation rate (more extreme events); the vertical acceleration PDFs shows just the opposite effect (twodimensionalization).
- Amplification of the memory time-scales of the fluid particle velocity and acceleration with rotation.
   Increase of horizontal and vertical length scales and suppression of vertical motion.
- Strong amplification of the transversal horizontal acceleration (*a<sub>th</sub>*) correlation, and mild increase of the longitudinal acceleration (*a<sub>l</sub>*) correlation.
   Direct and indirect effects, respectively, of the Coriolis acceleration.

![](_page_62_Picture_4.jpeg)