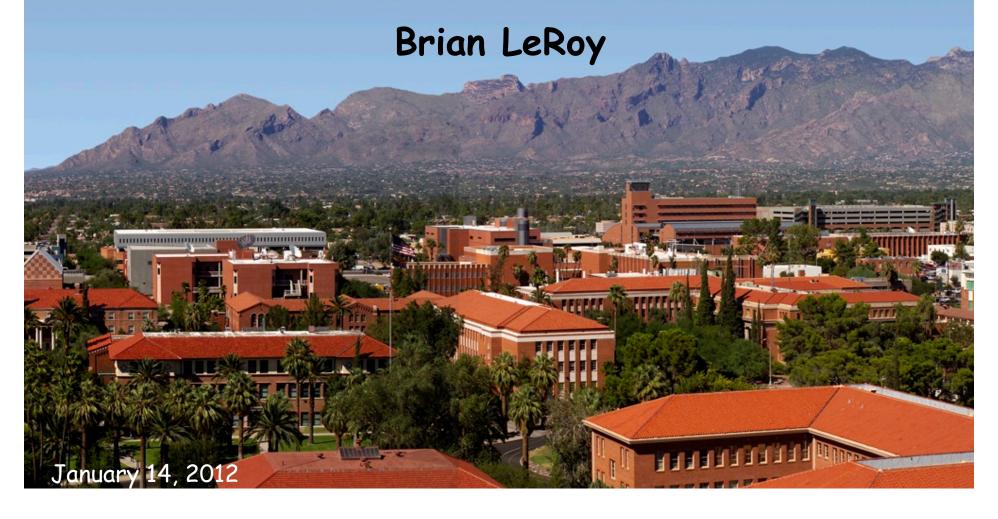


# Scanning tunneling spectroscopy of graphene on BN



#### Outline

Scanning tunneling microscopy

Topography Moiré patterns

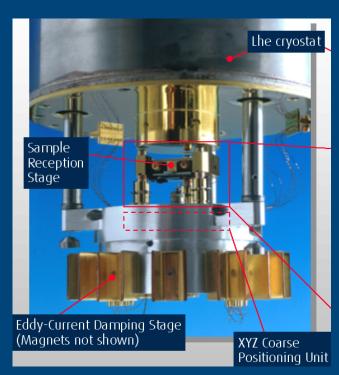
Spectroscopy
Charged impurities
Scattering from edges
Periodic potentials

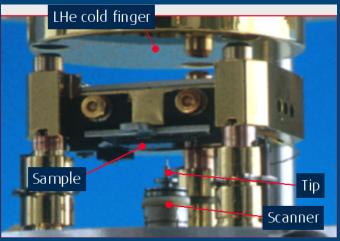
Conclusions

## Scanning Tunneling Microscope



Spatial Resolution 1  $^{\text{A}}$  Pressure < 3  $\times$  10<sup>-11</sup> mbar Temperature 4.5 K





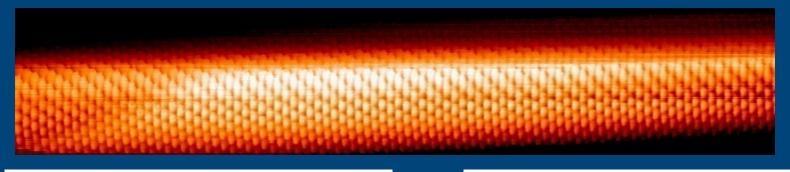
### Scanning Tunneling Microscopy

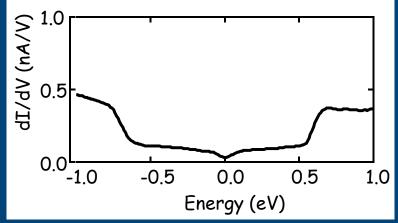
Topography

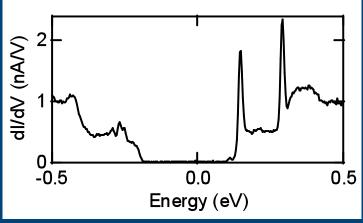
$$I \propto e^{-Z/Z_0} \int_0^{eV} dE \, LDOS(E, r)$$

Spectroscopy

$$\partial I/\partial V \propto LDOS(eV,r)$$







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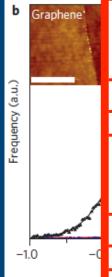
#### Why use hBN?

h-BN: hexagonal lattice, with ~same lattice constant as graphene (1.8% mismatch) atomically smooth surface no dangling bonds / charge traps insulator, with bandgap ~6eV

Boron nitride substrates for high-quality graphene electronics

C. R. Dean<sup>1,2</sup>\*, A. F. Young<sup>3</sup>, I. Meric<sup>1</sup>, C. Lee<sup>4,5</sup>, L. Wang<sup>2</sup>, S. Sorgenfrei<sup>1</sup>, K. Watanabe<sup>6</sup>, T. Taniguchi<sup>6</sup>, P. Kim<sup>3</sup>, K. L. Shepard<sup>1</sup> and J. Hone<sup>2</sup>\*

Dean et al, Nature Nano **5**, 722 (2010)

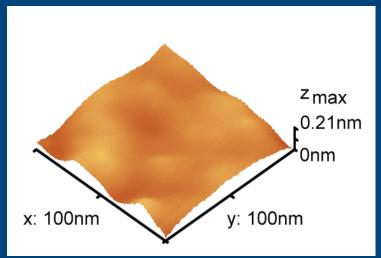


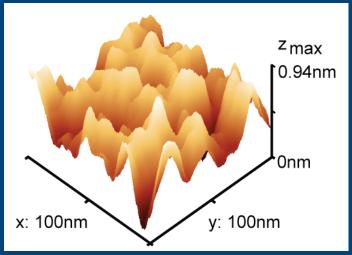
It has been proposed that a bandgap would be induced in graphene aligned to an h-BN substrate<sup>20</sup>. In our experiment the graphene had a random crystallographic orientation to the substrate, and thus we did not expect the necessary sublattice symmetry breaking to occur. Indeed, the temperature dependence of  $\sigma_{\min}$  observed here does not follow the simply activated behaviour that would be indicative of an energy gap. Although we cannot rule out the possibility of locally gapped regions resulting from symmetry breaking over finite length scales, we see no evidence from transport measurements that an appreciable gap is present in this randomly stacked graphene on h-BN.

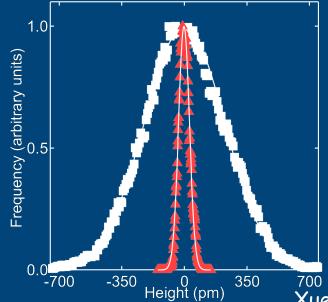
AFM topography, h(r)

Weakly shifted away from  $V_g=0$ 

## Topography Measurements hBN SiO<sub>2</sub>



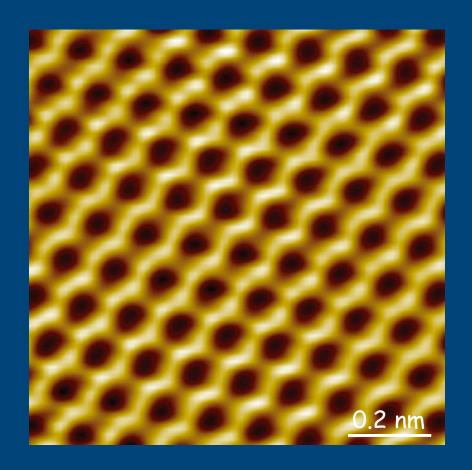




Standard deviations are:  $SiO_2$ : 224.5±0.9 pm BN:  $30.2 \pm 0.2$  pm

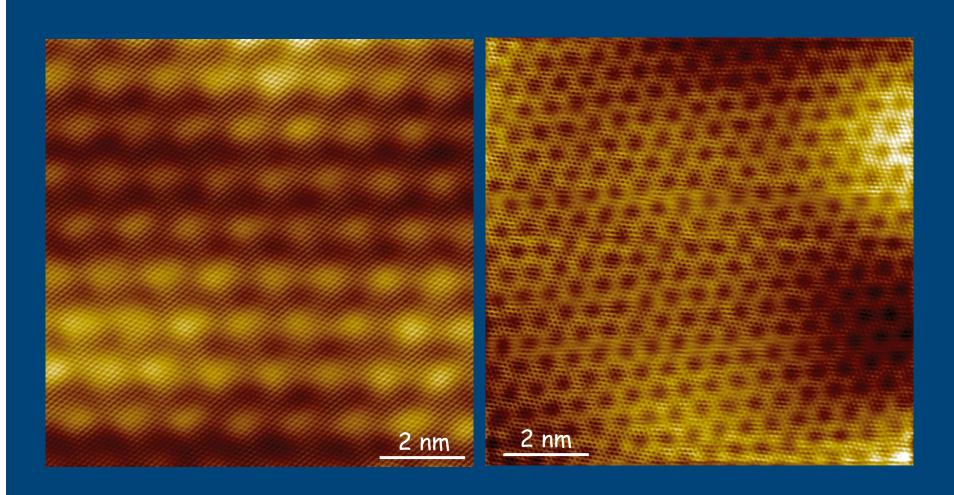
Xue et al., Nat. Mater. 10, 282 (2011)

#### Atomic Resolution



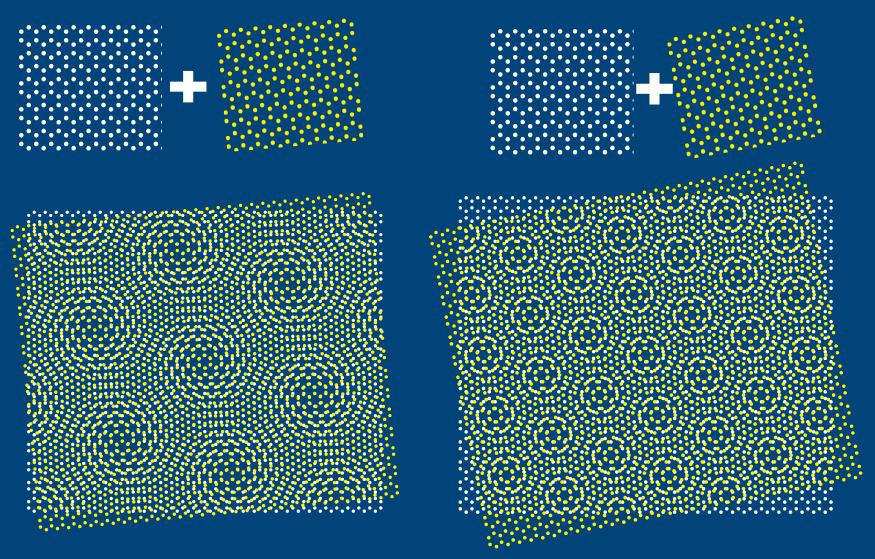
Atomic resolution image shows hexagonal lattice

### Topography Measurements



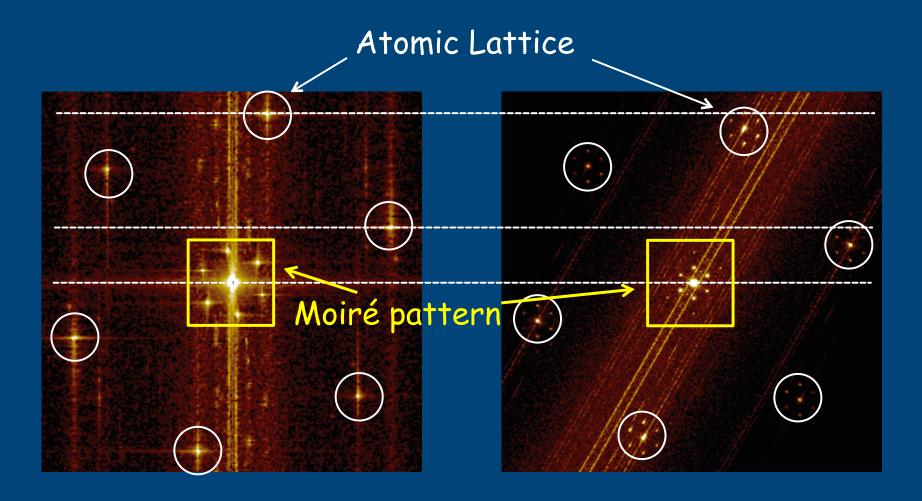
Large scale images show moiré pattern Different areas of the same graphene flake

#### Moiré Patterns



Moiré patterns arise from rotation between graphene and hBN

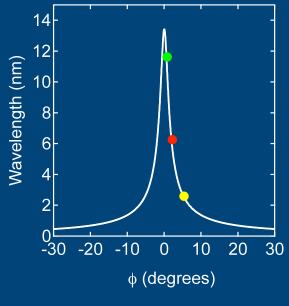
#### Fourier Transforms

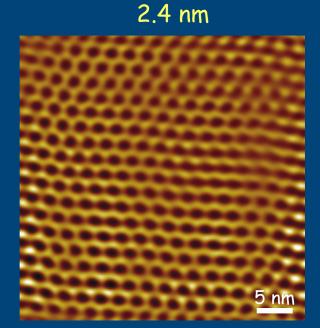


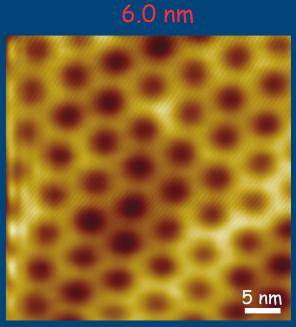
Different size moiré patterns Graphene lattice rotated between two images Moiré Wavelength

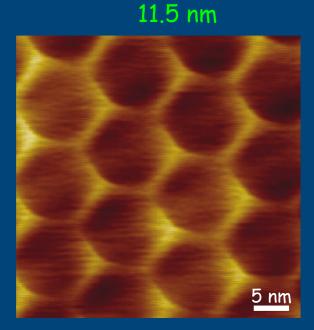
$$\lambda = \frac{(1+\delta)\alpha}{\sqrt{2(1+\delta)(1-\cos\phi)+\delta^2}}$$

- a: graphene lattice constant
- $\delta$ : mismatch of h-BN
- φ: Angle between lattices









#### Outline

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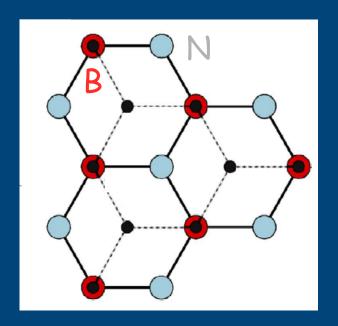
## Consequence of moiré pattern: No energy gap

Previous theoretical prediction<sup>1</sup>

No moiré pattern.

Broken symmetry of A-B carbon atoms.

50 meV energy gap.

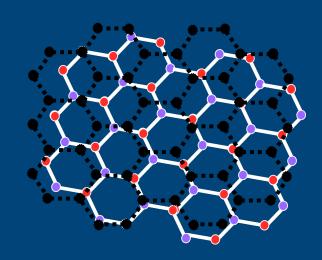


**Experimental observation:** 

Has moiré pattern.

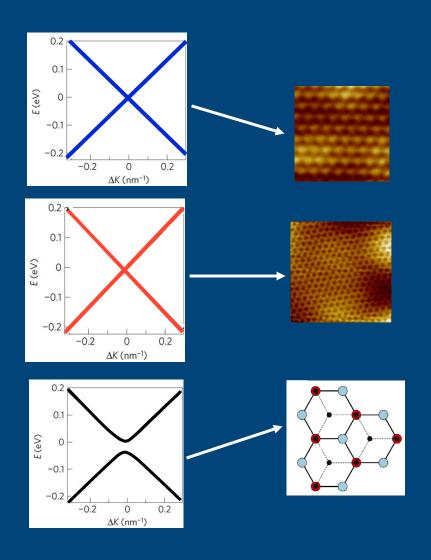
Symmetry of A-B carbon atoms is restored.

No energy gap.

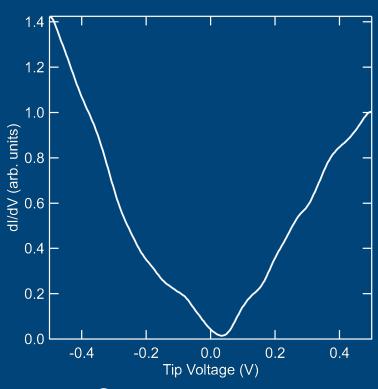


1: G. Giovannetti et al., PRB (2007)

#### Spectroscopy Measurements



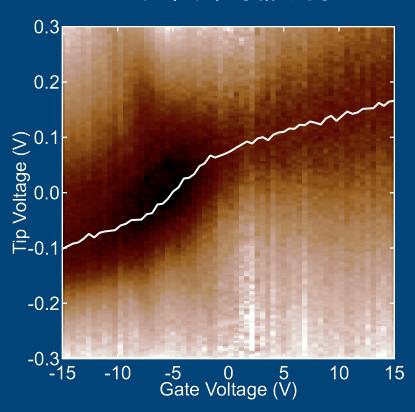
Calculated dispersion relations for three different configurations.

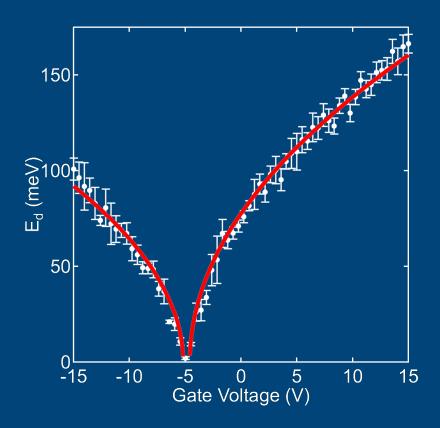


Spectroscopy measured by STM No energy gap.

#### Gate Dependence

#### dI/dV curves

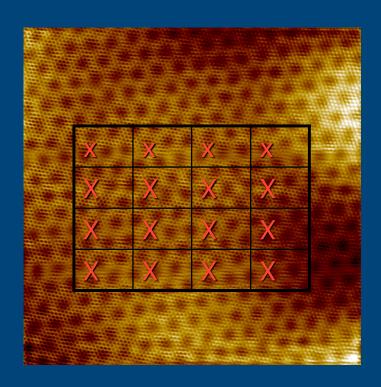


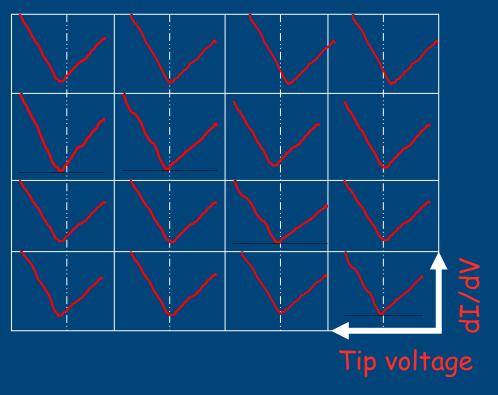


Dirac point follows expected energy for linear dispersion

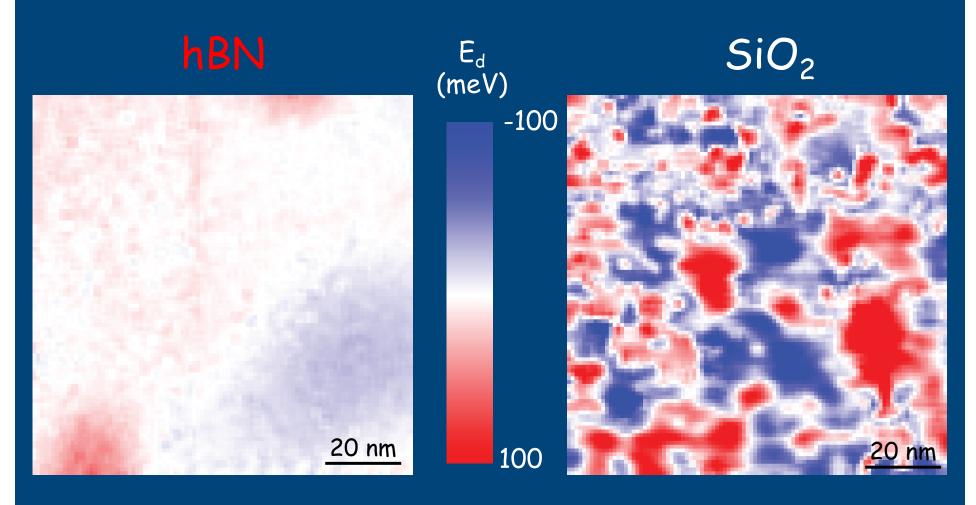
#### Spectroscopy Map

- Spectroscopy performed on a 1 nm x 1 nm grid
- Measure tip voltage corresponding to Dirac point



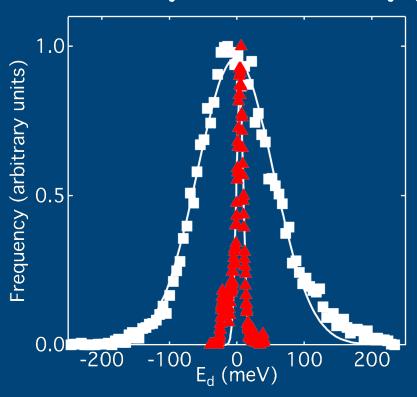


#### Spectroscopy Map

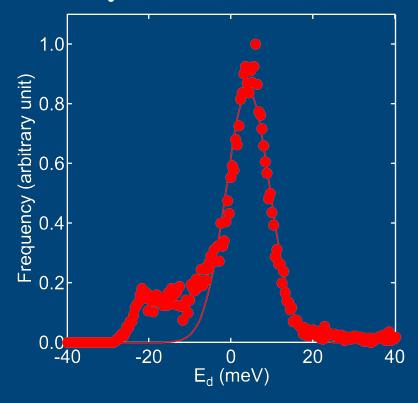


Potential fluctuations much smaller on hBN Spatial extent of puddles is larger on hBN

#### Spectroscopy Comparison



 $SiO_2$ : 55.6 ± 0.7 meV



BN: 5.4±0.1 meV

Distribution on hBN is 10 times narrower Extra bump in distribution observed in most samples

Xue et al., Nat. Mater. 10, 282 (2011)

#### Outline

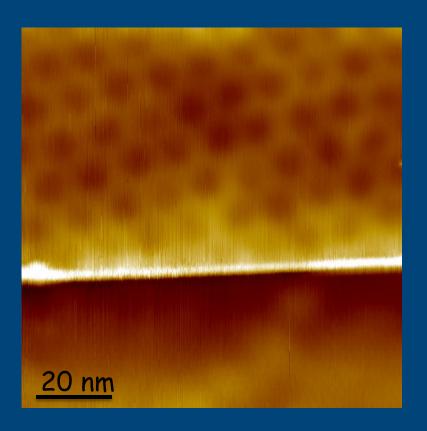
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## Topography of step edge

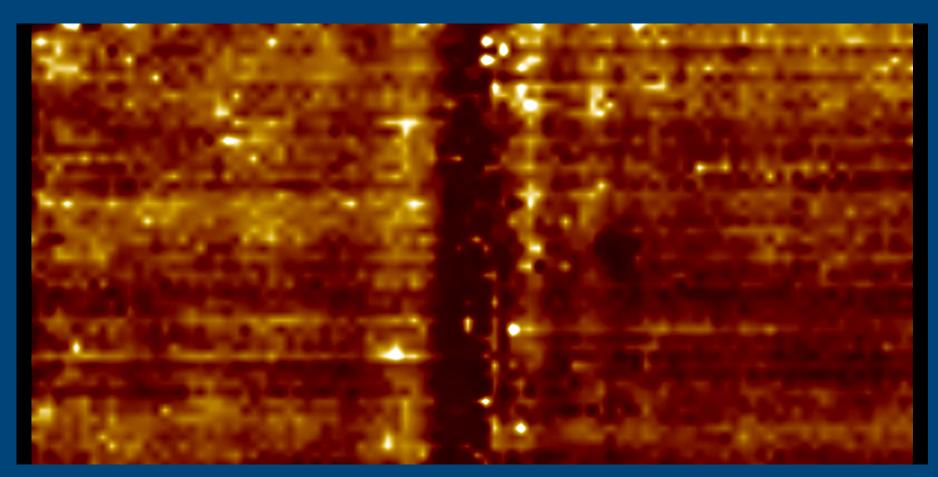


Step is about 0.6 nm high

Graphene lattice has the same orientation above and below the step

## Density of States Versus Energy

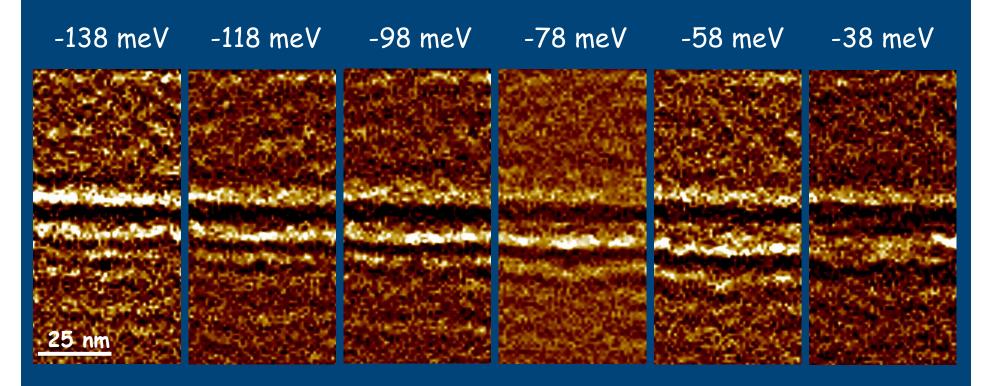
Step edge running vertically through images Series of gap voltages (energies)



Xue et al., PRL 108, 016801 (2012)

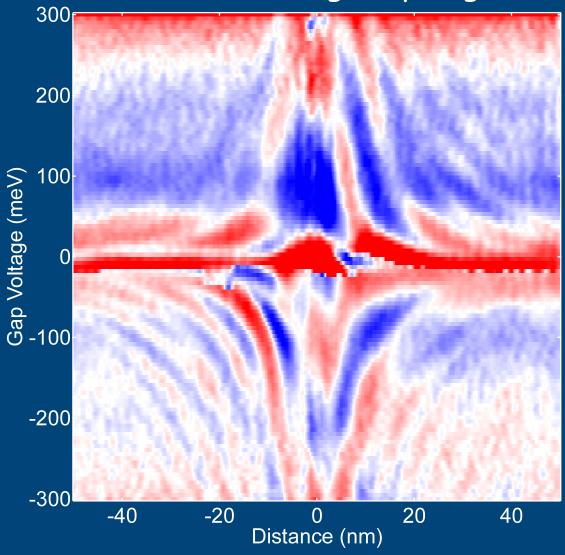
#### Density of States Images

Step edge running horizontally through images
Series of gap voltages (energies)



Wavelength increases with decreasing energy

## Density of States Images All data in x-direction (along step edge) averaged



Color scale is change in density of states

#### Distance Dependence

Barrier



Constant Energy
Contour



Assume barrier along x-axis

$$k_x = k_x' \qquad k_y = -k_y'$$

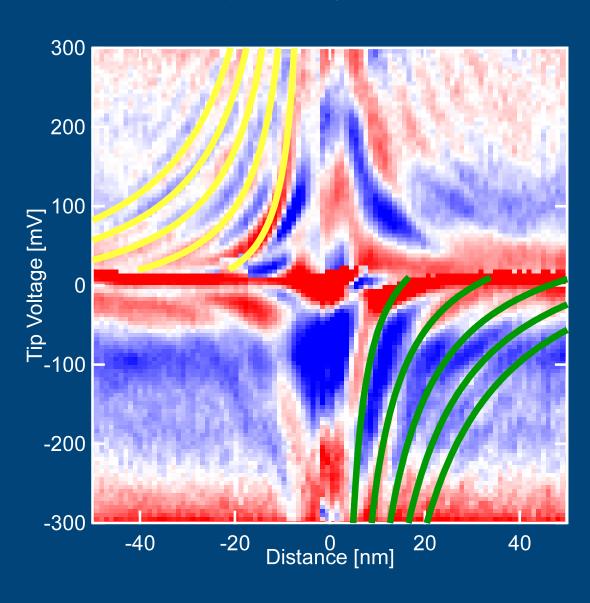
$$\rho(E, y) \propto \oint_{CEC} \left| \psi(k_x, k_y) + r \, \psi(k_x, -k_y) \right|^2 dk$$

$$\delta \rho(E, y) \propto \oint_{CEC} \cos(2k_{\mu}y) \sin\theta_{k} dk \qquad \tan\theta_{k} = \frac{k_{x}}{k_{y}}$$

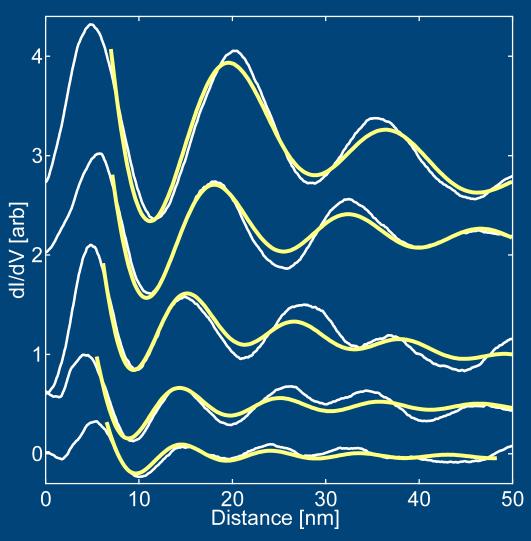
$$\delta \rho(E, y) \propto \frac{\cos(2ky - 3\pi/4)}{(ky)^{3/2}}$$

Faster decay than normal metal

## Energy Dependence



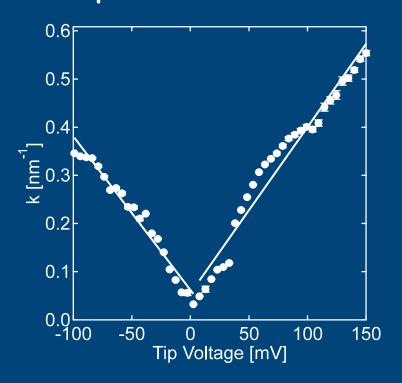
### Distance Dependence



$$\delta\rho(E,y) \propto \frac{\cos(2ky-3\pi/4)}{(ky)^{3/2}}$$

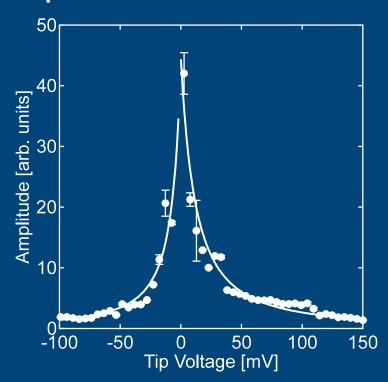
## Analysis

#### Dispersion Relation



 $v_F = 0.50 \pm 0.05 \times 10^6 \text{ m/s}$ 

#### Amplitude of oscillations



Xue et al., PRL 108, 016801 (2012)

#### Outline

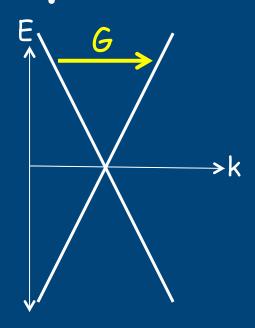
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#### Graphene in Periodic Potential



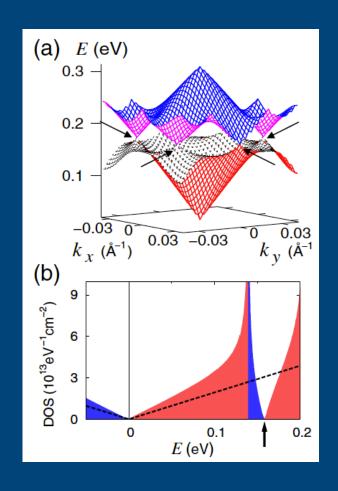
Schrödinger

Band gap opens

Dirac

New superlattice Dirac point

$$E = \hbar v_F G / 2$$

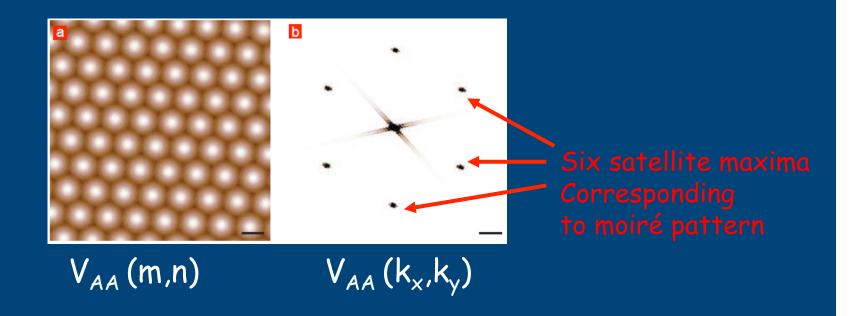


Park et al., *Nat. Phys.* **4**, 213 (2008); *PRL* **101**, 126804 (2008); Barbier et al., *PRB* **77**, 115446 (2008) Brey and Fertig, *PRL* **103**, 046809 (2009); Barbier et al., *PRB* **80**, 205415 (2009); Sun et al., *PRL* **105**, 156801 (2010); Burset et al., *PRB* **83**, 195434 (2011); Ortix et al., arXiv:1111:0399 (2011)

#### Graphene on hBN

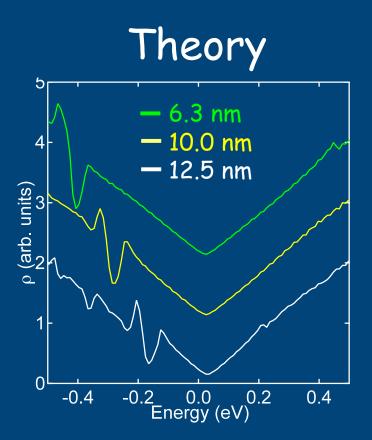
- $\triangleright$  Calculate interlayer hopping from  $\mu$ =A,B carbon sites to  $\nu$ =Boron, Nitrogen sites
- >Keep nearest neighbors and next n.n. -> four different hoppings

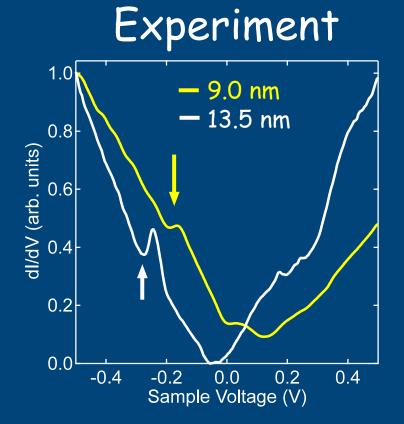
$$V_{\mu\nu}(m,n) = \gamma_{\perp} \exp[-|\mathbf{r}_{1\mu}(m,n) - \mathbf{r}_{2\nu}(m',n')|/\xi]$$



Graphene experiences periodic potential given by moiré pattern

#### Spectroscopy

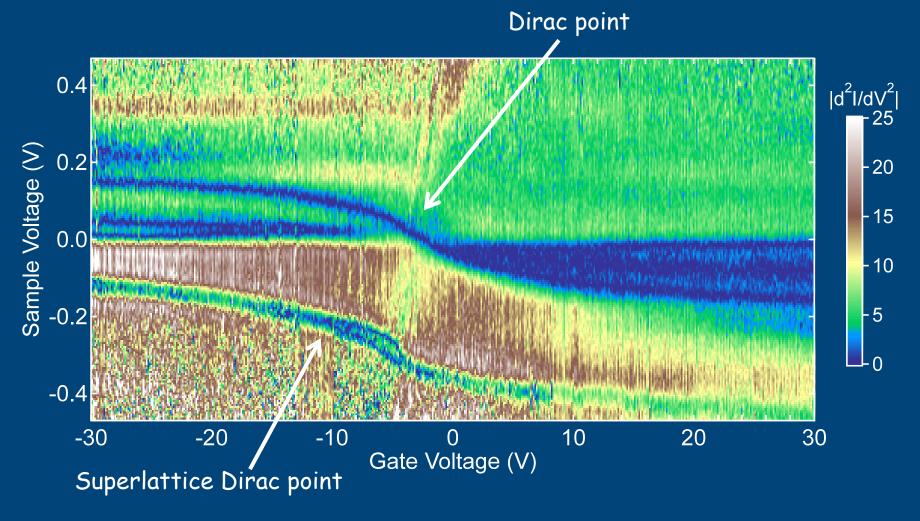




Dips in density of states due to superlattice potential

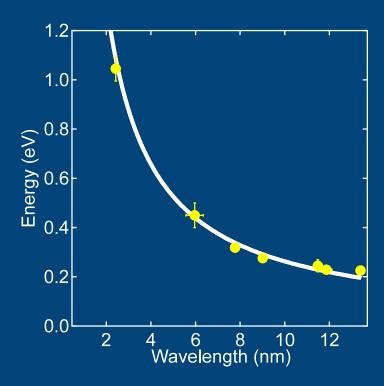
#### Gate Dependence

9.0 nm moiré



Dips move together with Fermi energy

#### Superlattice Dirac Point

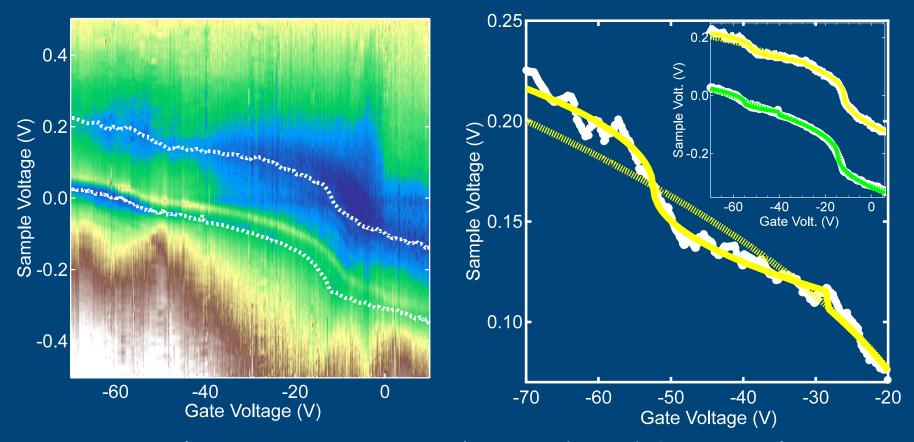


Energy of superlattice Dirac points determined by wavelength of potential

$$E = \hbar v_F G / 2 = 2\pi \hbar v_F / \sqrt{3}\lambda$$
$$v_F = 1.1 \times 10^6 \text{ m/s}$$

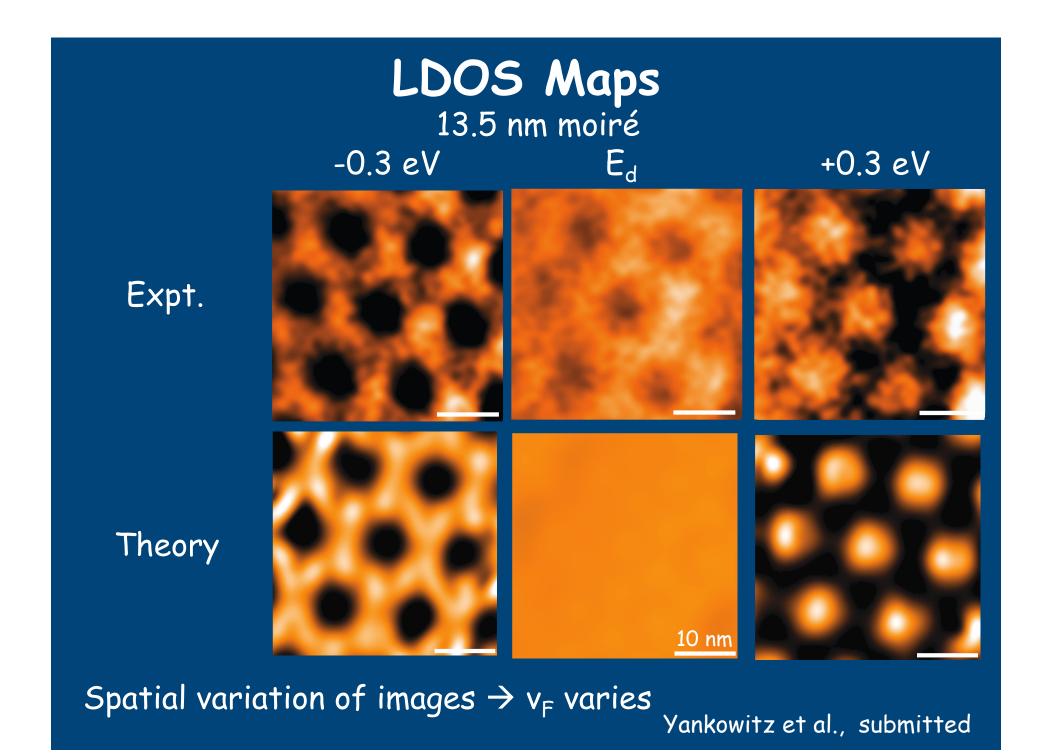
### Superlattice Dirac Point

13.5 nm moiré



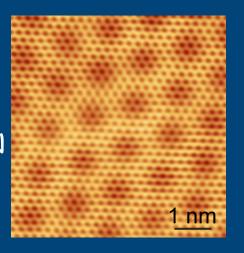
Superlattice Dirac point has reduced Fermi velocity

$$v_F^* / v_F = 0.5 - 0.7$$

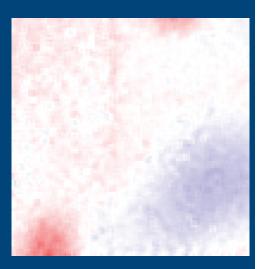


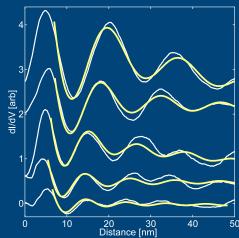
#### Conclusions

Moiré pattern observed for graphene on hBN

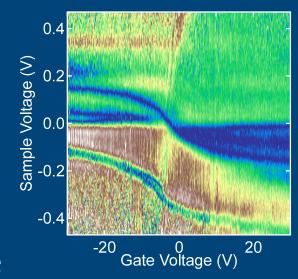


Electron and hole puddles reduced on hBN





LDOS oscillations near step edges



Superlattice Dirac point due to periodic potential

www.physics.arizona.edu/~leroy

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#### LeRoy group

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Sam Silva (Undergrad)
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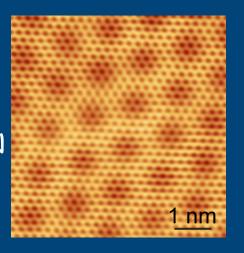




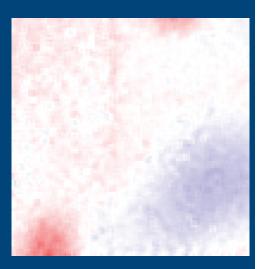


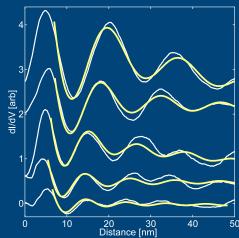
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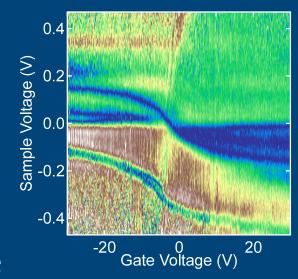


Electron and hole puddles reduced on hBN





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