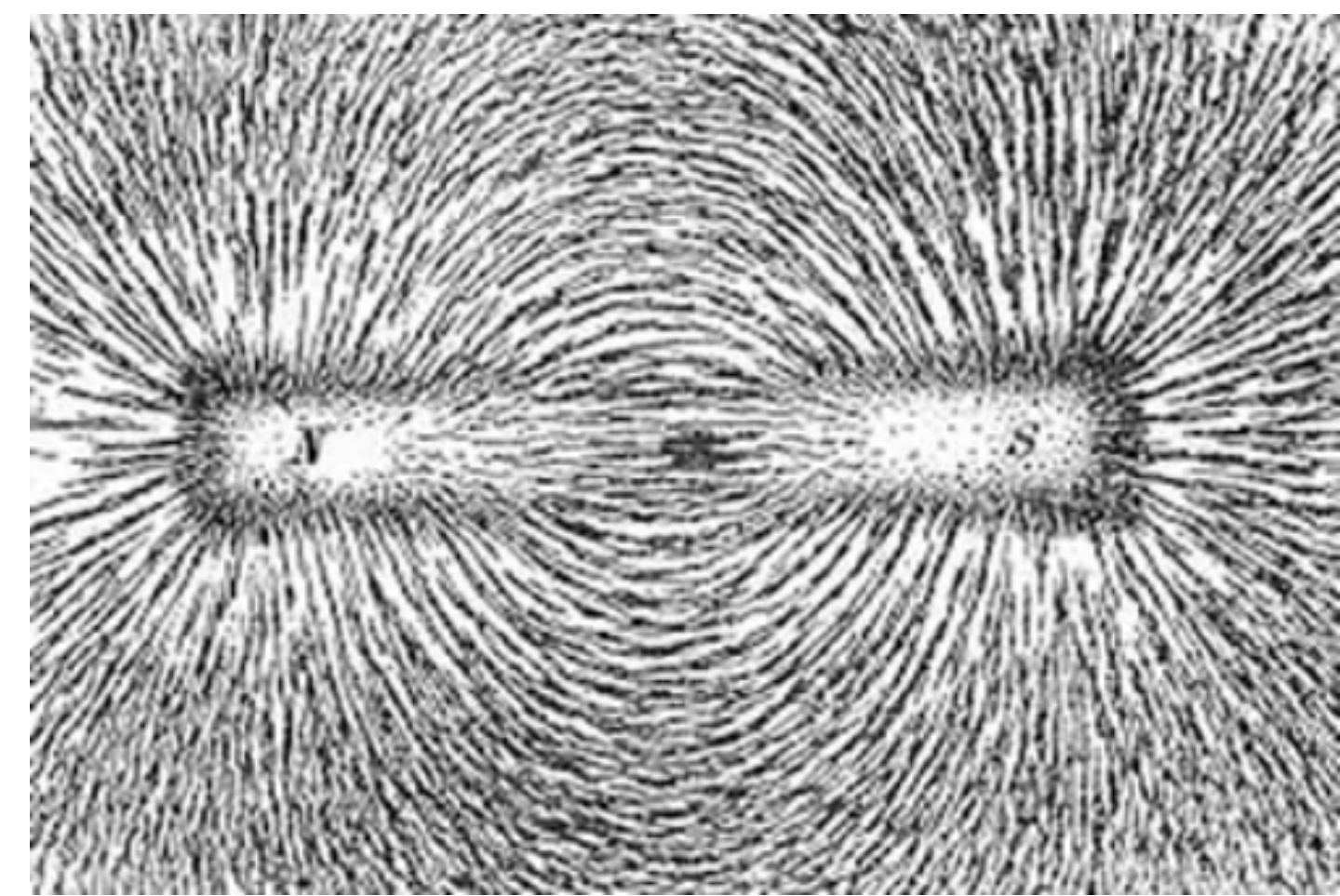




Non-thermal Processes in the CGM

Peng Oh (UC Santa Barbara)



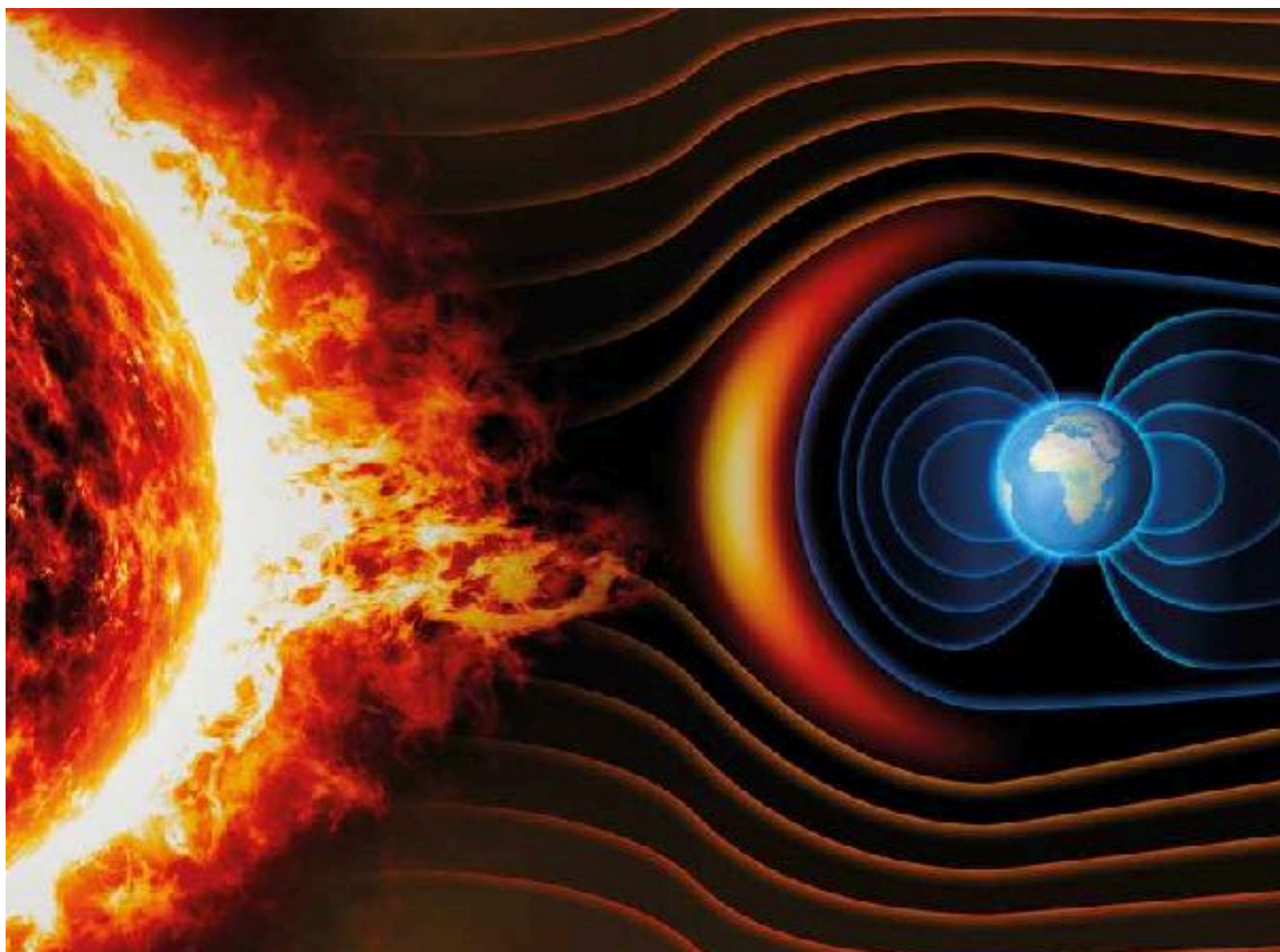
Non-thermal Processes: the Last Refuge of Scoundrels?



Cosmic Rays



Turbulence



Magnetic Fields



Why care about life non-thermal?

ASK ME



IF I CARE

© Paws, Inc.

1. Democracy

A CALL TO
#DEFENDDEMOCRACY



In ISM, $P_{gas} \sim P_B \sim P_{turb} \sim P_{CR}$ WHY???

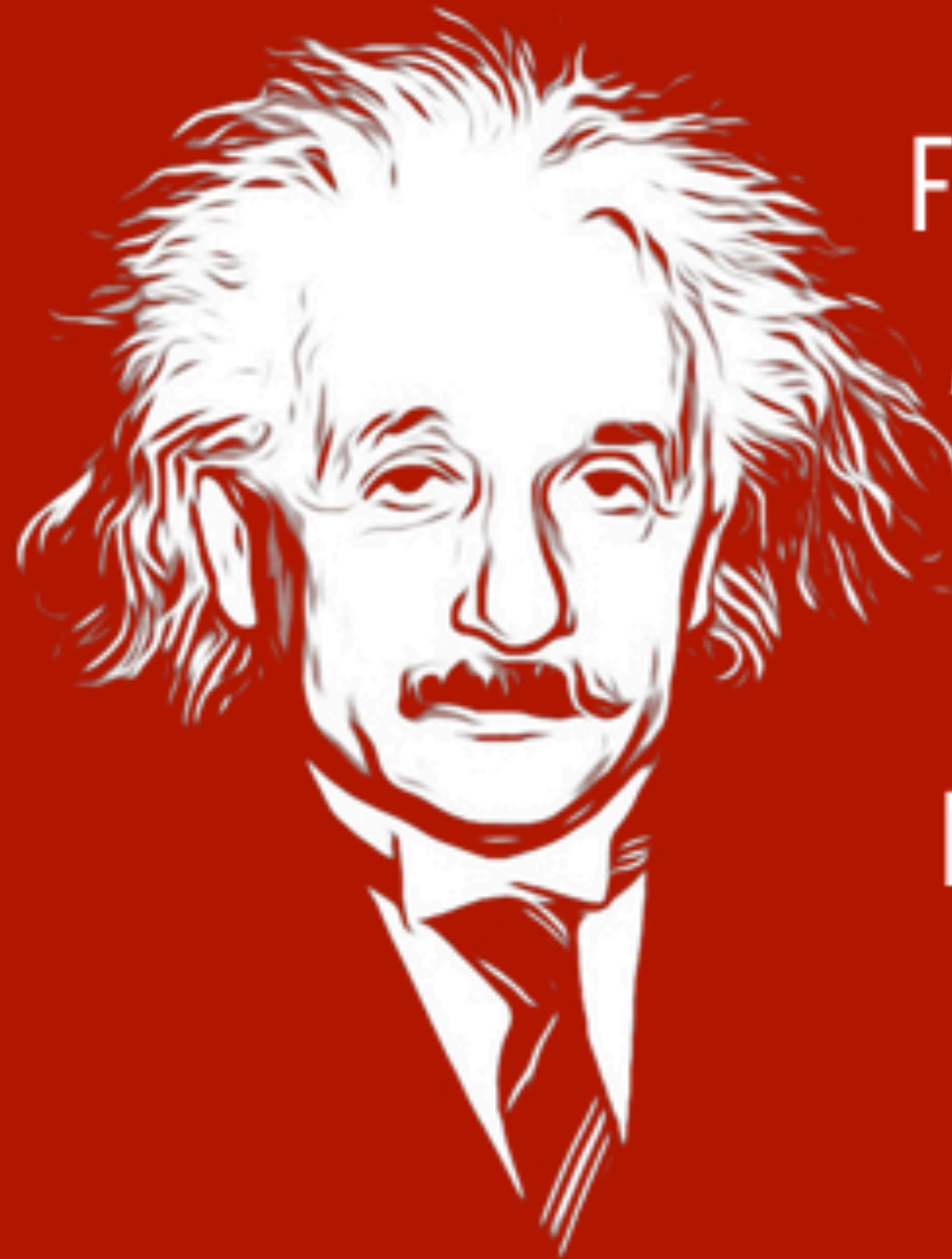
In galaxy clusters (ICM) $P_{gas} \gg P_B, P_{turb}, P_{CR}$

In CGM, ?????



2. Rich Physics

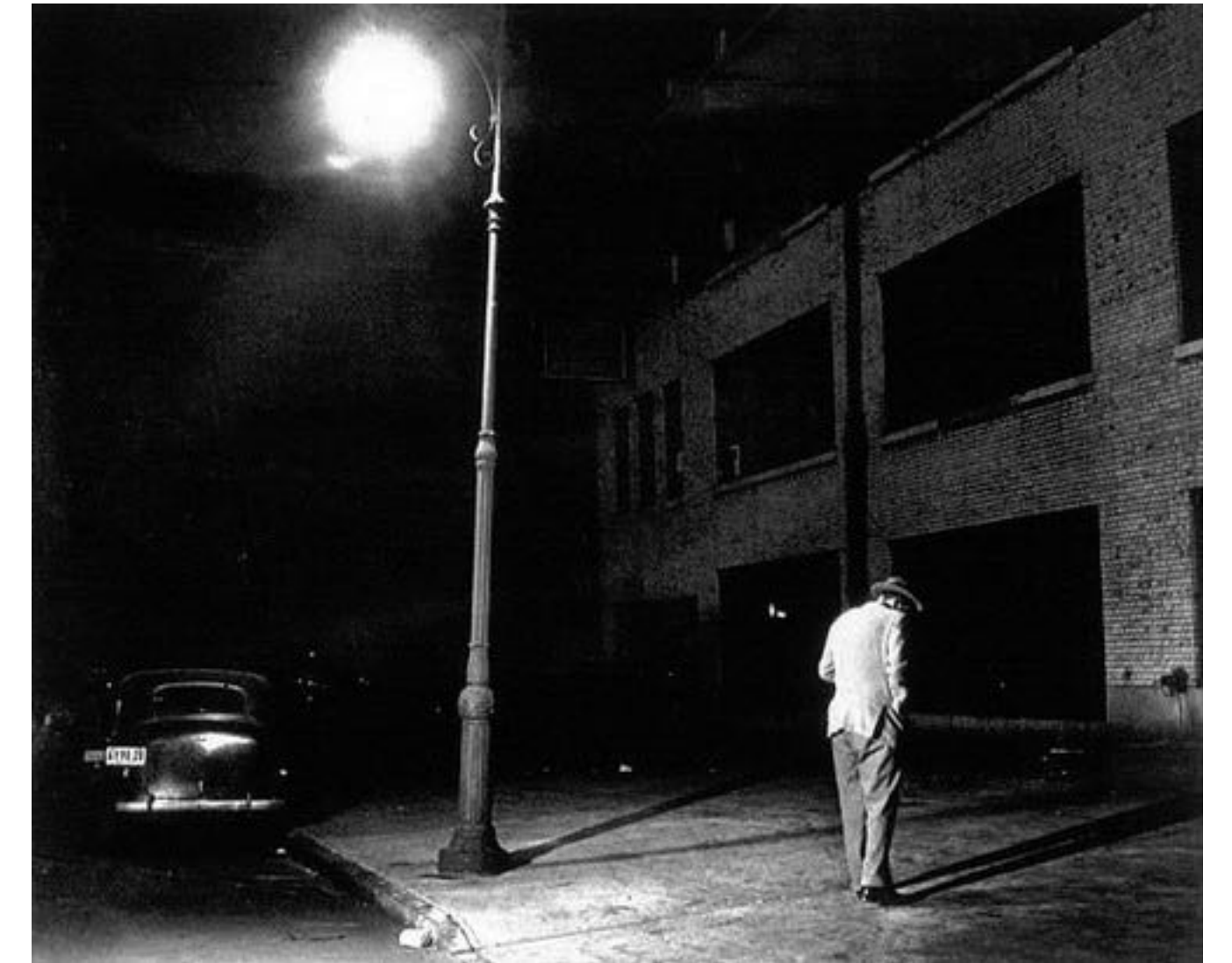
General Approach



“I HAVE LITTLE PATIENCE FOR SCIENTISTS WHO TAKE A BOARD OF WOOD, LOOK FOR ITS THINNEST PART, AND DRILL A GREAT NUMBER OF HOLES WHEN THE DRILLING IS EASY.”

— ALBERT EINSTEIN, 1941

#WEDNESDAYWISDOM



Hmm...that's exactly what I do for a living.

It is better to be lucky than to be smart

And the CGM is a good place to get lucky!



More interesting



CGM theory!



Great if you are Einstein



No jobs

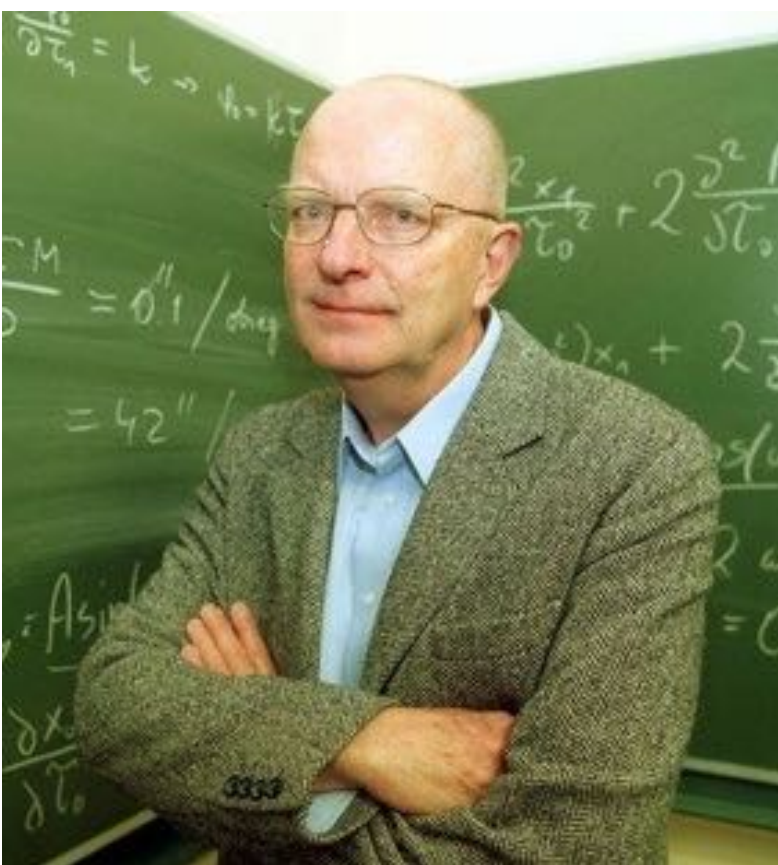


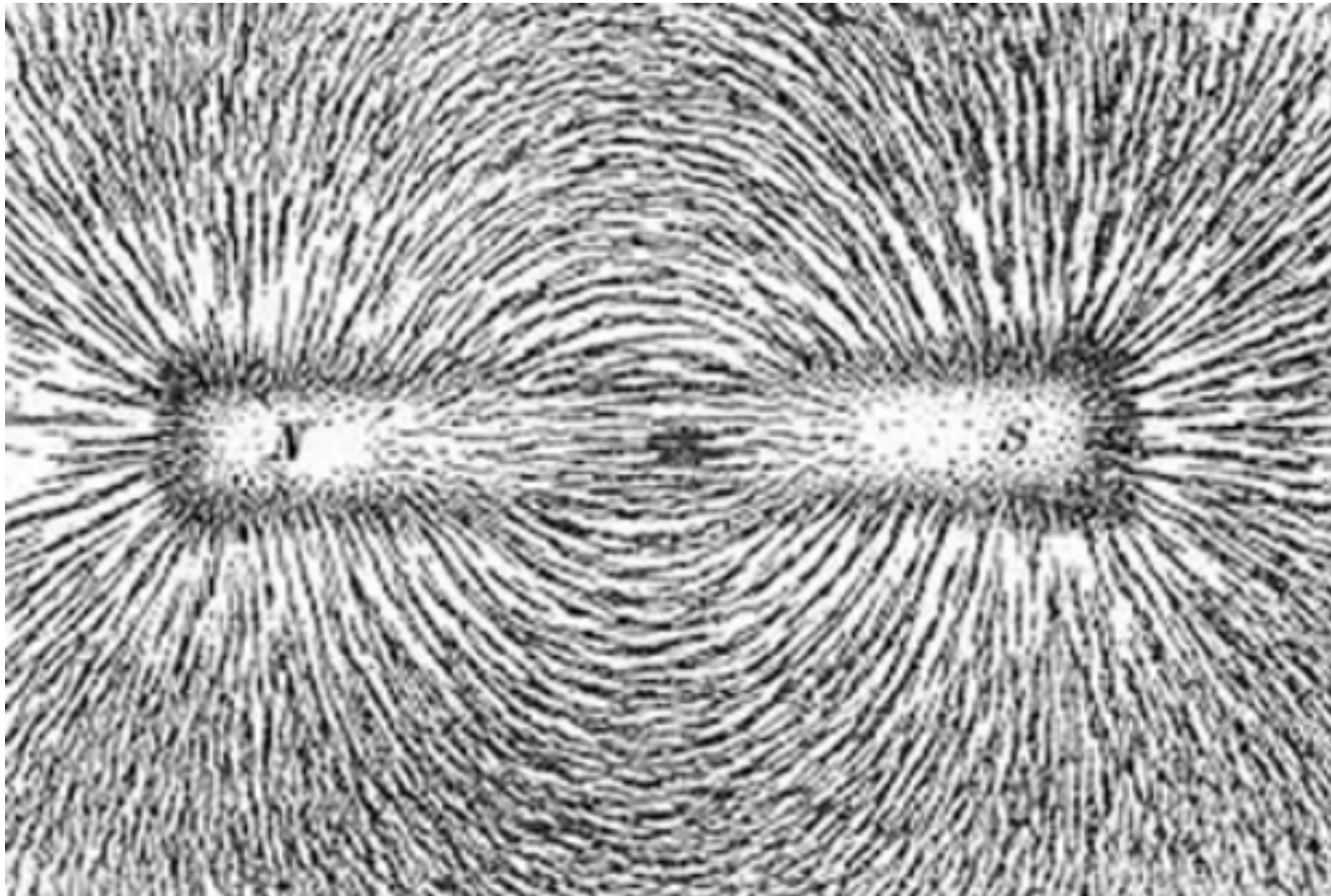
Nervous breakdown

Harder

“A lot of people confuse difficult with interesting.”

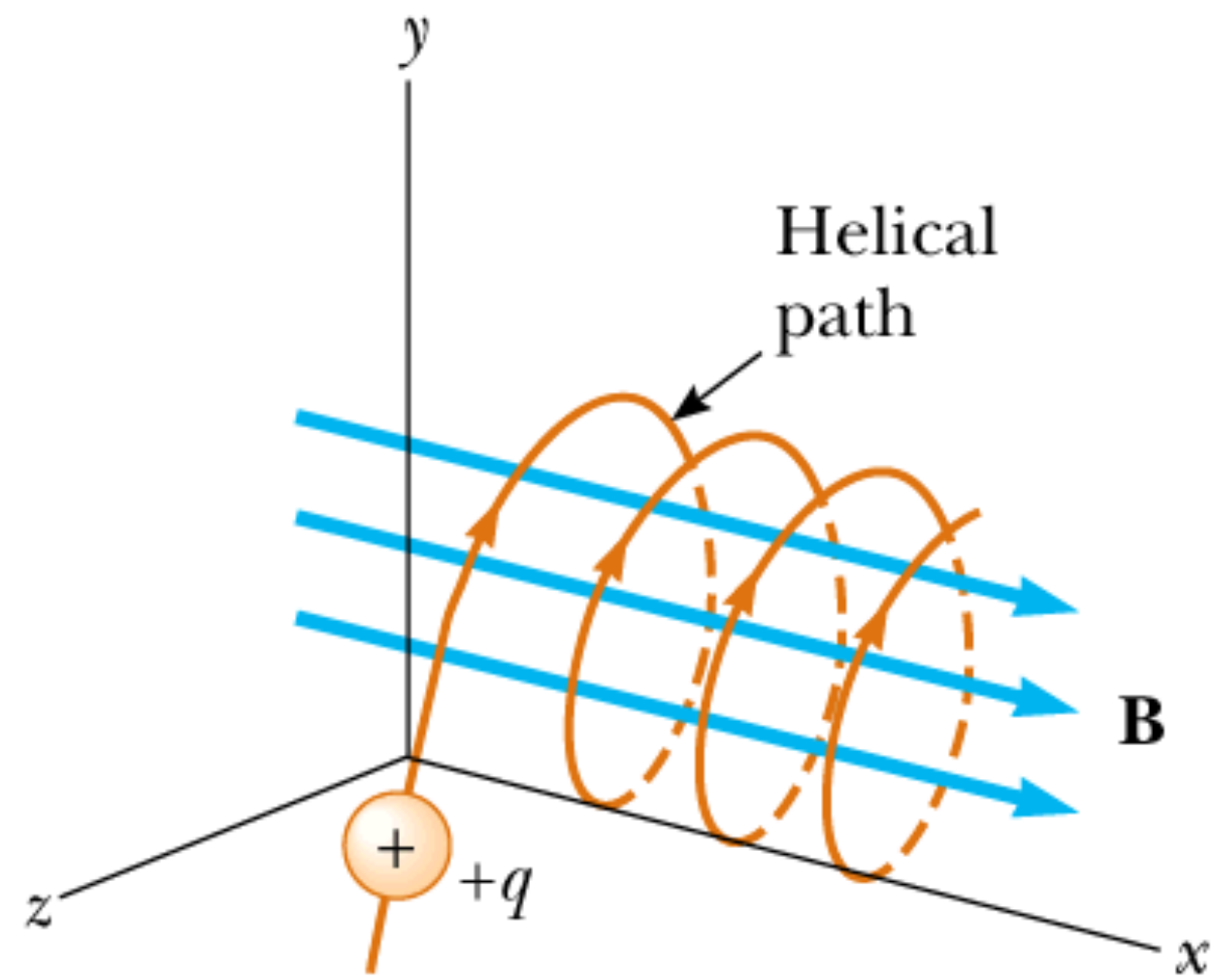
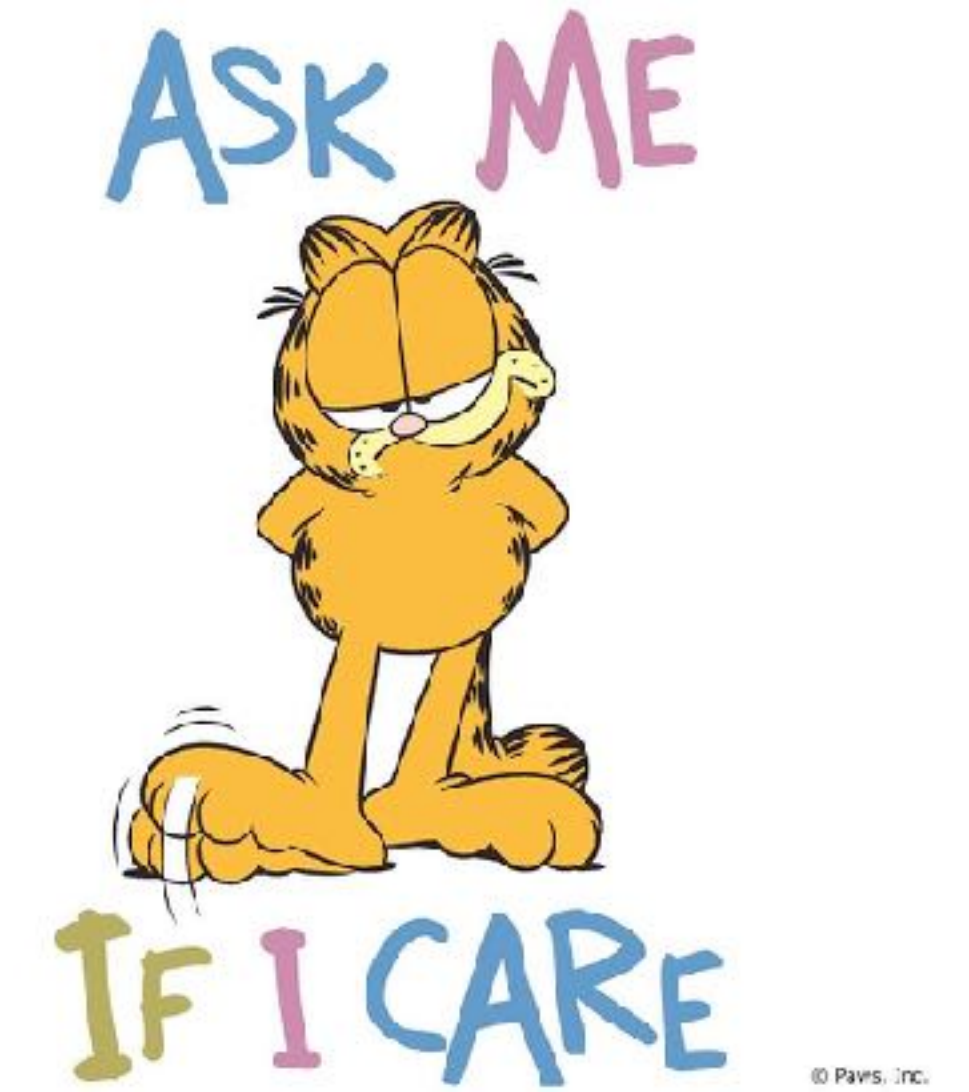
— Bohdan Paczynski





Magnetic Fields

Why care about magnetic fields?

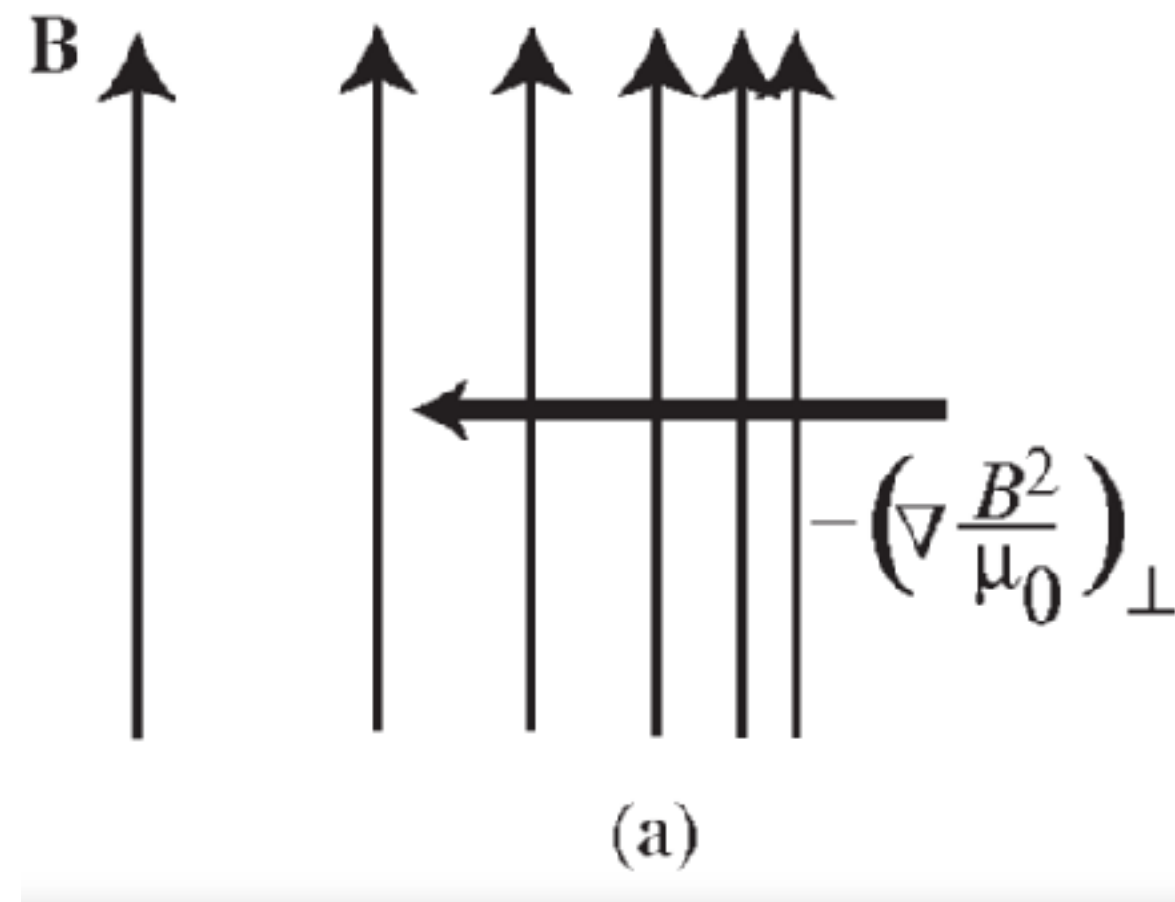


1. Anisotropic Transport

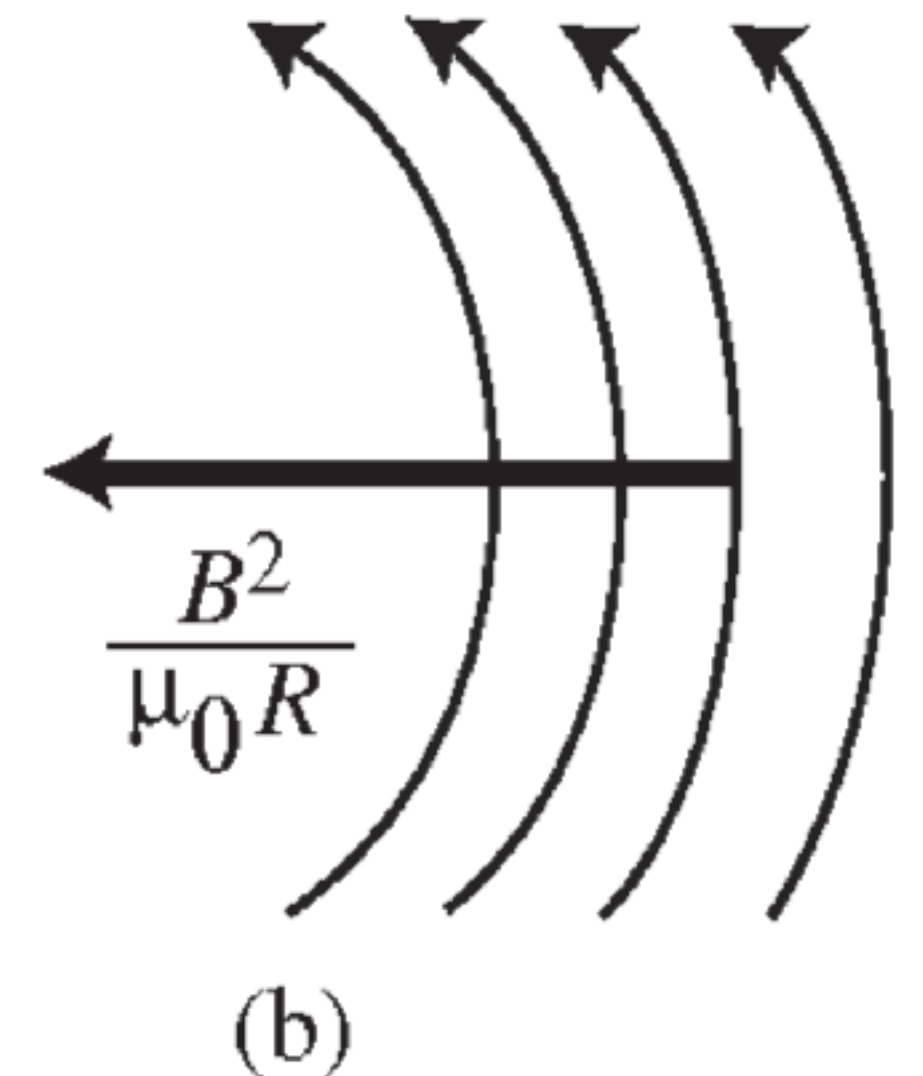
Of particles, momentum, heat...

2. MHD forces

Magnetic pressure
(B-fields don't like to squeeze)

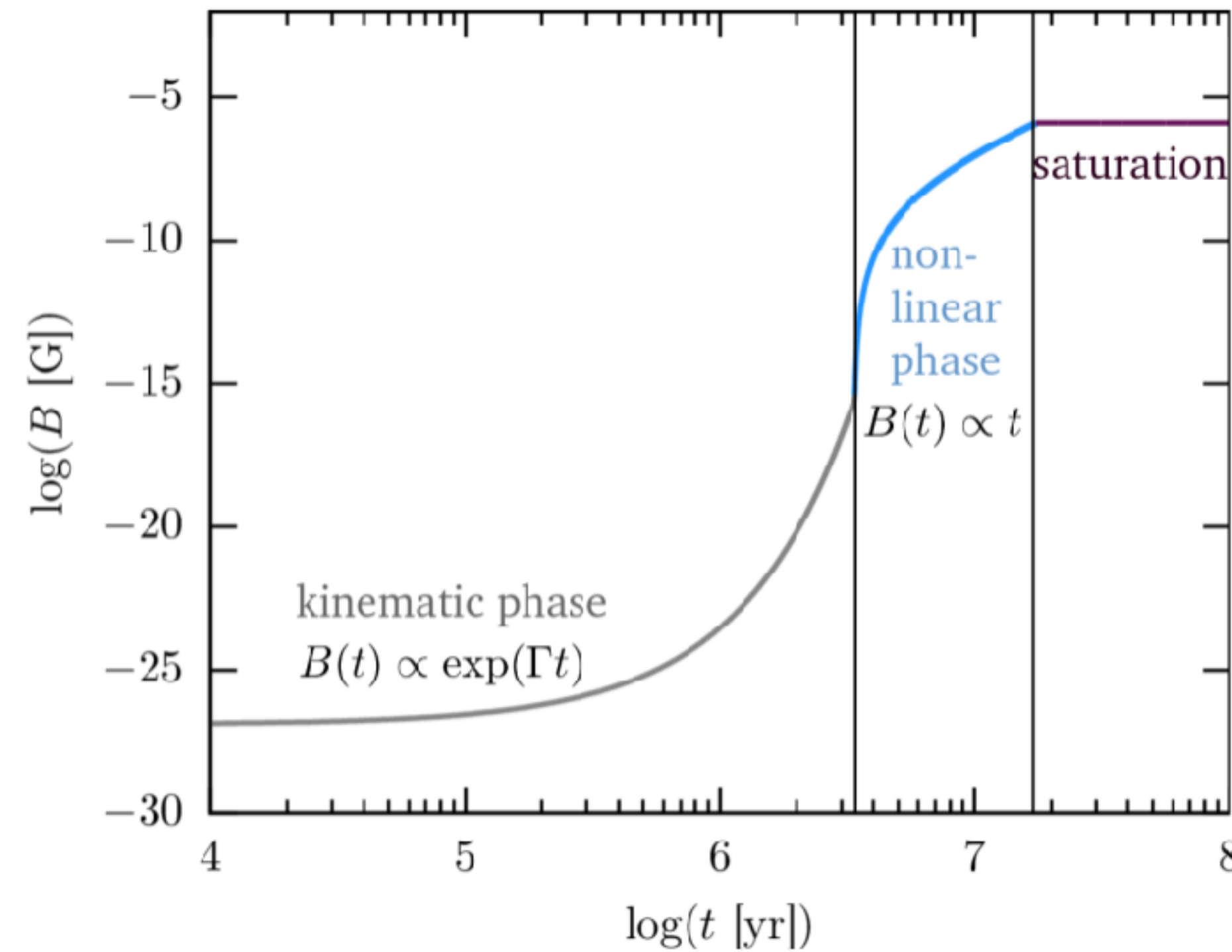
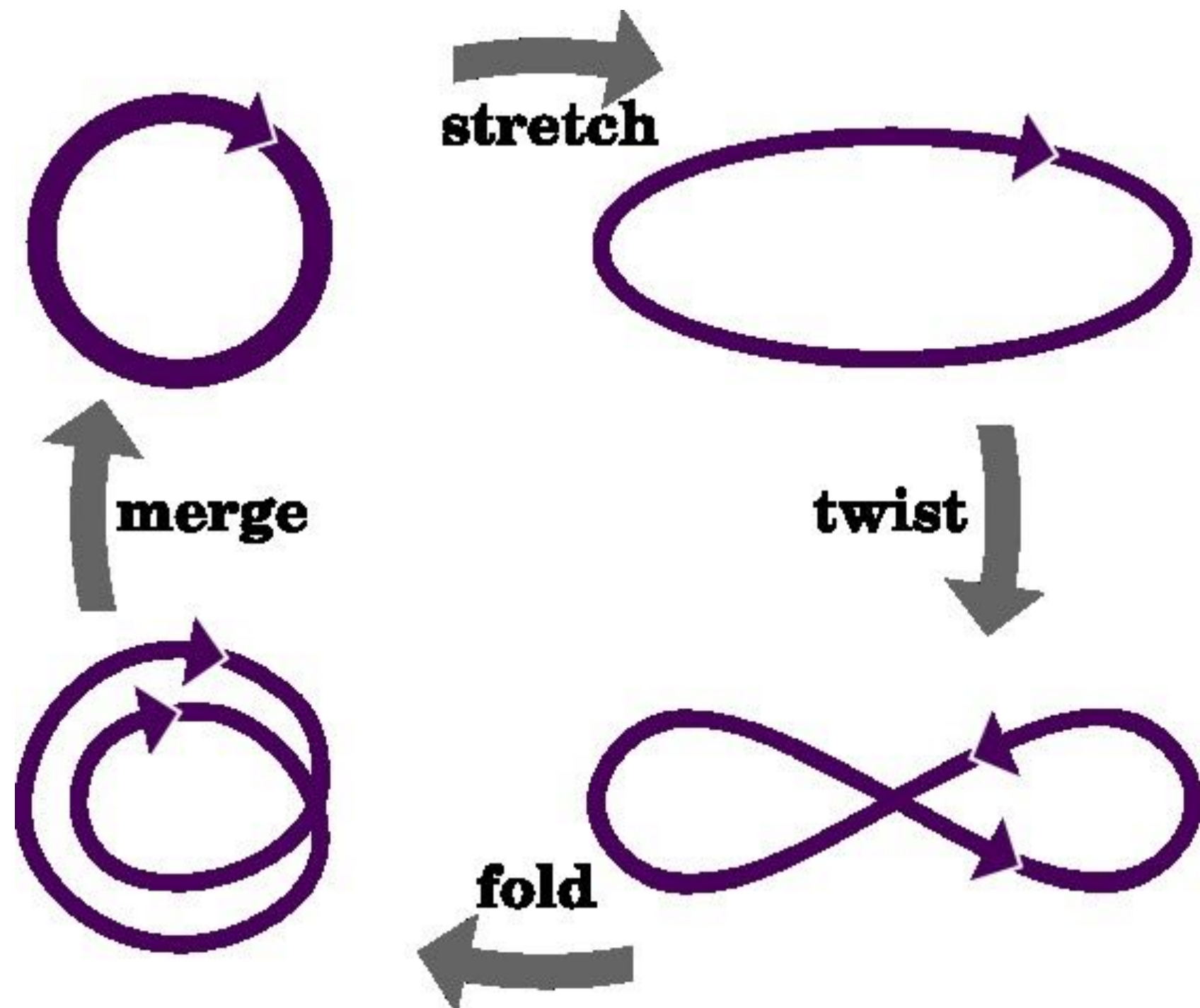


Magnetic tension
(B-fields don't like to bend)



B-fields grow in many ways

Turbulent dynamo is probably most relevant to CGM



Saturates when reaches equipartition with turbulence

Fig credit: J. Schober

What is the B-field in the CGM?

Constraints very uncertain...

FARADAY ROTATION FROM MAGNESIUM II ABSORBERS TOWARD POLARIZED BACKGROUND RADIO SOURCES

J. S. FARNES^{1,2}, S. P. O'SULLIVAN^{1,2}, M. E. CORRIGAN¹, AND B. M. GAENSLER^{1,2}

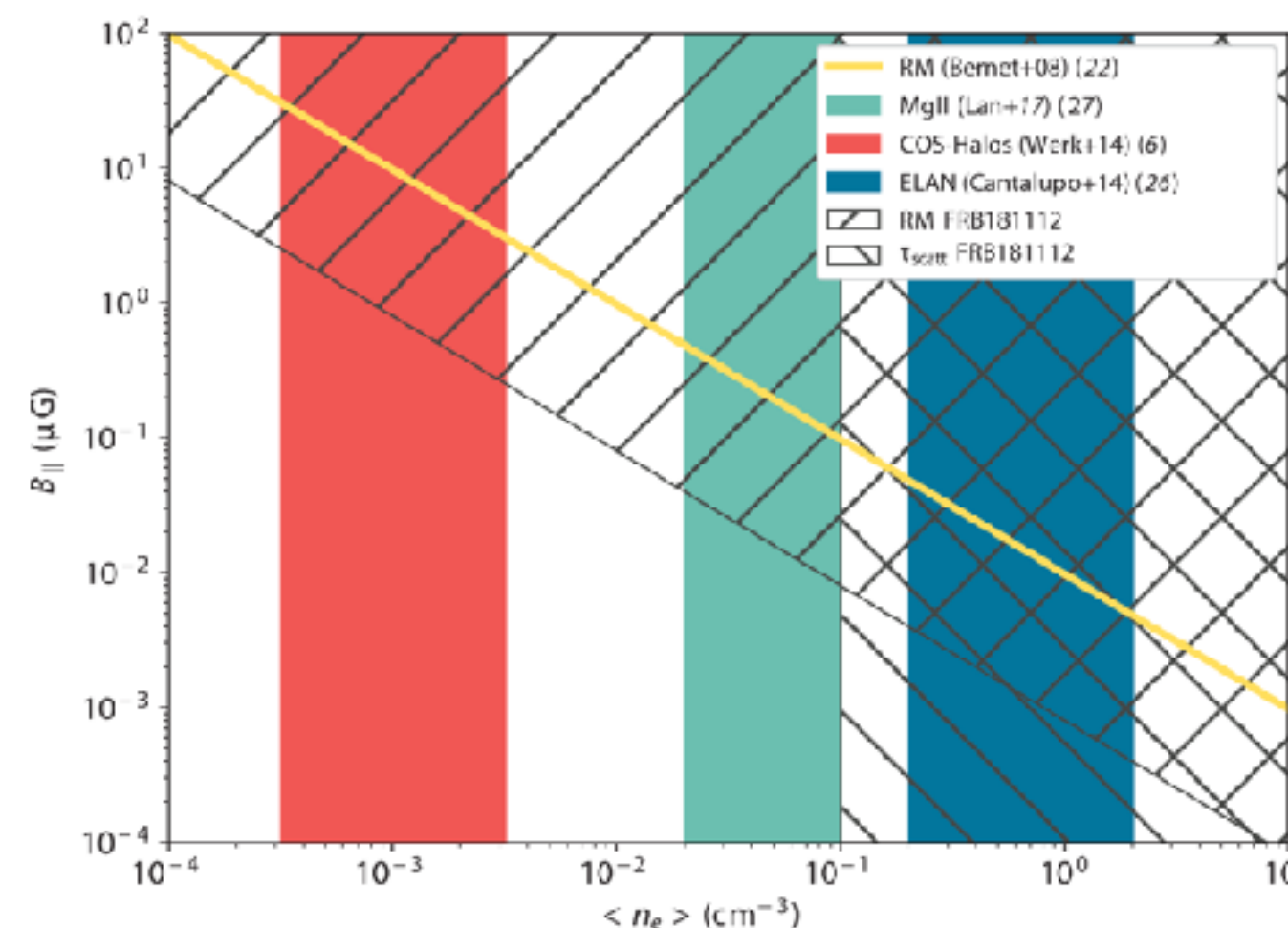
¹ Sydney Institute for Astronomy, School of Physics, The University of Sydney, NSW 2006, Australia; jamie.farnes@sydney.edu.au

² ARC Centre of Excellence for All-sky Astrophysics (CAASTRO), 44 Rosehill Street, Redfern, NSW 2016, Australia

Received 2014 June 2; accepted 2014 September 9; published 2014 October 14

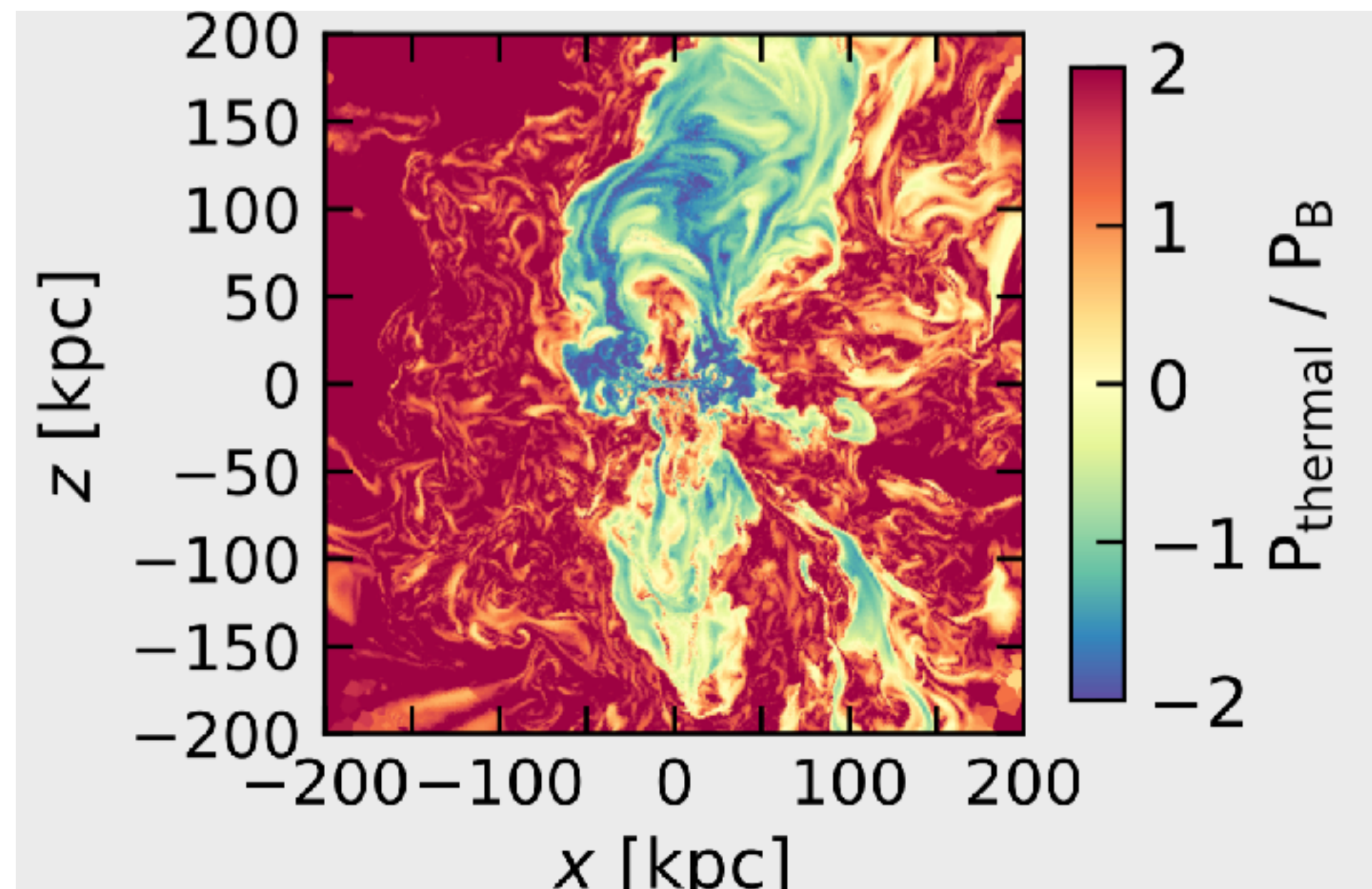
ABSTRACT

Strong singly ionized magnesium (Mg II) absorption lines in quasar spectra typically serve as a proxy for intervening galaxies along the line of sight. Previous studies have found a correlation between the number of these Mg II absorbers and the Faraday rotation measure (RM) at ≈ 5 GHz. We cross-match a sample of 35,752 optically identified non-intrinsic Mg II absorption systems with 25,649 polarized background radio sources for which we have measurements of both the spectral index and RM at 1.4 GHz. We use the spectral index to split the resulting sample of 599 sources into flat-spectrum and steep-spectrum subsamples. We find that our flat-spectrum sample shows significant ($\sim 3.5\sigma$) evidence for a correlation between Mg II absorption and RM at 1.4 GHz, while our steep-spectrum sample shows no such correlation. We argue that such an effect cannot be explained by either luminosity or other observational effects, by evolution in another confounding variable, by wavelength-dependent polarization structure in an active galactic nucleus, by the Galactic foreground, by cosmological expansion, or by partial coverage models. We conclude that our data are most consistent with intervenors directly contributing to the Faraday rotation along the line of sight, and that the intervening systems must therefore have coherent magnetic fields of substantial strength ($\bar{B} = 1.8 \pm 0.4 \mu\text{G}$). Nevertheless, the weak nature of the correlation will require future high-resolution and broadband radio observations in order to place it on a much firmer statistical footing.



Prochaska+20

Parts of CGM could be magnetically dominated



van der Voort+20

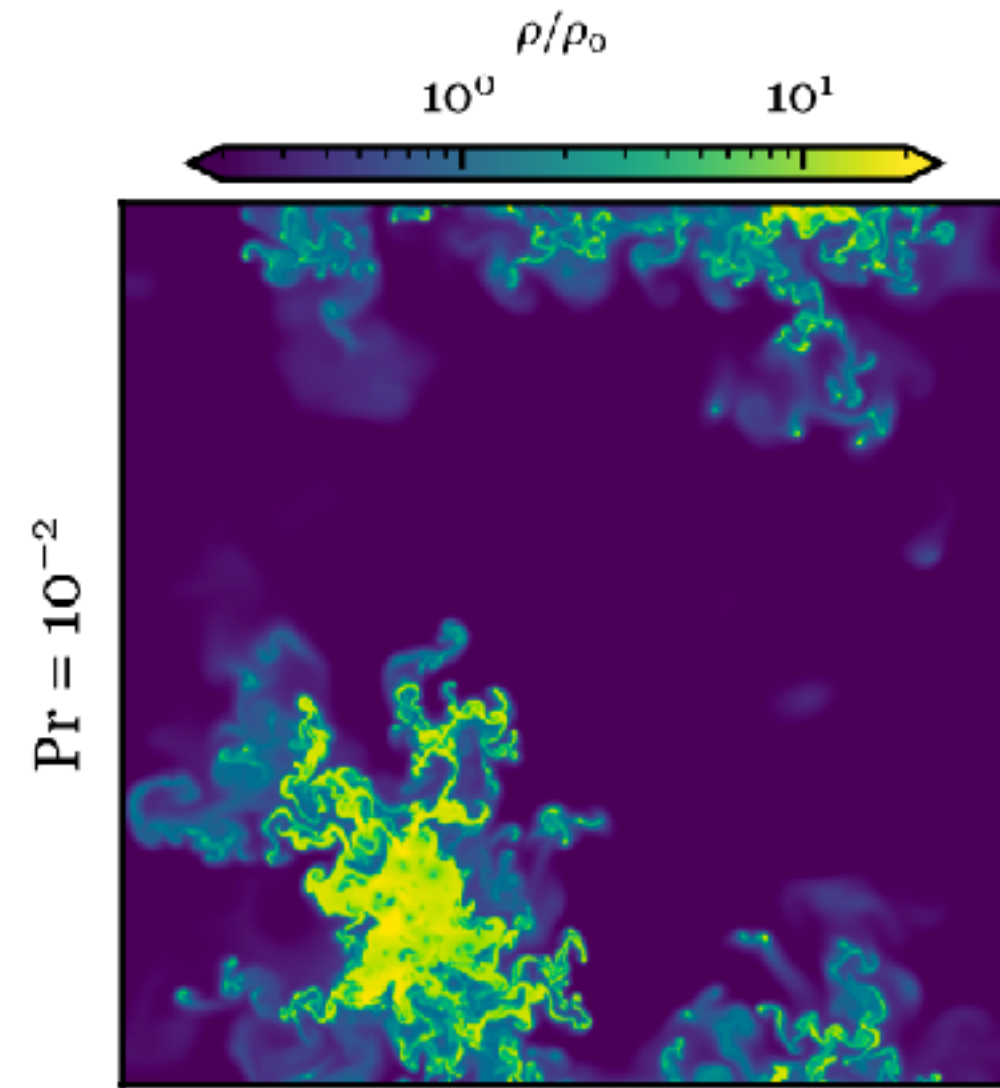
Sims:

$$\beta \ll 1$$

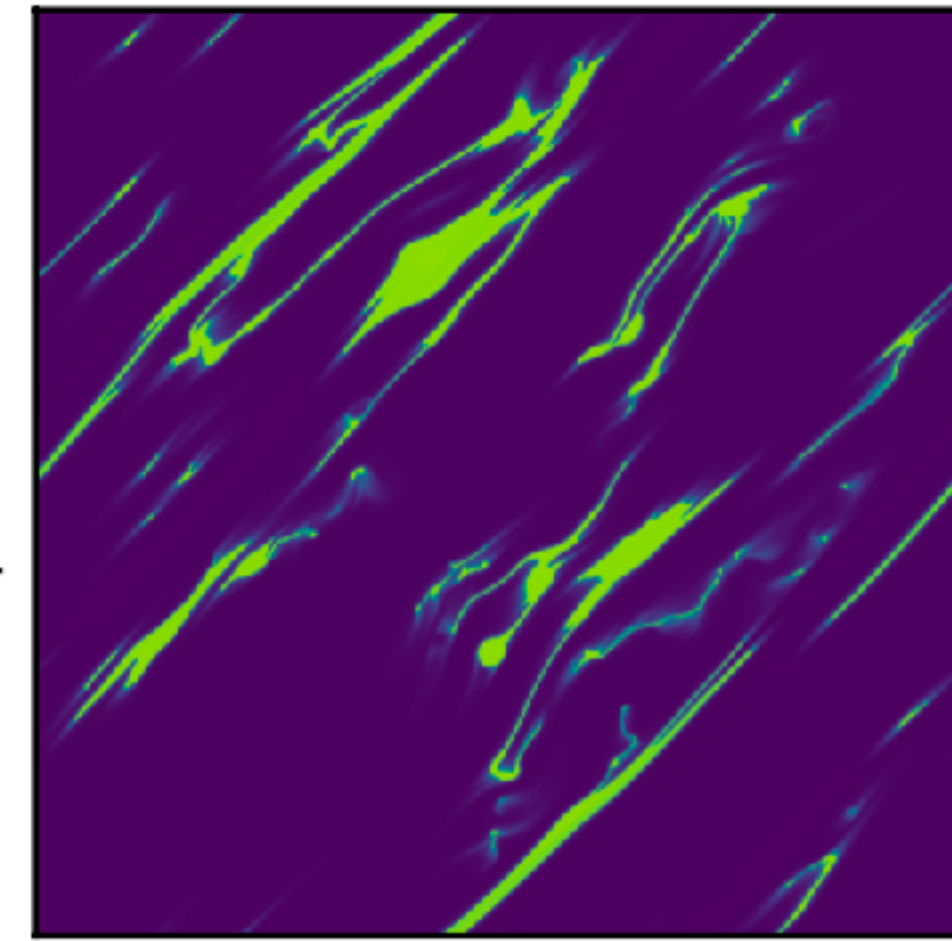
in biconical outflows!

Some CGM implications

1. Very different cloud morphology



Hydro

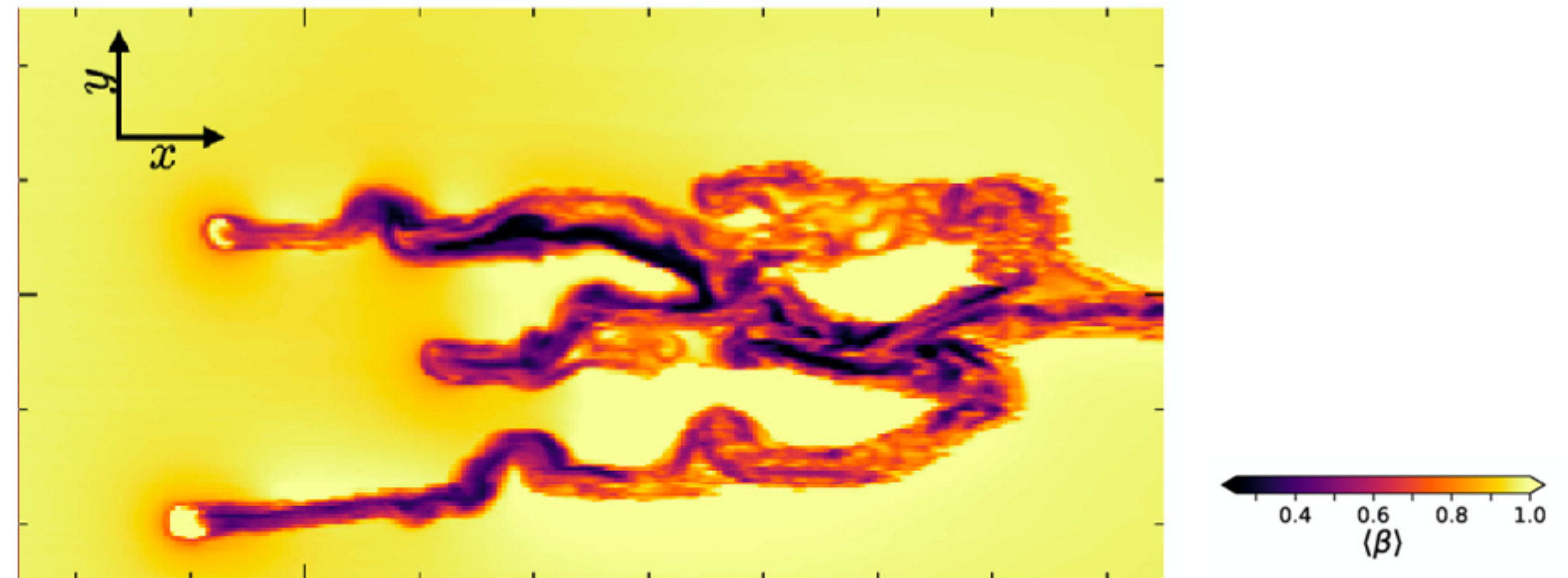


MHD

$$\beta = \frac{P_{\text{gas}}}{P_B}$$

Jennings & Li 2020

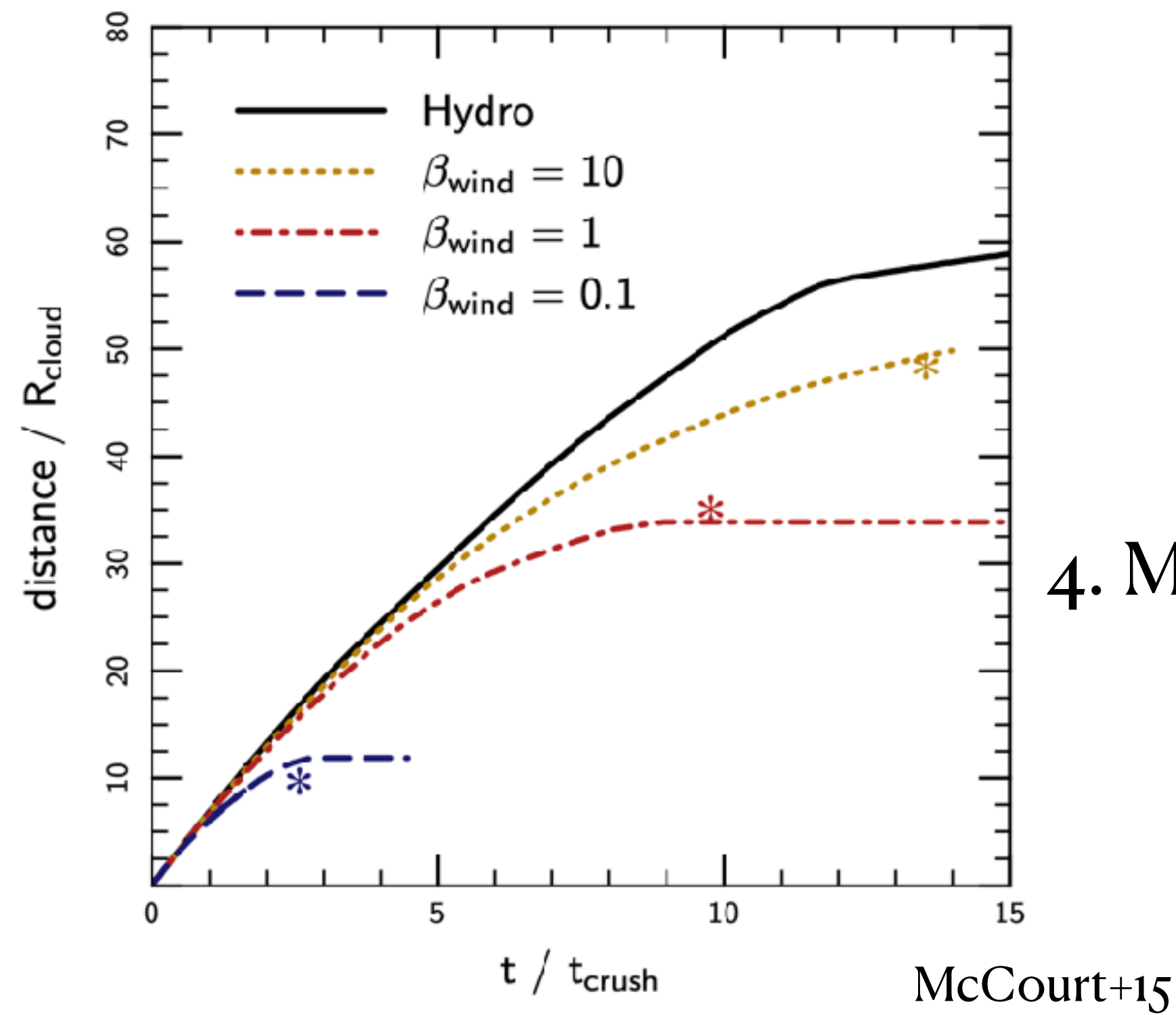
2. Magnetic pressure support in cold gas



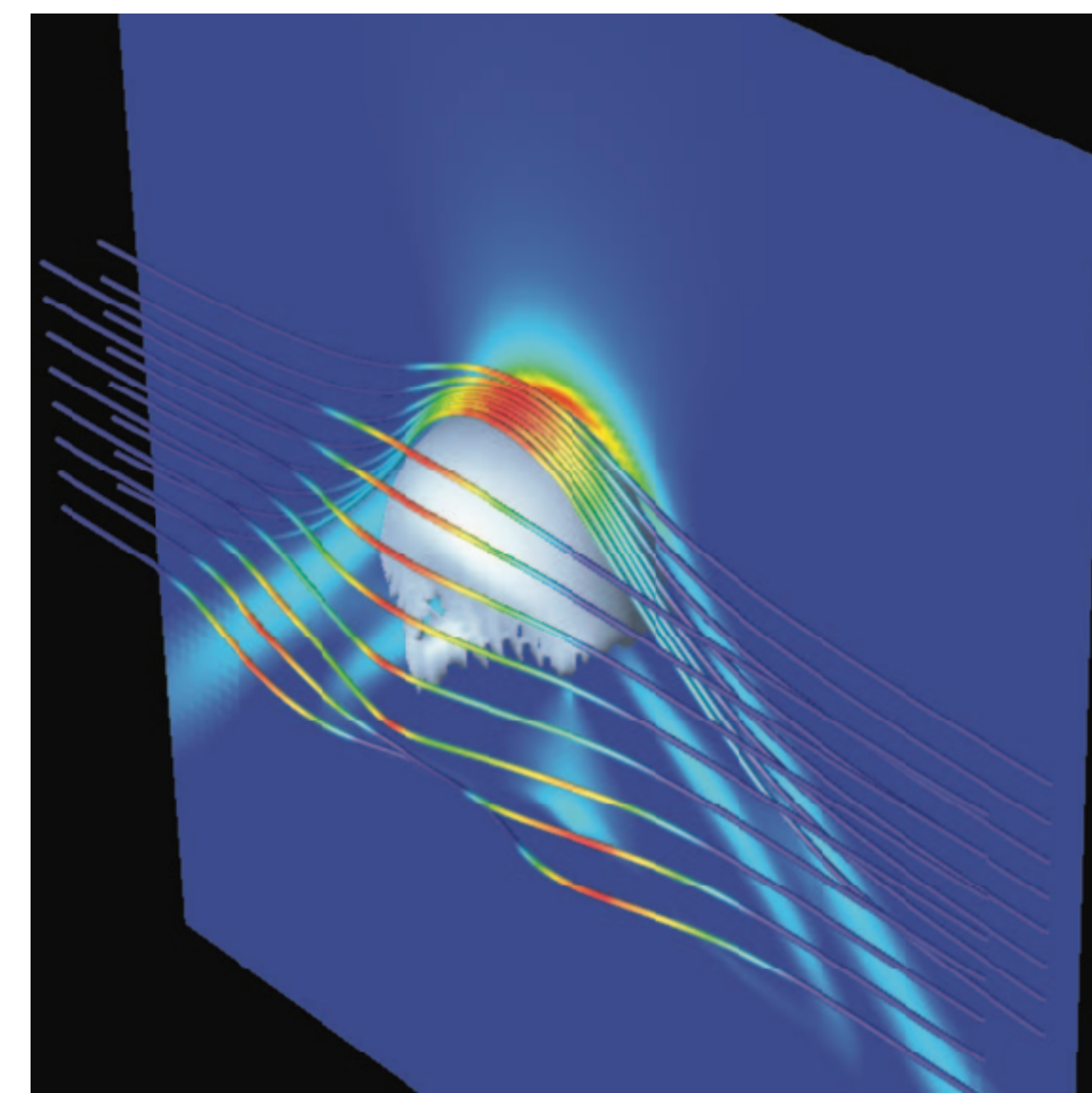
Gronke & Oh 2020

3. Magnetic draping can suppress hydro instabilities (via magnetic tension)

Distance travelled before cold gas comoves with hot wind
Shorter distance means better coupling



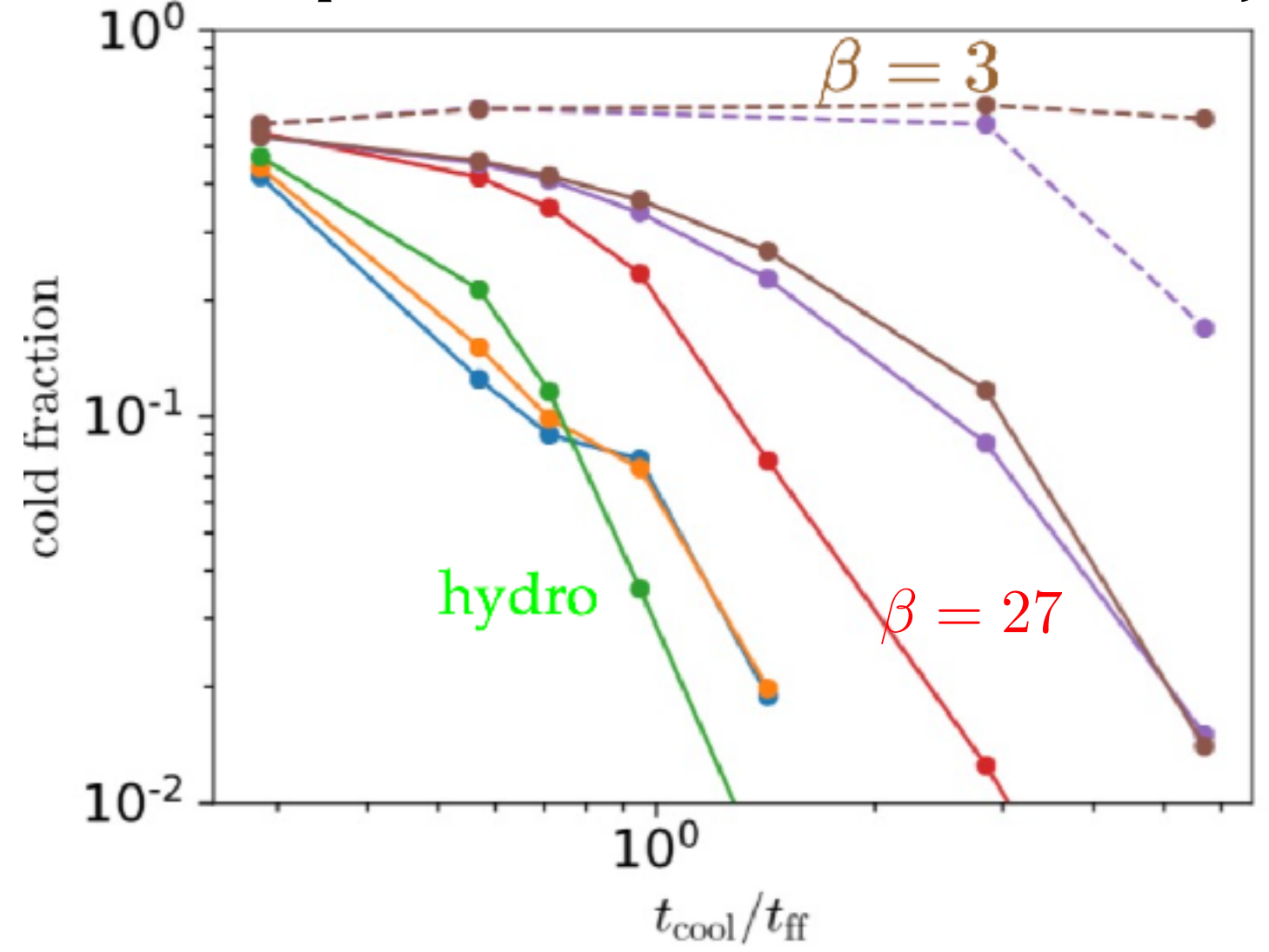
4. Magnetic drag increases momentum coupling of hot and cold phases



Dursi & Pfrommer 2008

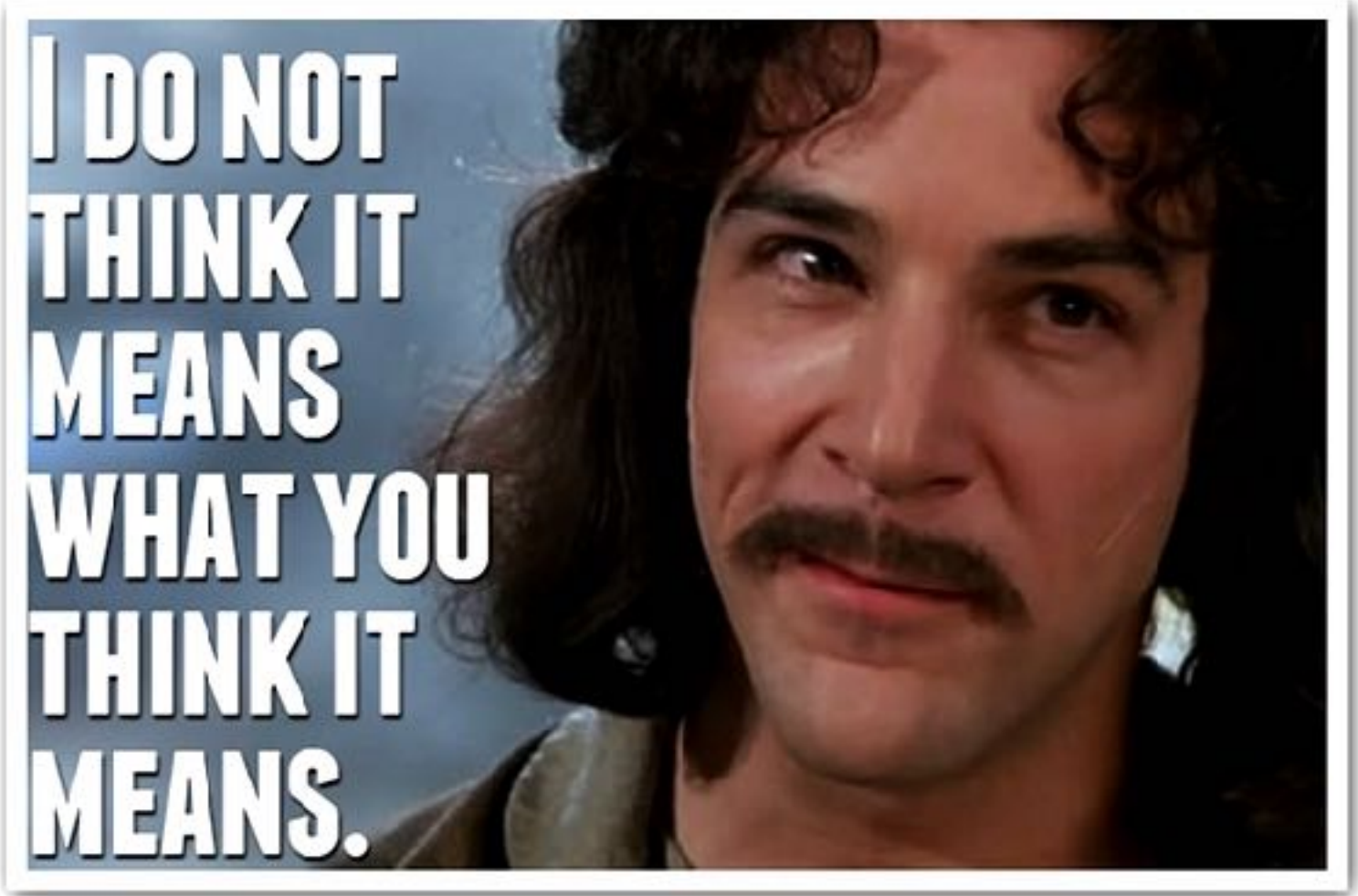
MHD forces can still surprise you

Example: Thermal instability with gravity + MHD



B-fields change threshold $\frac{t_{cool}}{t_{ff}}$ for a multiphase medium

Can be independent of $\frac{t_{cool}}{t_{ff}}$

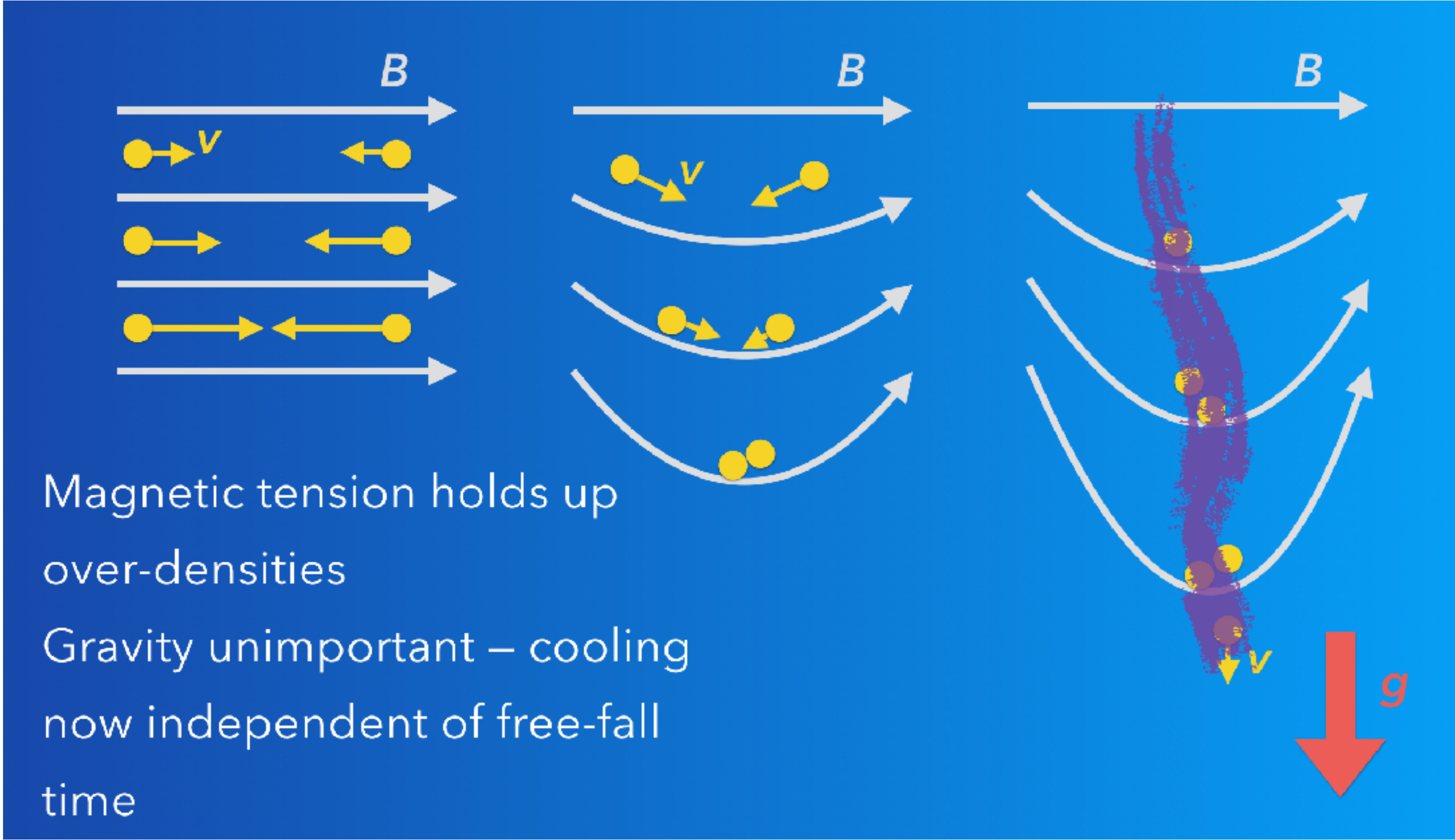


Even very weak fields matter

Because you must compare perturbed forces, not background forces



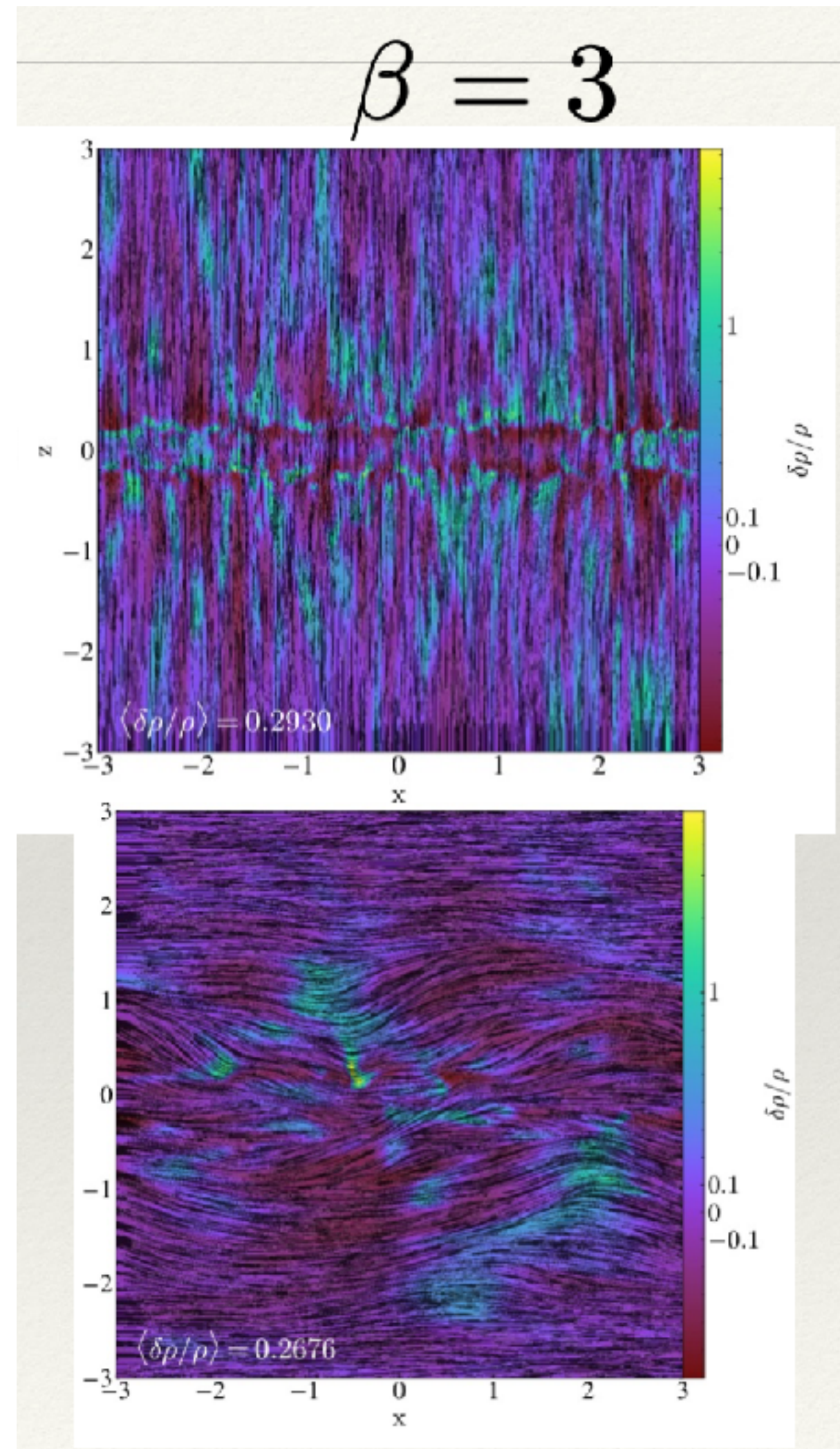
So far, so good.



here's the weird thing...

horizontal and vertical fields give the same cold gas mass!

even though they look very different

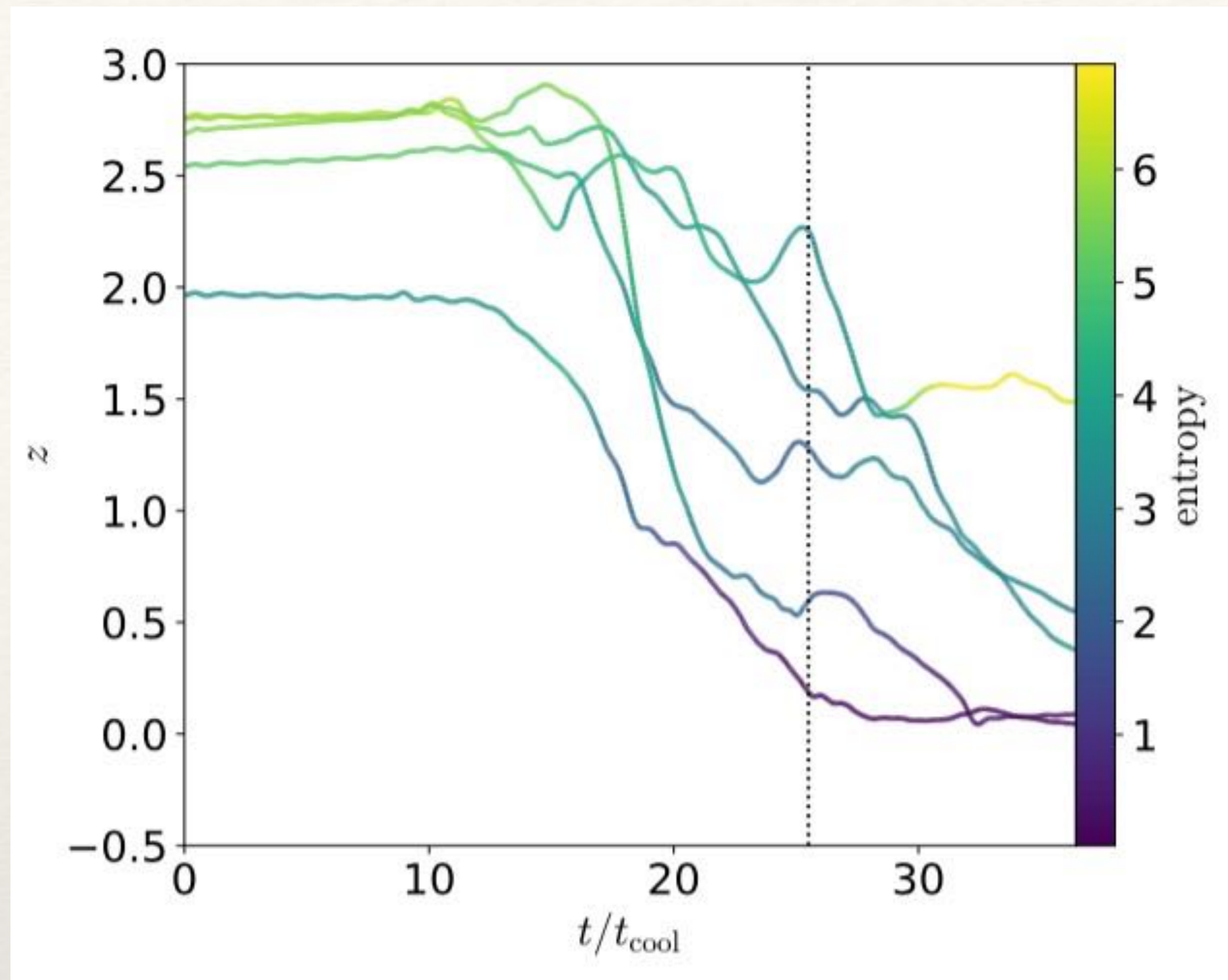


vertical field

$$\frac{\delta\rho}{\rho} = 0.29$$

horizontal field

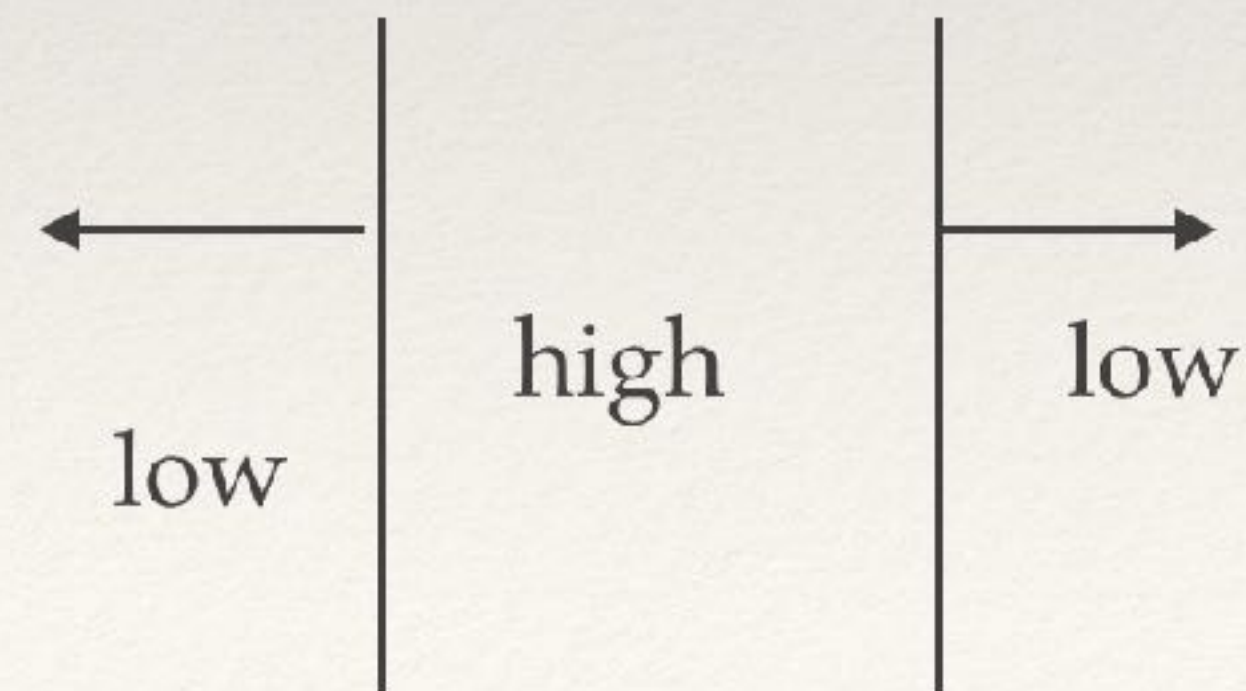
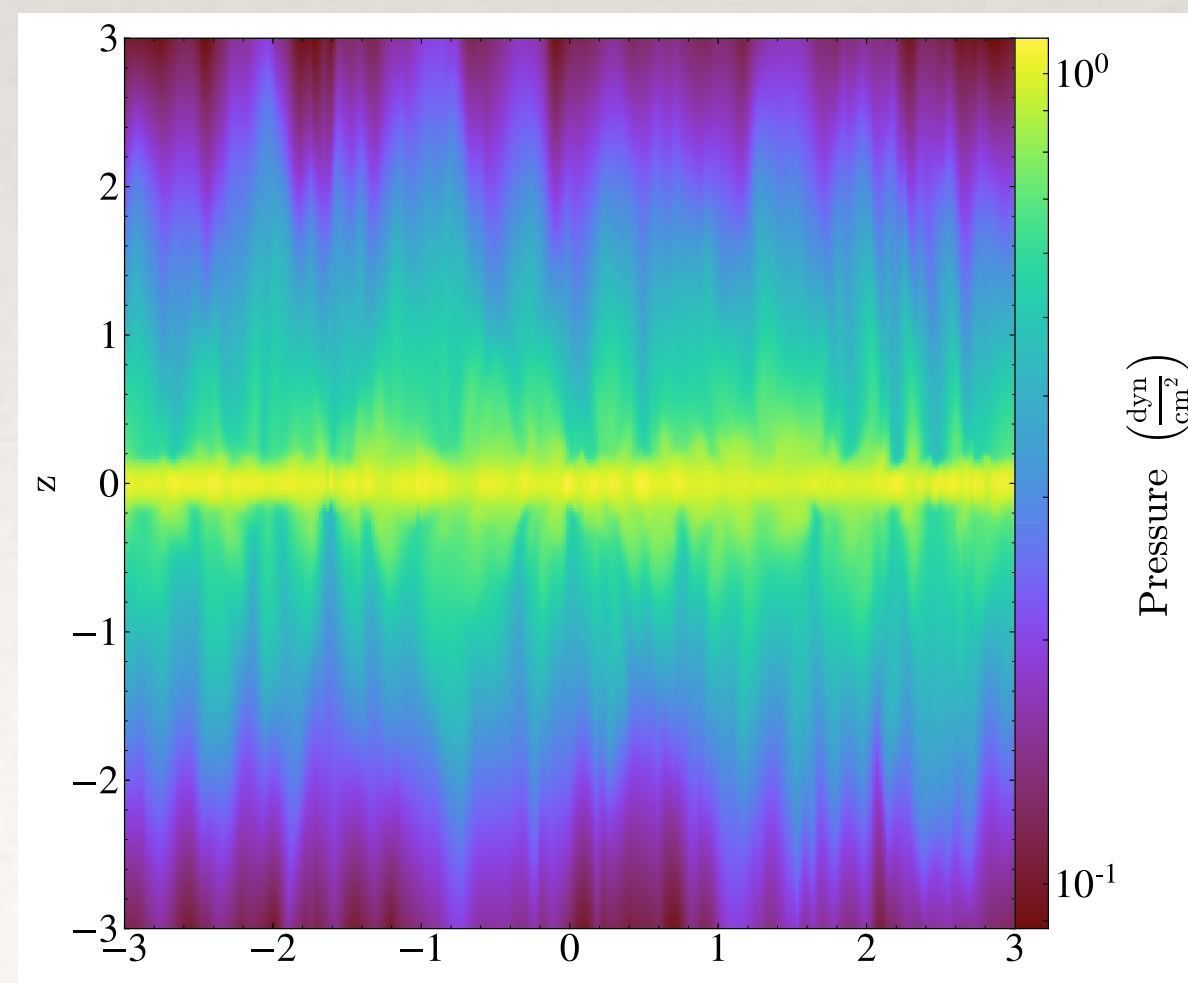
$$\frac{\delta\rho}{\rho} = 0.27$$



$$\beta = 3$$

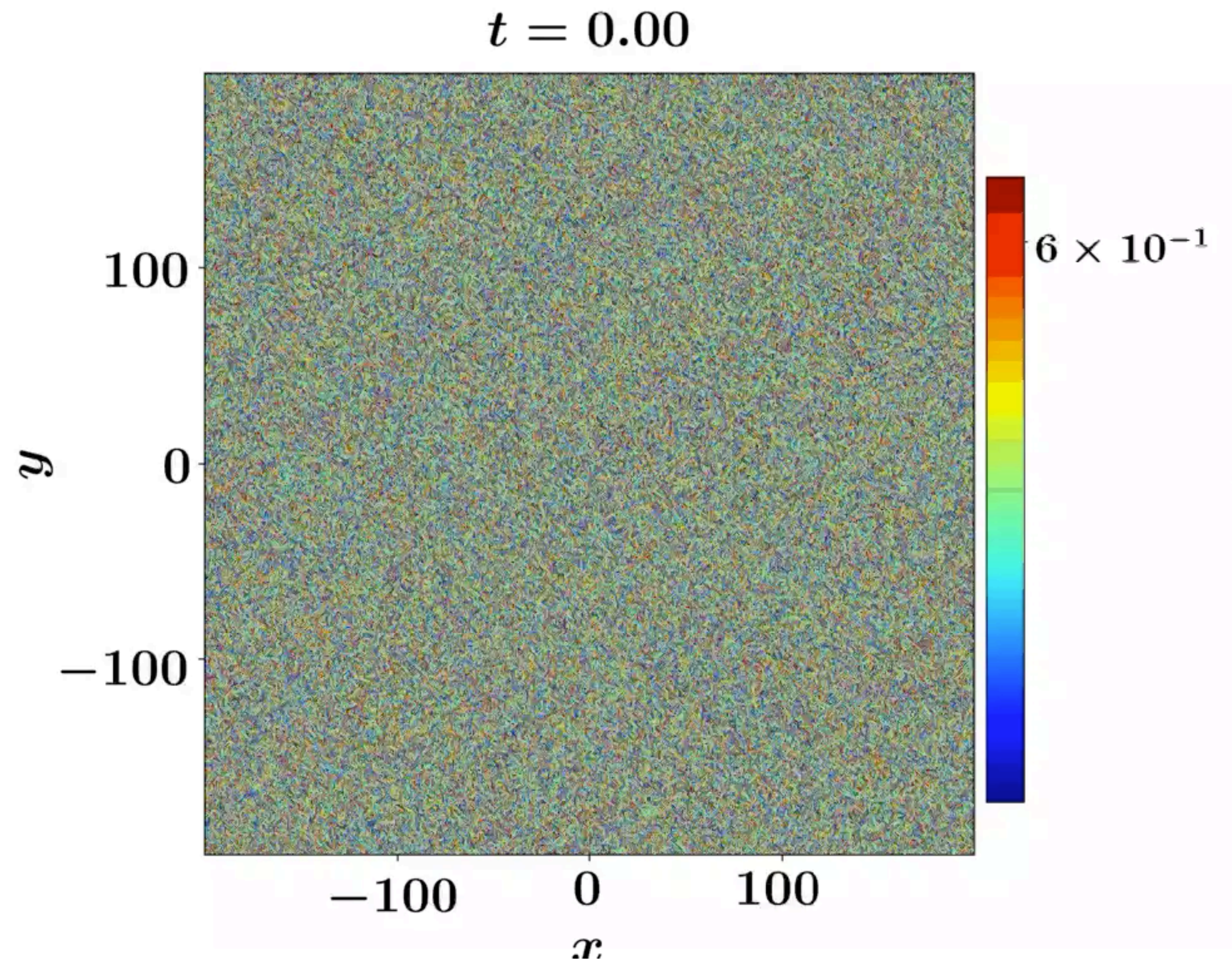
Particles also don't move!
Buoyancy also suppressed

Supported by high pressure regions ...which develop because of confining tension



Here's another head-scratcher

Outcome of thermal instability w/ B-fields. Solitons??



Jiang & Oh 2021

Turbulence



When I meet God, I'm going to ask him two questions: why relativity?
And why turbulence? I really believe he'll have an answer for the first.

—Werner Heisenberg

What's Turbulence?

...more than just non-thermal motions

- energy cascades (usually from large to small scales)
- comes from non-linear term in hydro equations
- Incompressible hydro: Kolmogorov turbulence

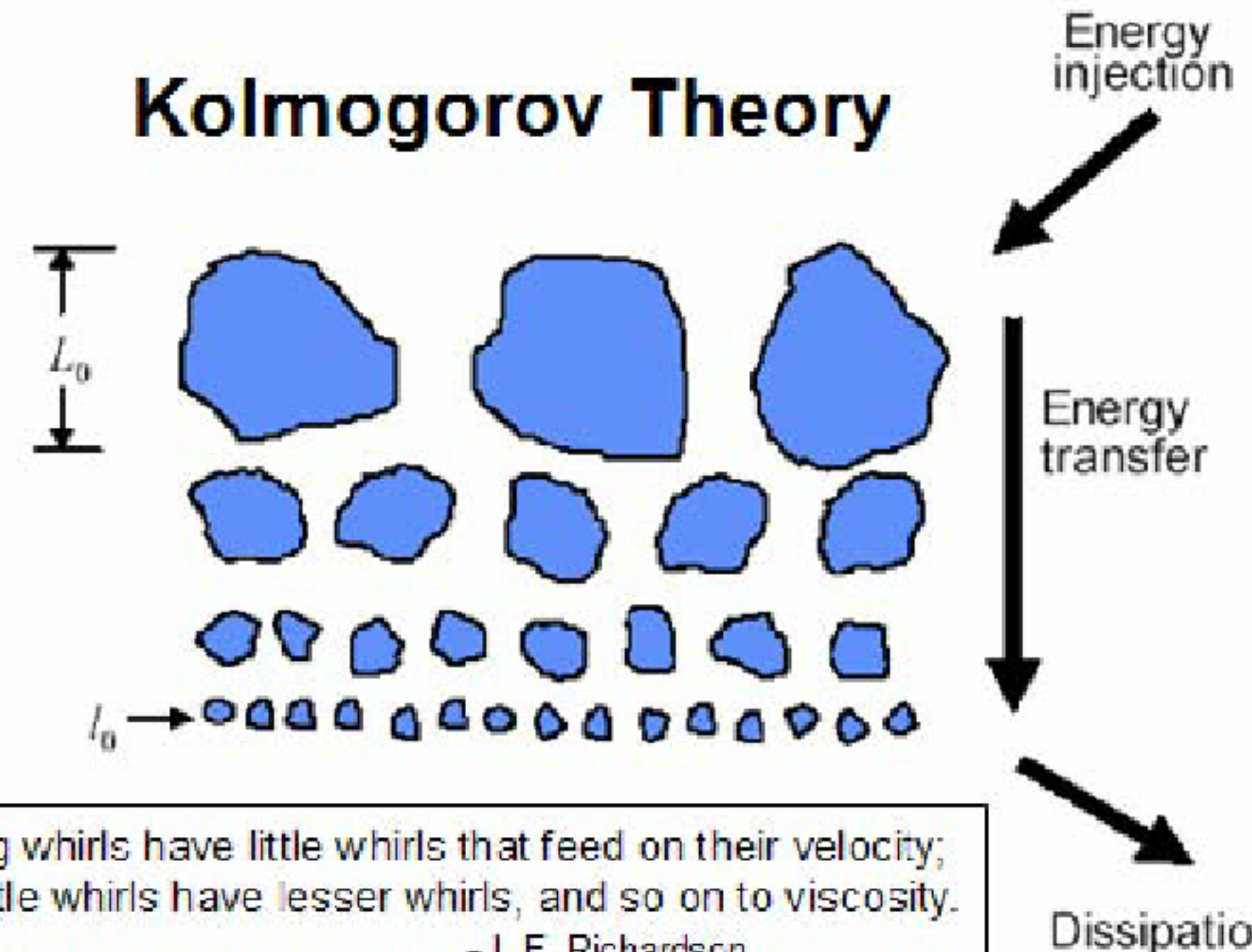
$$\mathbf{v} \cdot \nabla \mathbf{v}$$



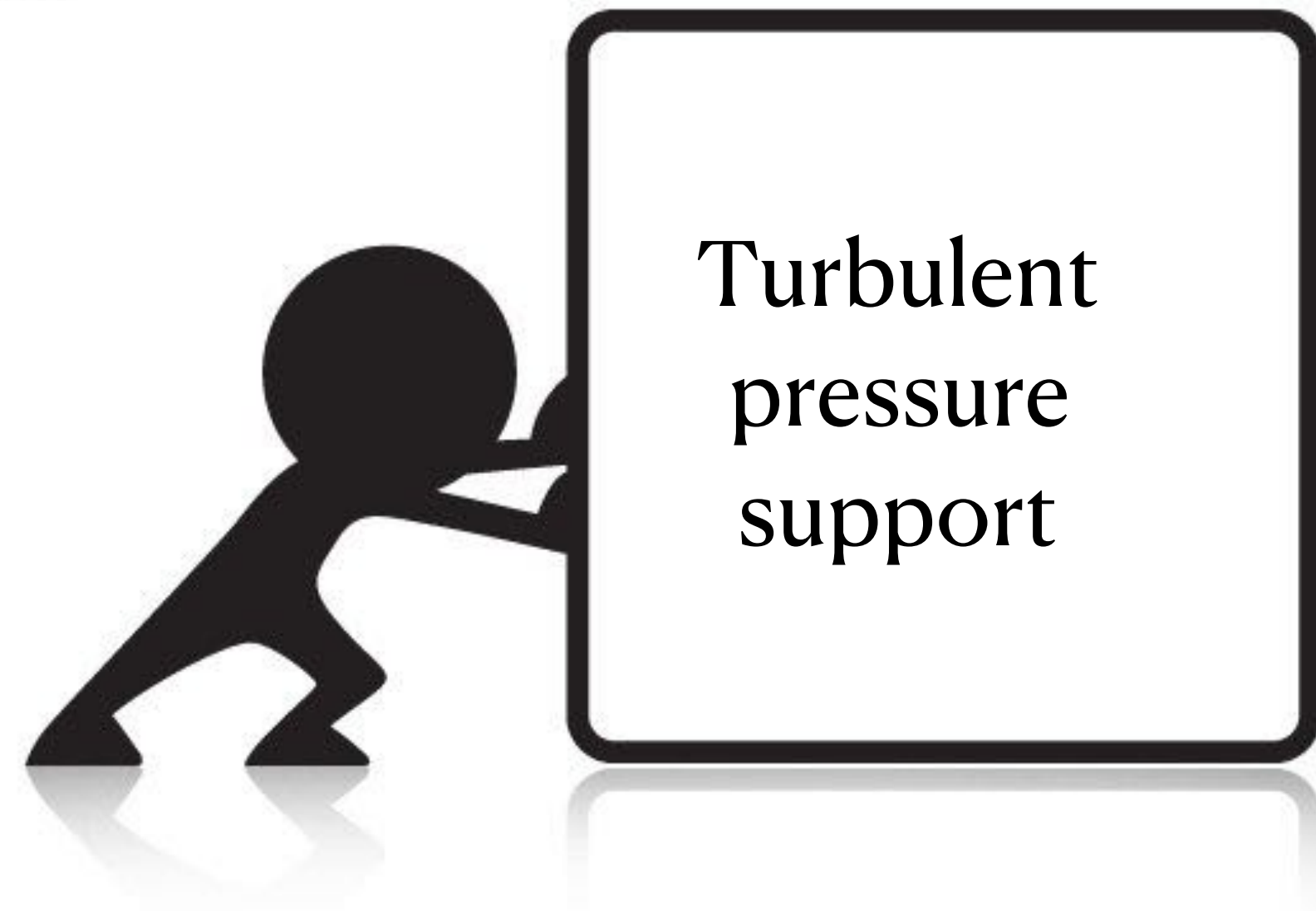
$$\frac{E_{\text{kinetic}}}{t_{\text{eddy}}} \sim \frac{\rho v^3}{l} \sim \text{const}$$

More to life than Kolmogorov:
compressible turbulence (jump directly to small scales!),
MHD (3 independent cascades!),
stratified turbulence (vertical motions suppressed)...
decaying turbulence, etc

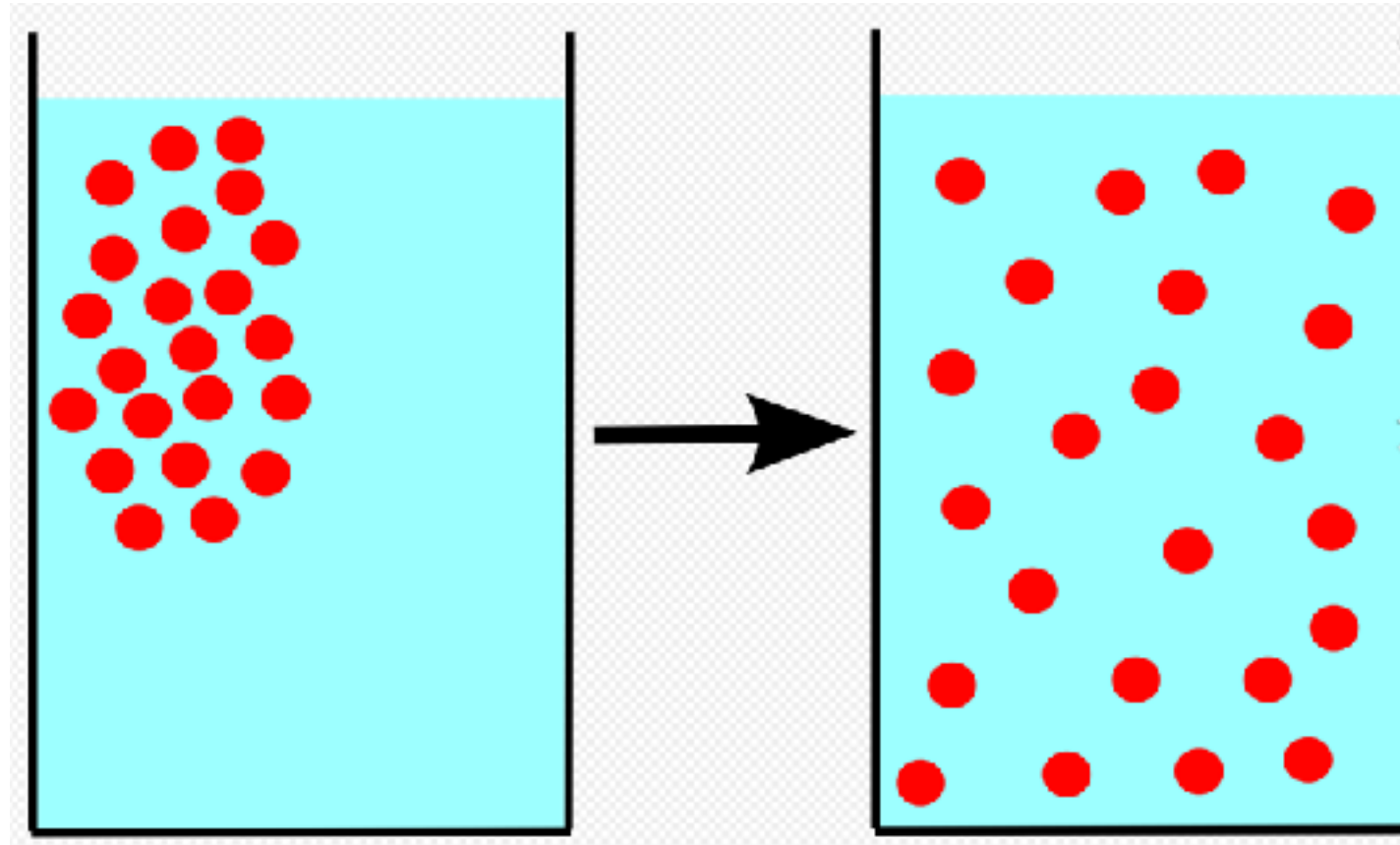
Kolmogorov Theory



Why should you care about turbulence?

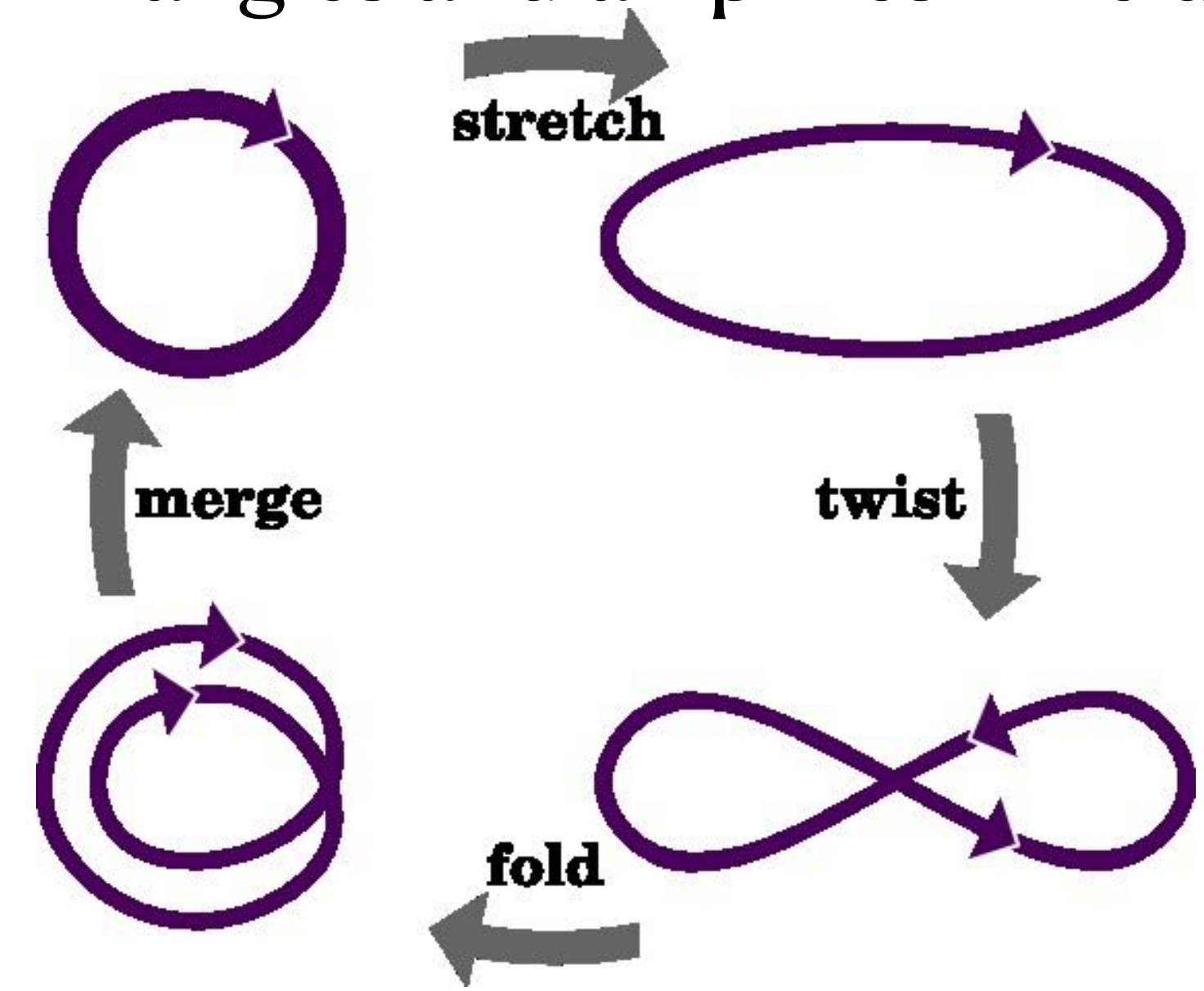


Diffuses (metals, entropy ...)

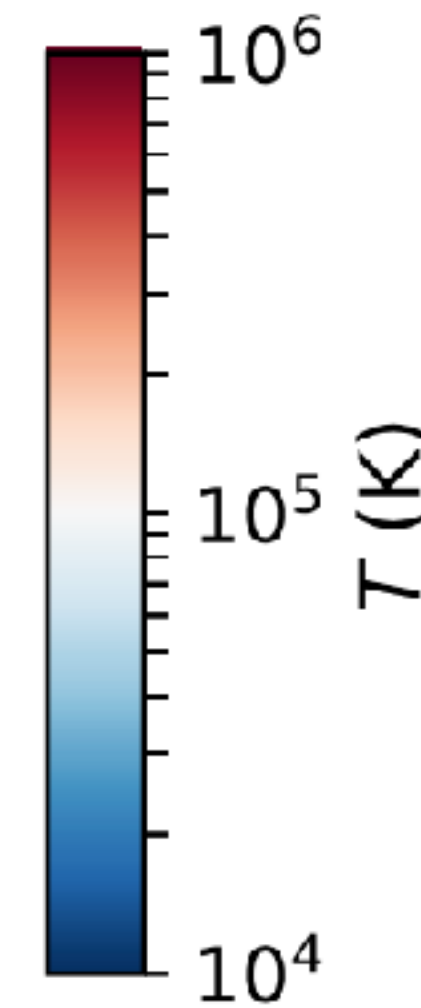
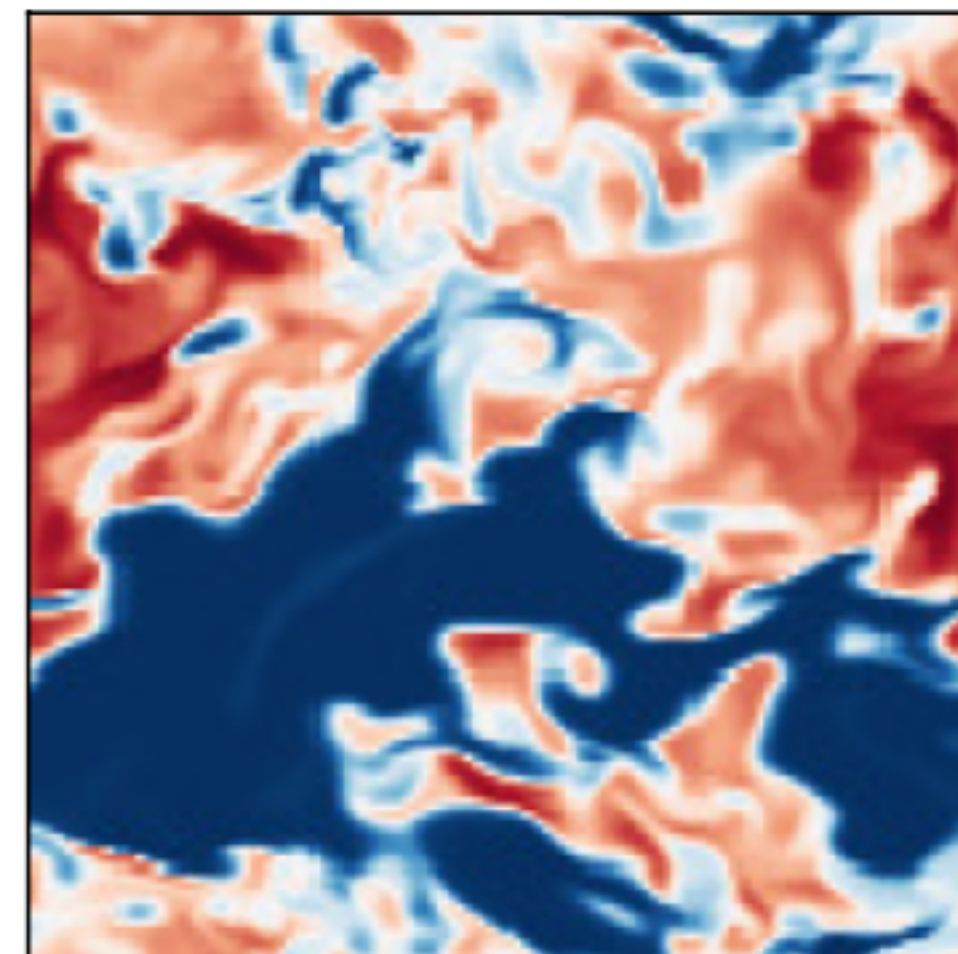


$$f_{\text{cold}} = 0.5$$

Tangles and amplifies B-fields



Heats

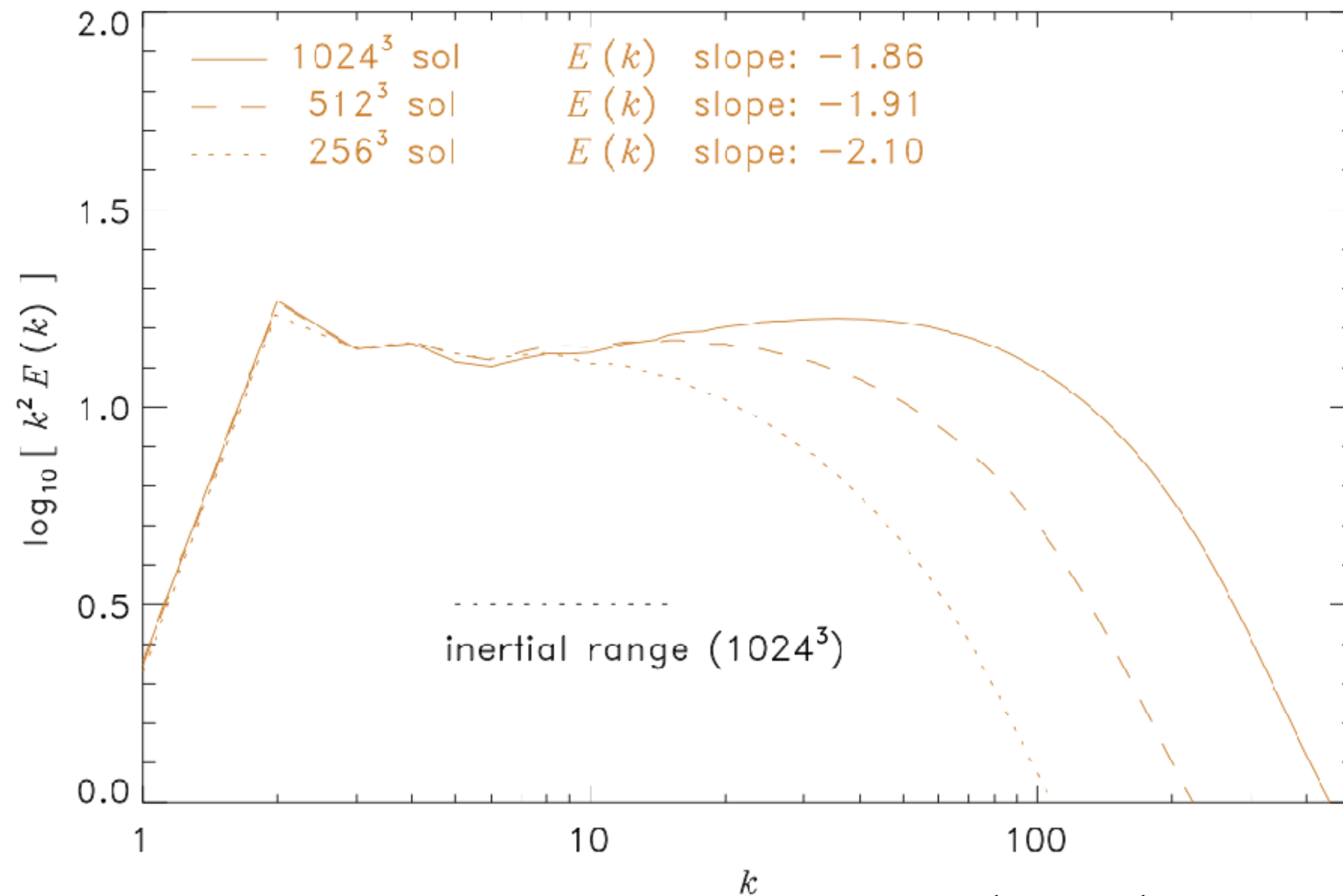


Changes multi-phase structure

$T_{\text{an}+20}$

An Inconvenient Truth

We barely resolve turbulence in our sims



Federrath+2010

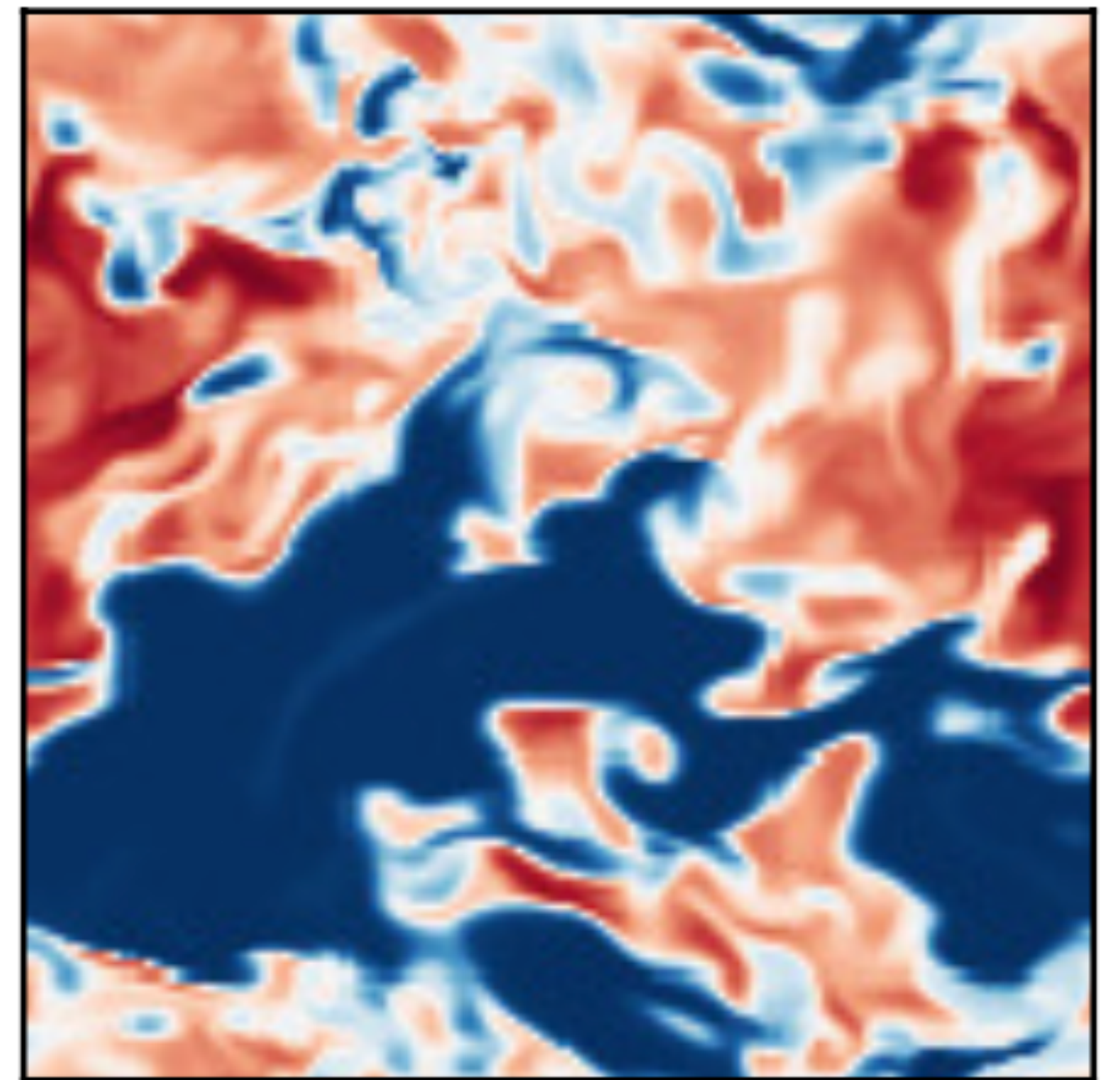


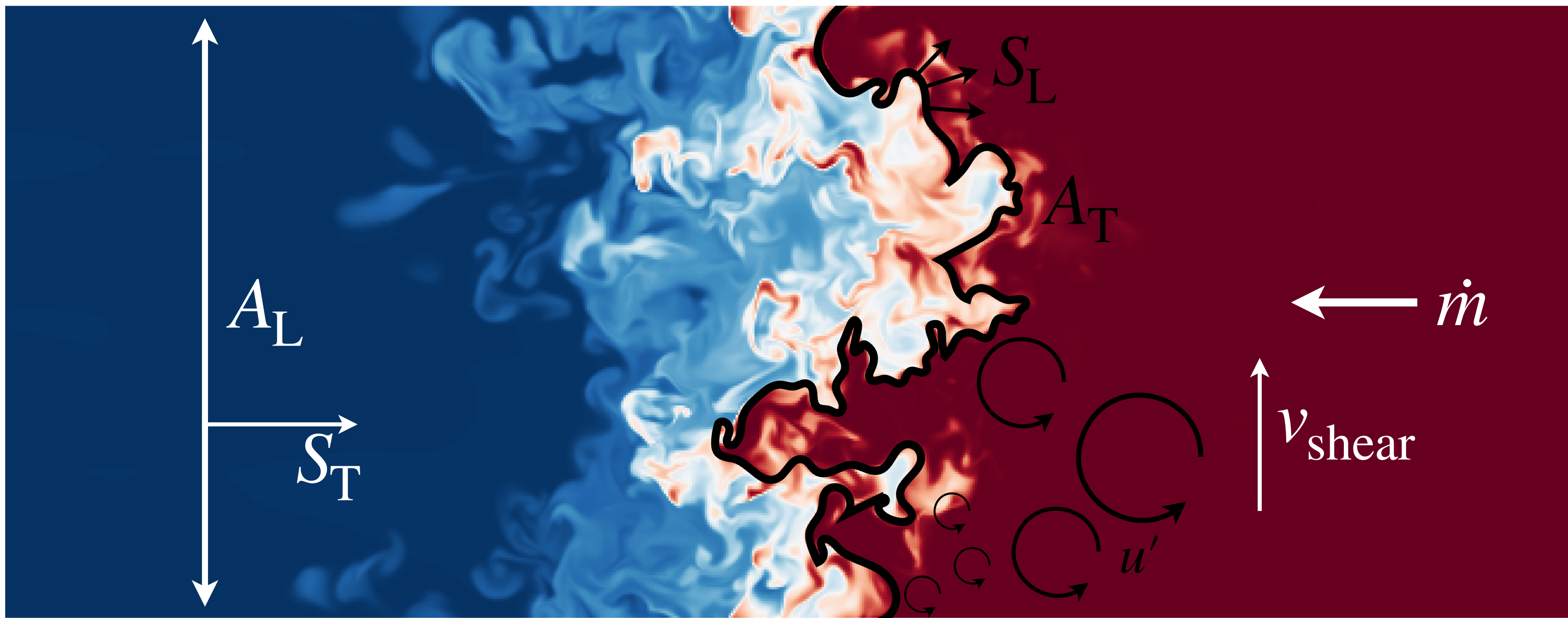
we're stirring honey...

Numerical viscosity extends up to ~ 20-30 times the grid scale

Turbulence & Radiative Cooling

$$f_{\text{cold}} = 0.5$$

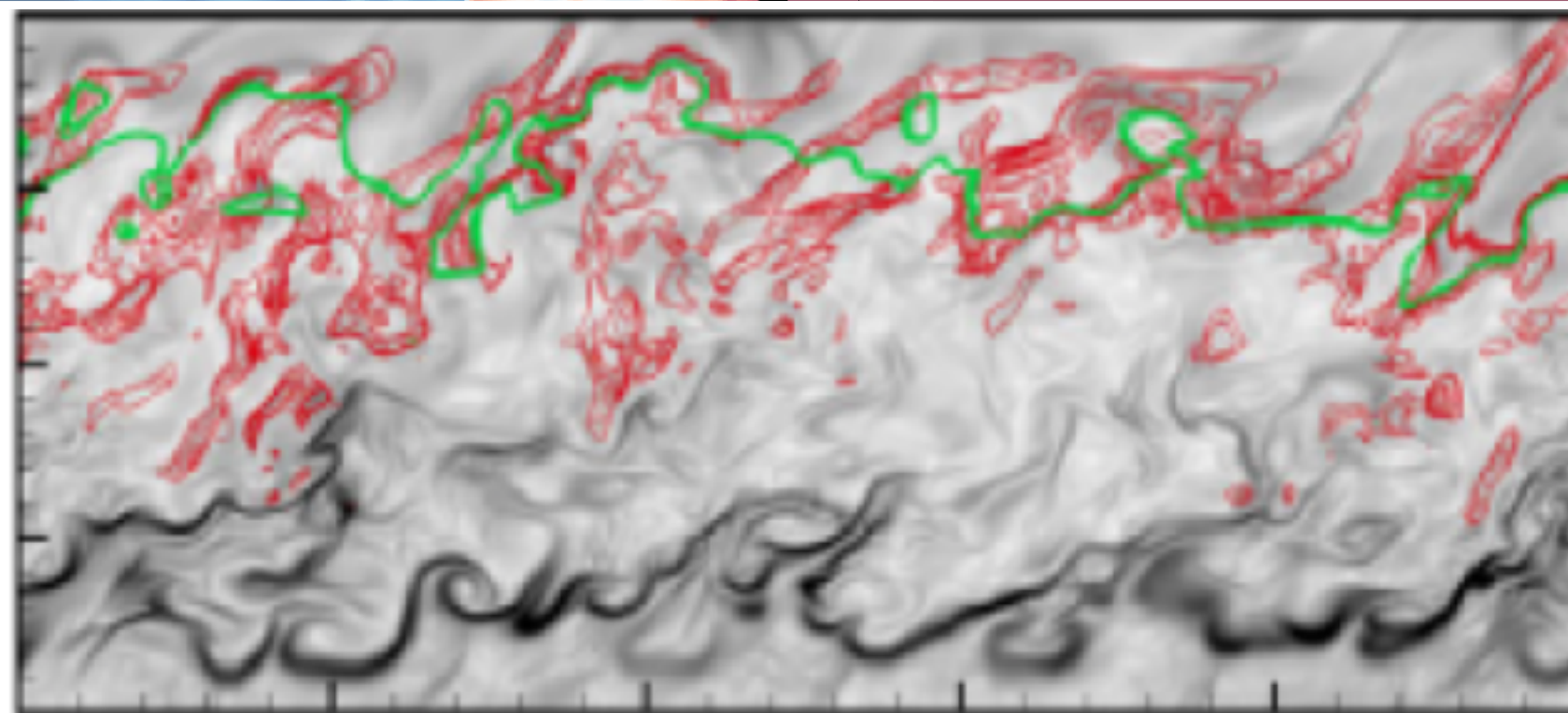
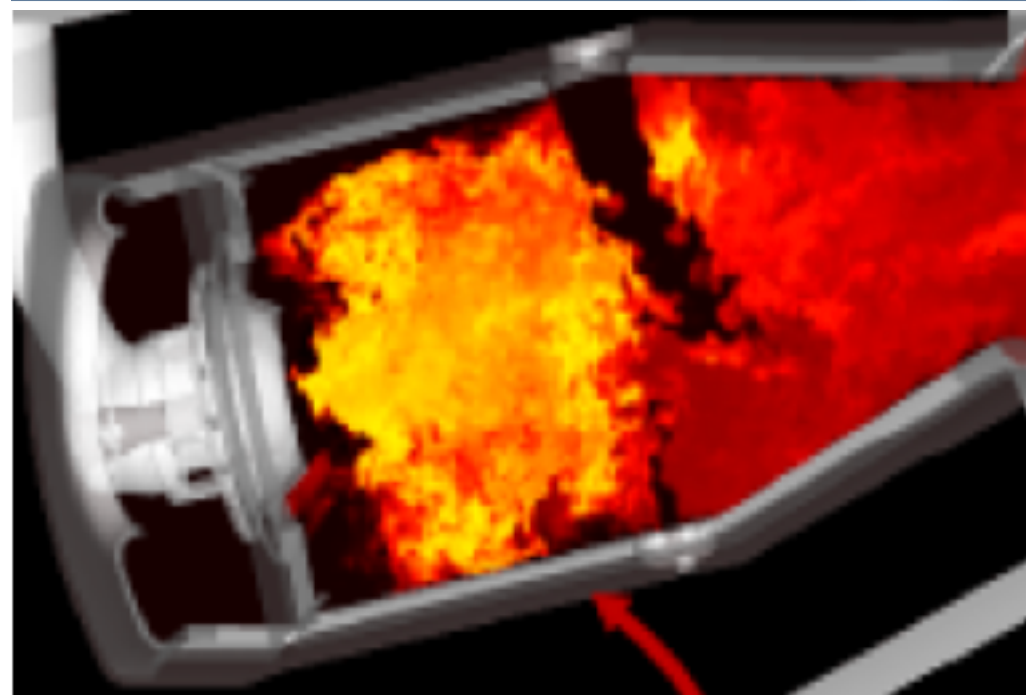




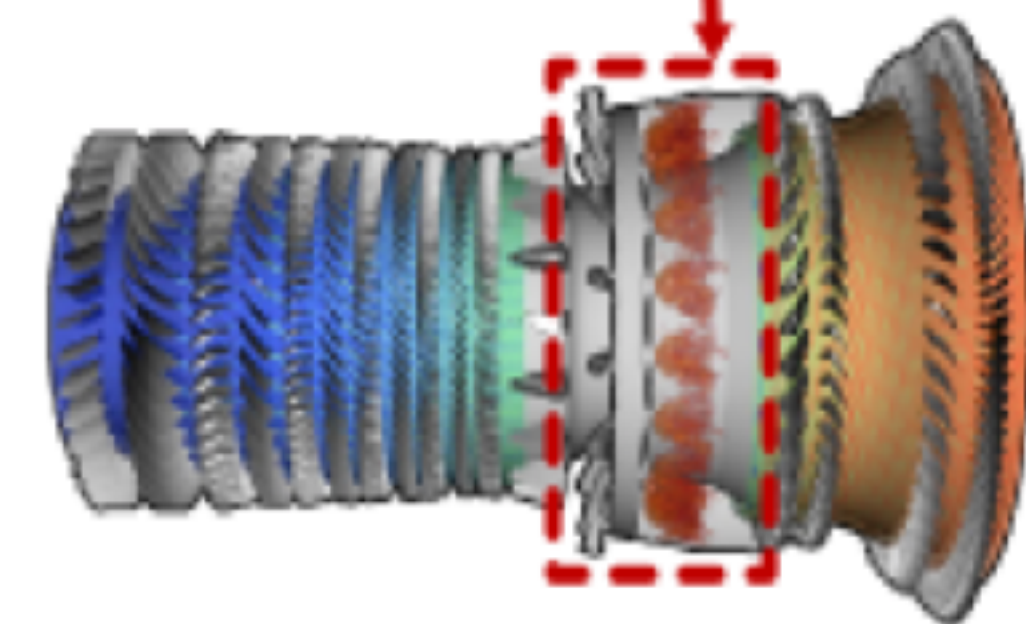
This..



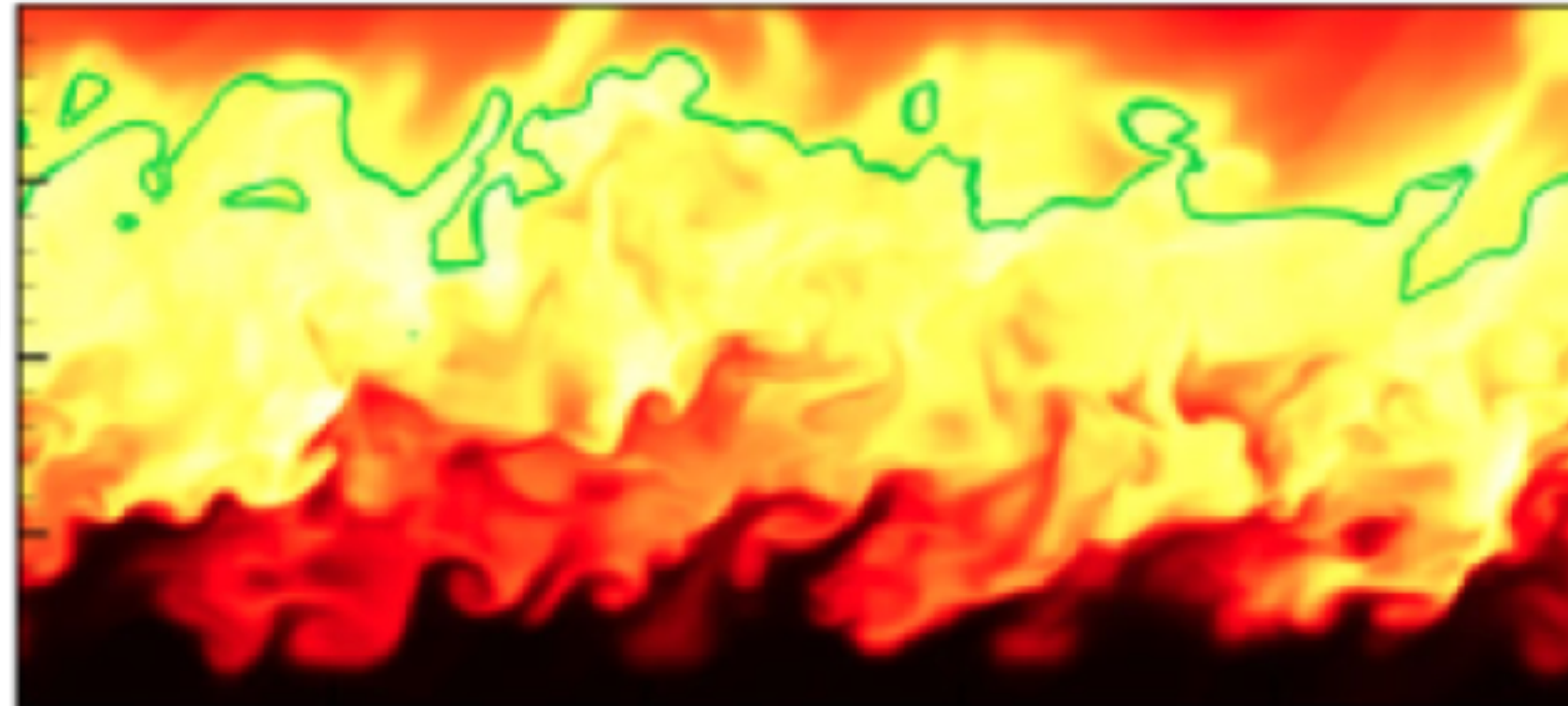
Brent Tan



looks like this



Reacting flow physics



Radiative Mixing Layers: Insights from Turbulent Combustion

Brent Tan^{1*}, S. Peng Oh¹, and Max Gronke^{1,2†}

¹University of California - Santa Barbara, Department of Physics, CA 93106-9530, USA

²Department of Physics & Astronomy, Johns Hopkins University, Bloomberg Center, 3400 N. Charles St., Baltimore, MD 21218, USA

Two highlights

Details of microscopic heat transfer between phases not important



Not sensitive to numerical diffusion

Don't need to resolve Field length



Effective cooling: geometric mean of elastic and inelastic collision times $\tilde{\tau}_{\text{cool}} \sim \sqrt{\frac{L}{u'} t_{\text{cool}}}$

Just like conduction

or radiative transfer

$$\lambda_F \sim \sqrt{\frac{\kappa T}{n^2 \Lambda(T)}} \sim \sqrt{\lambda_e v_e t_{\text{cool}}}$$

$$\tau_* \sim \sqrt{\tau_a \tau_s}$$

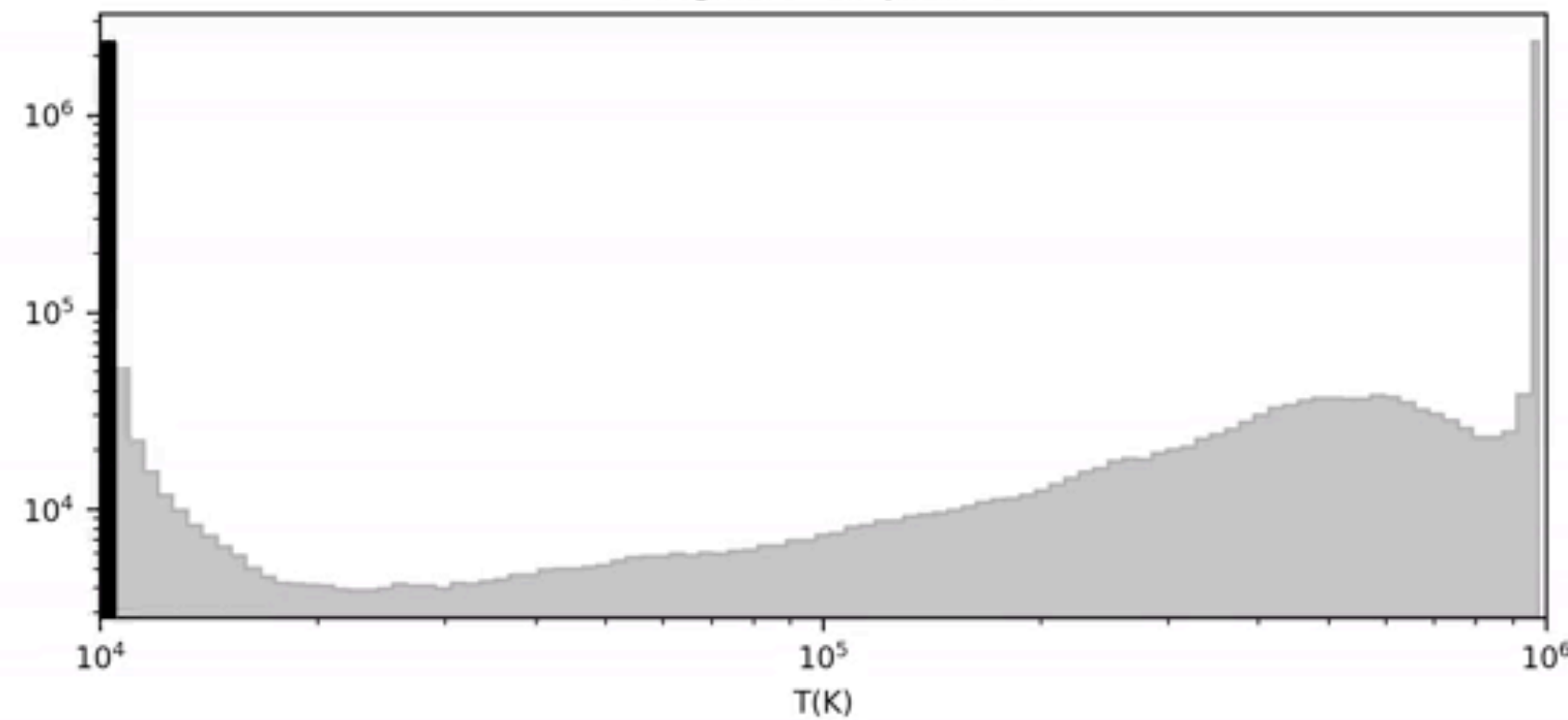
You want a subgrid model...

Analytic model for temperature PDFs

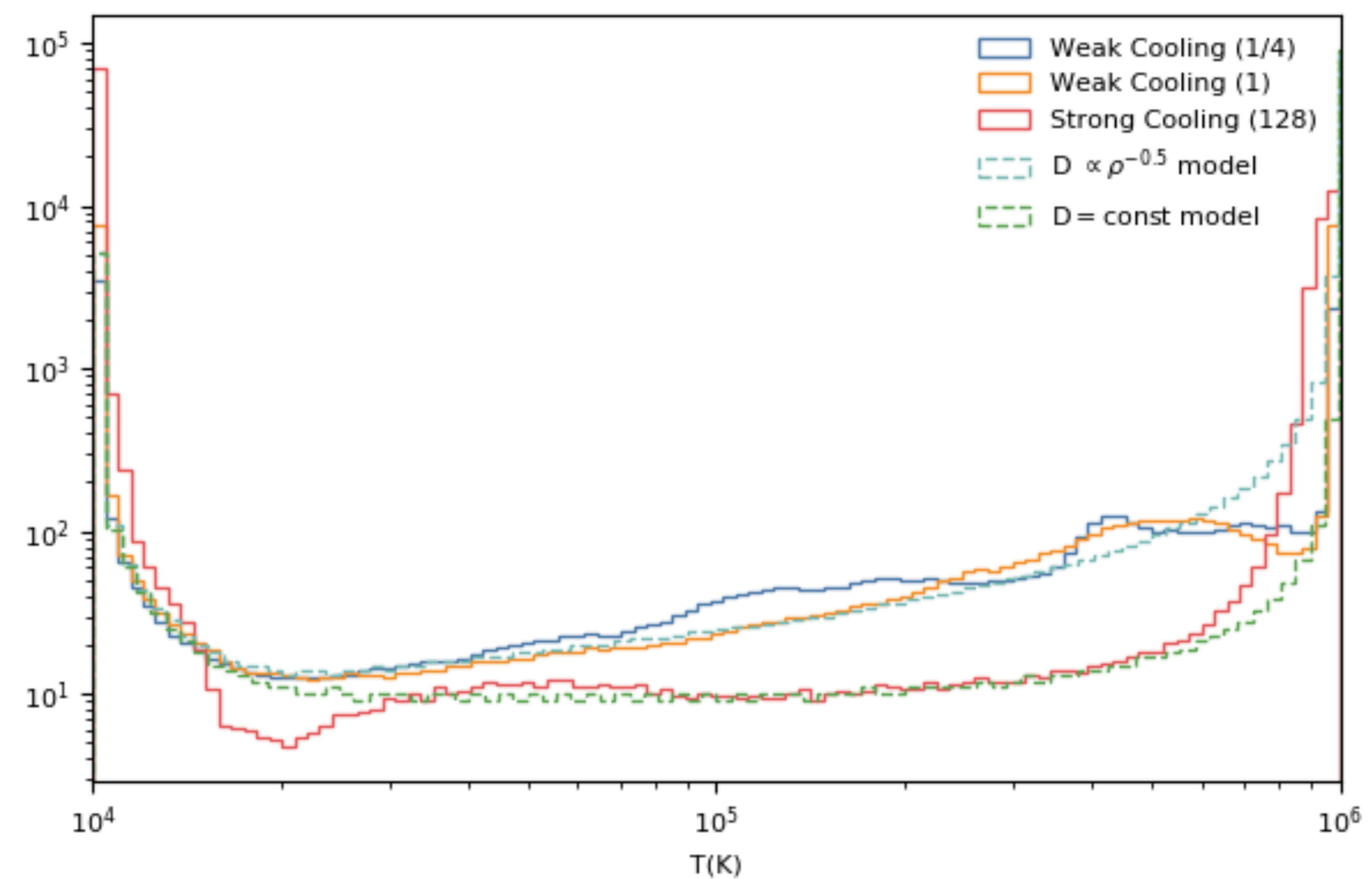
Input: cooling time, turbulent diffusivity



Histogram of Temperatures



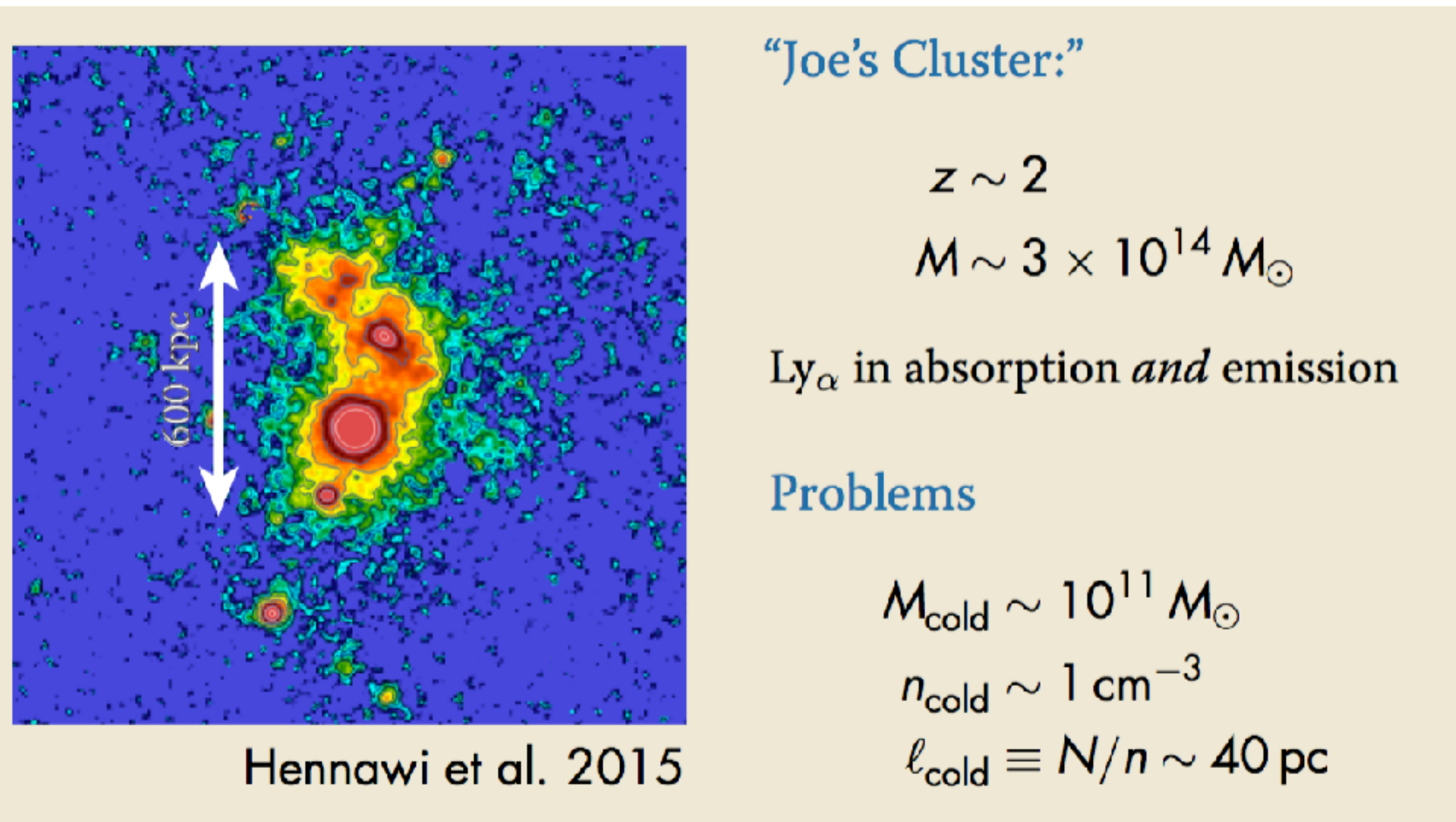
Histogram of Temperatures



By density or emission weighting, get absorption/emission line ratios from analytic model which matches sims

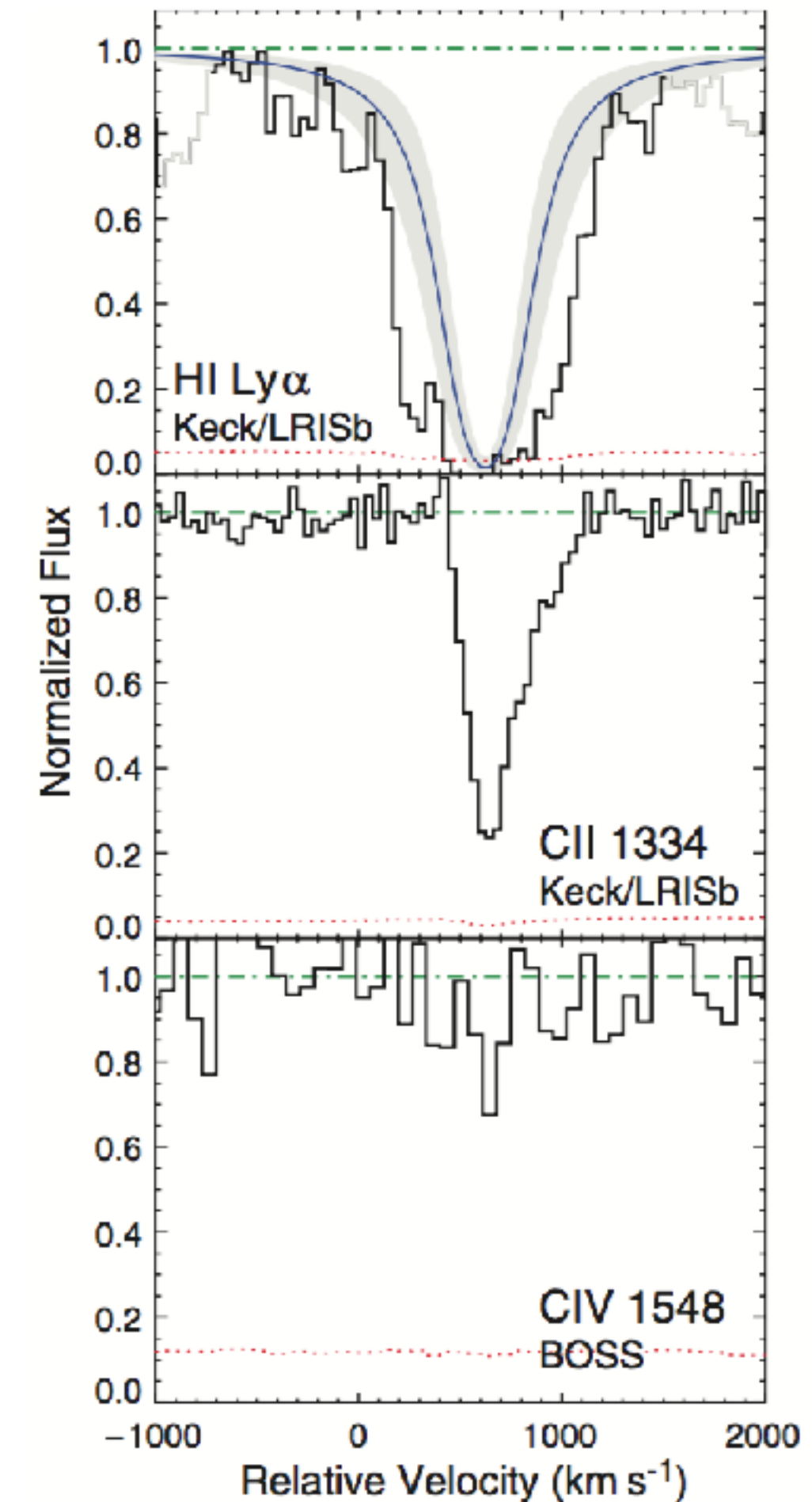
'Foggy' cold gas is **observationally** driven

Theorists are just catching up...



Highly
suprathemal
line widths

$$\mathcal{M} \sim 100!$$



Large area covering fraction $f_A \sim 1$
despite low volume filling fraction $f_V \sim 10^{-4}$
Need a very thin shell which surrounds the halo??

All makes sense if have a fog: intersect **many** cloudlets along line of sight

OK, you still want to talk sizes

Mike vs. Max

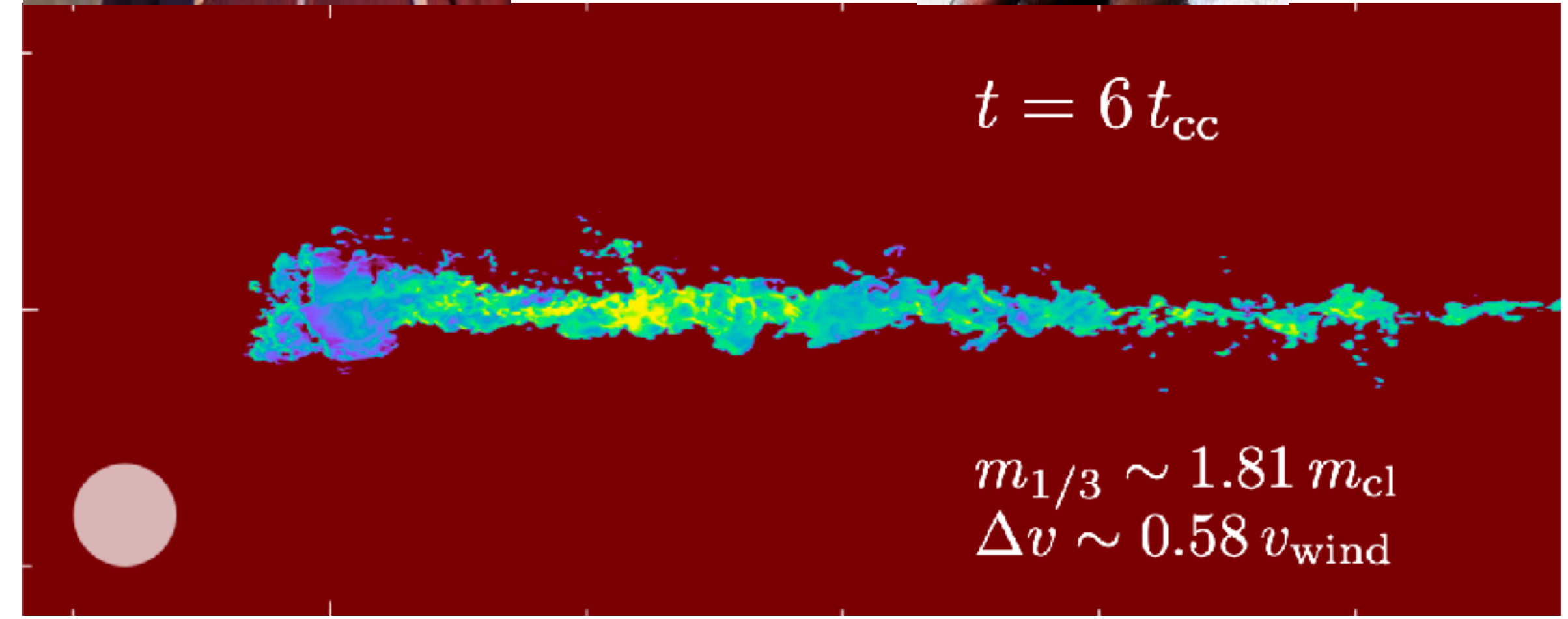
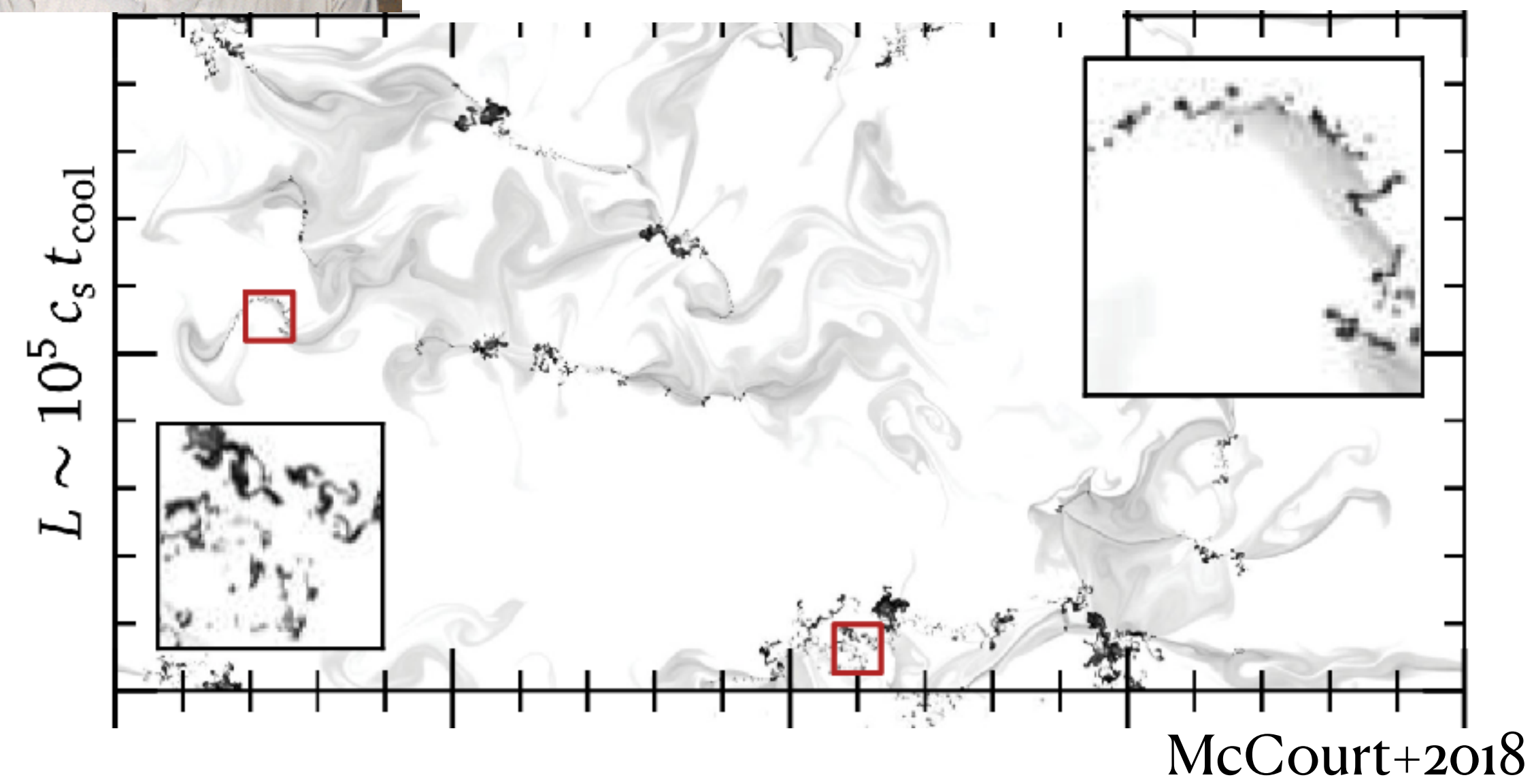


small is beautiful

Mike McCourt



Max Gronke



Gronke & Oh 2018

Typical size $\lambda_{Mc} \sim c_s t_{cool}(T \sim 10^4 K)$

Trash talk: 'Big clouds will just shatter'

Minimum size $\lambda_G \sim c_s t_{cool,mix}(T \sim 10^5 K) \gg \lambda_{Mc}$

Trash talk: 'Mist will just dissolve'

Who is right?

Both?



Instead of this...

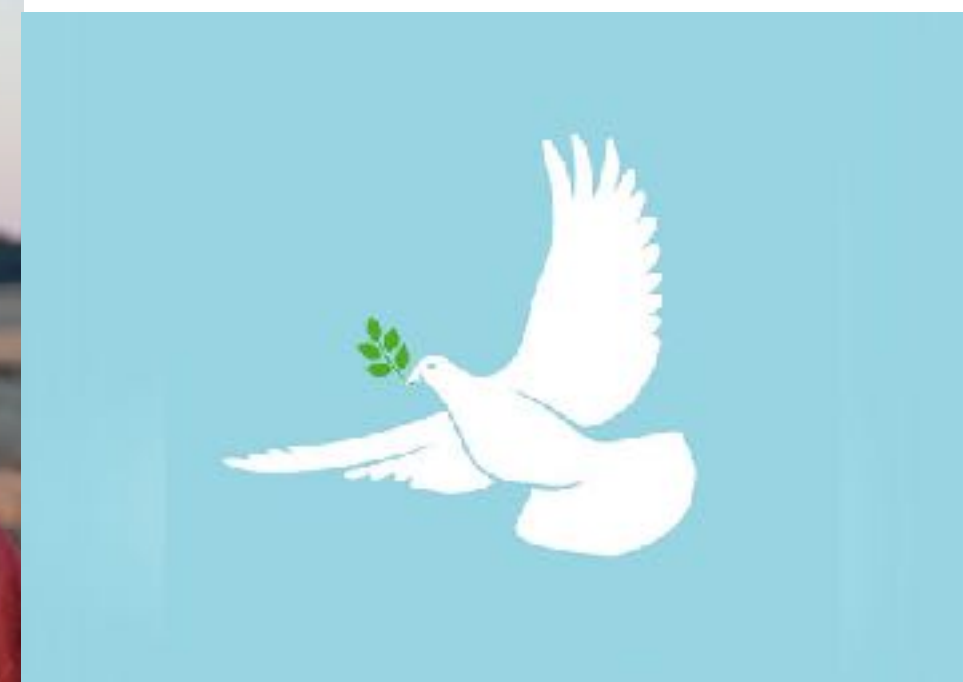
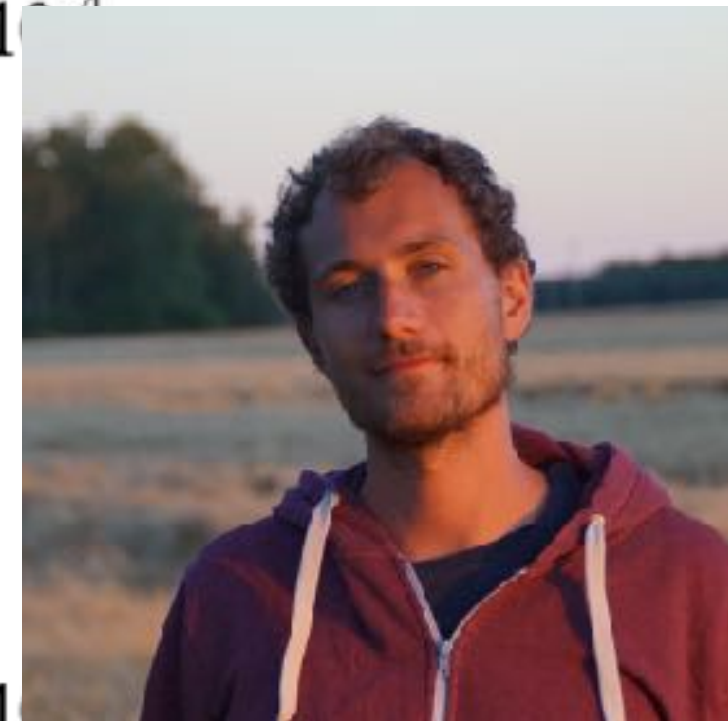
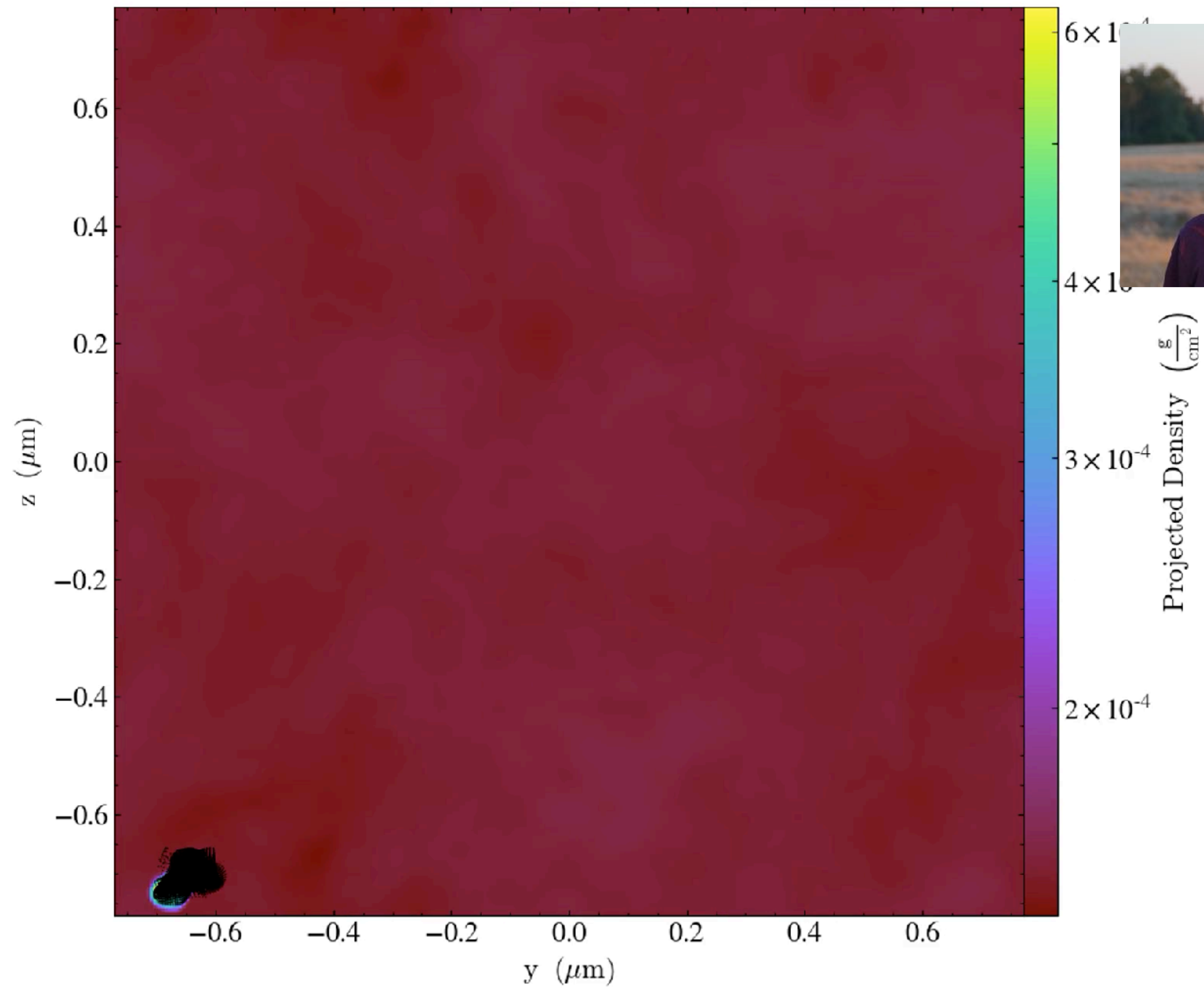


Laminar Flow

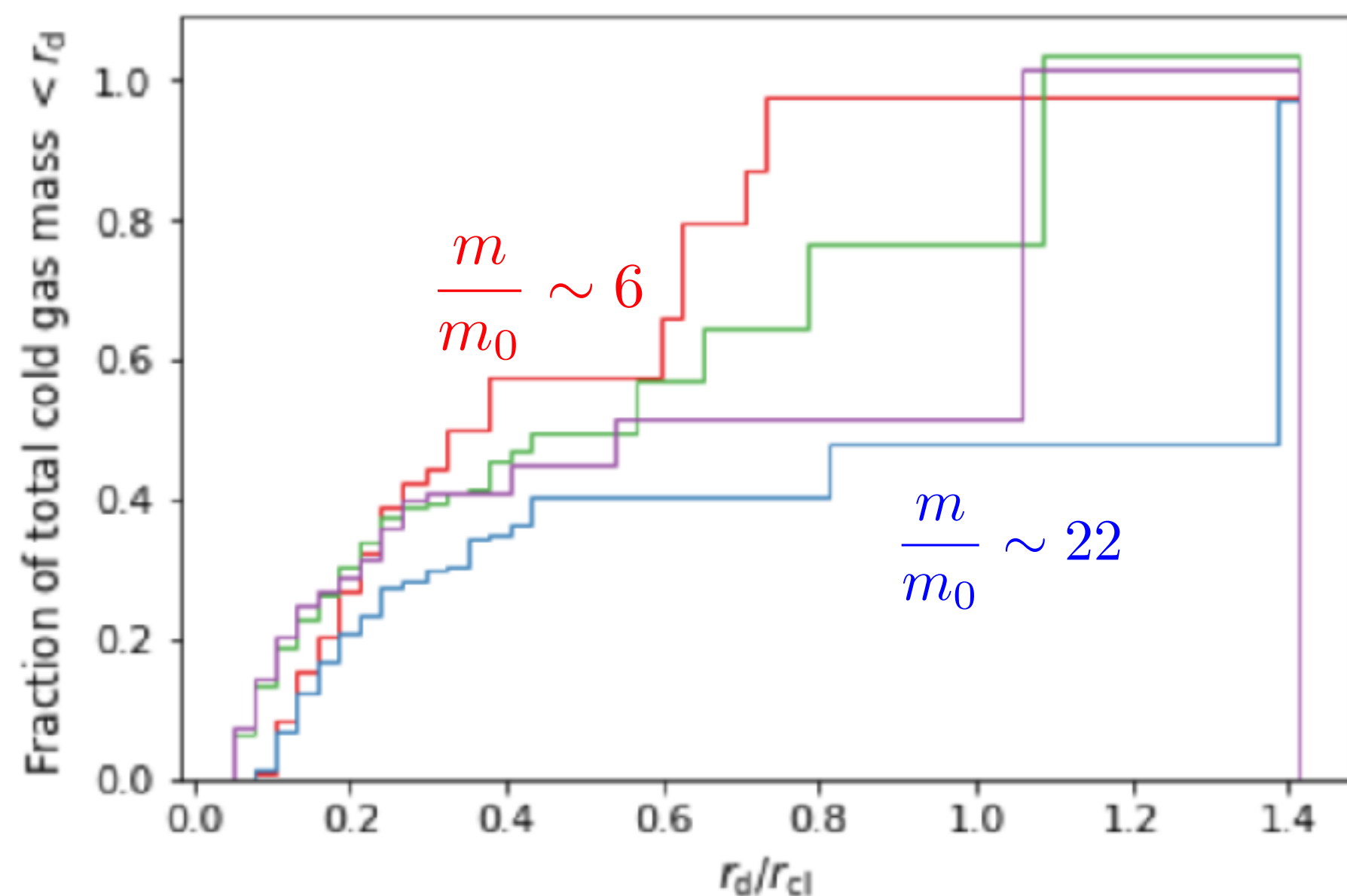
Do this.



Turbulent Flow



Lots of cool features

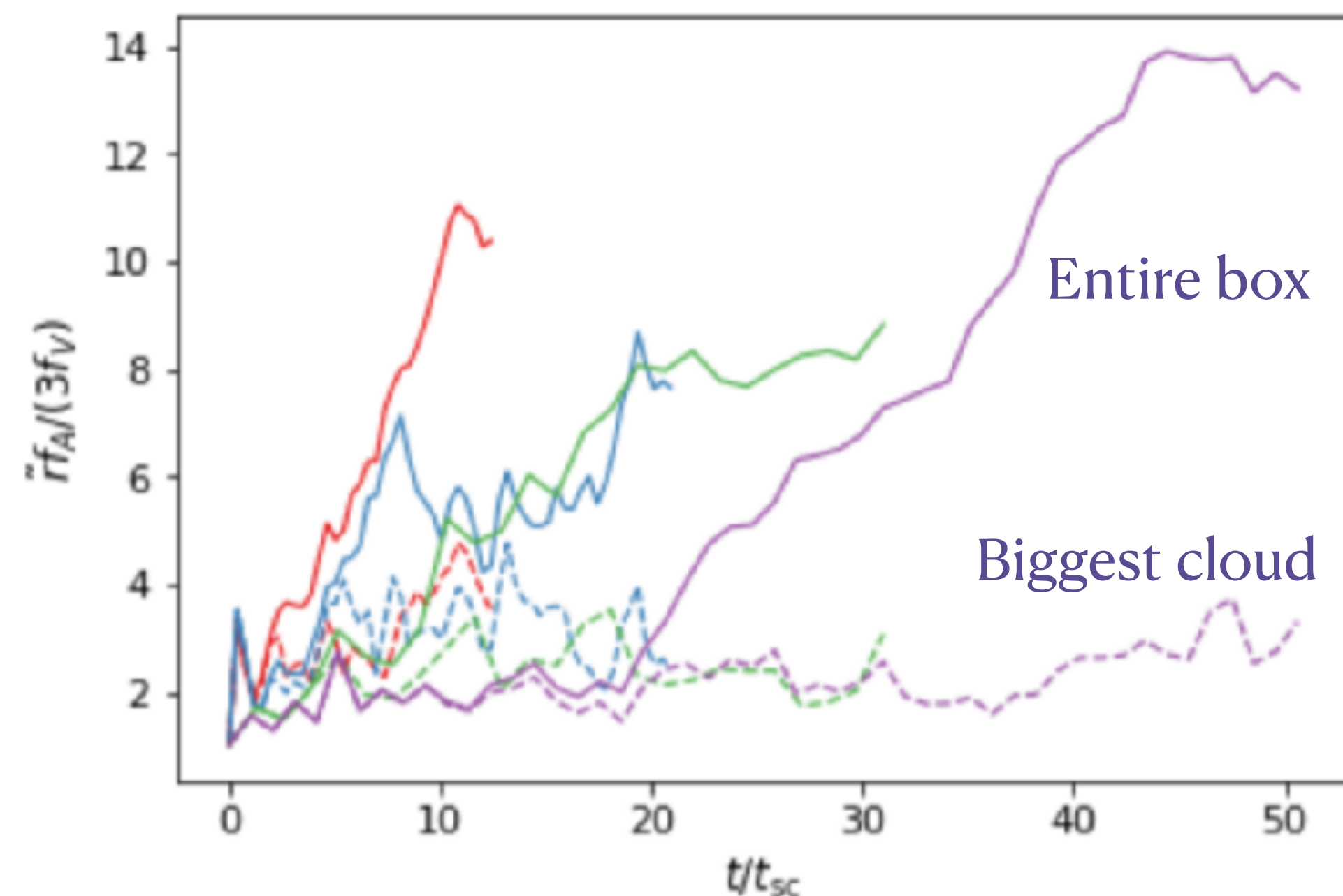


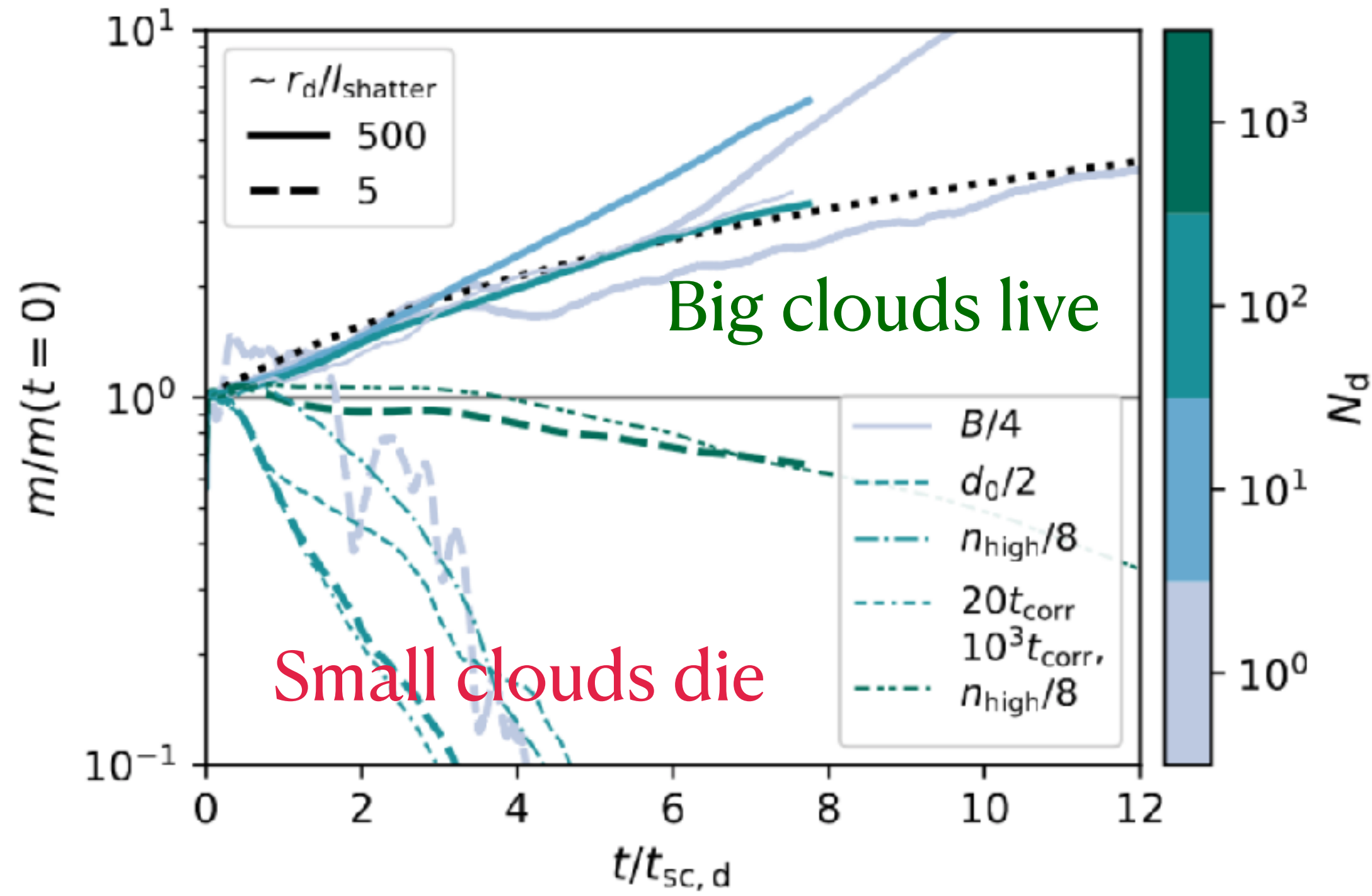
**Mass broadly distributed
across range of scales**

($\sim 1/2$ in a big cloud, $\sim 1/2$ in much smaller clouds)

But area dominated by small cloudlets

Along a random line of sight, you are *much* more likely to pierce the 'fog' than a 'cloud'





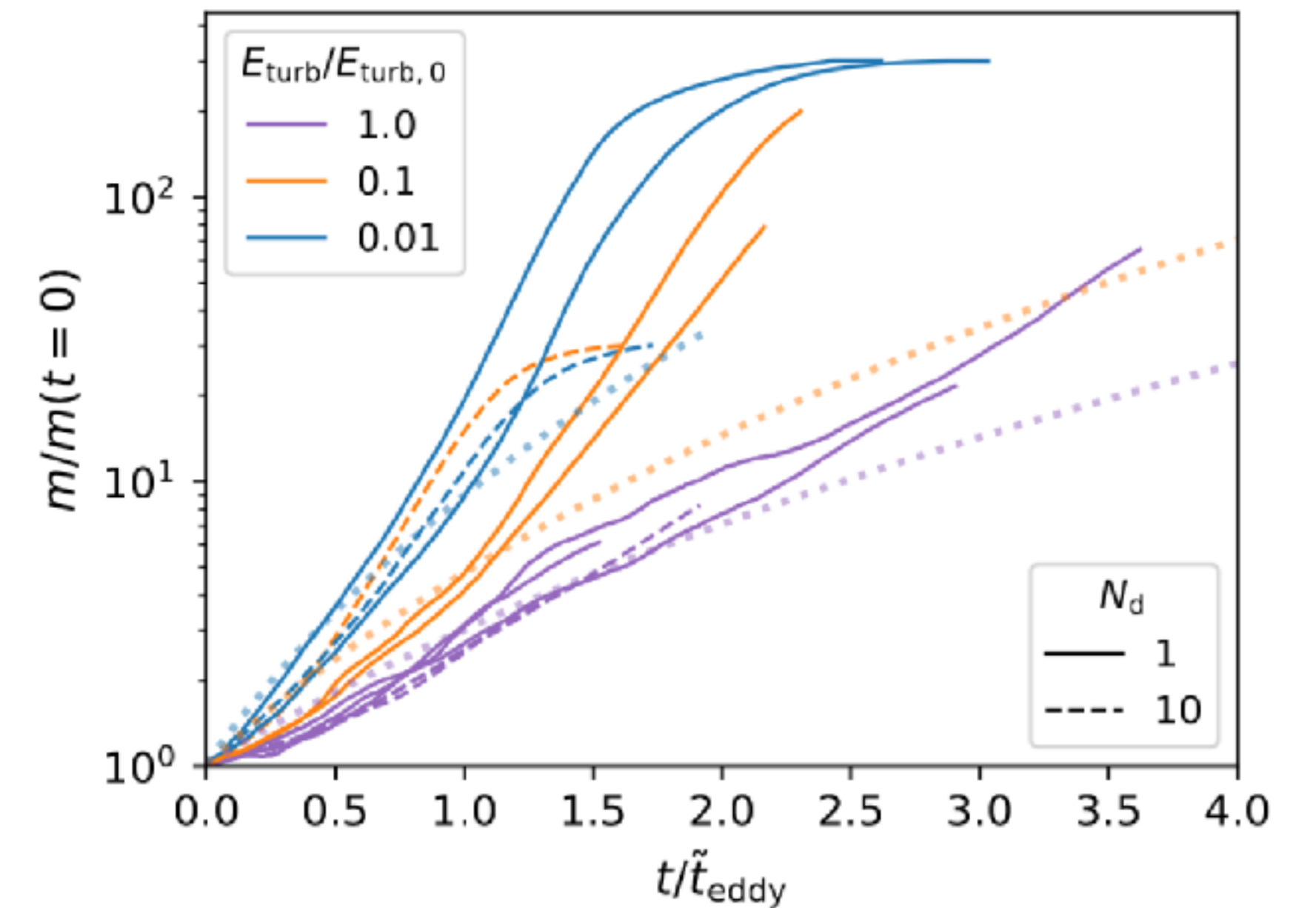
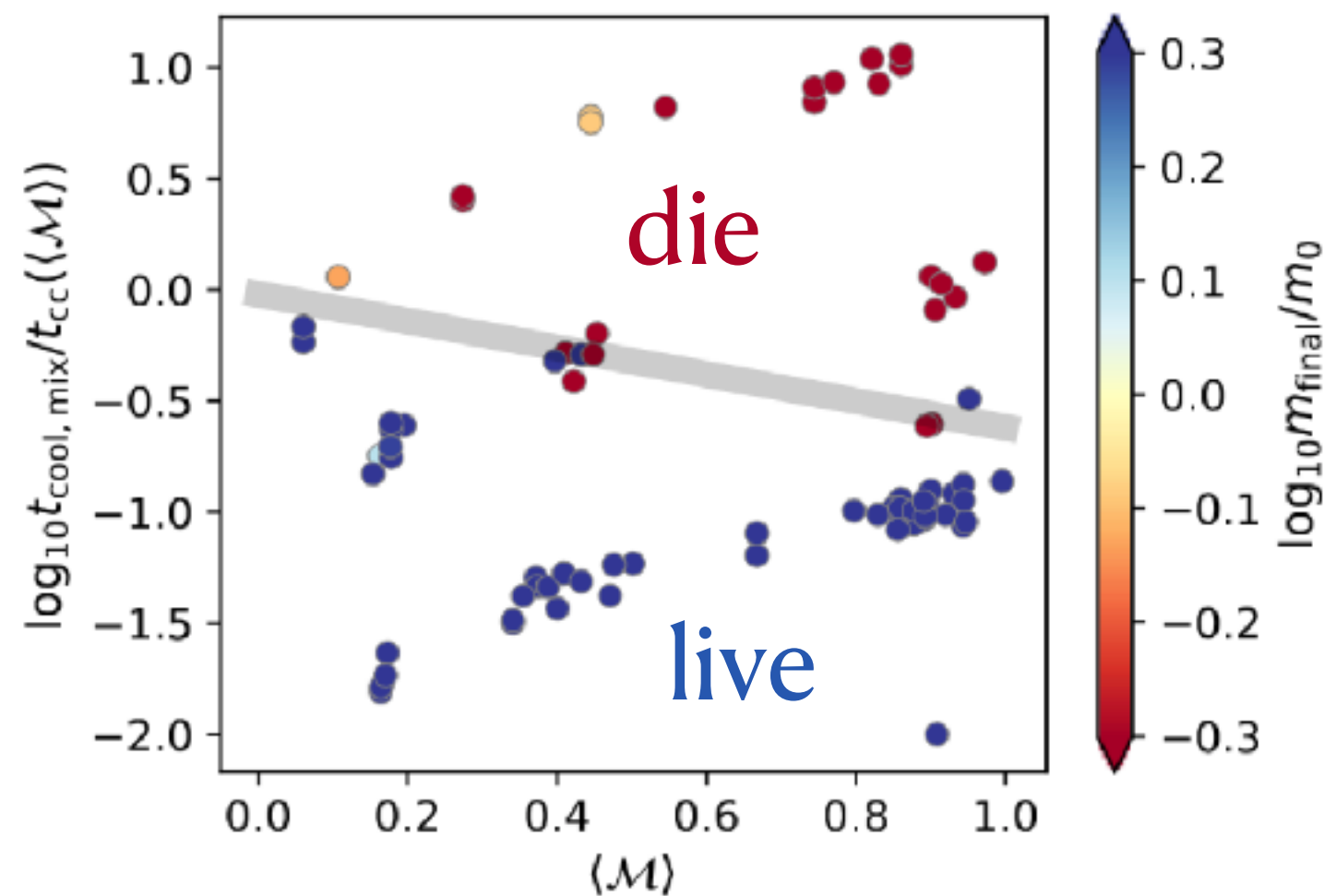
Large clouds are *required* for survival

Equivalent mass in small droplets doesn't work
Coagulation too slow to ensure survival

Growth roughly fits analytic models

but transition to exponential growth when
small clouds dominate area

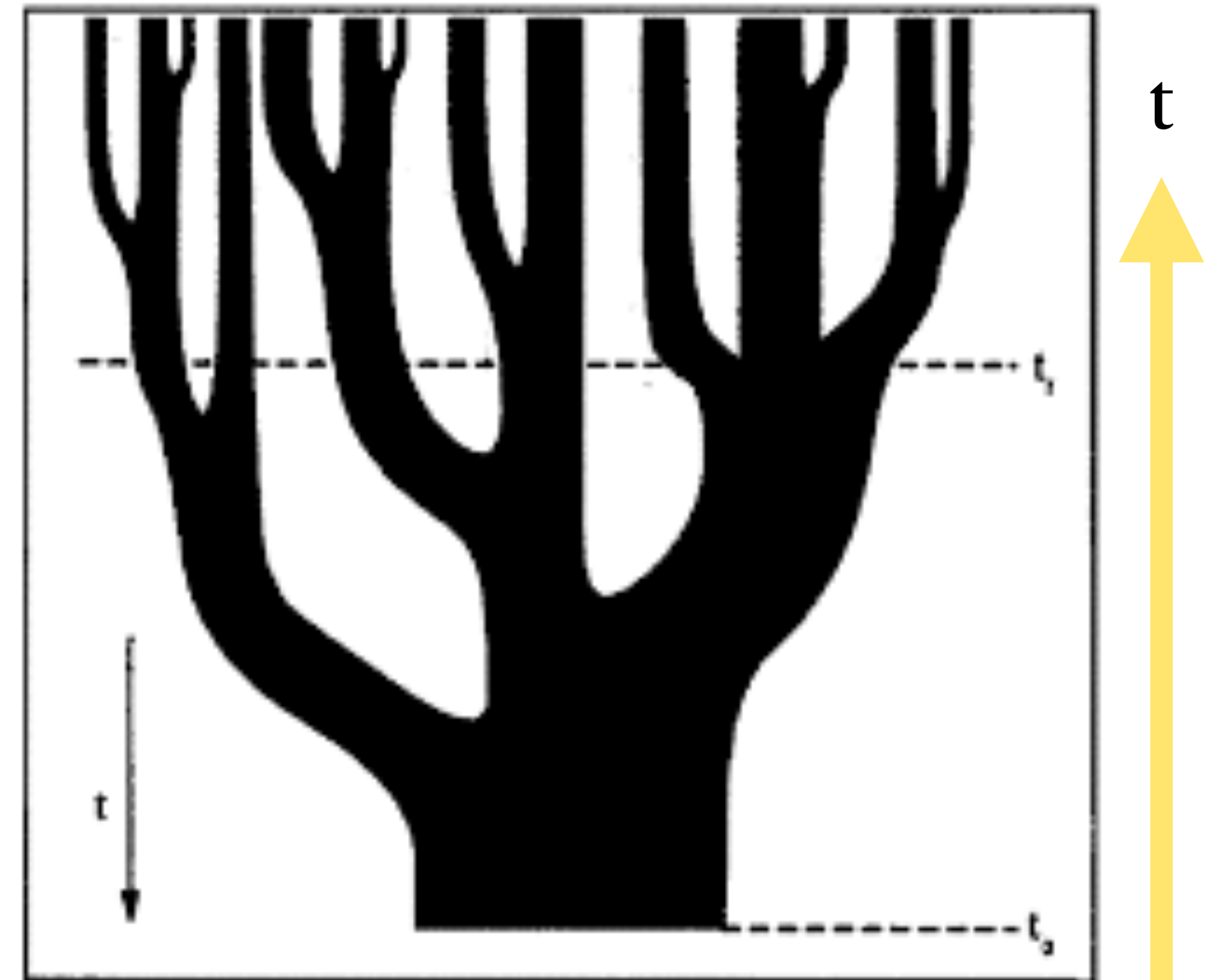
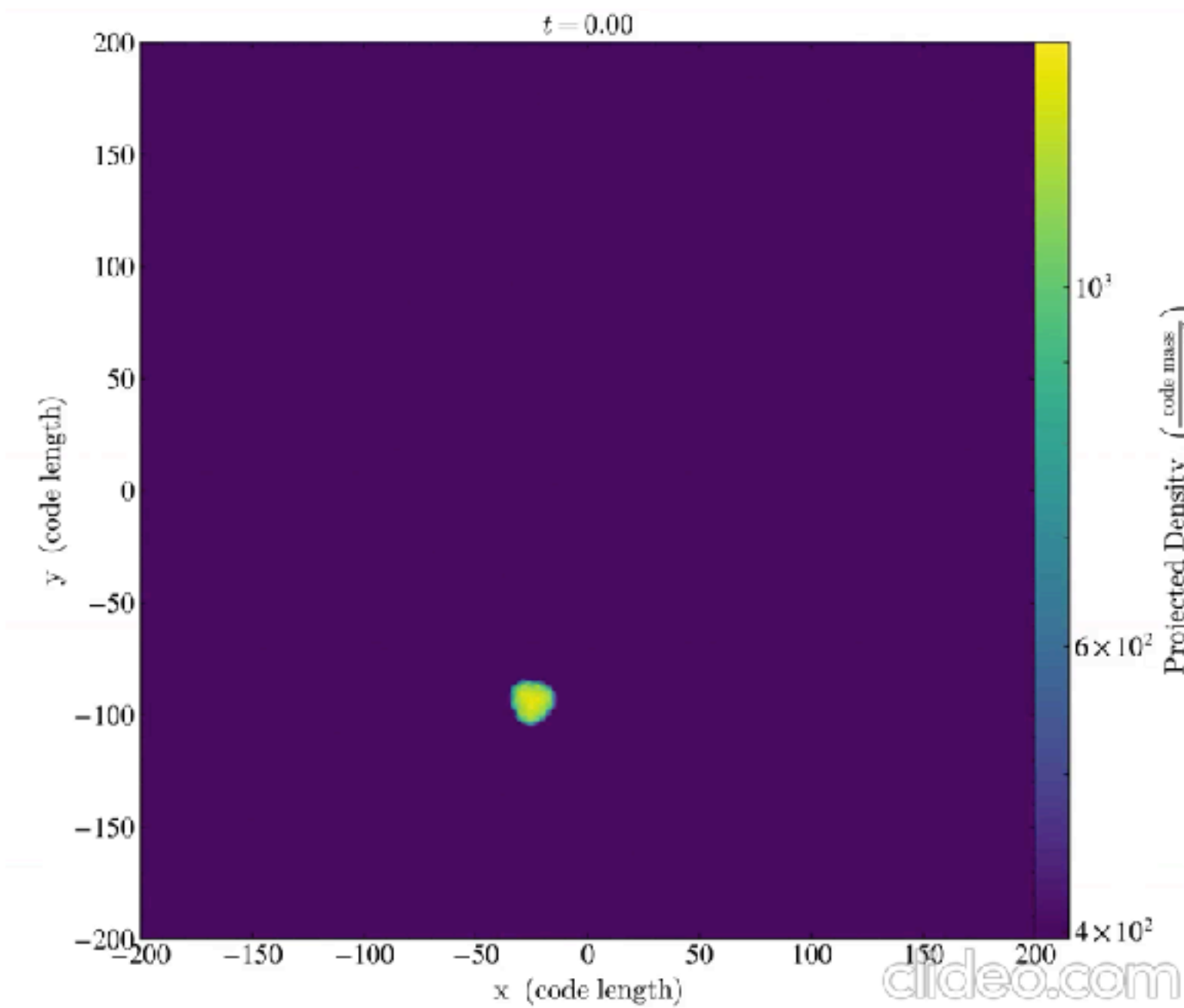
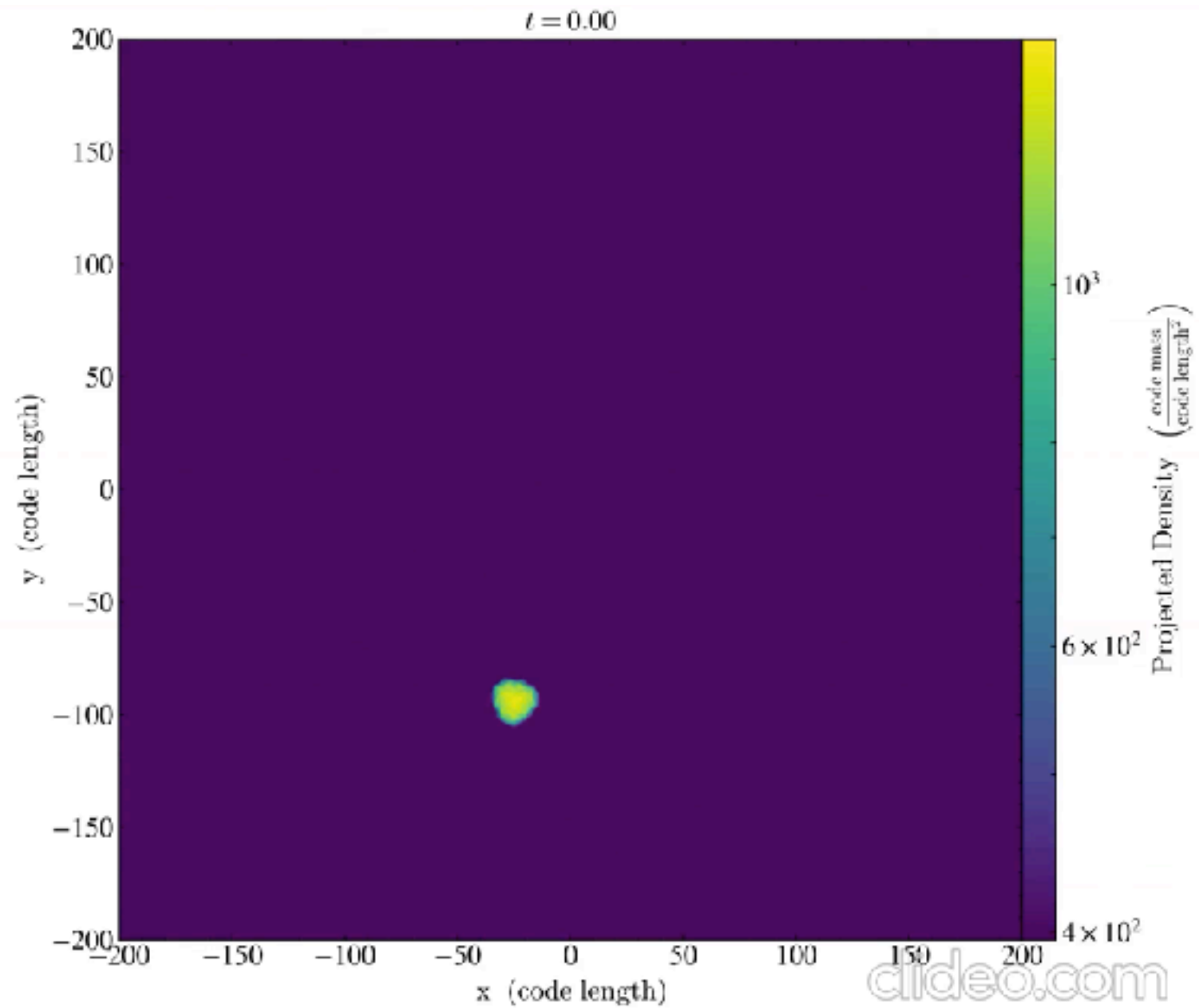
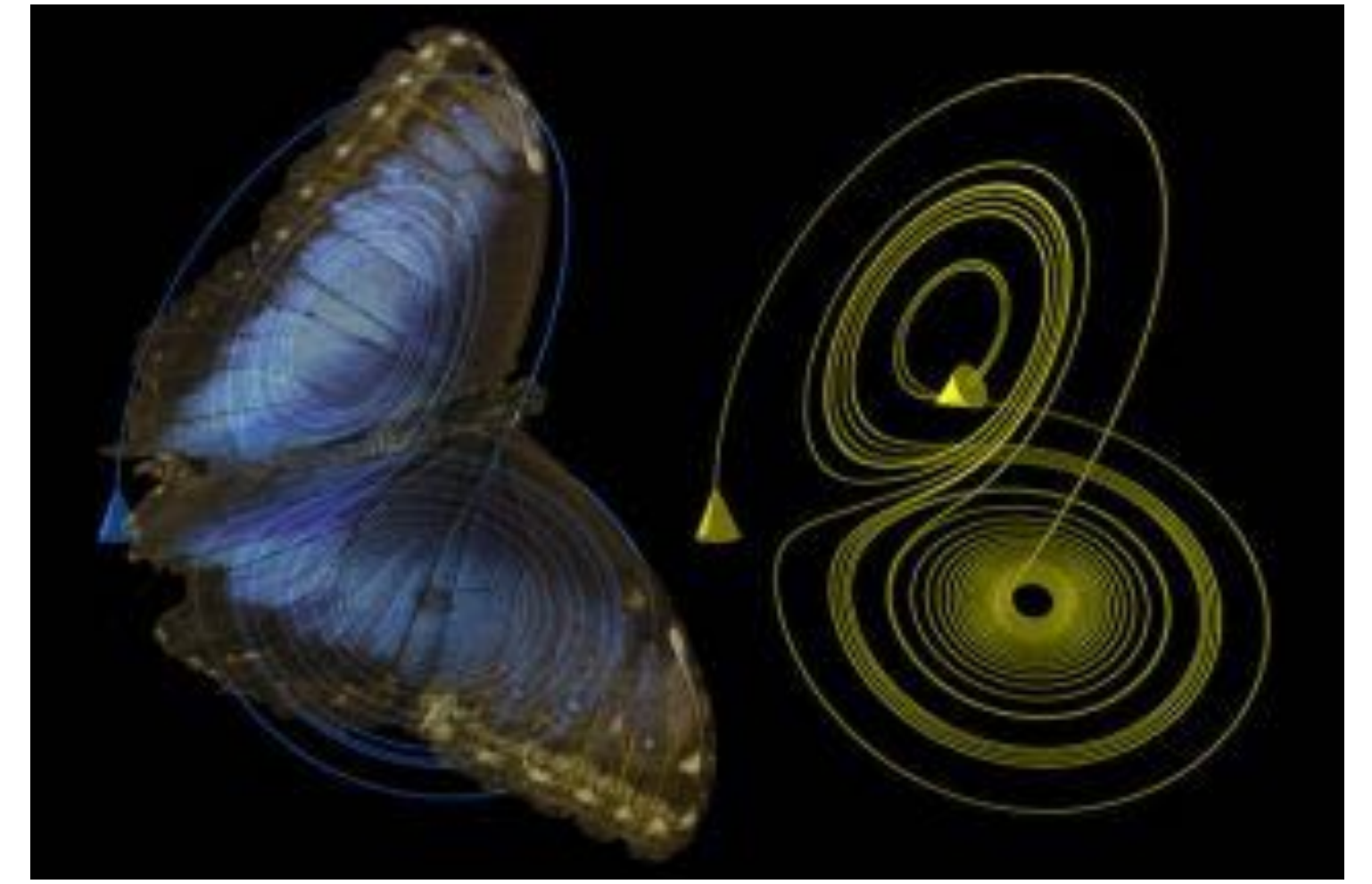
**Mass threshold
similar to
wind problem
(Gronke & Oh
2018)**



Outcome is highly stochastic

... a lot of 'cosmic variance'

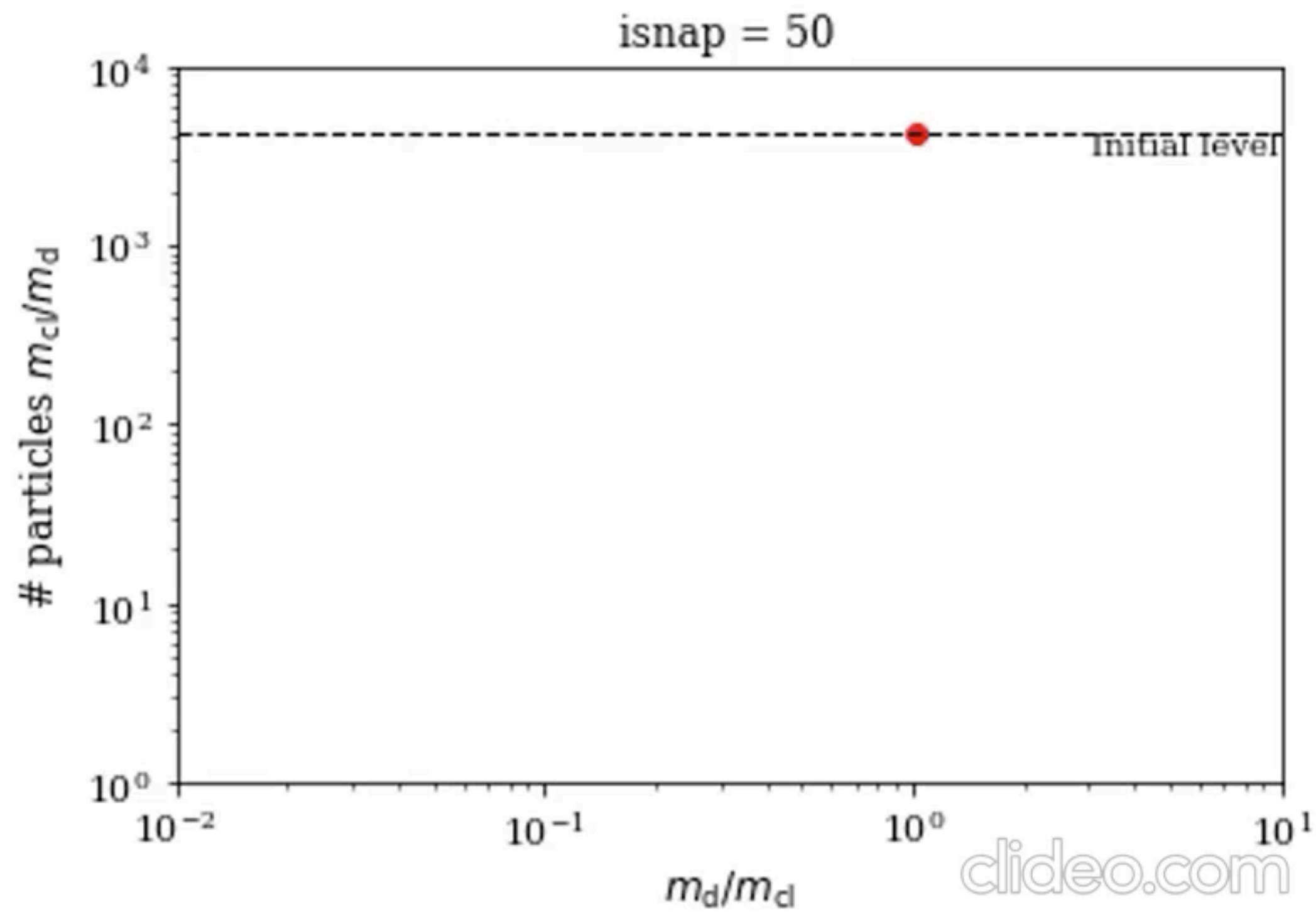
Need to ensemble average!



Monte-Carlo approach ('shattering' tree) probably useful

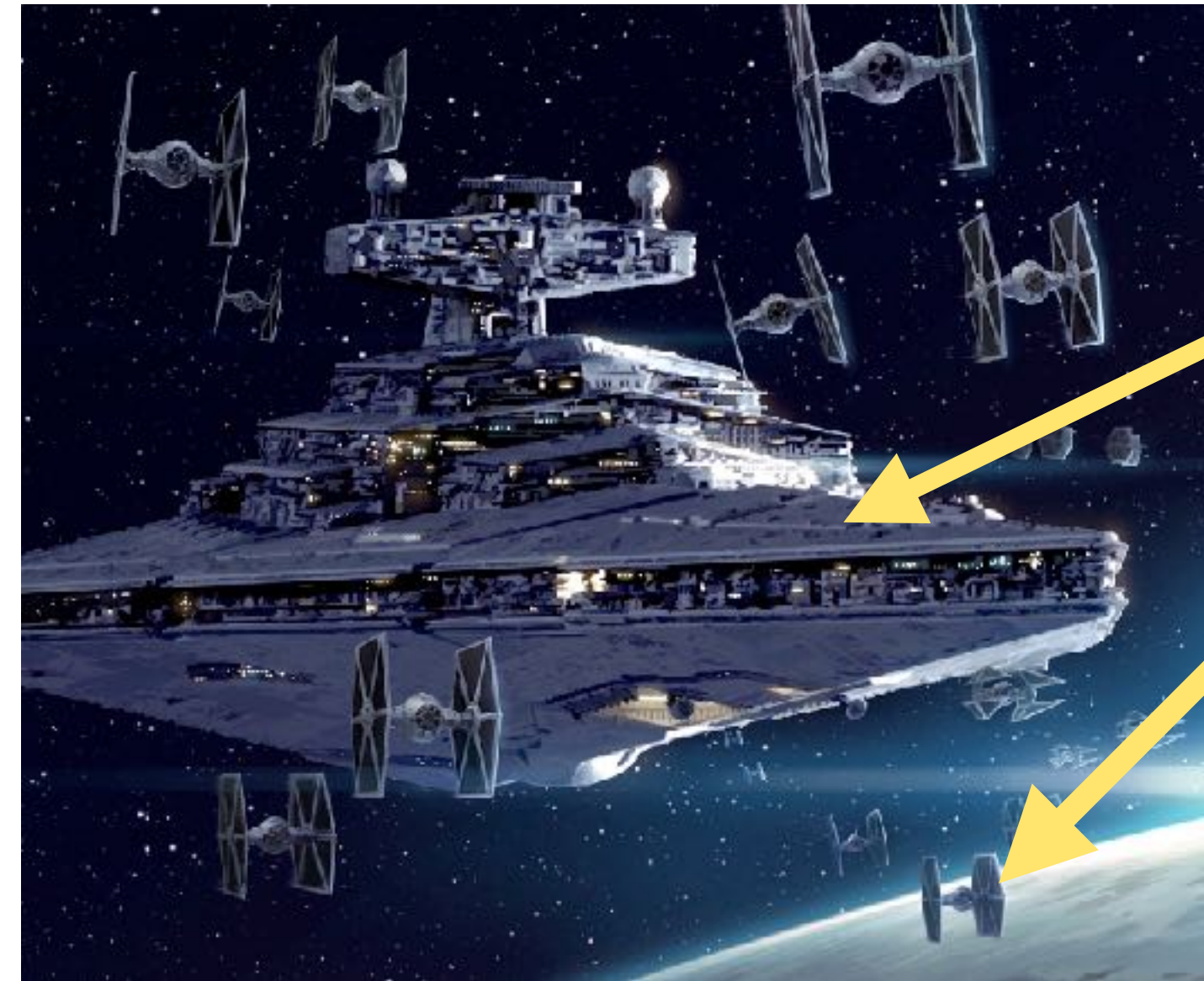
Invert merger tree

The Nerd Version



Continuous growth and breakup
on all scales

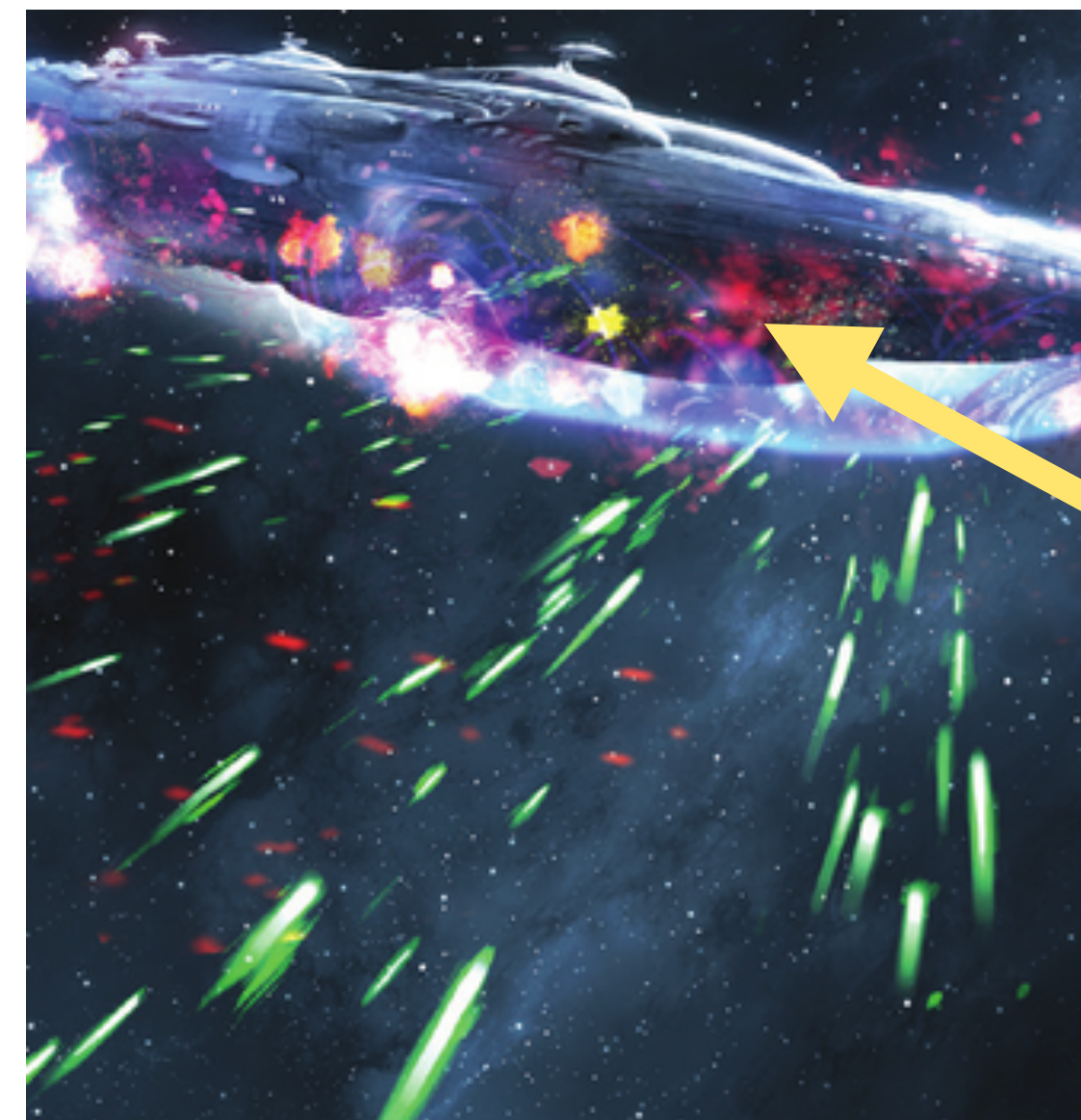
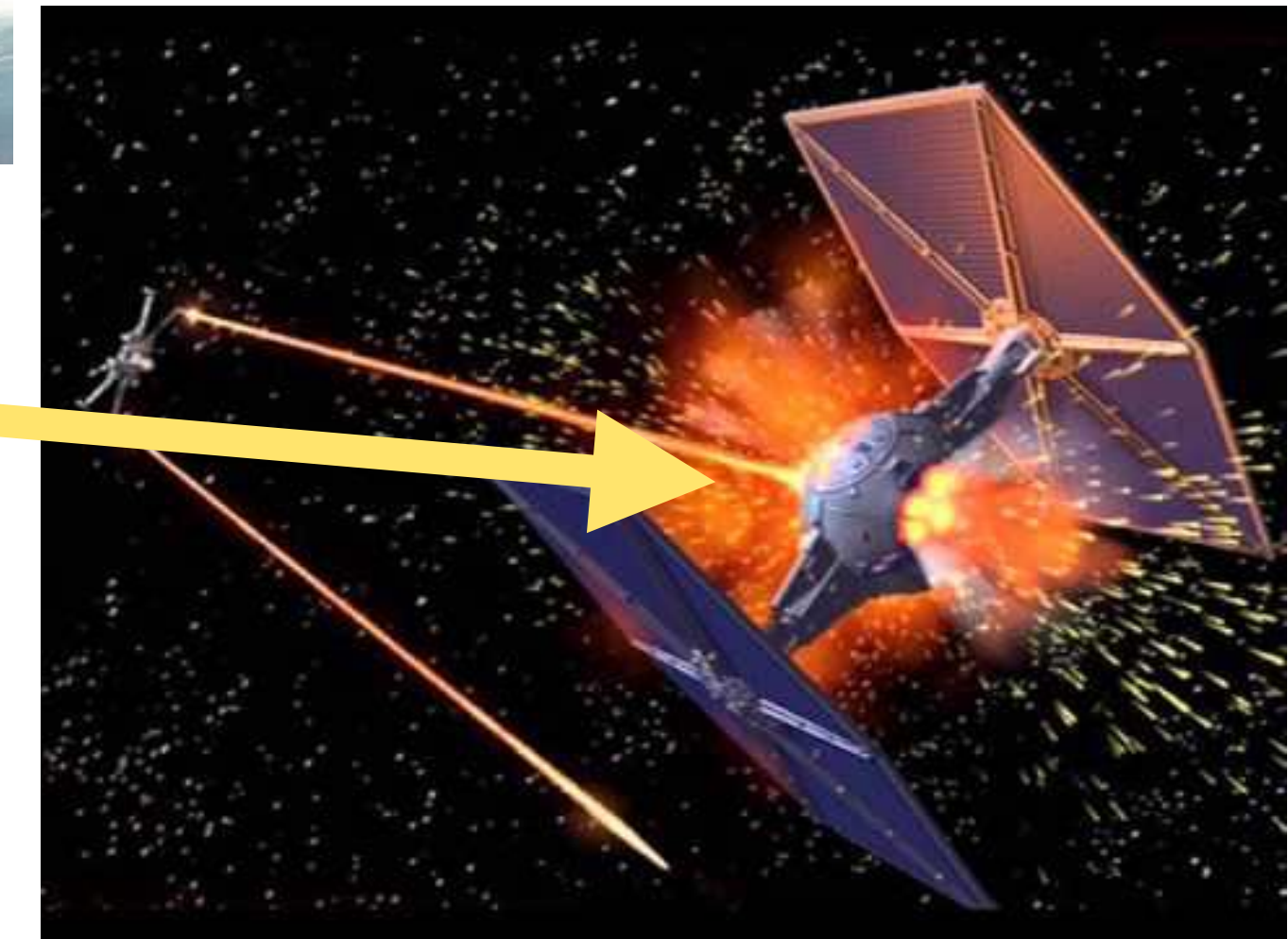
The Hollywood Version



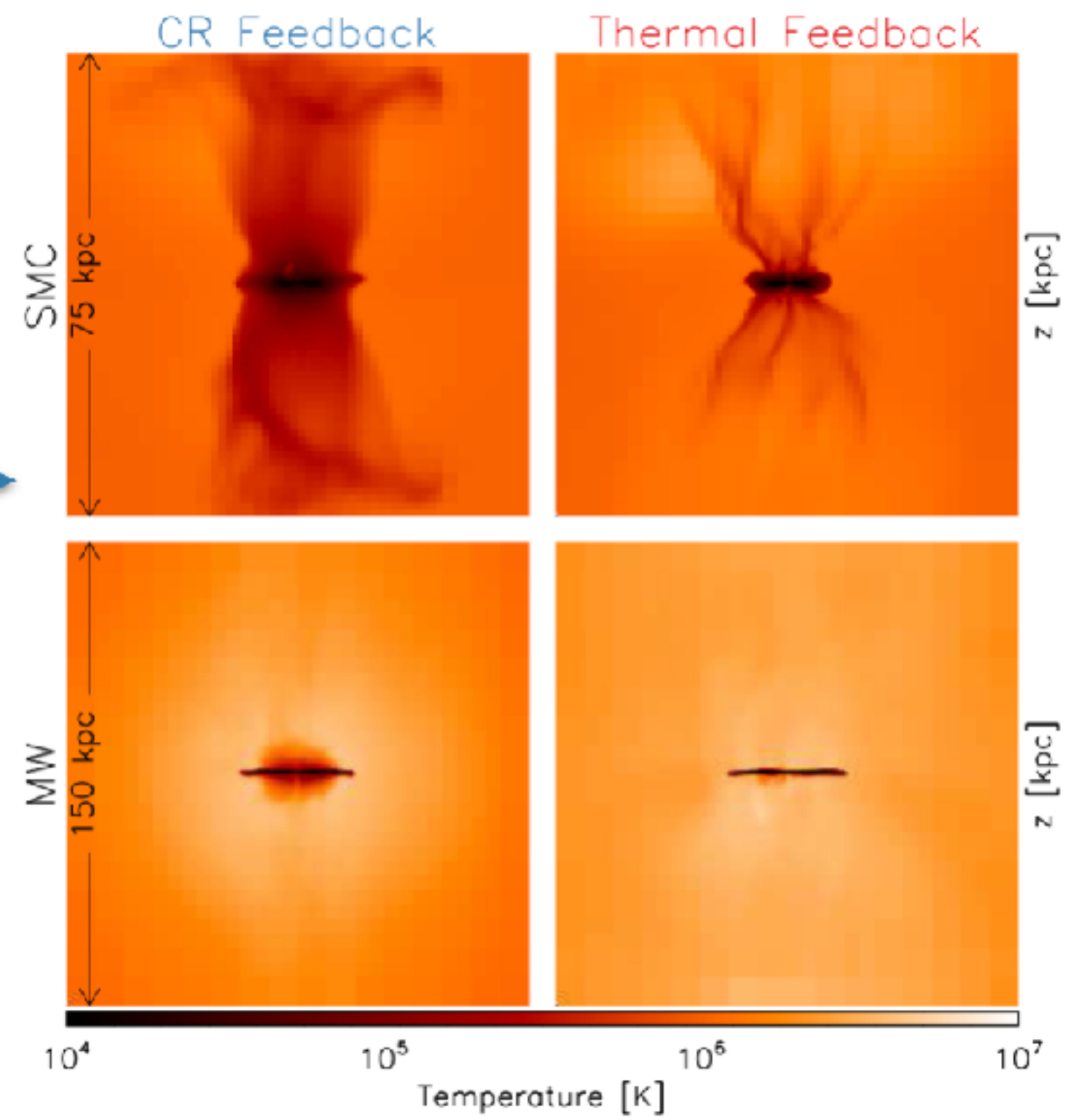
The big cloud...

..launches the droplets

which do battle and get
blown up

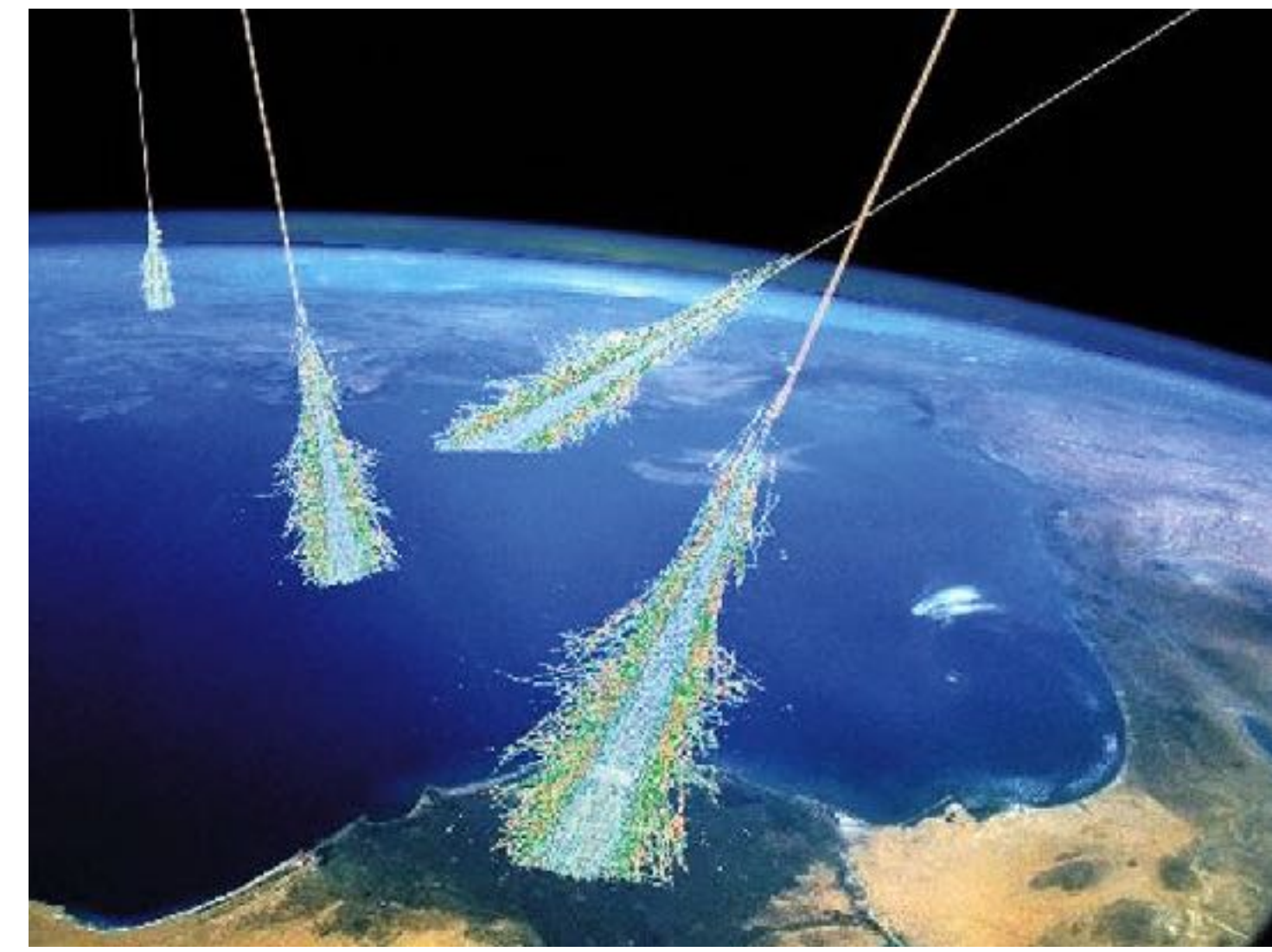


but the mothership survives
and keeps launching more

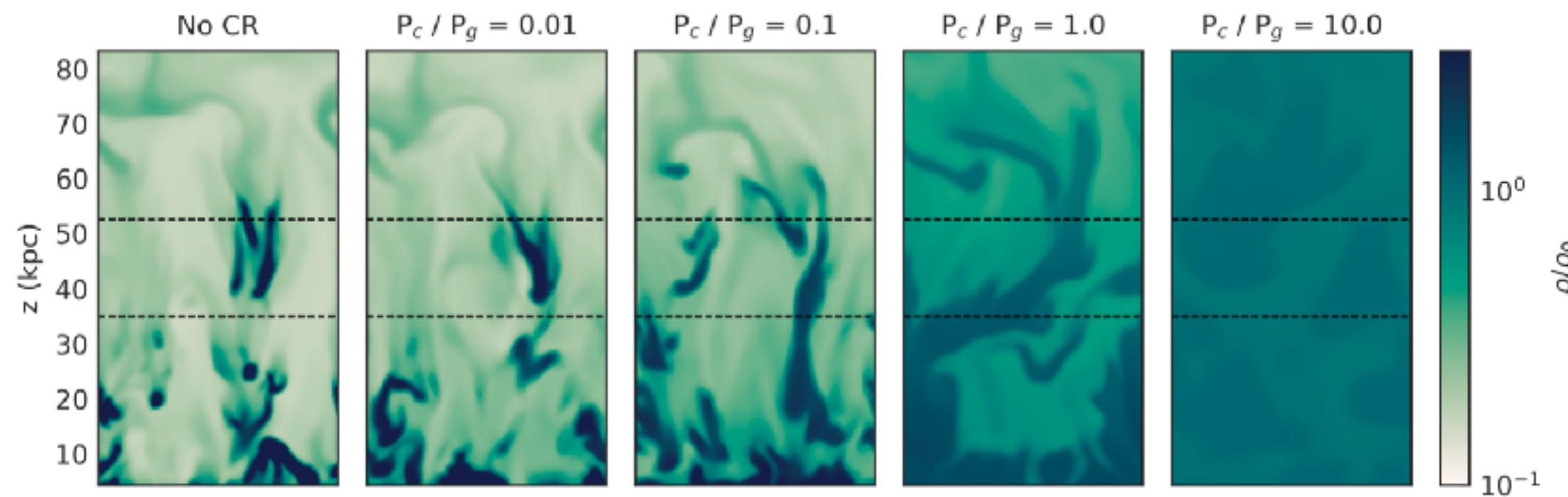


Booth et al 2013

Cosmic Rays

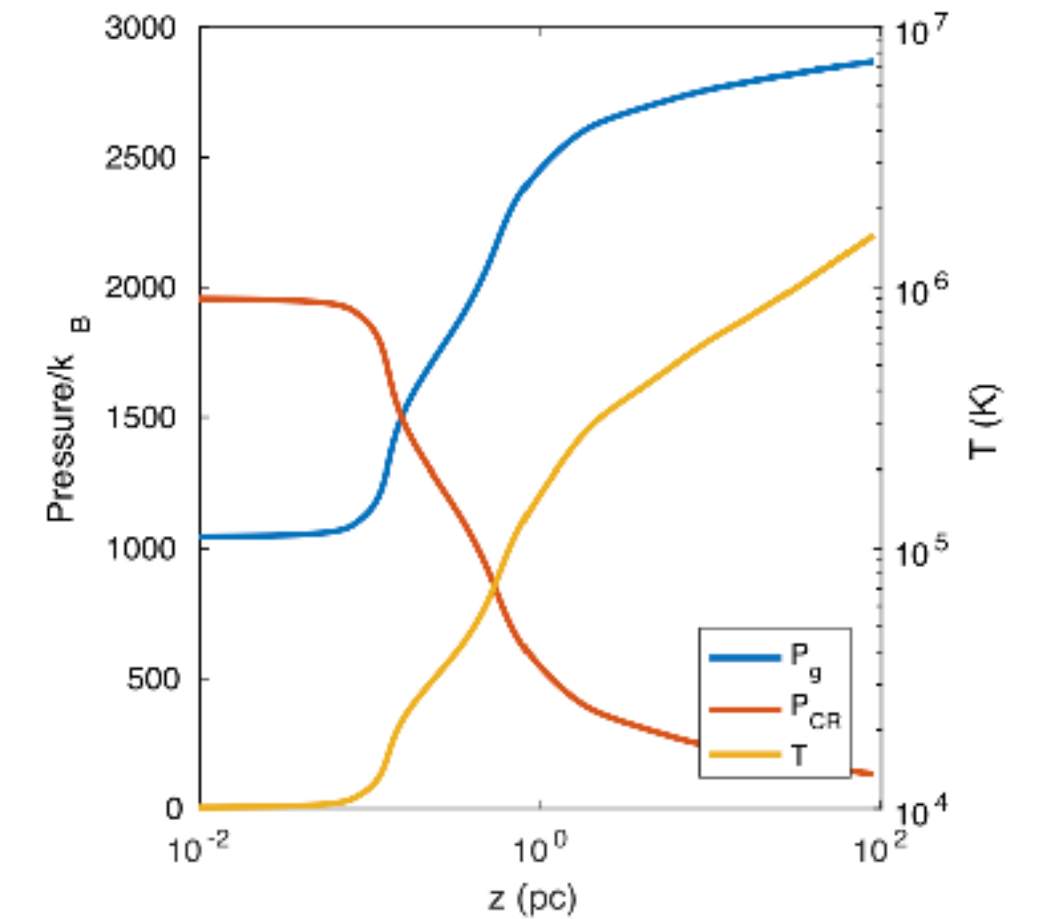


Why should you care about Cosmic Rays?



They provide
non-thermal pressure
support

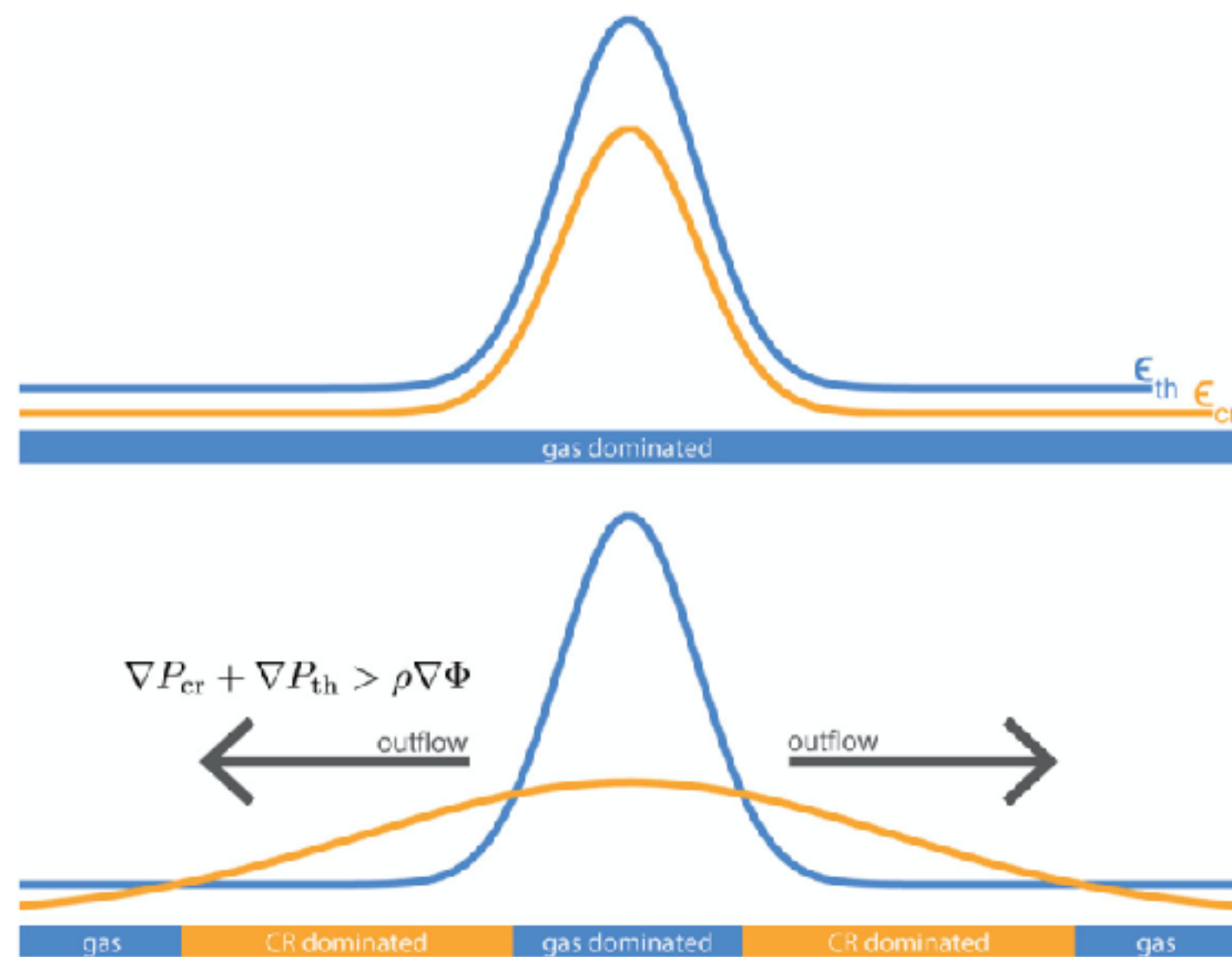
Butsky+20



ASK ME

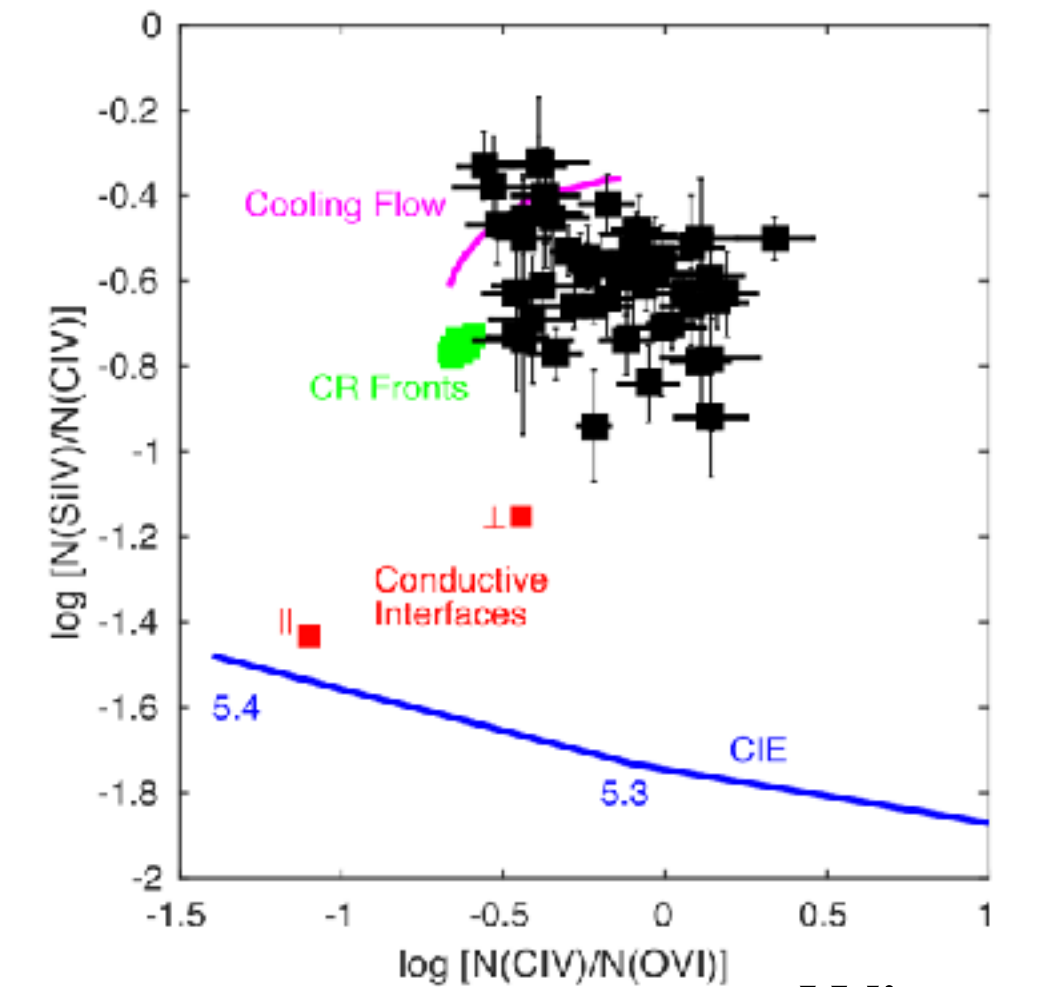


IF I CARE



They drive
Galactic winds
and dominate in
the halo

Salem & Bryan 2014



Wiener+17

They heat and alter
thermal interfaces

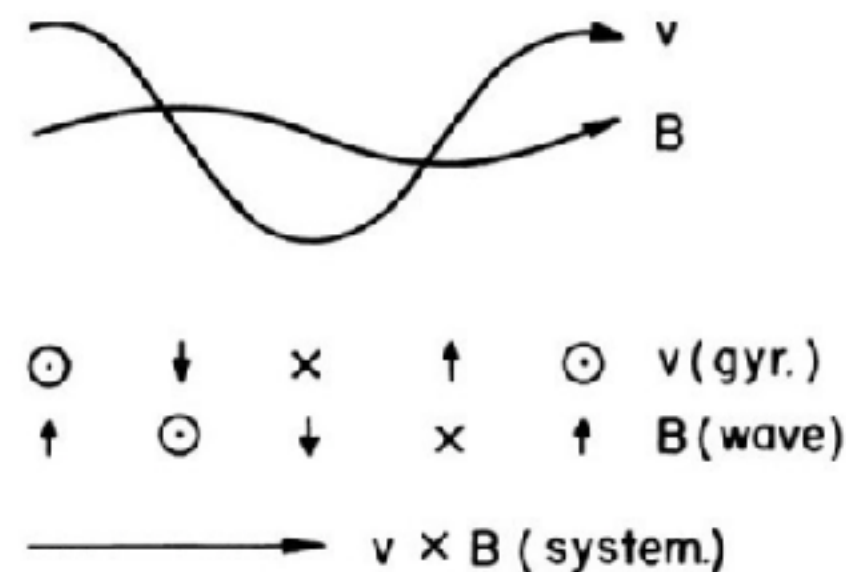
Cosmic Ray Physics in a Nutshell

See Zweibel 2013, 2017 for reviews and Ellen's talk on Thurs!

CRs are relativistic... but they live in our Galaxy for $\sim 3 \times 10^7 \text{ yr} \gg \frac{L_{\text{gal}}}{c}$

Why??

They are **self-confined** by the streaming instability



Wentzel 1972

Steady Lorentz force

Scatter off magnetic fluctuations
Random walk out

$$\lambda \sim r_L / (\delta B / B)^2 \sim 1 \text{ pc}$$

Scatter in pitch angle $\delta\theta \sim \pm \frac{\delta B}{B}$



Can write two-fluid hydrodynamic equations.

CRs **stream** at Alfvén velocity $v_A = \frac{B}{\sqrt{4\pi\rho}}$

As $\nabla P_c \rightarrow 0$ CRs decouple and stream at speed of light

They can also **diffuse** relative to Alfvén wave frame, with flux: $F_c = -\kappa_{\parallel} \nabla P_c$

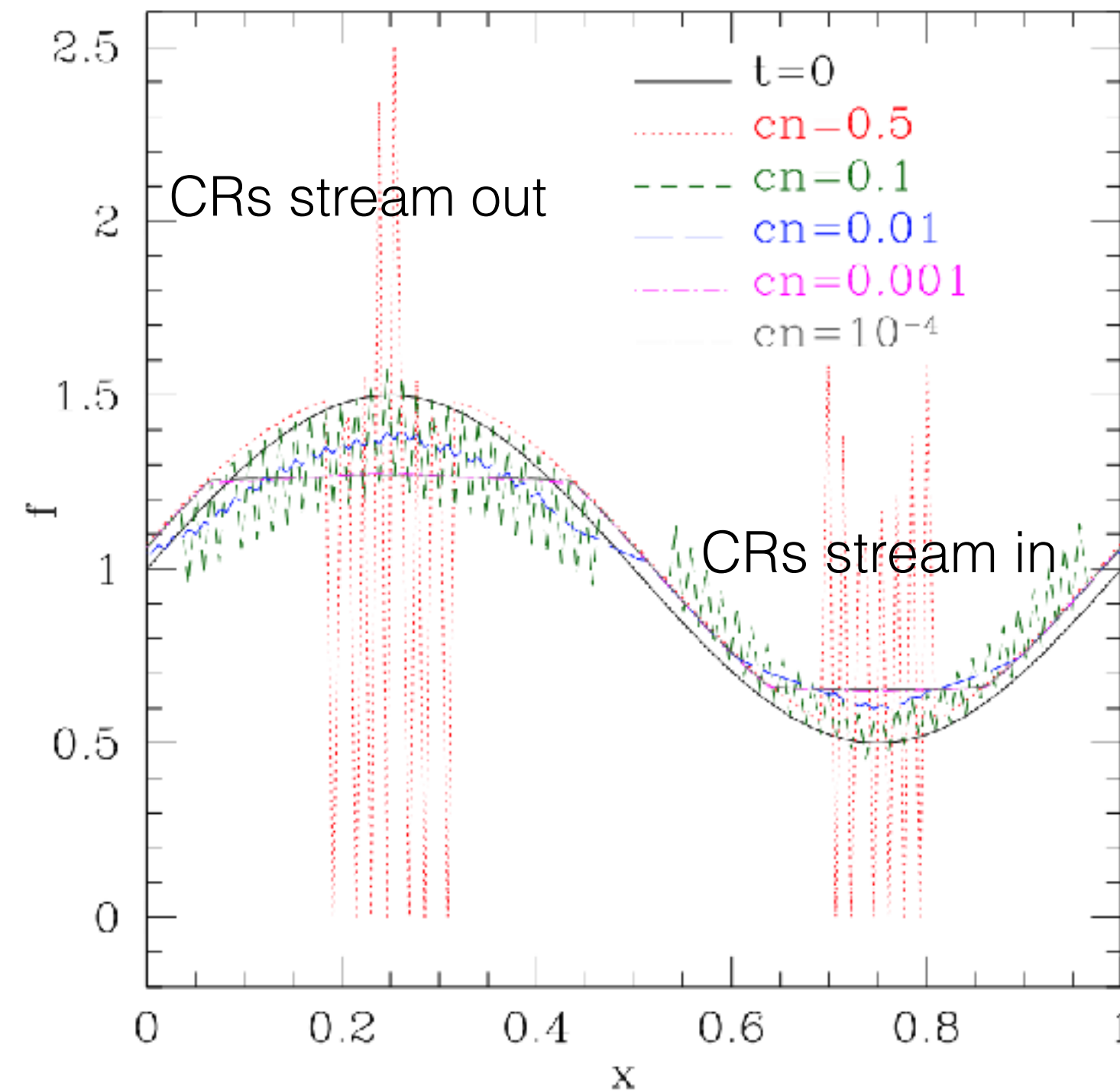
Tangled magnetic fields can also cause effectively diffusive behavior

CRs **push** gas with force ∇P_c

CRs **heat** gas at rate $v_A \cdot \nabla P_c$

CRs are a relativistic fluid like photons

We should use radiative transfer methods on them



Sharma et al 2010

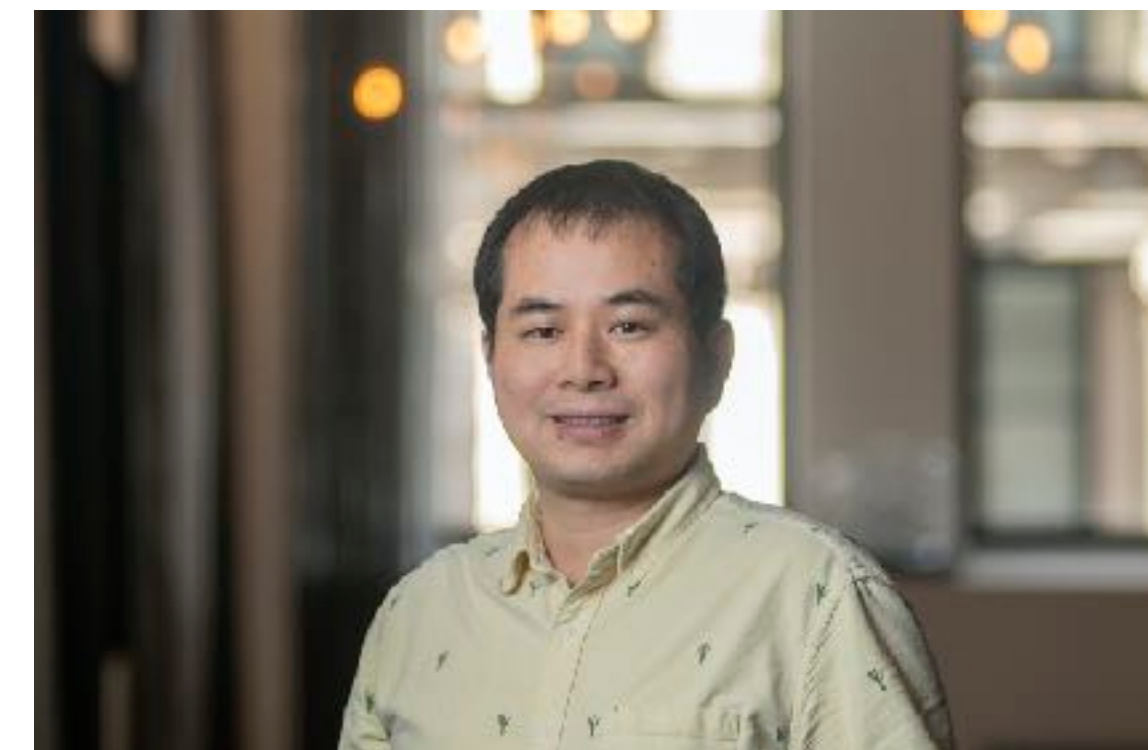
CRs can only stream down gradient: **grid scale instabilities!**

Have to add artificial diffusion: uncertain + expensive

Cured by two-moment method (Jiang & Oh 2018)

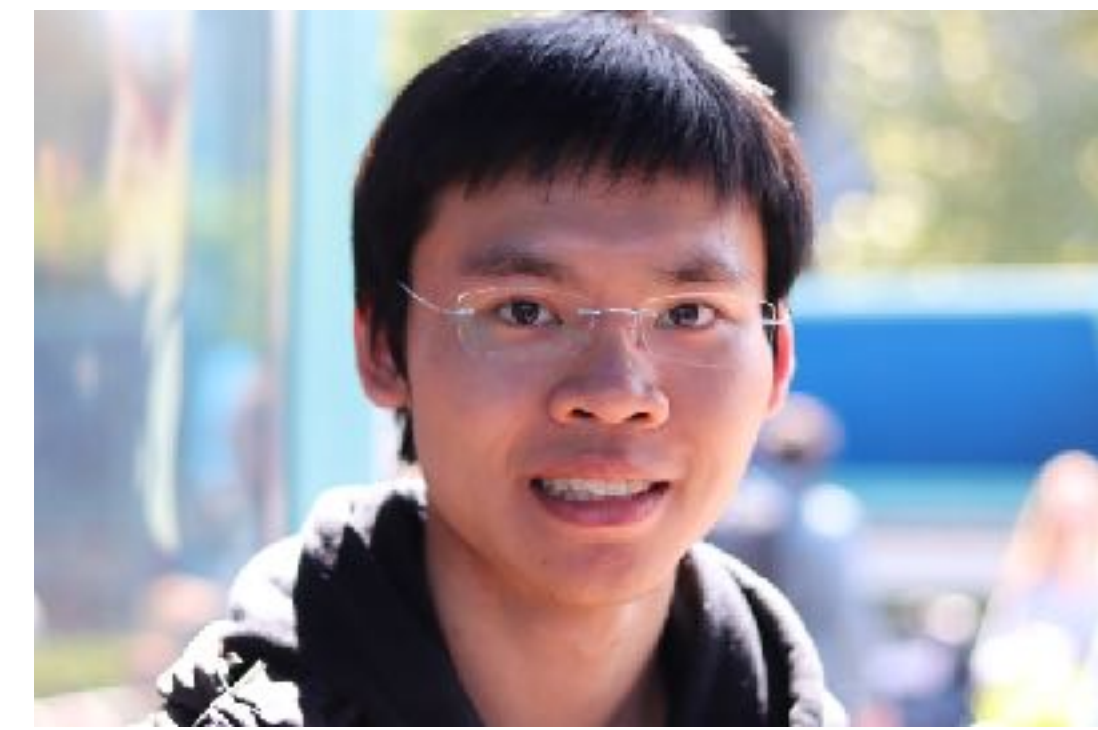
Enforces decoupling when $\nabla P_c \rightarrow 0$

Streaming + diffusion is now standard



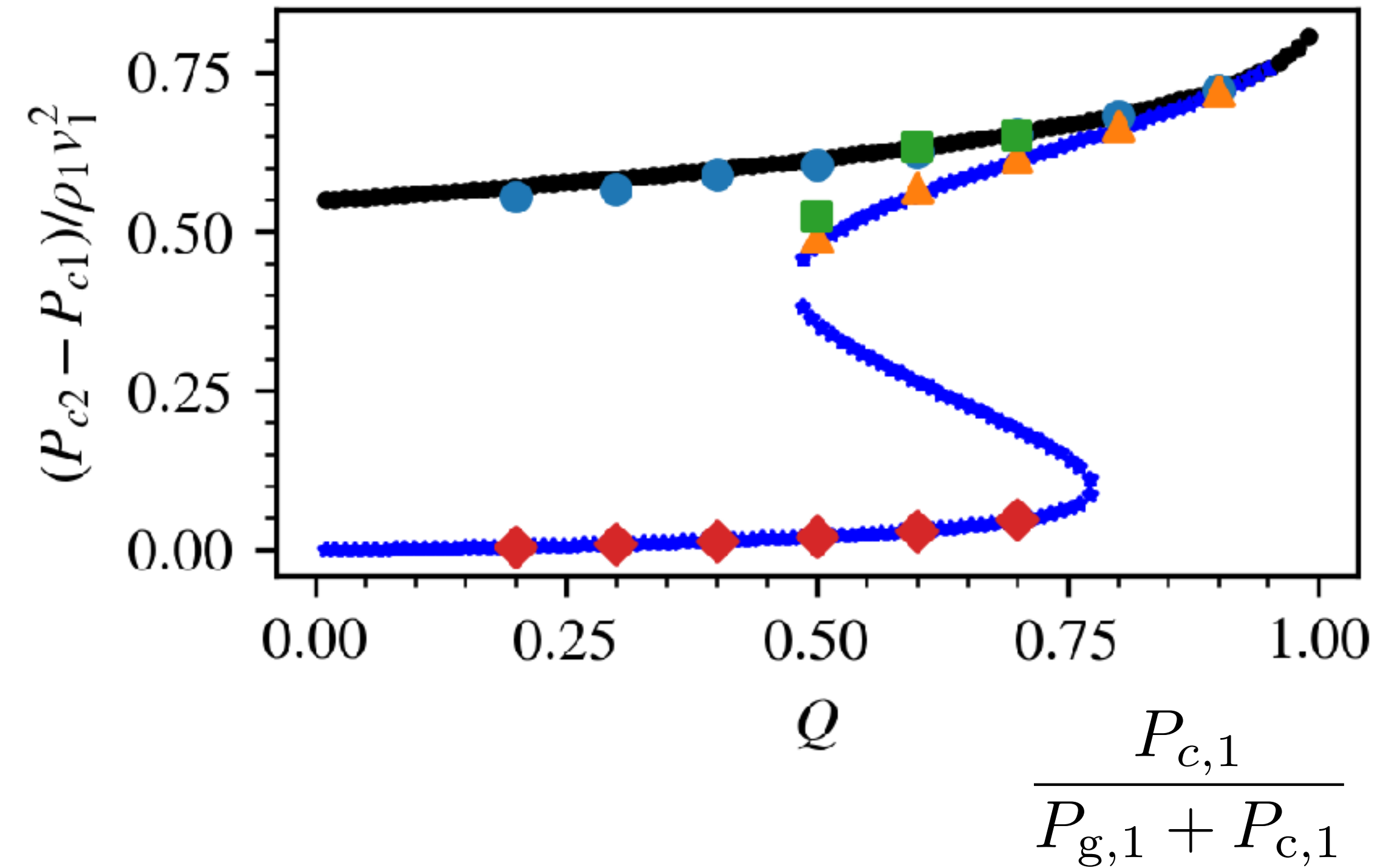
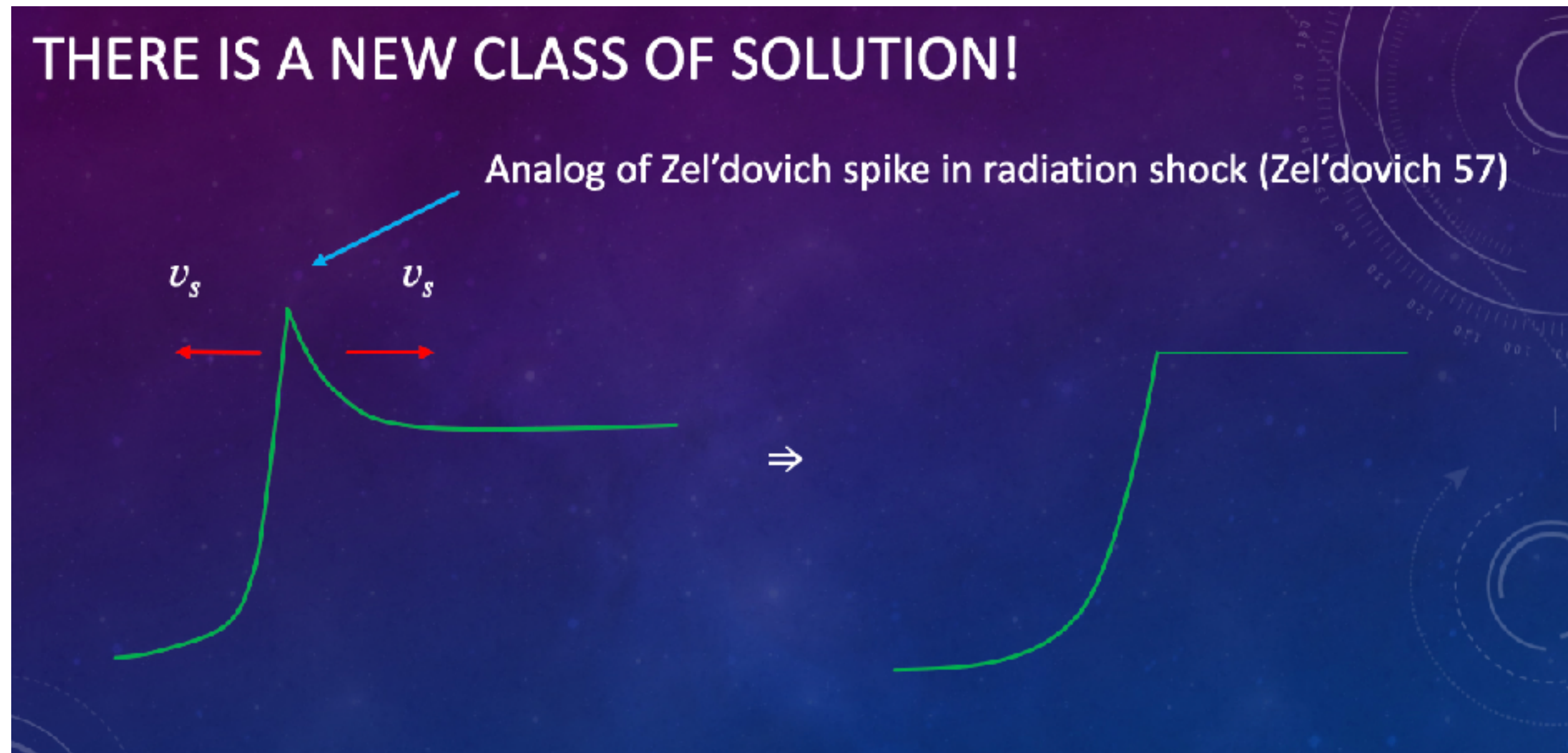
Tsung+20

CRs + shocks: most demanding test



Previously impossible to simulate with streaming

$$\beta = 1, \mathcal{M} = 20$$



A very demanding test of numerics

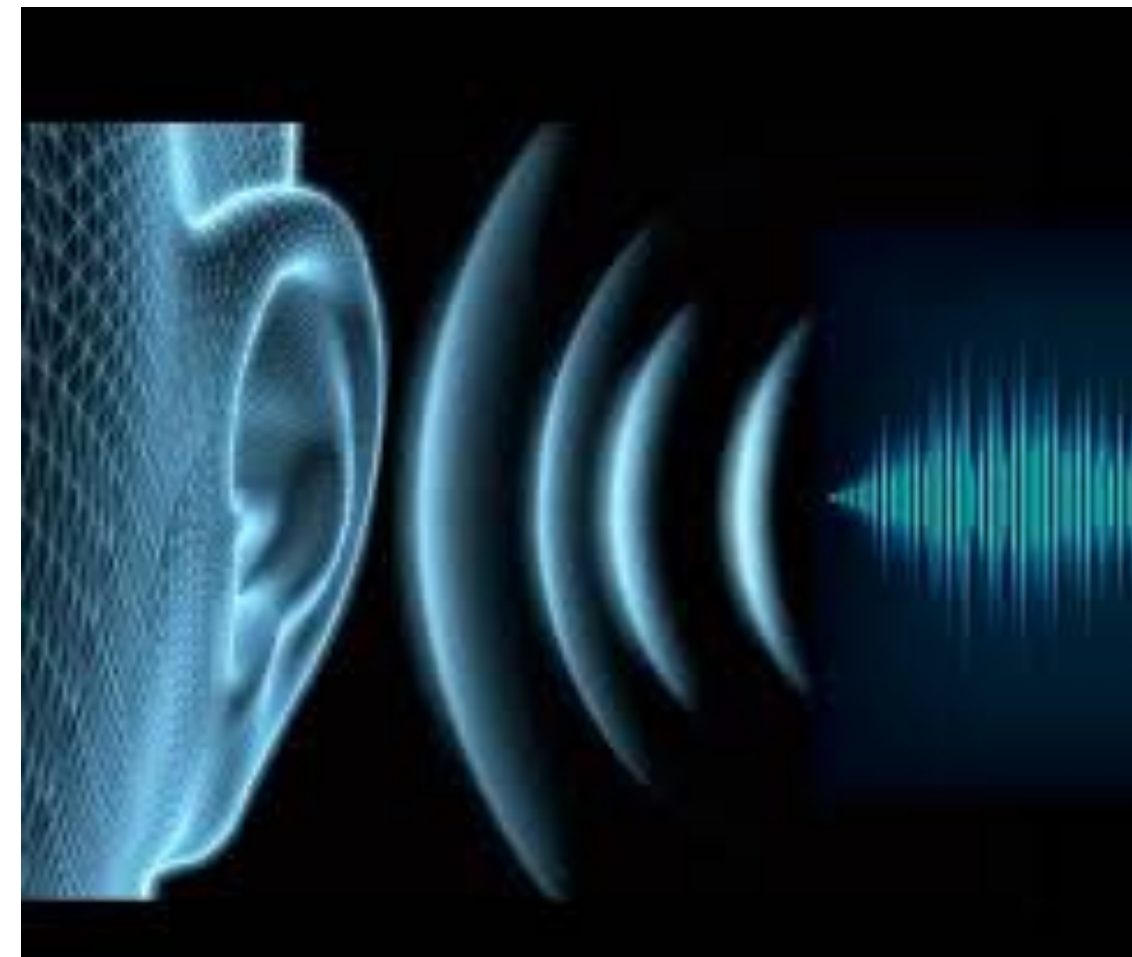
CR shocks equilibration times are very long (1000's of diffusion times)

OK, we got our elephant gun. What else to go after?



Turbulence!

Sound waves!

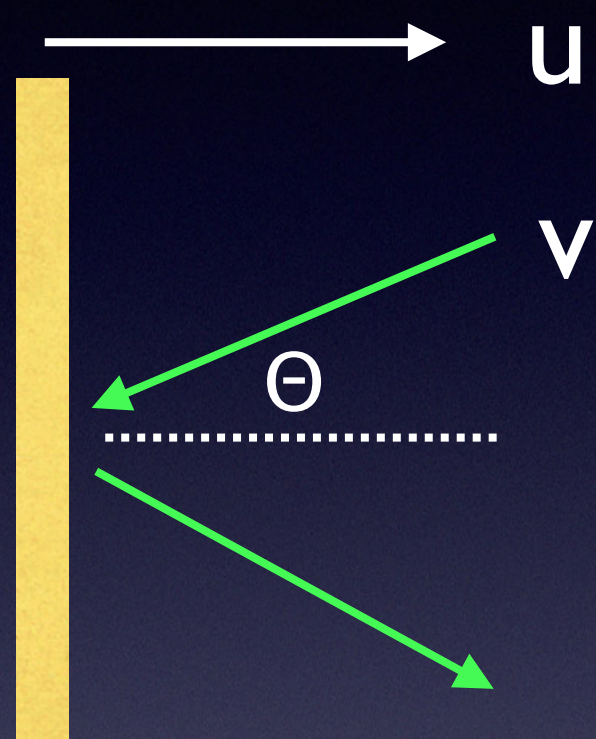


Just add Cosmic Rays and stir...



Where do CRs come from?

Fermi Acceleration



$$\frac{\Delta E}{E} = 4 \frac{u}{v} \cos \theta + 4 \frac{u^2}{v^2}$$

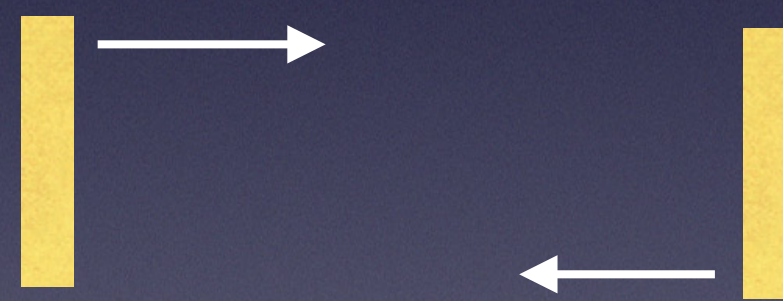
(interested in rel. particles with $v \gg u$)

random scatterers:

$$\langle \cos \theta \rangle = 0$$

$$\Delta E/E \sim u^2/v^2$$

(2nd order Fermi accel)



converging flow:

$$\langle \cos \theta \rangle = 1$$

$$\Delta E/E \sim u/v$$

(1st order Fermi accel)

Cosmic Rays + Turbulence: what happens?

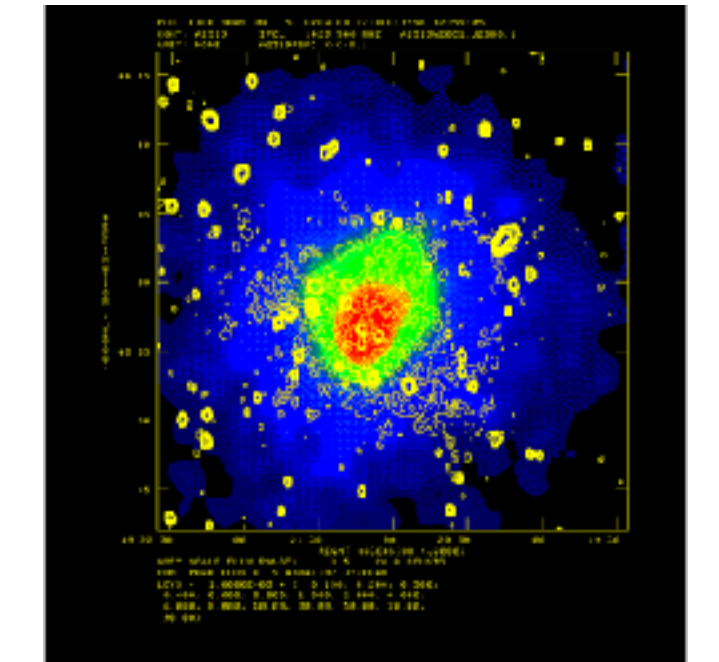


PLAYS WELL WITH OTHERS

2nd order Fermi: the original. Faster than you think!

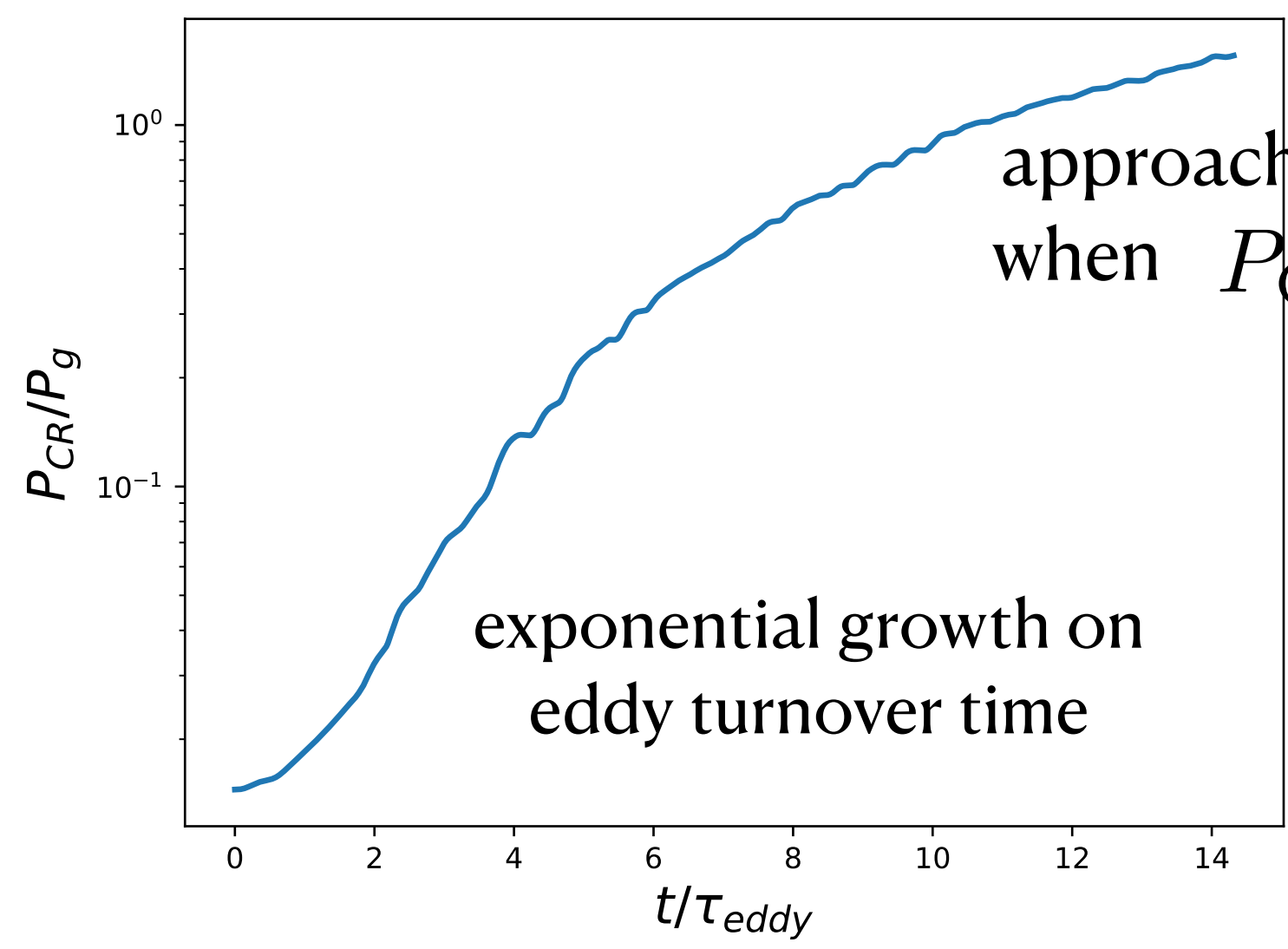
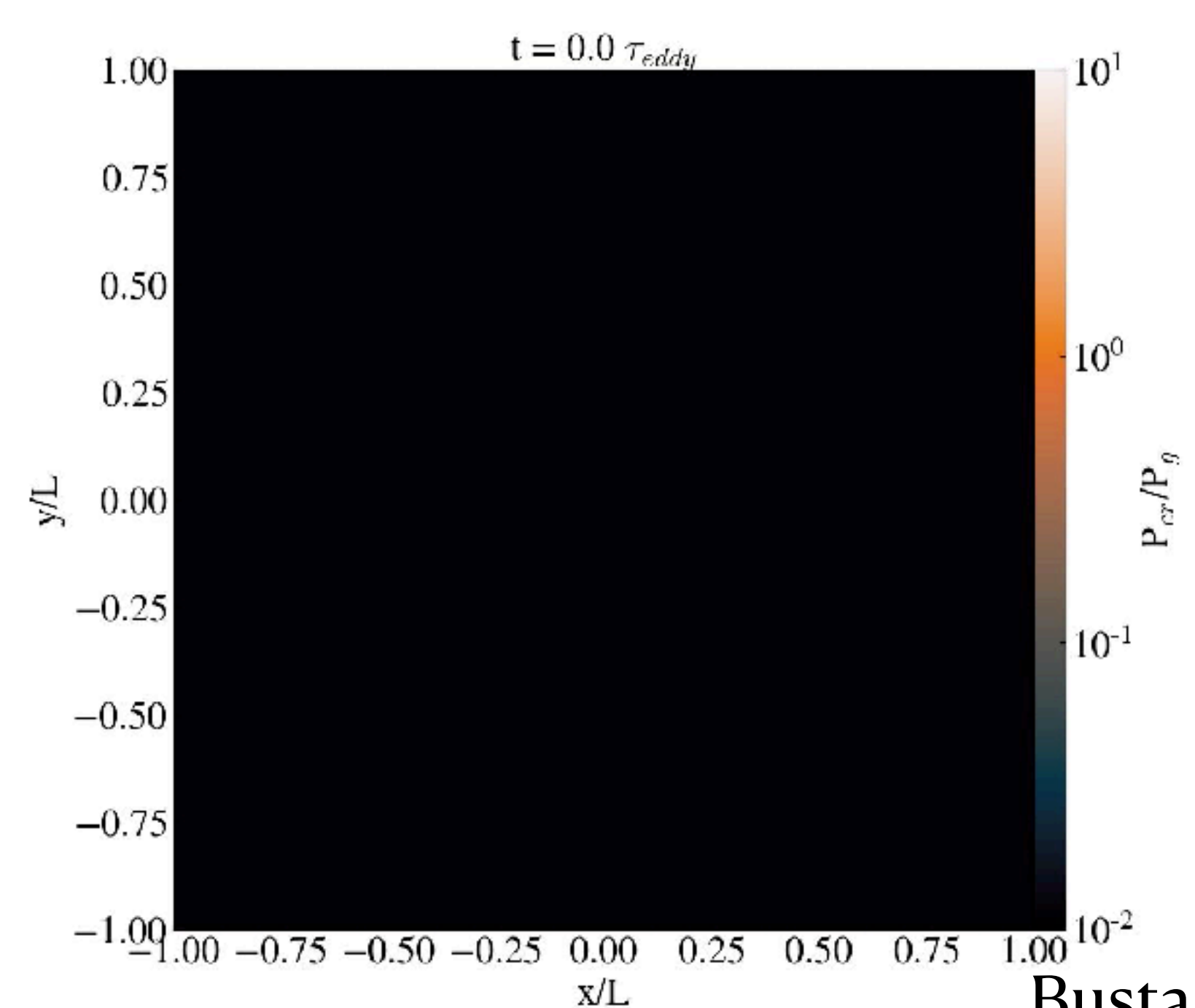
Resonant (transit time damping): large CR mfp, needs kinetic treatment, relevant to ICM

Non-resonant: small CR mfp, fluid approximation good, relevant to ISM + CGM



cluster radio halos

Neglected, not explored in simulations



Turbulence accelerates the CRs!

Bustard & Oh 2021

Chad Bustard

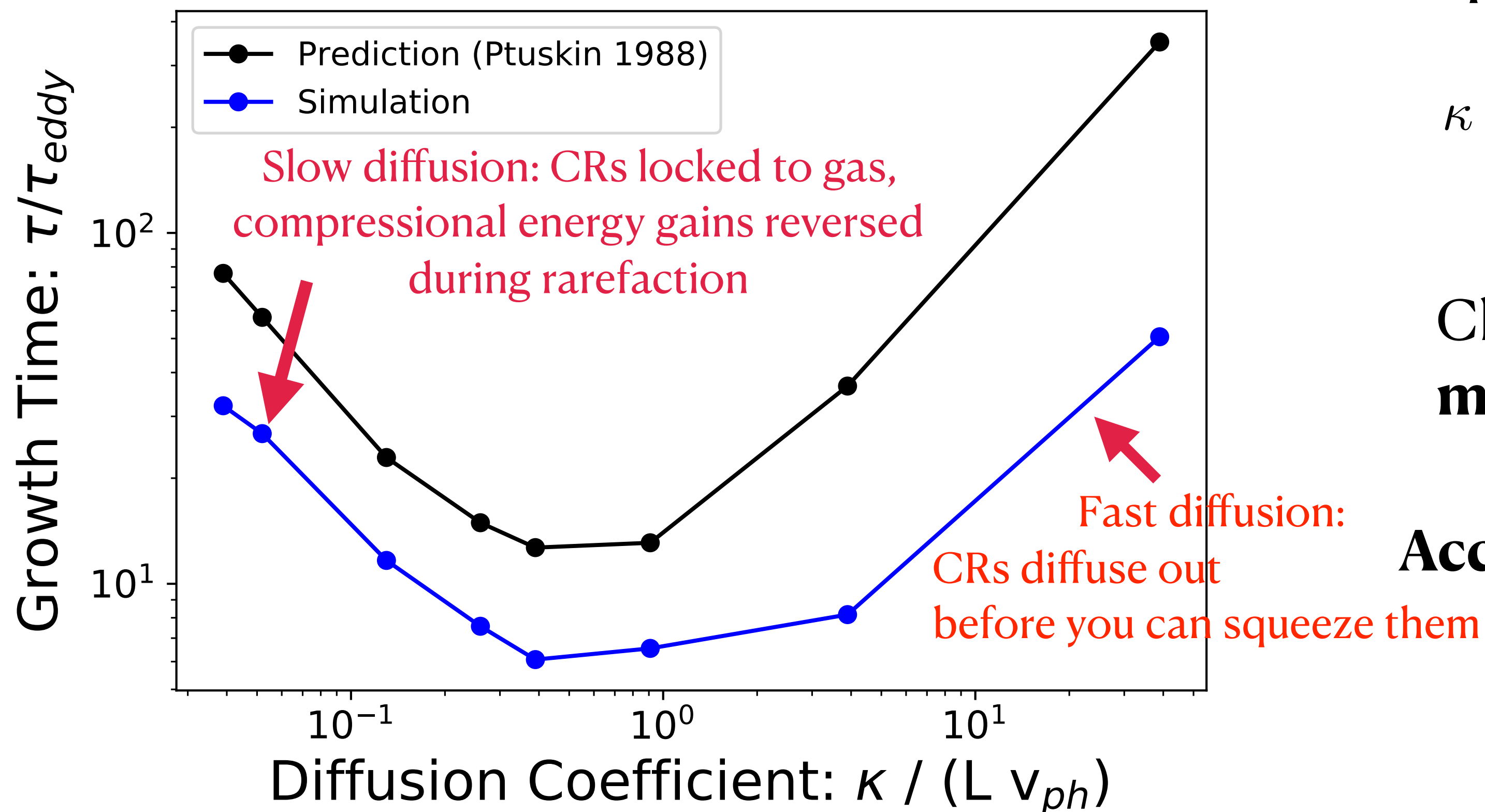




There is a sweet spot for acceleration

It's really just secular adiabatic heating

Isothermal, $\mathcal{M} \sim 0.5$ Turbulence



Happens when

$$\kappa \sim 3 \times 10^{29} \text{ cm}^2 \text{ s}^{-1} \left(\frac{L}{10 \text{ kpc}} \right) \left(\frac{v}{100 \text{ km s}^{-1}} \right)$$

Close to Milky Way values! Coincidence?... maybe not.

Acceleration time is short

$$\tau_{\text{grow}} \sim \tau_{\text{eddy}} \sim 10^8 \text{ yr} \left(\frac{L}{10 \text{ kpc}} \right) \left(\frac{v}{100 \text{ km s}^{-1}} \right)^{-1}$$

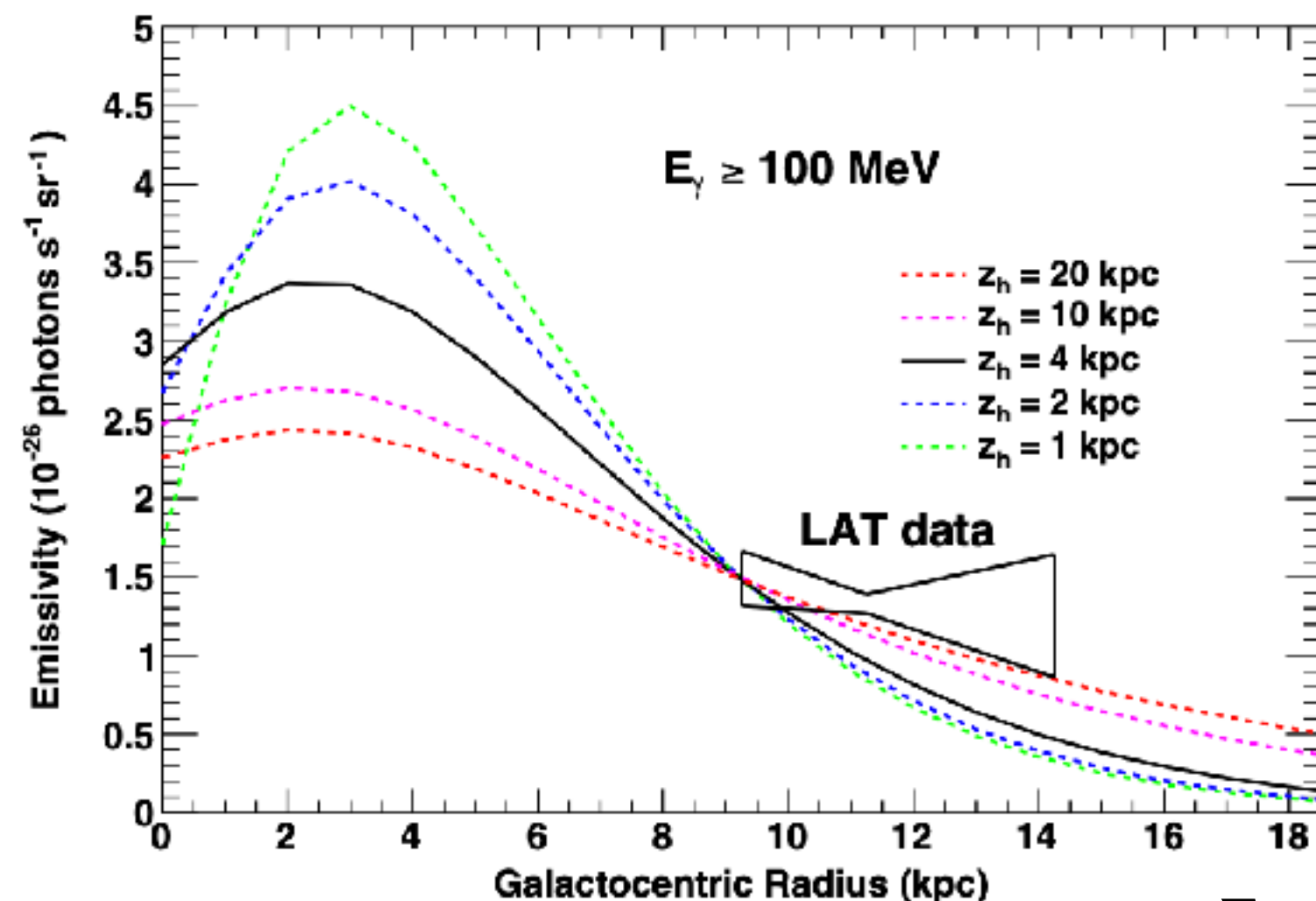


Nature *and* Nurture matter

CRs gain energy from turbulence in the halo



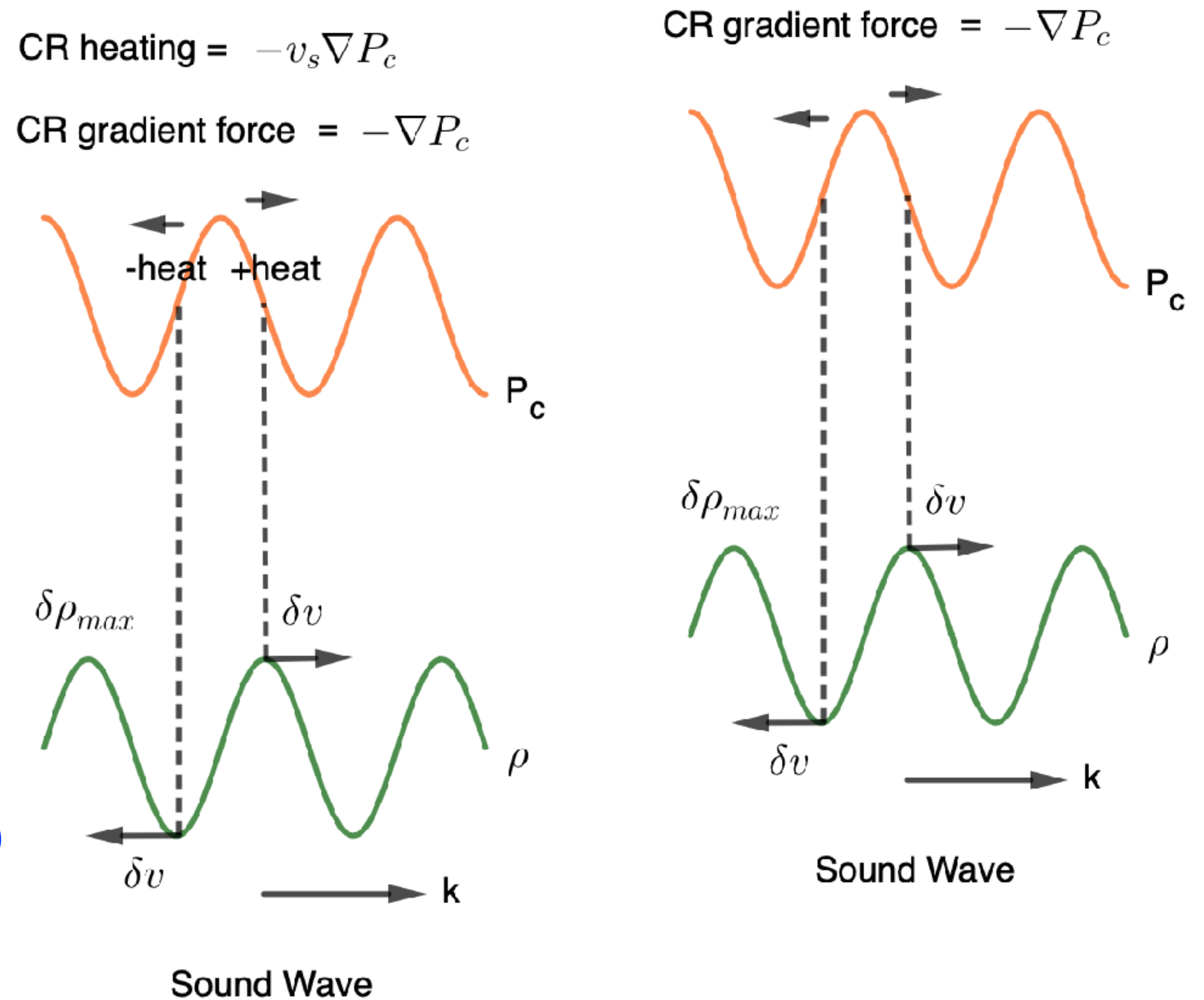
- CRs revived in halo — changing CR profile and wind solutions
- Turbulent ‘dynamo’ for CRs. Potentially enforces $P_{\text{CR}} \sim P_{\text{gas}}$?
- Affects gamma-ray luminosity vs. wind-driving, CR gradient problem, etc etc...



CR gradient problem: Cosmic Ray profile in Milky Way declines more slowly away from sources than expected from standard models

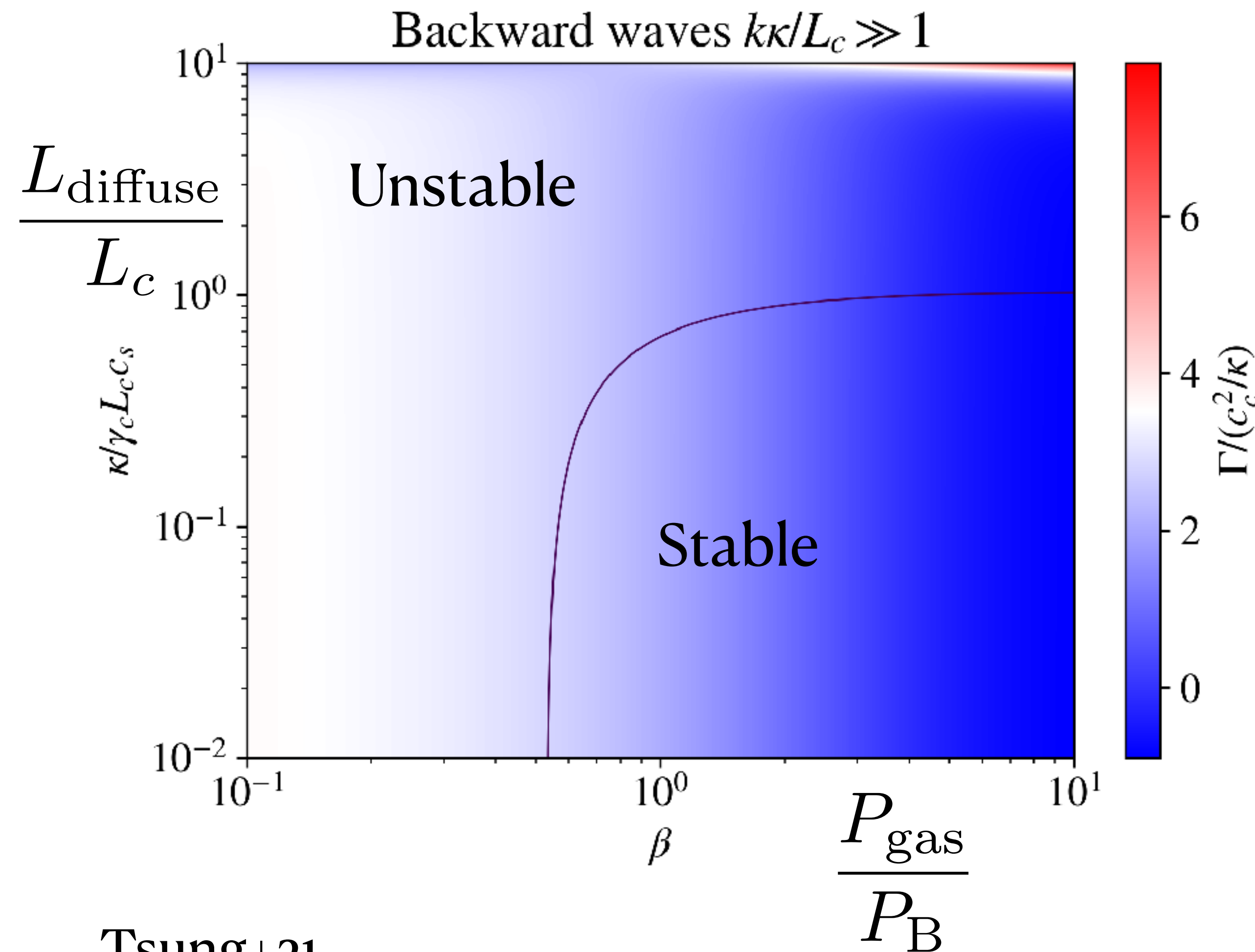
CRs + Sound Waves: what happens?

- Sound waves are harmonic oscillator
- Direct driving: CRs provide forcing which can be in phase (growth) or out of phase (damping) [Drury & Falle 1986](#)
- Indirect driving: CR streaming heats gas which also changes gas pressure forces. [Begelman & Zweibel \(1994\)](#)



Growth time is short! (of order wave period)

Instability with streaming, diffusion + background gradient



Tsung+21

For instability:

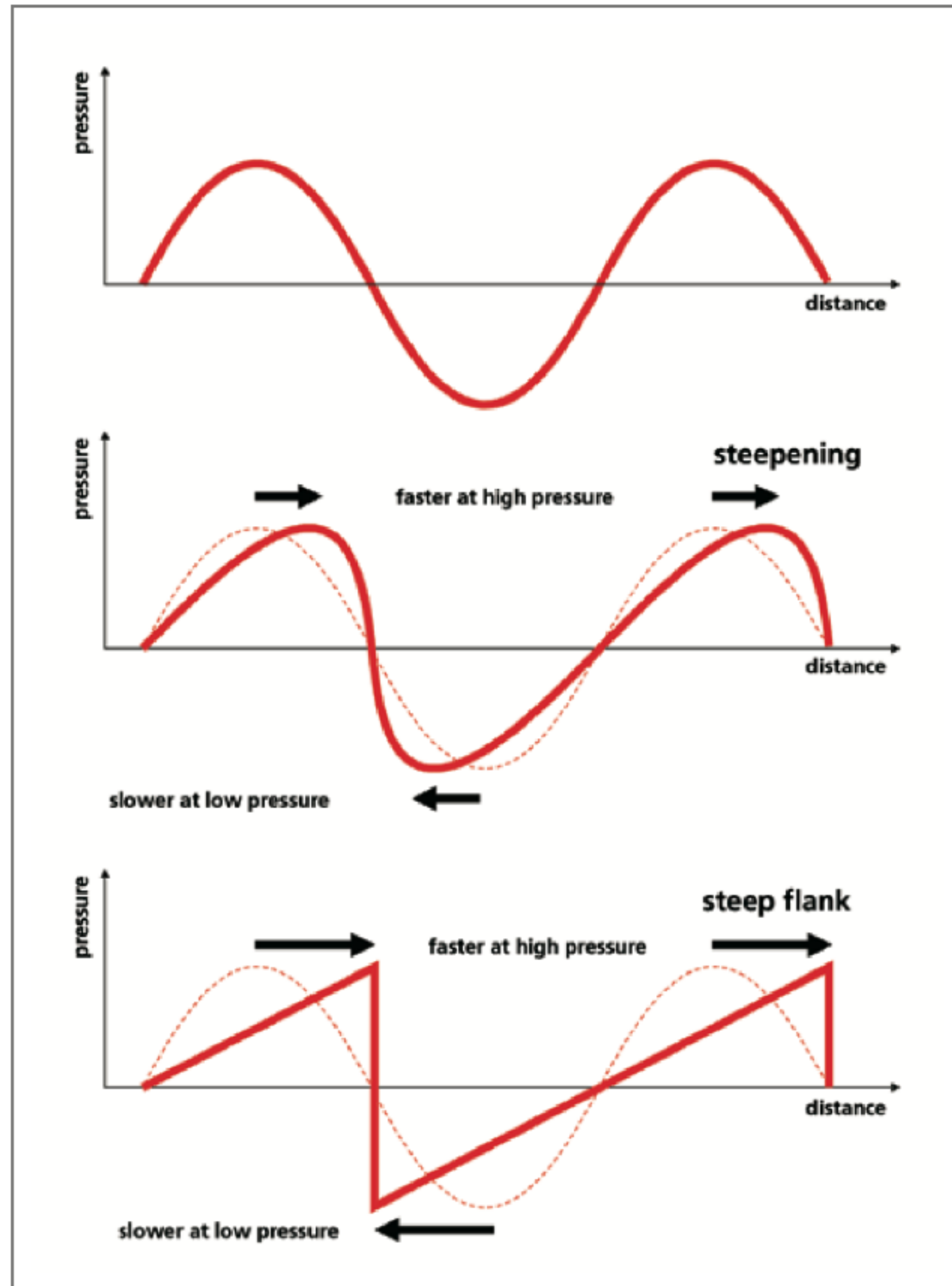
Short CR scale height (unstable regardless of B-field)

or

Strong B-fields (unstable regardless of CR scale height)

In very weak B-field (ICM) regime, potential CR instability driven by pressure anisotropy (Kempski & Quataert 2020)

What happens to a growing sound wave?



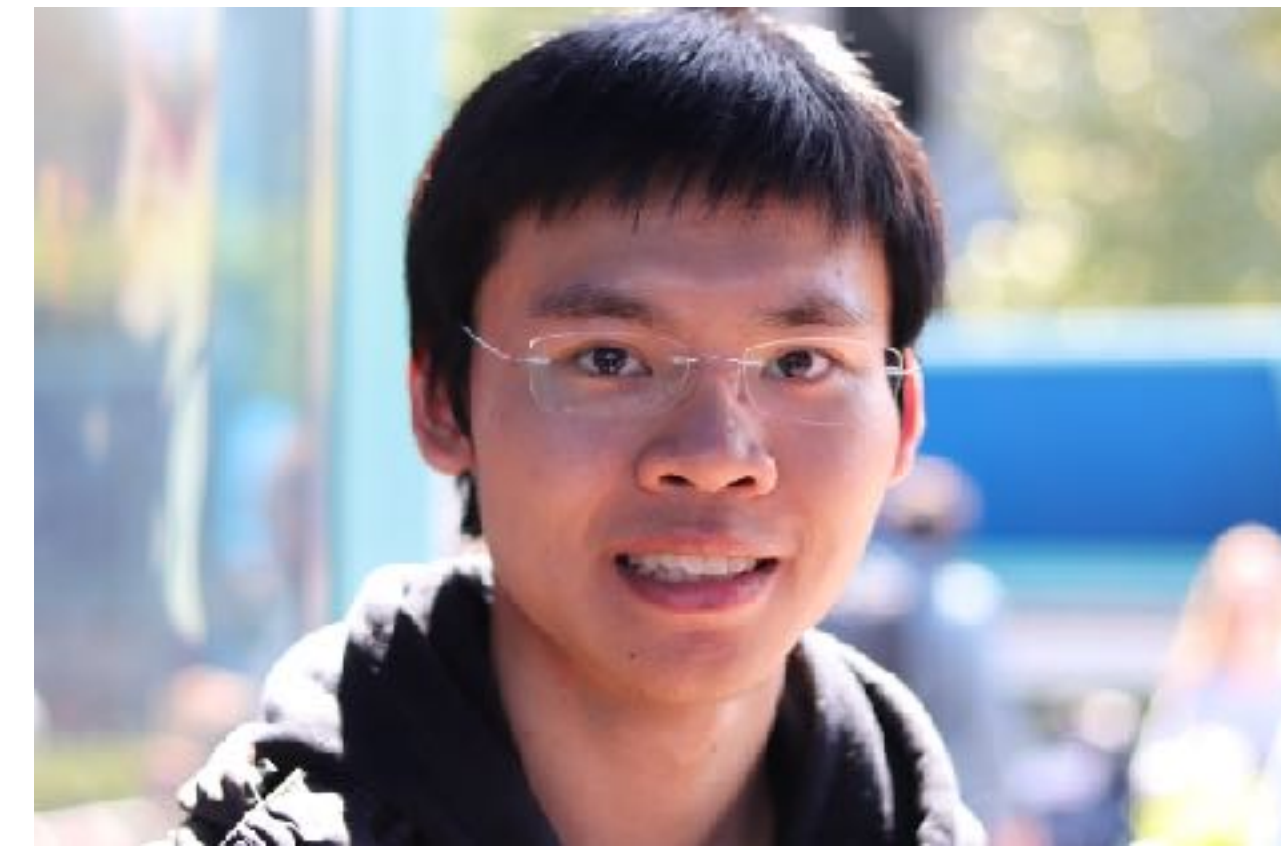
It goes non-linear, steepens and shocks

Figure from Storz 2007

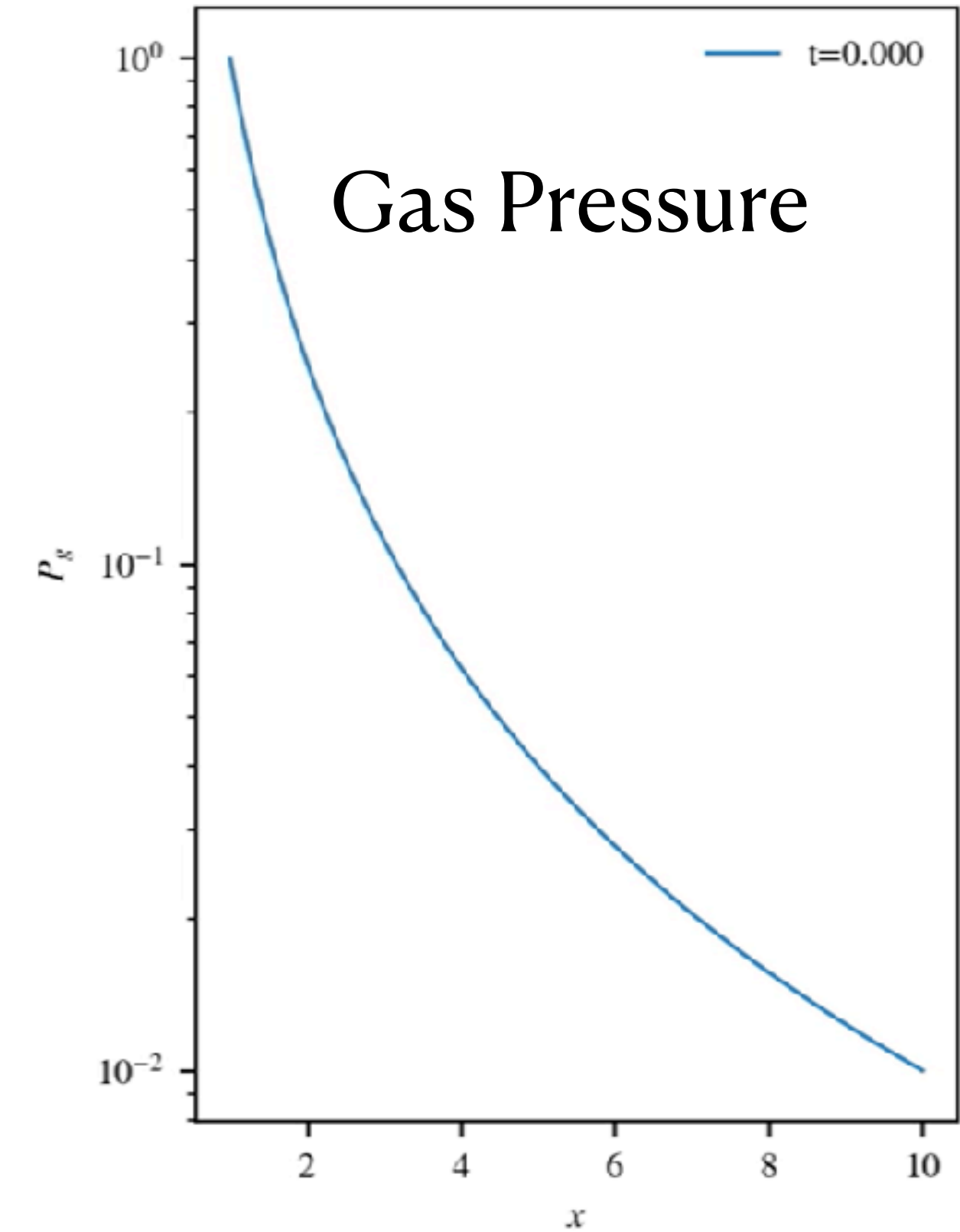
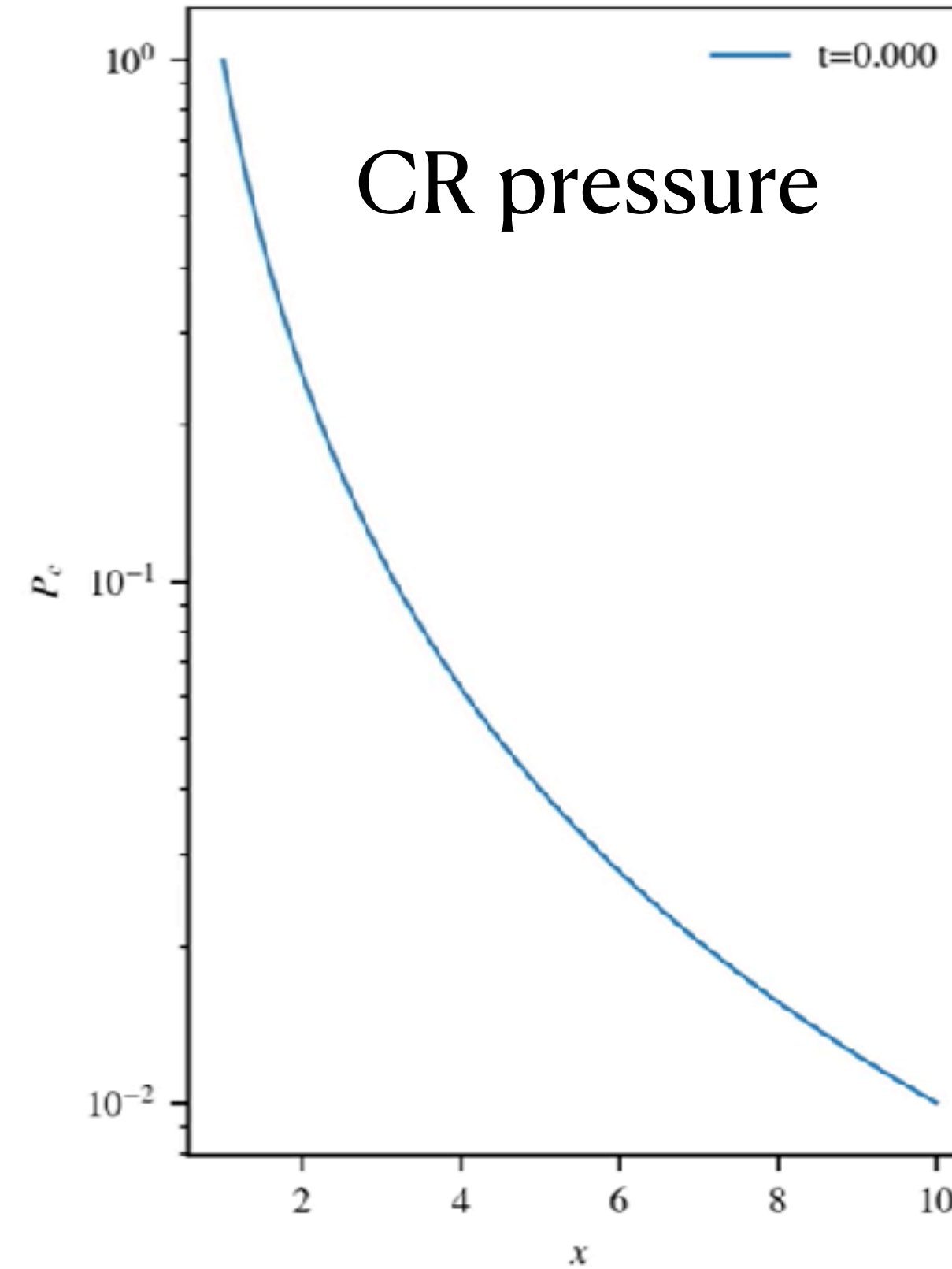
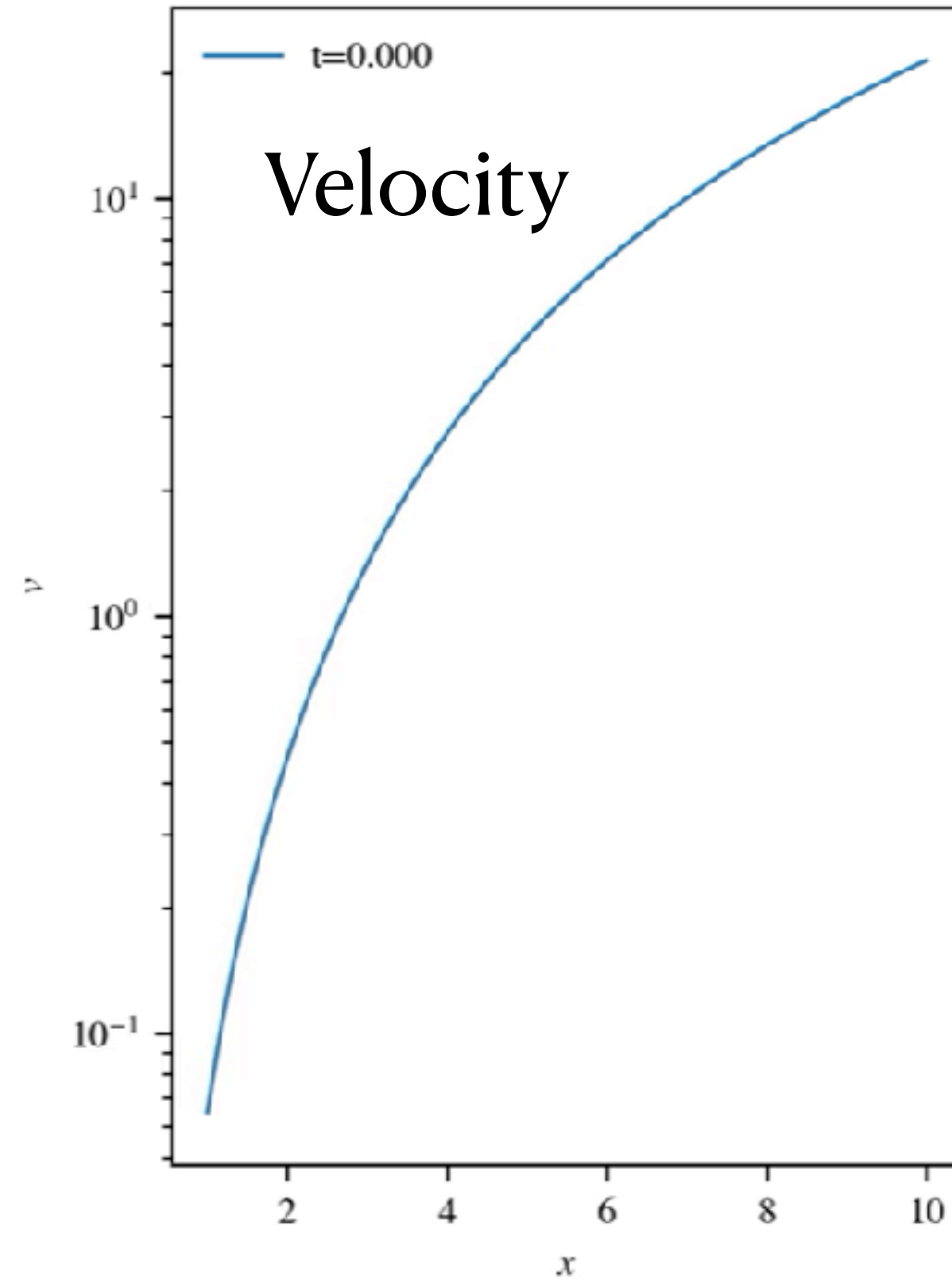
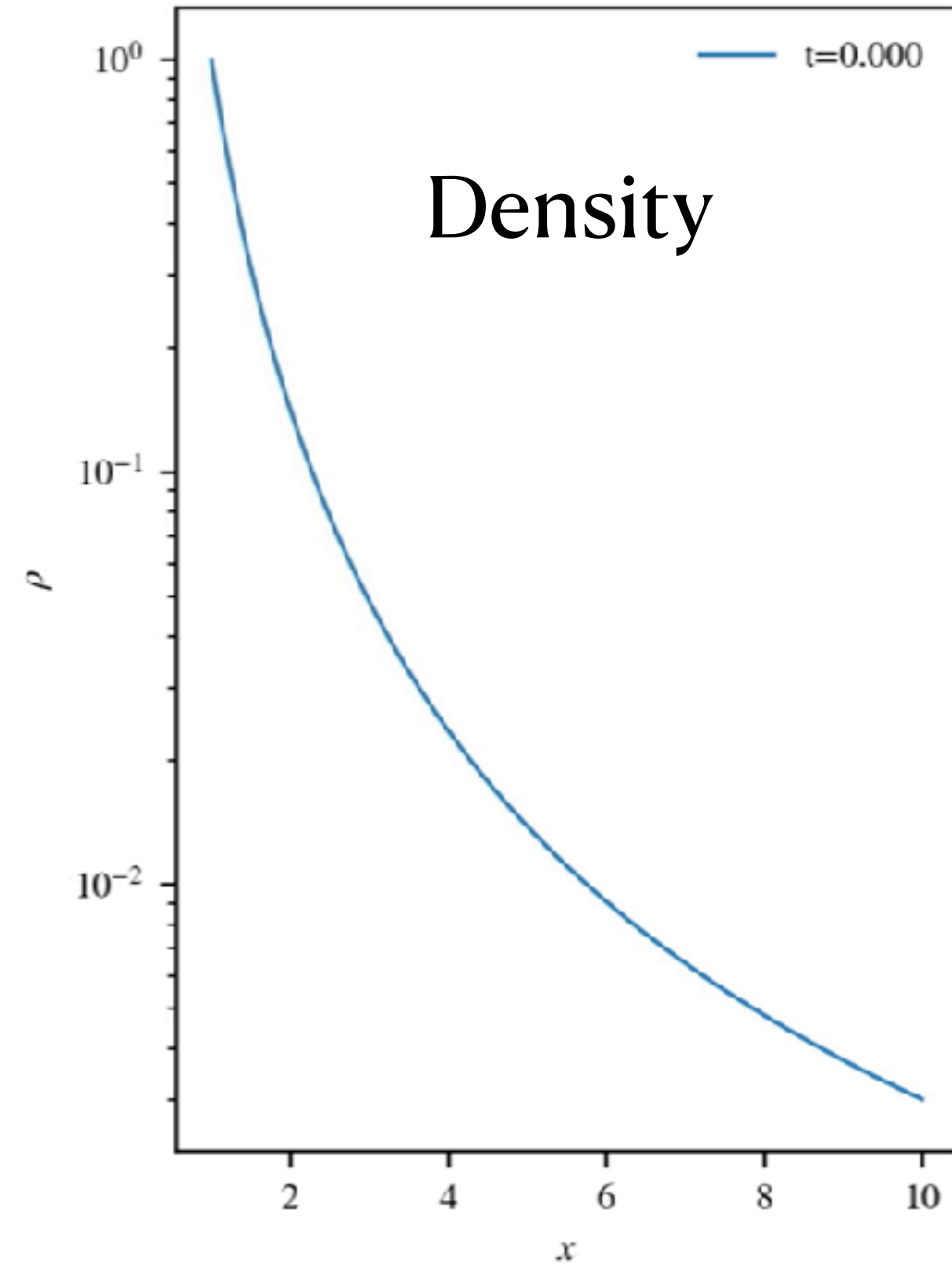
Bring it On

This staircase structure
is **new**: has never been seen
before in CR sims!

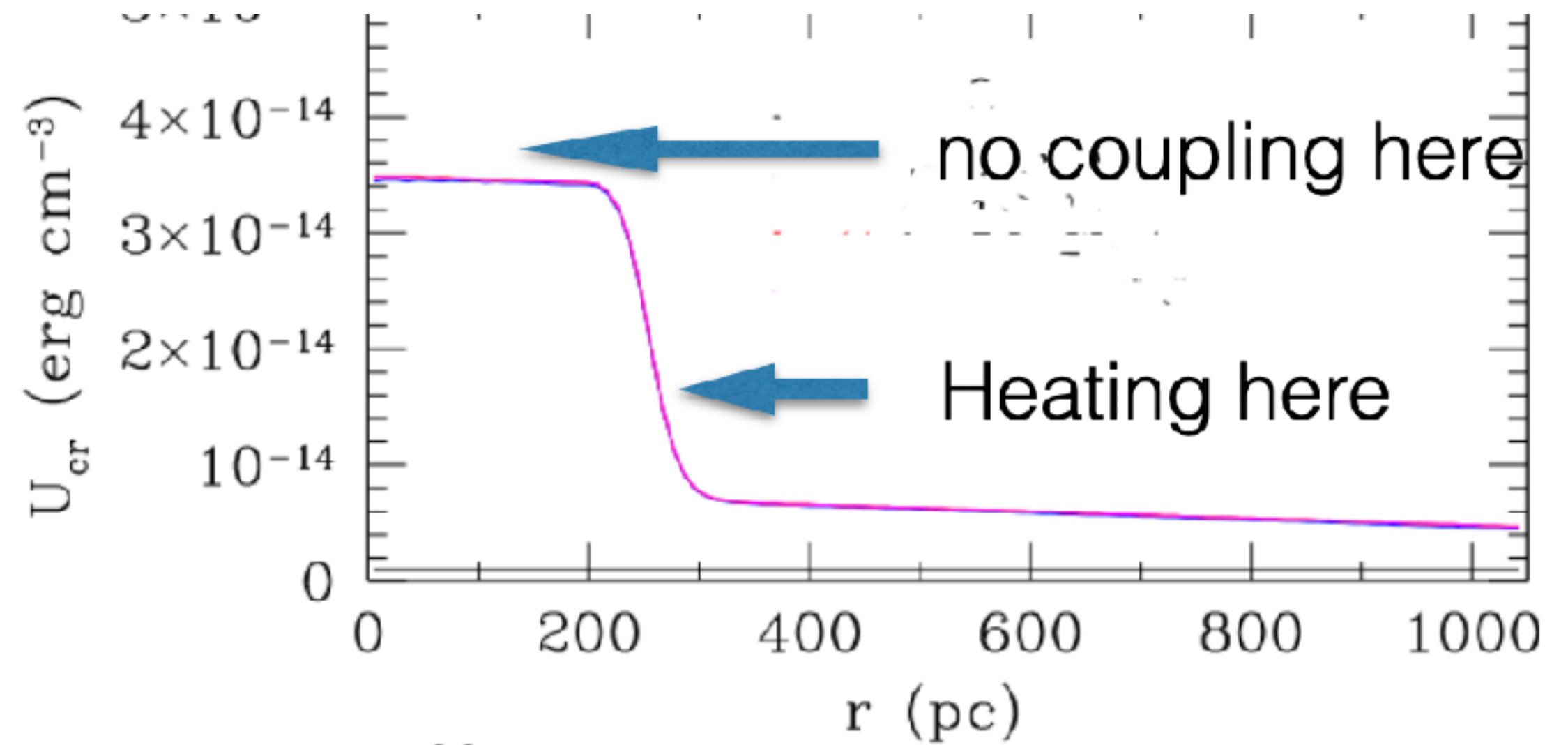
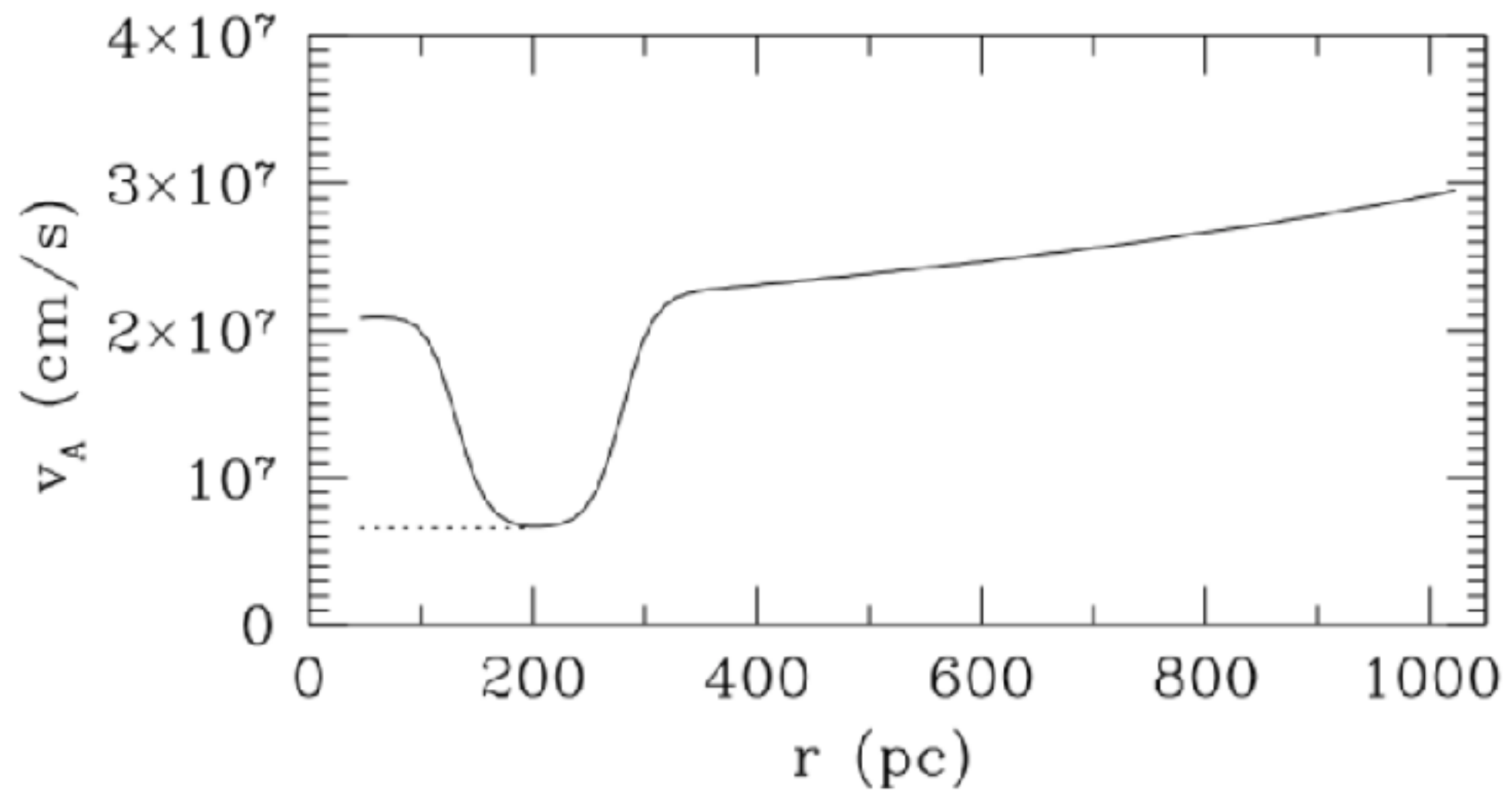
Weak shocks
separated by uncoupled
regions (free-streaming
CRs)



Tsung+21



CR Staircase is due to 'bottleneck effect'



Hydro turned off

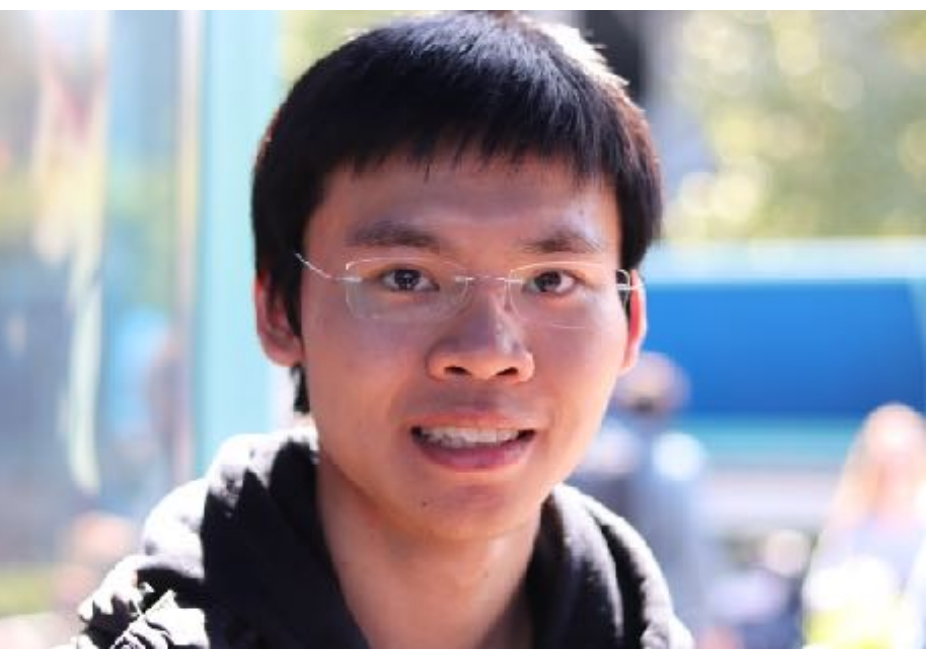
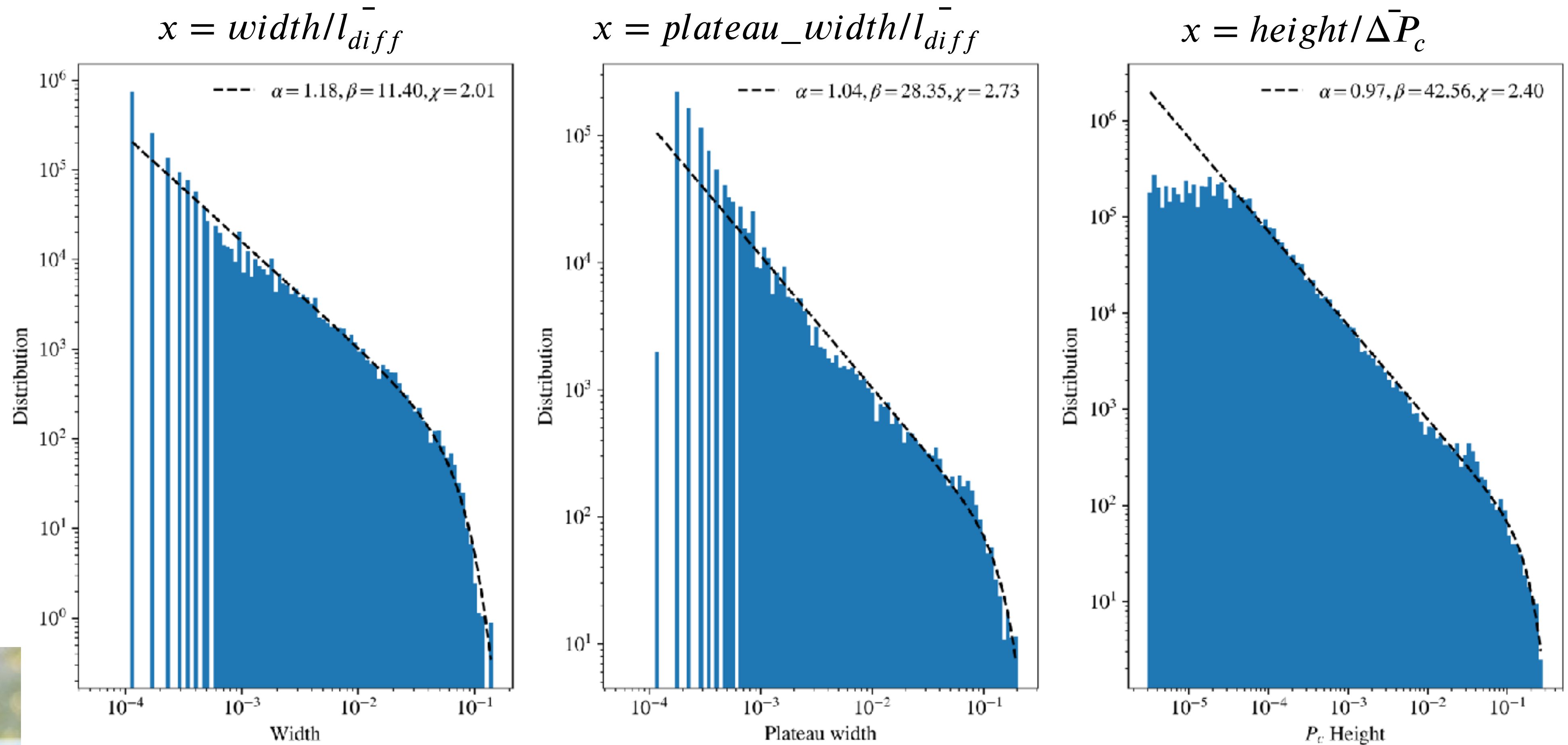
[Wiener, Oh & Zweibel 2017 \(see also Skilling 1971\)](#)

Minimum in Alfvén speed ('bottleneck') causes CRs to decouple

System reaches a statistical steady state

Fitting function $f \propto x^{-\alpha} \exp\left(-\left(\frac{x}{\beta}\right)^\chi\right)$

Statistics



Distribution of widths and height follows a Schechter-like function!
Possible to understand characteristic lengthscales

Reminiscent of adhesion model!

Extension of Zeldovich approximation — solves Burger's equation

Burgers' equation, Devil's staircases and the mass distribution for large-scale structures

M. Vergassola¹, B. Dubrulle², U. Frisch¹, and A. Noullez¹

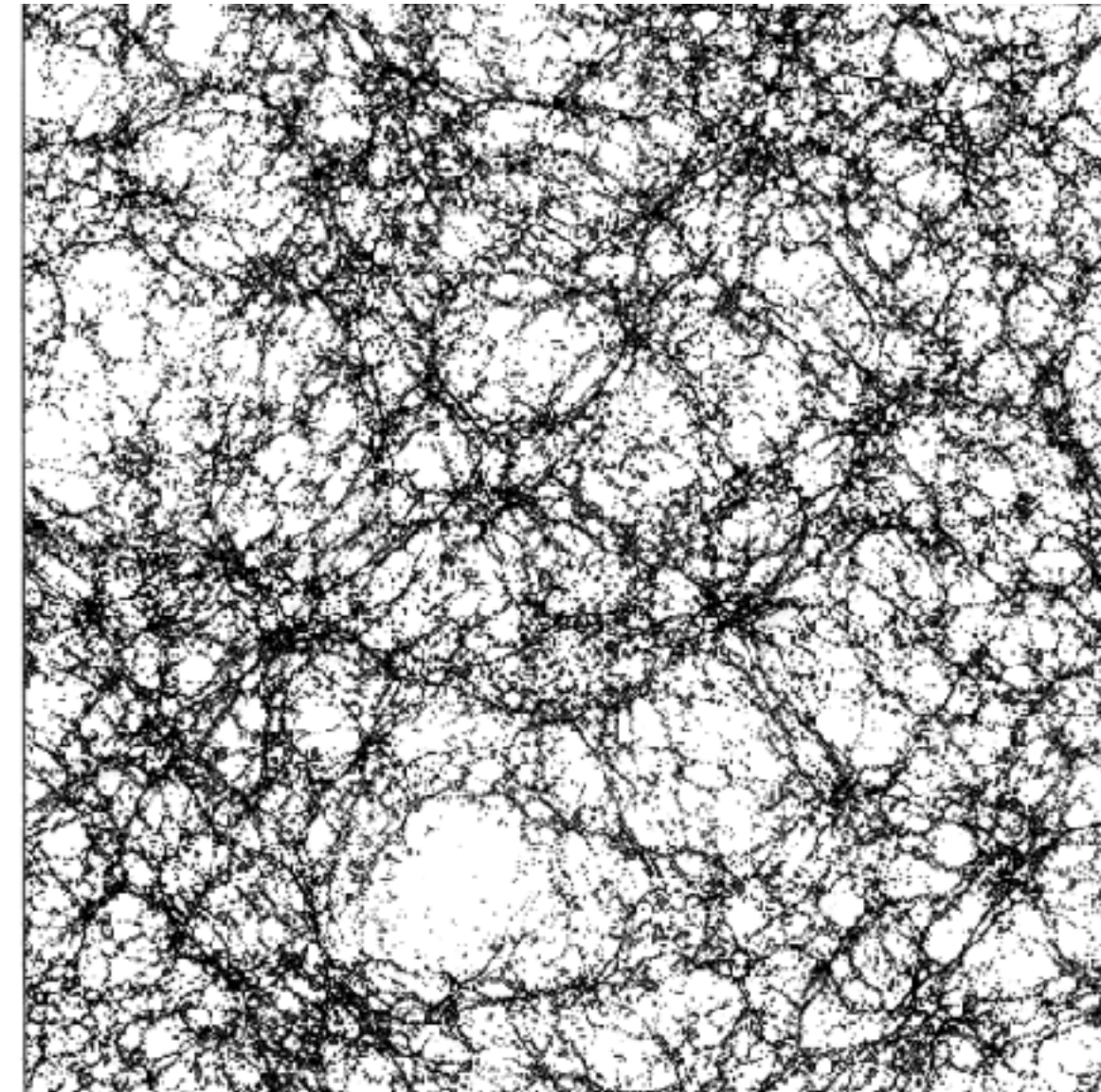
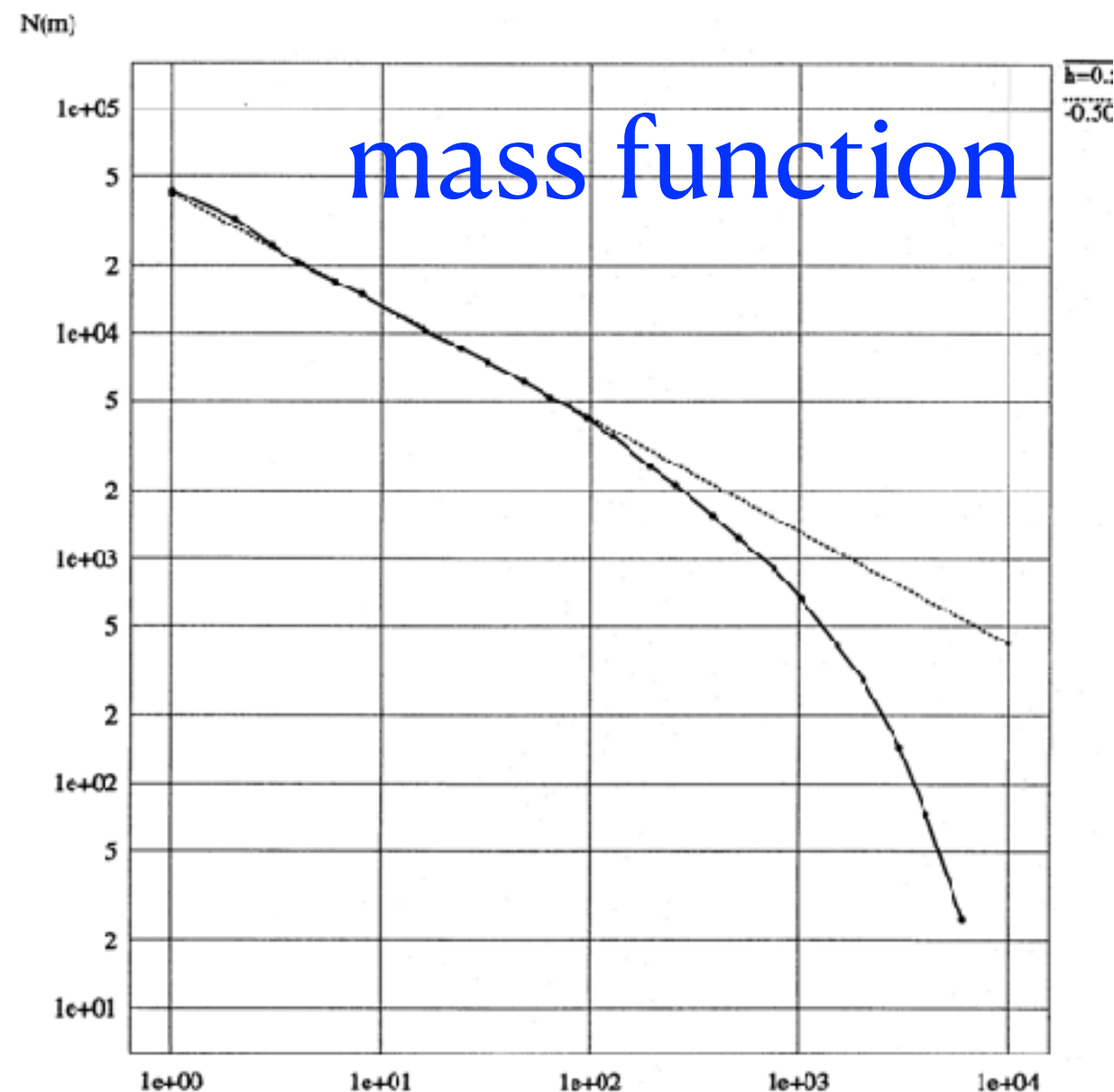
¹ CNRS, URA 1362, Observatoire de la Côte d'Azur, BP 229, F-06304 Nice Cedex 4, France

² CNRS, URA 285, Observatoire Midi Pyrénées, 14 av. E. Belin, F-31400 Toulouse, France

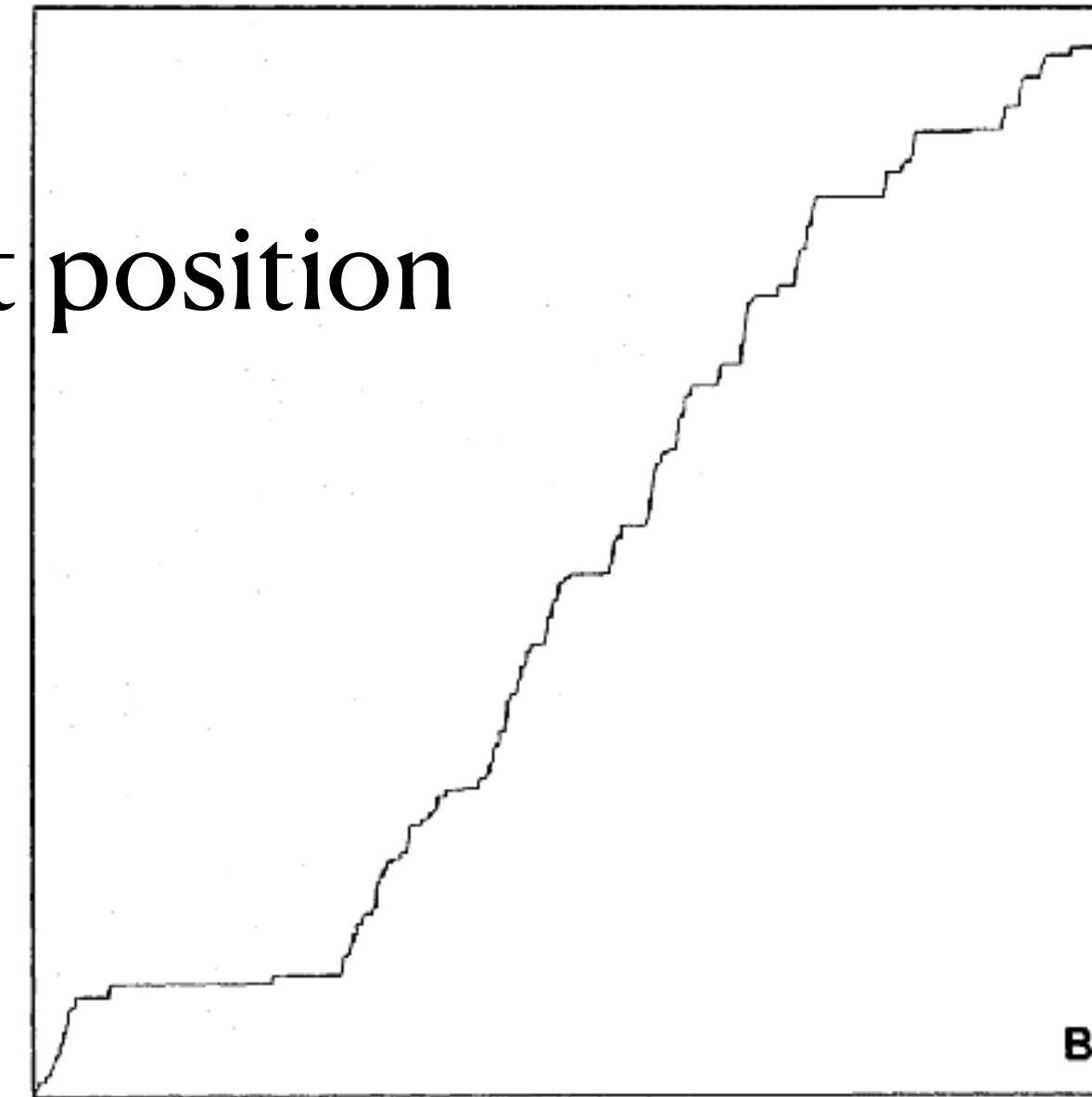
Received 27 October 1993 / Accepted 8 February 1994

Schechter-like

mass function



Current position



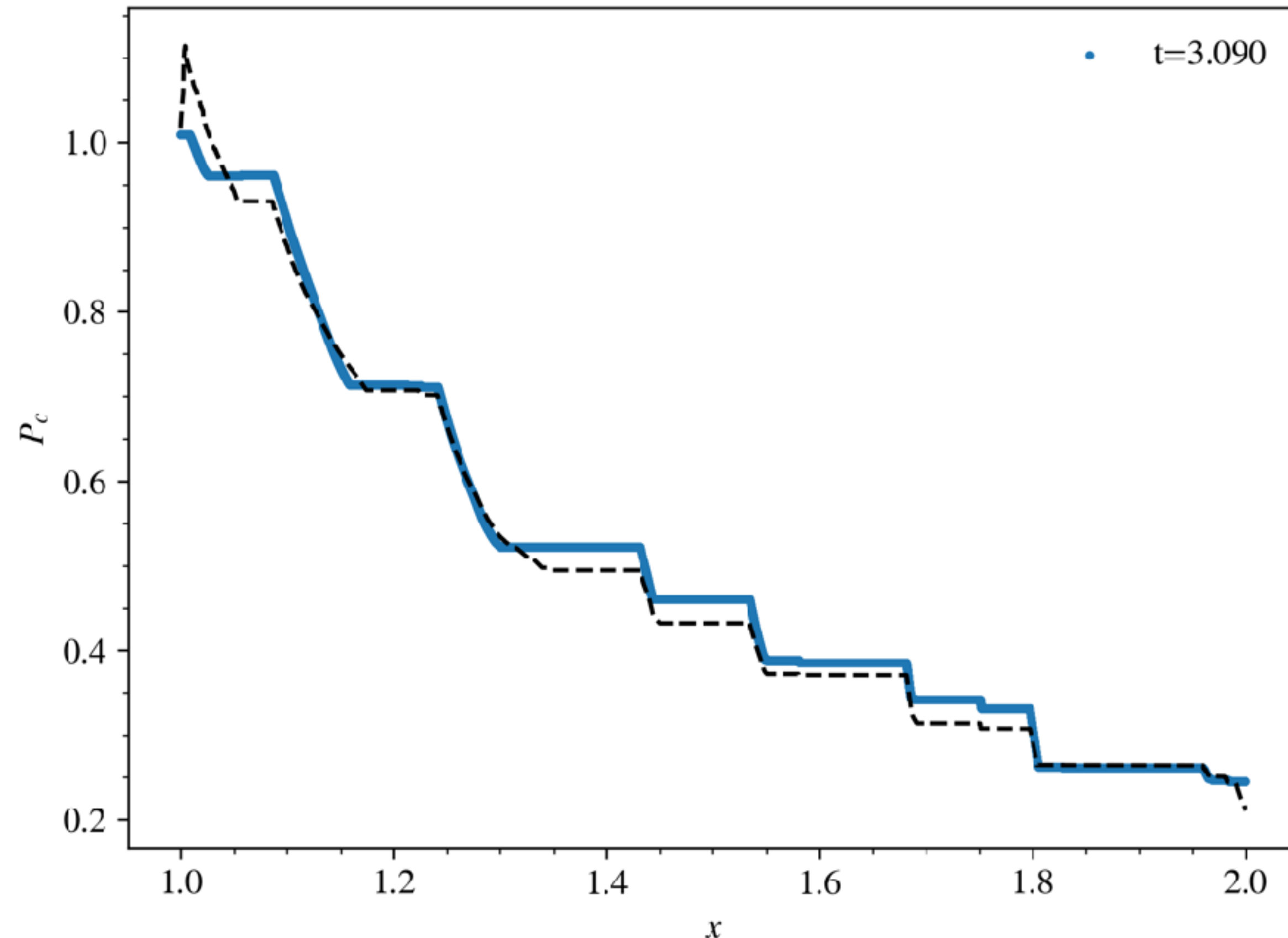
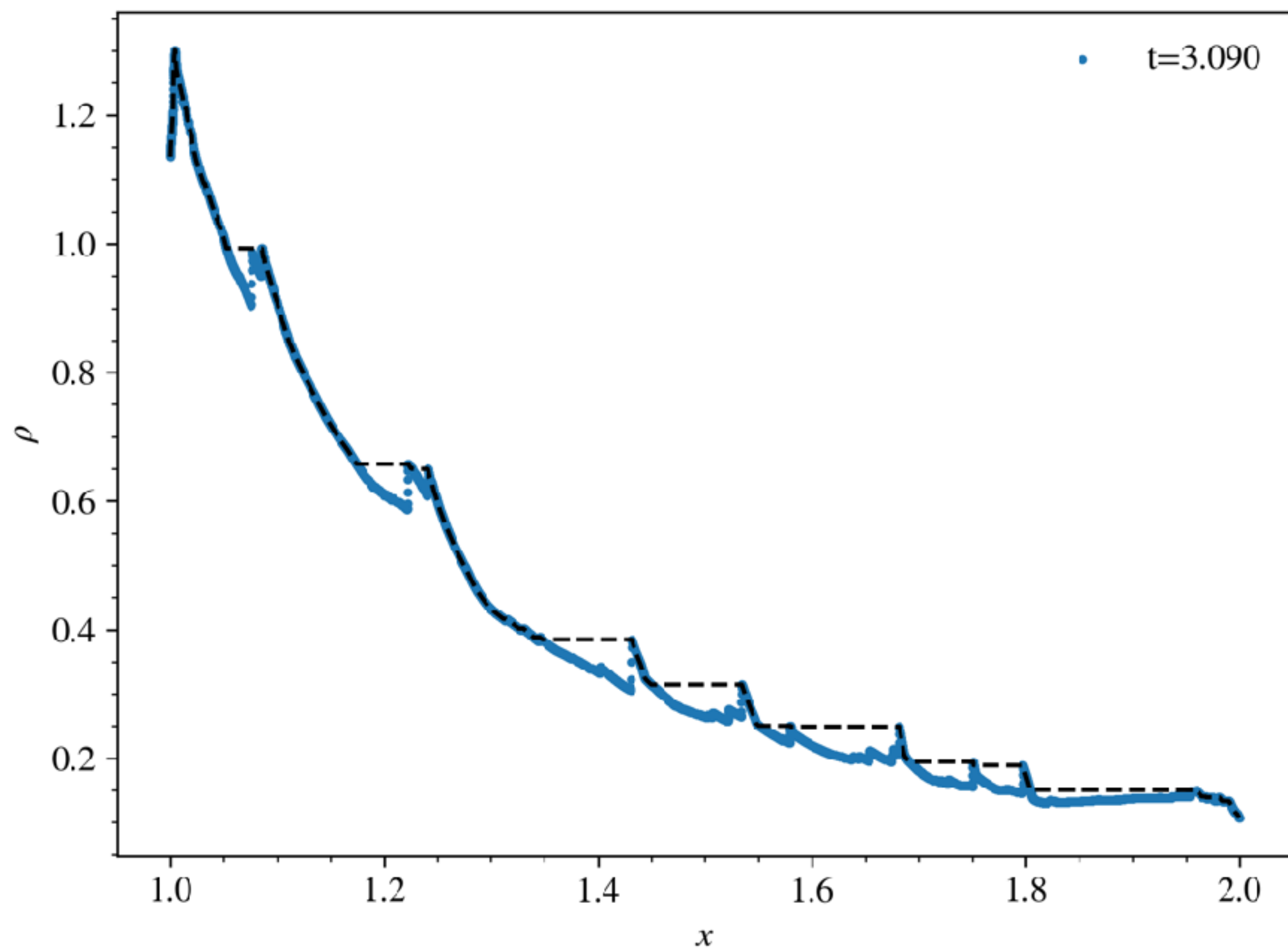
initial position

Lagrangian map
is a
“Devil's staircase”
with fractal
dimension

Reasonable density distribution!

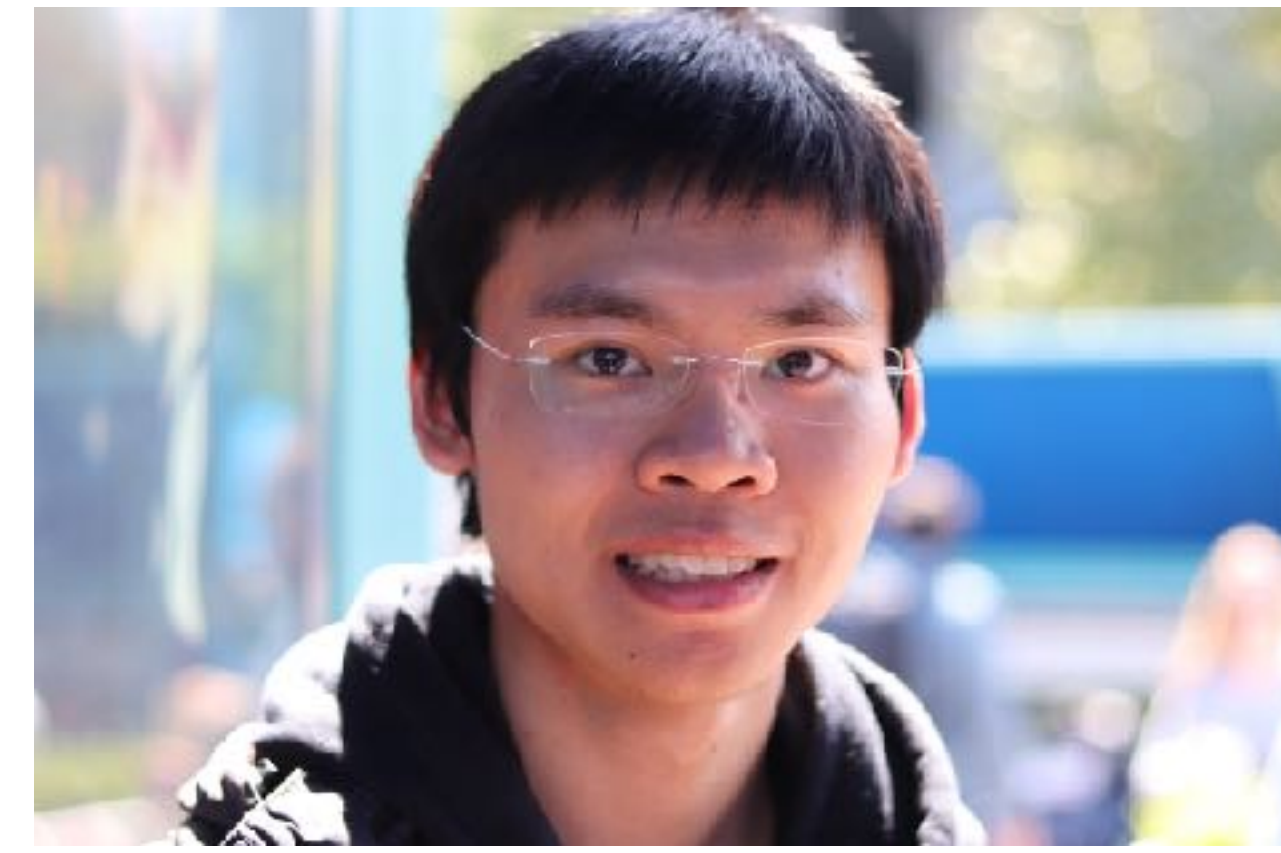
Many beautiful properties, with deep roots
in Legendre transform

Scheme works well

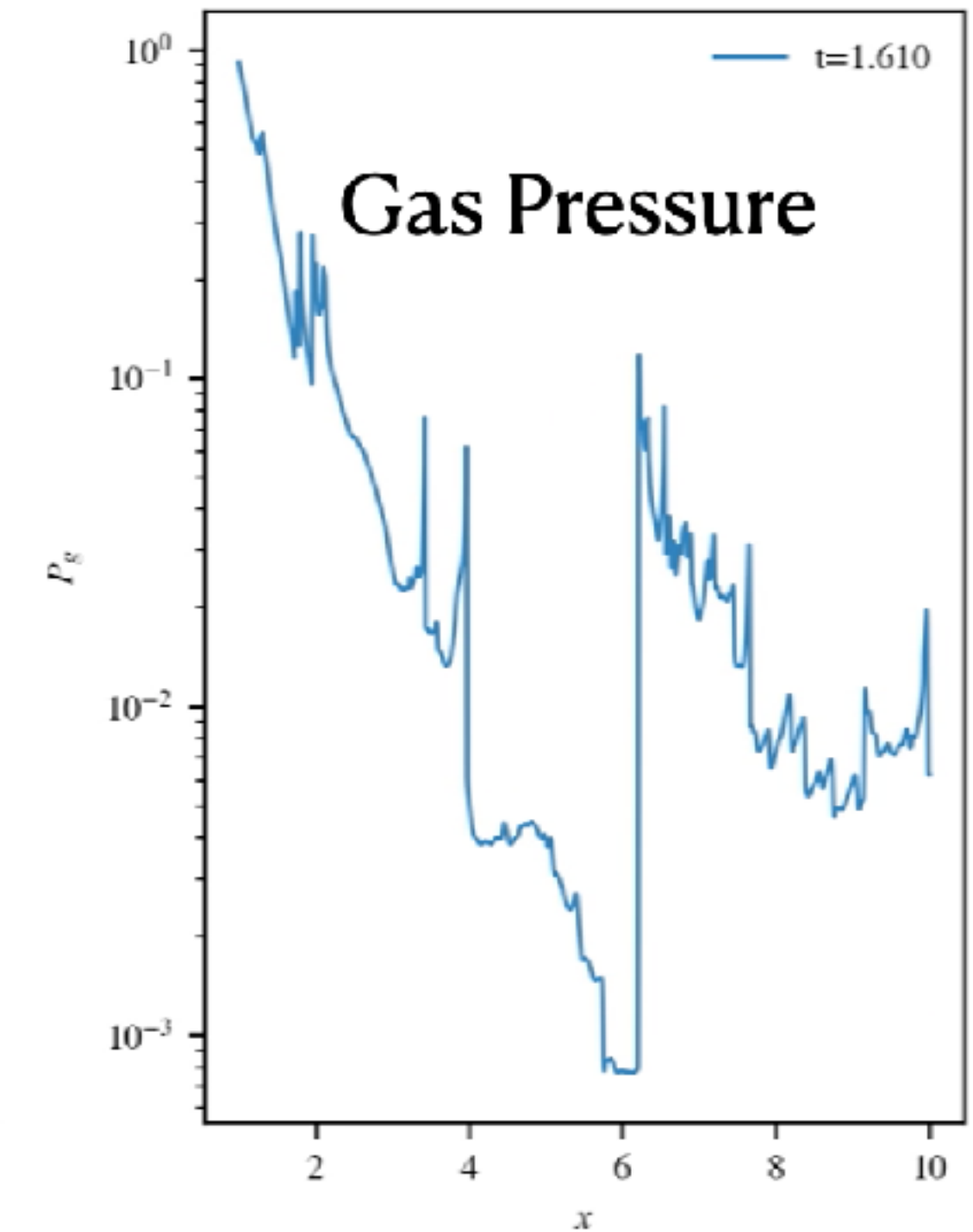
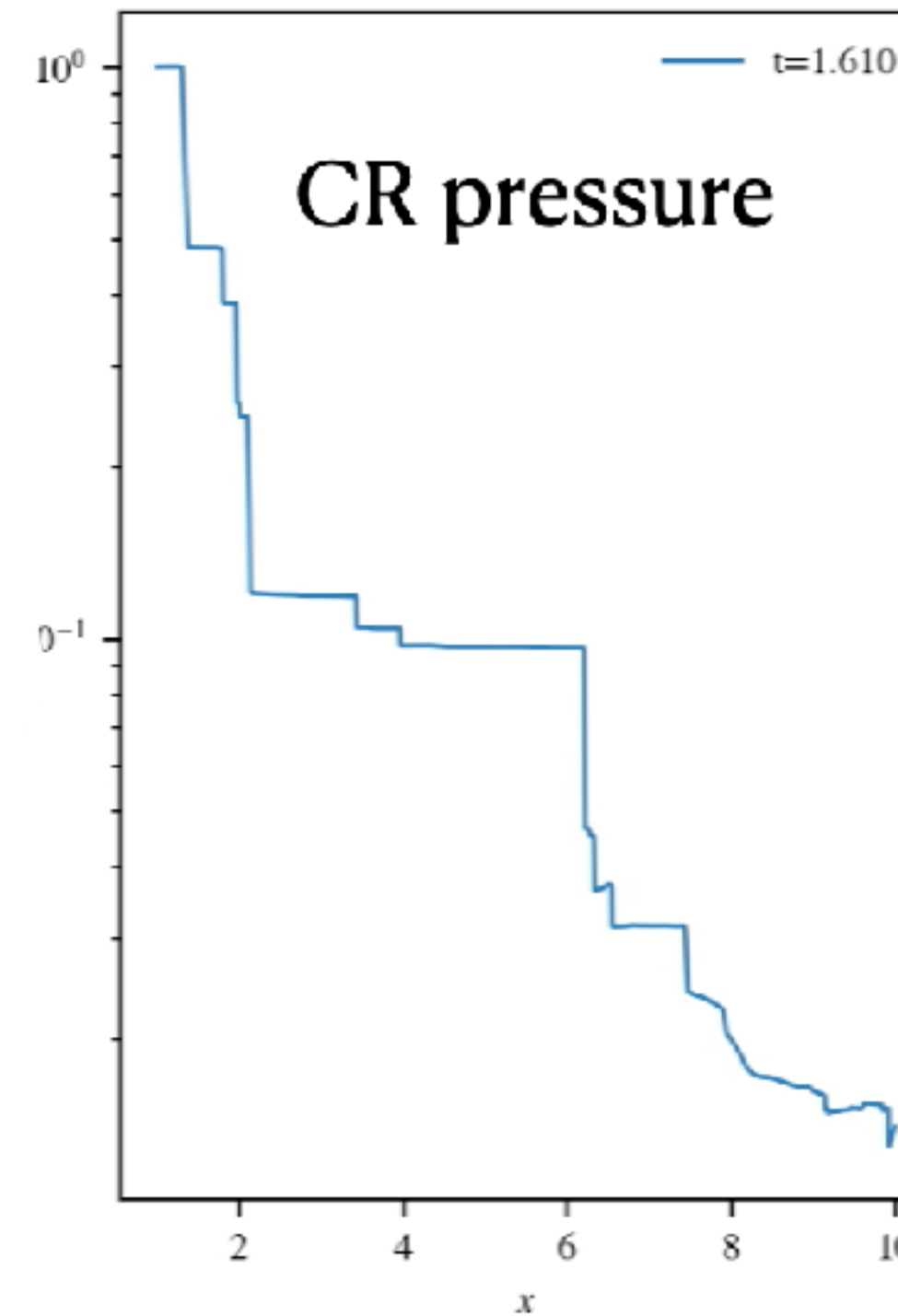


Black dotted line represents convex hull construction

OK, so what?



- Time-averaged momentum and energy transfer appear unchanged
- But strong transient density and velocity fluctuations! (visible w/ FRBs?)
- **Strong spatial and temporal pressure fluctuations** — assuming a fixed pressure will be a poor approximation



Many cool and wacky surprises in store



There are more things in heaven
and earth, Horatio, than are
dreamt of in your philosophy.

- William Shakespeare