



Michael Trott,



Papers discussed:

Espinosa, Muhlleitner, Grojean, Trott arXiv:1202.3697

Espinosa, Muhlleitner, Grojean, Trott arXiv: 1205.6790

Espinosa, Muhlleitner, Grojean, Trott arXiv:1207.1717

Espinosa, Grojean, Sanz, Trott arXiv:1207.7355

Preview:

Espinosa, Muhlleitner, Grojean, Trott

Grojean, Jenkins, Manohar, Trott



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Content (for ref):

Basic fit EFT, limits, χ^2 , 7 TeV data

Invisible width fits, PDF developments

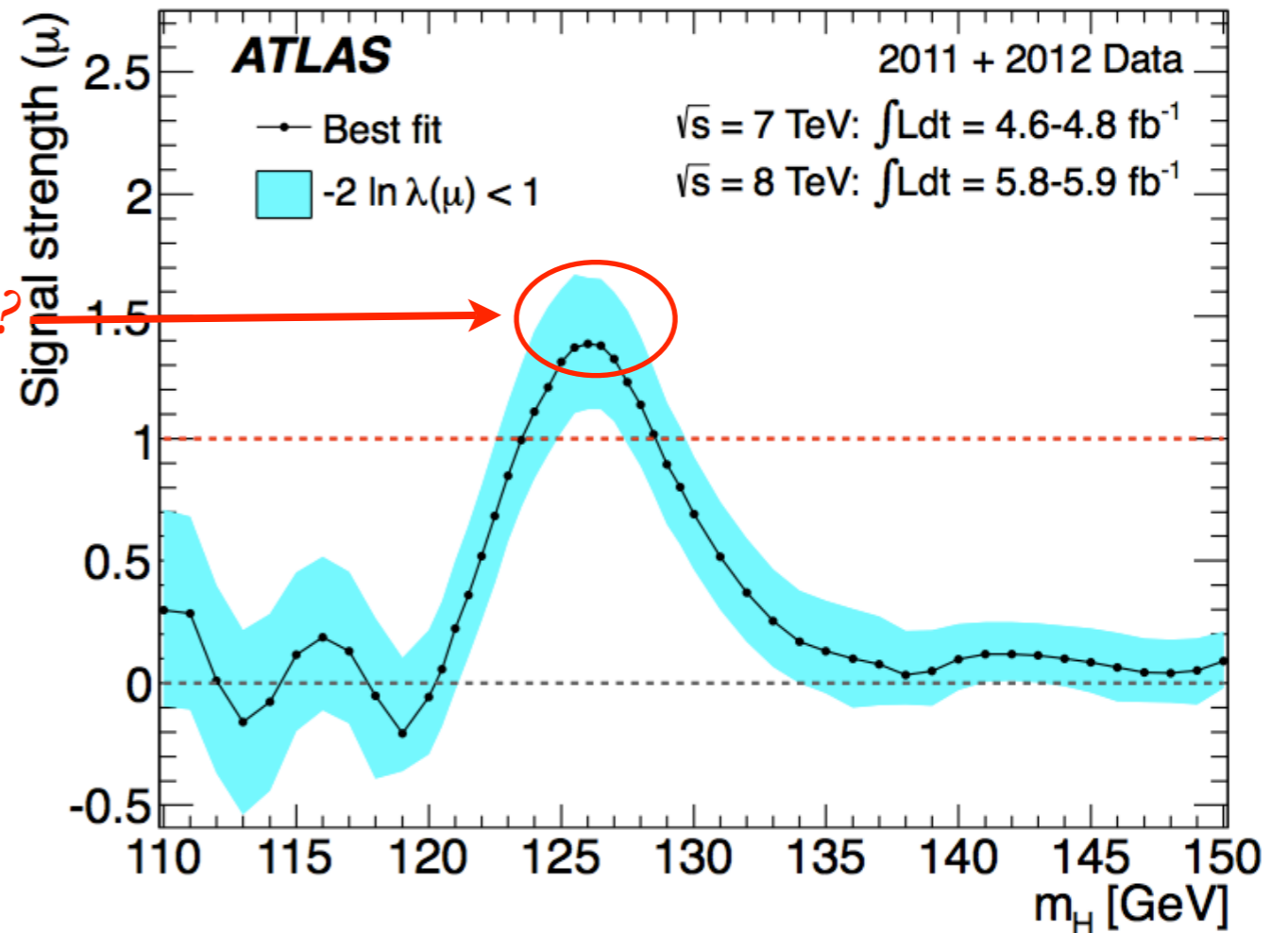
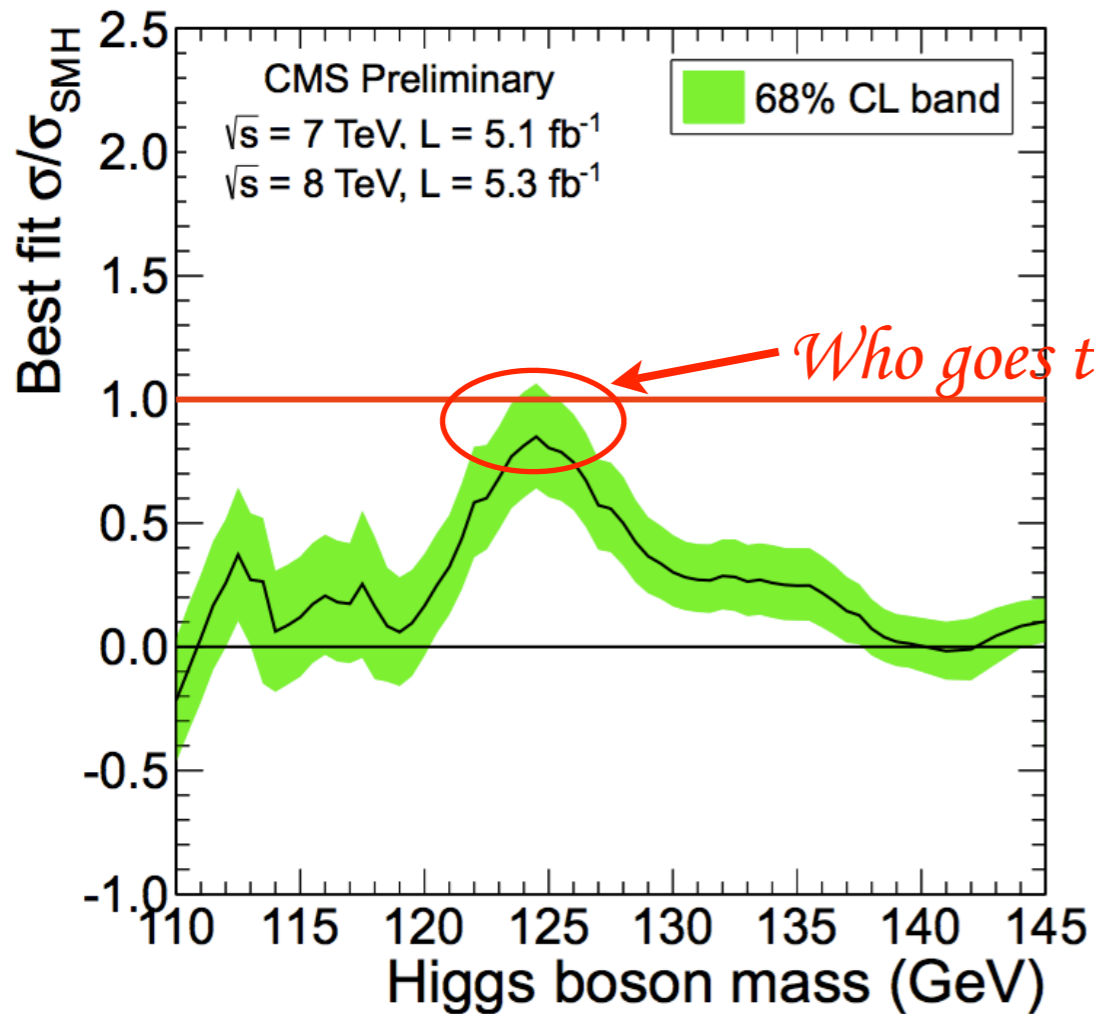
“Tension”, ICHEP data (updates)

NSUSY

yet more data analysis... soon

some loops....sooner

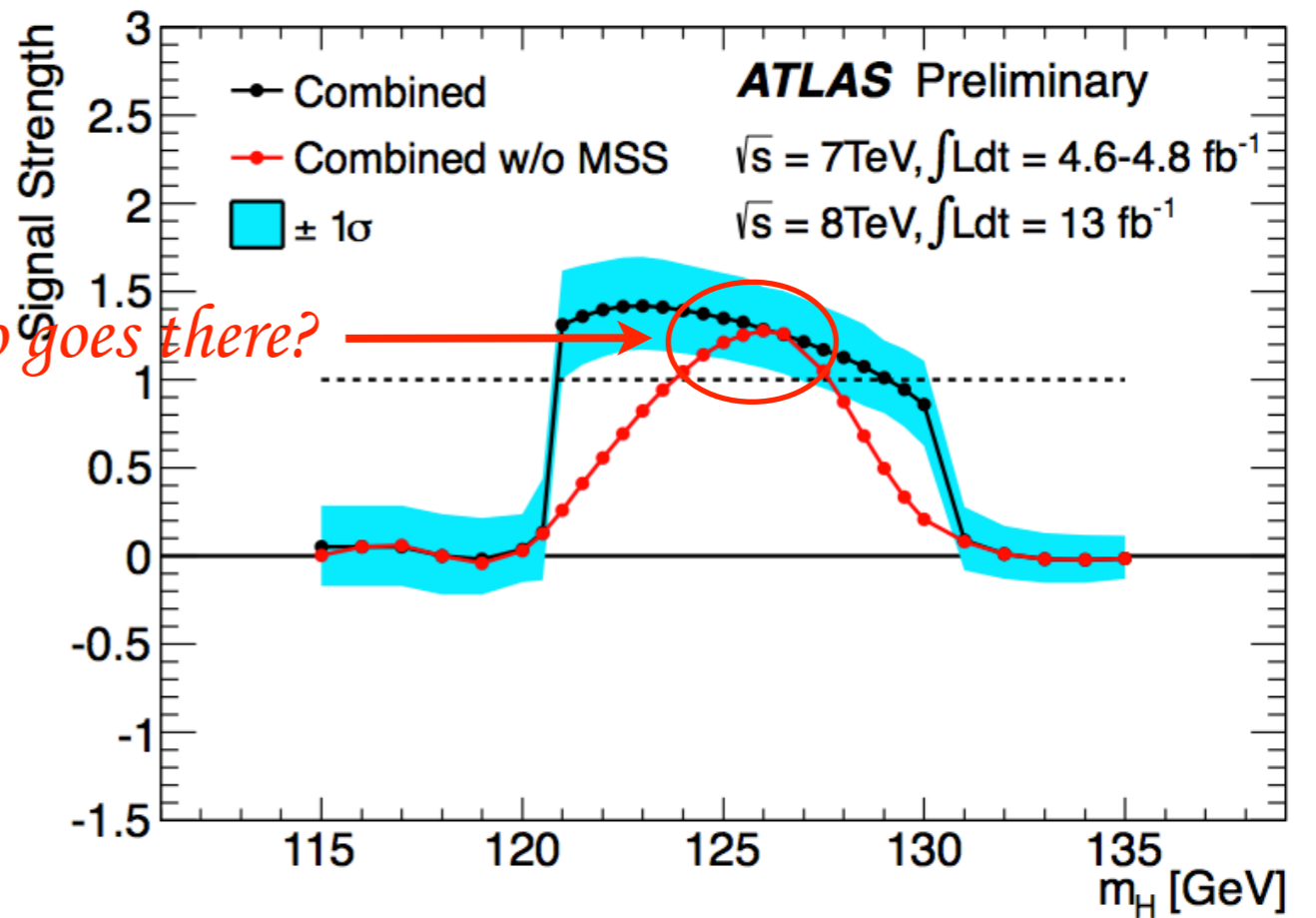
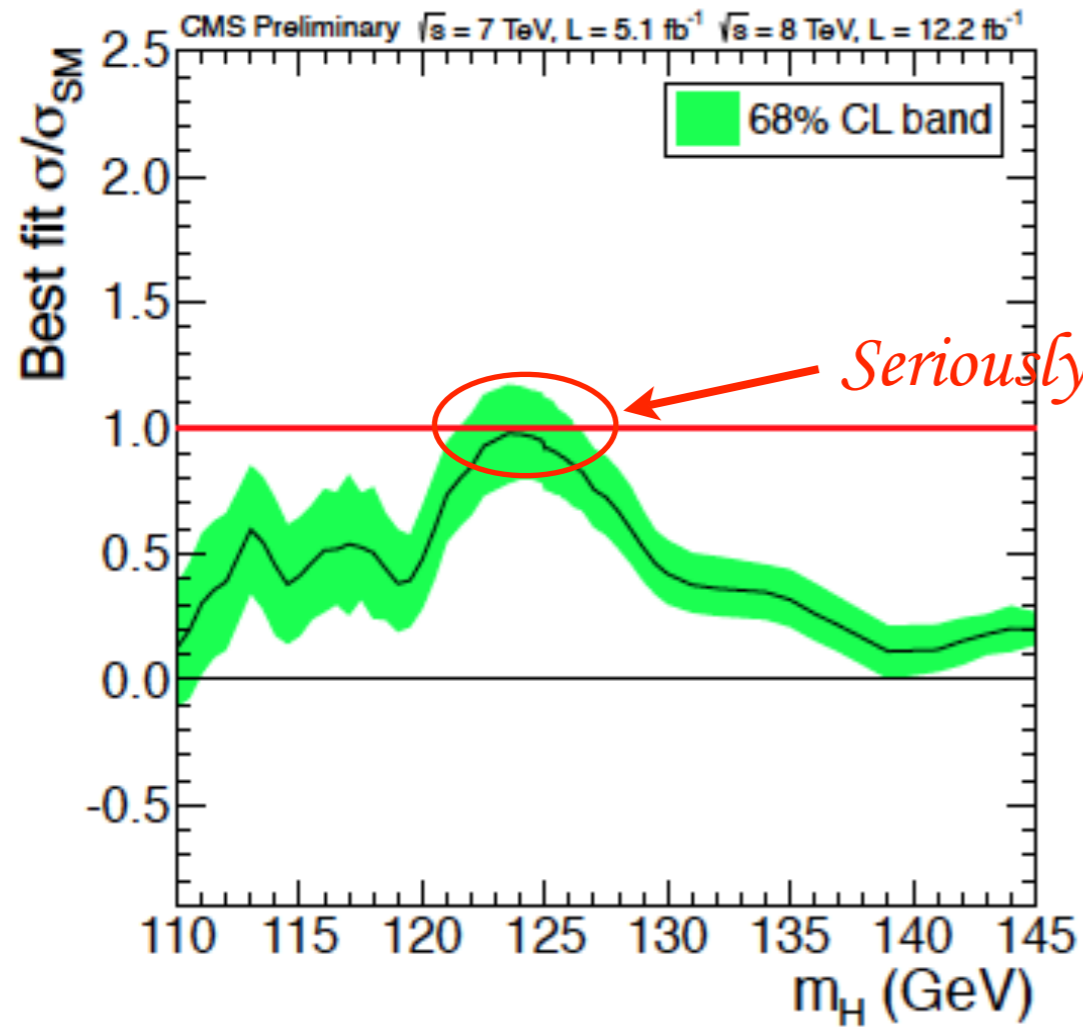
Discovery Physics!



“As a layman, I would say: I think we have it.” - the D.G.

Typical theorist mental state then.

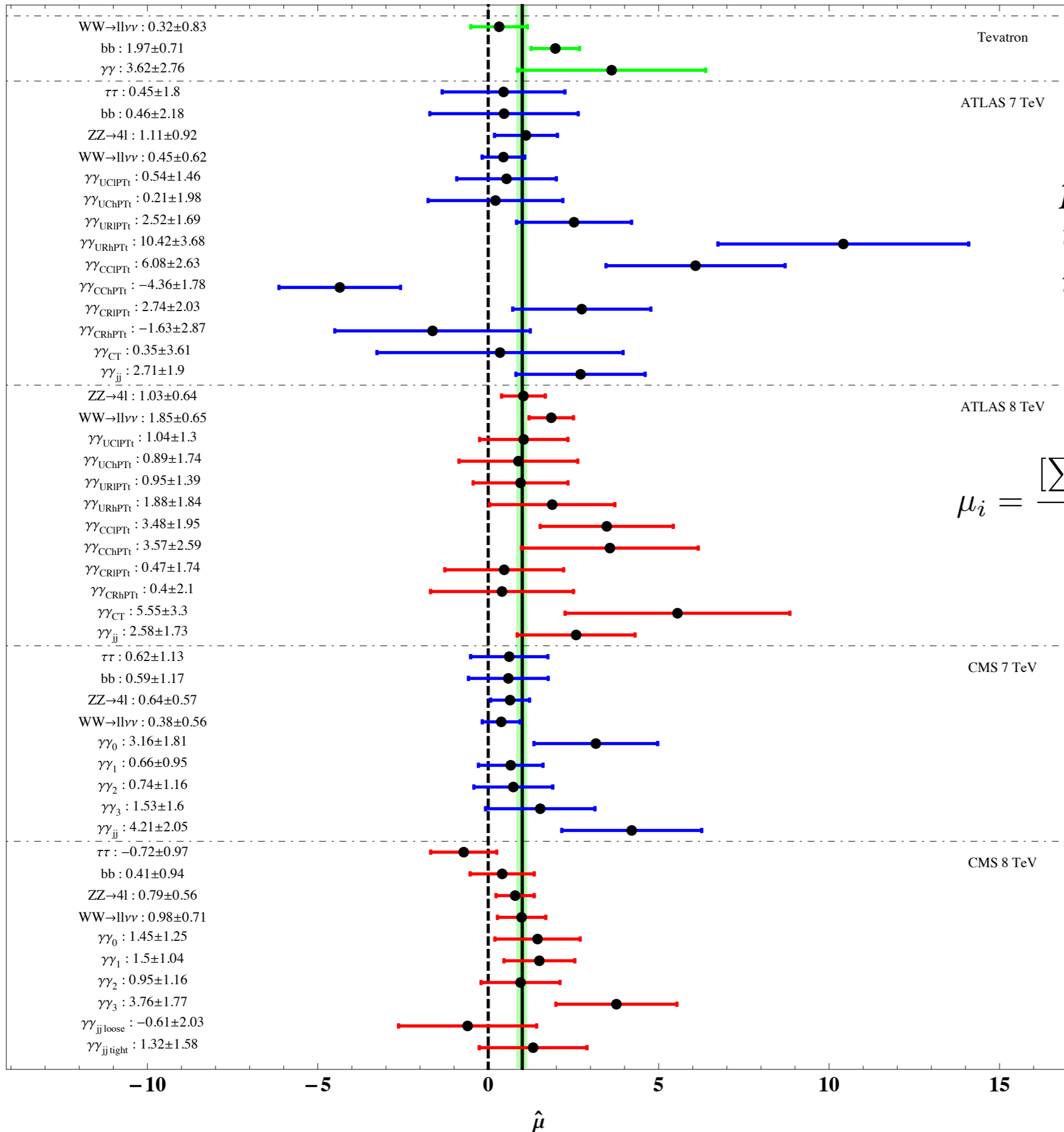
Post-Discovery Physics...



"Yeah, yeah, where is SUSY?"

Typical theorist mental state now.

All Data (48 channels)



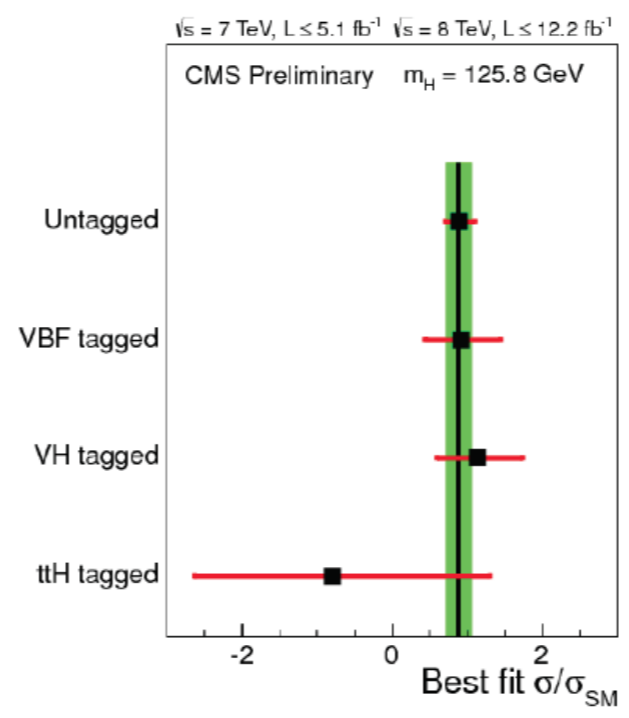
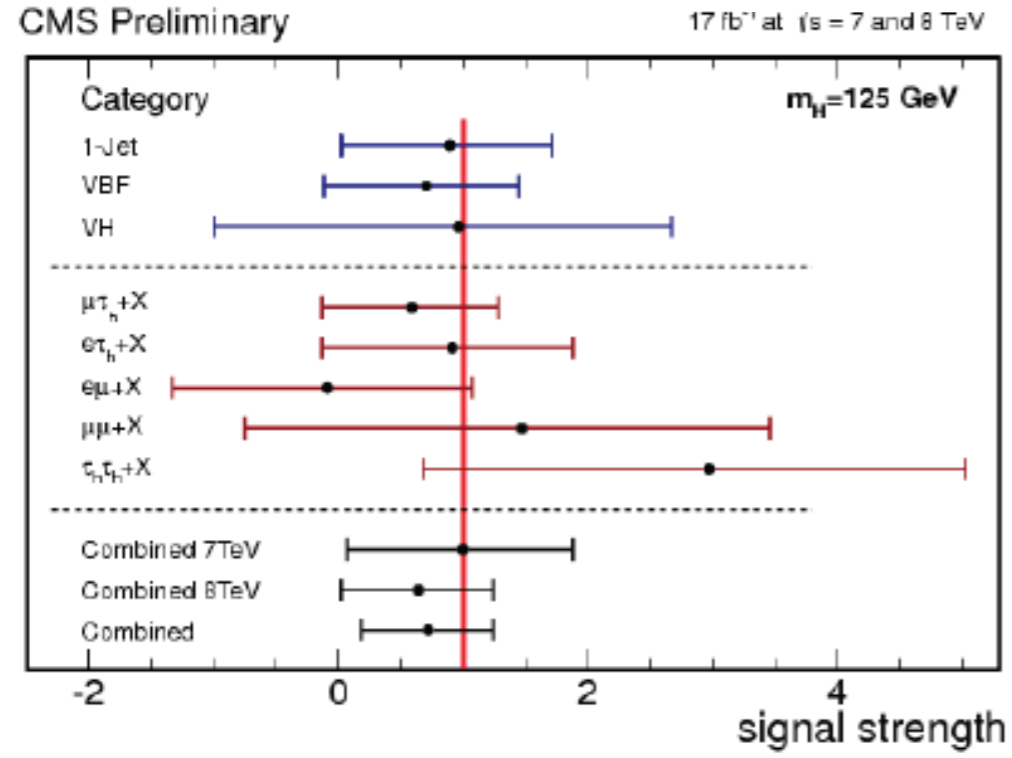
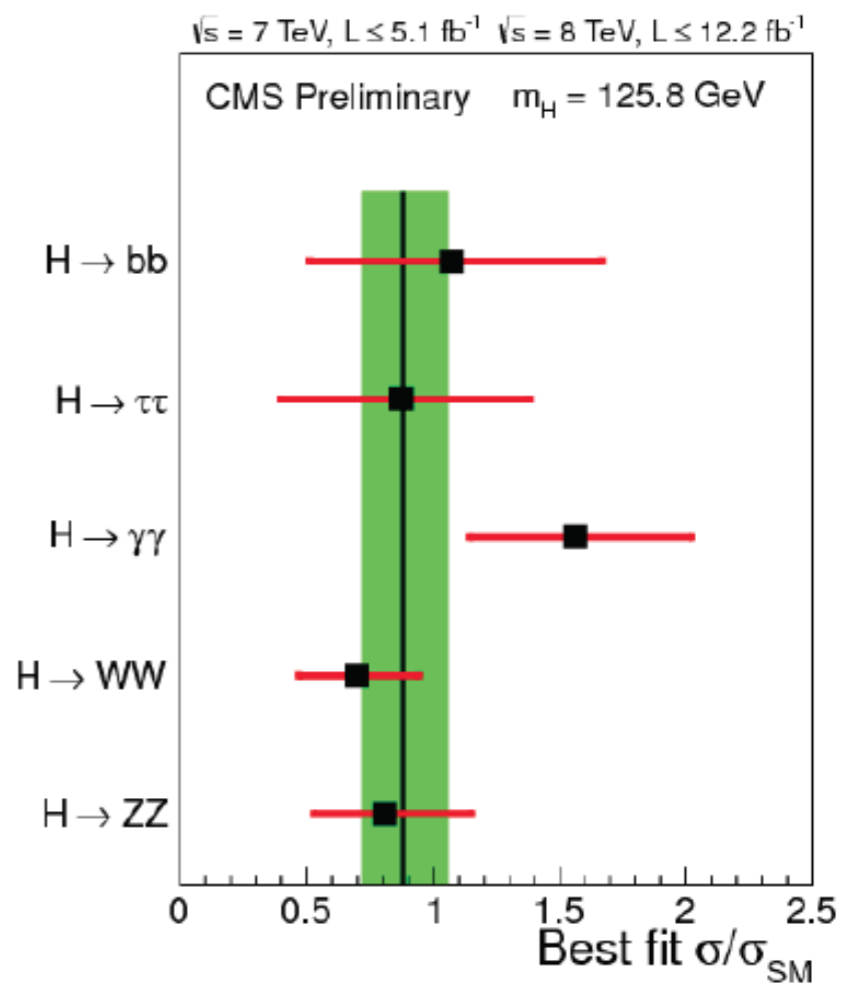
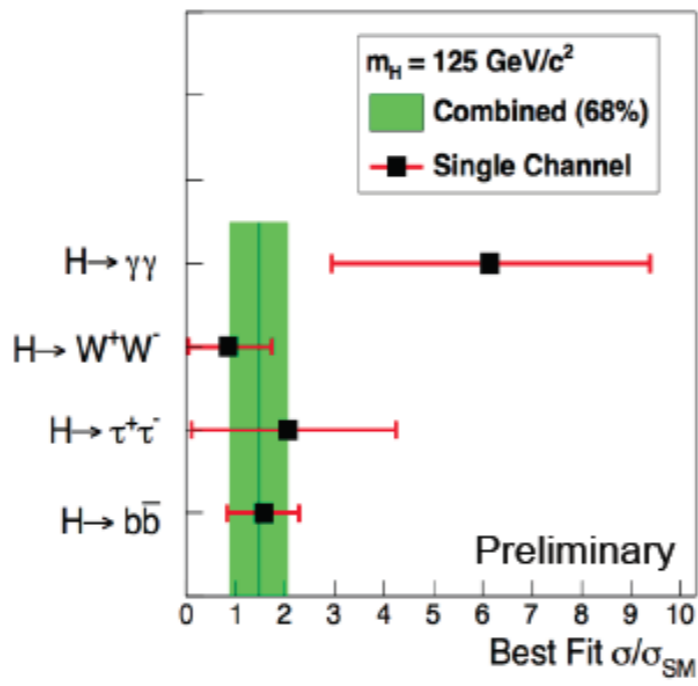
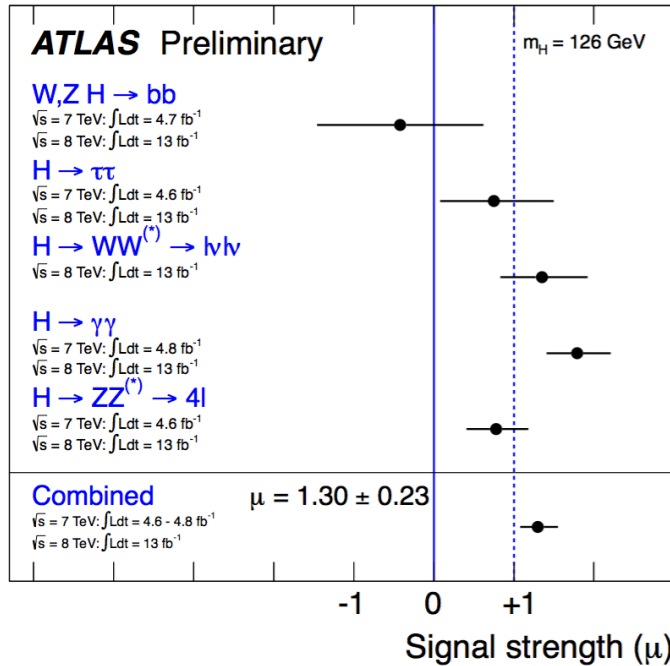
“IT”.

pre-HCP data. Warning the 7 tev has changed in a non public way in some signals.

Signal Strengths.

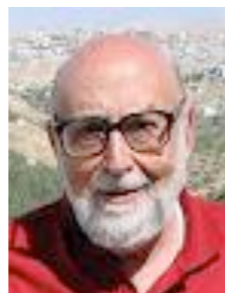
$$\mu_i = \frac{[\sum_j \sigma_{j \rightarrow h} \times \text{Br}(h \rightarrow i)]_{\text{observed}}}{[\sum_j \sigma_{j \rightarrow h} \times \text{Br}(h \rightarrow i)]_{SM}}$$

More on "IT".



WHAT IS "IT"?

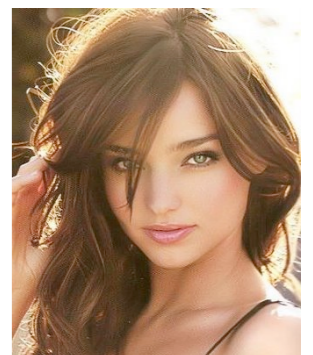
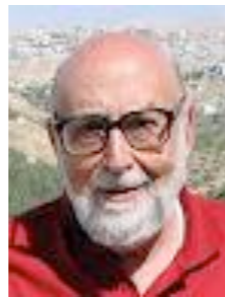
Is it a SM "Higgs"?



Good luck sorting the credit out noble committee!

WHAT IS "IT"?

Is it a SM "Higgs"?



Good luck sorting the credit out noble committee!

Is it a SM "Higgs" and the impact of NP? A 2 for 1?

WHAT IS "IT"?

Is it a SM "Higgs"?



Good luck sorting the credit out noble committee!

Is it a SM "Higgs" and the impact of NP?

Have to assume something. - assume a scalar field - Consider more exotic possibilities AFTER broad scalar EFT attempts fail.



Consistency Problems requiring some New Physics

A scalar field is already strongly implied by the problem to solve.


The massive W, Z indicate that there is a consistency issue at high energies:

$$\mathcal{L}_{eff} = m_W^2 W^+ W^- + \frac{m_Z^2}{2} Z Z + \dots \quad \begin{array}{l} W_L^+ W_L^- \rightarrow W_L^+ W_L^- : A \propto \frac{s}{v^2} \\ \psi \bar{\psi} \rightarrow W_L^+ W_L^- : A \propto \frac{s}{v^2} \end{array}$$

The way to think of the theory we have actually been probing till now is:

$$\mathcal{L} = m_W^2 W W + \frac{m_Z^2}{2} Z Z = \frac{v^2}{4} \text{Tr} (D^\mu \Sigma^\dagger D_\mu \Sigma)$$

Goldstones of broken $SU_L(2) \times SU_R(2)/SU_V(2)$ give mass to W and Z

 *grouped in the nonlinear chiral EW Lagrangian as $\Sigma = \exp(i \sigma_a \pi^a / v)$*

Consistency Problems requiring some New Physics

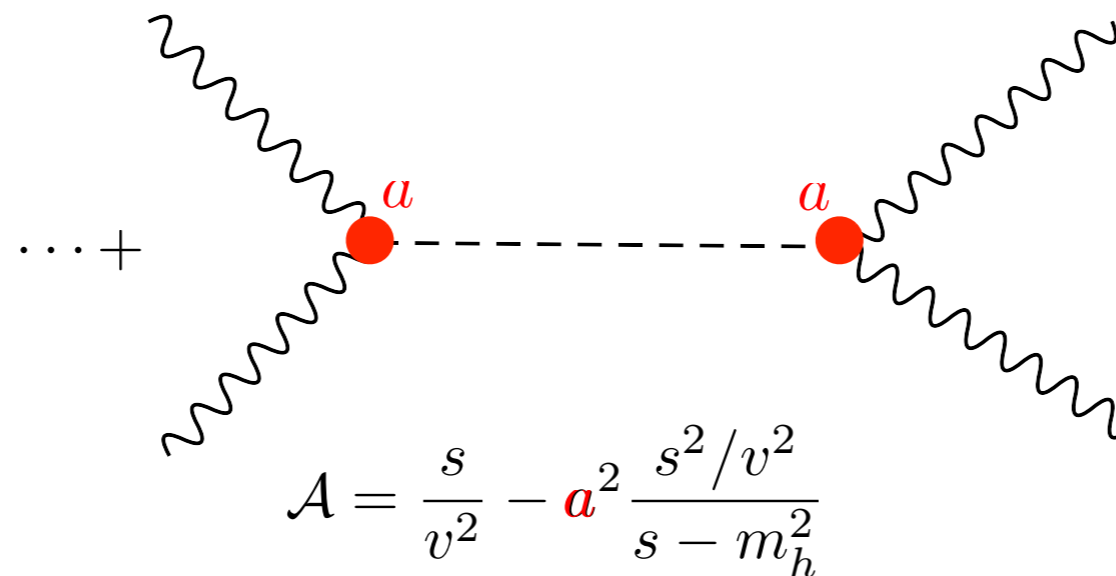
A scalar field is already strongly implied by the problem to solve.

The massive W, Z indicate that there is a consistency issue at high energies:

$$\mathcal{L}_{eff} = m_W^2 W^+ W^- + \frac{m_Z^2}{2} Z Z + \dots$$

$$W_L^+ W_L^- \rightarrow W_L^+ W_L^- : A \propto \frac{s}{v^2}$$

$$\psi \bar{\psi} \rightarrow W_L^+ W_L^- : A \propto \frac{s}{v^2}$$



*Couplings within 10%
of the SM,
cut off scale 7 tev...*

Cut off scale of the EFT: $\Lambda = 4 v \pi$..raised to... $\Lambda = 4 v \pi / \sqrt{|1 - a^2|}$

Fairly suggestive that a scalar field of some form will be involved in the UV completion.

Nonlinear Chiral EW Lagrangian + scalar

Leading terms in the EFT, there is a systematic derivative expansion to exploit:

$$\begin{aligned} \mathcal{L} &= \frac{1}{2}(\partial_\mu h)^2 - V(h) + \frac{v^2}{4} \text{Tr}(D_\mu \Sigma^\dagger D^\mu \Sigma) \left[1 + 2 a_{W,Z} \frac{h}{v} + b_{Z,W} \frac{h^2}{v^2} + b_{3,Z,W} \frac{h^3}{v^3} + \dots \right], \\ &\quad - \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c_i^{u,d} \frac{h}{v} + c_{2,j}^{u,d} \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u & u_R^j \\ y_{ij}^d & d_R^j \end{pmatrix} + h.c., \\ V(h) &= \frac{1}{2} m_h^2 h^2 + \frac{d_3}{6} \left(\frac{3 m_h^2}{v} \right) h^3 + \frac{d_4}{24} \left(\frac{3 m_h^2}{v^2} \right) h^4 + \dots \end{aligned}$$

Also higher dimensional operators: (hats -dual fields)

$$\begin{aligned} \mathcal{L}_{HD}^5 &= -\frac{c_g g_3^2}{32 \pi^2 v} h G_{\mu\nu}^A G^{A\mu\nu} - \frac{c_W g_2^2}{32 \pi^2 v} h W_{\mu\nu}^a W^{a\mu\nu} - \frac{c_B g_1^2}{32 \pi^2 v} h B_{\mu\nu} B^{\mu\nu} \\ &\quad - \frac{\hat{c}_g g_3^2}{32 \pi^2 v} h \hat{G}_{\mu\nu}^A G^{A\mu\nu} - \frac{\hat{c}_W g_2^2}{32 \pi^2 v} h \hat{W}_{\mu\nu}^a W^{a\mu\nu} - \frac{\hat{c}_B g_1^2}{32 \pi^2 v} h \hat{B}_{\mu\nu} B^{\mu\nu} + \mathcal{O}(h^2) \end{aligned}$$

Also higher dimensional derivative operators in the chiral EFT...

TOO MANY DAMN PARAMETERS!

Nonlinear Chiral EW Lagrangian + scalar

EFT gives model independence \mathcal{L} is a systematically improvable Lagrangian approach. ALSO LETS ONE USE SYMMETRY TO REDUCE PARAMETERS.

Assuming custodial sym breaking as in the SM and consistent with MFV:

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 + \frac{v^2}{4}\text{Tr}(D_\mu \Sigma^\dagger D^\mu \Sigma) \left[1 + 2a \frac{h}{v}\right] - \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c^{u,d} \frac{h}{v}\right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + h.c.,$$

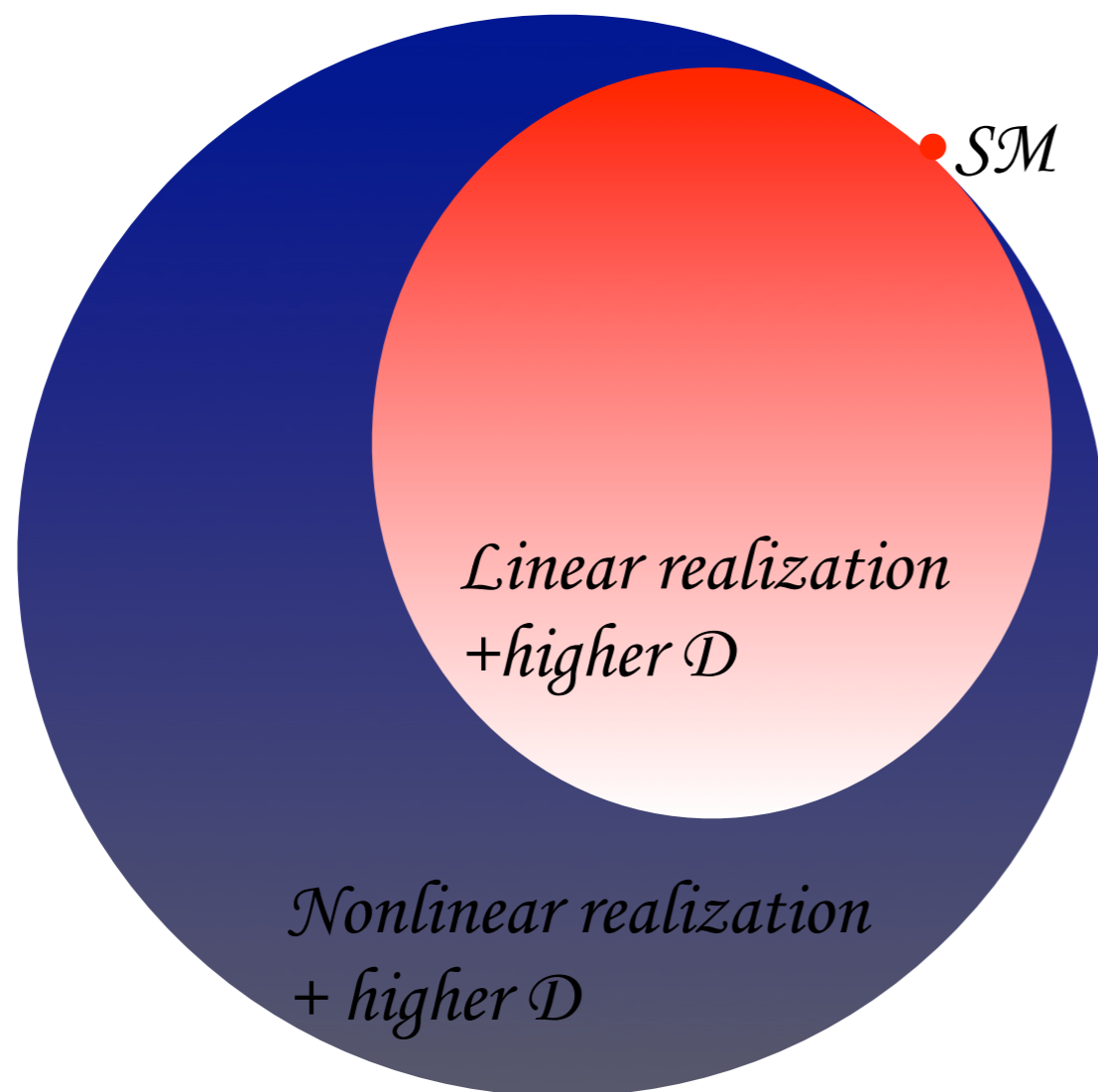
Also higher dimensional operators: - assuming no large BSM CP violation

$$\mathcal{L}_{HD}^5 = -\frac{c_g g_3^2}{32 \pi^2 v} h G_{\mu\nu}^A G^{A\mu\nu} - \frac{c_W g_2^2}{32 \pi^2 v} h W_{\mu\nu}^a W^{a\mu\nu} - \frac{c_B g_1^2}{32 \pi^2 v} h B_{\mu\nu} B^{\mu\nu}$$

Reasonable coupling space, can draw physical conclusions for sym theories with current data. Still have degeneracies. **ONLY THE START OF THIS PROGRAM.**

What is the right approach to use?

It is obvious that we will be doing Higgs Effective Field theory for LHC and linear colliders. But WHICH EFT to use?



We want a robust field theory generalization of the SM scalar predictions to use to prove FROM THE DATA (not our brains or guts) that the SM is preferred.

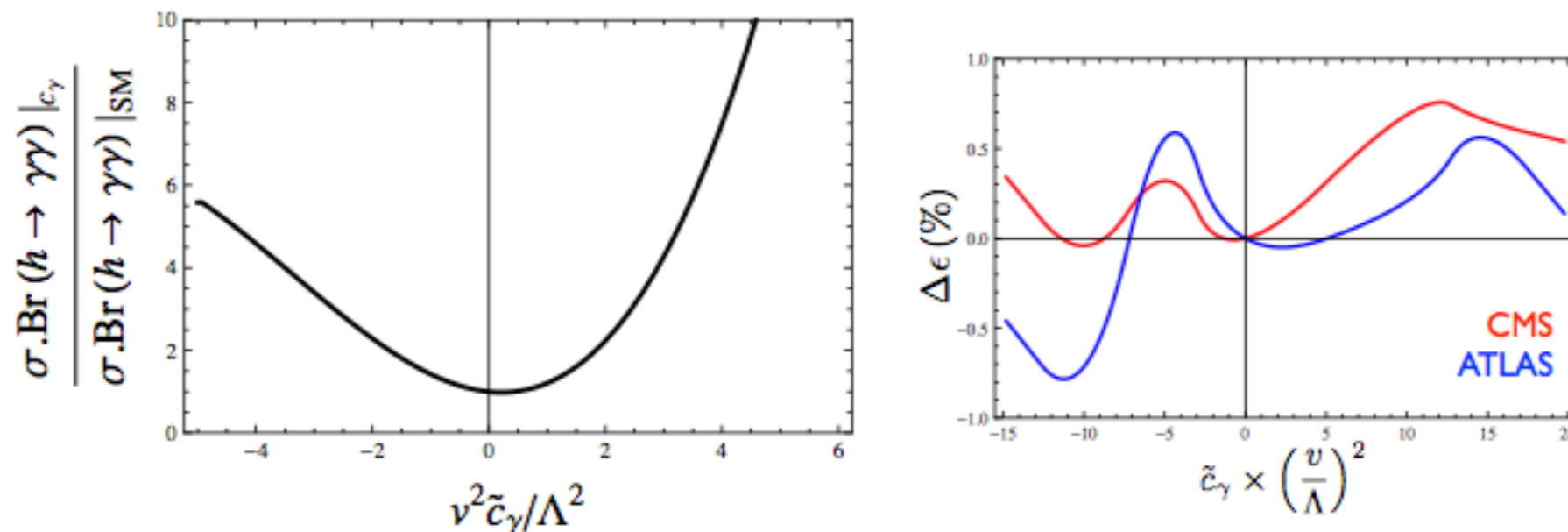
What is the right approach to use?

This is the framework that leads to generalizing the SM predictions with tree level rescalings of the cross section and branching ratios:

$$\mu_i = \frac{[\sum_j \sigma_{j \rightarrow h} \times \text{Br}(h \rightarrow i)]_{\text{observed}}}{[\sum_j \sigma_{j \rightarrow h} \times \text{Br}(h \rightarrow i)]_{\text{SM}}}, \quad \chi^2(\mu_i) = \sum_{i=1}^{N_{ch}} \frac{(\mu_i - \hat{\mu}_i)^2}{\sigma_i^2}$$

This modifies $\mu_{SM}^i \rightarrow \mu^i(a, c)$ but what about efficiency corrections?

Event rates will only change if a kinematic distribution is changed significantly we have checked that in a number of cases as well eff corrections can be safely neglected:



History of a discovery in Fit space

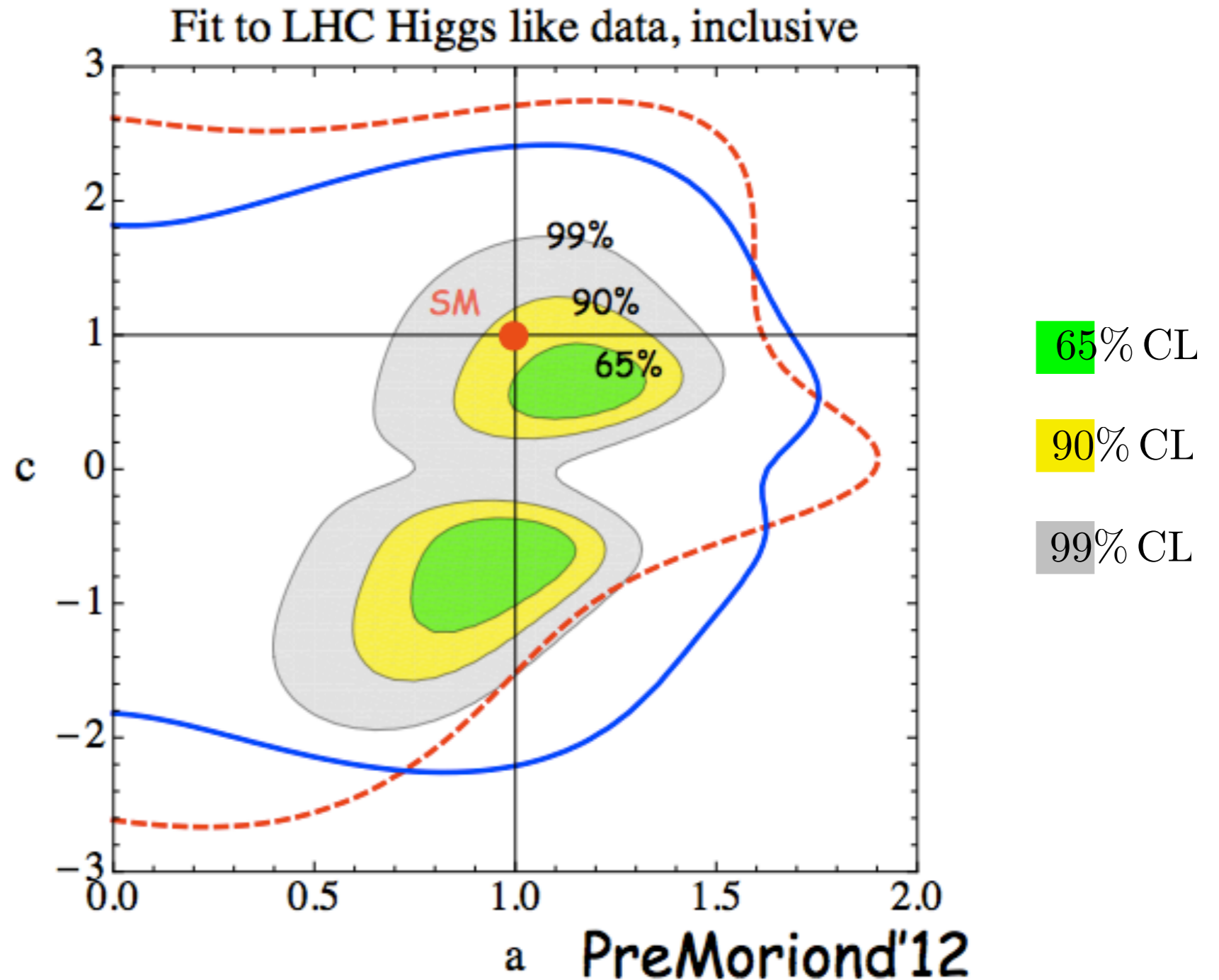
SM 82%CL
away from
best fit point

Two minima:

$(a,c)=(1.13,0.58)$
 $\chi^2=2.86$

$(a,c)=(0.96,-0.64)$
 $\chi^2=1.96$

- Atlas 95%CL exclusion
- CMS 95%CL exclusion
- Tevatron 95%CL exclusion
- combined 95%CL exclusion



Espinosa, Grojean, Muhlleitner, Trott
JHEP 1205 (2012) 097 arxiv:1202.3697

History of a discovery in Fit space

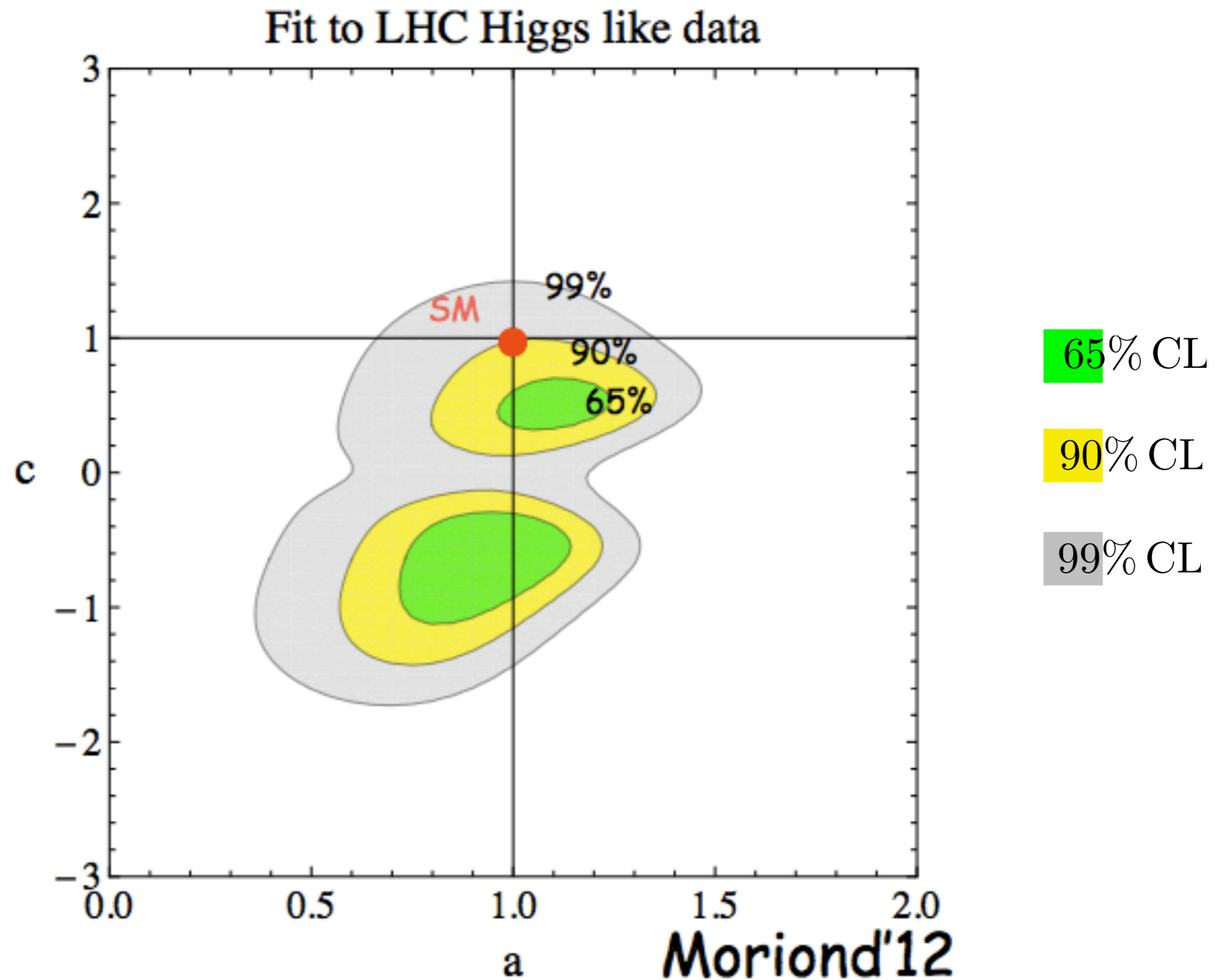
SM 88%CL
away from
best fit point
($\sim 2\sigma$)

Two minima:

(a,c)=(1.18,0.55)
 $\chi^2=7.5$

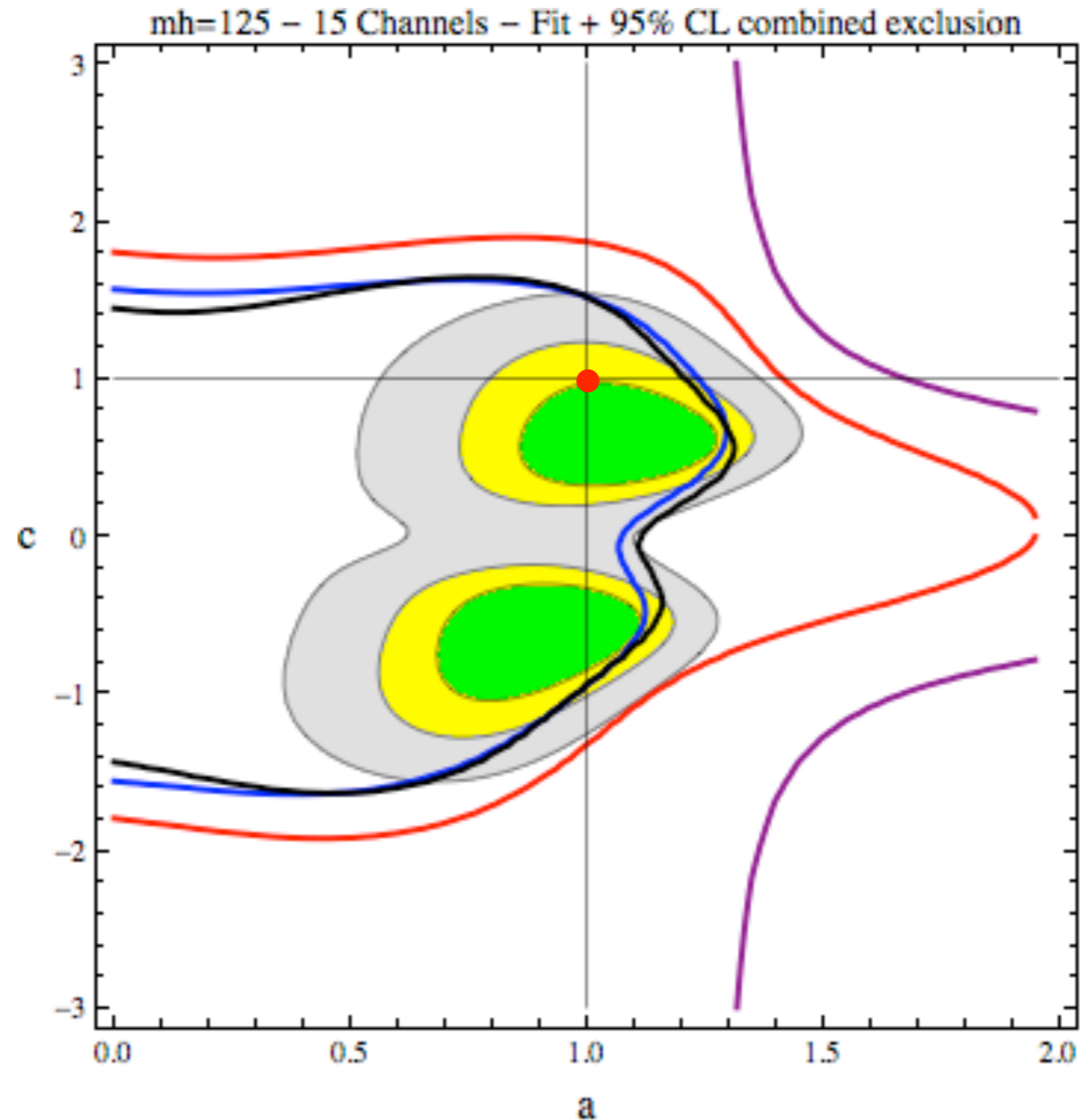
(a,c)=(0.99,-0.64)
 $\chi^2=6.3$

- Atlas 95%CL exclusion
- CMS 95%CL exclusion
- Tevatron 95%CL exclusion
- combined 95%CL exclusion



Espinosa,Grojean,Muhlleitner,Trott
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History of a discovery in Fit space



65% CL

90% CL

99% CL

--- Atlas 95%CL exclusion

— CMS 95%CL exclusion

— Tevatron 95%CL exclusion

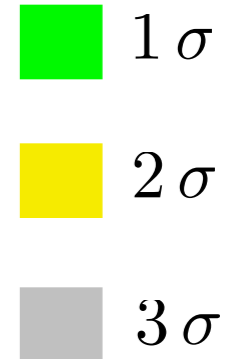
— combined 95%CL exclusion

Post-Moriond/Pre-ICHEP

Espinosa, Grojean, Muhlleitner, Trott
JHEP 1205 (2012) 097 arxiv:1202.3697

History of a discovery in Fit space

Notice scale change, errors coming down!

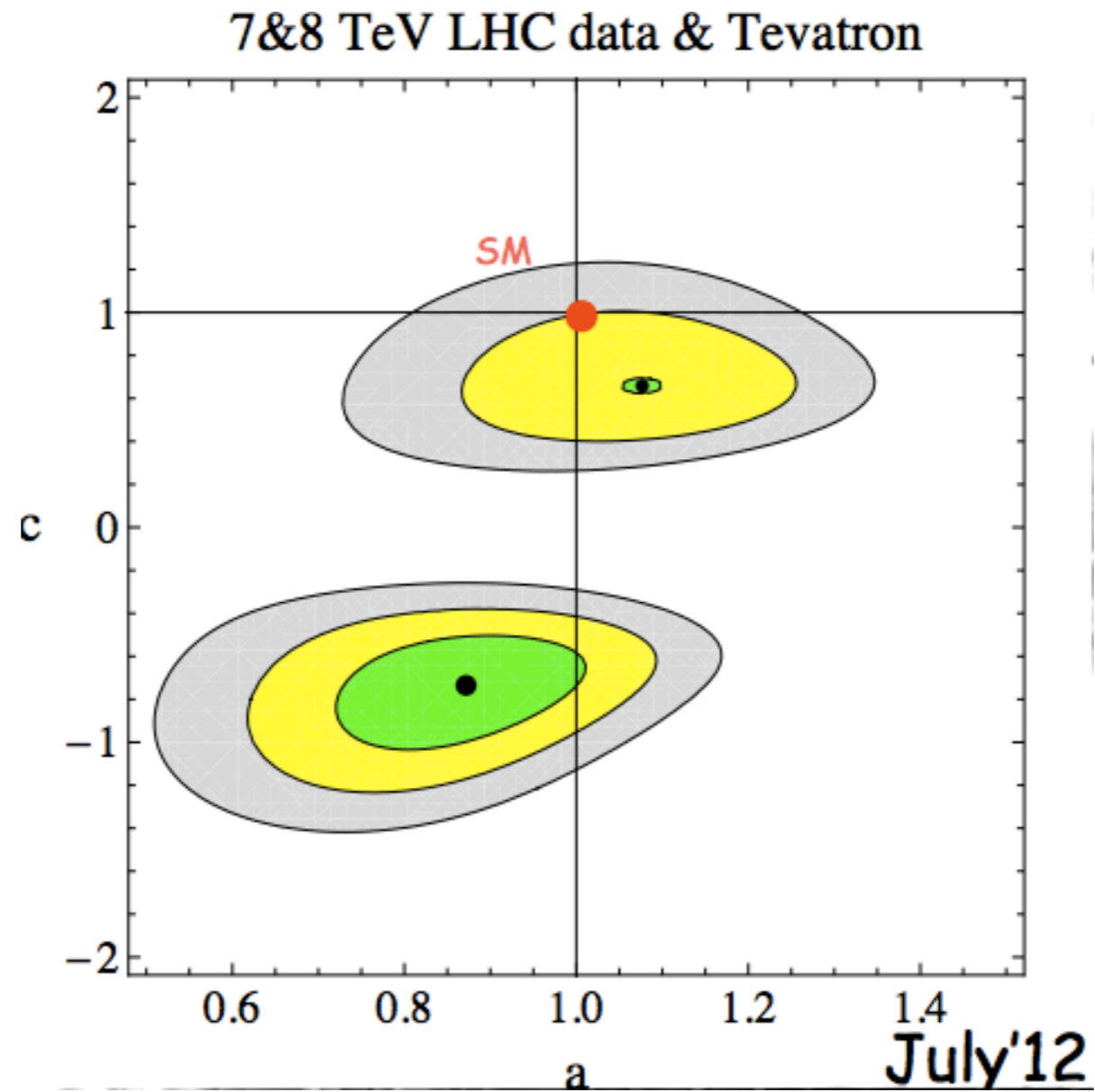


SM 93%CL
away from
best fit point
($\sim 1.8\sigma$)

Two minima:

(a,c)=(0.86,-0.64)
 $\chi^2=41$

(a,c)=(1.05,0.63)
 $\chi^2=41.4$



Need to rescale
7, 8 TeV data
independently.

Data released day of
largely 7+8 TeV
combinations.

- Atlas 95%CL exclusion
- CMS 95%CL exclusion
- Tevatron 95%CL exclusion
- combined 95%CL exclusion

Data day of discovery (released).

History of a discovery in Fit space

**SM 88%CL
away from
best fit point**

Two minima:

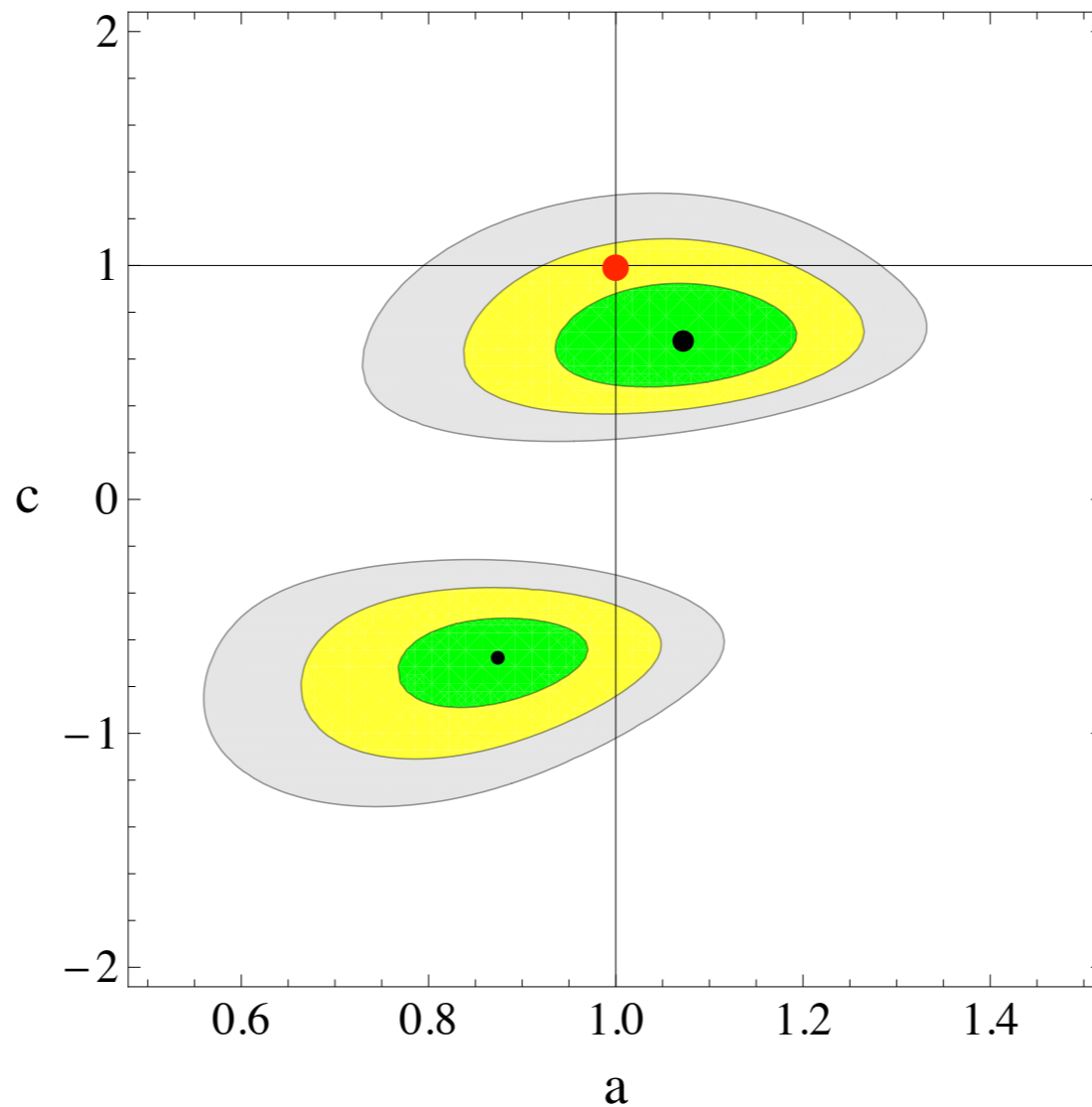
$(a,c)=(1.07, 0.68)$
 $\chi^2=43.6$

$(a,c)=(0.87,-0.68)$
 $\chi^2=44.3$

- - - Atlas 95%CL exclusion
- CMS 95%CL exclusion
- Tevatron 95%CL exclusion
- combined 95%CL exclusion

Notice scale change, errors coming down!

7&8 TeV LHC data & Tevatron



- 1σ
- 2σ
- 3σ

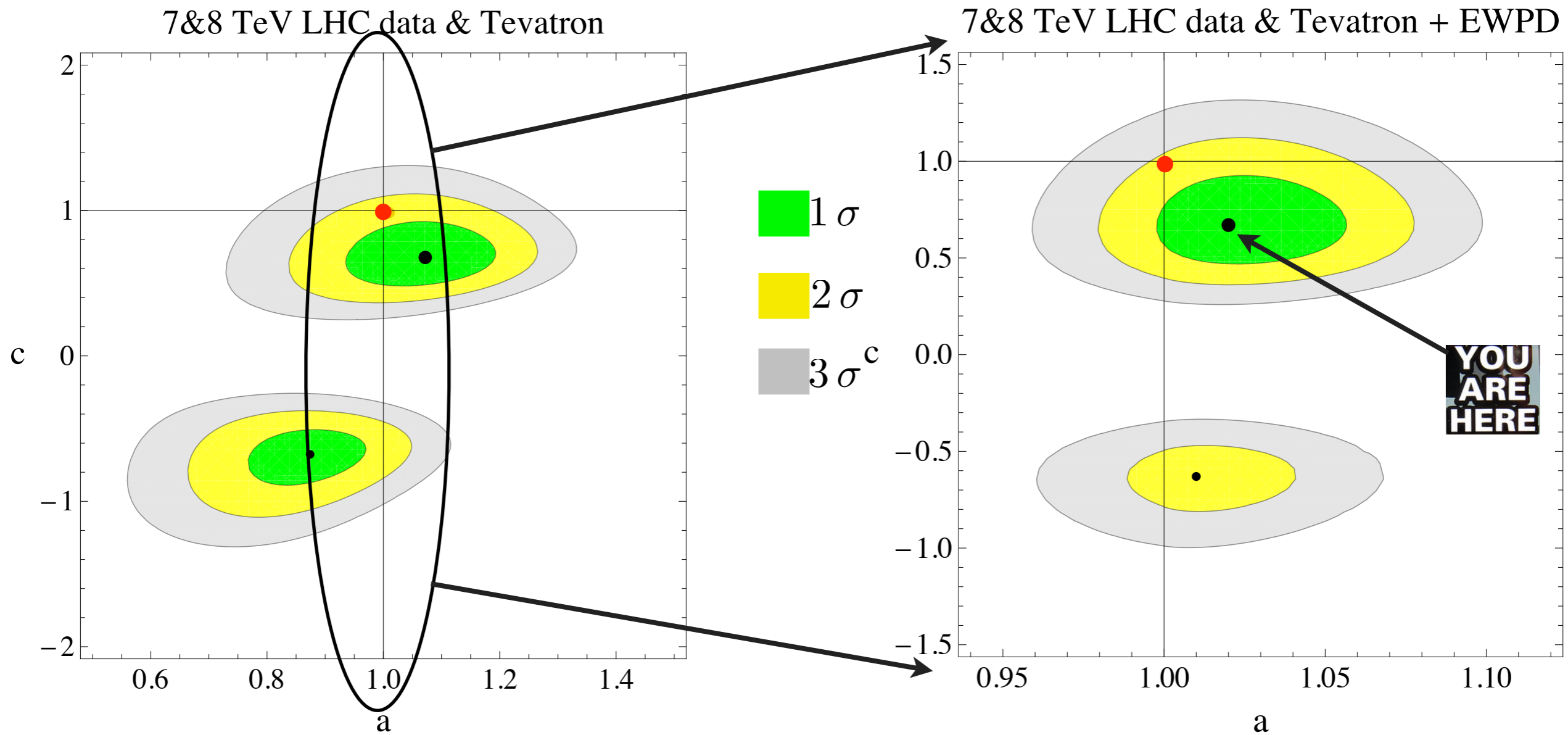
*Most 7,8
data now split.*

*Atlas photon
subcategories
and WW largest
changes.*

Post ICHEP Data - no HCP.

Espinosa, Grojean, Muhlleitner, Trott
JHEP 1205 (2012) 097 arxiv:1207.1717

Add in EWPD to the Fit



Espinosa, Grojean, Muhlleitner, Trott
 JHEP 1205 (2012) 097 arxiv:1207.1717

Notice "a" scale changed significantly

Here we use the log dependence on a not 1 in EWPD determined in

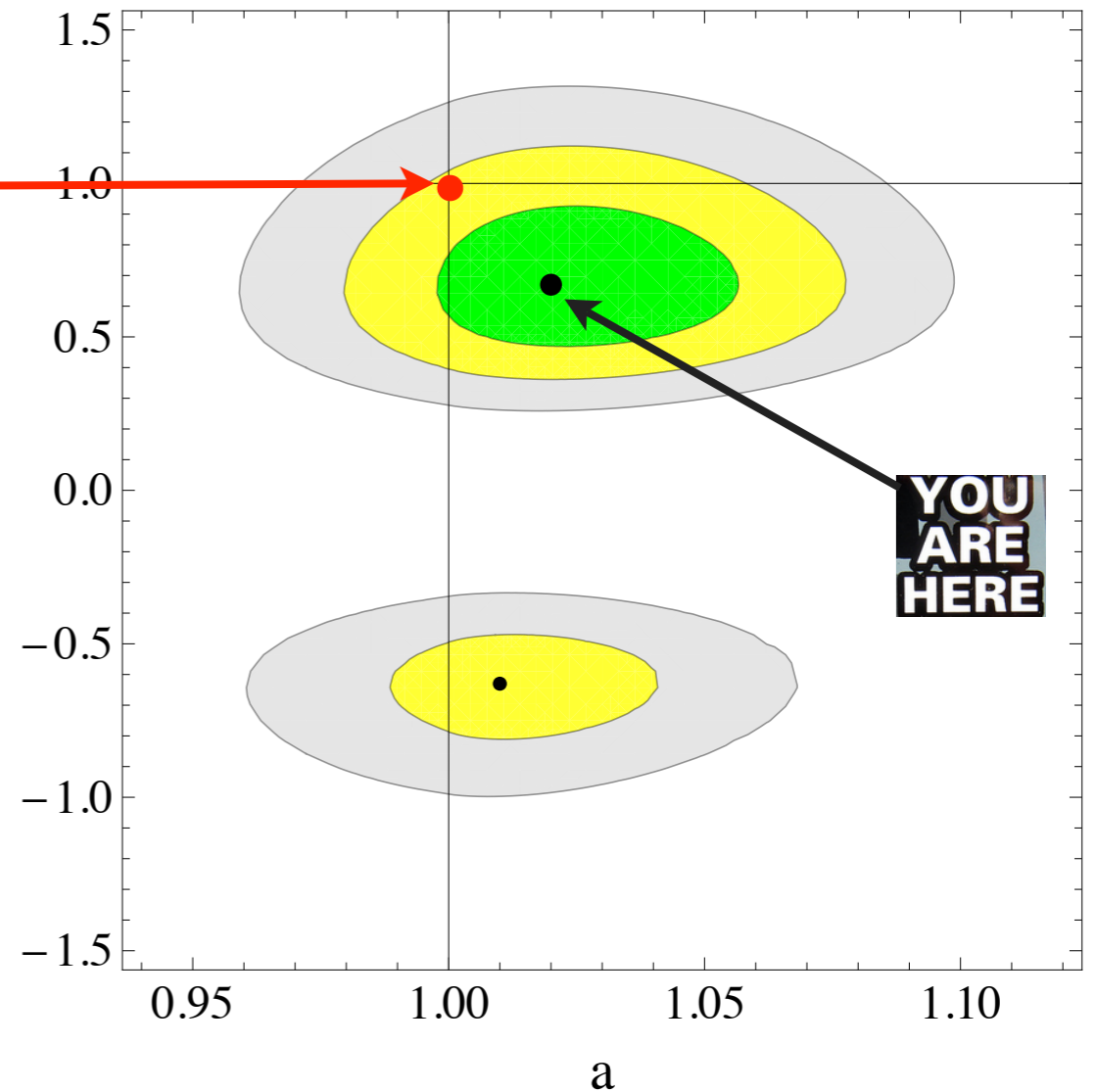
Barbieri, Bellazzini, Rychkov, Varagnolo arXiv:0706.0432

Add in EWPD to the Fit



GWS is here, will we converge?

7&8 TeV LHC data & Tevatron + EWPD



Notice “a” scale changed significantly

Here we use the log dependence on a not 1 in EWPD determined in

Barbieri, Bellazzini, Rychkov, Varagnolo arXiv:0706.0432

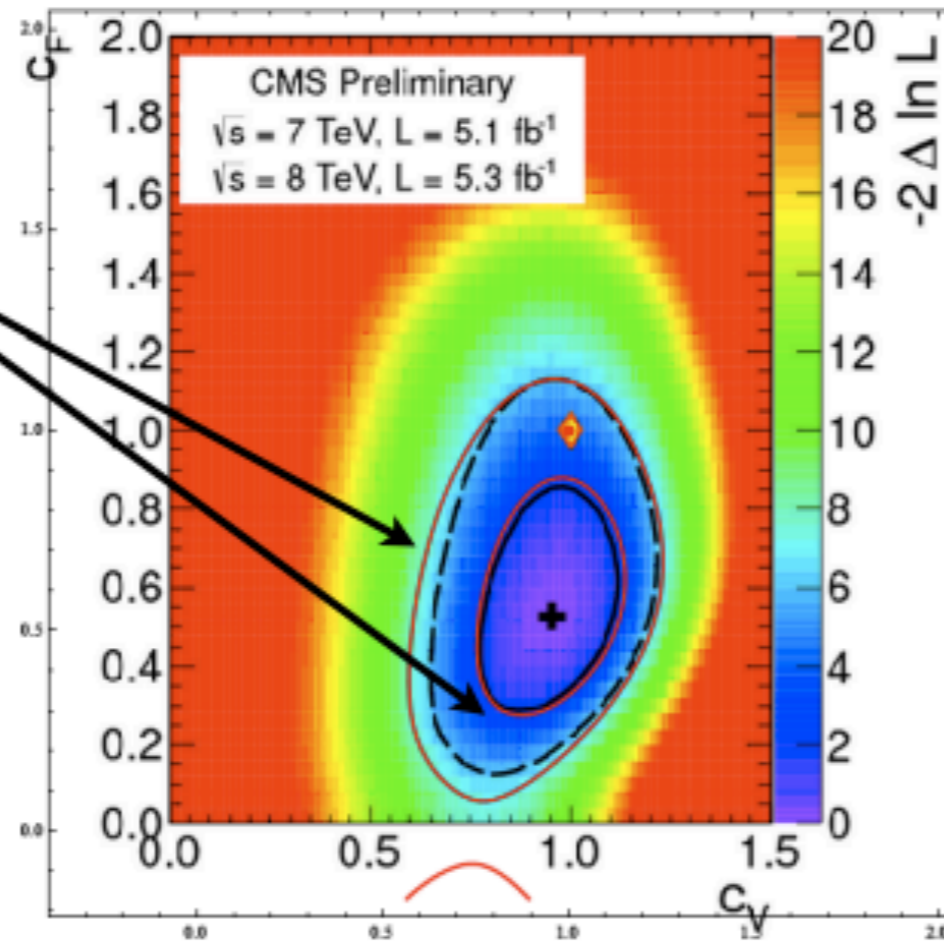
Can you trust a theorist to do this?

Comparison with CMS "official" fit

*Your χ^2 is too damn good!
-our friendly competition*

CMS imposed a prior $c > 0$
(it doesn't affect χ^2 , but it modifies $\Delta\chi^2$)

Our contours



This means that:

a) We are not badly screwing up.

b) correlations do not matter (summer)

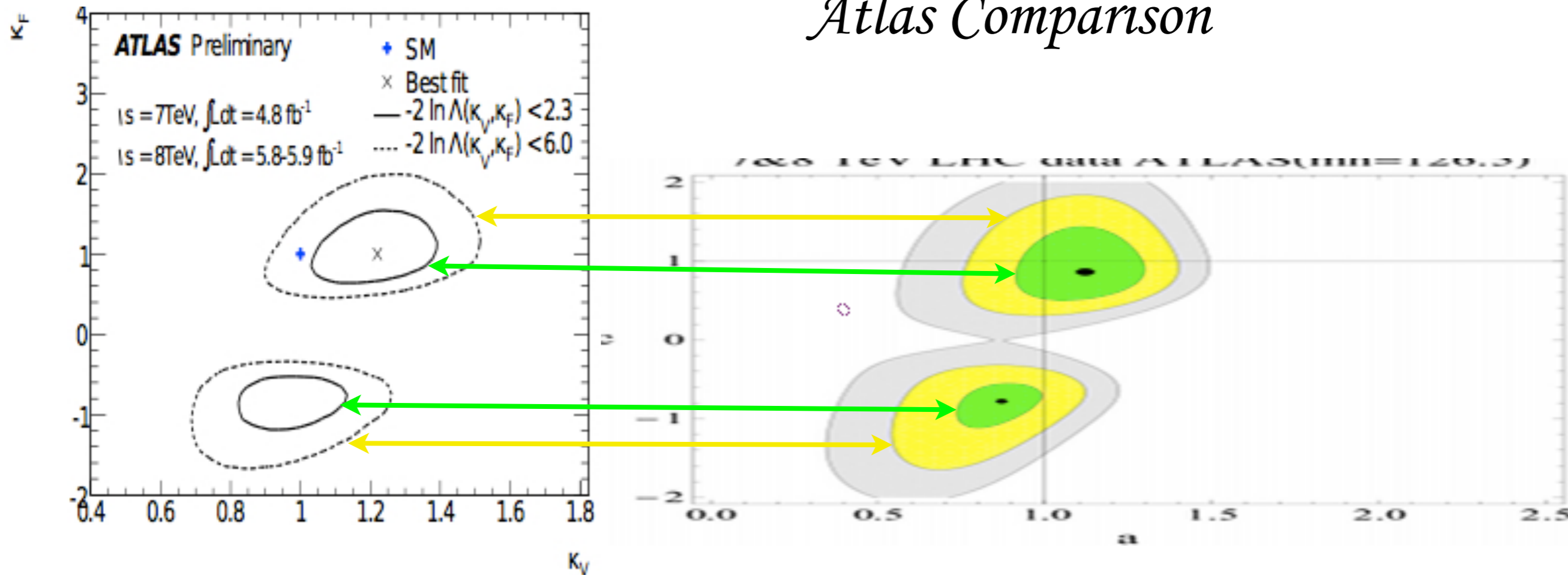
or

b) they do matter but CMS is as lost on estimating them correctly as we are.

Conclusion: You can trust some theorists to do this (for now).

More Movement by Atlas and CMS in this direction

Atlas Comparison



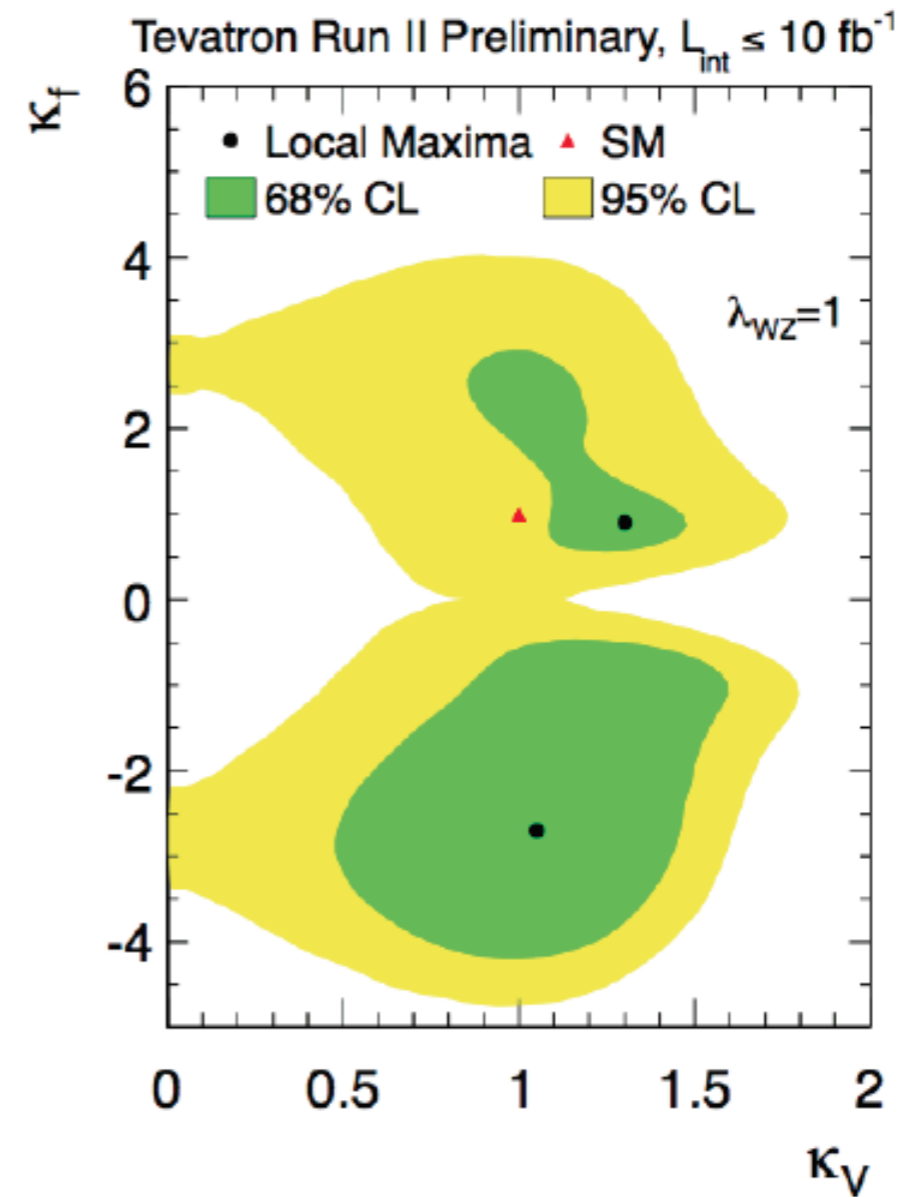
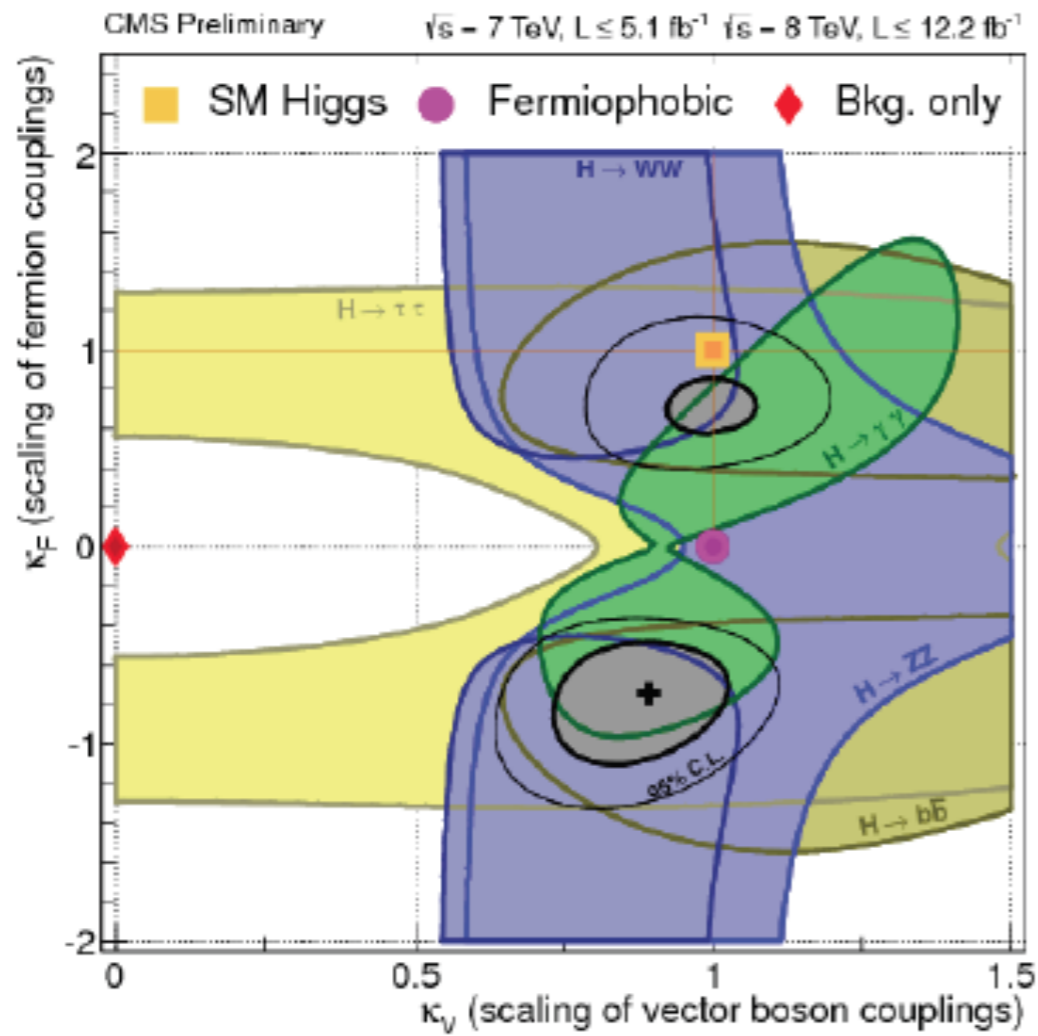
The contours of comparison are pretty good. However, more of a shift than CMS.

- 1) WW (0,1,2 jet) sub-channel treatment, they have more info, can use the sub-channels.
- 2) gamma gamma correlations might matter here due to the way the data was sliced up

Need more info from ATLAS for a more direct comparison.

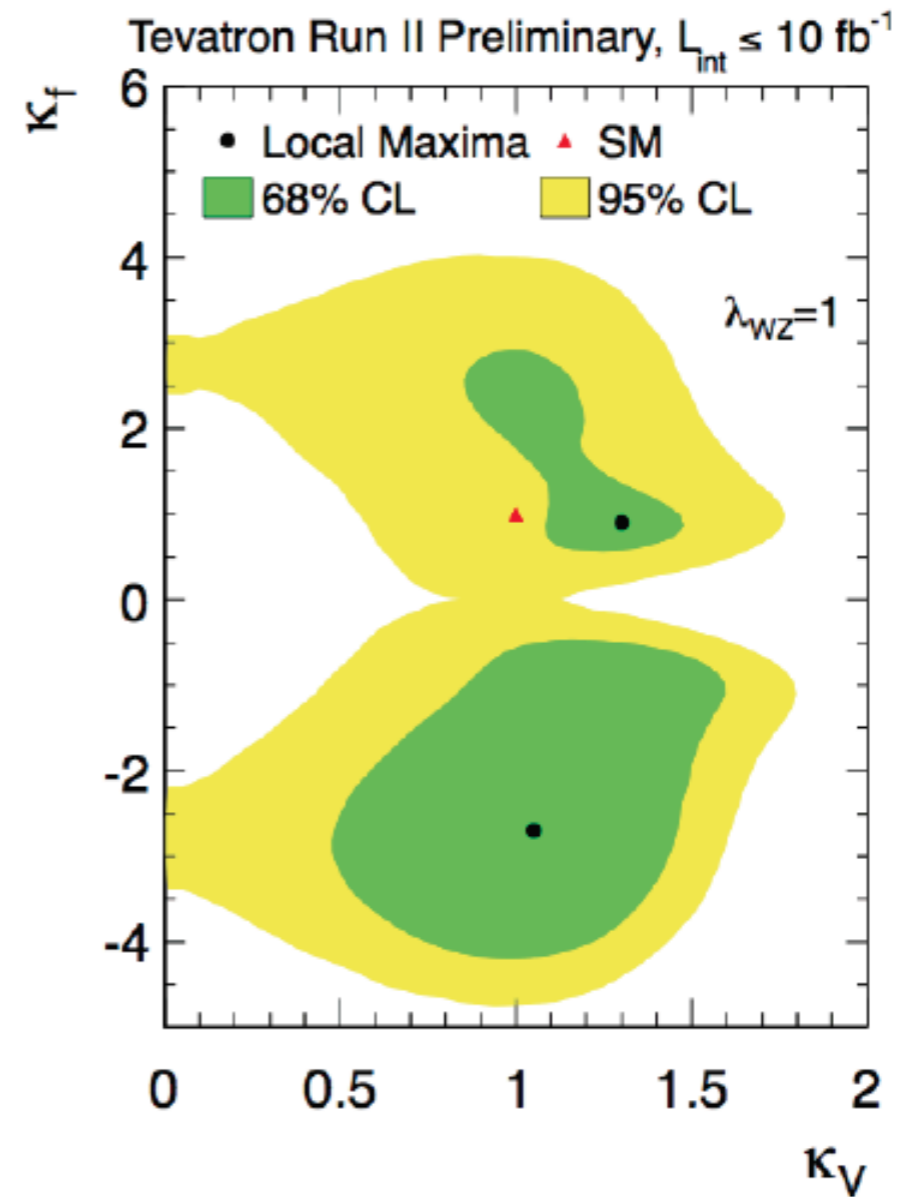
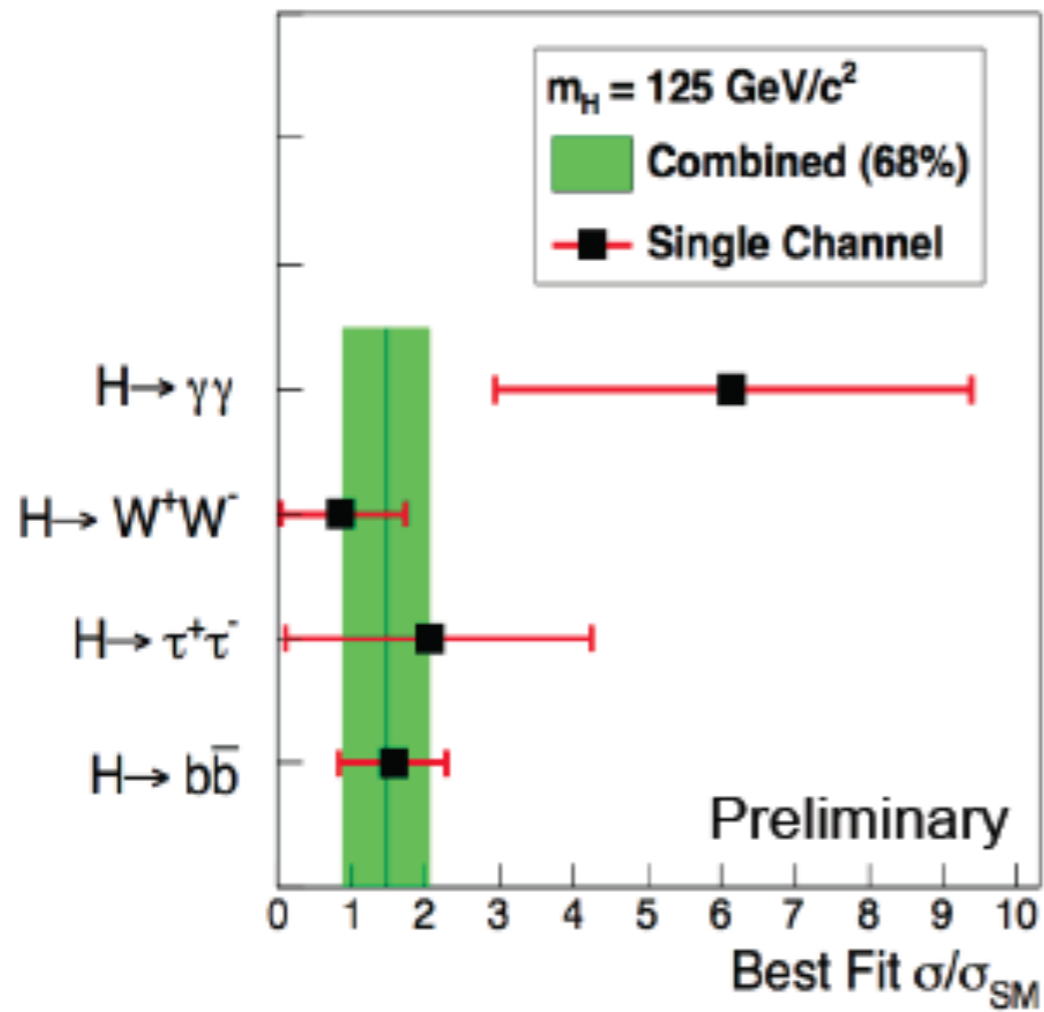
The hand off...

“Via con Dios” experimental fitters...



Wait a minute...

Public Tevatron data:

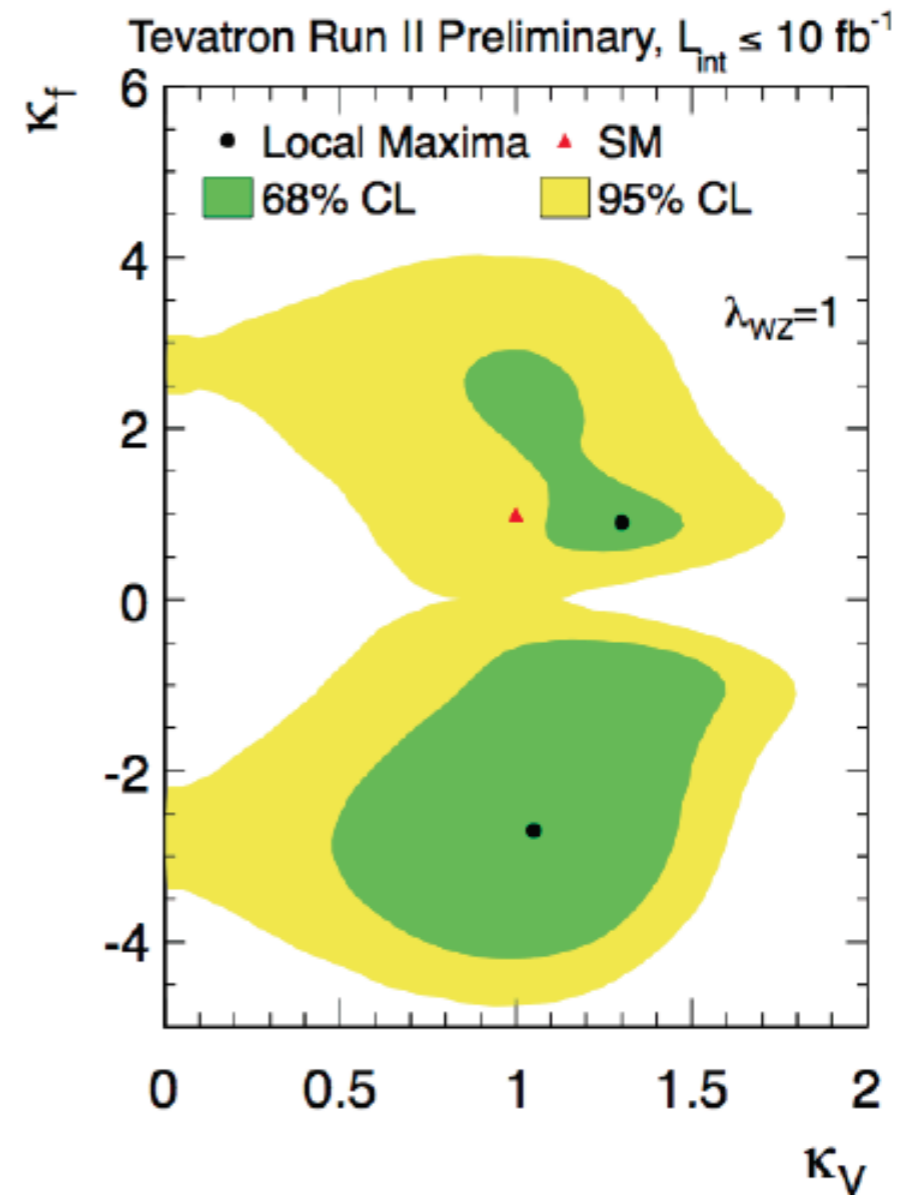
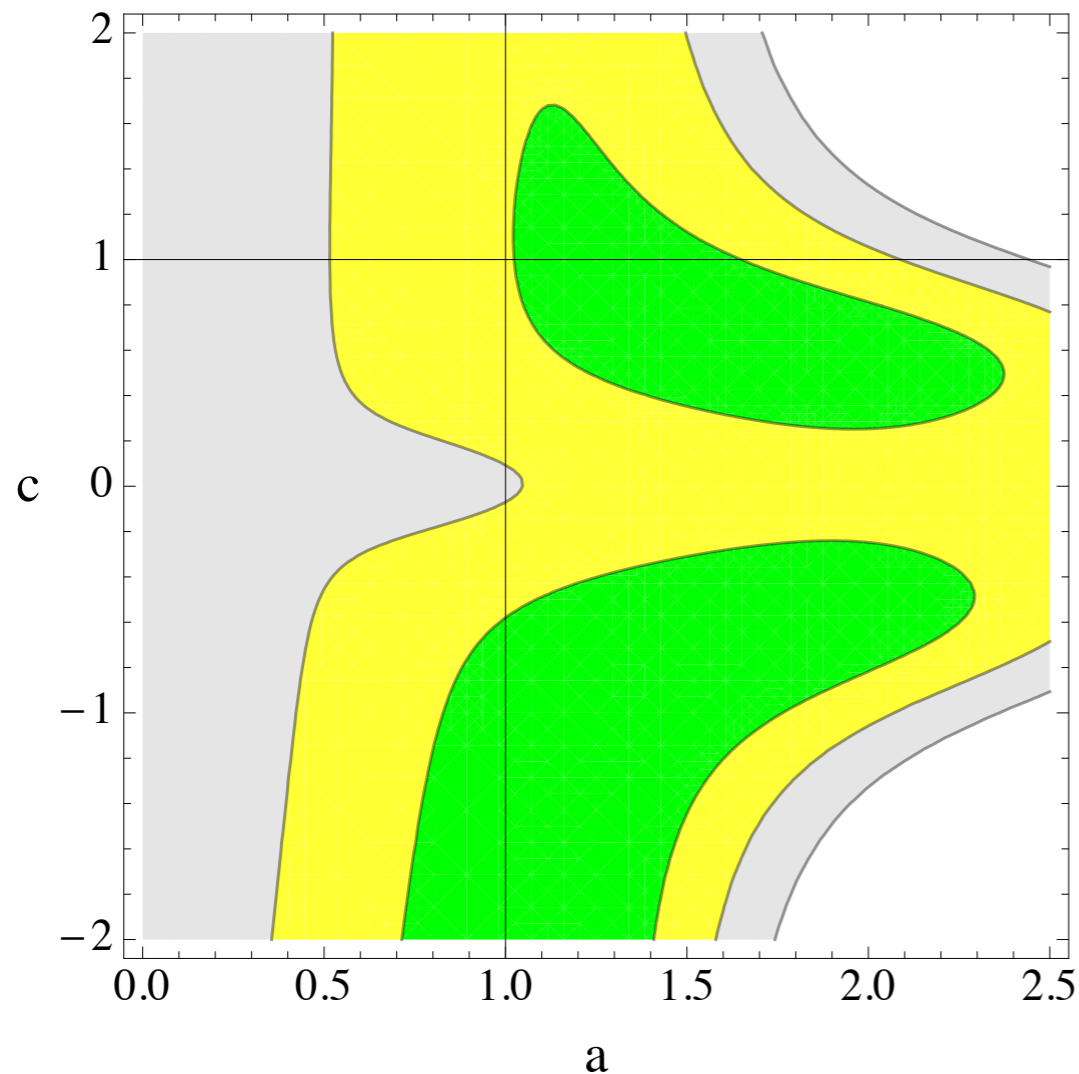


Wait a minute...

Public Tevatron data, gives:

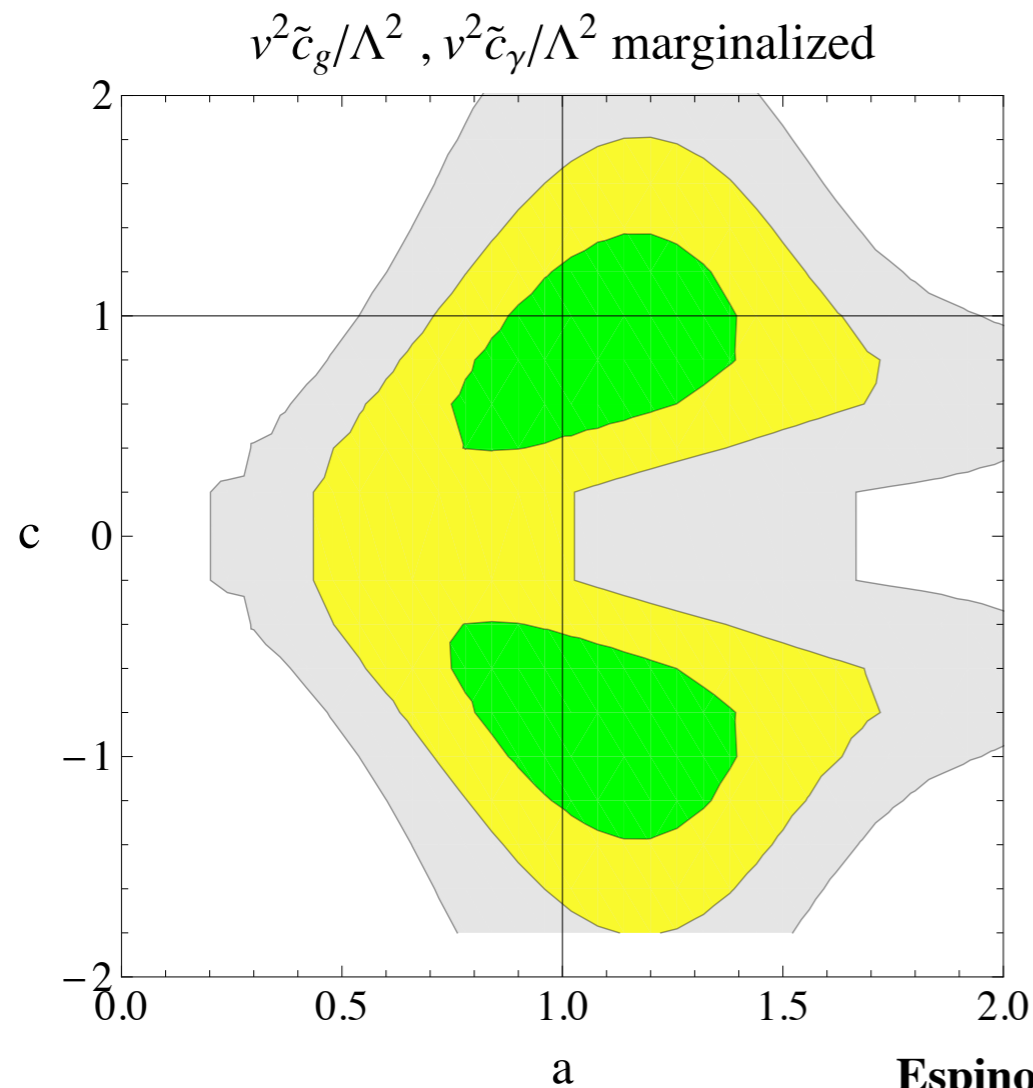


Tevatron Data, HCP update, $m_h = 125$ GeV



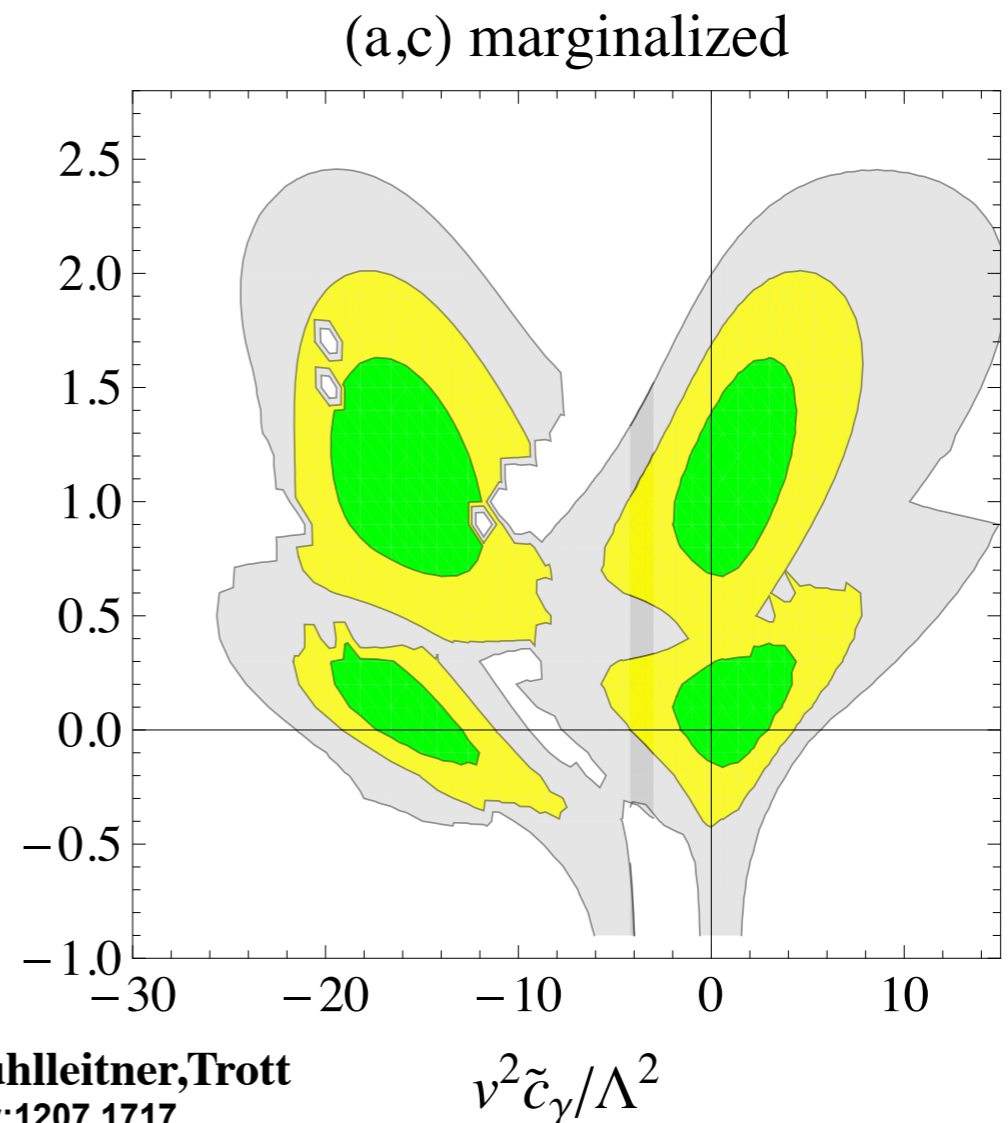
.....different scales....but still the public info not sufficient. Info will have to be further resolved so that broad physics conclusions can be drawn by theorists from Higgs data.

Marginalization Games



Espinosa, Grojean, Muhlleitner, Trott
JHEP 1205 (2012) 097 arxiv:1207.1717

$v^2 \tilde{c}_g / \Lambda^2$



$v^2 \tilde{c}_\gamma / \Lambda^2$

Very interesting that the SM higgs hypothesis test is improved in the context of NP in this way. Need more data.

A way of seeing that the existence of the $\gamma\gamma$ "excess" depends upon the Yukawa couplings being SM-like. Need more data.

Limit/Tension Methodology:

*What is this statistics # \$ @ ! * \$ in your paper, it is just equivalent to a damn χ^2 as far as i am concerned!* *-another charitable physicist.*

Each signal strength measurement can be approx:

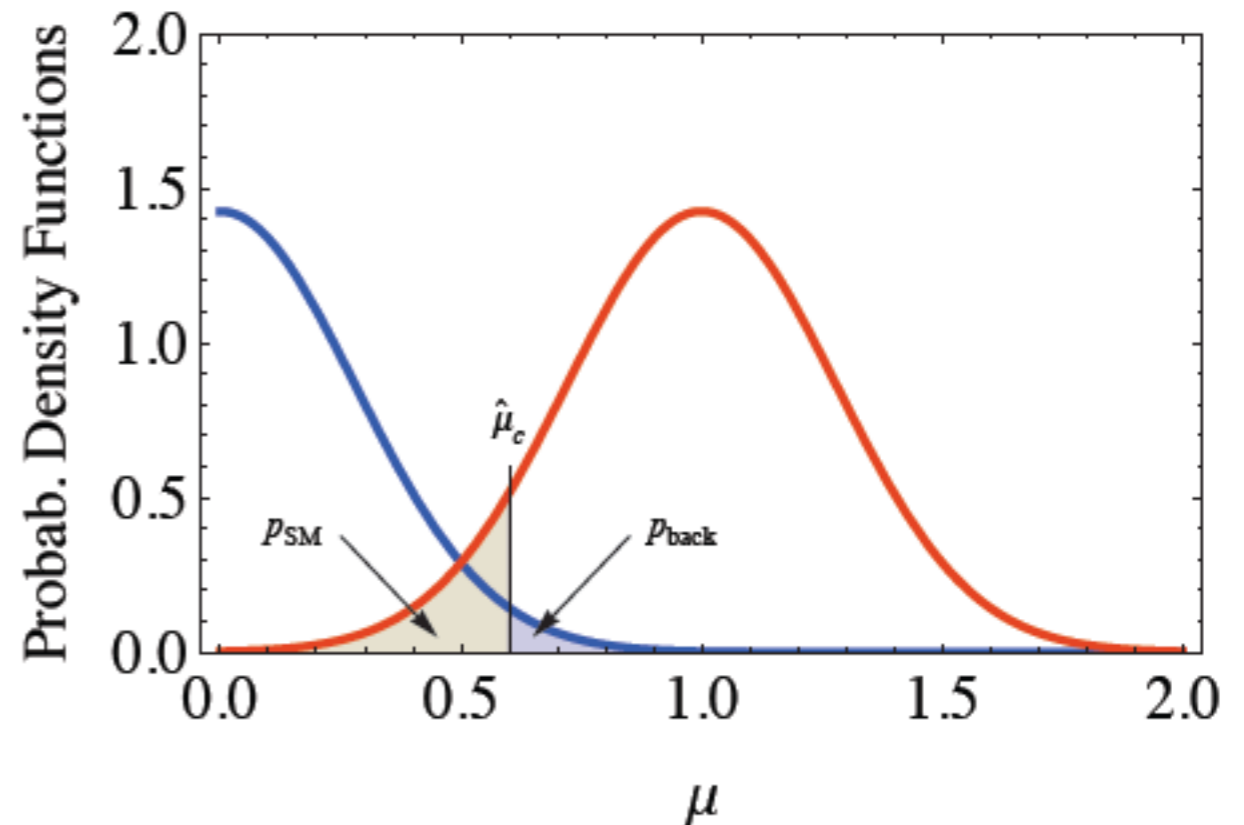
$$pdf_i(\mu, \hat{\mu}_i, \sigma_i) \approx e^{-(\mu - \hat{\mu}_i)^2 / (2\sigma_i^2)}$$

The PDF's can be combined to get global PDF's

$$pdf(\mu, \hat{\mu}_c, \sigma_c) \propto \prod_i^{N_{ch}} pdf_i(\mu, \hat{\mu}_i, \sigma_i) = \mathcal{N}_c e^{-(\mu - \hat{\mu}_c)^2 / (2\sigma_c^2)}$$

Where you have the combination variables:

$$\frac{1}{\sigma_c^2} = \sum_i^{N_{ch}} \frac{1}{\sigma_i^2}, \quad \hat{\mu}_c = \sum_i^{N_{ch}} \frac{\hat{\mu}_i}{\sigma_i^2}$$

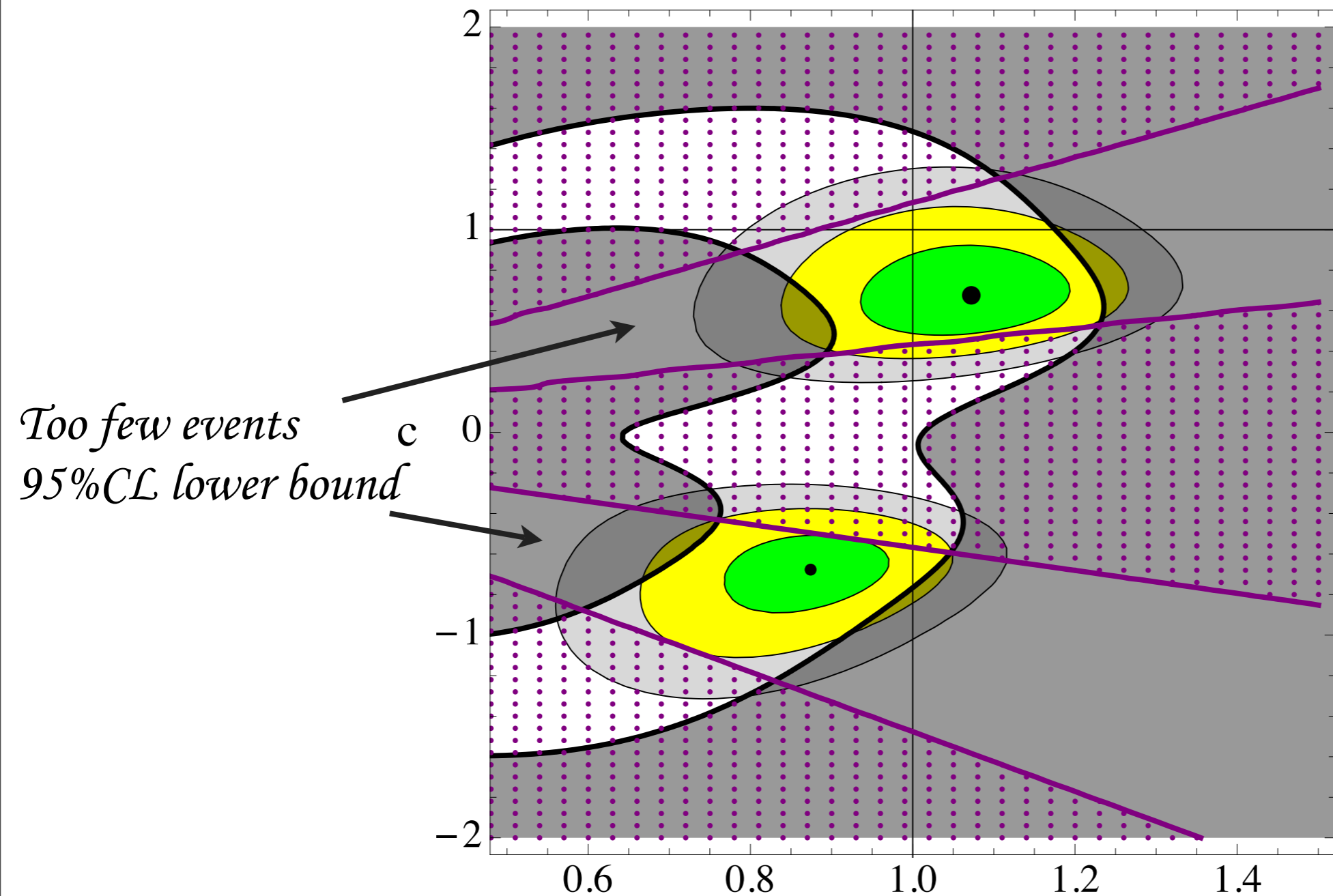


(note: correlations neglected here)

PDF's make clear one can set upper, lower and consistency limits on signal strength values.

Ultimate PDF usage to check the data set

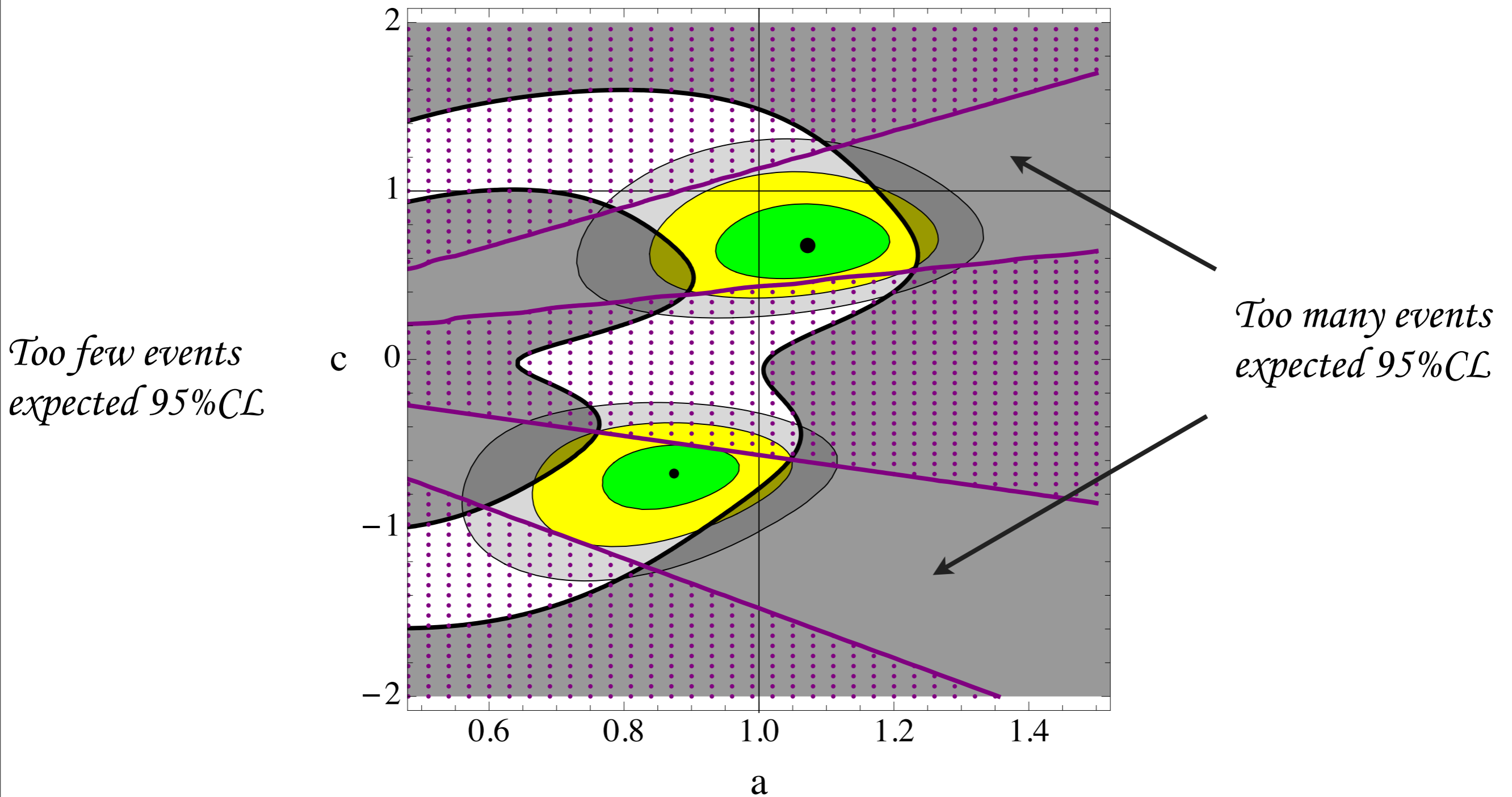
7&8 TeV LHC & Tevatron data



Espinosa, Grojean, Muhlleitner, Trott
JHEP 1205 (2012) 097 arxiv:1207.1717

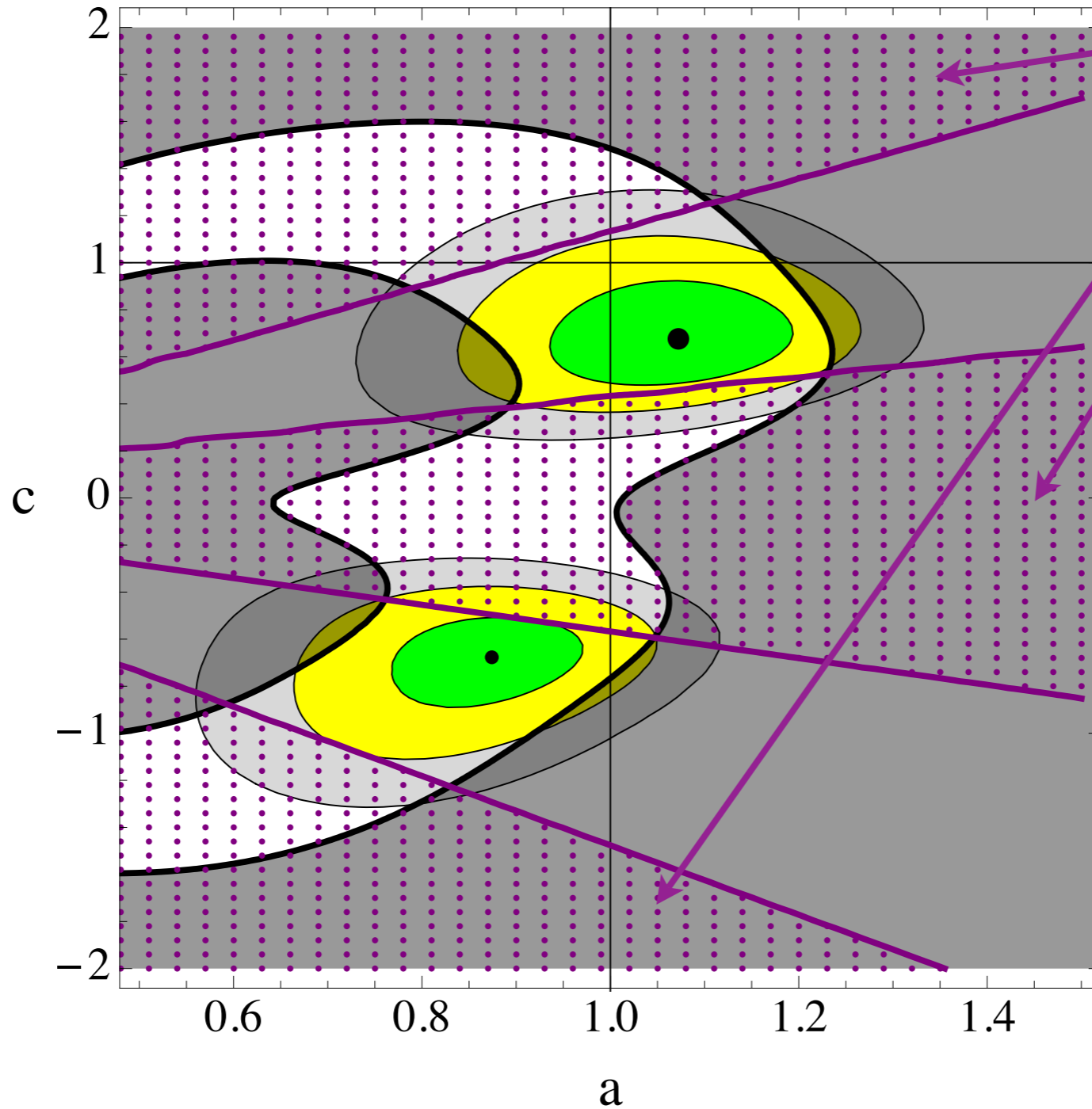
Ultimate PDF usage to check the data set

7&8 TeV LHC & Tevatron data



Ultimate PDF usage to check the data set

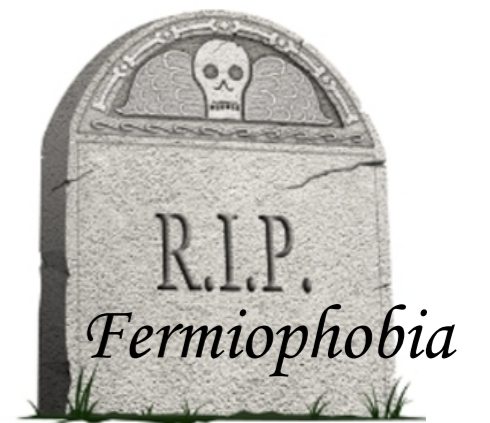
7&8 TeV LHC & Tevatron data



Inconsistencies driven by individual channels (2 sigma)

Too many events expected 95%CL

Too few events expected 95%CL



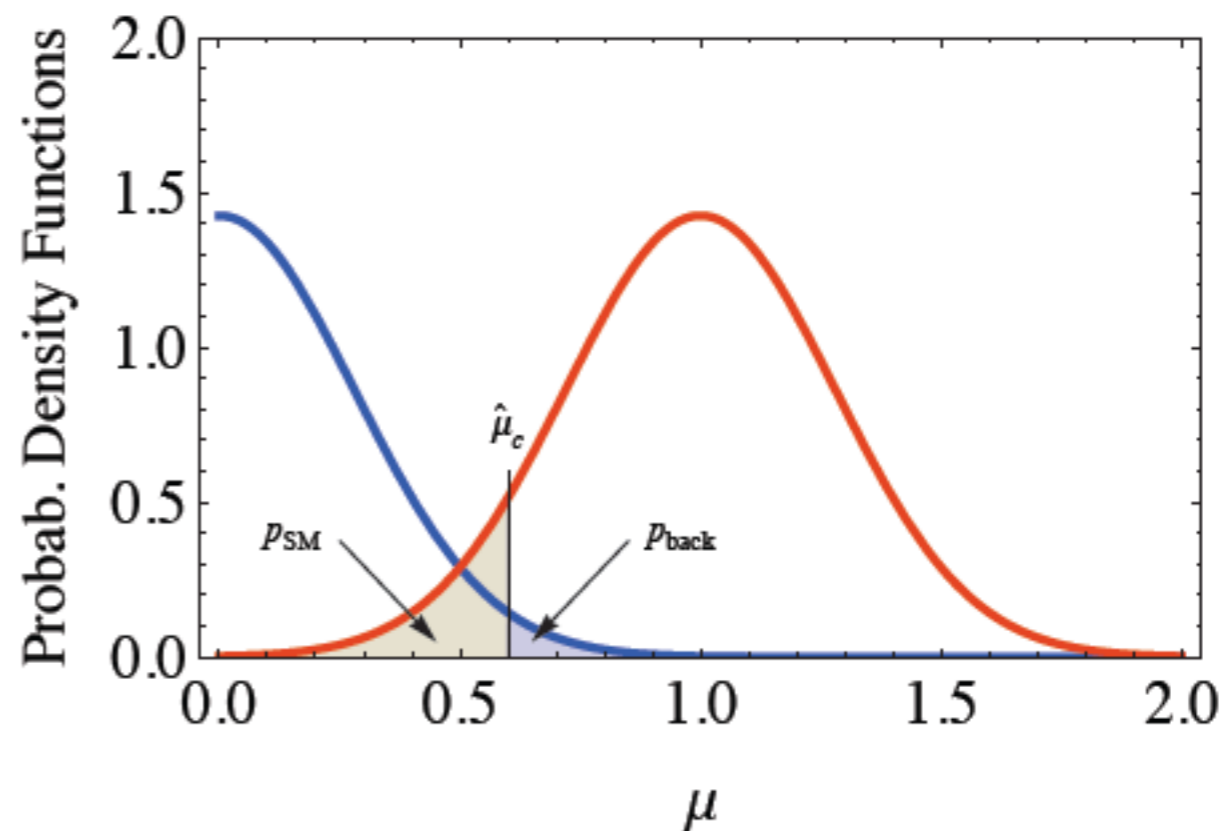
Branching Ratio Invisible

The invisible branching ratio is great as it is a universal shift on signal strengths.

$$\text{Br}(h \rightarrow f) \equiv \frac{\Gamma(h \rightarrow f)}{\Gamma_{\text{SM}} + \Gamma_{\text{inv}}} = (1 - \text{Br}_{\text{inv}}) \times \text{Br}_{\text{SM}}(h \rightarrow f).$$

In terms of the gaussian combination variables

$$\frac{1}{\sigma_c^2} = \sum_i^{N_{ch}} \frac{1}{\sigma_i^2}, \quad \hat{\mu}_c = \sum_i^{N_{ch}} \frac{\hat{\mu}_i}{\sigma_i^2}.$$

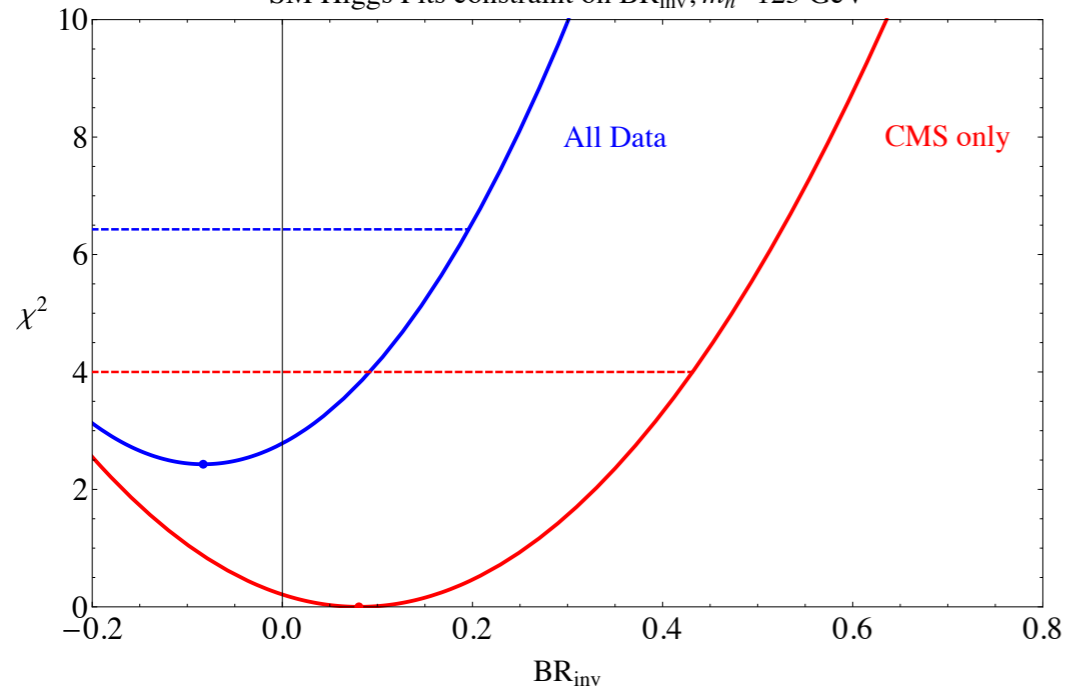


The invisible branching ratio is expressed as: $\text{Br}_{\text{inv}} = 1 - \hat{\mu}_c.$

One can fit to it using the SUPPLIED COMBINED SIGNAL STRENGTHS

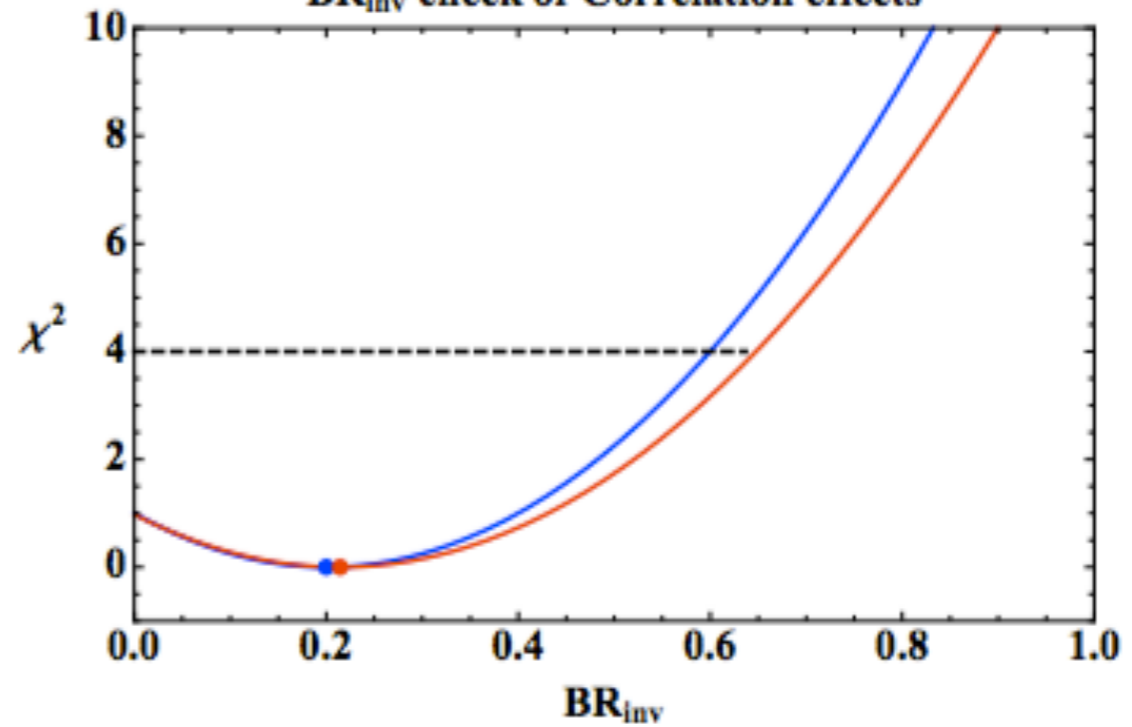
Branching Ratio Invisible - latest data

SM Higgs Fits constraint on BR_{inv} , $m_h=125$ GeV



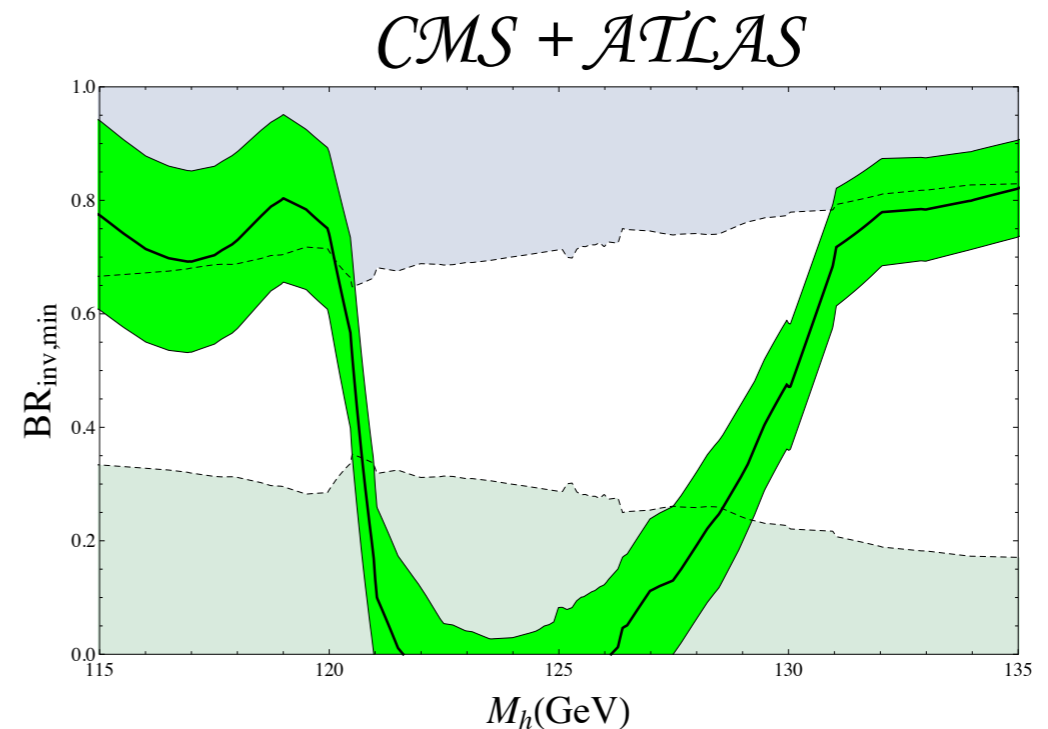
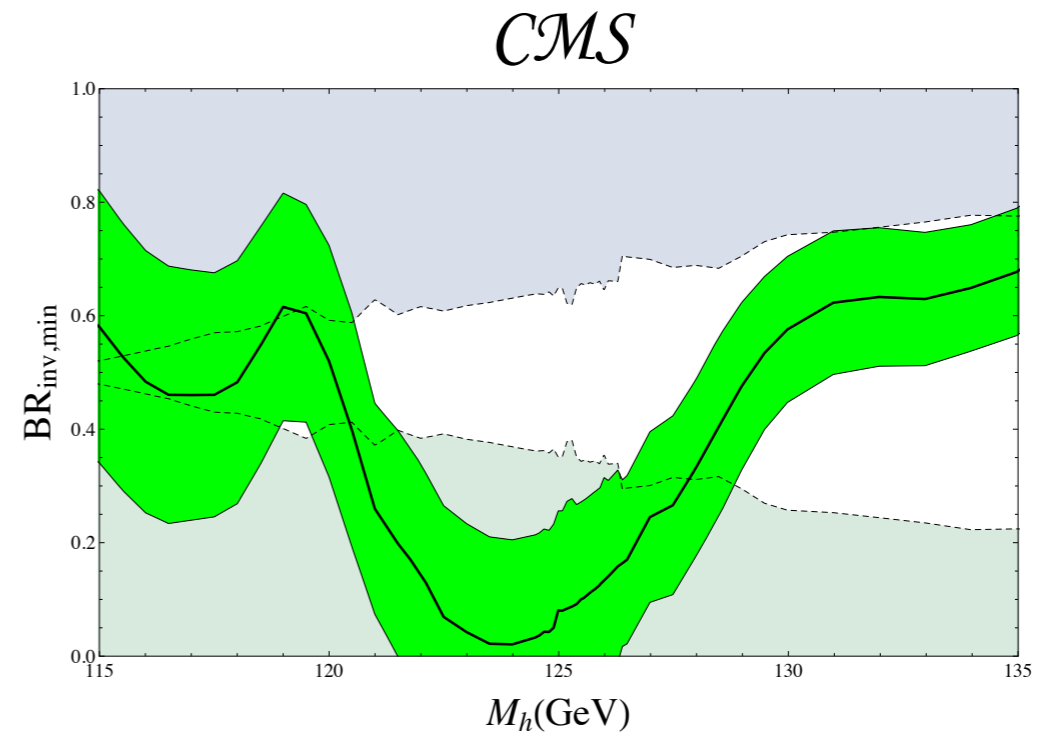
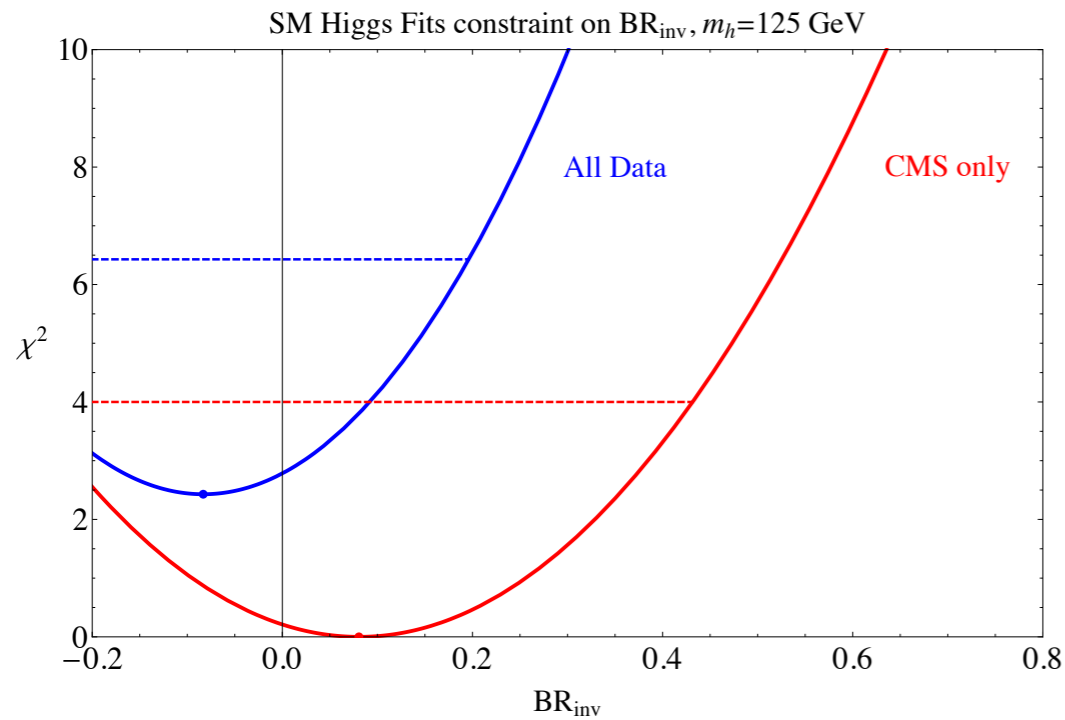
*Any global evidence of Br_{inv} ?
Nope.*

BR_{inv} check of Correlation effects



*Any global evidence that correlations
are so important that we are full of it?
Nope.*

Branching Ratio Invisible - latest data



Current data on supplied combined signal strengths, so all correlations included.

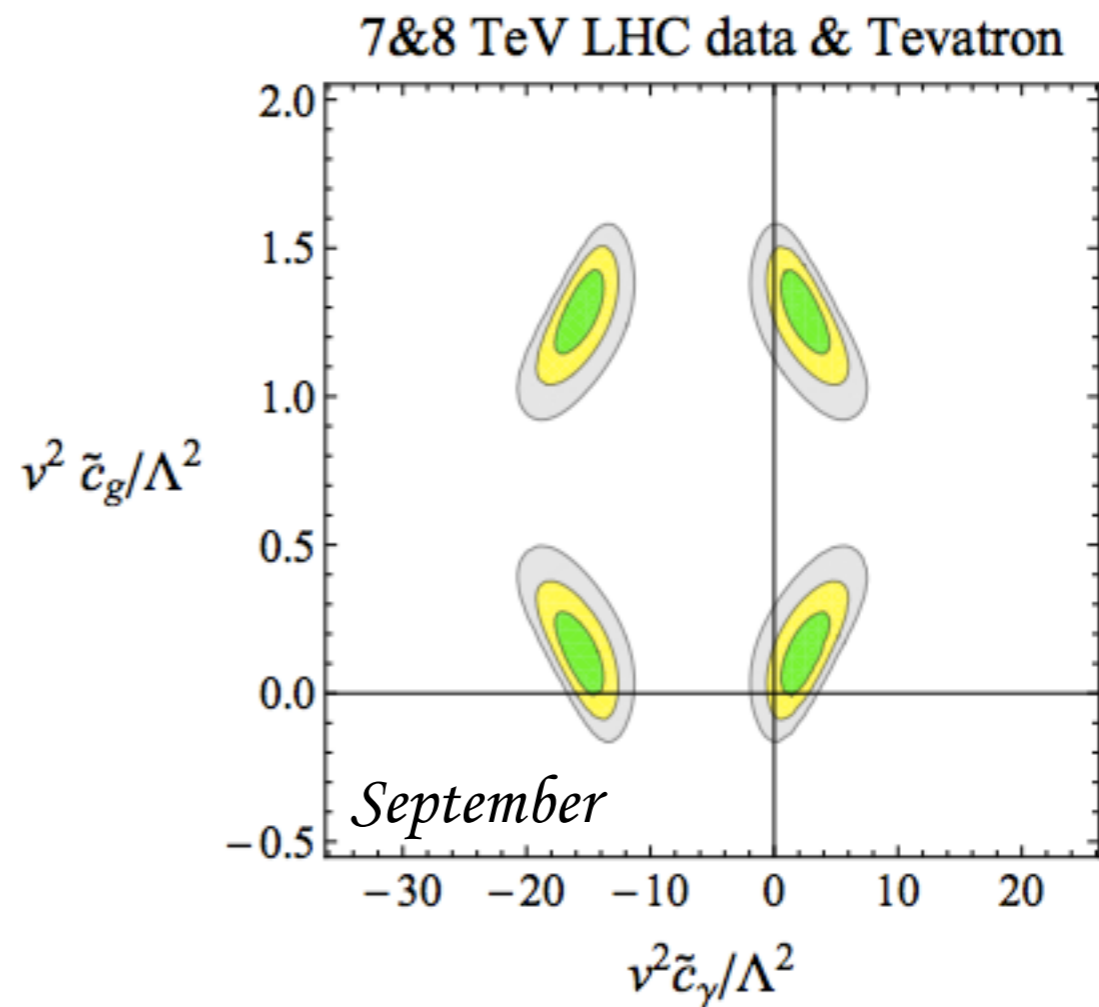
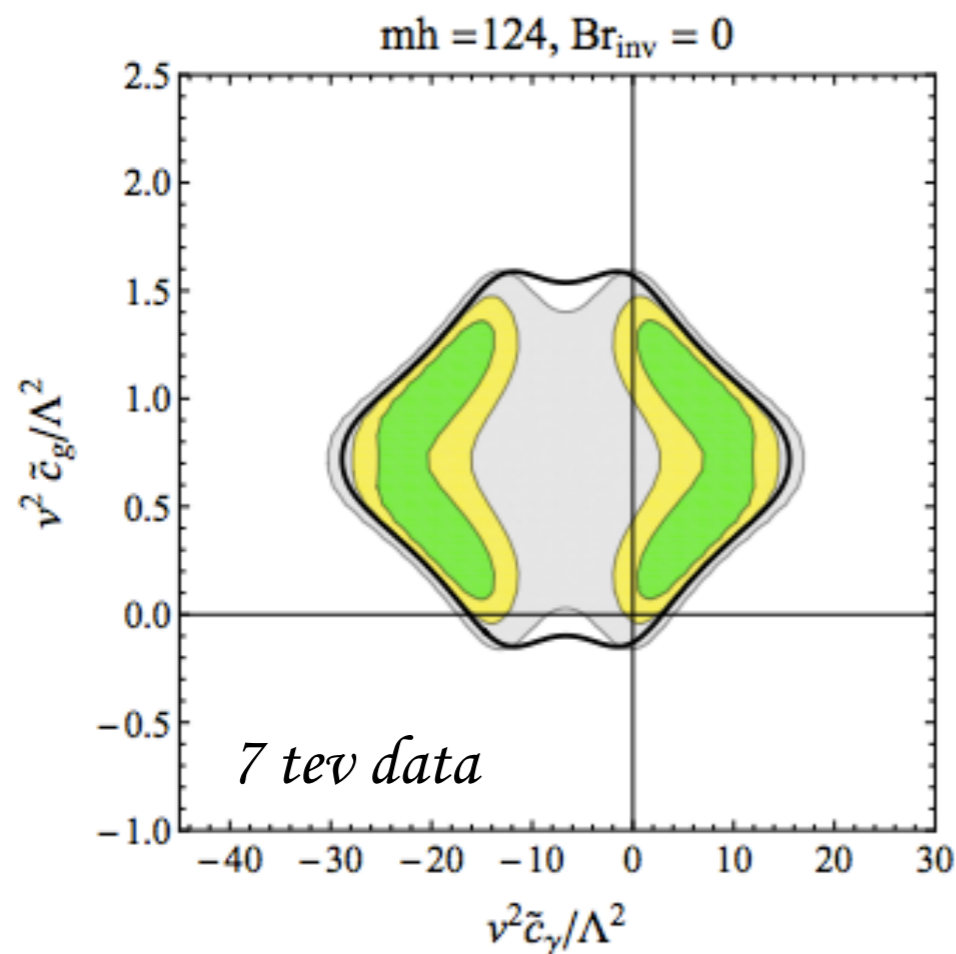
CMS is still low on the combined signal strength.

No Smoking gun for dark matter as yet.

Constraints are scaling.

If this is the right effective Lagrangian, we are learning a lot!

$$\mathcal{L}_{HD} = -\frac{c_g g_3^2}{2\Lambda^2} H^\dagger H G_{\mu\nu}^A G^{A\mu\nu} - \frac{c_W g_2^2}{2\Lambda^2} H^\dagger H W_{\mu\nu}^a W^{a\mu\nu} - \frac{c_B g_1^2}{2\Lambda^2} H^\dagger H B_{\mu\nu} B^{\mu\nu} - \frac{c_{WB} g_1 g_2}{2\Lambda^2} H^\dagger \tau^a H B_{\mu\nu} W^{a\mu\nu},$$



Natural Susy

Consider minimal spectrum of stops, left handed sbottom, charginos and neutralinos with large gaugino mass.

Stop contributions to Higgs production:

$$\frac{\sigma(gg \rightarrow h)}{\sigma^{SM}(gg \rightarrow h)} \simeq \frac{\Gamma(h \rightarrow gg)}{\Gamma^{SM}(h \rightarrow gg)} \simeq |1 + r_g|^2 \quad r_g = \frac{C_g(\alpha_s) F_g(m_{\tilde{t}_1}, m_{\tilde{t}_2}, \theta_{\tilde{t}})}{F_g^{SM}(m_t, m_b \dots)}$$

Stop contributions to interesting decays:

$$\frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma^{SM}(h \rightarrow \gamma\gamma)} \simeq |1 + r_\gamma|^2, \quad r_\gamma = \frac{N_c Q_{\tilde{t}}^2 C_\gamma(\alpha_s) F_g(m_{\tilde{t}_1}, m_{\tilde{t}_2}, \theta_{\tilde{t}})}{F_\gamma^{SM}(m_t, W, m_b \dots)}$$

QCD squark
matching
correction

Same loop functions, non abelian nature of QCD irrelevant at leading order in matching.

Natural Susy

Consider minimal spectrum of stops, left handed sbottom, charginos and neutralinos with large gaugino mass.

Stop contributions to Higgs production, In terms of operators:

$$\sigma_{gg \rightarrow h} \approx \sigma_{gg \rightarrow h}^{SM} \left| 1 + \frac{2}{F_g^{SM}} \frac{v^2 \tilde{c}_g}{\Lambda^2} \right|^2, \quad \Gamma_{h \rightarrow \gamma\gamma} \approx \Gamma_{h \rightarrow \gamma\gamma}^{SM} \left| 1 + \frac{1}{F_\gamma^{SM}} \frac{v^2 \tilde{c}_\gamma}{\Lambda^2} \right|^2.$$

Matching with no running: $\frac{v^2 \tilde{c}_g}{\Lambda^2} \simeq C_g(\alpha_s) \frac{F_g}{2}, \quad \frac{v^2 \tilde{c}_\gamma}{\Lambda^2} \simeq N_c Q_{\tilde{t}}^2 C_\gamma(\alpha_s) F_g$

$$C_g(\alpha_s) = 1 + \frac{25 \alpha_s}{6 \pi}, \quad C_\gamma(\alpha_s) = 1 + \frac{8 \alpha_s}{3 \pi}.$$

This is a predictive scenario for the wilson coefficients of the higher d ops:

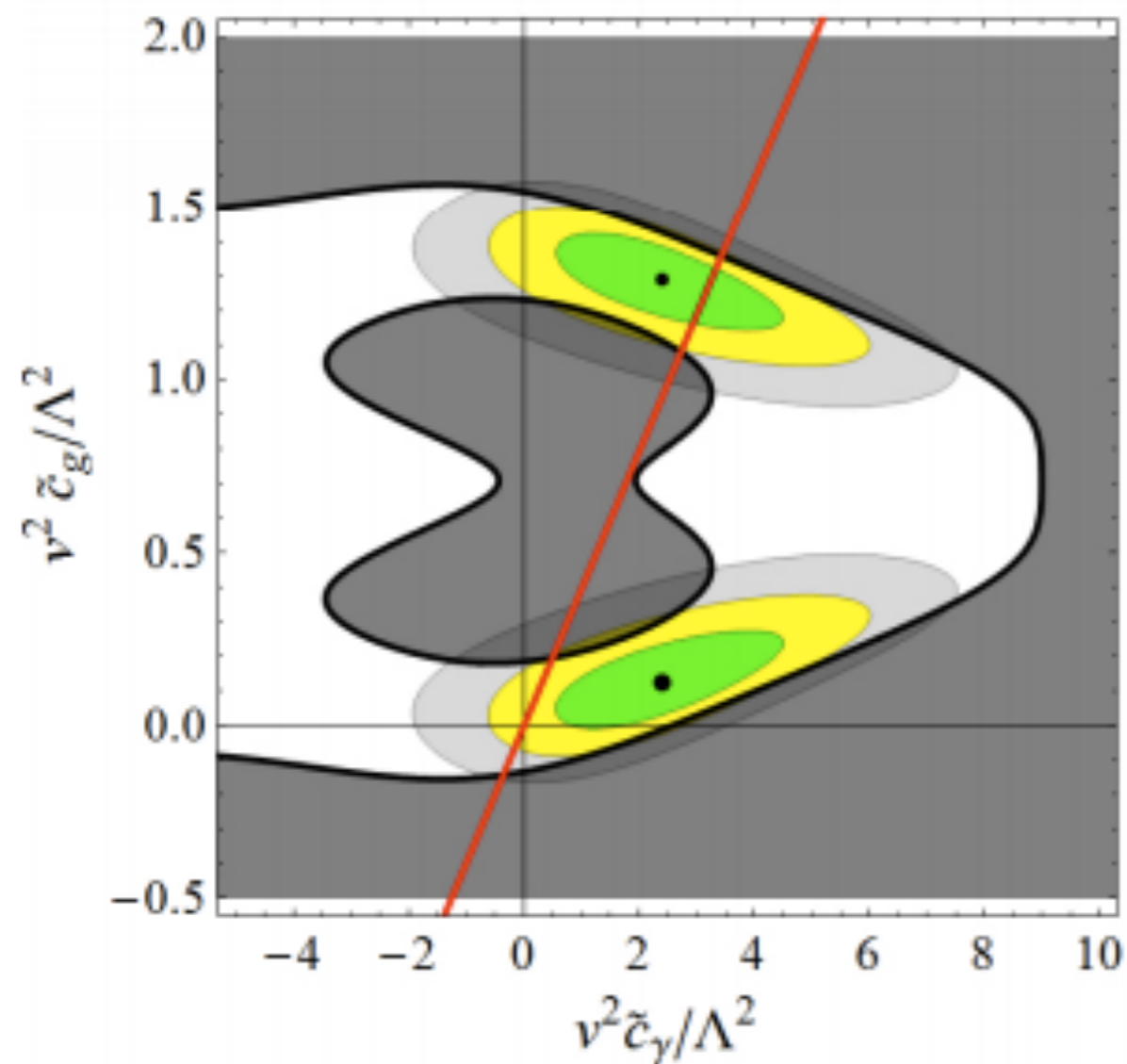
$$\frac{\tilde{c}_g}{\tilde{c}_\gamma} = \frac{1}{2N_c Q_{\tilde{t}}^2} \frac{C_g(\alpha_s)}{C_\gamma(\alpha_s)} = \frac{3}{8} \left(1 + \frac{3 \alpha_s}{2 \pi} \right)$$

Natural Susy

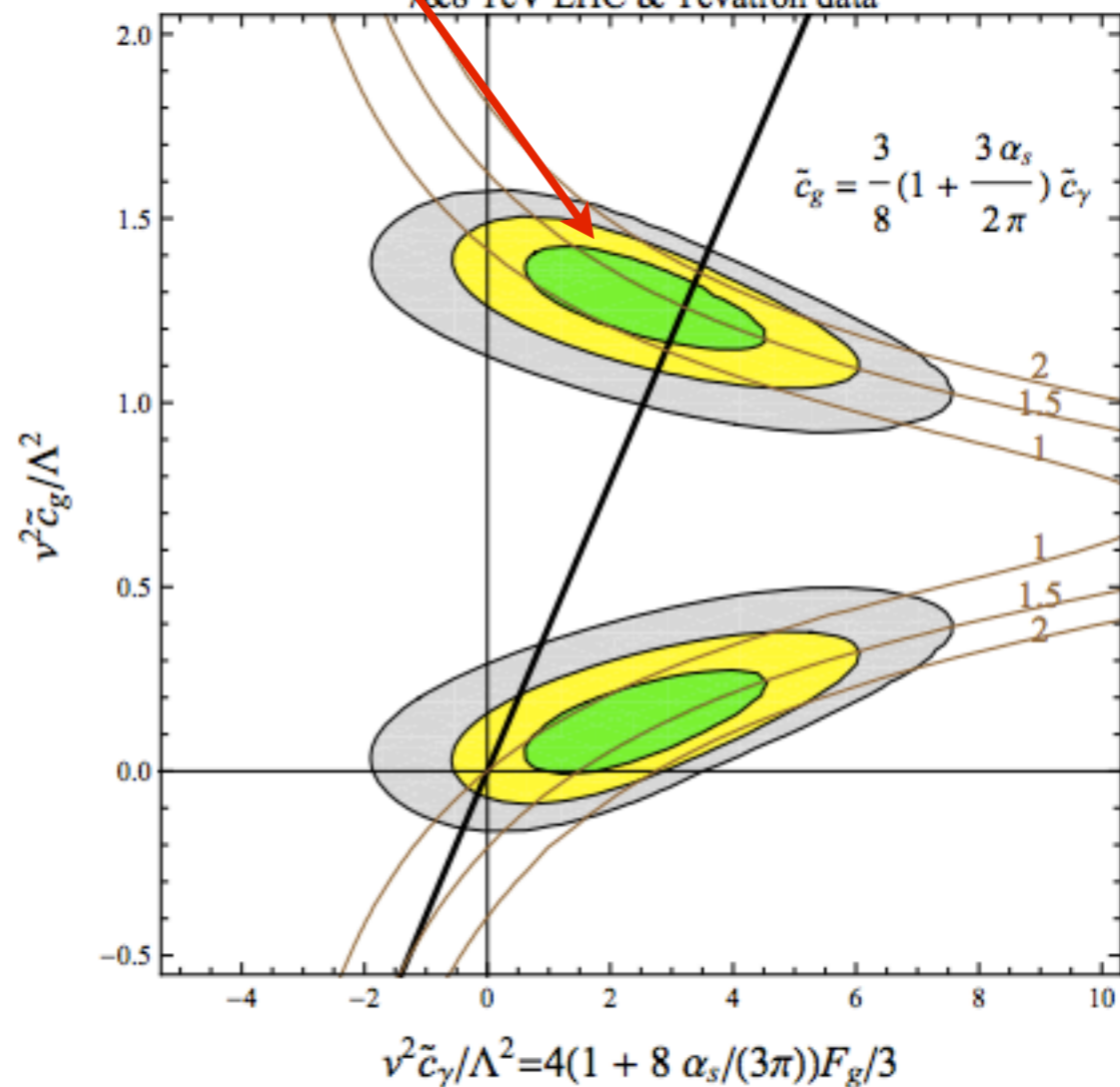
This is a predictive scenario for the wilson coefficients of the higher d ops:

*monophoton missing E_t exclusions
limit parameter space, $t\bar{t}$ spin correlations, etc
sanz et al. arXiv:1205.1463*

7&8 TeV LHC data & Tevatron

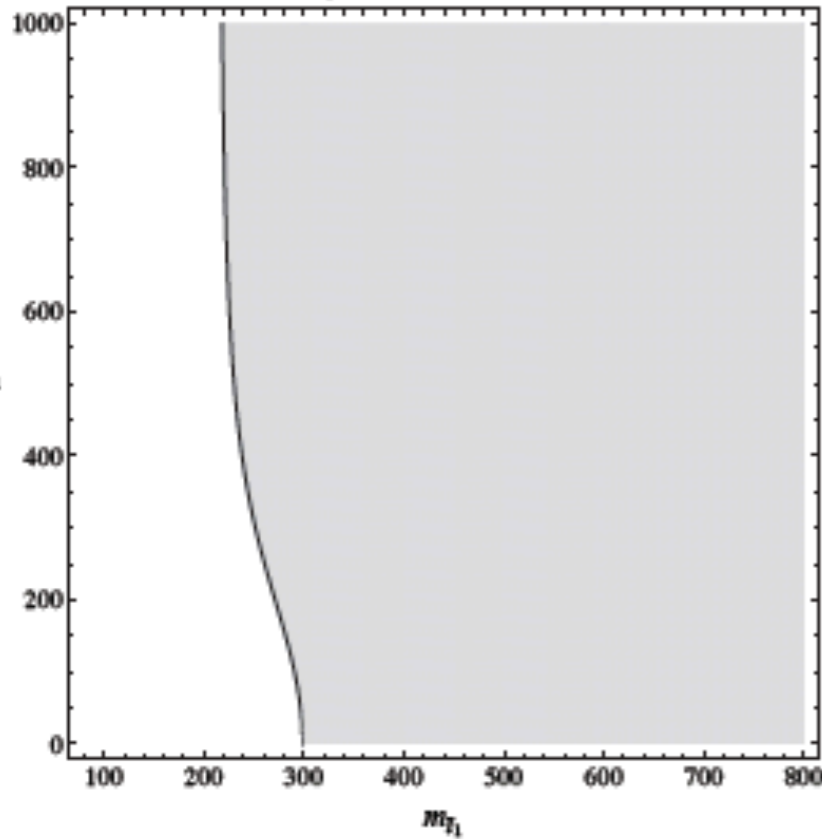


7&8 TeV LHC & Tevatron data

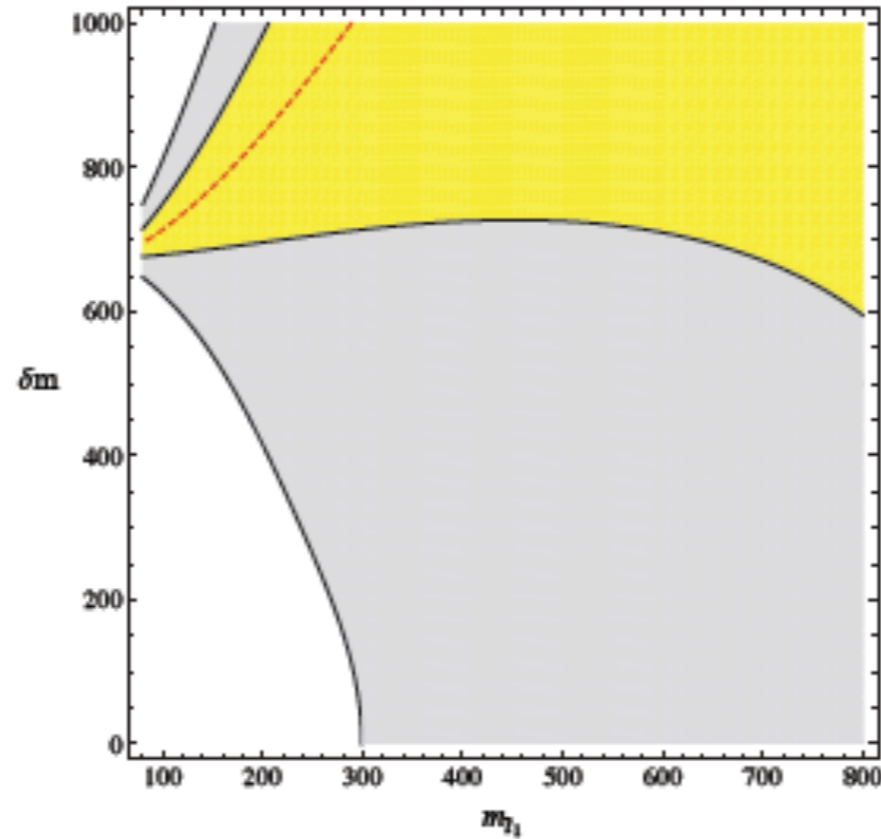


Translate to stop space

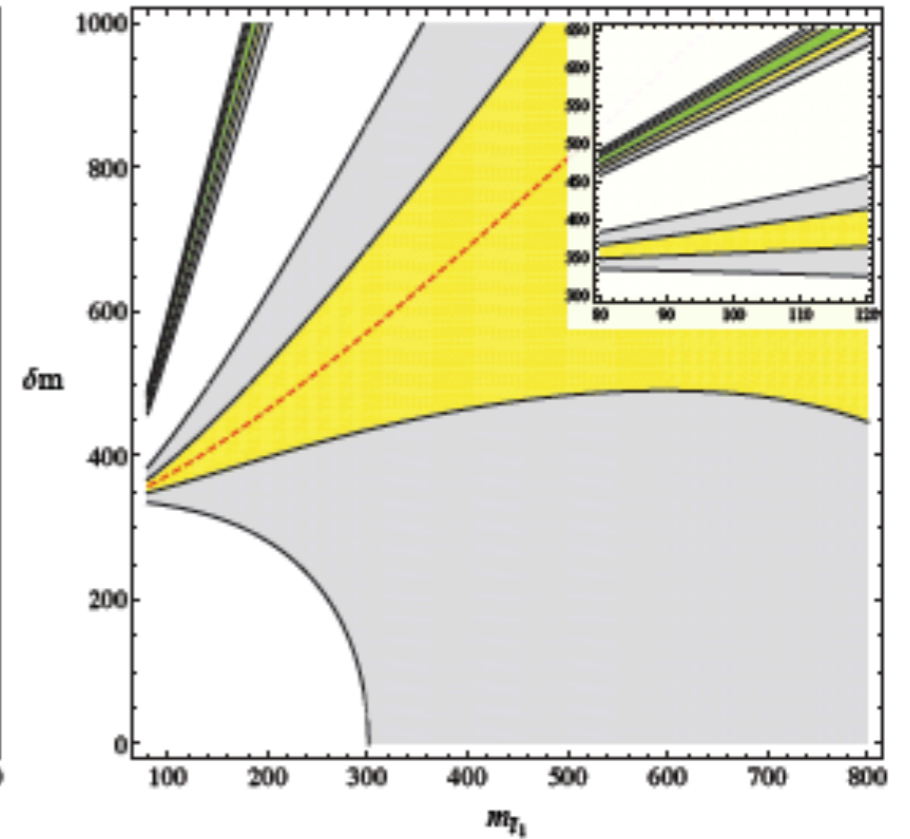
$\theta_t=0$ and $\tan\beta=2-20$



$\theta_t=\frac{\pi}{12}$ and $\tan\beta=2-20$



$\theta_t=\frac{\pi}{4}$ and $\tan\beta=2-20$



- 1σ
- 2σ
- 3σ

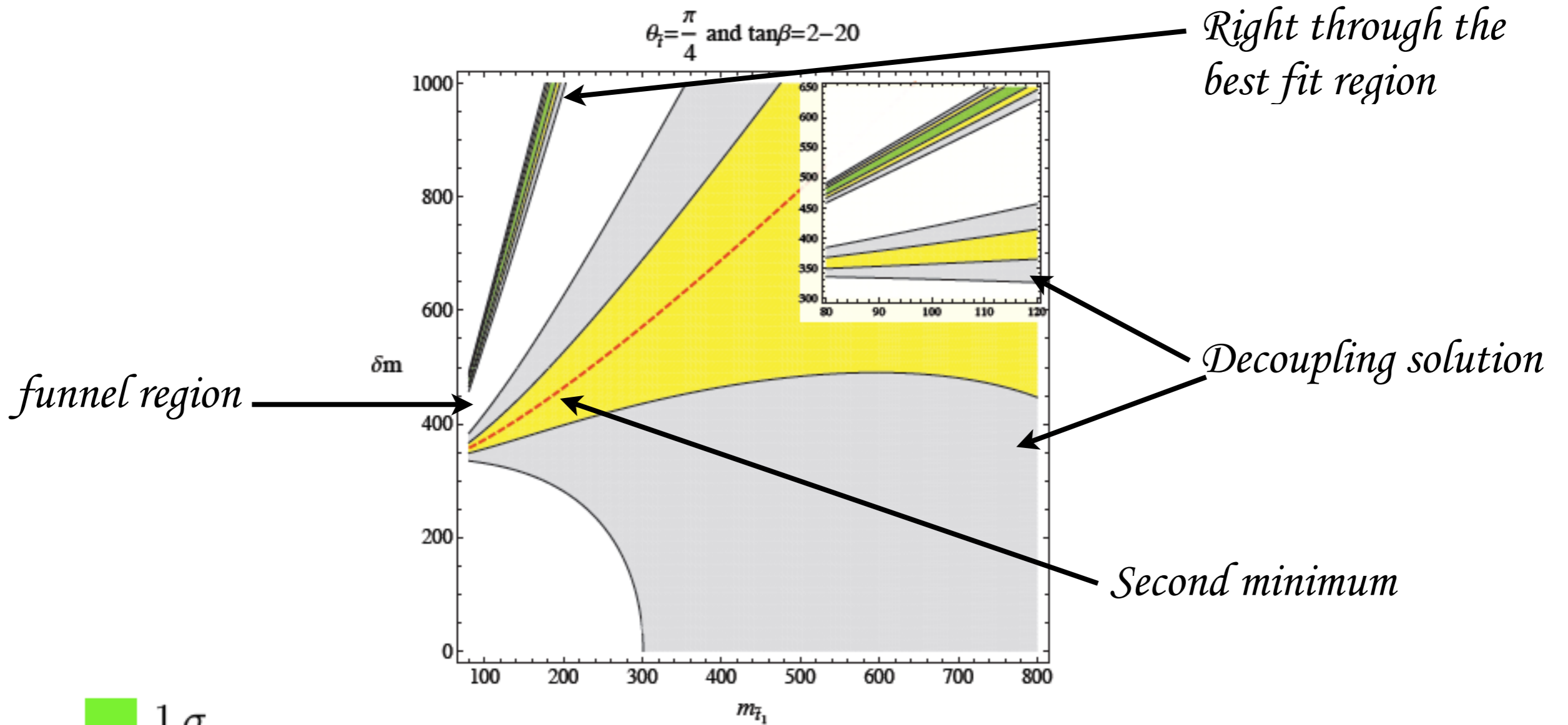
Map to stop space:

$$F_g = -\frac{1}{3} \left[\frac{m_t^2}{m_{\tilde{t}_1}^2} + \frac{m_t^2}{m_{\tilde{t}_2}^2} - \frac{1}{4} \sin^2(2\theta_t) \frac{\delta m^4}{m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} \right]$$

Light unmixed stops in bad shape, large mixing preferred in the data.

Translate to stop space

$$\theta_{\tilde{t}} = \frac{\pi}{4} \text{ and } \tan\beta = 2-20$$



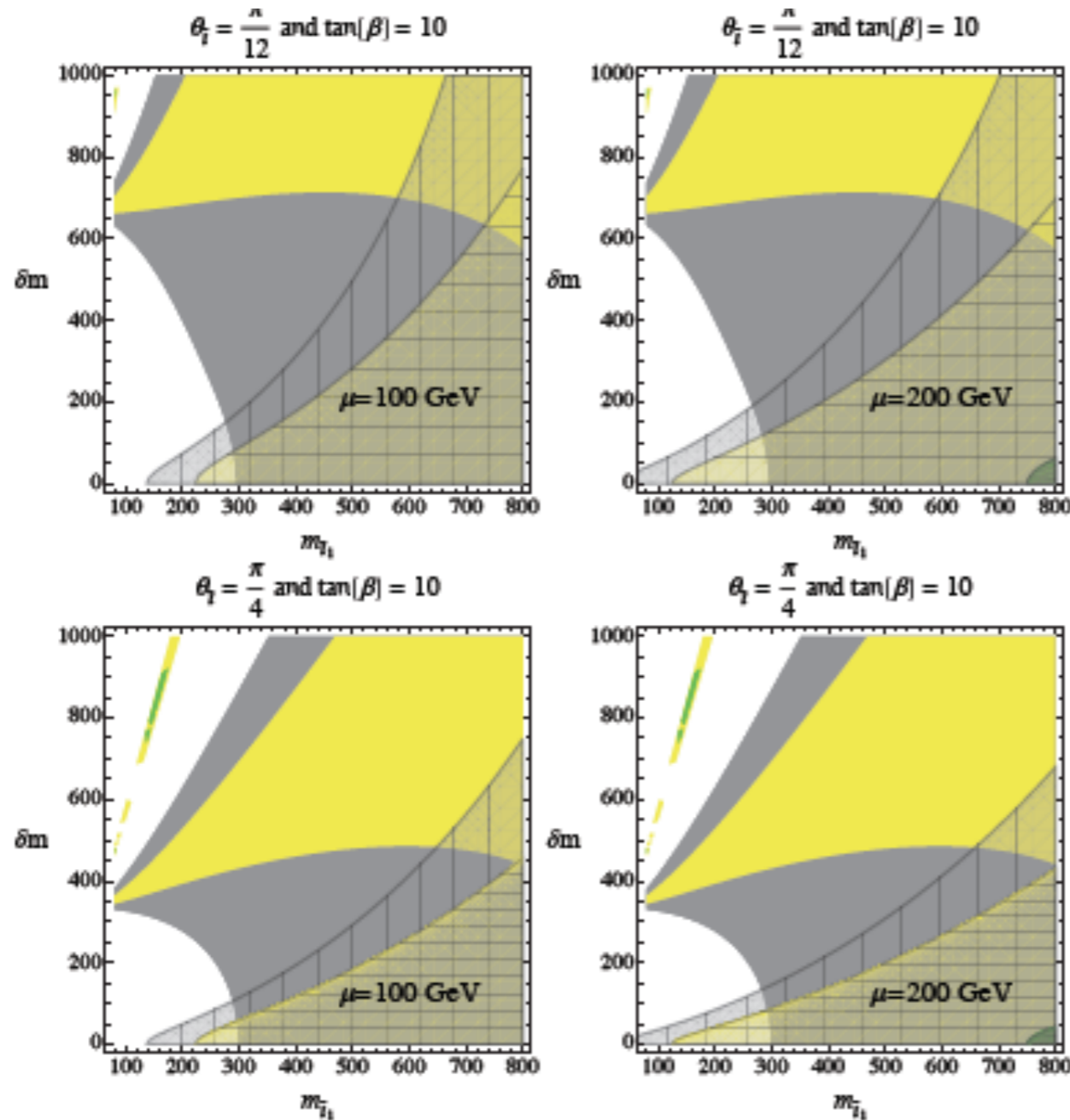
- 1σ
- 2σ
- 3σ

Map to stop space:

$$F_g = -\frac{1}{3} \left[\frac{m_t^2}{m_{\tilde{t}_1}^2} + \frac{m_t^2}{m_{\tilde{t}_2}^2} - \frac{1}{4} \sin^2(2\theta_t) \frac{\delta m^4}{m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} \right]$$

Light unmixed stops in bad shape, large mixing preferred in the data.

Consistency of B to s gamma



*Not so good news...
 μ dependance as the
 light chargino
 contributes here in
 NSUSY*

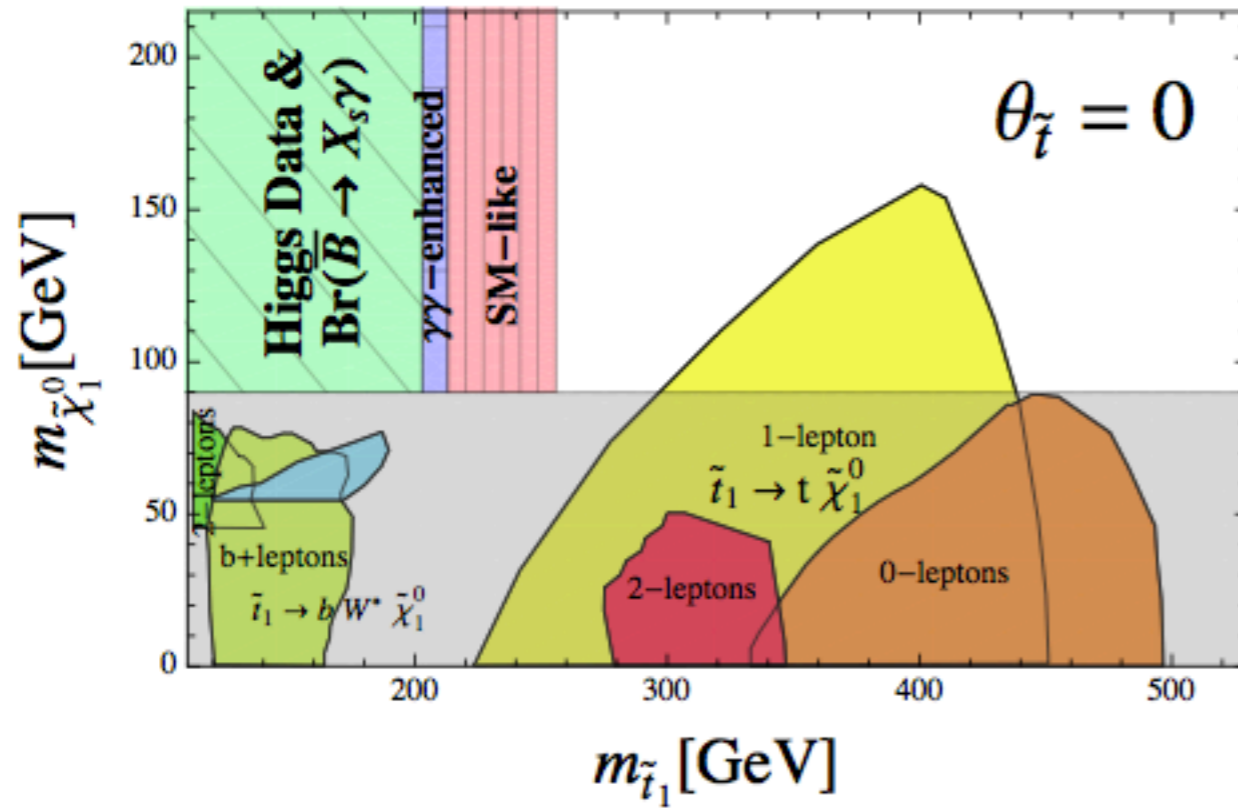
*$\text{Br}(\bar{B} \rightarrow X_s \gamma)$ basically wants
 degenerate stops in NSUSY
 spectrum we consider*

$$\text{BR}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6 \text{ GeV}} = [(3.15 \pm 0.23) - 8.0 \Delta C_7(\mu_0) - 1.9 \Delta C_8(\mu_0)] \times 10^{-4}$$

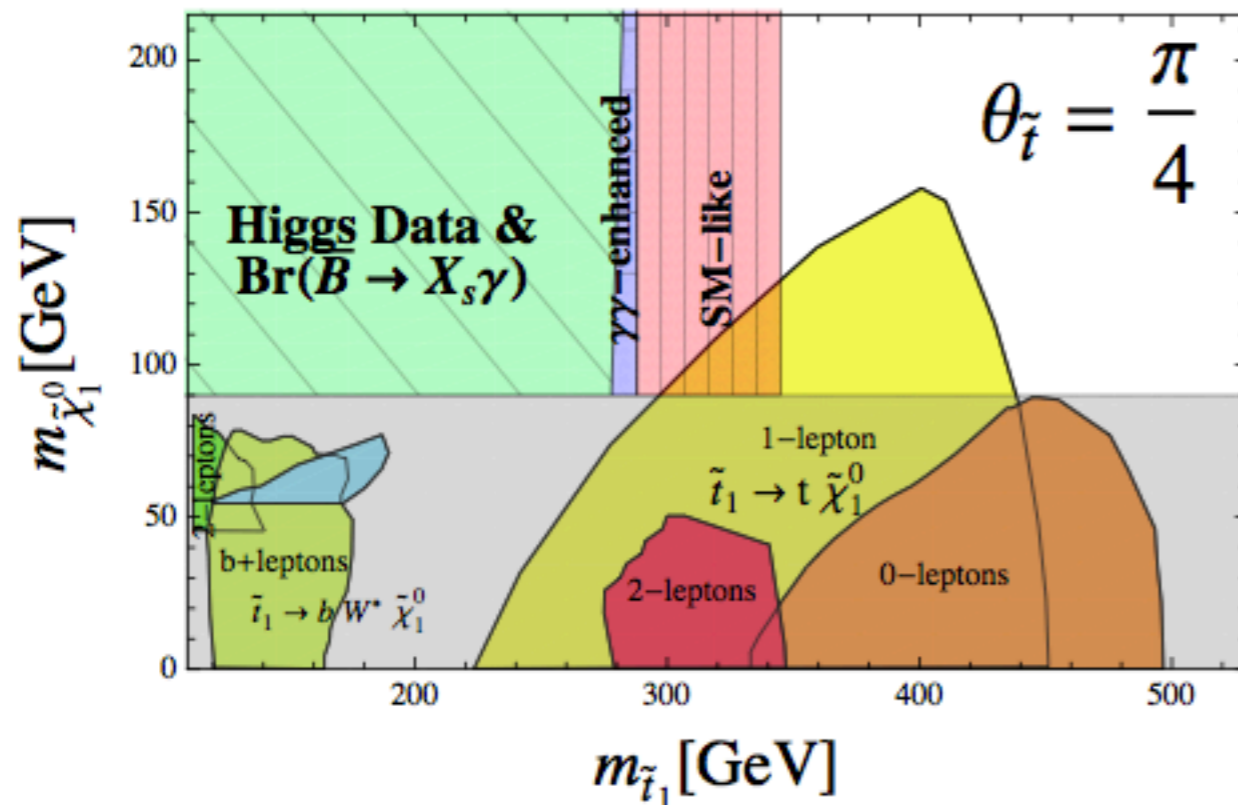
Theory prediction: B. Grzadkowski and M. Misiak, Phys. Rev. D **78**, 077501 (2008) [hep-ph/0802.1413].

$$\text{BR}(\bar{B} \rightarrow X_s \gamma)_{E_\gamma > 1.6 \text{ GeV}} = [3.43 \pm 0.21 \pm 0.07] \times 10^{-4}, \quad \text{Latest HFAAG combo number}$$

Indirect Exclusion currently and prospects for stops



Combined exclusion with demand that the higgs properties 95 % CL exclude and $\text{Br}(\bar{B} \rightarrow X_s \gamma)$ be within 2 sigma of its experimental value



Interpret with caution!

Indirect probes for stops are powerful tools. Dedicated experimental study warranted and underway now in ATLAS

Conclusions



Need more data.