General bounds on shot noise in the absence of currents: charge and heat transport

Janine Splettstößer



Applied Quantum Physics, MC2, Chalmers University of Technology

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Short introduction: – Charge and heat currents and noise.

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- Results: Charge and heat shot noise at zero currents in constant-transmission conductors.

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- Example for a concrete conductor.

Introduction - Charge and heat currents and noise

Introduction: transport and noise



Heat current:

$$J_{\mathsf{L}} = \frac{2}{h} \int dE \, (E - \mu_{\mathsf{L}}) D(E) [f_{\mathsf{L}}(E) - f_{\mathsf{R}}(E)]$$

G. Benenti, G. Casati, K. Saito, R. S. Whitney: Phys. Rep. 694, 1 (2017)

Introduction: transport and noise



Noise

- Further spectroscopy
- Important for "non-macroscopic" thermodynamics
- TUR, FR,...



Shot noise = partition noise



Ya. M. Blanter, M. Büttiker: Phys. Rep. 336, 1 (2000)

Introduction: Shot noise (in the absence of currents)

Shot and thermal noise:

$$S_{\text{th}}^{X} = \int dE \frac{4x^{2}}{h} D(E) \sum_{\alpha = \mathsf{L},\mathsf{R}} f_{\alpha}(E) [1 - f_{\alpha}(E)] \implies \text{Persists at equilibrium}$$
$$S_{\text{sh}}^{X} = \int dE \frac{4x^{2}}{h} D(E) [1 - D(E)] [f_{\mathsf{L}}(E) - f_{\mathsf{R}}(E)]^{2} \implies \text{Nonequilibrium effect}$$

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Nonequilibrium does not necessary imply an average current flow!!!

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Delta-T-noise

E. V. Sukhorukov, D. Loss: Phys. Rev. B 59, 13054 (1999)

• Recently measured!!

O. S. Lumbroso, L. Simine, A. Nitzan, D. Segal, O. Tal: Nature 562, 240 (2018)

E. Sivre, H. Duprez, A. Anthore, A. Aassime, F. D. Parmentier, A. Cavanna, A. Ouerghi, U. Gennser, F. Pierre: Nat. Commun. 10, 1 (2019)

• Limit for $T_{\rm L} \gg T_{\rm R} \approx 0$

$$S_{\mathsf{sh}}^l/S_{\mathsf{th}}^l
ightarrow \left(1-D^0
ight)\left(2\ln 2-1
ight)$$

S. Larocque, E. Pinsolle, C. Lupien, B. Reulet: Phys. Rev. Lett. 125, 106801 (2020)

Challenge to address



- \Rightarrow Find general statements on shot noise at zero currents
- \Rightarrow Particularly interesting for thermoelectrics

Zero-current charge and heat shot noise in constant-transmission conductors

Zero-current charge and heat shot noise at constant transmission



Zero-current charge and heat shot noise at constant transmission



Limit of $T_L \gg T_R \approx 0$

$$\frac{S_{sh}^{I}}{S_{th}^{I}}\Big|_{I=0} = R_{I} \to R_{I}^{0}$$

= (1 - D) (2 ln 2 - 1) \approx 0.4
$$\frac{S_{sh}^{J}}{S_{th}^{J}}\Big|_{I=0} = R_{J} \to R_{J}^{0}$$

= 3(1 - D)A($\pi/\sqrt{3}$)/ $\pi^{2} \approx 0.45$

Zero-current charge and heat shot noise at constant transmission



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Can not be exceeded at any temperature difference if D = const.

J. Eriksson, M. Acciai, L. Tesser, J. Splettstoesser: arXiv:2102.12988 (2021); accepted for publication in Phys. Rev. Lett.

General bounds on zero-current charge and heat shot noise

Charge transport Which D(E) maximizes S_{sh}^{l} ?

$$S_{sh}^{I} = \frac{4e^{2}}{h} \int dE D(E) [1 - D(E)] [f_{L}(E) - f_{R}(E)]^{2}$$

Maximal at constant transmission D = 1/2 — even if $I \equiv 0$ imposed!

Use variational principle as in: R. S. Whitney: Phys. Rev. B 91, 115425 (2015)

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Maximum of S_{sh}^{l} compared to S_{th}^{l} at I = 0?

$$S_{\mathsf{sh}}^{l} \leq rac{4e^2}{h} \int dED(E) f_\mathsf{L}(E) [1 - f_\mathsf{R}(E)] \leq S_{\mathsf{th}}^{l}$$

 $\Rightarrow R_I \equiv \left[S_{\rm sh}^I/S_{\rm th}^I\right]_{I=0} \le 1$

- Relies only on $0 \le D(E) \le 1$ and $0 \le f_{\alpha}(E) \le 1$
- Valid also for multi-channel contacts

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$$S_{\mathsf{sh}}^{\prime} \leq rac{4e^2}{h}\int dED(E)f_\mathsf{L}(E)[1-f_\mathsf{R}(E)] \leq S_{\mathsf{th}}^{\prime}$$

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 $\Rightarrow \mathsf{bound}(R_I) > (1-D)(2\ln 2 - 1) = R_I^0$

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Which D(E) maximizes R_I ? Can the bound be reached?Any $D(E) \neq 0$ within an energy interval δ , with $R_I \lesssim 1$ $max(T_L, T_R) \gg \delta/k_B \gg min(T_L, T_R)$ and $D(E) \ll 1$.

Heat shot noise is not bounded with respect to heat thermal noise!

Counter example:



- Counter example with finite transmission in separate energy intervalls!
 - ⇒ Major contributions to shot and thermal noise can be separately tuned!
- Different transport channels ⇔ different energy!

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 $R_J = \frac{S_{sh}^{\prime}}{S_{th}^{\prime}}\Big|_{J=0}$ can become arbitrarily large!

J. Eriksson, M. Acciai, L. Tesser, J. Splettstoesser: arXiv:2102.12988 (2021); accepted for publication in Phys. Rev. Lett.

Concrete example: Conductor with resonant (Lorentzian-shaped) transmission

Concrete example: Lorentzian transmission

Zero-current charge noise ratio:





Concrete example: Lorentzian transmission



Zero-current heat noise ratio:



J. Eriksson, M. Acciai, L. Tesser, J. Splettstoesser: arXiv:2102.12988 (2021); accepted for publication in Phys. Rev. Lett.

Concrete example: Lorentzian transmission



Zero-current heat noise ratio:



$R_J \leq 1$ for this specific transmission probability!

Test more complex transmissions to reach $R_J \ge 1!$ e.g. E. Hailloo, P. T. Alonso, N. Dashti, L. Arrachea,

e.g. F. Hajiloo, P. T. Alonso, N. Dashti, L. Arrachea, J. Splettstoesser: Phys. Rev. B **102**, 155434 (2020)

Conclusions

• Zero-current shot noise extended to heat transport!

Generic nonequilibrium conditions (beyond Delta-T noise), arbitrary transmission probabilities

- Simple expressions for limits of $R = S_{sh}/S_{th}$ at constant D
- General bounds on shot noise to thermal noise ratio at zero current Zero-current charge shot noise bounded by thermal noise; heat shot noise behaves fundamentally differently!!

Lorentzian transmission as example

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Lorentzian transmission as example

J. Eriksson, M. Acciai, L. Tesser, J. Splettstoesser: arXiv:2102.12988 (2021); accepted for publication in Phys. Rev. Lett.

Outlook

• Many relevant cases in quantum thermoelectrics where I, J = 0 but $S_I, S_J \neq 0$

R. Sánchez, J. Splettstoesser, R. S. Whitney: Nonequilibrium System as a Demon. Phys. Rev. Lett. 123, 216801 (2019)

Fluctuation dissipation theorem out of equilibrium?

B. Altaner, M. Polettini, M. Esposito: Fluctuation-Dissipation Relations Far from Equilibrium. Phys. Rev. Lett. 117, 180601 (2016)

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 Shot noise contribution from upper part:

$$S_{\mathsf{sh}}^{J} = rac{1}{h} \int_{\epsilon - \Delta/2}^{\epsilon + \Delta/2} dE E^{2} [f_{L} - f_{R}]^{2} \ pprox rac{\epsilon^{2} \Delta}{h} = \epsilon J^{+}$$

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Thermal noise contribution from lower part: Independent of *ϵ*!!

$$S_{\text{th}}^{\prime} pprox rac{4}{h} \int_{-\infty}^{0} dE E^2 f_L [1 - f_L]$$

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• Thermal noise contribution from lower part: Independent of *ϵ*!!

$$S_{\rm th}^J \approx \frac{4}{h} \int_{-\infty}^0 dE E^2 f_L [1 - f_L]$$

R_J can become arbitrarily large!

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