

Time-reversal symmetry breaking and Polar Kerr effect from charge order in cuprates

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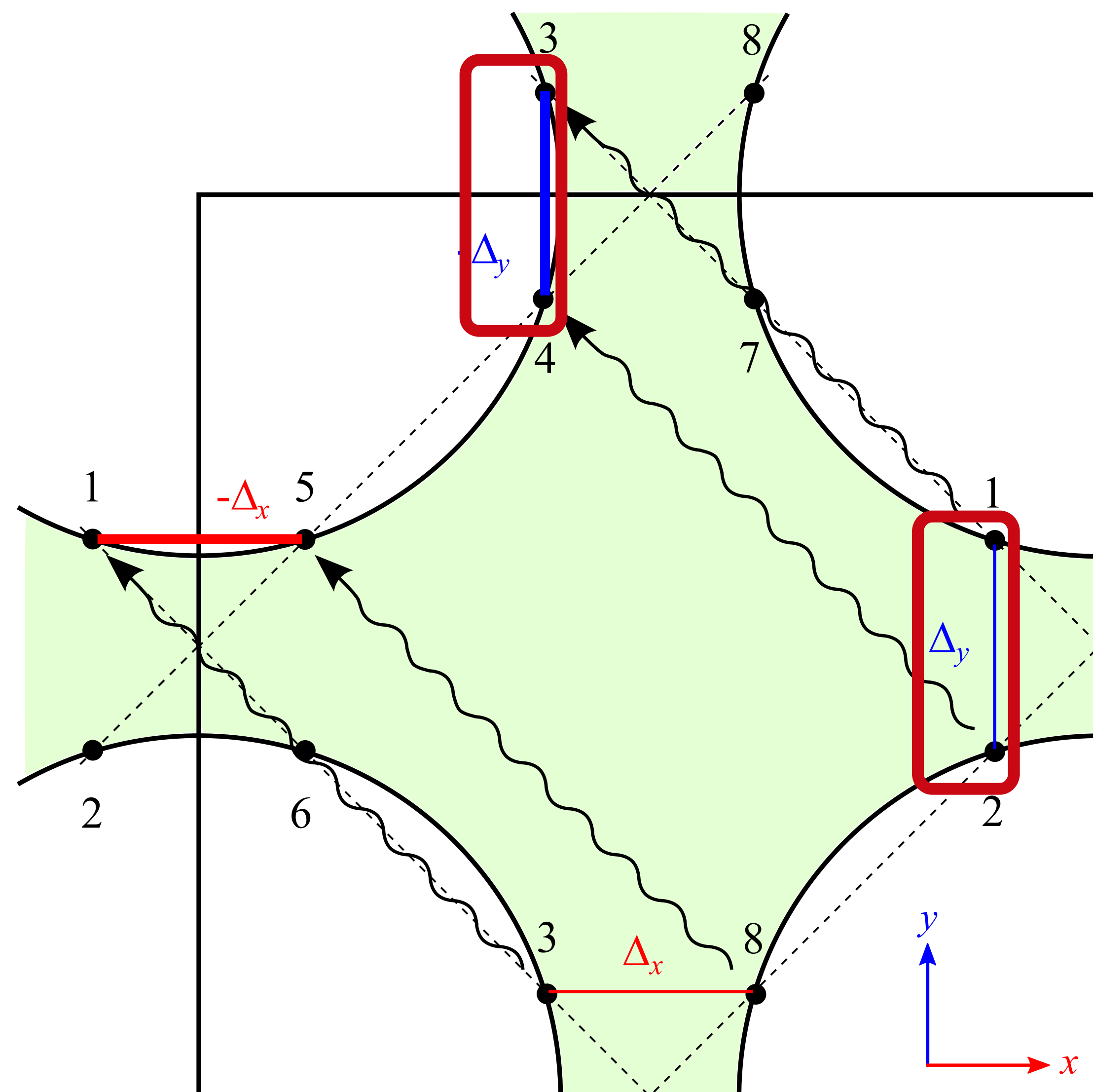
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Motivation

- Charge order has been ubiquitously found in hole doped cuprates.
- Polar Kerr effect has been observed in hole-doped cuprates. The most natural explanation is time-reversal symmetry breaking.
- Intra-unit-cell magnetic order has been observed in polarized neutron scattering, which indicates time-reversal symmetry is broken.

Model



Spin-fermion interaction (represented by wavy lines in the figure above)

$$\mathcal{S} = - \int_k^\Lambda G_0^{-1}(k) \psi_{k,\alpha}^\dagger \psi_{k,\alpha} + \frac{1}{2} \int_q^\Lambda \chi_0^{-1}(q) \mathbf{S}_q \cdot \mathbf{S}_{-q} + g \int_{k,q}^\Lambda \psi_{k+q,\alpha}^\dagger \sigma_{\alpha\beta} \psi_{k,\beta} \cdot \mathbf{S}_{-q}.$$

$$\Delta_{k_0}^Q(\Omega_m) = \frac{3\bar{g}T}{4\pi^2} \sum_{m',k} G_1(\omega_m, k) G_2(\omega_m, k) \frac{\Delta_{k_\pi}^Q(\Omega_{m'})}{k_x^2 + k_y^2 + \gamma|\Omega_m - \Omega_{m'}|}$$

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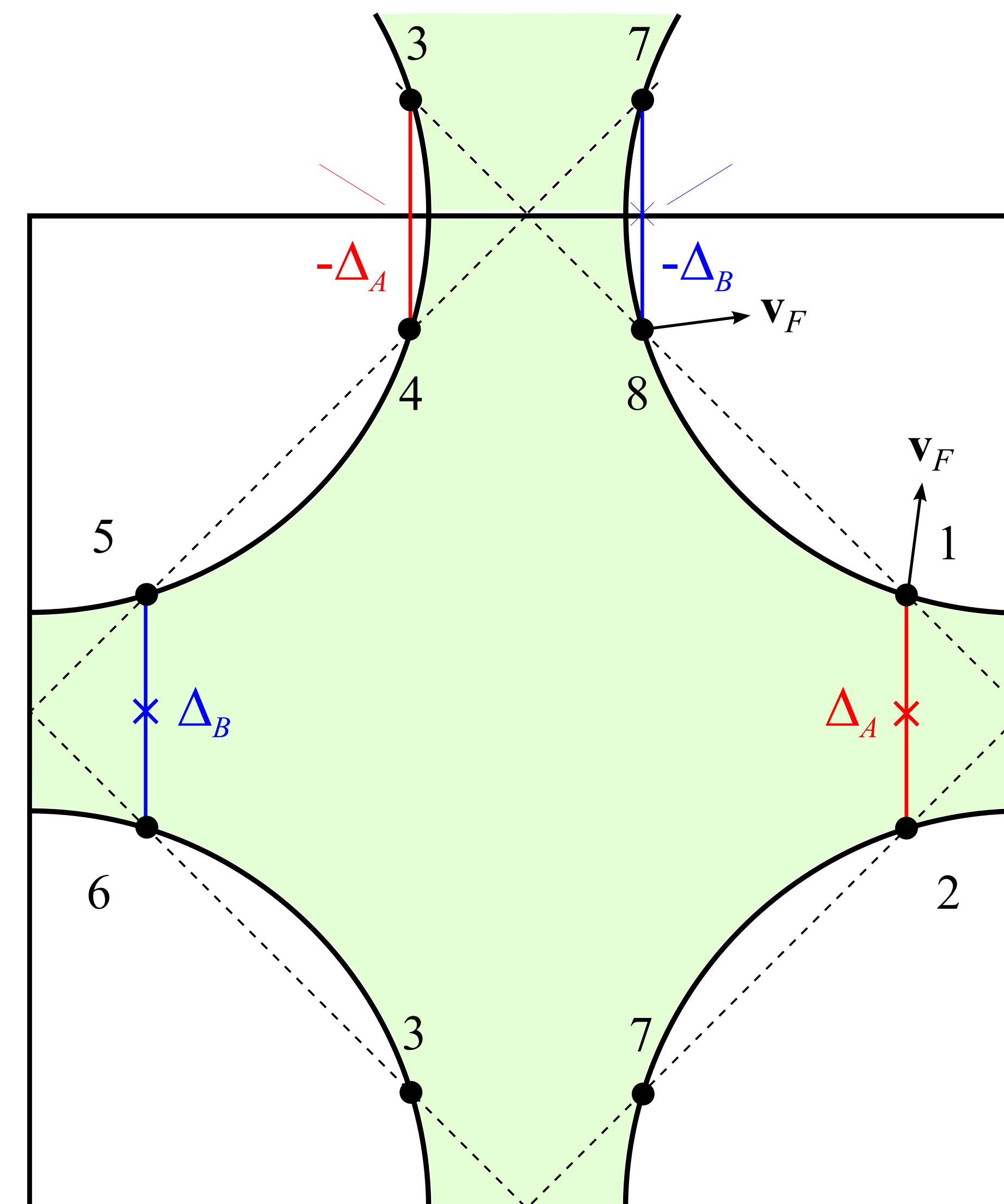
Instability comes from bosonic propagator — Landau damping term. Moving away from QCP the CDW will ultimately vanish.

Nematic transition

$$\mathcal{S}_{\text{eff}} = \alpha(|\Delta_x|^2 + |\Delta_y|^2) + \beta_1(|\Delta_x|^4 + |\Delta_y|^4) + 2\beta_2|\Delta_x|^2|\Delta_y|^2$$

- $\beta_2 > \beta_1$, so charge order develops in either x or y direction, but not in both (stripes).

TRSB transition



Δ_A and Δ_B are related by time reversal. Going beyond hot-spot analysis we find a coupling term between them

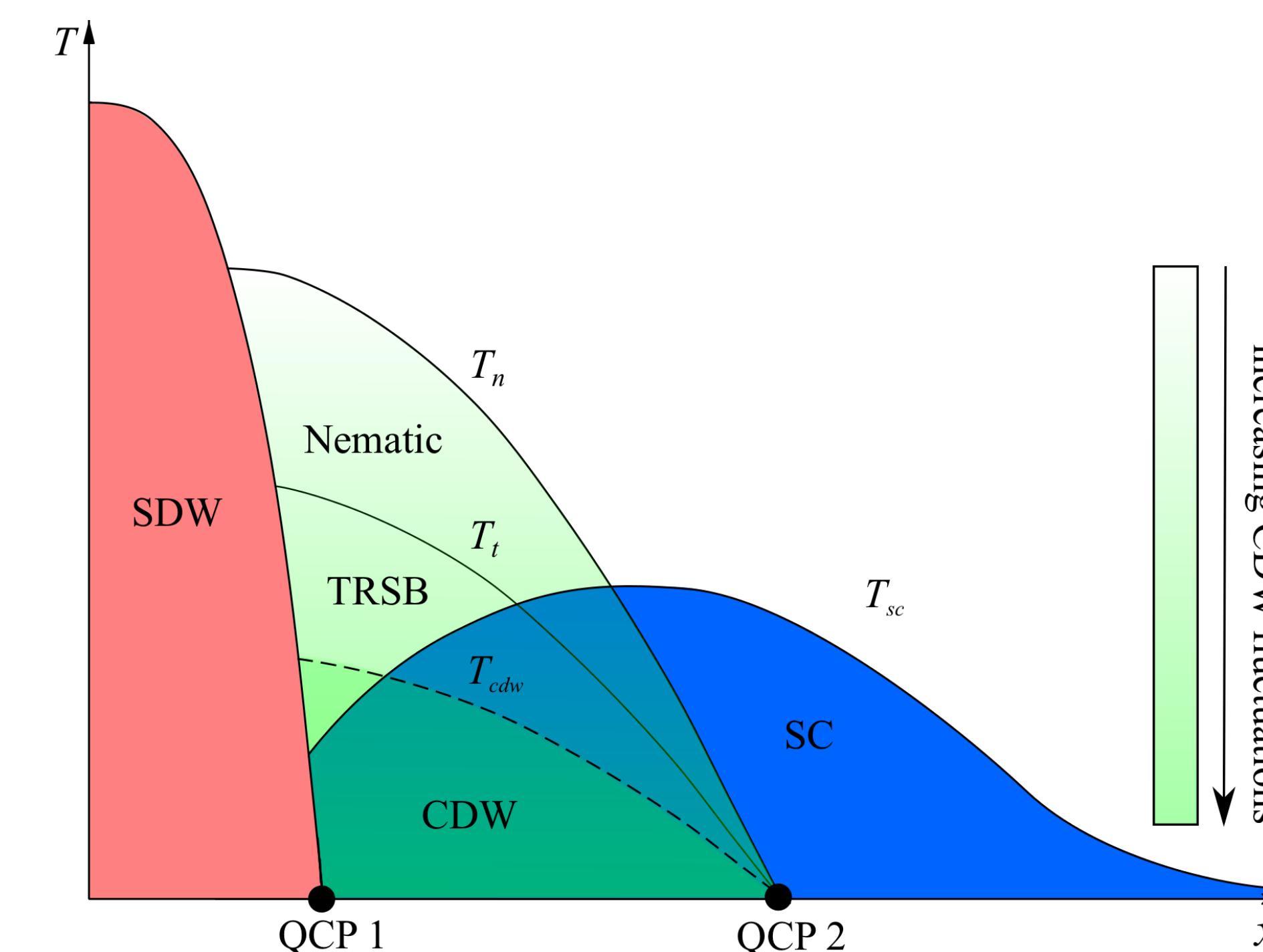
$$\mathcal{S}_{\text{eff}} = \alpha(|\Delta_A|^2 + |\Delta_B|^2) + \beta_1(|\Delta_A|^4 + |\Delta_B|^4) + \beta'[i(\Delta_A \Delta_B^* - \Delta_A^* \Delta_B)]^2$$

$\beta' > 0$, Δ_A and Δ_B develop with a relative $\pi/2$ phase difference. Time-reversal symmetry is broken. $\langle \Delta_A \Delta_B^* - \Delta_A^* \Delta_B \rangle \equiv \pm i\Upsilon$

Preemptive orders

Going beyond mean-field analysis we find that nematic transition and TRSB transition occurs at a higher temperature than the primary CDW transition. Even before CDW develops, C_4 rotation symmetry and time-reversal symmetry is already broken.

Phase diagram

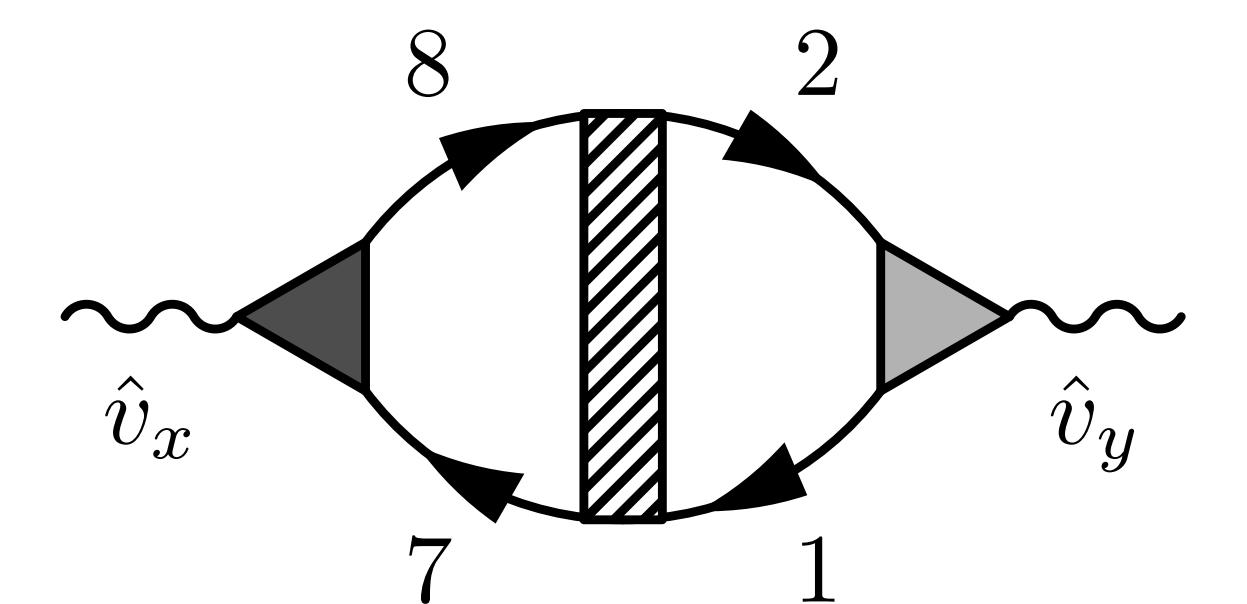


- The charge order terminates at a finite SDW correlation length (controlled by doping).
- Two preemptive orders, one nematic and one time-reversal symmetry breaking, set in prior to the primary charge density wave order.
- The charge-density-wave order is enhanced by preemptive orders, while SC is suppressed.

Polar Kerr effect from composite order

- Polar Kerr effect requires time-reversal symmetry and both mirror symmetries to be broken. → Satisfied!

	\mathcal{M}_x	\mathcal{M}_y	\mathcal{T}
Δ_A	Δ_A^*	Δ_B	Δ_B
Δ_B	Δ_B^*	Δ_A	Δ_A
$\Upsilon = -i(\Delta_A \Delta_B^* - \Delta_B \Delta_A^*)$	$-\Upsilon$	$-\Upsilon$	$-\Upsilon$

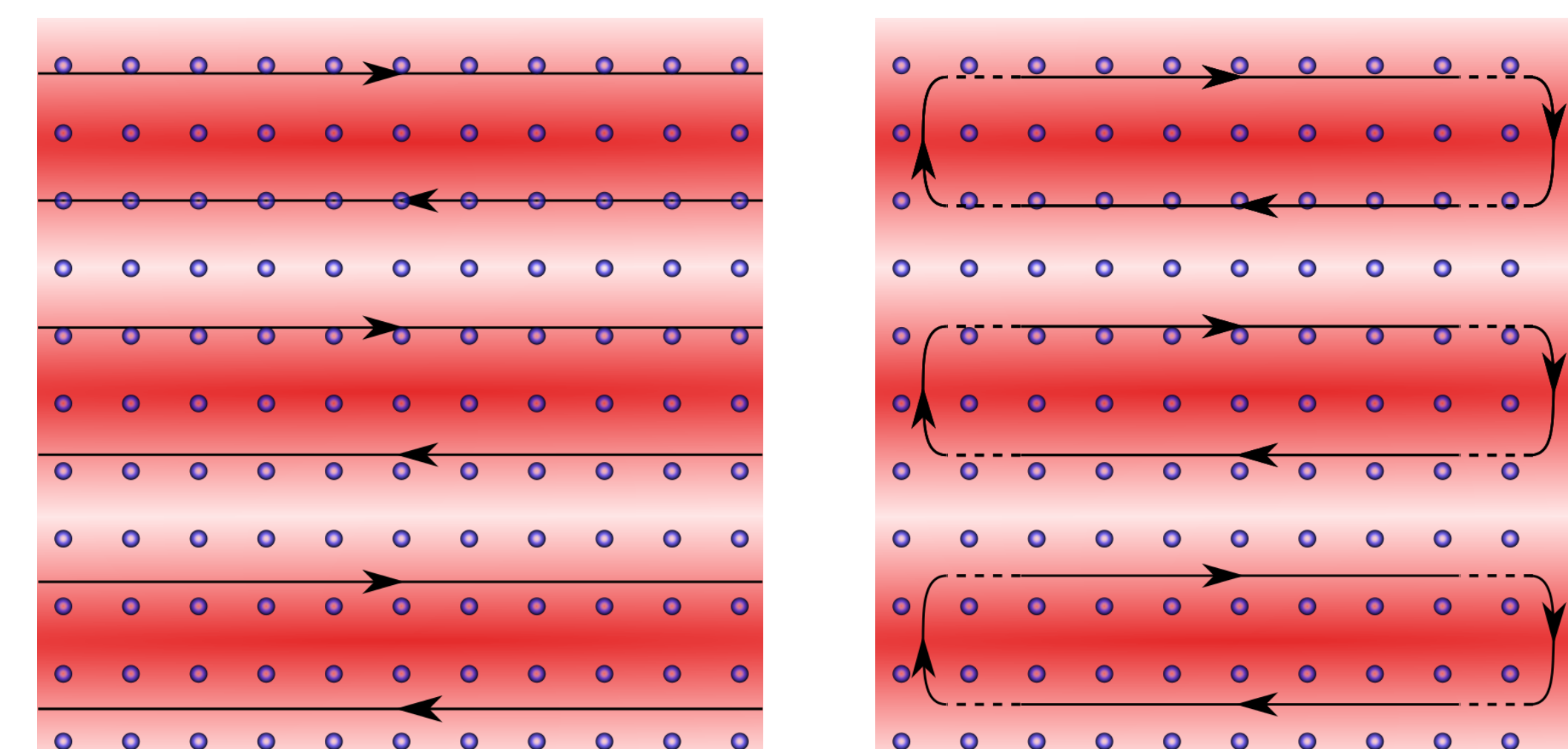


For a general system that is not rotationally invariant, the Kerr angle can be expressed by the antisymmetric component of the Hall conductivity.

$$\theta_K = \frac{\lambda}{c} \text{Im} \left[\frac{\sigma_{xy} - \sigma_{yx}}{n(n^2 - 1)} \right]$$

- If one uses linear dispersion for fermions then the interaction can be fulfilled by skew scattering by disorder, which is odd in frequency.
- Gaussian-type disorder and spin fluctuations can also give nonzero result, but one has to invoke particle-hole asymmetry (nonlinear dispersion).

The bulk current $\langle \Delta_A - \Delta_B \rangle$ has a finite momentum and does not couple to a uniform magnetic field. Symmetry-wise, the composite order Y can couple to a uniform magnetic field. However, we argue this coupling is via a surface current, which vanishes in thermodynamical limit.



References

- YW and AC, Phys Rev B **90** 035149 (2014)
- YW, AC and RN, arXiv:1409.5441 (2014)