# Lepton and Quark masses from Top 

loops

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Bogdan Dobrescu to appear...

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## Loopy masses for leptons and quarks

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## Important high energy physics questions?

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# $\mathrm{Q}:$ What sucks in the Standard Model?? 

## A:The Higgs

-James Wells

## Standard Model Higgs

Responsible for $\mathrm{W}, \mathrm{Z}$ mass and (charged) fermion masses
Associated hierarchies:

Gauge hierarchy
$m_{W} \ll M_{p l}$

Yukawa hierarchy
$y_{e} \ll y_{t}$

## Yukawa hierarchy

Technically natural but would still like an explanation
Symmetries (Froggatt Nielsen Models)

$$
Y_{i j}\left(\frac{\phi}{M}\right)^{q_{i}+q_{j}+q_{H}} H \bar{\psi}_{i} \psi_{j}
$$

$$
Y_{i j}^{S M}=Y_{i j} \epsilon^{q_{i}+q_{j}+q_{H}} \quad \epsilon=\frac{\langle\phi\rangle}{M}
$$

Charge the SM fermions differently

## Geography (Extra dim~n-:-nal models)



$$
Y_{i j}^{S M}=\int d x_{5} \psi_{i}\left(x_{5}\right) \psi_{j}\left(x_{5}\right) h\left(x_{5}\right)
$$



Place the SM fermions in different places
-The SM is coupled to a strongly coupled CFT

- SM fields get large anomalous dimensions
- Enters approximate fixed point at scale $\mu$ and leaves at scale $\mu_{0}$

$$
Y_{i j}^{S M}(\mu)=Y_{i j}\left(\mu_{0}\right)\left(\frac{\mu}{\mu_{0}}\right)^{\frac{1}{2}\left(\gamma_{i}+\gamma_{j}+\gamma_{H}\right)}
$$

SM fermions have different couplings

- Many clever mechanisms exist but must treat SM fermions separately.
- Convert small differences to large differences
-Example where SM fermions all charged the same way but get differences in Yukawas?


## Quantum mechanics

Masses are generated through quantum effects

Electron mass from muon mass?
Georgi and Glashow, '73

Work in the ` 80 's, mainly one and two loop mass generation

Babu and Ma, ' 89

## Quantum mechanics

Masses are generated through quantum effects

Electron mass from muon mass?

Work in the ` 80 's, mainly one and two loop mass generation

Babu and Ma, ' 89

Naively all masses at approximately the same
loop order



## More ambitious attempt



Loop-level where
mass is generated

PJF and Dobrescu



Loop-level where
mass is generated

## More likely to fail...?



Loop-level where
mass is generated

## More likely to fail...?



Loop-level where
mass is generated

## Top is clearly special

## So,

assume only the top has a tree level Yukawa

$$
y_{t} H \bar{u}_{R}^{3} Q_{L}^{3}
$$

## Top is clearly special

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$$
y_{t} H \bar{u}_{R}^{3} Q_{L}^{3}
$$

## Charge the top?

## Top is clearly special

## So,

assume only the top has a tree level Yukawa


Instead charge Higgs under an extra $U(1)$
$U(1)$ broken by the vev of a SM singlet $\phi$ of charge - I

Introduce a vector like pair of fermions with quantum numbers of left handed quarks, also charged under $U(1)$

Yukawas:

$$
m_{i j} \propto \tilde{c}_{i} c_{j}
$$

But lh top and rh top only appear linearly in couplings Redefine couplings so only one lh and one rh couple Call these the top

Mass matrix is rank I

## Only the top gets a tree level mass

## Chiral symmetries

$$
U(3)_{Q} \times U(3)_{u} \times U(3)_{d} \rightarrow U(1)_{t} \times U(2)_{Q} \times U(2)_{u} \times U(3)_{d}
$$

Need to break remaining chiral symmetries

Introduce a scalar leptoquark $\tilde{r}:(3,2,+7 / 6)$

Most general interactions

$$
\lambda_{i j} \tilde{r} \bar{u}_{R}^{i} L_{L}^{j}+\lambda_{i j}^{\prime} \tilde{r} \bar{Q}_{L}^{i} e_{R}^{j}+\text { Н.c. }
$$

## Chiral symmetries

$$
U(3)_{Q} \times U(3)_{u} \times U(3)_{d} \rightarrow U(1)_{t} \times U(2)_{Q} \times U(2)_{u} \times U(3)_{d}
$$

Need to break remaining chiral symmetries
Introduce a scalar leptoquark $\tilde{r}:(3,2,+7 / 6)$
(charge 0 under extra $U(1)$ )
Most general interactions

$$
\lambda_{i j} \tilde{r} \bar{u}_{R}^{i} L_{L}^{j}+\lambda_{i j}^{\prime} \tilde{r} \bar{Q}_{L}^{i} e_{R}^{j}+\text { Н.с. }
$$

$$
y_{t} \neq 0
$$

$$
U(3)_{Q} \times U(3)_{u} \times U(3)_{d} \rightarrow U(1)_{t} \times U(2)_{Q} \times U(2)_{u} \times U(3)_{d}
$$

$$
\begin{aligned}
& \lambda \neq 0 \\
& \xrightarrow{\lambda^{\prime} \neq 0} U(1)_{u} \times U(3)_{d} \\
& U(3)_{L} \times U(3)_{e} \xrightarrow[\substack{\lambda \neq 0 \\
\lambda^{\prime} \neq 0}]{\longrightarrow} U(1)_{L}
\end{aligned}
$$

With this breaking of chiral symmetries up type quarks and charged leptons can get a mass at some loop order

$$
y_{t} \neq 0
$$

$$
U(3)_{Q} \times U(3)_{u} \times U(3)_{d} \rightarrow U(1)_{t} \times U(2)_{Q} \times U(2)_{u} \times U(3)_{d}
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\end{aligned}
$$

With this breaking of chiral symmetries up type quarks and charged leptons can get a mass at some loop order

## But what loop order?



Linear couplings

## Redefine fields:

$$
\left(\begin{array}{lll}
\lambda_{11} & \lambda_{12} & \lambda_{13} \\
\lambda_{21} & \lambda_{22} & \lambda_{23} \\
\lambda_{31} & \lambda_{32} & \lambda_{33}
\end{array}\right)
$$



Linear couplings
Redefine fields:


Linear couplings
Redefine fields:
-Define $L_{3}$ so it only couples only to $u_{3}$

$$
\left(\begin{array}{ccc}
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\lambda_{21} & \lambda_{22} & \lambda_{23} \\
0 & 0 & \lambda_{33}
\end{array}\right)
$$

$$
\lambda_{i j} \tilde{r} \bar{u}_{R}^{i} L_{L}^{j}+\lambda_{i j}^{\prime} \tilde{r} \bar{Q}_{L}^{i} e_{R}^{j}+\text { Н.с. }
$$

Linear couplings

## Redefine fields:

-Define $L_{3}$ so it only couples only to $u_{3}$

- $u_{2}$ couples only to $L_{2}$ and $L_{3}$

$$
\left(\begin{array}{ccc}
\lambda_{11} & \lambda_{12} & \lambda_{13} \\
0 & \lambda_{22} & \lambda_{23} \\
0 & 0 & \lambda_{33}
\end{array}\right)
$$



Linear couplings
Redefine fields:
-Define $L_{3}$ so it only couples only to $u_{3}$

- $u_{2}$ couples only to $L_{2}$ and $L_{3}$
-Rotation of $u_{1}$ and $u_{2}$

$$
\left(\begin{array}{ccc}
\lambda_{11} & \lambda_{12} & 0 \\
0 & \lambda_{22} & \lambda_{23} \\
0 & 0 & \lambda_{33}
\end{array}\right)
$$

$$
\lambda_{i j} \tilde{r} \bar{u}_{R}^{i} L_{L}^{j}+\lambda_{i j}^{\prime} \tilde{r} \bar{Q}_{L}^{i} e_{R}^{j}+\text { Н.с. }
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Linear couplings
Redefine fields:
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-Rotation of $u_{1}$ and $u_{2}$

$$
\left(\begin{array}{ccc}
\lambda_{11} & \lambda_{12} & 0 \\
0 & \lambda_{22} & \lambda_{23} \\
0 & 0 & \lambda_{33}
\end{array}\right)
$$

$\lambda_{i j}, \lambda_{i j}^{\prime}$
can be made real and positive

## One loop tau mass



$$
m_{\tau} \simeq \lambda_{33} \lambda_{33}^{\prime} m_{t} \underbrace{\frac{N_{c}}{16 \pi^{2}} \ln \left(\frac{\Lambda^{2}}{M_{\tilde{r}}^{2}}\right)}
$$

$$
\approx 0.09 \text { for } \Lambda \approx 10 M_{\tilde{r}}
$$

$\lambda_{33} \lambda_{33}^{\prime} \approx(0.36)^{2}$ for correct $m_{\tau} / m_{t}$ ratio

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## One loop tau mass



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$$

$\approx 0.09$ for $\Lambda \approx 10 M_{\tilde{r}}$
$\lambda_{33} \lambda_{33}^{\prime} \approx(0.36)^{2}$ for correct $m_{\tau} / m_{t}$ ratio

Two loop charm mass - a "rainbow" diagram


$$
\begin{aligned}
M_{u}[\tilde{r} \tilde{r}] & =\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & \lambda_{23}^{\prime} \lambda_{23} & \lambda_{33}^{\prime} \lambda_{23} \\
0 & \lambda_{23}^{\prime} \lambda_{33} & \lambda_{33}^{\prime} \lambda_{33}
\end{array}\right) \lambda_{33}^{\prime} \lambda_{33} m_{t} \epsilon_{\tilde{r}}^{(2)} \\
m_{c} & =\lambda_{23}^{\prime} \lambda_{23} m_{\tau} \frac{1}{16 \pi^{2}} \log \frac{\Lambda^{2}}{M_{\tilde{r}}^{2}}
\end{aligned}
$$

$\lambda_{23} \lambda_{23}^{\prime} \approx(3.3)^{2}$ for correct $m_{c} / m_{\tau}$ ratio

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Two loop charm mass - a "rainbow" diagram


$$
\begin{aligned}
M_{u}[\tilde{r} \tilde{r}] & =\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & \lambda_{23}^{\prime} \lambda_{23} & \lambda_{33}^{\prime} \lambda_{23} \\
0 & \lambda_{23}^{\prime} \lambda_{33} & \lambda_{33}^{\prime} \lambda_{33}
\end{array}\right) \lambda_{33}^{\prime} \lambda_{33} m_{t} \epsilon_{\tilde{r}}^{(2)} \\
m_{c} & =\lambda_{23}^{\prime} \lambda_{23} m_{\tau} \frac{1}{16 \pi^{2}} \log \frac{\Lambda^{2}}{M_{\tilde{r}}^{2}}
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Two loop charm mass - a "rainbow" diagram


$$
\begin{aligned}
M_{u}[\tilde{r} \tilde{r}] & =\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & \lambda_{23}^{\prime} \lambda_{23} & \lambda_{33}^{\prime} \lambda_{23} \\
0 & \lambda_{23}^{\prime} \lambda_{33} & \lambda_{33}^{\prime} \lambda_{33}
\end{array}\right) \lambda_{33}^{\prime} \lambda_{33} m_{t} \epsilon_{\tilde{r}}^{(2)} \\
m_{c} & =\lambda_{23}^{\prime} \lambda_{23} m_{\tau} \frac{1}{16 \pi^{2}} \log \frac{\Lambda^{2}}{M_{\tilde{r}}^{2}}
\end{aligned}
$$

$\lambda_{23} \lambda_{23}^{\prime} \approx(3.3)^{2}$ for correct $m_{c} / m_{\tau}$ ratio

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## Three loop muon mass




The diagram with no name $\sim N_{C}$

Three loop muon mass


The diagram with no name $\sim N_{C}$

## Three loop muon mass



$$
m_{\mu} \approx \lambda_{22}^{\prime} \lambda_{22} m_{c}(1+x) \frac{N_{c}}{16 \pi^{2}} \log \frac{\Lambda^{2}}{M_{\bar{r}}^{2}}
$$

$\lambda_{22} \lambda_{22}^{\prime}(1+x) \approx(1.5)^{2}$ for correct $m_{\mu} / m_{c}$ ratio

## Three loop muon mass



$$
m_{\mu} \approx \lambda_{22}^{\prime} \lambda_{22} m_{c}(1+x) \frac{N_{c}}{16 \pi^{2}} \log \frac{\Lambda^{2}}{M_{\bar{r}}^{2}}
$$

$\lambda_{22} \lambda_{22}^{\prime}(1+x) \approx(1.5)^{2}$ for correct $m_{\mu} / m_{c}$ ratio

## Four loop up quark mass



## +4 other diagrams

Muon mass implies: $\quad \# \lambda_{12} \lambda_{12}^{\prime} \approx(0.6)^{2}$

## Four loop up quark mass

## Three loop muon mass

## +4 other diagrams

Muon mass implies: $\quad \# \lambda_{12} \lambda_{12}^{\prime} \approx(0.6)^{2}$

Five loop electron mass
If only source of electron mass will determine $\lambda_{11} \lambda_{11}^{\prime}$

$$
\text { Only input: } \tilde{r}:(3,2,+7 / 6) \quad \lambda_{i j} \tilde{r} \bar{u}_{R}^{i} L_{L}^{j}+\lambda_{i j}^{\prime} \tilde{r} \bar{Q}_{L}^{i} e_{R}^{j}+\text { H.c. }
$$



Loop-level where mass is generated

## Down quark masses

Need to break the remaining chiral symmetries

$$
U(3)_{d} \times U(1)_{u} \times U(1)_{L}
$$

Have choices diquarks, leptoquarks...

$$
\begin{aligned}
H_{8}: & (8,2,-1 / 2) \\
\tilde{q}: & (3,2,1 / 6) \\
\tilde{d}_{6}: & (\overline{6}, 1,-1 / 3) \\
\tilde{d}: & (3,1,-1 / 3)
\end{aligned}
$$

## New field content

|  | $\phi$ | $\psi_{L}, \psi_{R}$ | $H$ | $r$ | $r^{\prime}$ | $H_{8}$ | $H_{8}^{\prime}$ | $\Phi_{3}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $S U(3)$ | 1 | 3 | 1 | 3 | 3 | 8 | 8 | $\overline{3}$ |
| $S U(2)$ | 1 | 2 | 2 | 2 | 2 | 2 | 2 | 2 |
| $U(1)_{Y}$ | 0 | $1 / 6$ | $1 / 2$ | $7 / 6$ | $7 / 6$ | $1 / 2$ | $-1 / 2$ | $-1 / 6$ |
| $U(1)^{\prime}$ | -1 | -1 | 1 | 0 | 2 | 1 | 1 | 0 |
| $\uparrow$ |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  | $\uparrow$ |  |

Up quarks and leptons
Down quarks

Most general couplings

$$
\begin{array}{r}
\kappa_{i} \Phi_{8} \bar{u}_{R}^{i} \Psi_{L}+\kappa^{\prime} \Phi_{8}^{\prime} \bar{d}_{R}^{3} \Psi_{L} \\
\quad \eta_{i j} \Phi_{3} \bar{d}_{R}^{i} L_{L}^{j}+\text { h.c. }
\end{array}
$$

break the remaining chiral symmetries

$$
U(3)_{d} \times U(1)_{u} \times U(1)_{L} \rightarrow U(1)_{L} \times U(1)_{Q}
$$

Most general couplings

$$
\begin{aligned}
& \kappa_{i} \Phi_{8} \bar{u}_{R}^{i} \Psi_{L}+\kappa^{\prime} \Phi_{8}^{\prime} \bar{d}_{R}^{3} \Psi_{L} \\
& \quad \eta_{i j} \Phi_{3} \bar{d}_{R}^{i} L_{L}^{j}+\text { h.c. }
\end{aligned}
$$

break the remaining chiral symmetries

$$
U(3)_{d} \times U(1)_{u} \times U(1)_{L} \rightarrow U(1)_{L} \times U(1)_{Q}
$$

Without altering up type and leptons have the freedom to rotate such that,

$$
\begin{gathered}
\eta=\left(\begin{array}{ccc}
\eta_{11} & \eta_{12} & 0 \\
\eta_{21} & \eta_{22} & \eta_{23} \\
\eta_{31} & \eta_{32} & \eta_{33}
\end{array}\right) \\
\kappa=\left(\kappa_{1}, \kappa_{2}, \kappa_{3}\right)
\end{gathered}
$$

Without altering up type and leptons have the freedom to rotate such that,

$$
\eta=\left(\begin{array}{ccc}
\eta_{11} & \eta_{12} & 0 \\
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\end{array}\right)
$$

Diagonal entries can be made real and positive

$$
\kappa=\left(\kappa_{1}, \kappa_{2}, \kappa_{3}\right)
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Without altering up type and leptons have the freedom to rotate such that,

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\eta=\left(\begin{array}{ccc}
\eta_{11} & \eta_{12} & 0 \\
\eta_{21} & \eta_{22} & \eta_{23} \\
\eta_{31} & \eta_{32} & \eta_{33}
\end{array}\right)
$$

Diagonal entries can be made real and positive

$$
\kappa=\left(\kappa_{1}, \kappa_{2}, \kappa_{3}\right)
$$

Entries can be made real and positive

## One loop bottom mass



$$
m_{b} \approx m_{t} \kappa^{\prime} \frac{N_{C}}{16 \pi^{2}}\left(\frac{\langle\phi\rangle}{M_{8}}\right)^{2} \log \frac{\Lambda^{2}}{M_{8}^{2}}
$$

## One loop bottom mass

$$
m_{b} \approx m_{t} \kappa^{\prime} \frac{N_{C}}{16 \pi^{2}}\left(\frac{\langle\phi\rangle}{M_{8}}\right)^{2} \log \frac{\Lambda^{2}}{M_{8}^{2}}
$$

## One loop bottom mass



Three loop strange mass


## Four loop down masses

## The down has a 4 loop mixed diagram (exercise for reader)

"Cross Talk"

There are also corrections to some of the states that have mass:

Charm gets a two loop correction
Up gets a four loop correction

Muon gets a three loop correction
Electron gets a five loop correction

Lepton and Quark masses at 1 TeV


Lepton and Quark masses at 1 TeV


Lepton and Quark masses at 1 TeV


Lepton and Quark masses at 1 TeV


Lepton and Quark masses at 1 TeV


Lepton and Quark masses at 1 TeV


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## CKM

$m_{u} \approx m_{t}\left(\begin{array}{ccc}\epsilon^{4} & \epsilon^{2} & \epsilon^{2} \\ \epsilon^{4} & \epsilon^{2} & \epsilon^{2} \\ \epsilon^{4} & \epsilon^{2} & 1\end{array}\right) \quad m_{d} \approx m_{t}\left(\begin{array}{ccc}\epsilon^{4} & \epsilon^{4} & \epsilon^{4} \\ \epsilon^{4} & \epsilon^{3} & \epsilon^{3} \\ \epsilon^{4} & \epsilon^{3} & \epsilon\end{array}\right)$

Resulting in

$$
V_{C K M} \approx\left(\begin{array}{ccc}
1-\epsilon^{2} & \epsilon & \epsilon^{3} \\
-\epsilon & 1-\epsilon^{2} & \epsilon^{2} \\
\epsilon^{3} & \epsilon^{2} & 1
\end{array}\right)
$$

Still to think about phases...

## James Wells:

## Q:What sucks in the Standard Model??

## A:The Higgs

## James Wells:

# $\mathrm{Q}:$ What sucks in the Standard Model?? 

## A:The Higgs

Q:Does the solution predict LHC physics?

The model contains extra fermions and scalar Leptoquarks
(Alternative realisation contains diquarks - easier to see at LHC than TeVatron)

## Constraints

Tree level exchange of leptoquark can lead to flavour changing processes e.g.

$$
\begin{aligned}
K^{+} & \rightarrow \pi^{0} \mu^{+} \mu^{-} \\
\tau^{+} & \rightarrow K^{0} e^{+} \\
\pi^{+} & \rightarrow e^{+} \nu \text { versus } \pi^{+} \rightarrow \mu^{+} \nu \\
\mu & \rightarrow e \text { conversion }
\end{aligned}
$$

$$
M_{\tilde{r}} \gtrsim 5-50 \mathrm{TeV}
$$

## Conclusions

-Fermions have complicated mass hierarchy

- Many attempts exist to explain it
- Top is probably special, perhaps only top mass has a tree level Yukawa
-With extra scalars coupling to fermions top mass is communicated at loop level
- Interesting structure of fermion mass spectrum arises
-Predicts flavour changing processes


## Conclusions

-Fermions have complicated mass hierarchy

- Many attempts exist to explain it
- Top is probably special, perhaps only top mass has a tree level Yukawa
-With extra scalars coupling to fermions top mass is communicated at loop level
- Interesting structure of fermion mass spectrum arises
-Predicts flavour changing processes
-Project X?

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