## Lepton and Quark masses from Top loops

Patrick Fox

## **Fermilab**

Bogdan Dobrescu to appear...



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# Loopy masses for leptons and quarks

## Patrick Fox

Bogdan Dobrescu to appear...



## Important high energy physics questions?



Important high energy physics questions?

## Q:What sucks in the Standard Model??

## A:The Higgs

-James Wells



Standard Model Higgs

Responsible for W, Z mass and (charged) fermion masses

Associated hierarchies:

Gauge hierarchy

 $m_W \ll M_{pl}$ 

Yukawa hierarchy

 $y_e \ll y_t$ 



Technically natural but would still like an explanation

Symmetries (Froggatt Nielsen Models)

$$Y_{ij}\left(\frac{\phi}{M}\right)^{q_i+q_j+q_H}H\bar{\psi}_i\psi_j$$

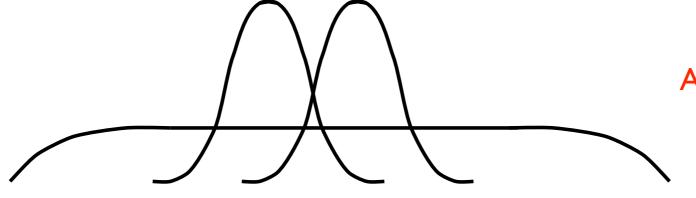
$$Y_{ij}^{SM} = Y_{ij} \,\epsilon^{q_i + q_j + q_H} \qquad \epsilon = \frac{\langle \varphi \rangle}{M}$$

Charge the SM fermions differently



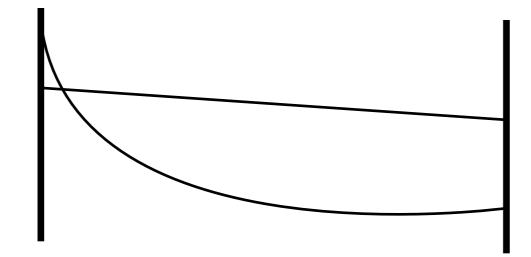
 $/ \perp$ 

## <u>Geography</u> (Extra dimanaional models)



Arkani-hamed, Schmaltz

$$Y_{ij}^{SM} = \int dx_5 \,\psi_i(x_5)\psi_j(x_5)h(x_5)$$



Place the SM fermions in different places



•The SM is coupled to a strongly coupled CFT •SM fields get large anomalous dimensions •Enters approximate fixed point at scale  $\mu$  and leaves at scale  $\mu_0$ 

$$Y_{ij}^{SM}(\mu) = Y_{ij}(\mu_0) \left(\frac{\mu}{\mu_0}\right)^{\frac{1}{2}(\gamma_i + \gamma_j + \gamma_H)}$$

SM fermions have different couplings



Many clever mechanisms exist but must treat SM fermions separately.

•Convert small differences to large differences

•Example where SM fermions all charged the same way but get differences in Yukawas?



Masses are generated through quantum effects

Electron mass from muon mass? Georgi and Glashow, `73

## Work in the `80's, mainly one and two loop mass generation

Babu and Ma, `89



Masses are generated through quantum effects

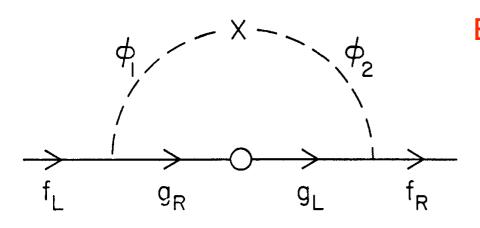
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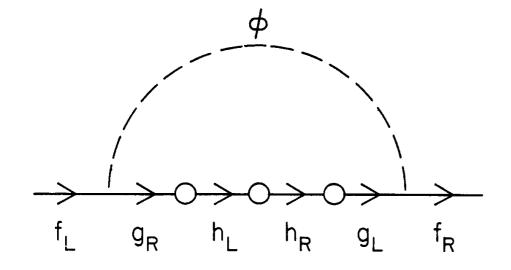
Work in the `80's, mainly one and two loop mass generation

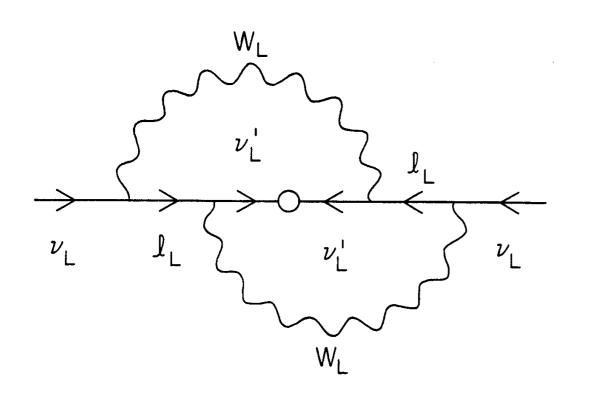
Babu and Ma, `89

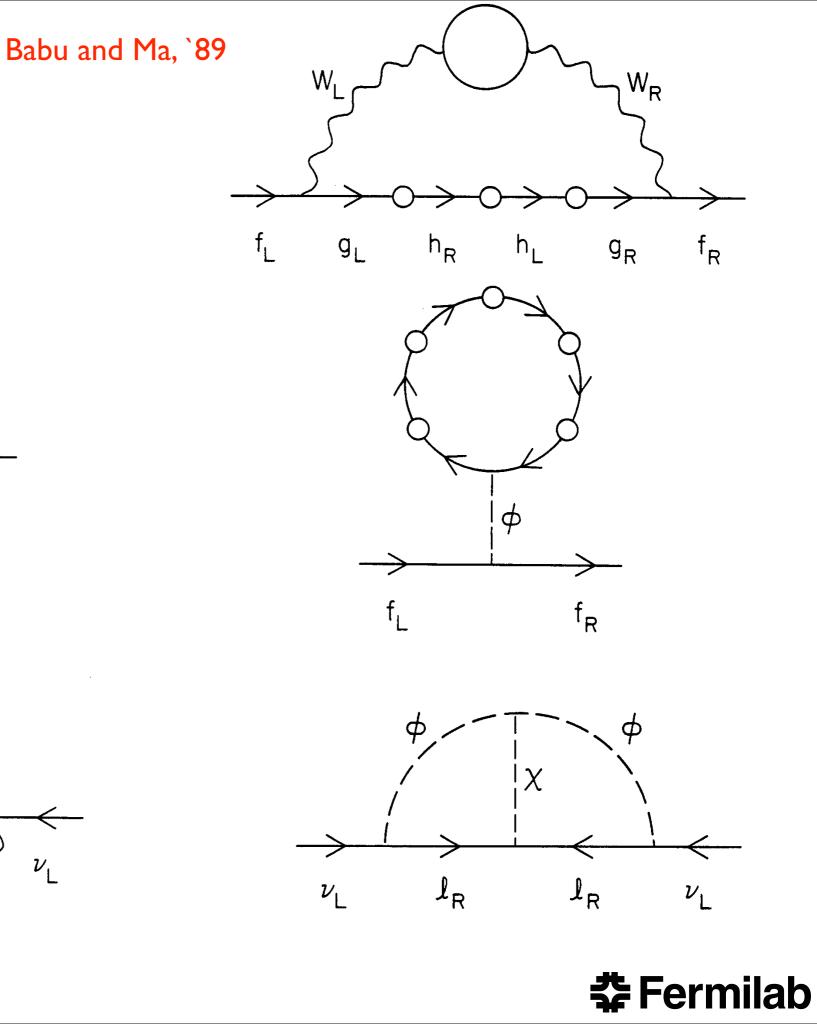
Naively all masses at approximately the same loop order











## More ambitious attempt

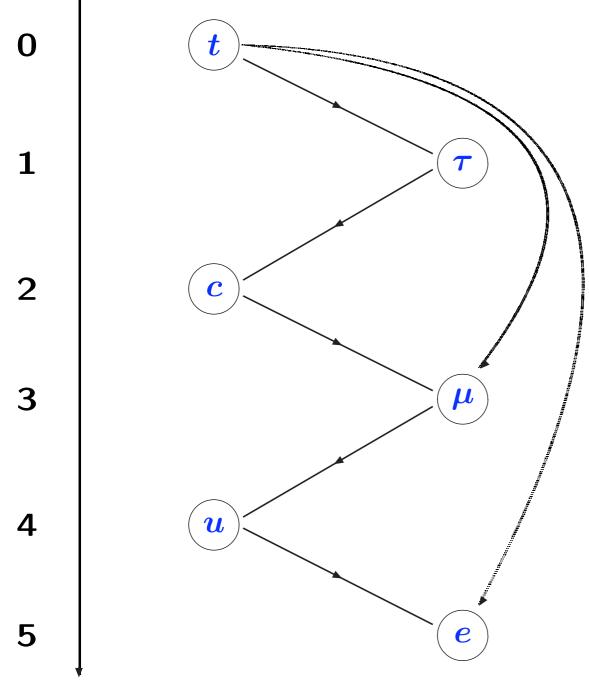
## 0 1 $\boldsymbol{ au}$ 2 *C* 3 $\boldsymbol{\mu}$ 4 U 5 e

Loop-level where mass is generated

#### PJF and Dobrescu



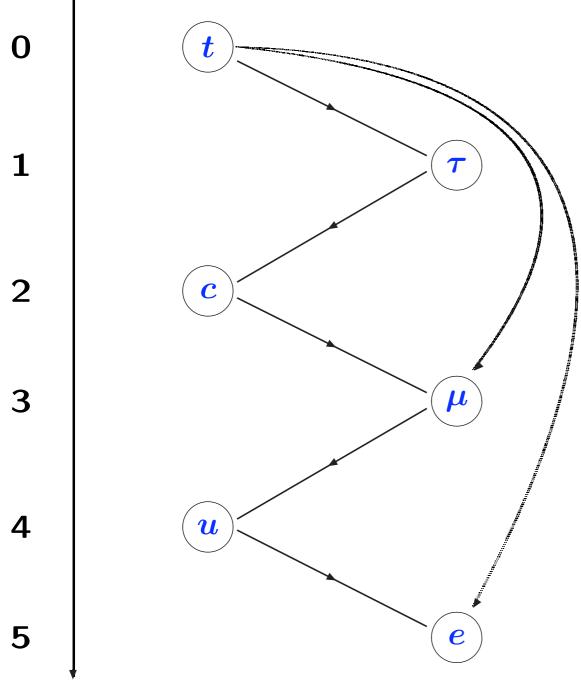
PJF and Dobrescu



Loop-level where mass is generated



## More likely to fail...?



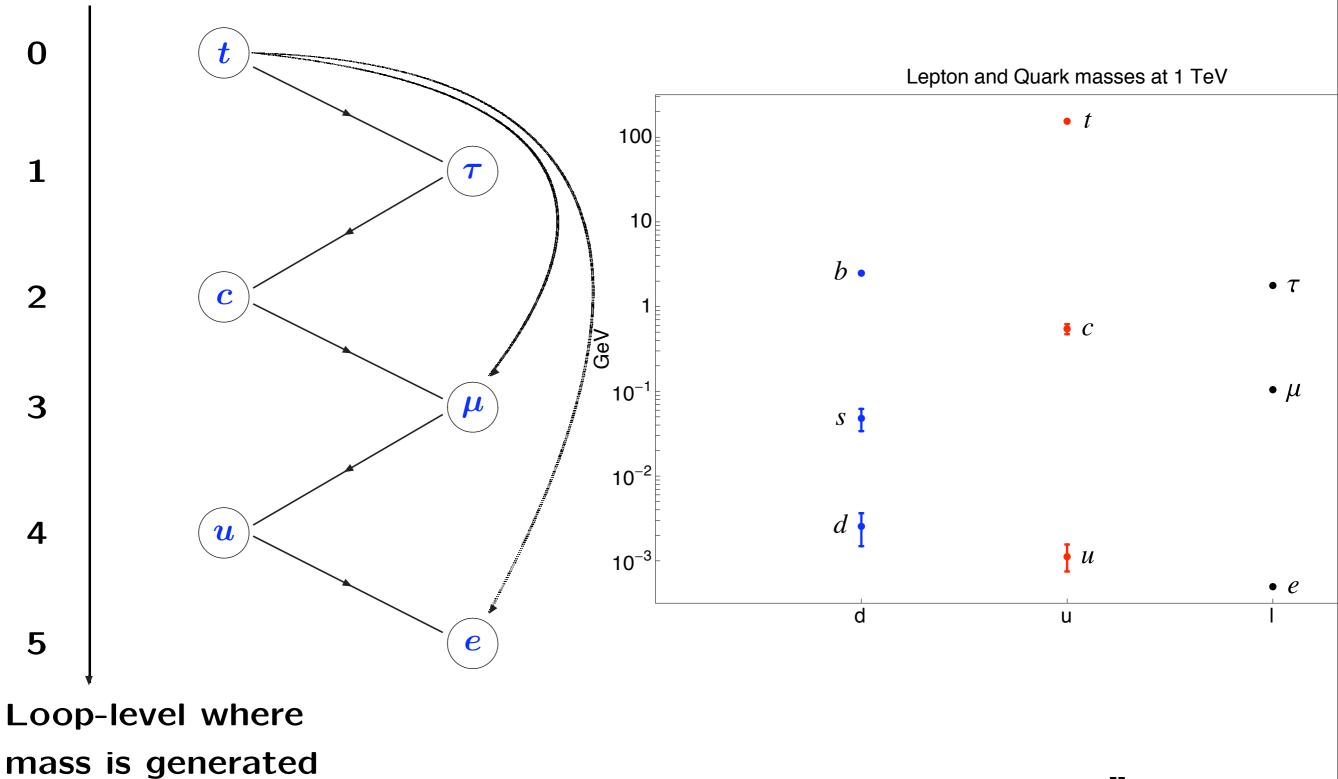
Loop-level where mass is generated





## More likely to fail...?

#### PJF and Dobrescu



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Top is clearly special

So,

## assume only the top has a tree level Yukawa

 $y_t H \bar{u}_R^3 Q_L^3$ 



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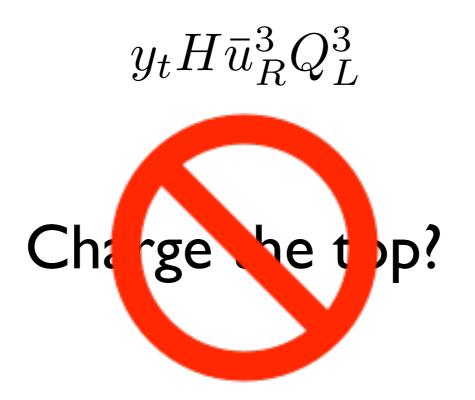
Charge the top?



Top is clearly special

So,

## assume only the top has a tree level Yukawa





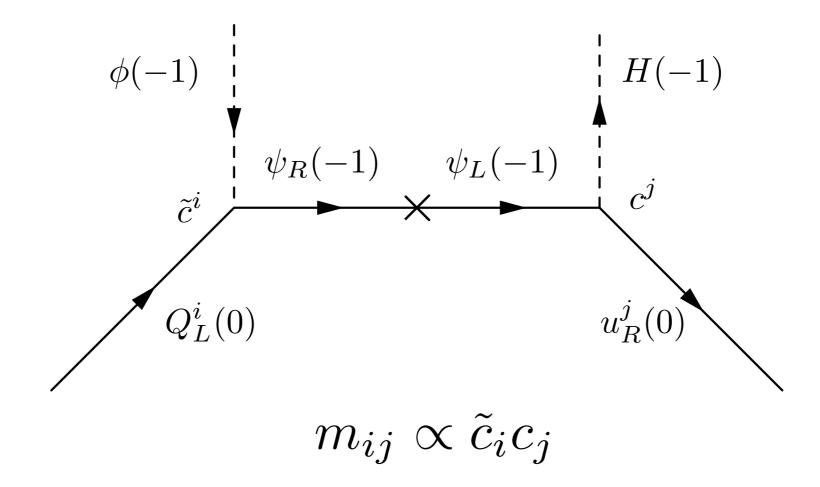
Instead charge Higgs under an extra U(1)

 $U(1)\,{\rm broken}$  by the vev of a SM singlet  $\phi\,$  of charge - I

Introduce a vector like pair of fermions with quantum numbers of left handed quarks, also charged under  $U(1)\,$ 



### Yukawas:



But Ih top and rh top only appear *linearly* in couplings Redefine couplings so only one Ih and one rh couple Call these the top

Mass matrix is rank I

Only the top gets a tree level mass



## Chiral symmetries

 $y_t \neq 0$  $U(3)_Q \times U(3)_u \times U(3)_d \to U(1)_t \times U(2)_Q \times U(2)_u \times U(3)_d$ 

Need to break remaining chiral symmetries

Introduce a scalar leptoquark

$$\tilde{r}: (3, 2, +7/6)$$

Most general interactions

$$\lambda_{ij} \tilde{r} \overline{u}_R^i L_L^j + \lambda'_{ij} \tilde{r} \overline{Q}_L^i e_R^j + \text{H.c.}$$



## Chiral symmetries

 $y_t \neq 0$  $U(3)_Q \times U(3)_u \times U(3)_d \to U(1)_t \times U(2)_Q \times U(2)_u \times U(3)_d$ 

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 $y_t \neq 0$ 

 $U(3)_Q \times U(3)_u \times U(3)_d \to U(1)_t \times U(2)_Q \times U(2)_u \times U(3)_d$ 

$$\stackrel{\lambda \neq 0}{\to} U(1)_u \times U(3)_d$$

$$U(3)_L \times U(3)_e^{\substack{\lambda \neq 0\\\lambda' \neq 0}} U(1)_L$$

With this breaking of chiral symmetries up type quarks and charged leptons can get a mass at some loop order



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 $U(3)_Q \times U(3)_u \times U(3)_d \to U(1)_t \times U(2)_Q \times U(2)_u \times U(3)_d$ 

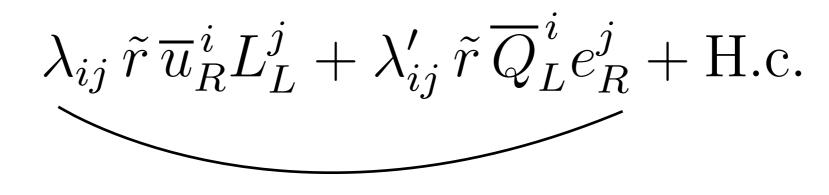
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## With this breaking of chiral symmetries up type quarks and charged leptons can get a mass at some loop order

But what loop order?



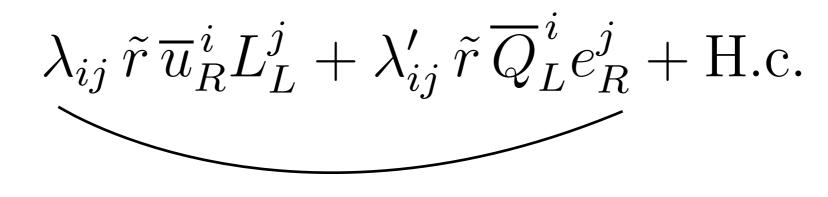


Linear couplings

Redefine fields:

$$\begin{pmatrix} \lambda_{11} & \lambda_{12} & \lambda_{13} \\ \lambda_{21} & \lambda_{22} & \lambda_{23} \\ \lambda_{31} & \lambda_{32} & \lambda_{33} \end{pmatrix}$$

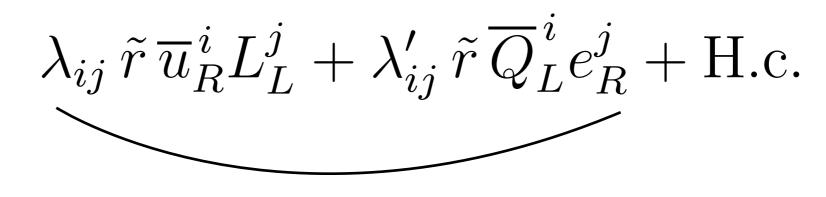




Linear couplings

Redefine fields:





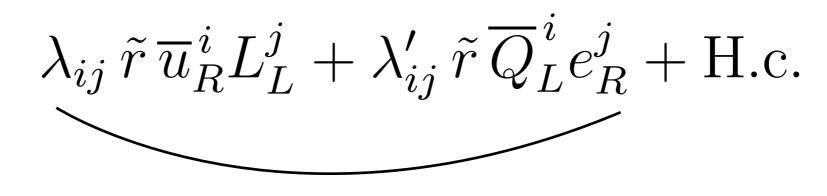
Linear couplings

Redefine fields:

•Define  $L_3$  so it only couples only to  $u_3$ 

$$\begin{pmatrix} \lambda_{11} & \lambda_{12} & \lambda_{13} \\ \lambda_{21} & \lambda_{22} & \lambda_{23} \\ 0 & 0 & \lambda_{33} \end{pmatrix}$$





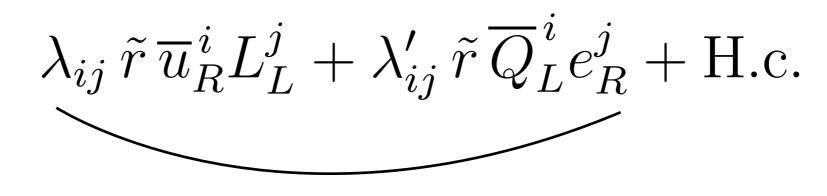
Linear couplings

Redefine fields:

- •Define  $L_3$  so it only couples only to  $u_3$
- $u_2$  couples only to  $L_2$  and  $L_3$

$$\begin{pmatrix} \lambda_{11} & \lambda_{12} & \lambda_{13} \\ 0 & \lambda_{22} & \lambda_{23} \\ 0 & 0 & \lambda_{33} \end{pmatrix}$$





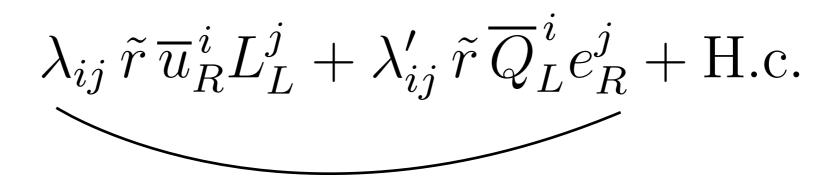
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- •Define  $L_3$  so it only couples only to  $u_3$
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- •Rotation of  $u_1$  and  $u_2$

$$\begin{pmatrix} \lambda_{11} & \lambda_{12} & 0 \\ 0 & \lambda_{22} & \lambda_{23} \\ 0 & 0 & \lambda_{33} \end{pmatrix}$$





Linear couplings

Redefine fields:

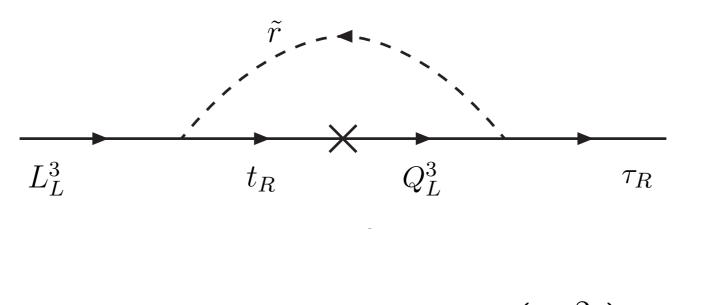
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$$\begin{pmatrix} \lambda_{11} & \lambda_{12} & 0 \\ 0 & \lambda_{22} & \lambda_{23} \\ 0 & 0 & \lambda_{33} \end{pmatrix}$$

 $\lambda_{ij}, \lambda'_{ij}$ can be made real and positive

#### **Fermilab**

## One loop tau mass

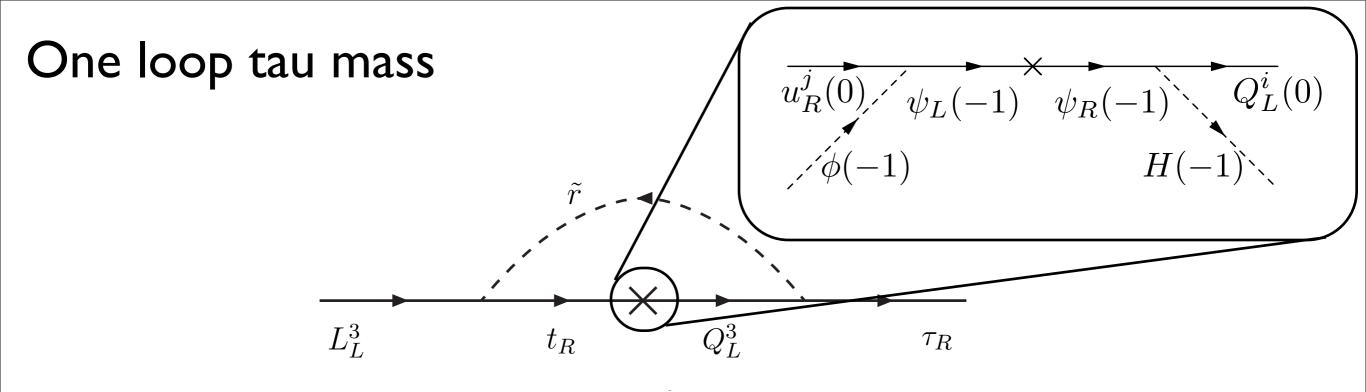


$$m_{\tau} \simeq \lambda_{33} \lambda_{33}' m_t \frac{N_c}{16\pi^2} \ln\left(\frac{\Lambda^2}{M_{\tilde{r}}^2}\right)$$

pprox 0.09 for  $\Lambda pprox 10 M_{ ilde{r}}$ 

 $\lambda_{33}\lambda'_{33} \approx (0.36)^2$  for correct  $m_{\tau}/m_t$  ratio





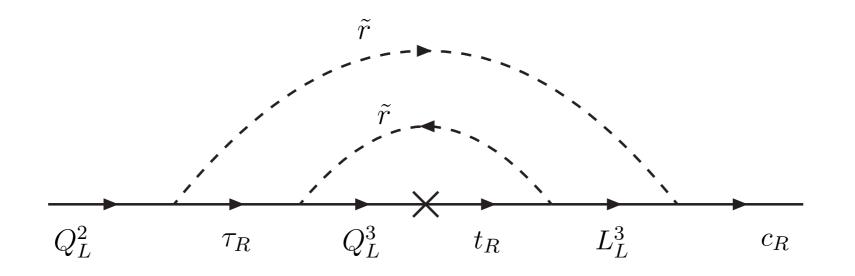
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## Two loop charm mass - a "rainbow" diagram



$$M_u[\tilde{r}\tilde{r}] = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \lambda'_{23}\lambda_{23} & \lambda'_{33}\lambda_{23} \\ 0 & \lambda'_{23}\lambda_{33} & \lambda'_{33}\lambda_{33} \end{pmatrix} \lambda'_{33}\lambda_{33} m_t \epsilon_{\tilde{r}}^{(2)}$$

$$m_c = \lambda'_{23} \lambda_{23} m_\tau \frac{1}{16\pi^2} \log \frac{\Lambda^2}{M_{\tilde{r}}^2}$$

 $\lambda_{23}\lambda_{23}' \approx (3.3)^2$  for correct  $m_c/m_{ au}$  ratio

#### **‡** Fermilab

## Two loop charm mass - a "rainbow" diagram

$$\widetilde{r}_{R}$$

$$\widetilde{One \ loop}_{L_{2}}$$

$$\widetilde{r}_{R}$$

$$\widetilde{l}_{L}^{3}$$

$$\widetilde{r}_{R}$$

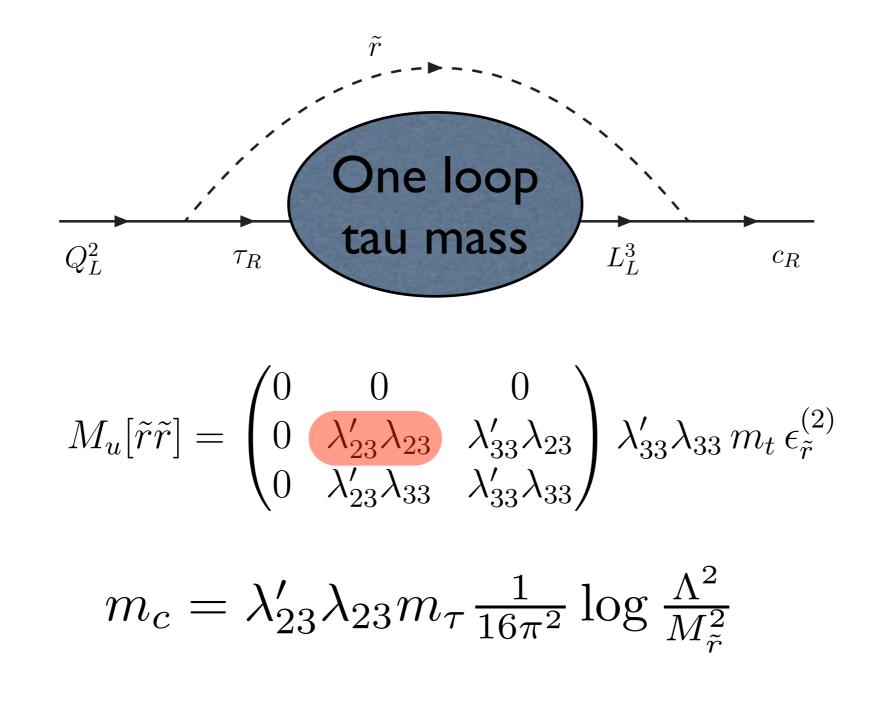
$$M_{u}[\widetilde{r}\widetilde{r}] = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \lambda'_{23}\lambda_{23} & \lambda'_{33}\lambda_{23} \\ 0 & \lambda'_{23}\lambda_{33} & \lambda'_{33}\lambda_{33} \end{pmatrix} \lambda'_{33}\lambda_{33} m_{t} \epsilon_{\widetilde{r}}^{(2)}$$

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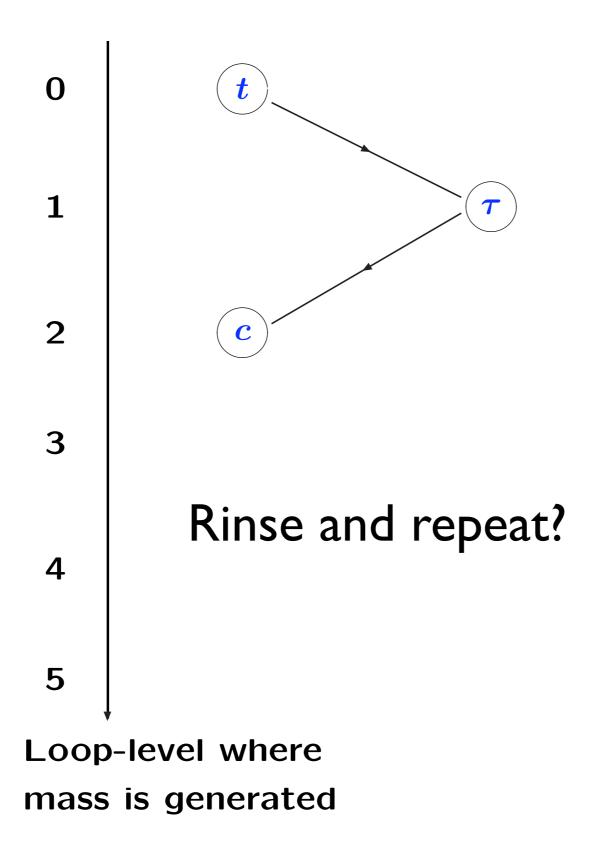
#### **Fermilab**

#### Two loop charm mass - a "rainbow" diagram

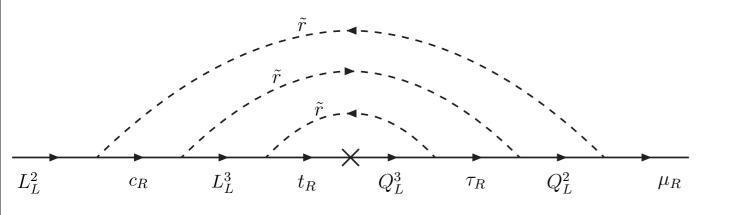


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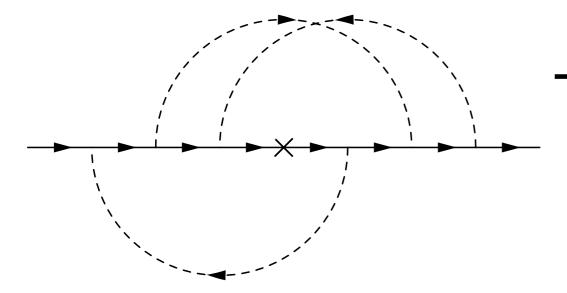






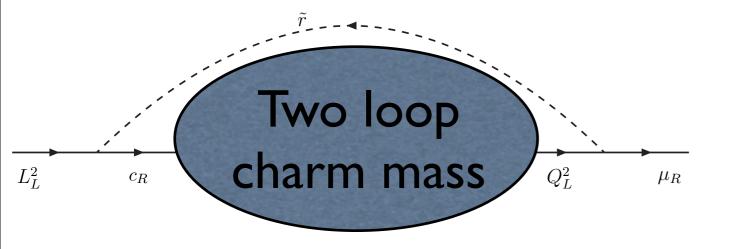






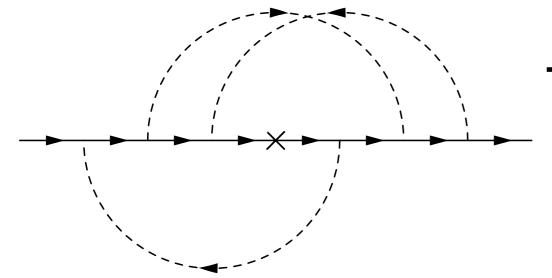
The diagram with no name  $~\sim N_C$ 





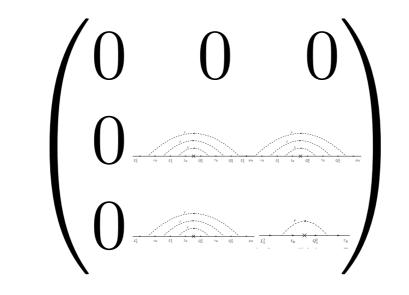






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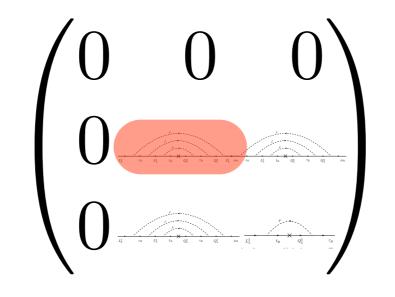




$$m_{\mu} \approx \lambda_{22}^{\prime} \lambda_{22} m_c (1+x) \frac{N_c}{16\pi^2} \log \frac{\Lambda^2}{M_{\tilde{r}}^2}$$

 $\lambda_{22}\lambda'_{22}(1+x) \approx (1.5)^2$  for correct  $m_{\mu}/m_c$  ratio



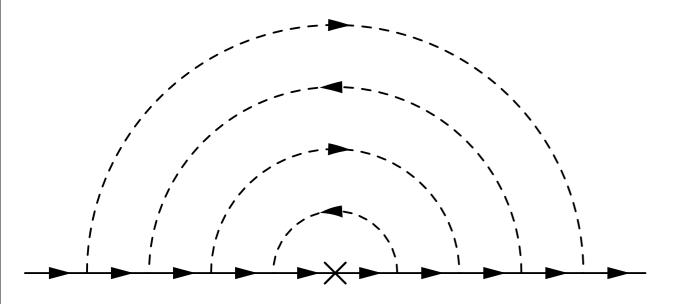


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#### Four loop up quark mass

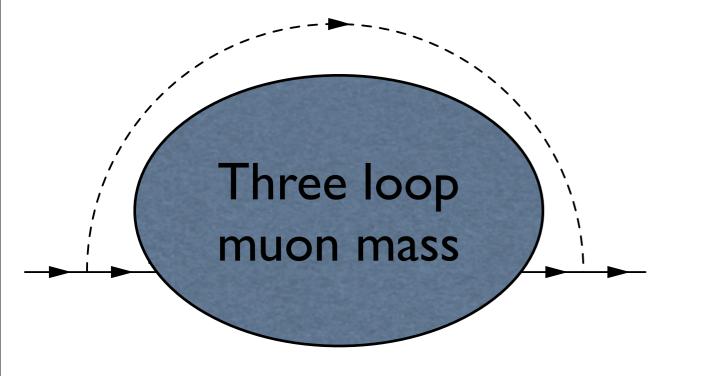


+4 other diagrams

Muon mass implies:  $\#\lambda_{12}\lambda'_{12} \approx (0.6)^2$ 



#### Four loop up quark mass



+4 other diagrams

Muon mass implies:  $\#\lambda_{12}\lambda'_{12} \approx (0.6)^2$ 



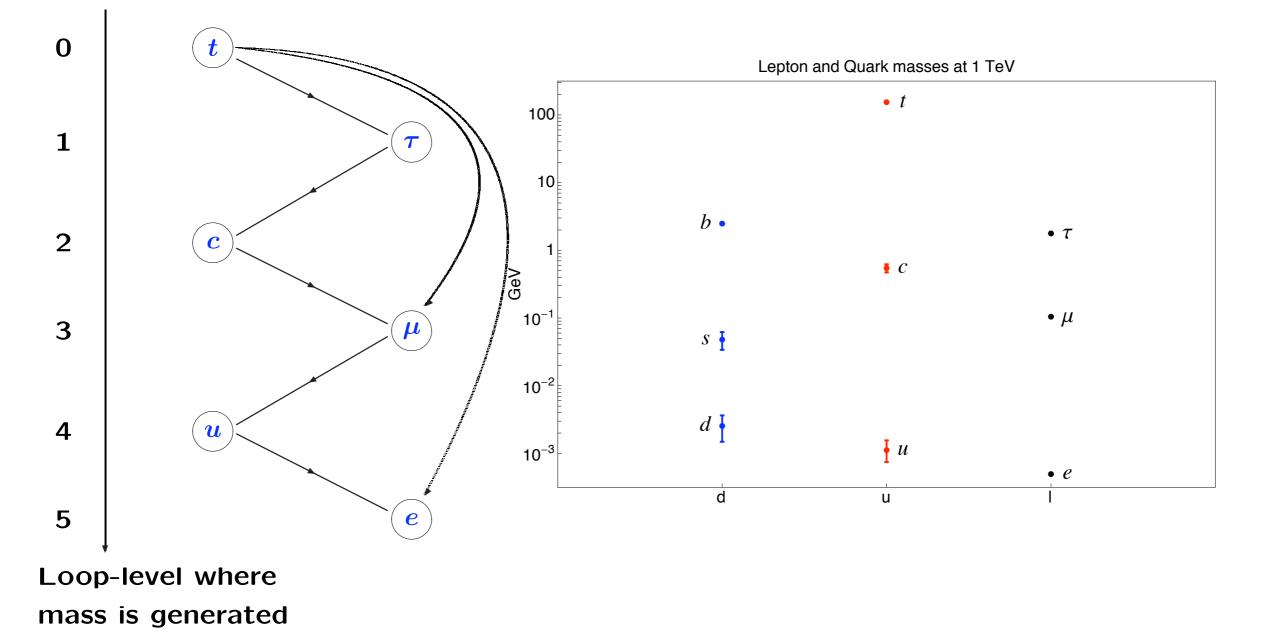
#### Five loop electron mass

If only source of electron mass will determine  $\lambda_{11}\lambda'_{11}$ 

Only input:

$$\tilde{r}: (3, 2, +7/6)$$

$$\lambda_{ij} \,\tilde{r} \,\overline{u}_R^i L_L^j + \lambda_{ij}' \,\tilde{r} \,\overline{Q}_L^i e_R^j + \text{H.c.}$$



#### Down quark masses

#### Need to break the remaining chiral symmetries

$$U(3)_d \times U(1)_u \times U(1)_L$$

#### Have choices diquarks, leptoquarks...



#### New field content

	$\phi$	$\psi_L,\psi_R$	Н	r	r'	$H_8$	$H'_8$	$\Phi_3$
SU(3)	1	3	1	3	3	8	8	$\overline{3}$
SU(2)	1	2	2	2	2	2	2	2
$U(1)_Y$	0	1/6	1/2	7/6	7/6	1/2	-1/2	-1/6
U(1)'	-1	-1	1	0	2	1	1	0

Up quarks and leptons Down quarks



Most general couplings

$$\kappa_i \Phi_8 \,\overline{u}_R^i \Psi_L + \kappa' \Phi_8' \,\overline{d}_R^3 \Psi_L$$

$$\eta_{ij} \Phi_3 \overline{d}_R^i L_L^j + \text{h.c.}$$

break the remaining chiral symmetries

$$U(3)_d \times U(1)_u \times U(1)_L \to U(1)_L \times U(1)_Q$$



Most general couplings

$$\kappa_i \, \Phi_8 \, \overline{u}_R^i \Psi_L + \kappa' \, \Phi_8' \, \overline{d}_R^3 \Psi_L$$
Only couples to b

$$\eta_{ij} \Phi_3 \overline{d}_R^i L_L^j + \text{h.c.}$$

break the remaining chiral symmetries

$$U(3)_d \times U(1)_u \times U(1)_L \to U(1)_L \times U(1)_Q$$



# Without altering up type and leptons have the freedom to rotate such that,

$$\eta = \begin{pmatrix} \eta_{11} & \eta_{12} & 0 \\ \eta_{21} & \eta_{22} & \eta_{23} \\ \eta_{31} & \eta_{32} & \eta_{33} \end{pmatrix}$$

$$\kappa = (\kappa_1, \kappa_2, \kappa_3)$$



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Diagonal entries can be made real and positive

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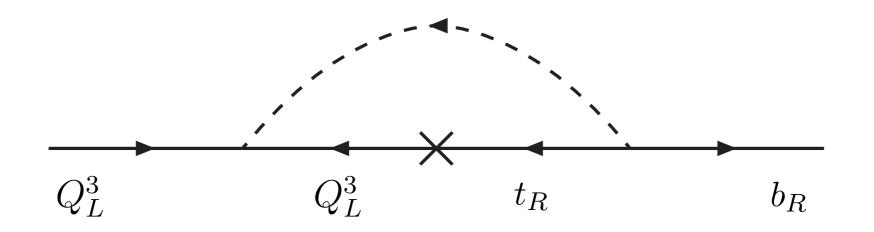
Diagonal entries can be made real and positive

$$\kappa = (\kappa_1, \kappa_2, \kappa_3)$$

#### Entries can be made real and positive



#### One loop bottom mass



$$m_b \approx m_t \kappa' \frac{N_C}{16\pi^2} \left(\frac{\langle \phi \rangle}{M_8}\right)^2 \log \frac{\Lambda^2}{M_8^2}$$

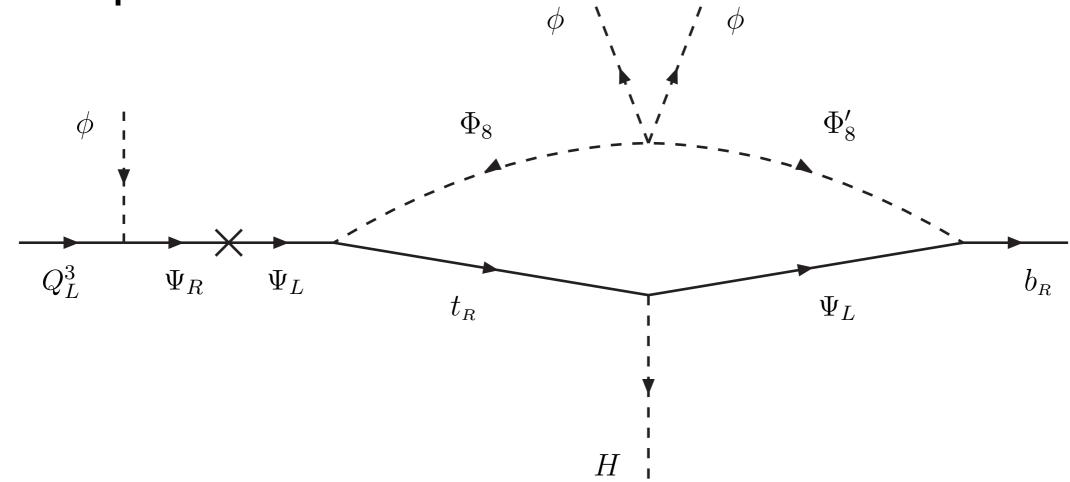


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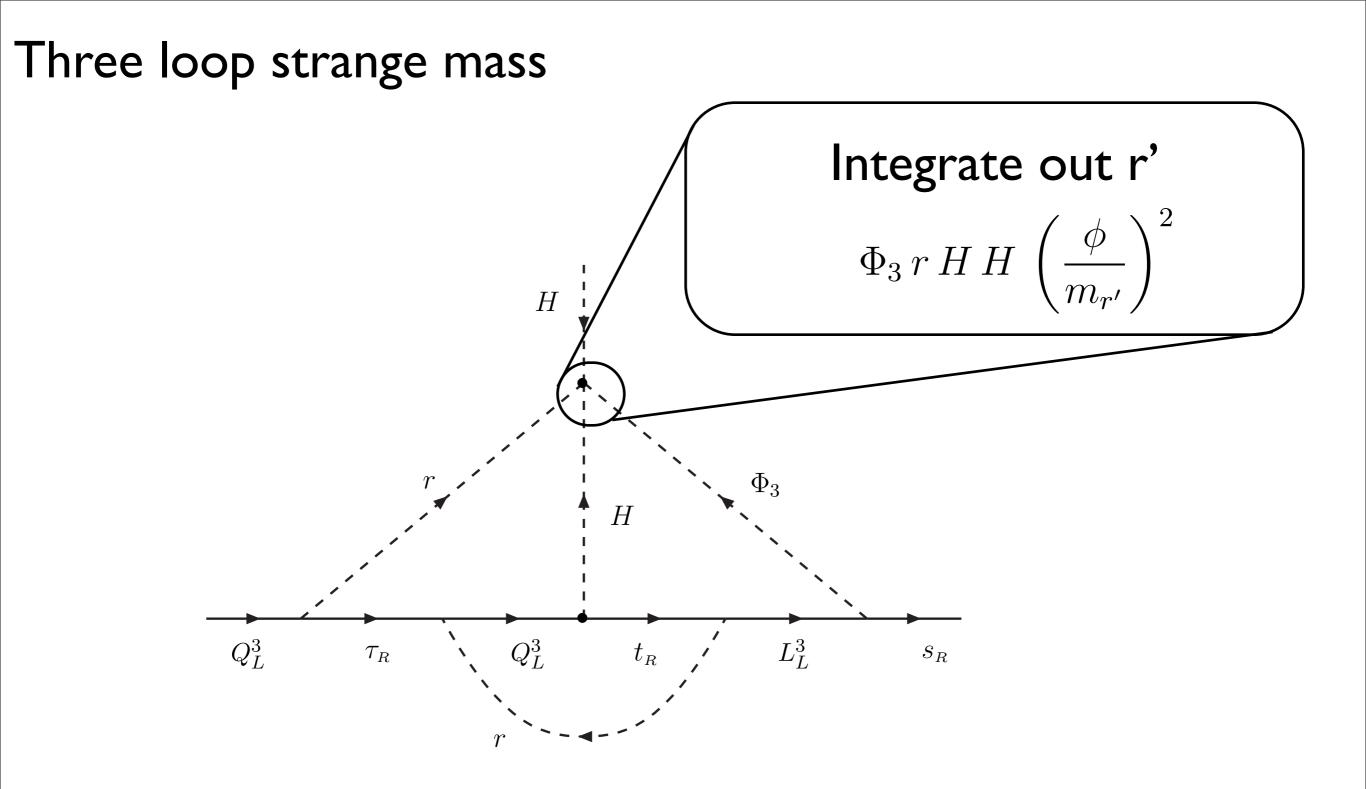


#### One loop bottom mass



$$m_b \approx m_t \kappa' \frac{N_C}{16\pi^2} \left(\frac{\langle \phi \rangle}{M_8}\right)^2 \log \frac{\Lambda^2}{M_8^2}$$







#### Four loop down masses

# The down has a 4 loop mixed diagram (exercise for reader)



"Cross Talk"

There are also corrections to some of the states that have mass:

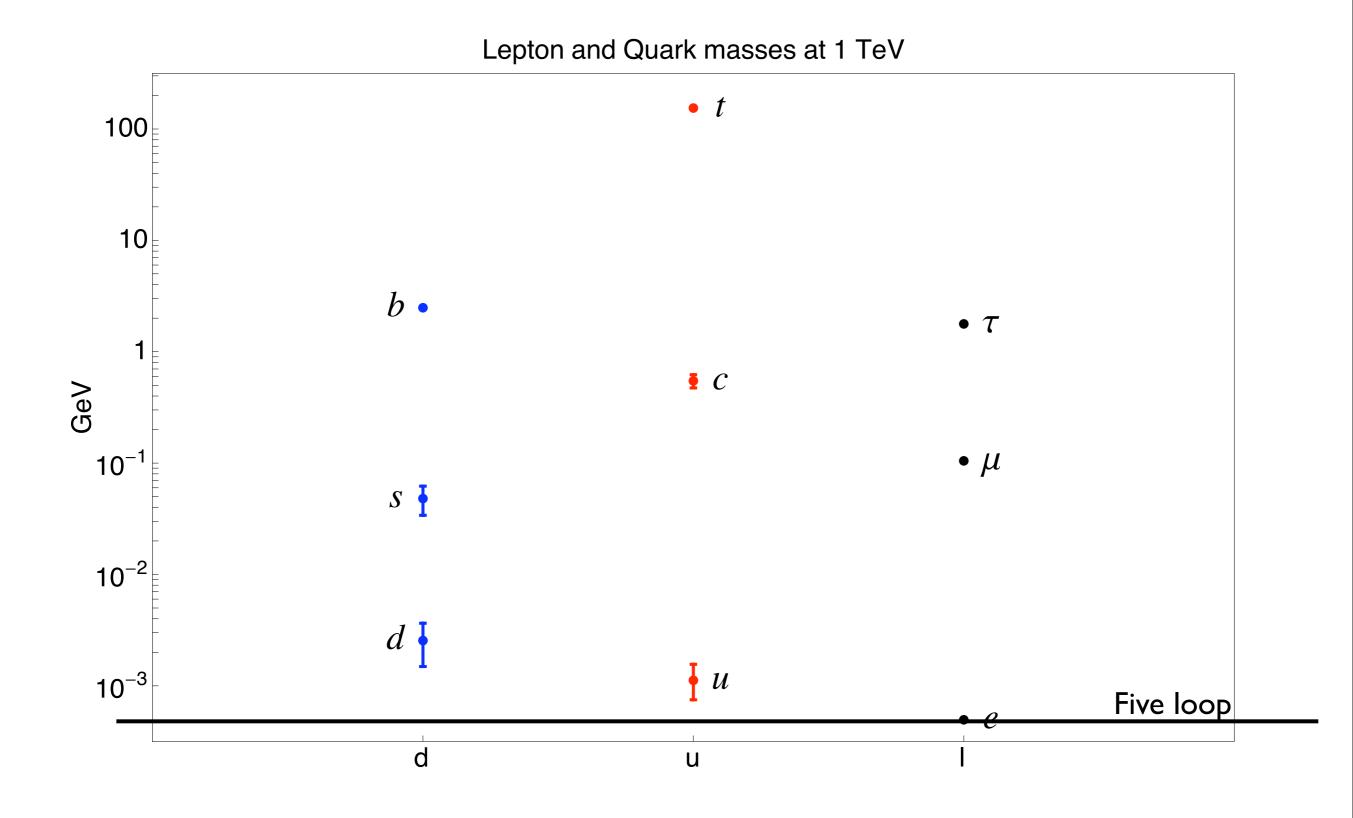
Charm gets a two loop correction

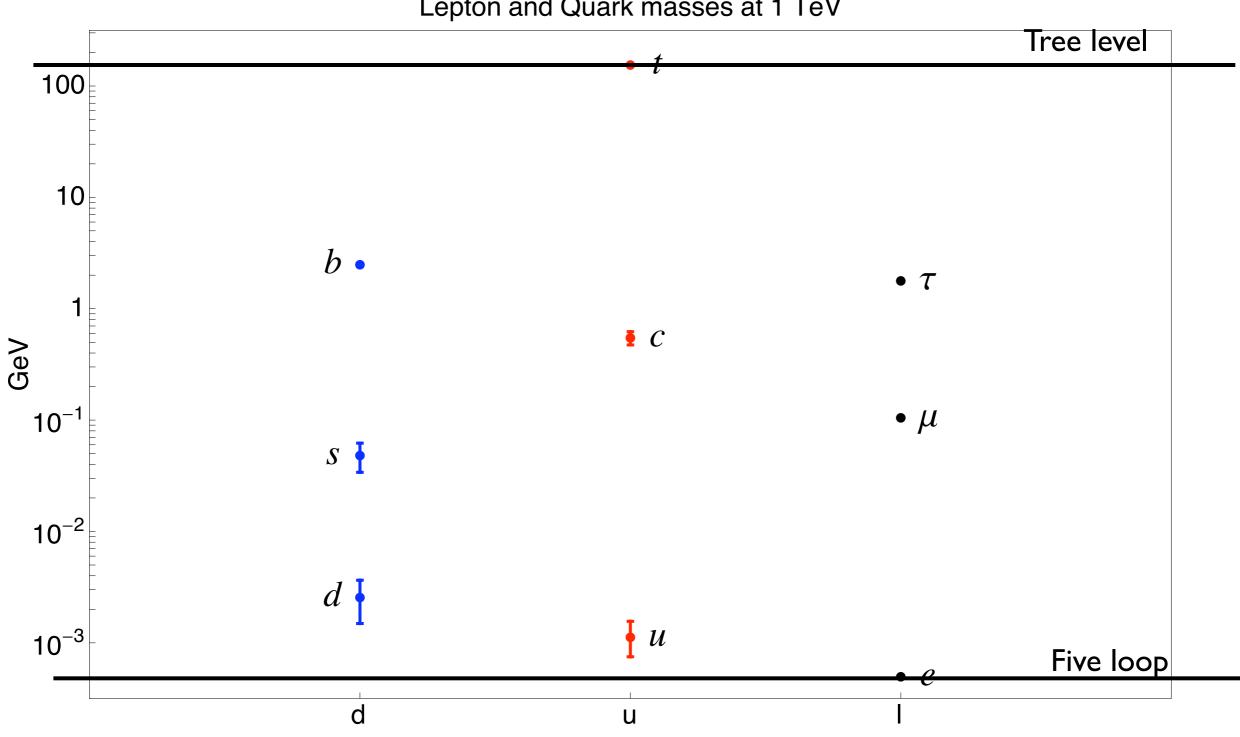
Up gets a four loop correction

Muon gets a three loop correction

Electron gets a five loop correction

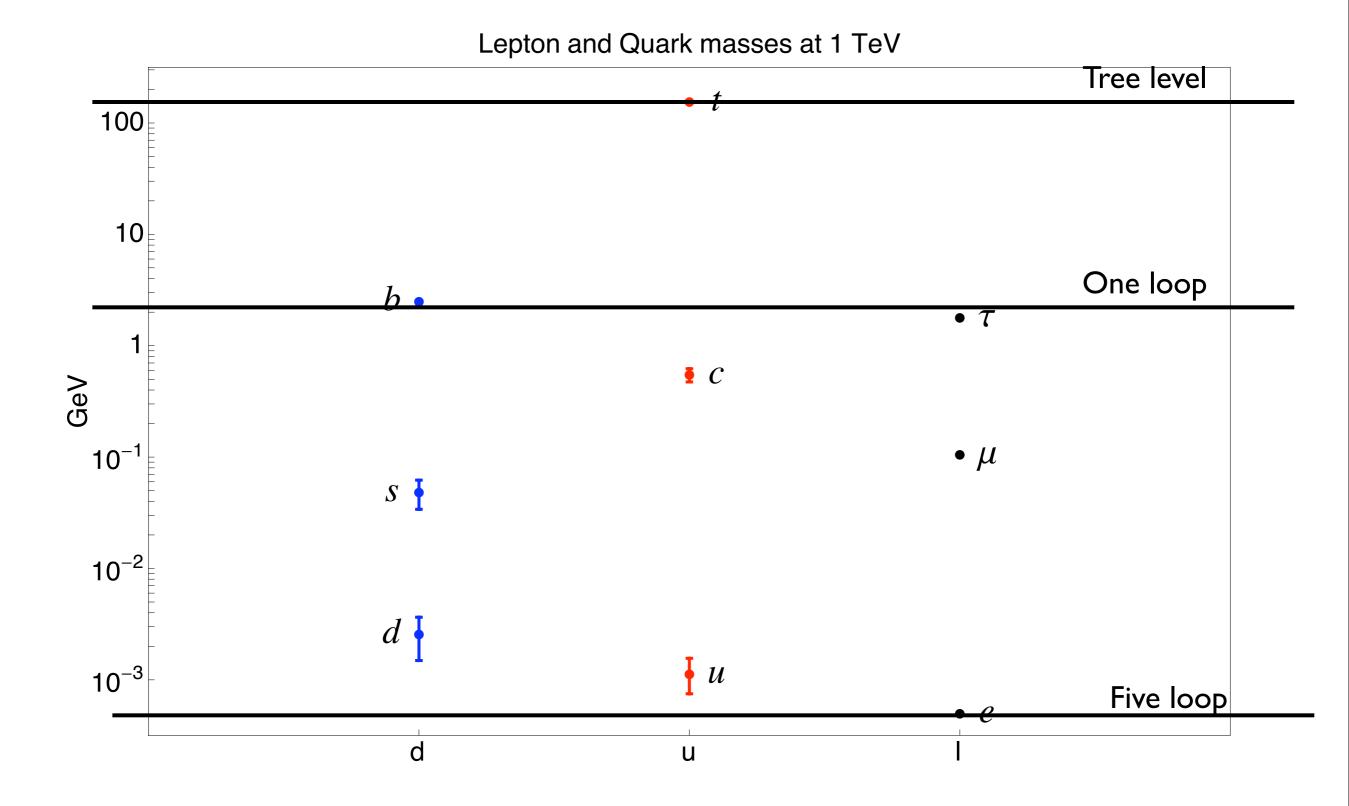




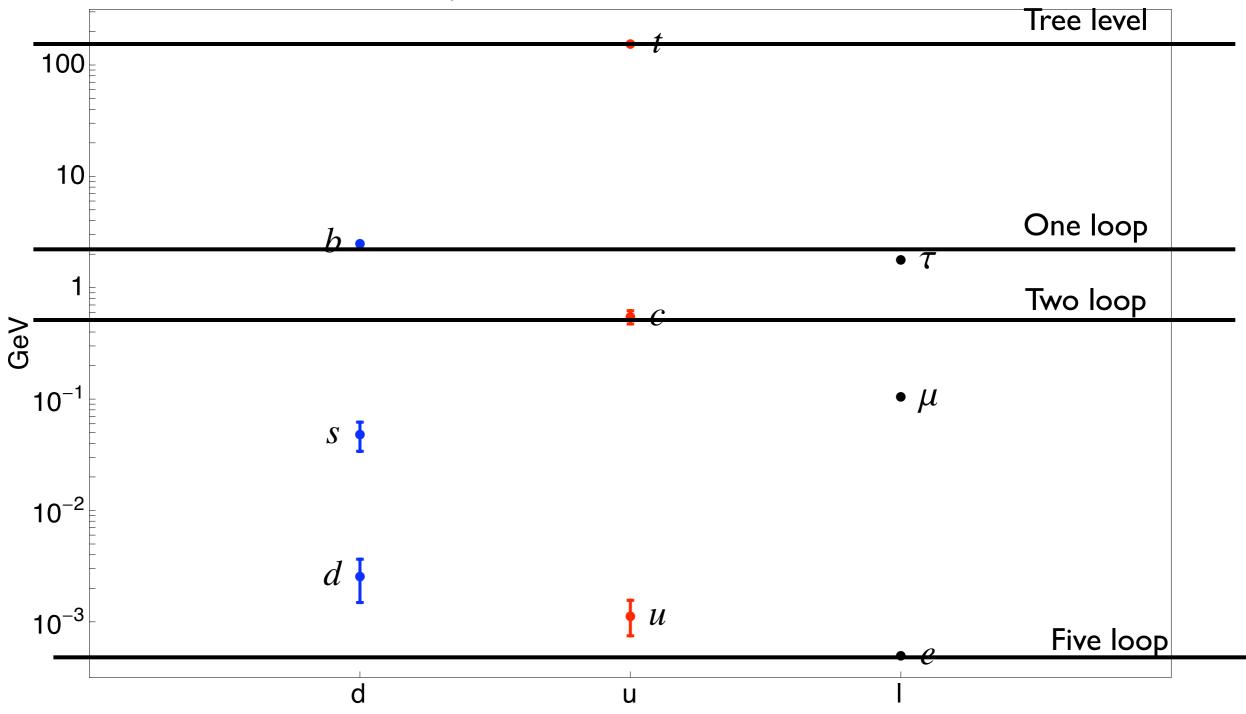


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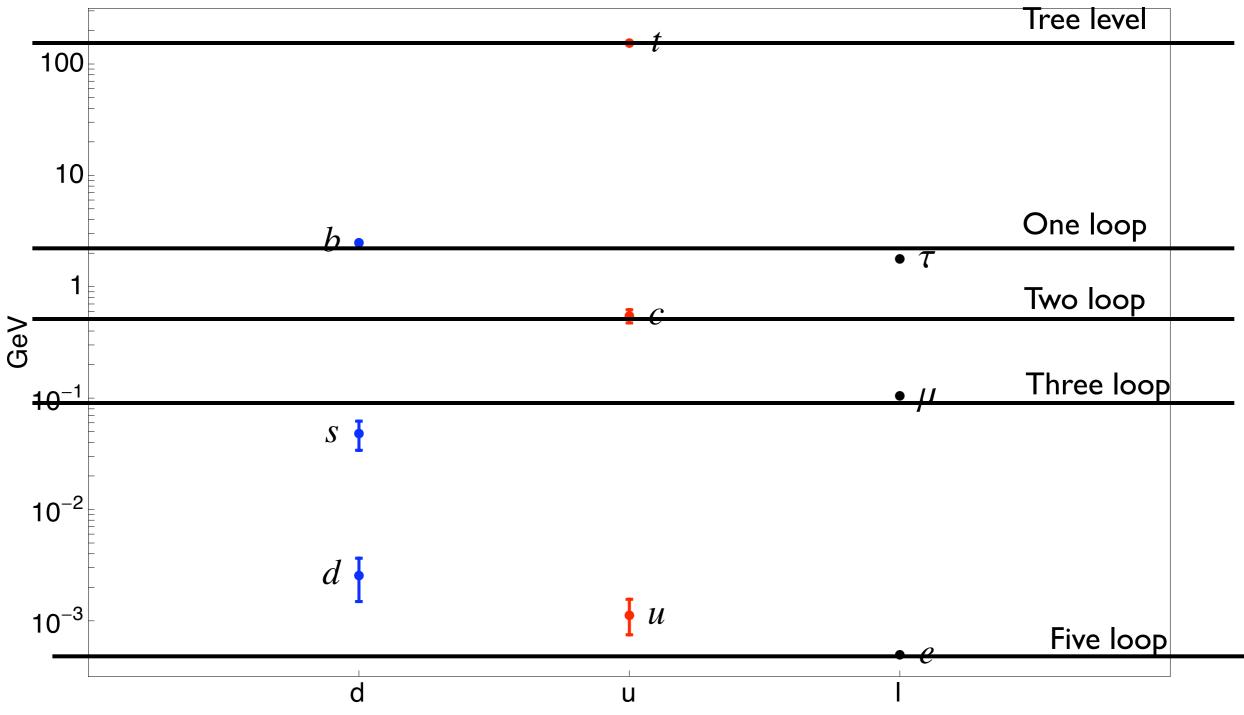
Lepton and Quark masses at 1 TeV



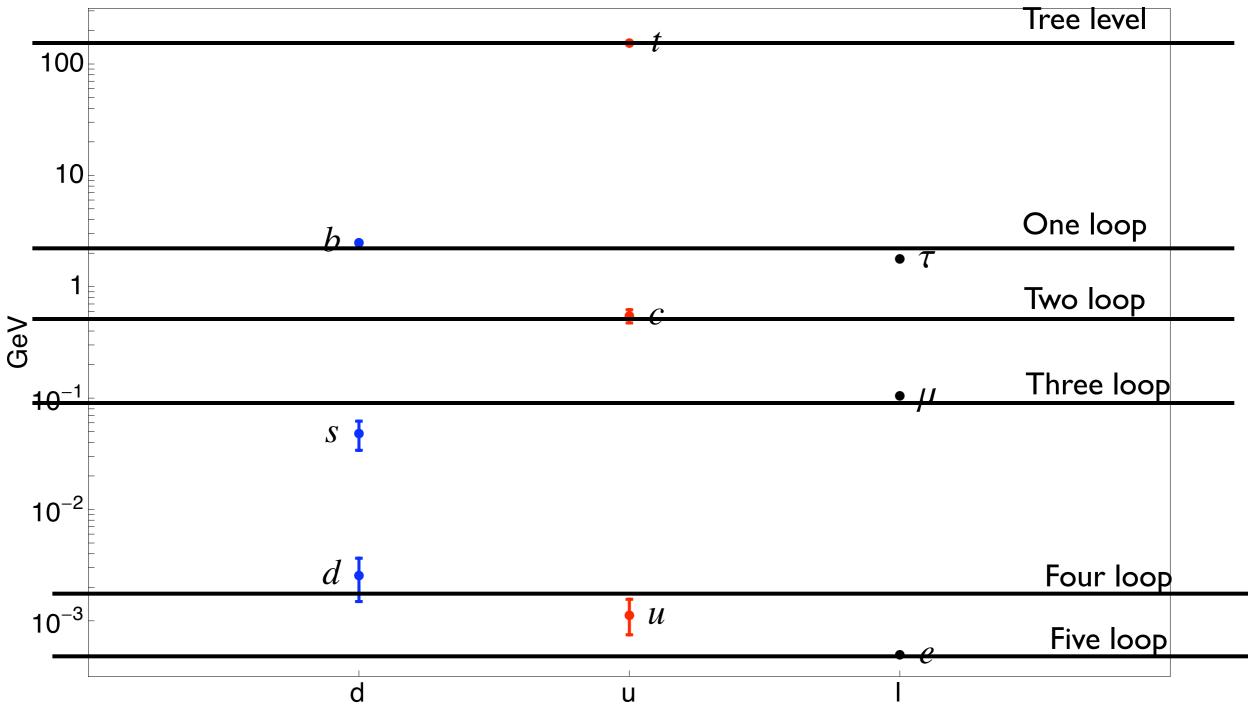
Lepton and Quark masses at 1 TeV



Lepton and Quark masses at 1 TeV



Lepton and Quark masses at 1 TeV





#### CKM

$$m_u \approx m_t \begin{pmatrix} \epsilon^4 & \epsilon^2 & \epsilon^2 \\ \epsilon^4 & \epsilon^2 & \epsilon^2 \\ \epsilon^4 & \epsilon^2 & 1 \end{pmatrix} \qquad m_d \approx m_t \begin{pmatrix} \epsilon^4 & \epsilon^4 & \epsilon^4 \\ \epsilon^4 & \epsilon^3 & \epsilon^3 \\ \epsilon^4 & \epsilon^3 & \epsilon \end{pmatrix}$$

#### Resulting in

$$V_{CKM} \approx \begin{pmatrix} 1 - \epsilon^2 & \epsilon & \epsilon^3 \\ -\epsilon & 1 - \epsilon^2 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix}$$

Still to think about phases...



James Wells:

### Q:What sucks in the Standard Model??

## A:The Higgs



James Wells:

### Q:What sucks in the Standard Model??

## A:The Higgs

# Q:Does the solution predict LHC physics?



#### The model contains extra fermions and scalar Leptoquarks

(Alternative realisation contains diquarks - easier to see at LHC than TeVatron)



#### Constraints

Tree level exchange of leptoquark can lead to flavour changing processes e.g.

$$K^+ \to \pi^0 \mu^+ \mu^-$$

$$\tau^+ \to K^0 e^+$$

$$\pi^+ \to e^+ \nu \text{ versus } \pi^+ \to \mu^+ \nu$$

 $\mu \rightarrow e$  conversion

$$M_{\tilde{r}} \gtrsim 5 - 50 \text{ TeV}$$



#### Conclusions

- •Fermions have complicated mass hierarchy
- •Many attempts exist to explain it
- •Top is probably special, perhaps only top mass has a tree level Yukawa
- •With extra scalars coupling to fermions top mass is communicated at loop level
- •Interesting structure of fermion mass spectrum arises
- •Predicts flavour changing processes



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- •Interesting structure of fermion mass spectrum arises
- •Predicts flavour changing processes
- •Project X?



