
KITP – LoopFest III — 3rd April 2004

*Automated resummation of QCD
final state observables*

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☞ *In collaboration with*

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Fermilab

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- ✗ The most famous example: the Thrust

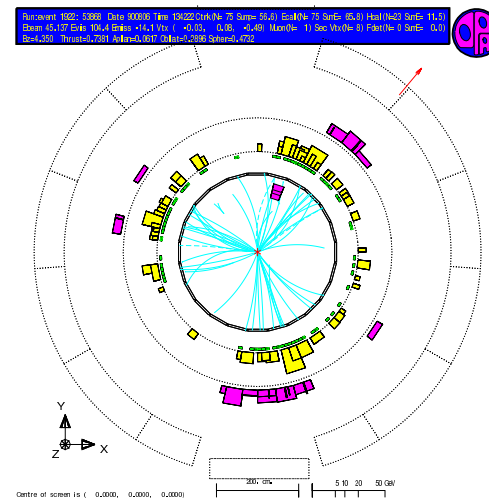
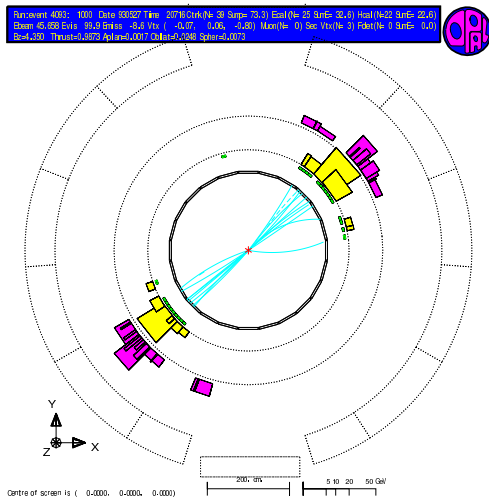
$$T \equiv \frac{1}{Q} \max_{\vec{n}_T} \sum_i |\vec{p}_i \cdot \vec{n}_T| = \frac{1}{Q} \sum_i |p_{iz}|$$

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Pencil-like event: $\tau \equiv 1 - T \ll 1$

Planar event: $T \simeq 2/3$



Perturbative QCD ingredients

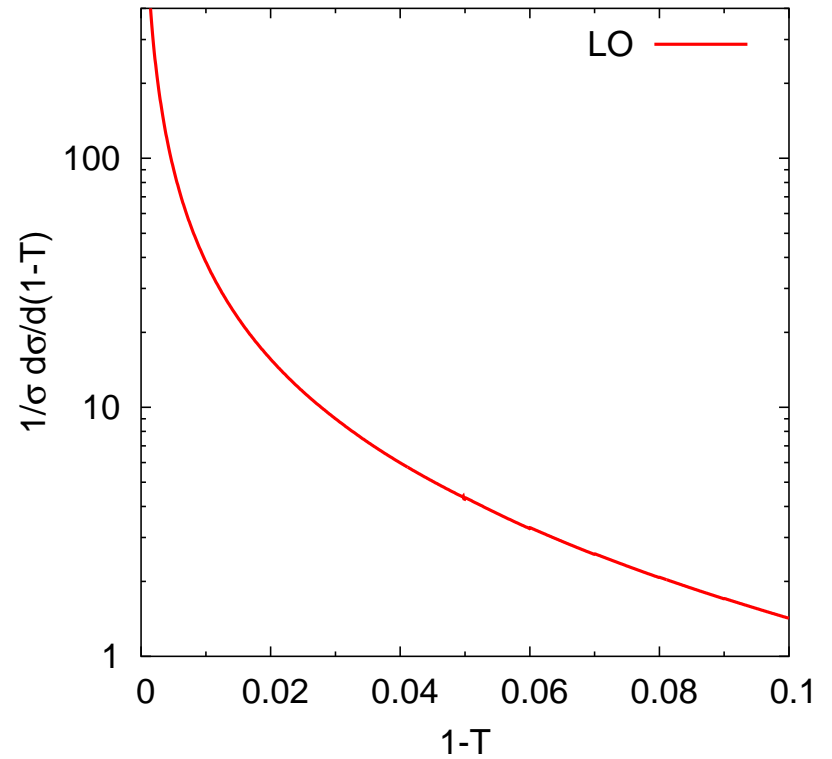
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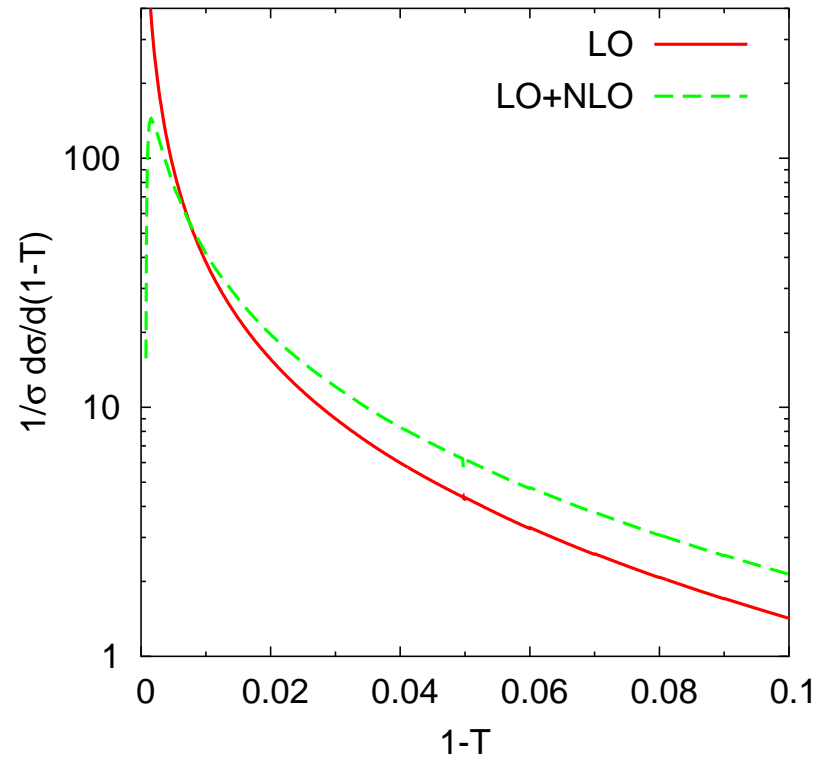
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- Usually only done numerically

[Event2, Disent, NLOJET++...]

LO, NLO, ... all *diverge* in two-jet region
($1 - T \rightarrow 0$)



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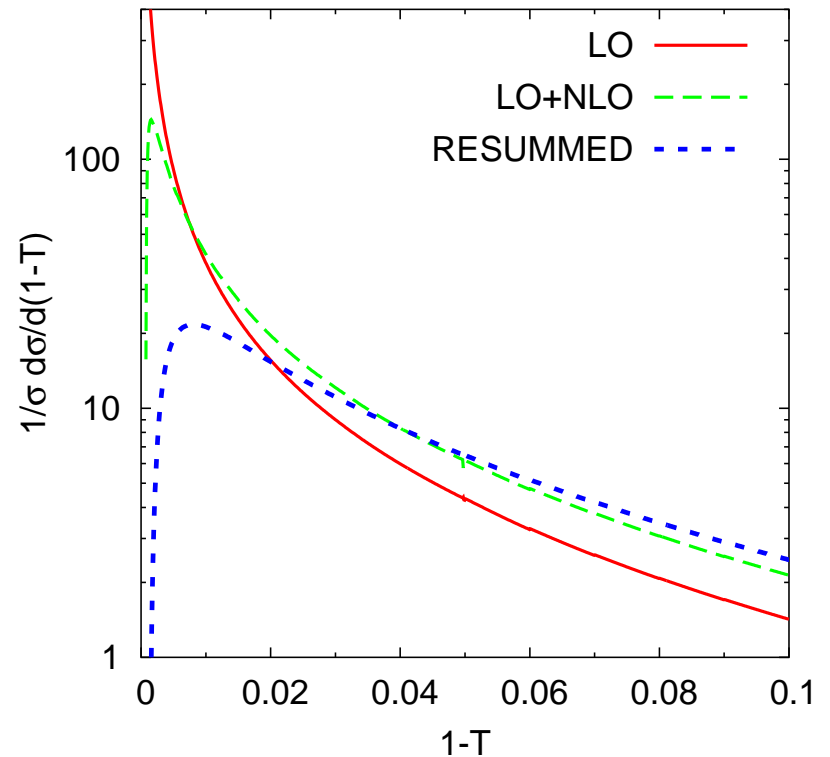
[Event2, Disent, NLOJET++...]

LO, NLO, ... all *diverge* in two-jet region
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Subject of this seminar is

FINAL-STATE RESUMMATION

i. e. *all-orders* description of the “exclusive” 2-jet limit.



Jet observables are a **good compromise** between

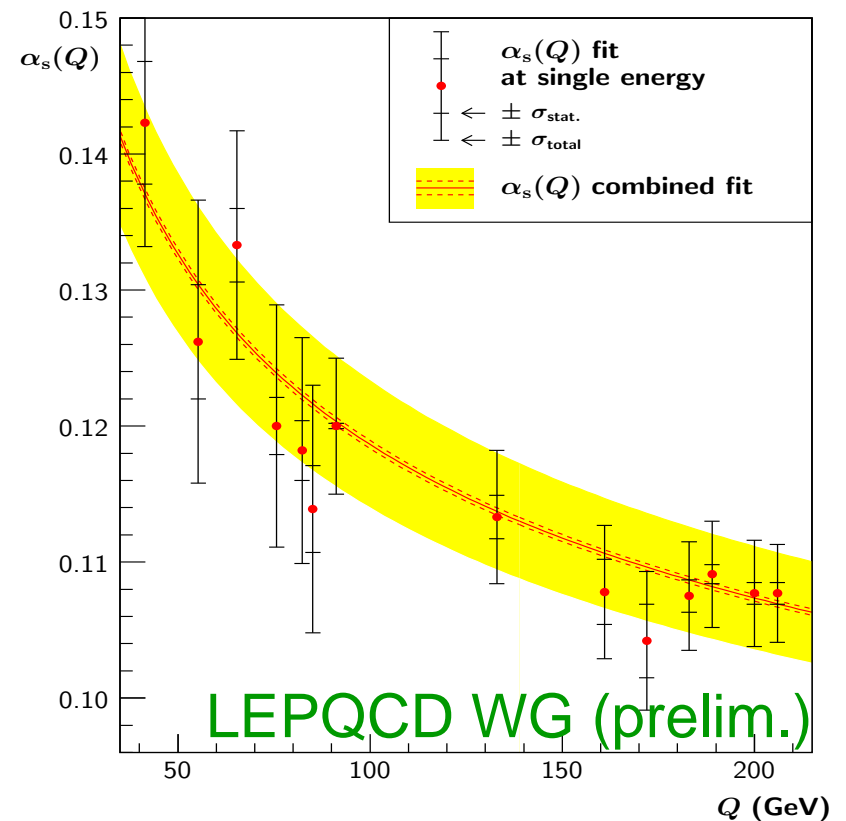
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- **sensitivity** to properties of QCD radiation

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- Measurements of the coupling α_s and its **renormalization group running**

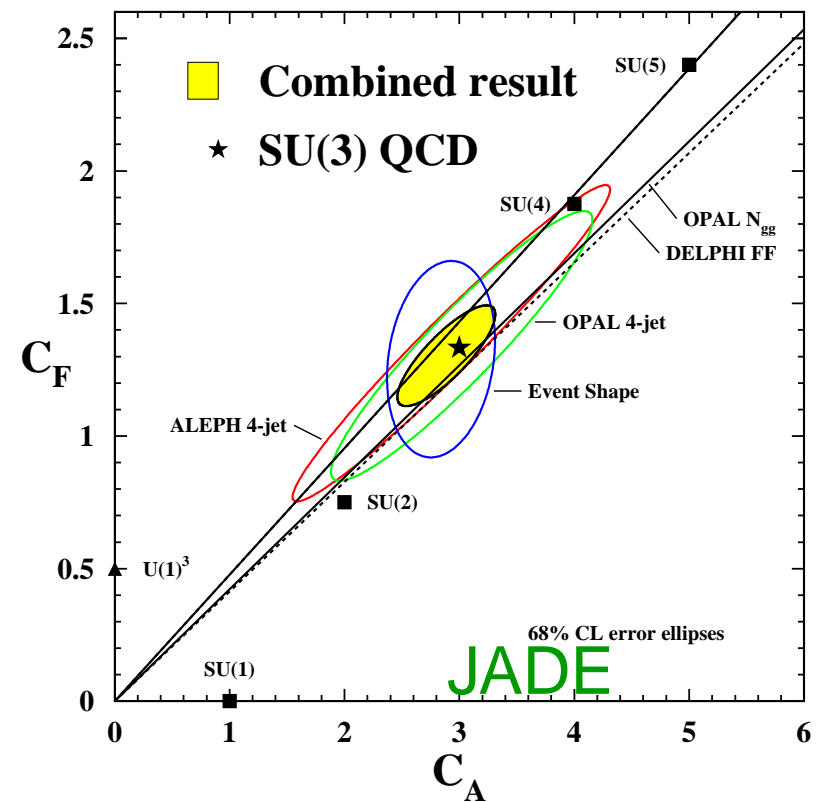


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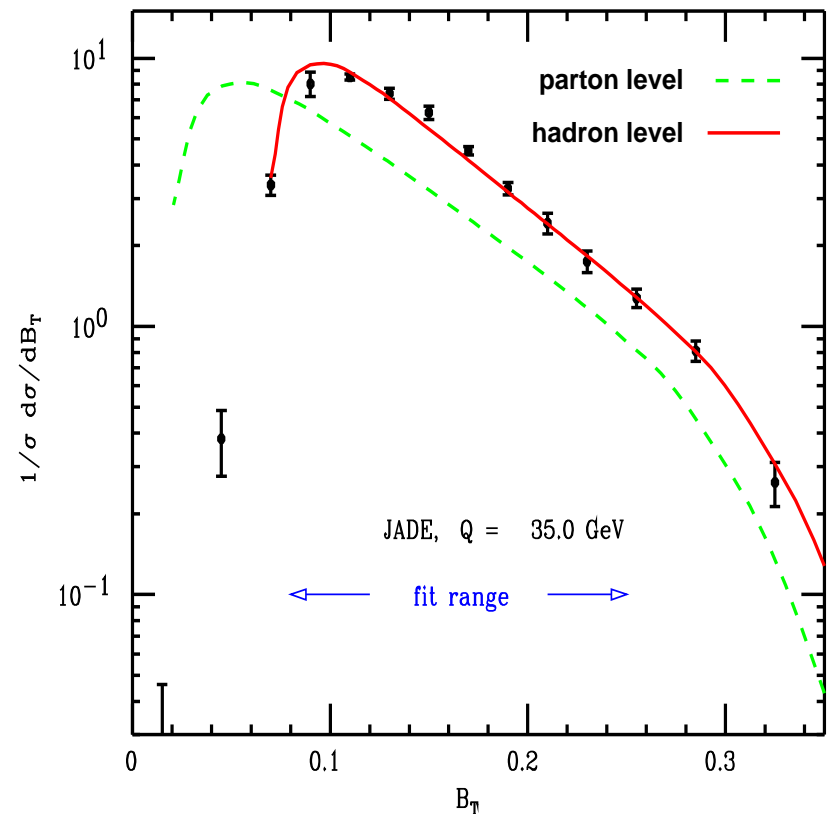


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Provide a wealth of information, e.g.:

- Measurements of the coupling α_s and its **renormalization group running**
- Measurements/cross checks of the values of the **colour factors** of QCD
- Studies of connection between **parton-level** (perturbative description of quarks and gluons) and **hadron-level** (the real)



Large Logarithms to all orders

Probability of “constrained” events, i. e. $V(k_1 \dots k_n) < v$, has a **divergent PT expansion**

$$\Sigma(v) \equiv \text{Prob}(V < v) = 1 + \sum_{m \leq 2n} R_{n,m} \alpha_s^n \text{Log}^m v + \dots$$

i. e. there is a soft & collinear divergence [\rightsquigarrow Log] for each emitted gluon

Today’s state-of-the art accuracy

- accounts for all **Leading (LL) and Next-to-Leading Logs (NLL)**

$$\Sigma(v) = \exp\left\{ \underbrace{Lg_1(\alpha_s L)}_{LL} + \underbrace{g_2(\alpha_s L)}_{NLL} + \dots \right\}$$

👉 NB:

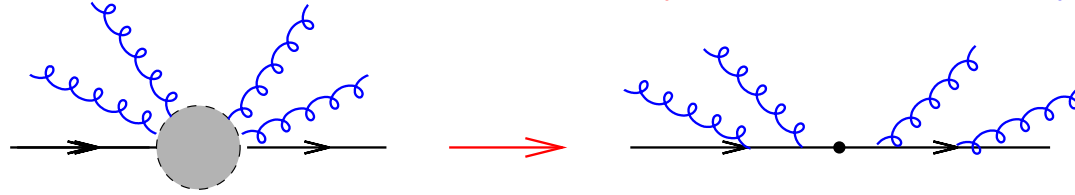
- LL means $\alpha_s^n L^{n+1}$ in $\ln \Sigma$, not just $\alpha_s^n L^{2n}$ in Σ
 - NLL means $(\alpha_s L)^n$ in $\ln \Sigma$, not just $\alpha_s^n L^{2n-1}$ in Σ
- furthermore resummed results are **matched to Fixed Order at NLO**

Basics of resummation: factorization

First half of the history: **Matrix elements and phase space**

exploit *angular ordering* \Rightarrow soft *independent emissions* (\Rightarrow QED)

e.g. $e^+e^- \rightarrow 2 \text{ jets} \Rightarrow w_{p\bar{p}}(k_1, \dots, k_n) = \frac{1}{n!} \prod_{i=1}^n w_{p\bar{p}}(k_i) \sim \frac{1}{n!} \prod_{i=1}^n \frac{\alpha_s C_F}{\pi} \frac{dE}{E} \frac{d\theta}{\theta}$

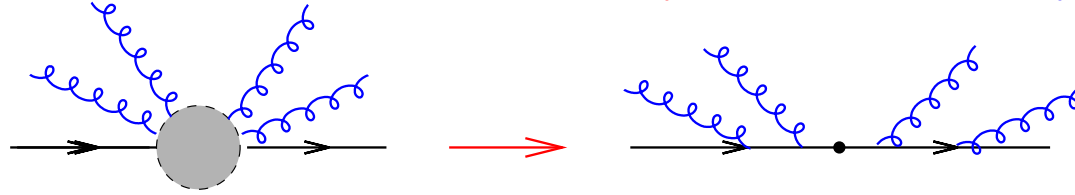


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Second half of the history: **The observable definition**

analyze the observable & use Mellin transforms

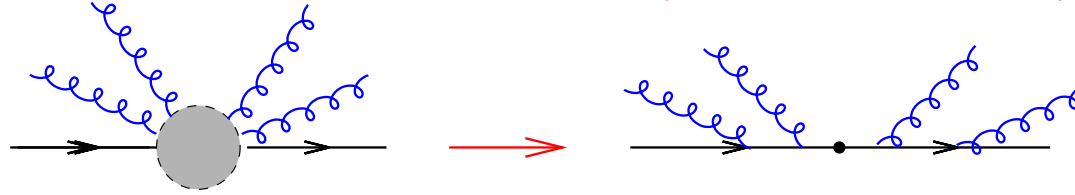
$$1 - T \simeq \frac{1}{Q} \sum_{i=1}^n \frac{E_i \theta_i^2}{2} \quad \longrightarrow \quad \Theta(1 - T < \tau) = \int \frac{d\nu}{2\pi i \nu} e^{\nu\tau} \prod_{i=1}^n e^{-\nu \frac{E_i \theta_i^2}{2Q}}$$

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THE ANSWER

$$\Sigma(\tau) \int \frac{d\nu}{2\pi i \nu} e^{\nu\tau} \exp \left[\int \frac{d\theta}{\theta} \frac{dE}{E} \frac{\alpha_s(E\theta) C_F}{\pi} \left(e^{-\nu \frac{E_i \theta_i^2}{2Q}} - 1 \right) \right]$$

An incomplete list of analytical NLL predictions

$e^+e^- \rightarrow 2$ jets

- ↪ S. Catani, G. Turnock, B. R. Webber and L. Trentadue, *Thrust distribution in e^+e^- annihilation*, Phys. Lett. B **263** (1991) 491.
- ↪ S. Catani, G. Turnock and B. R. Webber, *Heavy jet mass distribution in e^+e^- annihilation*, Phys. Lett. B **272** (1991) 368.
- ↪ S. Catani, Yu. L. Dokshitzer, M. Olsson, G. Turnock and B. R. Webber, *New clustering algorithm for multi-jet cross-sections in e^+e^- annihilation*, Phys. Lett. B **269** (1991) 432.
- ↪ S. Catani, L. Trentadue, G. Turnock and B. R. Webber, *Resummation of large logarithms in e^+e^- event shape distributions*, Nucl. Phys. B **407** (1993) 3.
- ↪ S. Catani, G. Turnock and B. R. Webber, *Jet broadening measures in e^+e^- annihilation*, Phys. Lett. B **295** (1992) 269.
- ↪ G. Dissertori and M. Schmelling, *An Improved theoretical prediction for the two jet rate in e^+e^- annihilation*, Phys. Lett. B **361** (1995) 167.
- ↪ Y. L. Dokshitzer, A. Lucenti, G. Marchesini and G. Salam, *On the QCD analysis of jet broadening*, JHEP **9801** (1998) 011
- ↪ S. Catani and B. R. Webber, *Resummed C-parameter distribution in e^+e^- annihilation*, Phys. Lett. B **427** (1998) 377
- ↪ S. J. Burby and E. W. Glover, *Resumming the light hemisphere mass and narrow jet broadening distributions in e^+e^- annihilation*, JHEP **0104** (2001) 029
- ↪ M. Dasgupta and G. Salam, *Resummation of non-global QCD observables*, Phys. Lett. B **512** (2001) 323
- ↪ C. F. Berger, T. Kucs and G. Sterman, *Event shape / energy flow correlations*, Phys. Rev. D **68** (2003) 014012

DIS 1+1 jet

- ↪ V. Antonelli, M. Dasgupta and G. Salam, *Resummation of thrust distributions in DIS*, JHEP **0002** (2000) 001
- ↪ M. Dasgupta and G. Salam, *Resummation of the jet broadening in DIS*, Eur. Phys. J. C **24** (2002) 213
- ↪ M. Dasgupta and G. Salam, *Resummed event-shape variables in DIS*, JHEP **0208** (2002) 032

e^+e^- , DY, DIS 3 jets

- ↪ A. Banfi, G. Marchesini, Y. L. Dokshitzer and GZ, *QCD analysis of near-to-planar 3-jet events*, JHEP **0007** (2000) 002
- ↪ A. Banfi, Y. L. Dokshitzer, G. Marchesini and GZ, *Near-to-planar 3-jet events in and beyond QCD perturbation theory*, Phys. Lett. B **508** (2001) 269
- ↪ A. Banfi, Y. L. Dokshitzer, G. Marchesini and GZ, *QCD analysis of D-parameter in near-to-planar three-jet events*, JHEP **0105** (2001) 040
- ↪ A. Banfi, G. Marchesini, G. Smye and GZ, *Out-of-plane QCD radiation in hadronic Z0 production*, JHEP **0108** (2001) 047
- ↪ A. Banfi, G. Marchesini, G. Smye and GZ, *Out-of-plane QCD radiation in DIS with high p(t) jets*, JHEP **0111** (2001) 066
- ↪ A. Banfi, G. Marchesini and G. Smye, *Azimuthal correlation in DIS*, JHEP **0204** (2002) 024
- ↪ C. F. Berger, T. Kucs and G. Sterman, *Energy flow in interjet radiation*, Phys. Rev. D **65**, 094031 (2002)

~ 1 observable per article

The current situation can be summarized as follows

- experimental studies **limited by availability** of theoretical calculations

- **error-prone business**, many subtle effects understood on the way

On the previous slide, *only 4 authors*, out of 21, can say that their results were always correct to the accuracy claimed [three of them quit physics. . .]

- there are **many phenomenological applications**

➔ **need to automate resummations** (as for fixed order)

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- there are **many phenomenological applications**

➔ **need to automate resummations** (as for fixed order)

On the other hand

- ▶ resumptions exploit always the same **standard factorization techniques** (for matrix element and observable)
- ▶ the **origin of logarithms is clearly the SAME** for all observables

➔ **automating the job seems feasible**



IDEA: Define a simpler observable

$$V(k_1, \dots, k_n) \implies V_s(k_1, \dots, k_n) \equiv \max\{V(k_1), \dots, V(k_n)\}$$

e. g.

$$B(k_1, \dots, k_n) \equiv \sum_i \frac{k_{ti}}{Q} \implies B_s(k_1, \dots, k_n) \equiv \max\left\{\frac{k_{ti}}{Q}\right\}$$

➡ With just one soft-collinear emission

$$V(k_1, \dots, k_n) = V_s(k_1, \dots, k_n)$$

⇒ same double logs and most of the single logs

➡ Simple factorization (no Mellin integrals)

$$\Theta(V_s - v) = \prod_i \Theta(V_i - v)$$

⇒ analytical resummation straightforward!

Fix a Born event and emit a soft gluon k collinear to a given hard leg ℓ .
We parametrize

$$V_s(k) \simeq d_\ell \left(\frac{k_t}{Q} \right)^{a_\ell} e^{-b_\ell \eta} g_\ell(\phi)$$

$k_t \Rightarrow$ transverse momentum wrt the leg

$\eta \Rightarrow$ rapidity wrt the leg

$\phi \Rightarrow$ azimuthal angle

✎ Σ_s known given the (automatically determined) quantities $a_\ell, b_\ell, d_\ell, g_\ell(\phi)$,
just exponentiating naively the one-gluon result

This account for *all double logs* and *single-logs due to*

- ✓ hard collinear effects
- ✓ soft, large angle emission
- ✓ inclusive gluon splitting

The computation of Σ_s is based on a veto on **single-emissions**

$$V(k_1, \dots, k_n) < v \quad \Longrightarrow \quad V_s \equiv \max[V(k_1), \dots, V(k_n)] < v$$

One then needs to relate the observable to *all secondary emissions*, i.e. account for the **observable specific mismatch between $V(k_1, \dots, k_n)$ and V_s**

- ➡ Physically one needs accurate understanding of the kinematics
- ➡ Mathematically this translates into performing Mellin integrals

We call these multiple emission effects.

How can these observable-specific effects be computed generally?

Aim: compute the mismatch between $\Sigma_s(v_s)$ and $\Sigma(v)$

The two distributions are related by a simple convolution

$$\frac{D(v)}{v} = \int \frac{dv_s}{v_s} D_s(v_s) P(v|v_s) \quad D(v) \equiv \frac{d\Sigma}{dL} \quad L = \text{Lnv}$$

→ $P(v|v_s)$ is the probability to have v given v_s

Since → $D_s(v_s) = e^{-R(v_s)} \Rightarrow$ known analytically

→ $v \sim v_s \Rightarrow$ same LL structure

→ expand and get $D_s(v_s) =_{NLL} D_s(v) e^{-R' \ln(v/v_s)} \quad R' \equiv dR/dL$

$$\rightarrow D(v) =_{NLL} D_s(v) \mathcal{F}(R') \quad \mathcal{F}(R') = \int \frac{dv_s}{v_s} e^{-R' \ln(v/v_s)} v P(v|v_s)$$

How to compute $\mathcal{F} \Leftrightarrow P(v|v_s)$ generally?

Fix a Born configuration and generate *decreasing soft-collinear (SC) emissions* according to phase space

① set $v(k_1) = v_s$ [START FROM: $V_s = v_s$]

② generate a formally infinite number of SC emissions

according to an independent emission pattern uniform in $\ln k_t, \eta, \phi$ such that on average the density of emissions per unit $\ln V$ from leg ℓ is R'_ℓ

➔ Finally compute $V(k_1, k_2, \dots, k_n) \equiv v$

☞ This gives the weighted probability of having $V = v$ given $V_s = v_s$ and allows so the computation of \mathcal{F} in a completely general way

Banfi , Salam, GZ *JHEP* 0201 (2002) 018

<http://www.ippp.dur.ac.uk/~zander/numsum.html>

The master formula

$$\Sigma(v) =_{NLL} \sum_{\text{sub.}} \int [d\Phi]_{\text{hard}} \Sigma_s(v) \cdot \mathcal{F}(R')$$

Banfi , Salam, GZ hep-ph/0304148

- ✓ Analytical resummation for the “easy” Σ_s : *pure LL and NLL terms*

$$\Sigma_s(v) = \prod_{\ell=1}^{n_{inc}} \underbrace{f_{\ell}(v^{\frac{2}{a+b_{\ell}}}, \mu_F^2)}_{\text{pdfs}} \otimes \prod_{\ell=1}^N \underbrace{J_{\ell}(L)}_{\text{jet function}} \cdot \underbrace{S(T(L/a))}_{\text{soft}}$$

- ✎ soft and collinear emission \Rightarrow **jet function $J_{\ell}(L)$**
(all LL Sudakov suppression and some NLL terms)
- ✎ hard collinear splitting \Rightarrow **evolution of the pdfs**
- ✎ soft large angle \Rightarrow **QCD coherence and geometry dependence in S**

- ✓ the observable-dependent “difficult” \mathcal{F} is computed numerically but is *by construction a pure NLL function*

Requirements on the observable

For the observable to be resummed automatically it should

- ✗ vanish in the Born limit and be positive defined
- ✗ behave as $V(k) \simeq d_\ell \left(\frac{k_t}{Q}\right)^{a_\ell} e^{-b_\ell \eta} g_\ell(\phi)$ for 1 SC gluon along leg ℓ
- ✗ be infrared and collinear safe
- ✗ be continuously global ($a_\ell = a \quad \forall$ hard legs ℓ)
- ✗ exponentiate (no JADE)

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- ✗ be **infrared and collinear safe**
- ✗ be **continuously global** ($a_\ell = a \quad \forall$ hard legs ℓ)
- ✗ **exponentiate** (no JADE)

While this might seem a long list

- ✚ practically the **limiting condition** is the requirement of **globalness**
[all other conditions are satisfied by all observables resummed so far]

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While this might seem a long list

- ☞ practically the limiting condition is the requirement of globalness
[all other conditions are satisfied by all observables resummed so far]
- ☞ the essential feature of the program is the ability to perform all checks automatically
[☞ use arbitrary precision to take asymptotic limits]

Bailey, RNR Technical Report RNR-94-013

Some observables have exponentiating double (and single) logs

$$P(v) = 1 - X \frac{\alpha_s C_F}{\pi} \ln^2 v + \frac{1}{2} X^2 \left(\frac{\alpha_s C_F}{\pi} \right)^2 \ln^4 v + \dots \Rightarrow e^{-X \frac{\alpha_s C_F}{\pi} \ln^2 v}$$

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others do not, e.g. Jade-algorithm jet rates:

$$P_{\text{Jade2-jet}}(y_{\text{cut}}) = 1 - \frac{\alpha_s C_F}{\pi} \ln^2 y_{\text{cut}} + \frac{1}{2} \cdot \frac{5}{6} \left(\frac{\alpha_s C_F}{\pi} \right)^2 \ln^4 y_{\text{cut}} + \dots$$

Brown and Stirling, Phys.Lett.B 252 (1990)

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Brown and Stirling, Phys.Lett.B 252 (1990)

➡ No one jet knows how to resum Double Logs, let alone what matrix-element ingredients are needed to achieve NLL accuracy!

Any automated approach to NLL resummation has better be able to establish whether an observables exponentiates

Consider n emissions $k_1(\lambda_1), \dots, k_n(\lambda_n)$ such that the **soft-collinear limit** corresponds to $\lambda_i \rightarrow 0$ and $V(k_i) = \lambda_i$. Then

Normal IRC safety implies

$$\lim_{\epsilon \rightarrow 0} V(k_1(\lambda_1), \dots, k_n(\lambda_n), k_{n+1}(\epsilon\lambda_{n+1})) = V(k_1(\lambda_1), \dots, k_n(\lambda_n))$$

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Recursive IRC safety adds two conditions

(a) $\lim_{\epsilon' \rightarrow 0} V(k_1(\epsilon' \lambda_1), \dots, k_n(\epsilon' \lambda_n)) / \epsilon' = \text{const.} (\neq 0)$

the SC scaling properties of V should be the same with just one or many emissions

(b) $\lim_{\epsilon \rightarrow 0} \lim_{\epsilon' \rightarrow 0} V(k_1(\epsilon' \lambda_1), \dots, k_n(\epsilon' \lambda_n), k_{n+1}(\epsilon \epsilon' \lambda_{n+1})) / \epsilon' = \text{same const.}$

i. e. the addition of a relatively much softer/more collinear parton should not change asymptotically the limit

This condition is *the formal requirement for exponentiation*

The condition of IRS safety allows one to translate

- a restriction on **an ensemble of emissions** $\Rightarrow V(k_1, \dots, k_n) < v$ into
- a restriction on **individual emissions** $\Rightarrow V(k_i) < v$ (modulo NLL terms in \mathcal{F})

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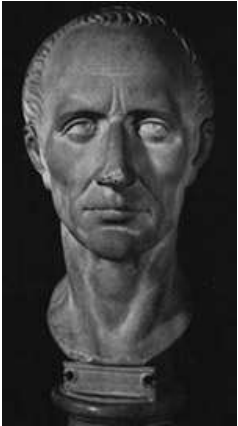
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Example of observables NOT satisfying the condition

- Jet rates in **Jade**-algorithm
- Combinations of “usual” event shapes $\tau \cdot B_T, B_T^3 / (1 - \tau), y_{3D} \cdot C \dots$



Computer Automated Expert Semi-Analytical Resummer

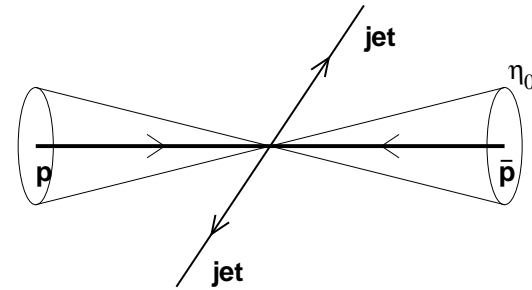


- currently limited to global observables
- tested against all known global, exponentiable event shapes
- results from an early version used by the LEP-QCD-WG for fits of α_s
- can be applied to
 - 2 & 3 jets in e^+e^-
 - [1+1] & [1+2] jets in DIS
 - Drell-Yan + 1 jet
 - hadron-hadron dijet events [\Leftarrow first resummations]

Observables in hadronic dijet production

Cut around the beam $|\eta| < \eta_0$

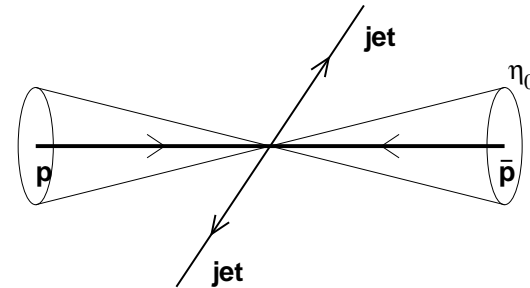
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Directly global observables: $\eta_0 > 1$

X Transverse thrust

$$T_T = \frac{1}{E_T} \max_{\vec{n}_T} \sum_i |\vec{p}_{ti} \cdot \vec{n}_T|$$

X Thrust minor

$$T_m = \frac{1}{E_T} \sum_i |p_i^{out}|$$

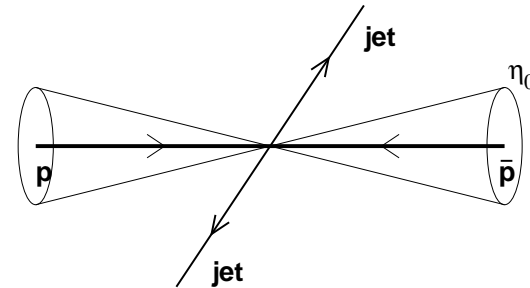
Predictions valid as long as

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Indirectly global observables: $\eta_0 = \mathcal{O}(1)$

✗ Transverse thrust

$$T_T = \frac{1}{E_{T,\eta_0}} \left(\max_{\vec{n}_T} \sum_{|\eta_i| < \eta_0} |\vec{p}_{ti} \cdot \vec{n}_T| - \left| \sum_{|\eta_i| < \eta_0} \vec{p}_{ti} \right| \right)$$

✗ Thrust minor

$$T_m = \frac{1}{E_{T,\eta_0}} \left(\sum_{|\eta_i| < \eta_0} |p_i^{out}| + \left| \sum_{|\eta_i| < \eta_0} \vec{p}_{ti} \right| \right)$$

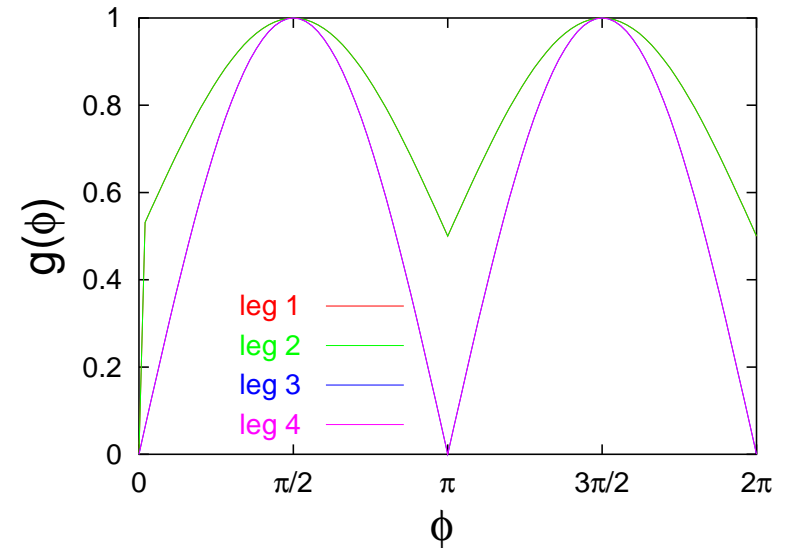
Predictions valid as usual,

but \mathcal{F} diverges at $R' = R'_c$

Sample output: the indirectly global thrust minor

X Tests on the observable

Test	result
check number of jets	T
observable positive	T
global	T
continuously global	T
additive	F
exponentiate	T
eliminate subleading effects	T
opt. probe region exists	T



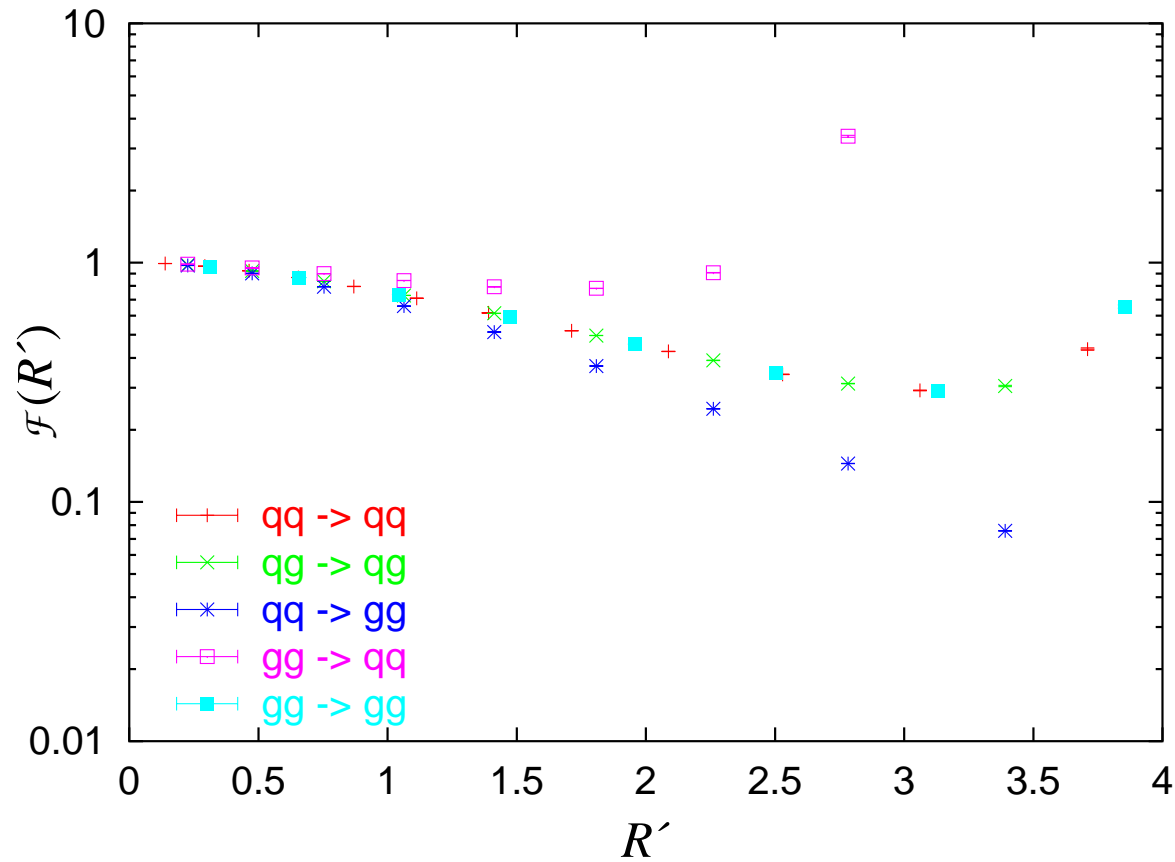
X Single emission properties

leg ℓ	a_ℓ	b_ℓ	$g_\ell(\phi)$	d_ℓ	$\langle \ln g_\ell(\phi) \rangle$
1	1	0	tabulated	2.0000	-0.2201
2	1	0	tabulated	2.0000	-0.2201
3	1	0	$\sin(\phi)$	2.0000	$-\ln(2)$
4	1	0	$\sin(\phi)$	2.0000	$-\ln(2)$

← Tables and plots generated automatically by CAESAR

$\mathcal{F}(R')$ for the indirectly global thrust minor

The multiple emission function $\mathcal{F}(R')$

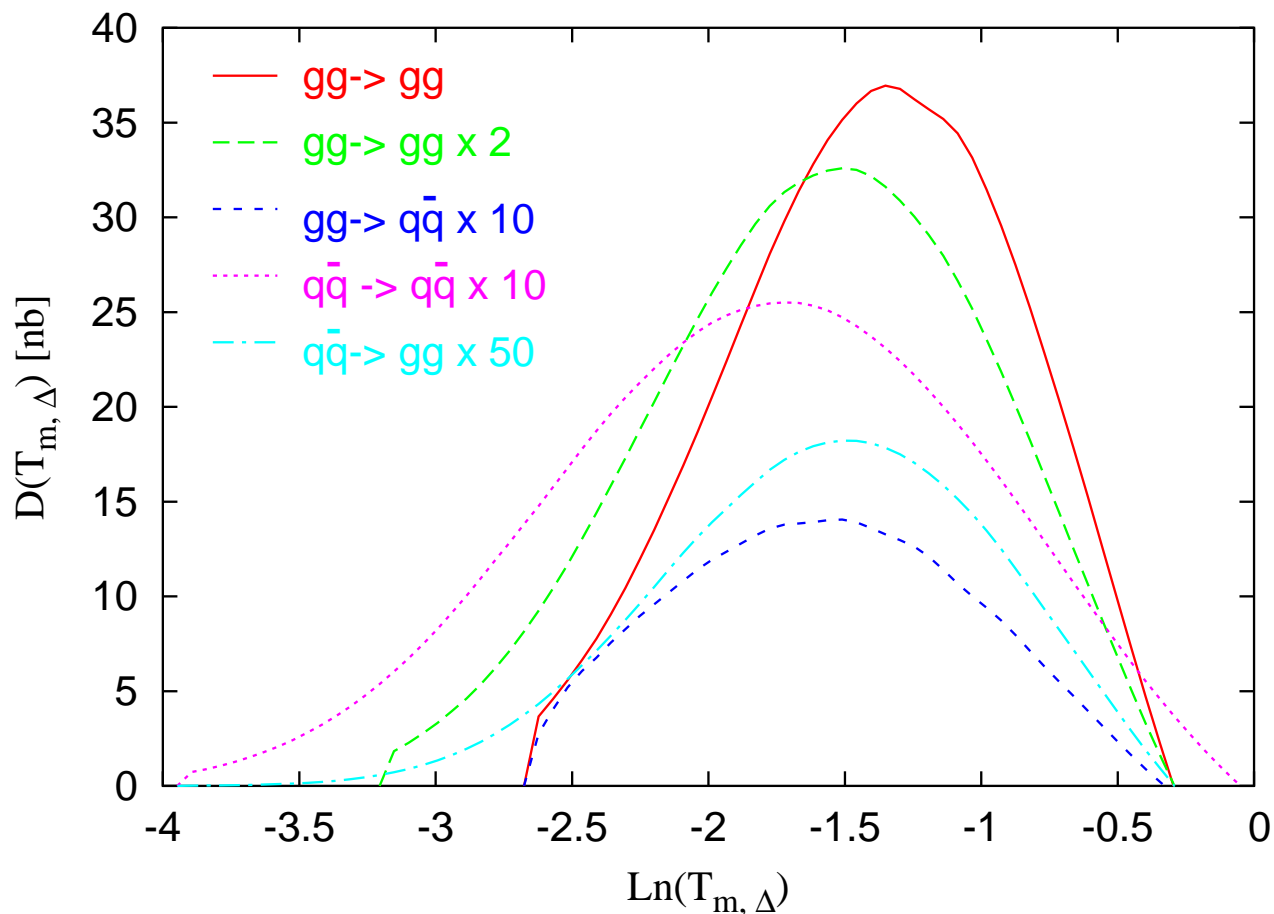


➡ Different result for different colour configurations

The indirectly global thrust minor

Dijets events at Tevatron run II regime

- ▶ run II regime $\sqrt{s} = 1.96 \text{ TeV}$
- ▶ cut on jet transverse energy $E_T > 50 \text{ GeV}$ and on rapidity $|\eta| < 1$



$$\alpha_s(M_Z) = 0.118$$

$$\mu_F = \mu_R = P_T$$

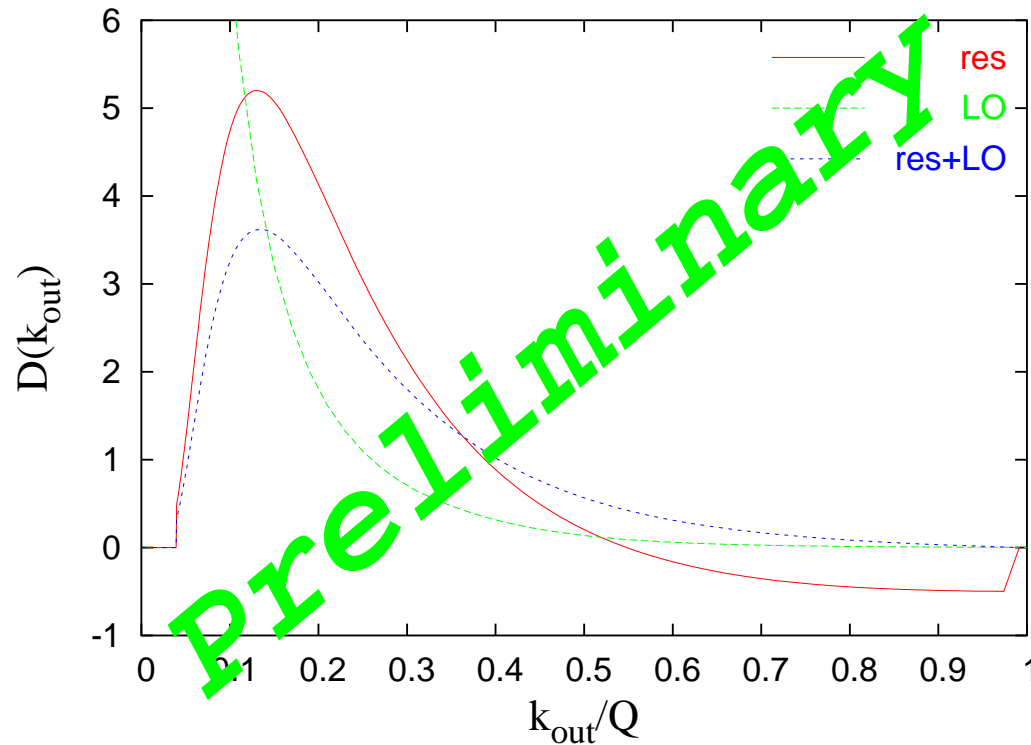
$$X_c = 1$$

PDFS: CTEQ6M

Out-of plane radiation in DIS [1+2] jet events

Dijets events at Hera

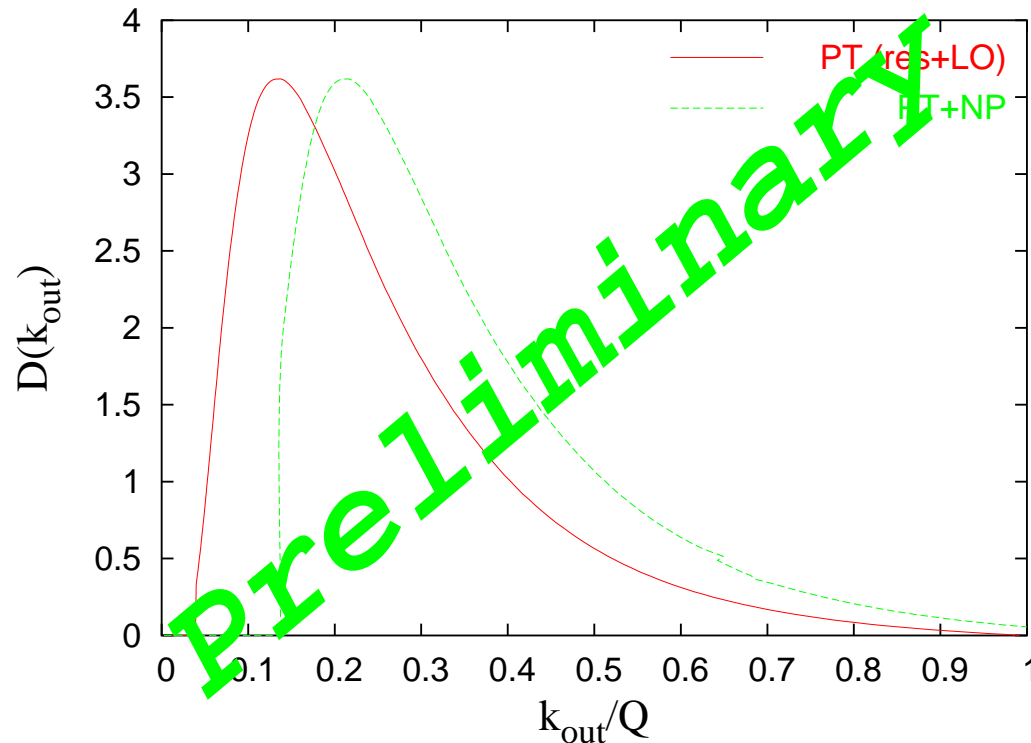
- ▶ Kinematical variables: $\sqrt{s} = 300 \text{ GeV}$ $Q = 36.7 \text{ GeV}$ $x_B = 0.056$
- ▶ Cuts: $y_{cut} = 0.1$ $\eta_{max} = 3$
- ▶ Scale choice and PDFs: $\alpha_s(M_Z) = 0.118$ $\mu_F = \mu_R = P_T$ PDFS: CTEQ6M



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NP-shift: Banfi , Dokshitzer, Marchesini, GZ, hep-ph/0111157

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Monte Carlo event generators (Herwig, Pythia . . .) do already embody many resummation ingredients (parton showering).

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- subleading effects always present and hard to estimate
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Our predictions

- do not contain subleading Logs \Rightarrow matching feasible
- purely perturbative, any hadronization model can be apply on top
- allow studies of factorization, renormalization scale dependencies
- are limited to a precise, well-defined class of observables

Conclusions & outlook

Main result: rigorous procedure to perform resummation semi-analytically

Banfi , Salam, GZ hep-ph/0304148

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X Most relevant applications

- *Theoretical*: criterion of recursive infrared and collinear safety
- *Experimental*: first NLL predictions in hadronic dijet events

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X To-do list and wish list

- automated matching of NLL with NLO(JET++)
- extension non-global observables and inclusion of mass effects