# Fermi node, Topological Phases and Domain Walls of Pyrochlore Iridates

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KITP , July 30, 2015





# Part 1 Fermi node and Topological Phases of Pyrochlore Iridates

UCSB

Leon Balents Eun-Gook Moon (now U. Chicago)





Japan

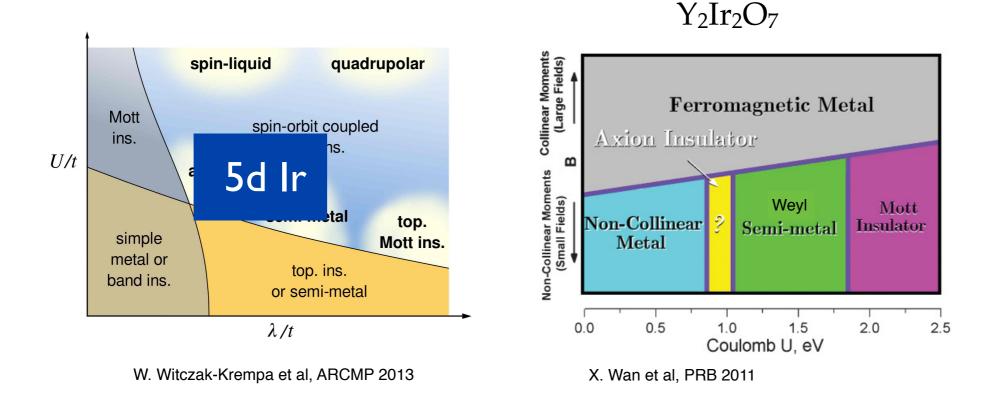
Takeshi Kondo Satoru Nakatsuji Shik Shin

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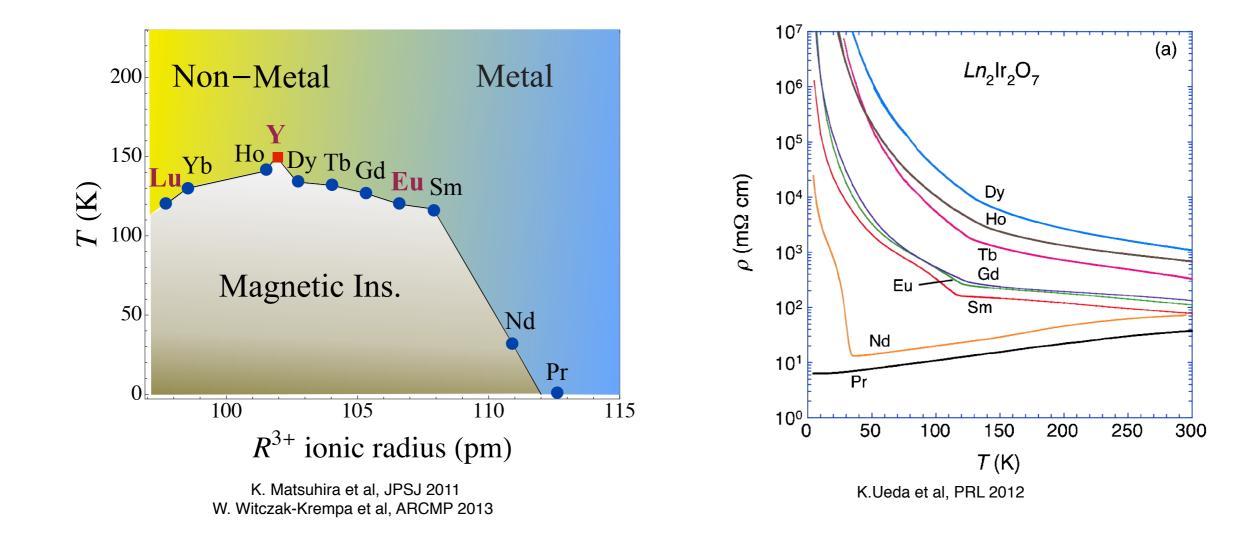




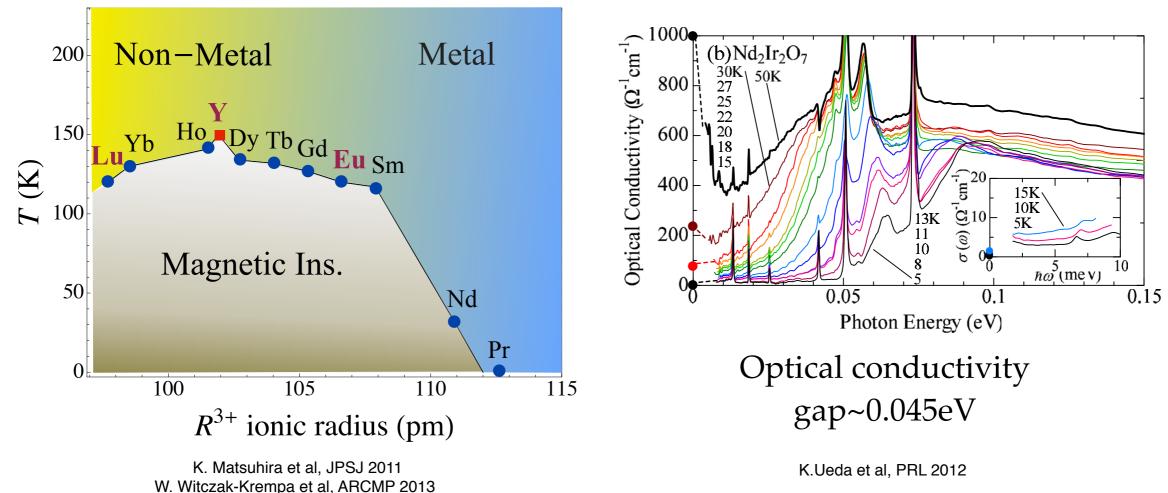
- Correlation physics: Mott insulator, superconductivity...
- Spin-orbit coupling: topological insulators, Majorana fermions...
- Coulomb interaction+spin-orbit coupling → new physics?
- Iridates: U(correlation)~W(band width)~ $\lambda$ (spin-orbit)



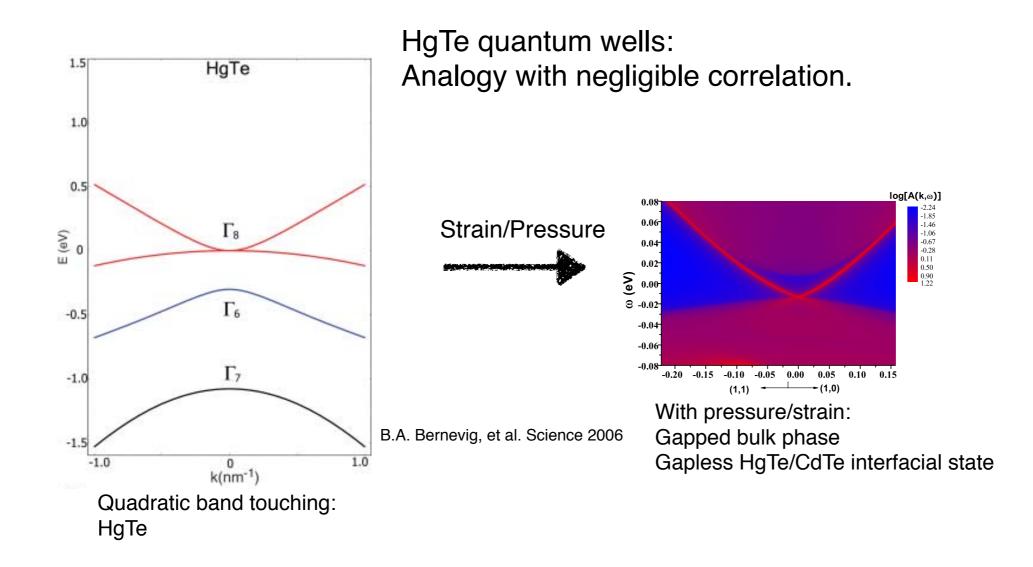
- Magnetic ordering transition, 2nd order transition
- Magnetic transition and metal-insulator transition coincides for relatively larger rare earth radius compound
- As rare earth ionic radius decreases, T<sub>MI</sub> increases and gap increases



- Magnetic ordering transition, 2nd order transition
- Magnetic transition and metal-insulator transition coincides for relatively larger rare earth radius compound
- As rare earth ionic radius decreases, T<sub>MI</sub> increases and gap increases
- Energy gap is small comparing to 3d materials



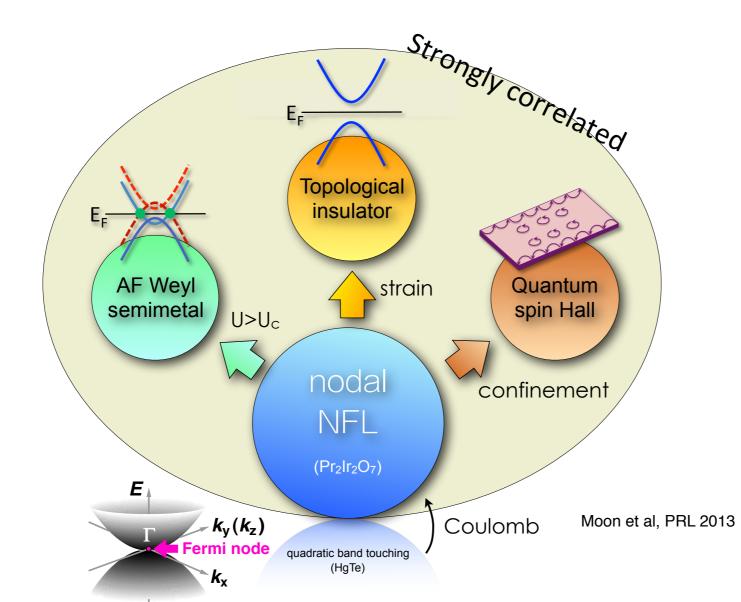
- What is the minimum low energy electronic structure of R<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>? Especially for more metallic ones.
- Quadratic Fermi node in Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> (*ab initio* electronic structure calculation and APRES data)



• Tune the ground state ? If so, how?

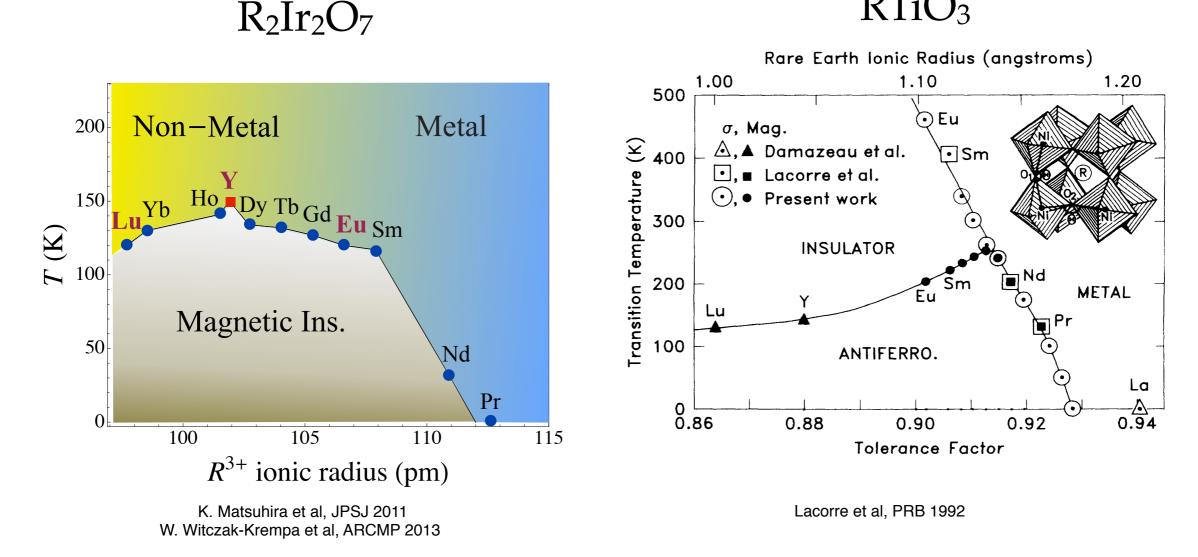
Goal: Topological insulators, quantum anomalous Hall effect, Weyl semimetal

• Growing techniques such as MBE and PLD, magnetic and pressure measurement have made tuning more accessible.

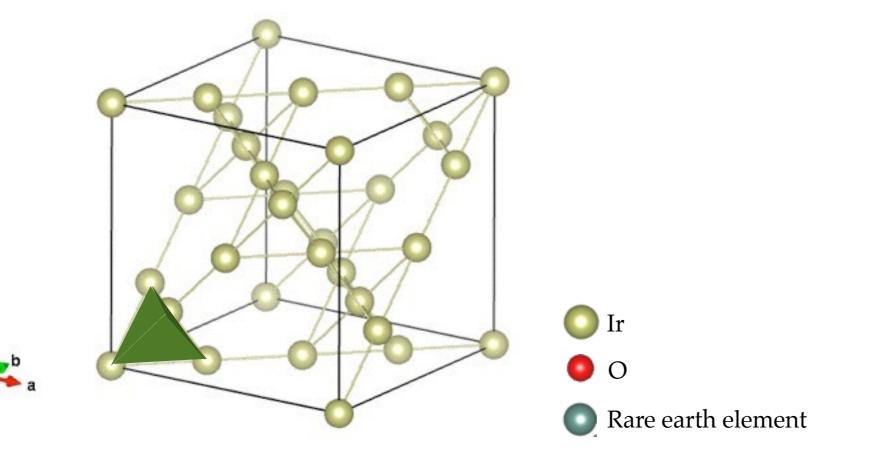


• What are the important microscopic parameters which drive the metal-insulator transition? What drives MIT?

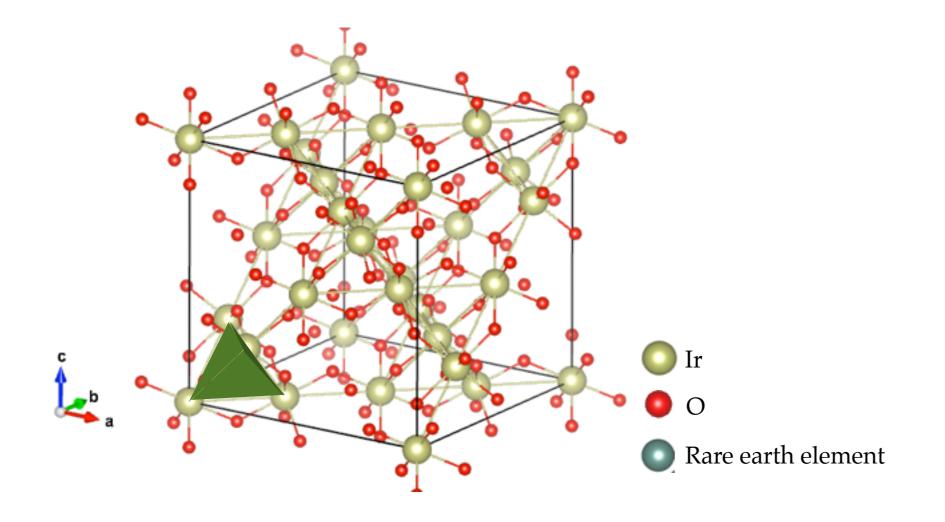
RNiO<sub>3</sub> RTiO<sub>3</sub>



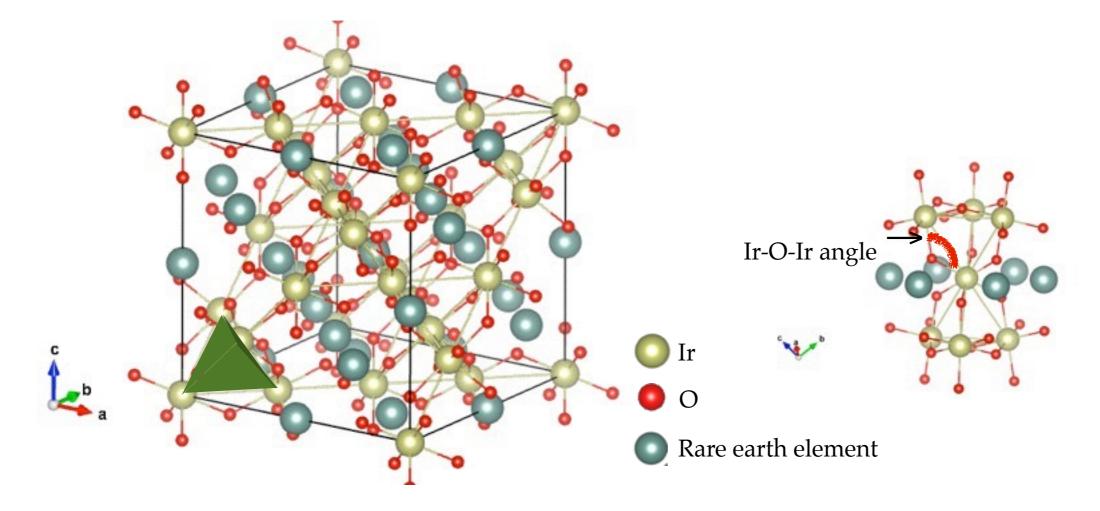
- Structure of pyrocholre iridate R<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> : FCC corner sharing tetrahedra
- Ir-O-Ir bonds are distorted, compressed along 111 direction



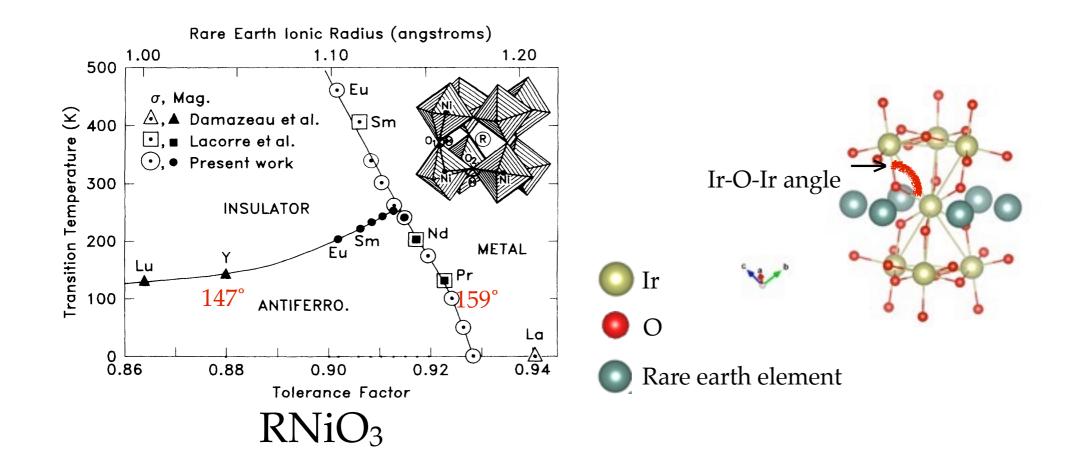
- Structure of pyrochlore(correct all!) iridate R<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>:
  FCC corner sharing tetrahedra
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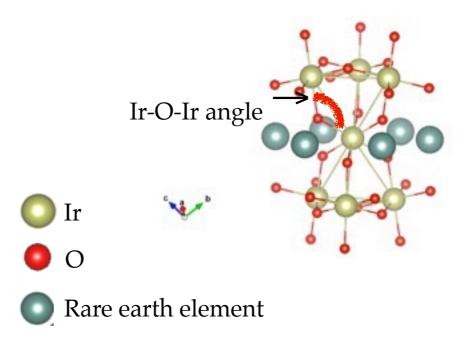
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- Ir-O-Ir bond angle Y: ~129°; Pr: ~132°



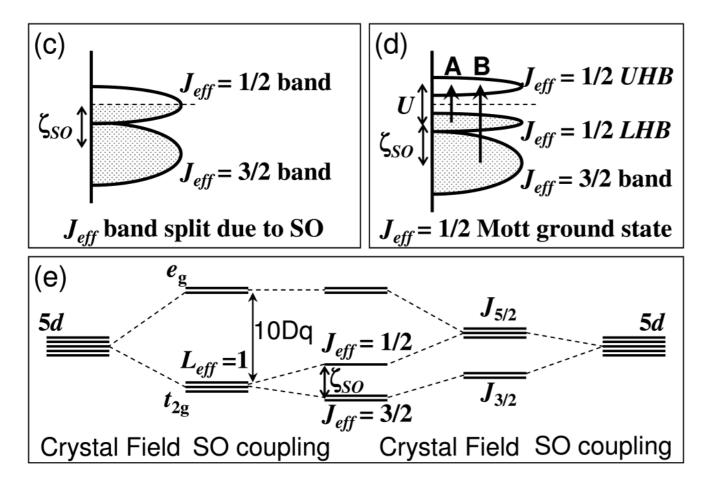
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- Structure of pyrocholre iridate R<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> : FCC corner sharing tetrahedra
- Ir-O-Ir bond are distorted, compressed along 111 direction
- Ir-O-Ir bond angle Y: ~129°; Pr: ~132°
- Ir-O bond length Y: 1.997Å ; Pr: 2.014Å
- Ir-Ir bond length Y: 3.599Å ; Pr: 3.677Å

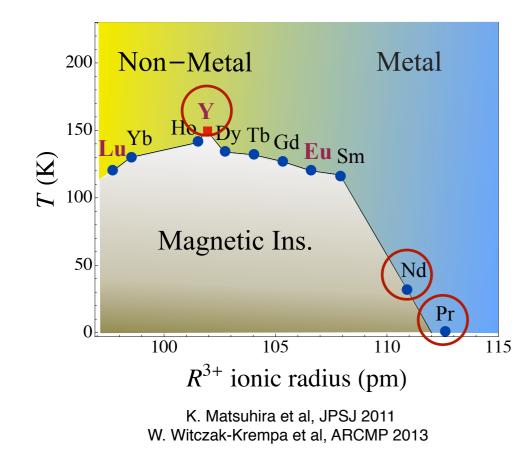


- Ionic picture: Ir<sup>4+</sup> 5d<sup>5</sup>
- J<sub>eff</sub>=1/2 doublet

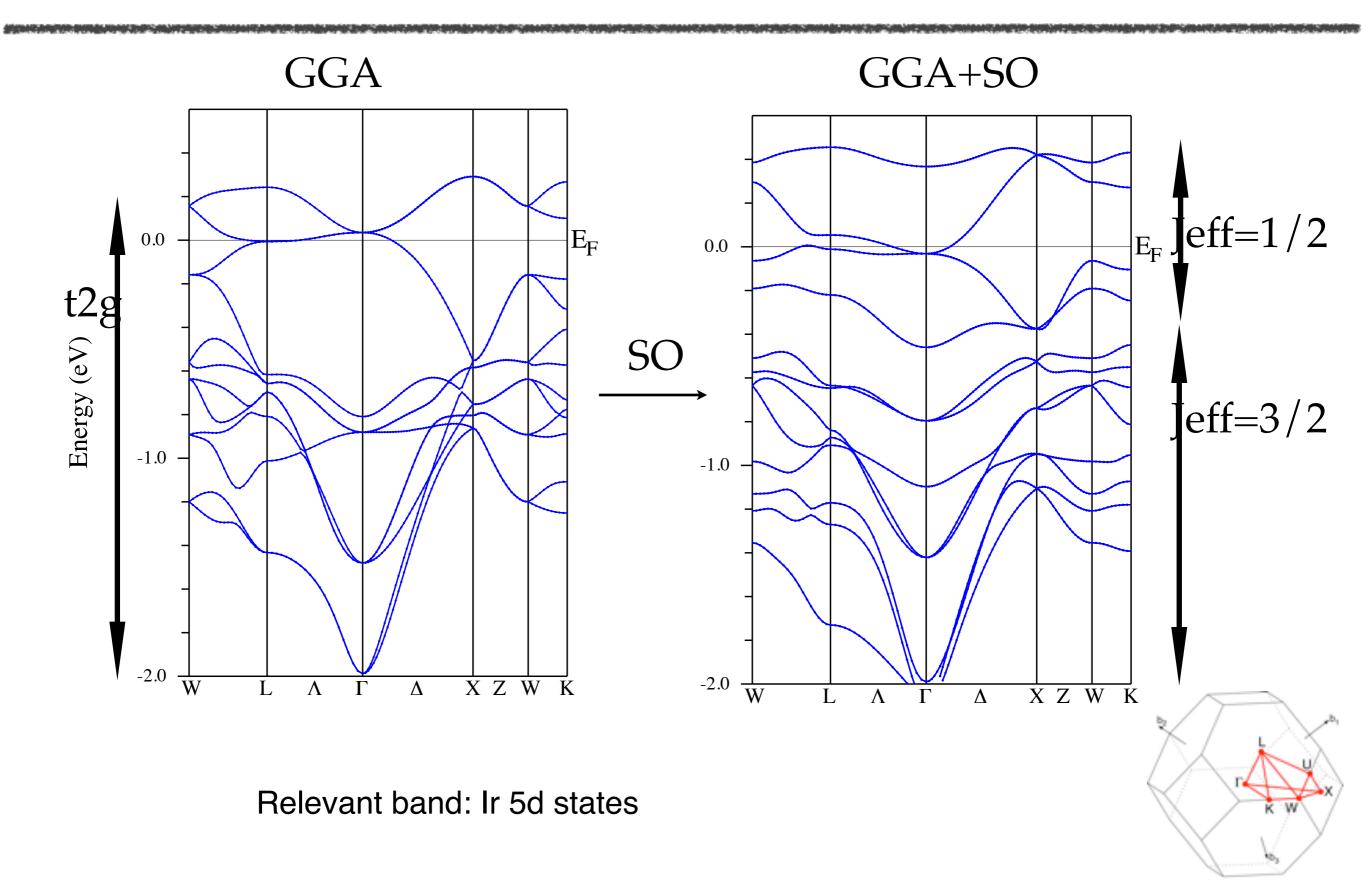


# Paramagnetic band structure calculation

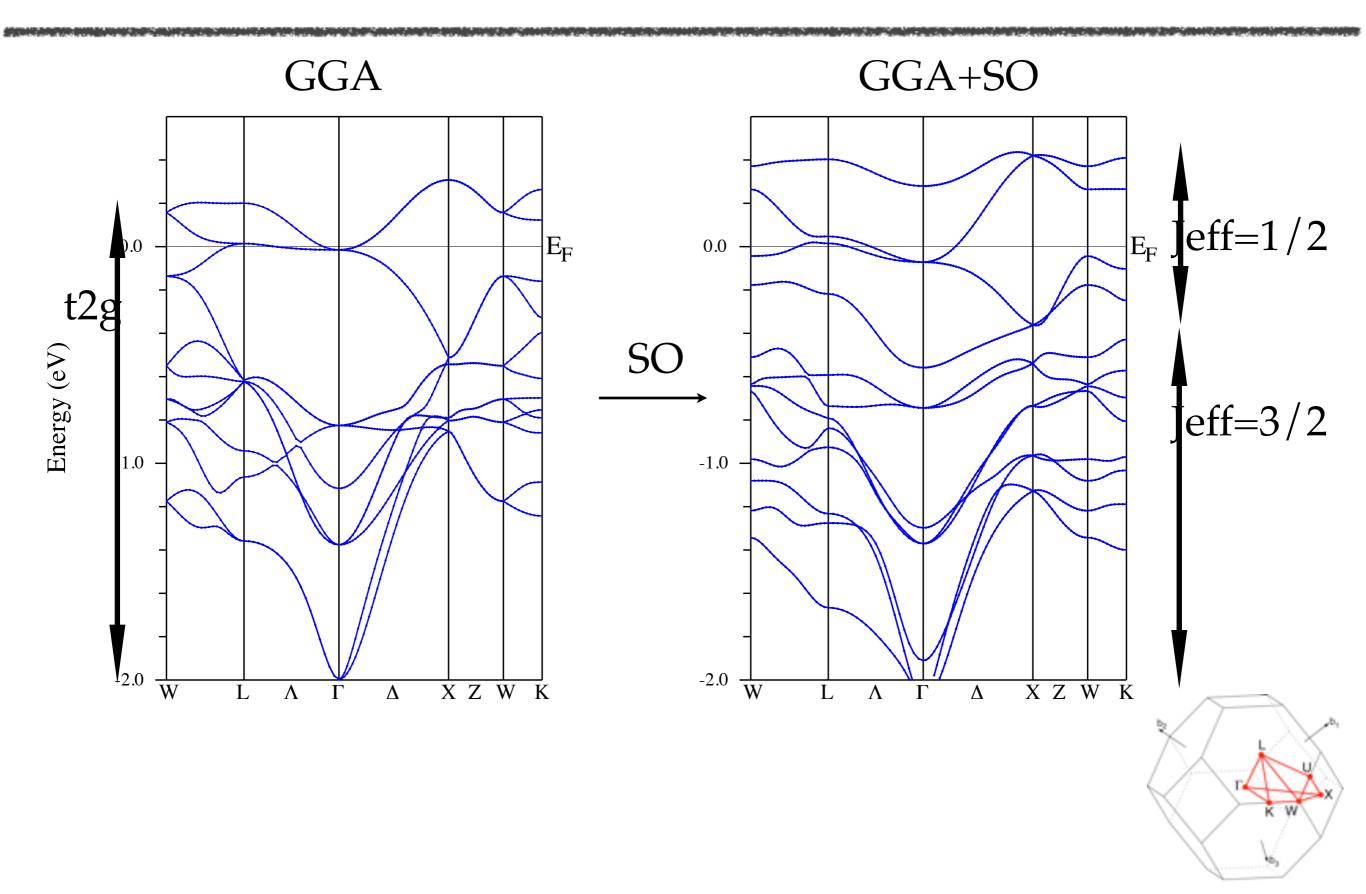
- Focus on Pr, Nd, Y
- Various functional to explore the trend: GGA+SO, LDA+SO, mBJ+SO
- All electronic potential Wien2k



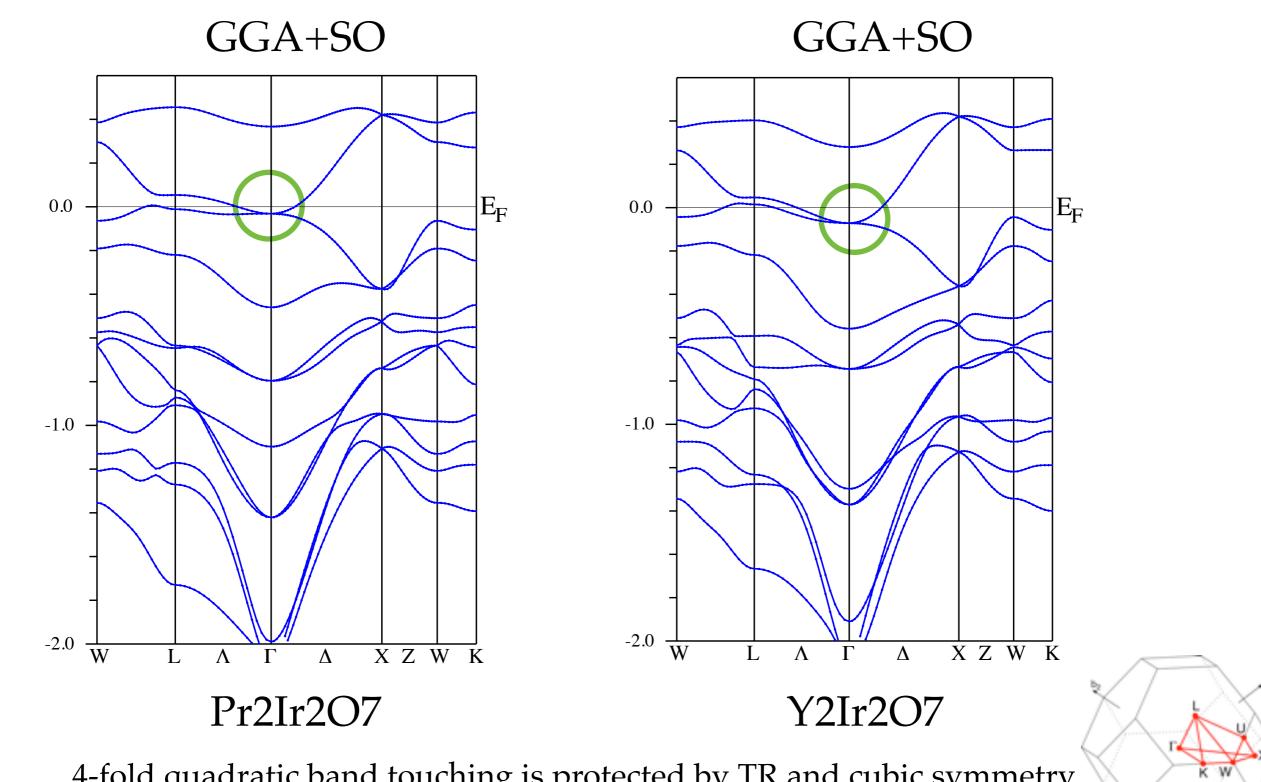
#### $Pr_2Ir_2O_7$



#### $Y_2Ir_2O_7$

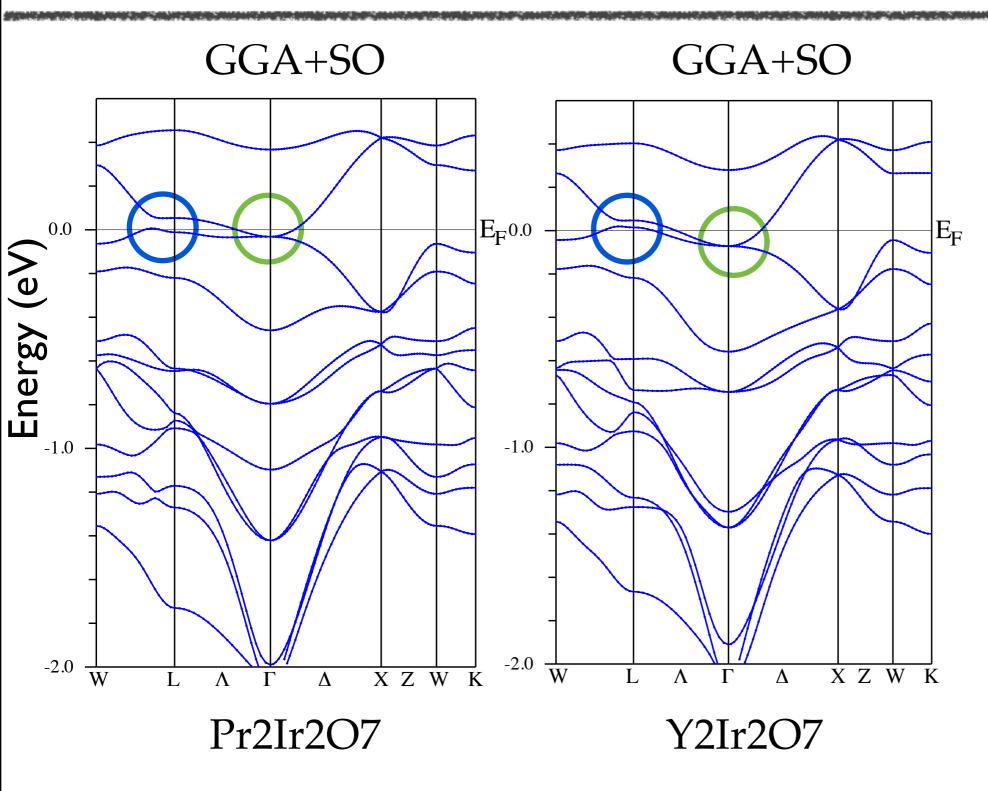


#### Comparison: Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> and Y<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>



4-fold quadratic band touching is protected by TR and cubic symmetry.

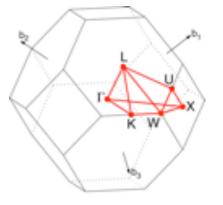
## Comparison: Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> and Y<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>



• Not band width tuned (hopping Ir-O-Ir, Ir-Ir competing)

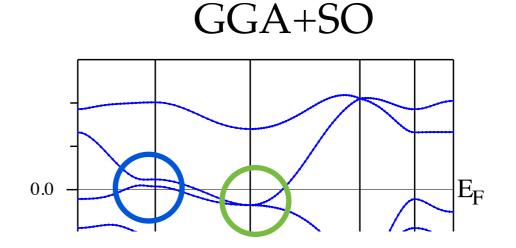
• Very subtle changes near Ef

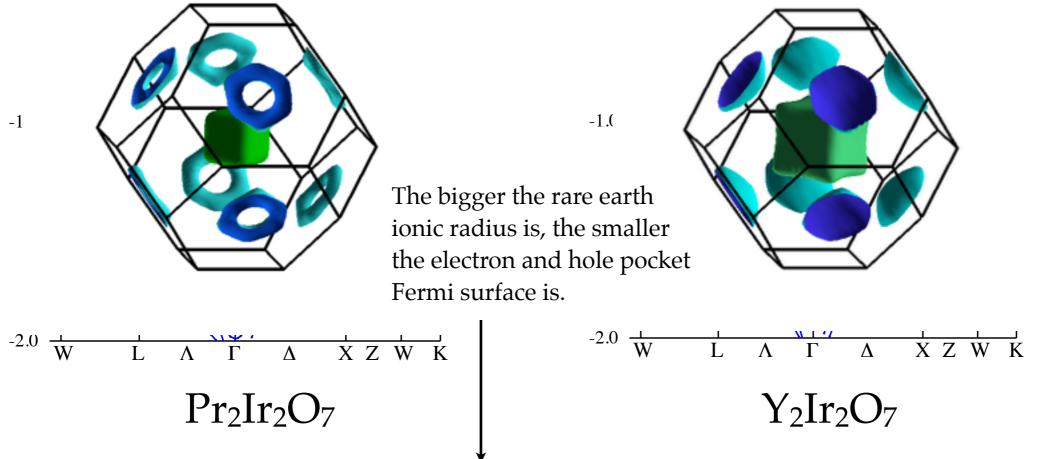
• Larger R is closer to quadratic band touching at Ef (R=Pr, Nd,Y)



#### Comparison: Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> and Y<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>

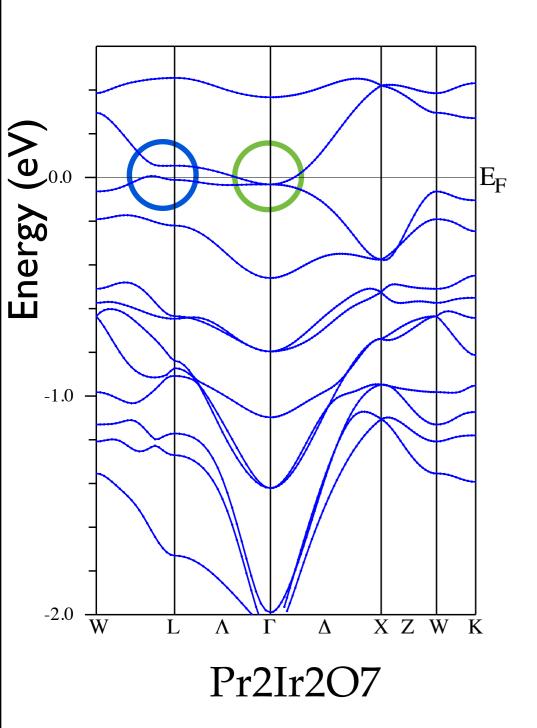
GGA+SO





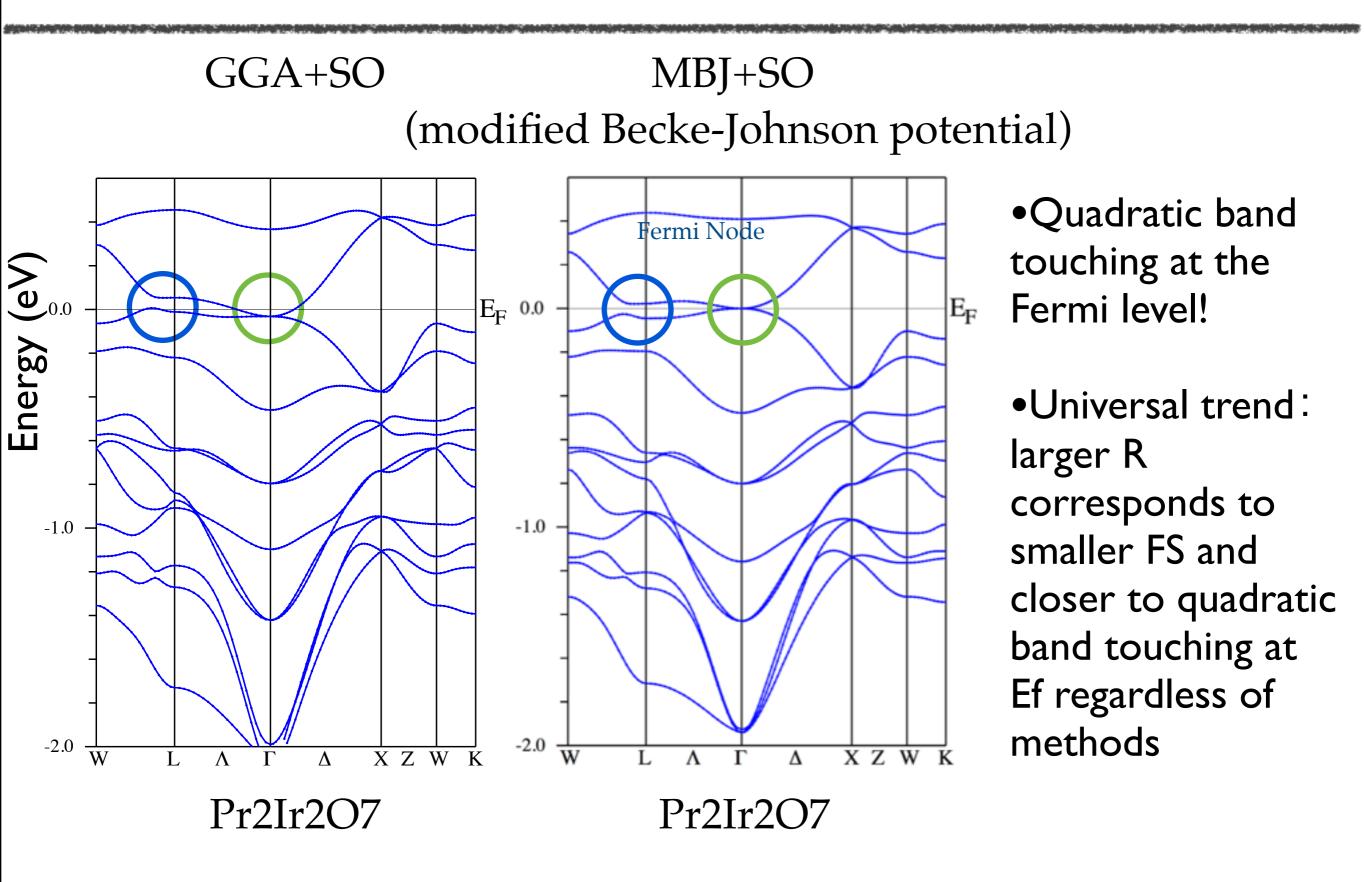
Why is Pr metallic and Y not?

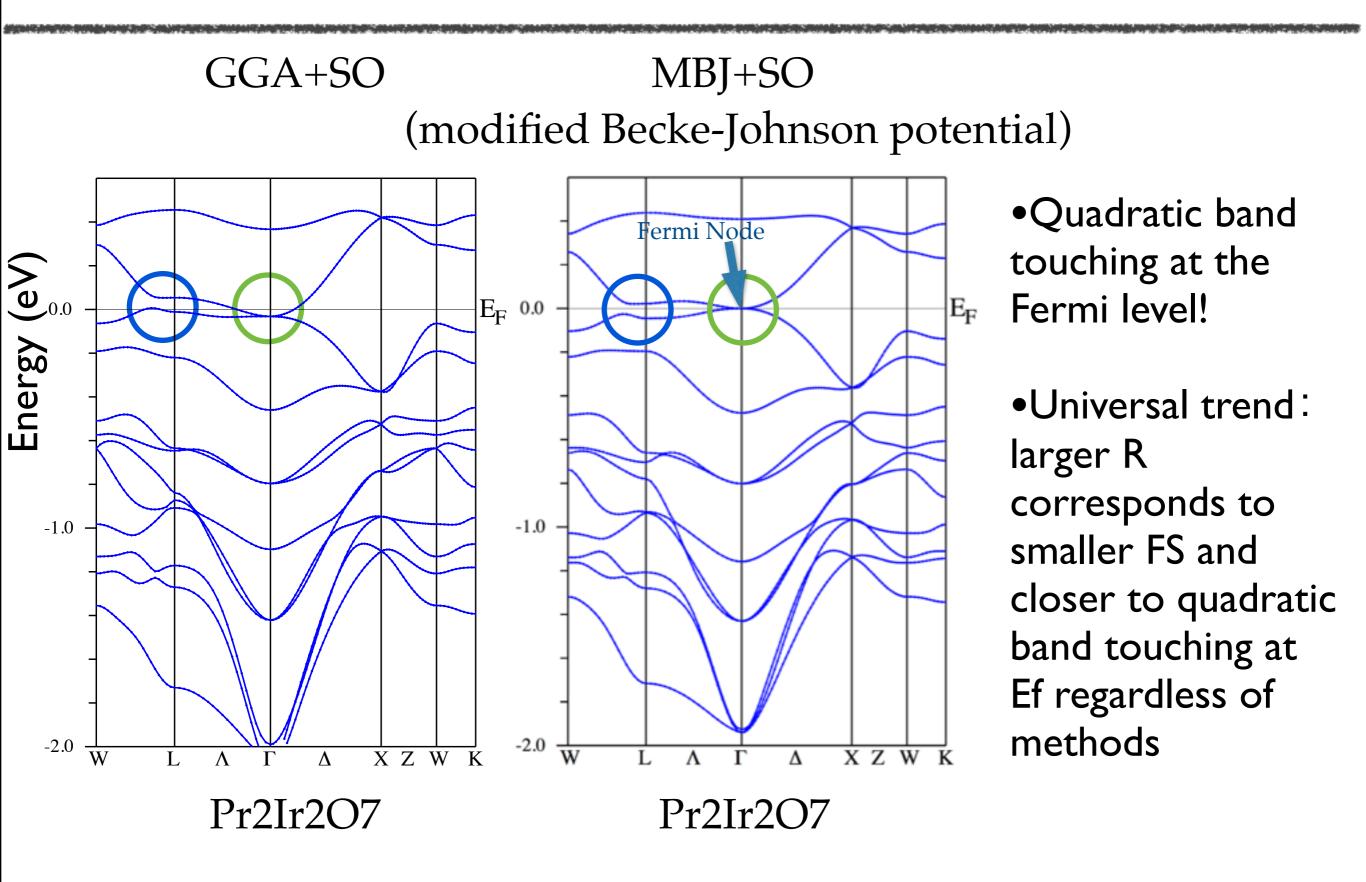
GGA+SO

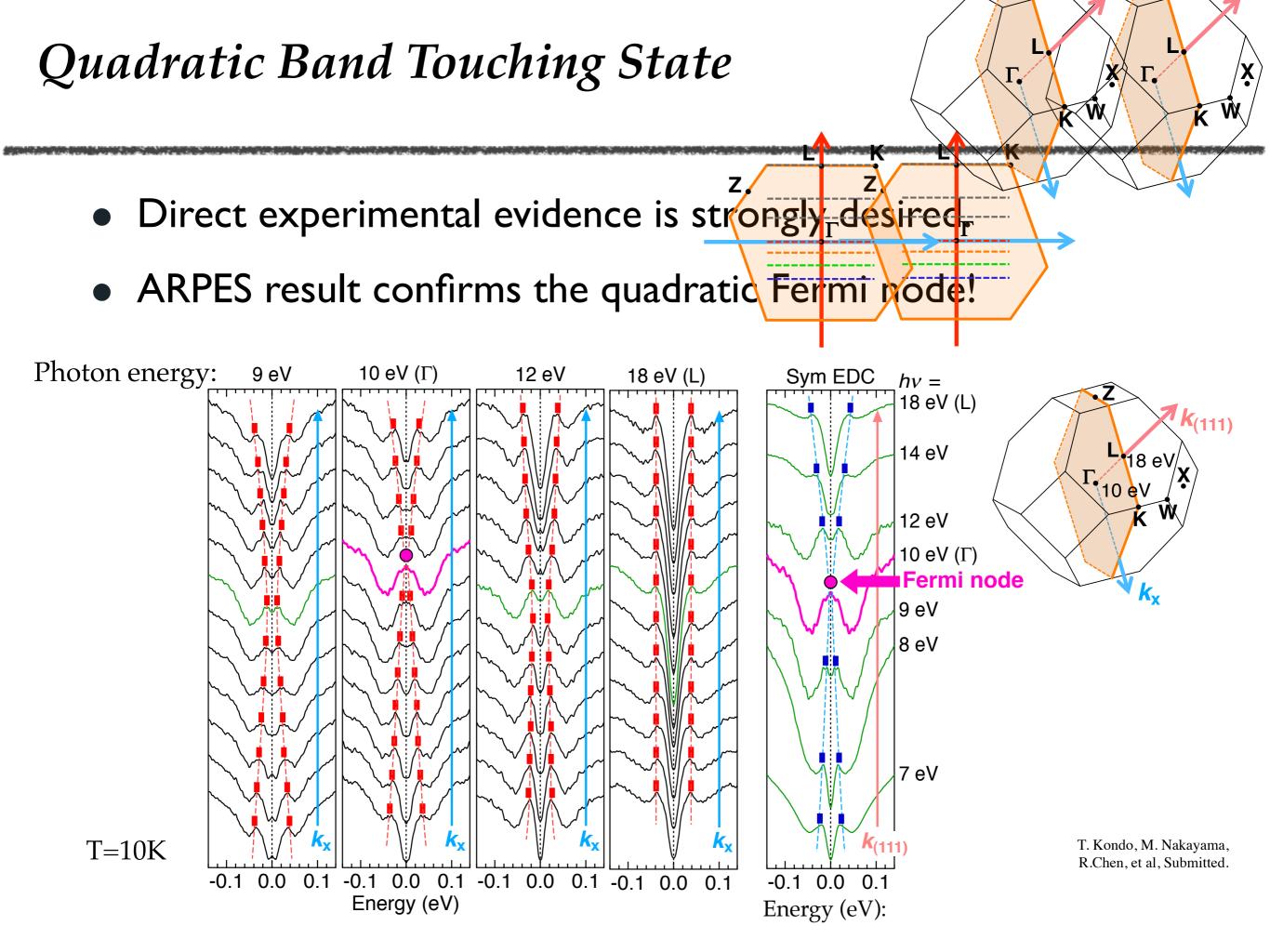


- GGA+SO calculations show Pr2Ir2O7 is very close to quadratic band touching at the Fermi level.
- 4-fold degeneracy at the  $\Gamma$  point is protected by the cubic lattice symmetry, everywhere else gapped.
- •GGA/LDA is well known for underestimating gaps.
- •Modified Becke-Johnson potential gives band gap very close to experiments/GW calculation.

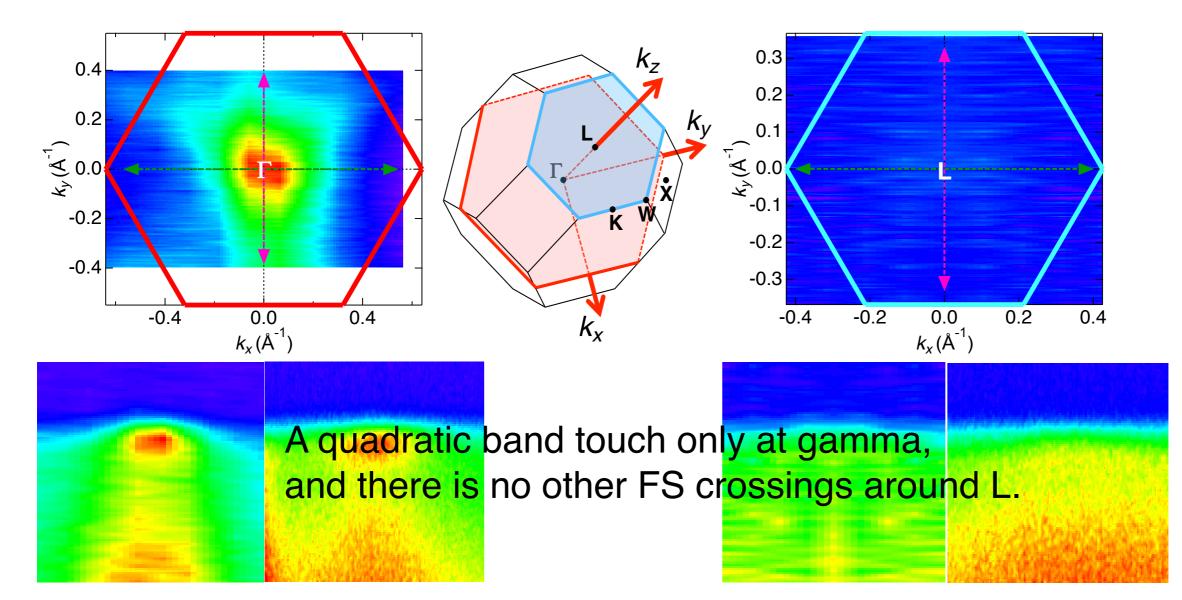
•Widely used for small gap semiconductor, topological insulator.







- Direct experimental evidence is strongly desired.
- ARPES result confirms the quadratic Fermi node!



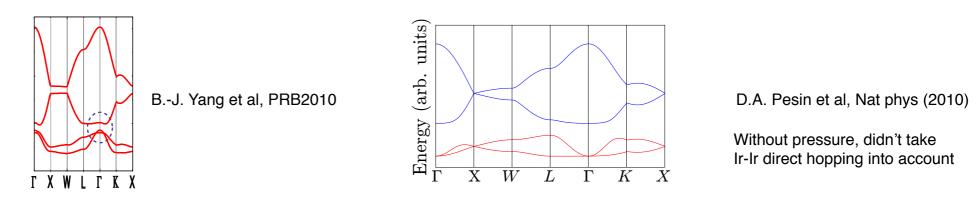
- k.p theory, Luttinger Hamiltonian around  $\Gamma$  point with cubic symmetry  $\mathcal{H}_0(k) = \frac{k^2}{2\tilde{M}_0} + \frac{\frac{5}{4}k^2 - (\vec{k}\cdot\vec{J})^2}{2m} - \frac{(k_x^2J_x^2 + k_y^2J_y^2 + k_z^2J_z^2)}{2M_c}$ E.G. Moon et al, PRL 2013
- Jeff= 3/2 matrix
- Similar to electronic structure of semiconductors with diamond structure and zinc blende structure in terms of the kp perturbation theory
- Effective mass fitting to the band structure:  $M_0 = 19.97$  me, m = 6.30 me ,  $M_c = 7.44$  me

- k.p theory, Luttinger Hamiltonian around  $\Gamma$  point with cubic symmetry  $\mathcal{H}_0(k) = \frac{k^2}{2\tilde{M}_0} + \frac{\frac{5}{4}k^2 - (\vec{k}\cdot\vec{J})^2}{2m} - \frac{(k_x^2J_x^2 + k_y^2J_y^2 + k_z^2J_z^2)}{2M_c}$ E.-G. Moon et al, PRL 2013
- Strain/pressure will break the cubic symmetry  $\mathcal{H}' = -\delta(J_z^2 \frac{5}{4})$
- Consider strain/pressure along (III) direction

$$\mathbf{k} = k\hat{n}_{(111)}$$
$$E_{\pm}(k) = \frac{k^2}{2M_0} \pm \sqrt{\left(\frac{k^2(m+M_c) + 2\delta m M_c}{2m M_c}\right)^2}$$

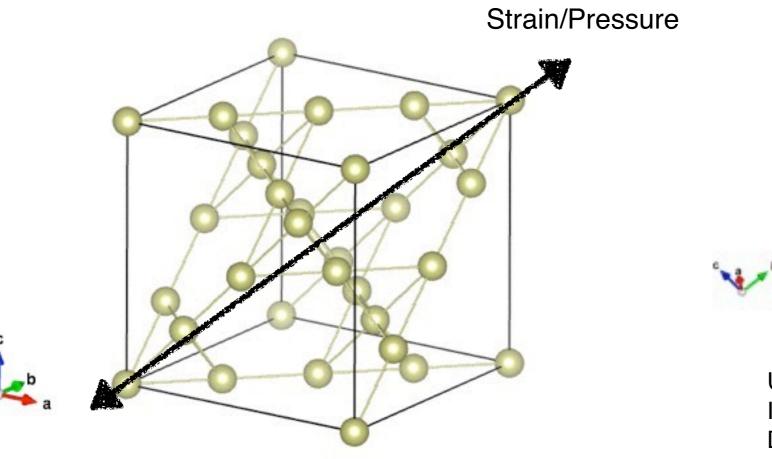
δ>0, opens up a gap at Γ point (topological insulator?)
 δ<0, bands crossing along (III) direction</li>

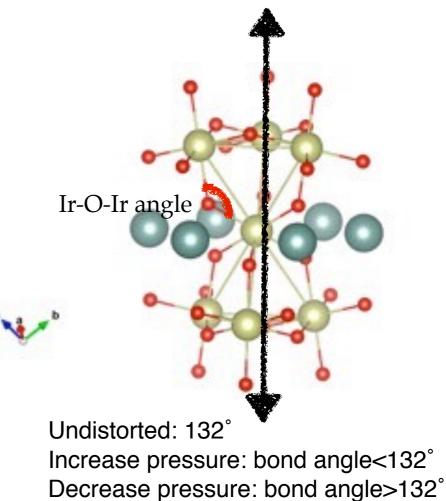
 Several tight-binding studies suggest that pressure would produce topological insulator.



- Ir 5d states are quite extended and a tight-binding model will require many parameters for a good fit. It will be interesting to see full ab initio calculation!
- Ab initio results show that Pr₂Ir₂O7 is closest to quadratic band touching at Fermi energy → it should be very sensitive to strain/pressure!

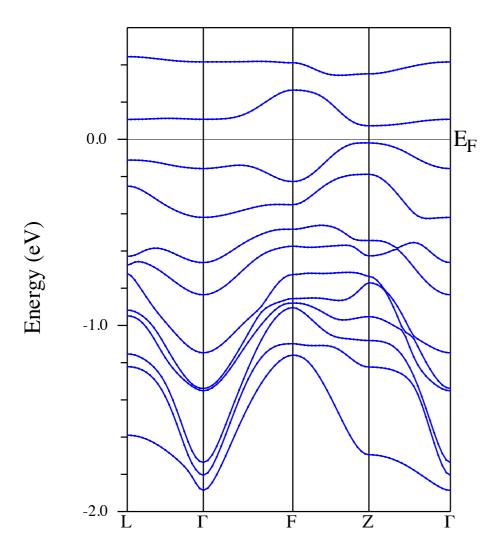
- Uniaxial pressure/Strain along (111) direction
- Cubic -> Rhombohedral symmetry
- Full lattice relaxation both on the lattice parameter and the atom positions.





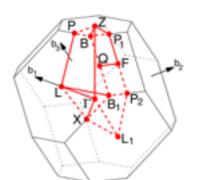
- Pressure/Strain along (111) direction
- Cubic -> Rhombohedral symmetry
- Compressive pressure (in-plane tensile strain)

c=0.96c<sub>0</sub>, a=1.05a<sub>0</sub>



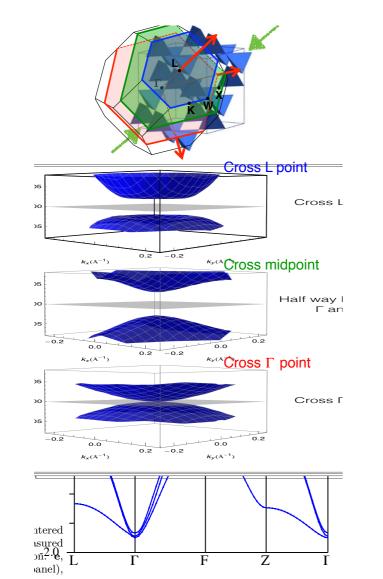
Full lattice relaxation: GGA

both GGA+SO and mBJ+SO calculation Insulator mBJ+SO gap=0.09eV



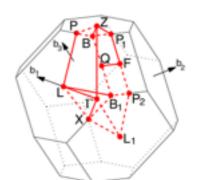
- Pressure/Strain along (111) direction
- Cubic -> Rhombohedral symmetry
- Compressive pressure (in-plane tensile strain)

 $c=0.99c_0, a=1.01a_0$ 



Full lattice relaxation: GGA

both GGA+SO and mBJ+SO calculation Insulator mBJ+SO gap>0eV

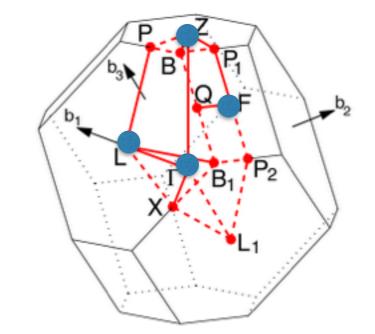


• Compressive pressure (tensile strain)

c=0.96c0, a=1.05a0

 $\delta_a = \prod_m \xi_m(\Lambda_a)$ Parity at time reversal invariant momenta:

Г	L	Z	F
+	-	-	-

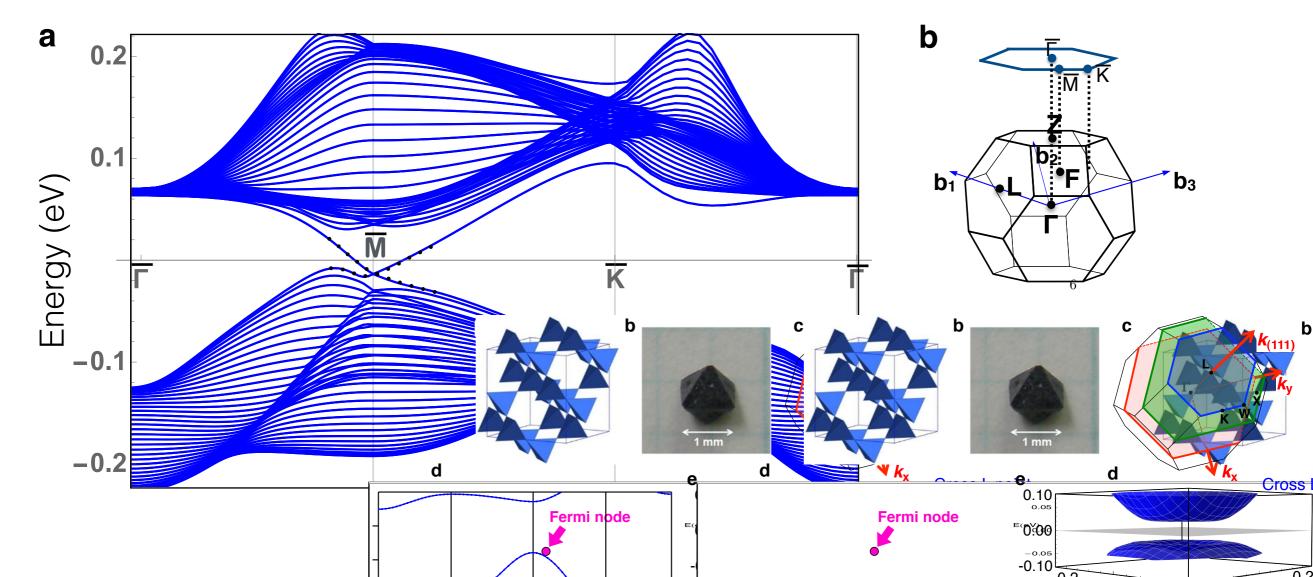


Z<sub>2</sub> invariant (1; 000)

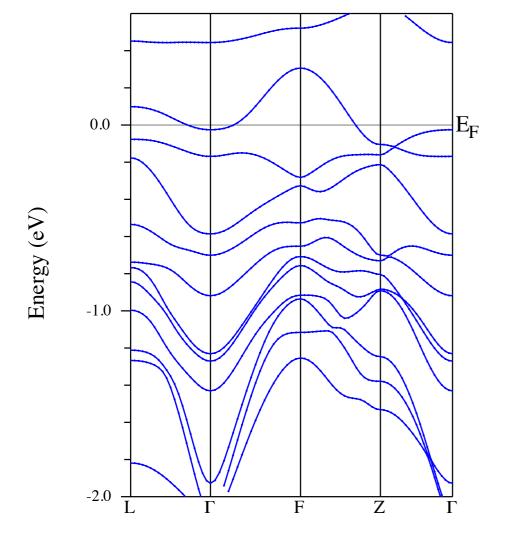
Compressive pressure/ tensile strain along (111) direction will make Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> a topological insulator!

- Compressive pressure (tensile strain)
  - c=0.96c0, a=1.05a0

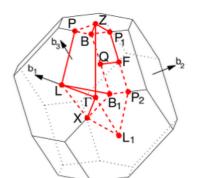




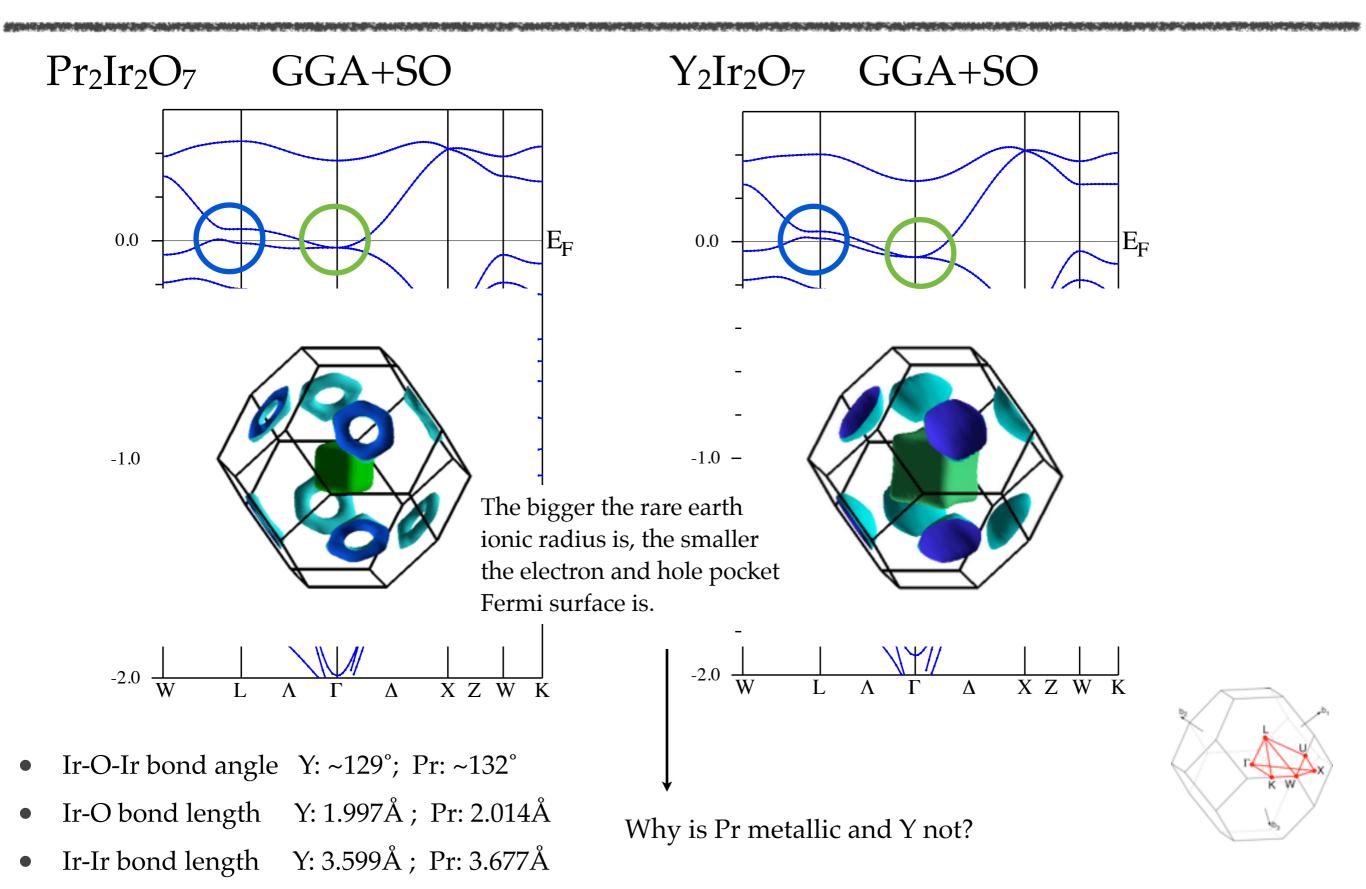
- Reducing pressure (in-plane compressive strain)
  - c=1.06c<sub>0</sub>, a=0.95a<sub>0</sub>



both GGA and mBJ+SO calculation Metallic Band crossing along (111) direction



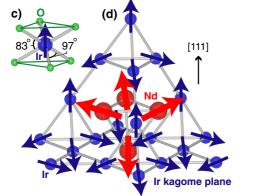
## Paramagnetic calculation: Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> and Y<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>



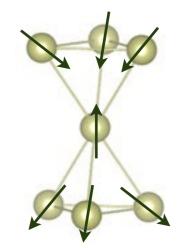
# Magnetic band structure calculation

- Focus on Pr, Nd, Y
- GGA+SO+U, U=2eV since 5d are more extended
- Magnetism in R 4f electrons and in Ir 5d electrons → Magnetic ground state can be complicated!
- non-collinear magnetism :Wien2k+WIENNCM

 $Nd_2Ir_2O_7$ 

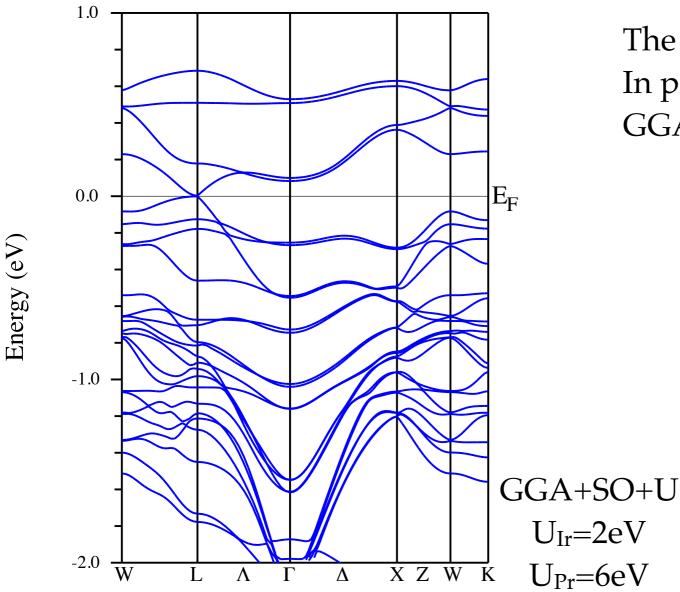


Neutron diffraction and inelastic scattering measurement  $\mathbf{q}_0 = (0,0,0)$  magnetic state: all-in all-out state

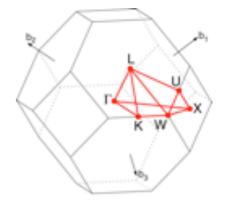


# *Pr2Ir2O7*

- Metallic/paramagnetic down to 0.3K
- Experimental: Pr 4f: 2-in 2-out Ir: paramagnetic
- Try both Pr and Ir all-in all-out order

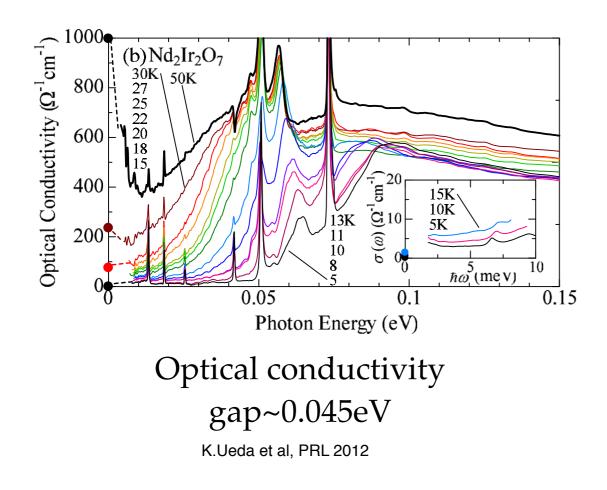


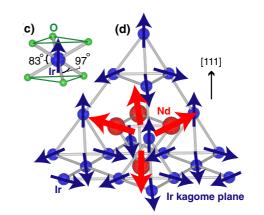
The gap barely closes at the L point In proximity to Weyl semimetal GGA/LDA+U may overestimate magnetism



# $Nd_2Ir_2O_7$

- T<sub>MI</sub>~33K, optical conductivity measurement infers the full opening of the charge gap.
- Experiment suggesting both Nd and Ir are all-in all-out



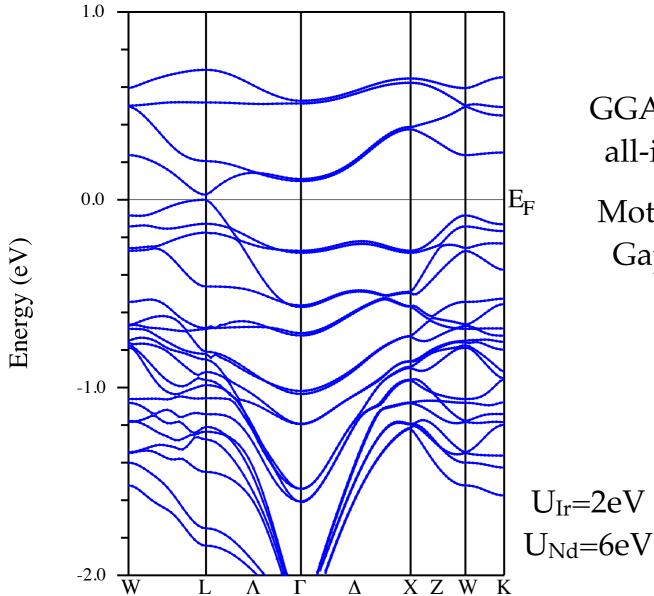


Neutron diffraction and inelastic scattering measurement  $\mathbf{q}_0 = (0,0,0)$ 

Keisuke Tomiyasu et al, JPSJ 2012

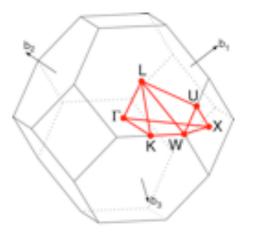
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- T<sub>MI</sub>~33K, optical conductivity measurement infers the full opening of the charge gap.
- Experiment suggesting both Nd and Ir are all-in all-out



GGA+SO+U calculation all-in all-out magnetic

Mott insulator Gap~0.03eV



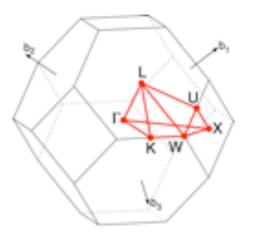
# $Y_2 Ir_2 O_7$

- T<sub>MI</sub>~150K
- Y: no f electron
- Theory and experiment: all-in all-out state X. Wan et al, PRB 2011

1.0 E<sub>F</sub> 0.0 Energy (eV) -1.0 -2.0 W XZWK Λ Г Δ L

GGA+SO+U all-in all-out magnetic Mott insulator Gap~0.12eV

U<sub>Ir</sub>=2eV



# Summary for Part 1

- Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>: Quadratic Fermi node
- Strain/pressure along (111) direction make Pr<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub> a strong topological insulator with correlation.
- Systematic gap decrease trend as rare earth ionic radius increases.









# Part 2 Domain walls in the insulating state of Pyrochlore Iridates

UCB & LBNL Ashvin Vishwanath Jeffrey Neaton

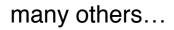


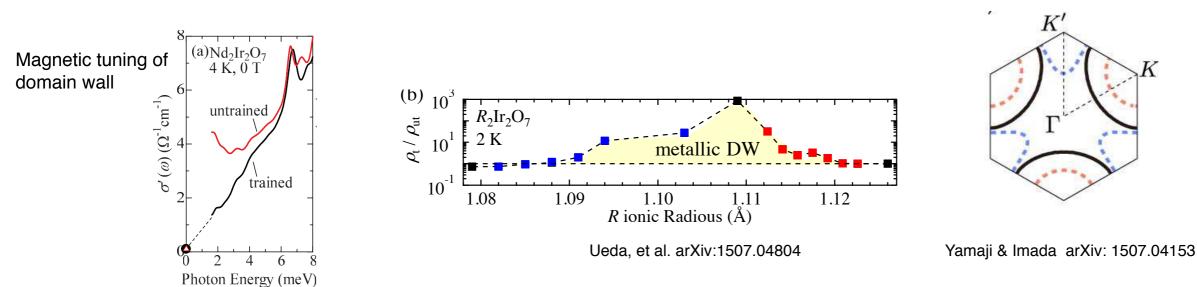


Conducting metallic domain wall even at the gapped bulk state!

Theory and experiments:

Yamaji & Imada Phys. Rev. X 2014, arXiv: 1507.04153 Ueda, Fujioka, Nagaosa, et al. arXiv: 1506.07336, arXiv:1507.04804 Fujita et al. Scientific Reports 2015

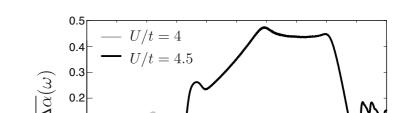




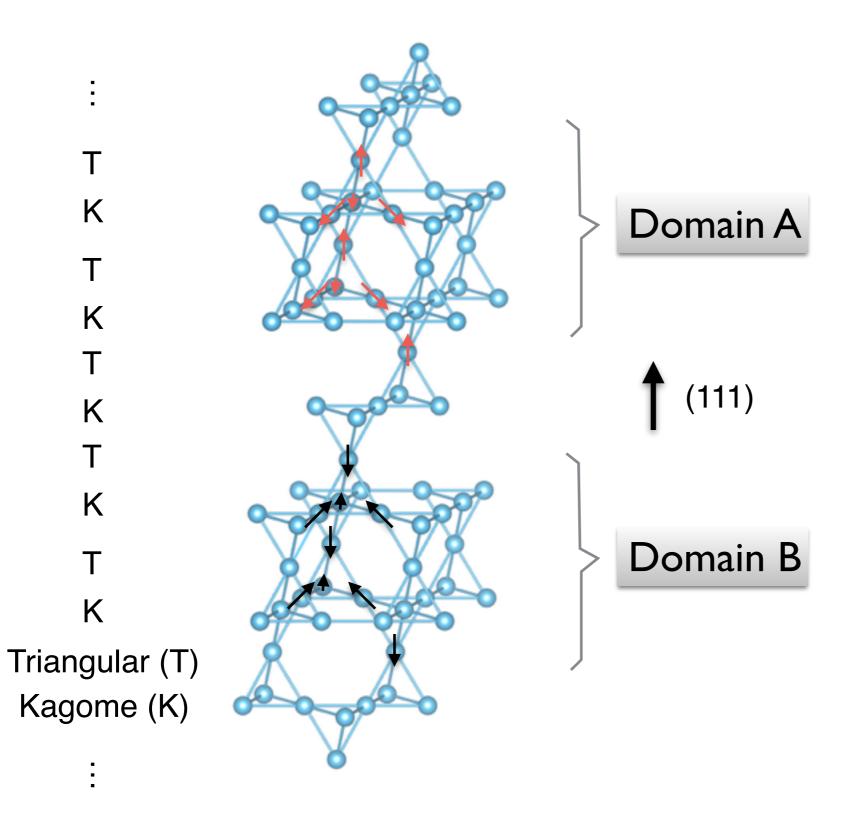
K. Ueda et al., Phys. Rev. B 89, 075127 (2014)

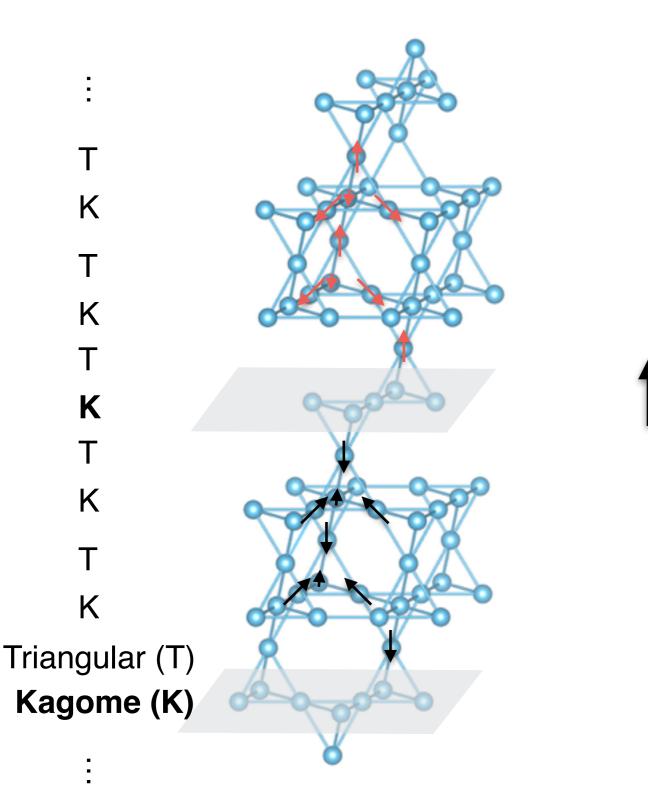
A full ab initio calculation is desirable.



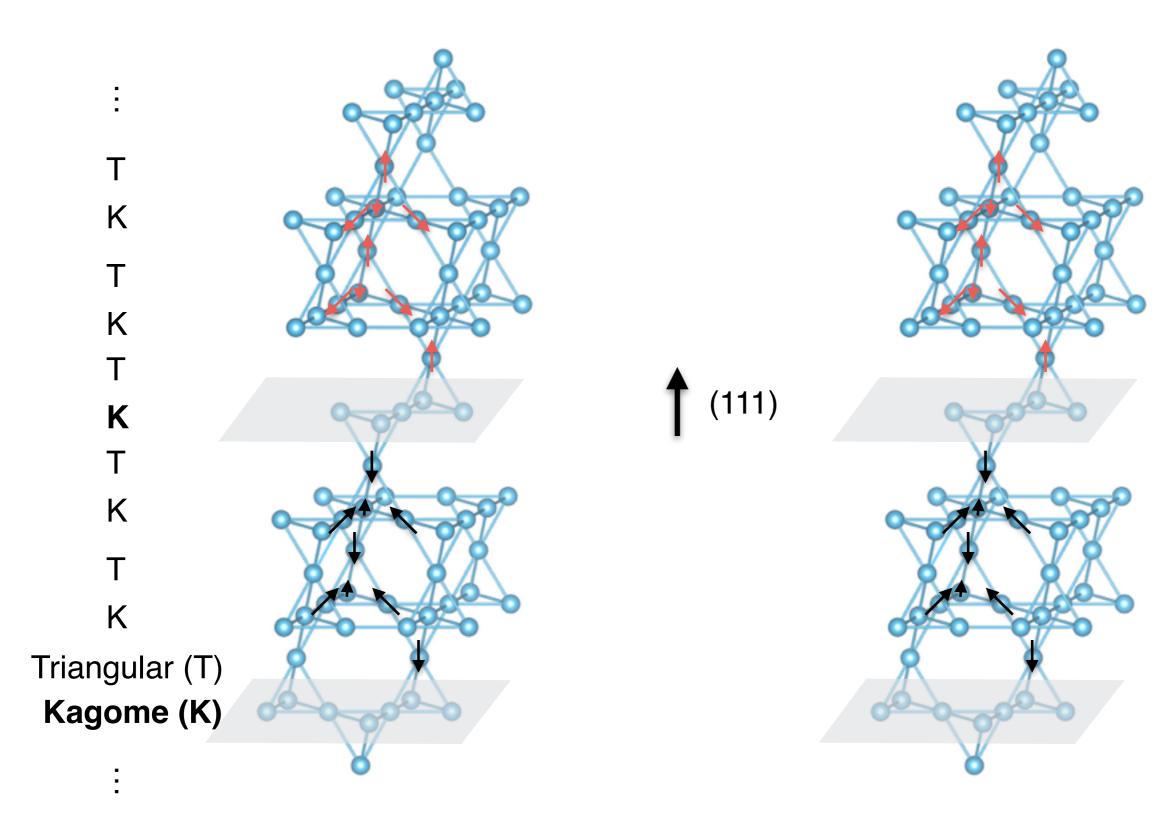


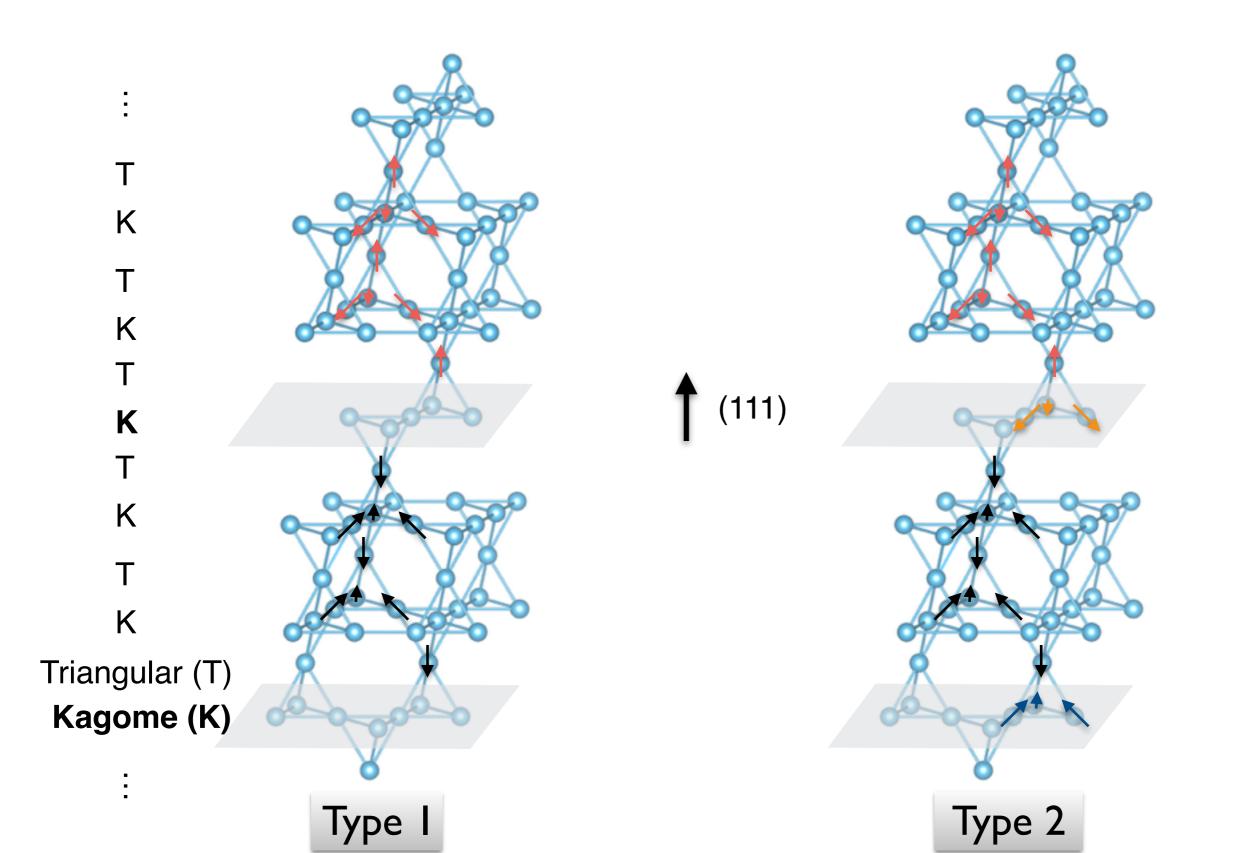
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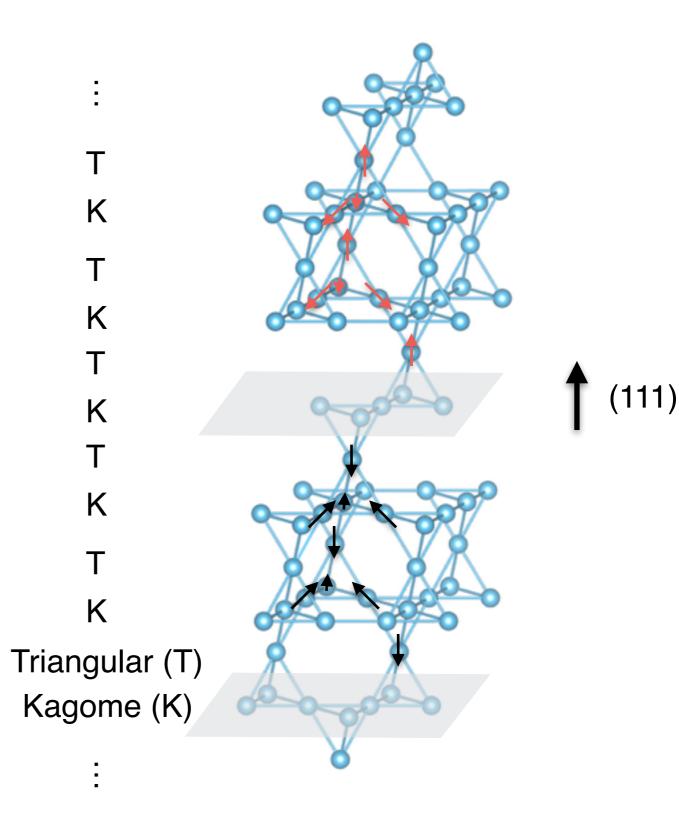
(111)





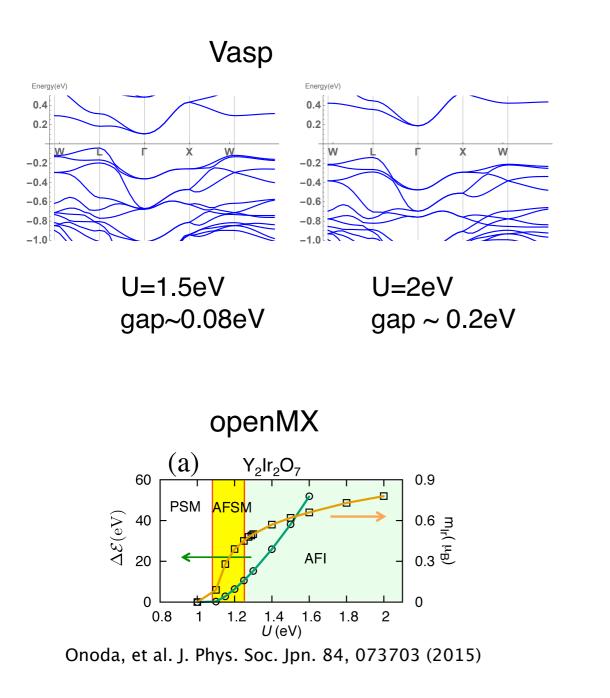
#### Method

- Vasp
- GGA+U+SOC
- Periodic domain wall condition
- Magnetic moment is fully relaxed.
- Convergence test:
  number of k points
  How big is domain region
- Type 1 and 2 domain wall as local minima, different initial magnetic moments.

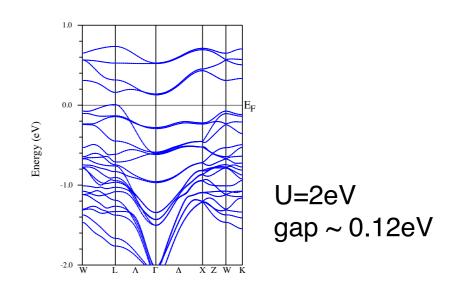


Method:

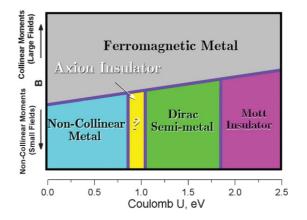
Band structure at different U deviates a little from different program. Overall qualitatively agrees.



Wien2k



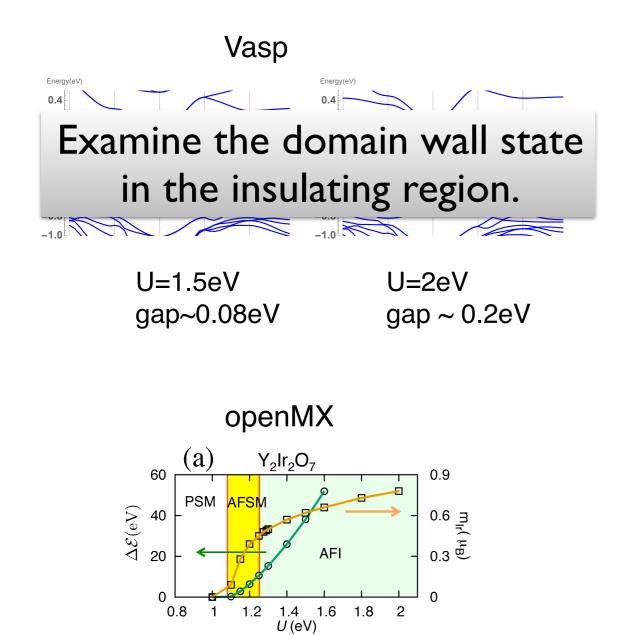
LmtART



X. Wan et al, PRB 2011

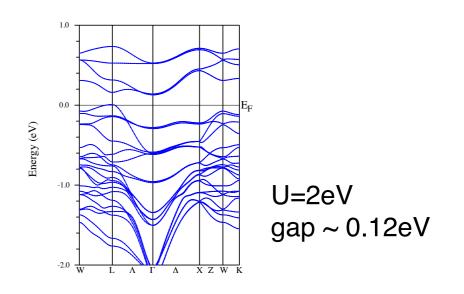
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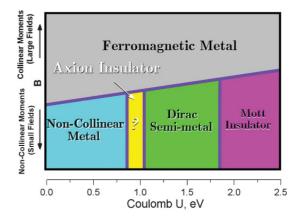


Onoda, et al. J. Phys. Soc. Jpn. 84, 073703 (2015)

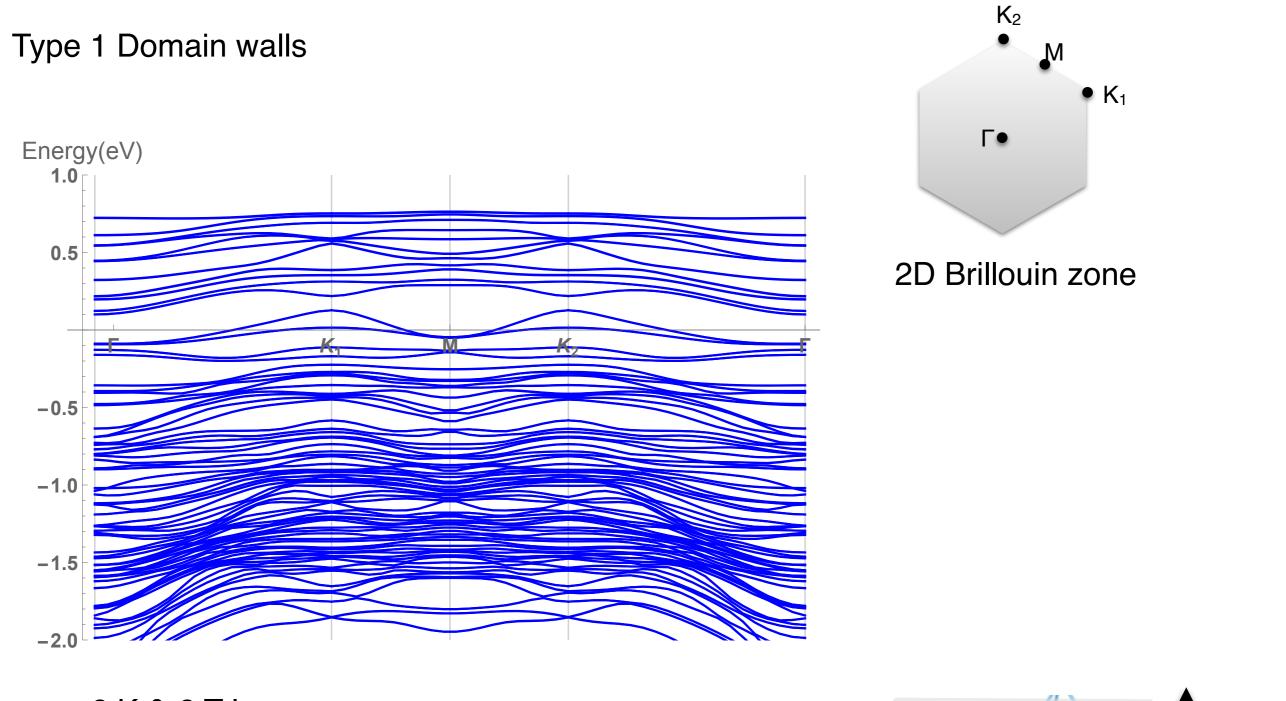
#### Wien2k



LmtART



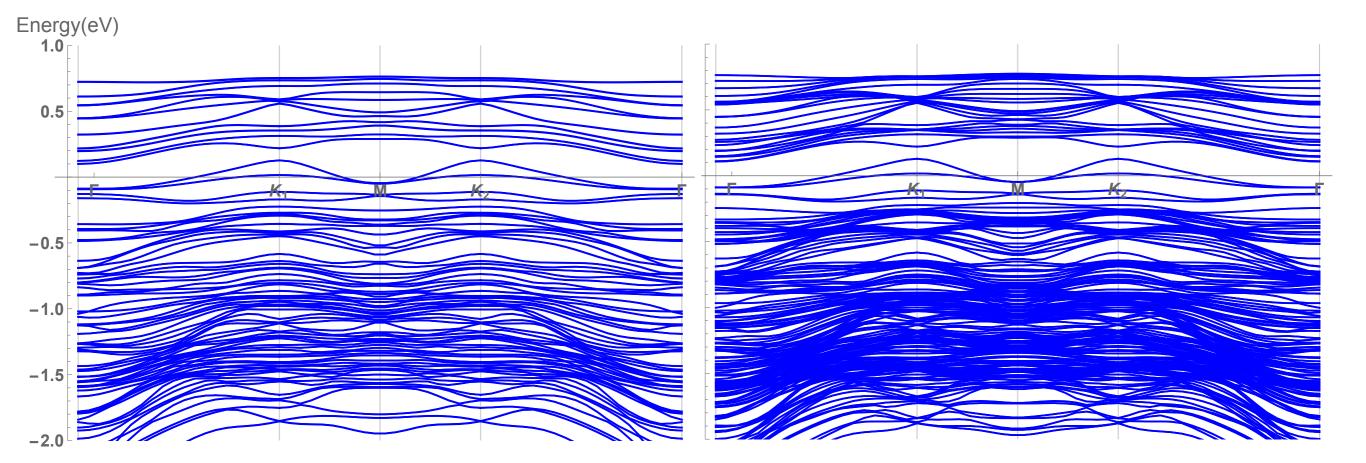
X. Wan et al, PRB 2011



6 K & 6 T layers GGA+SO+U Ulr=2eV





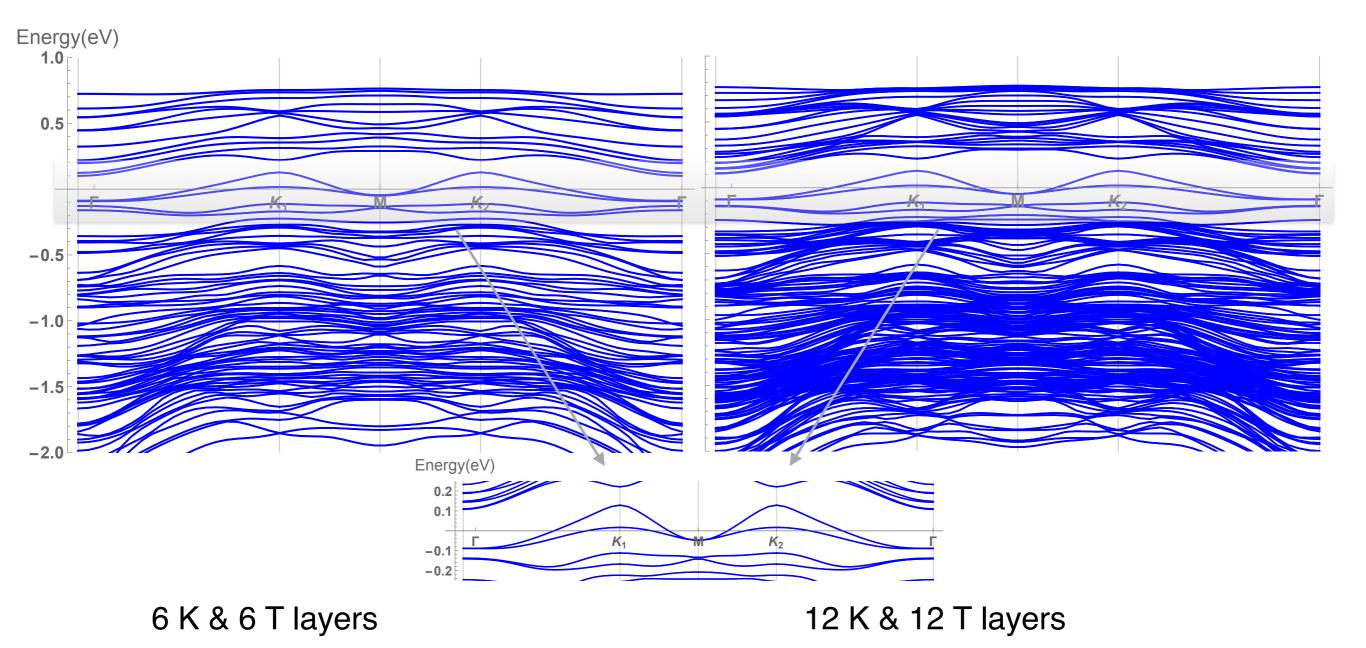


#### 6 K & 6 T layers

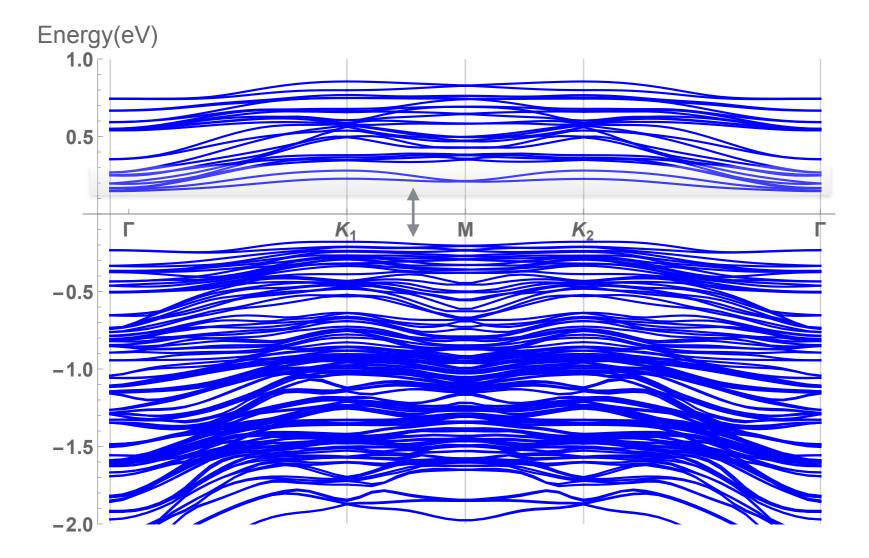
12 K & 12 T layers

Is 6K & 6T superlattice large enough?

Yes for both type 1 and 2 domain walls.



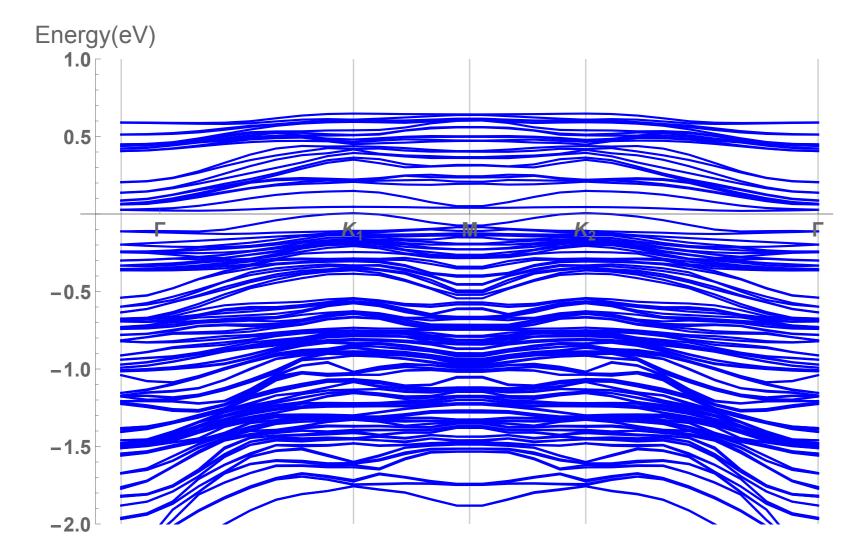
Type 2 Domain walls



Gapped out domain wall states

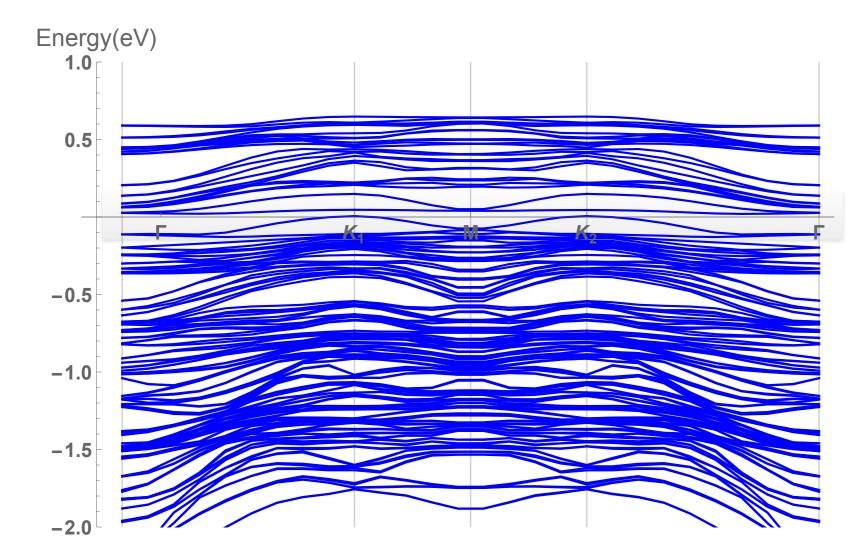
6 K & 6 T layers GGA+SO+U Ulr=2eV

Type 2 Domain walls



6 K & 6 T layers GGA+SO+U U<sub>Ir</sub>=1.5eV

Type 2 Domain walls



Metallic magnetic domain wall states

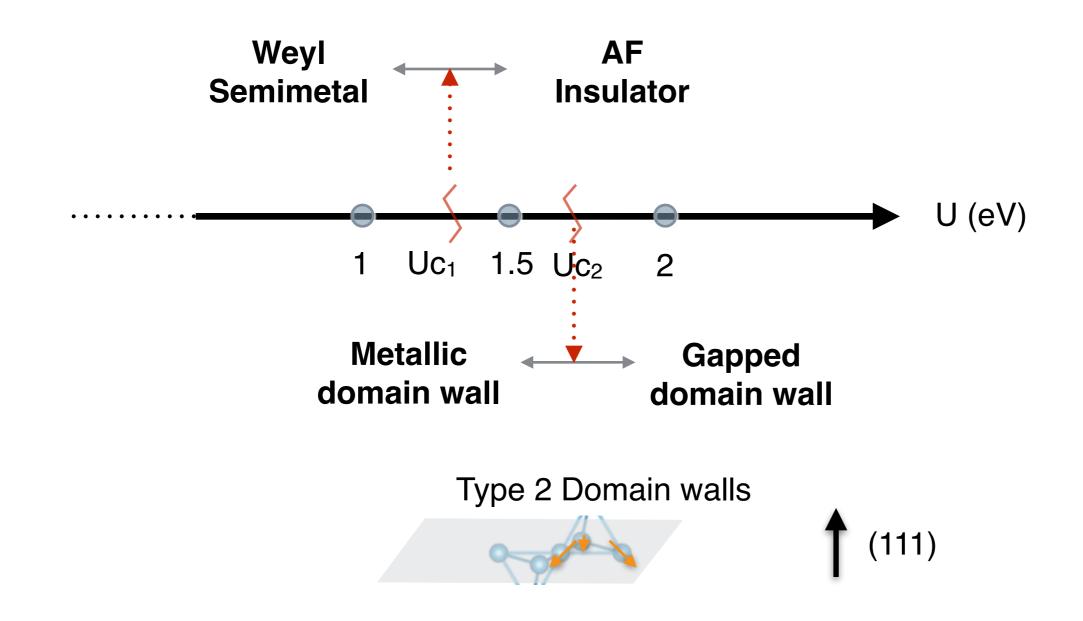
(111)

6 K & 6 T layers GGA+SO+U U<sub>Ir</sub>=1.5eV

Total energy per domain wall (2D Kagome plane)	U = I.5 eV	U = 2 eV
Type 1 Domain walls	0.05eV	0.25eV
Type 2 Domain walls	0	0

Total energy per domain wall (2D Kagome plane)	U = 1.5 eV	U = 2 eV
Type 1 Domain walls	0.05eV	0.25eV
Type 2 Domain walls energetically favorable!	0	0

Tentative DFT phase diagram



# Summary for Part 2

- From ab initio calculation, type 2 domain wall (magnetic) is energetically favored
- Within the AF insulating phase of Y<sub>2</sub>Ir<sub>2</sub>O<sub>7</sub>, there're two phases, one with conducting domain wall, with stronger correlation, the conducting domain wall is gapped out.





