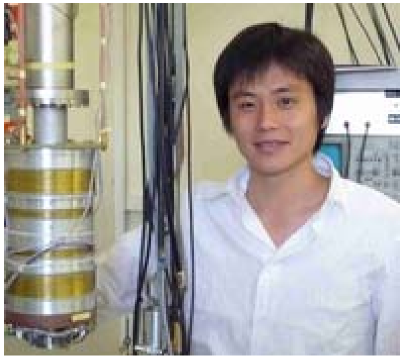
NMR study of the pressure induced Mott transition to superconductivity in the two phases of Cs_3C_{60}

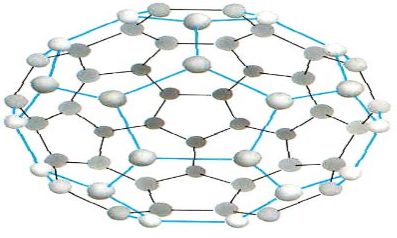


Y. Ihara H. Alloul P. Wzietek

Laboratoire de Physique des Solides, Université Paris XI, Orsay, France.

Daniele Pontiroli, Marcello Mazzani, Mauro Riccò.

CNISM and Dipartimento di Fisica, Universita di Parma



V.Brouet

V. Brouet, H. Alloul *et al*
PRL, ***82***,*2131 (1999);* ***86***,*4680(2001);*
PR ***B***, ***66***, *155122 , 15123, 15124(2002).* .

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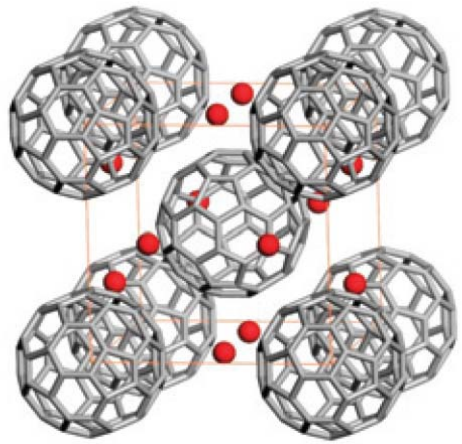
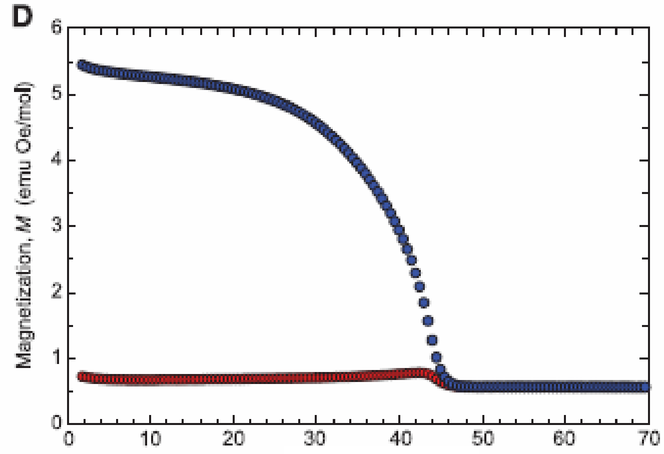
LPS ORSAY

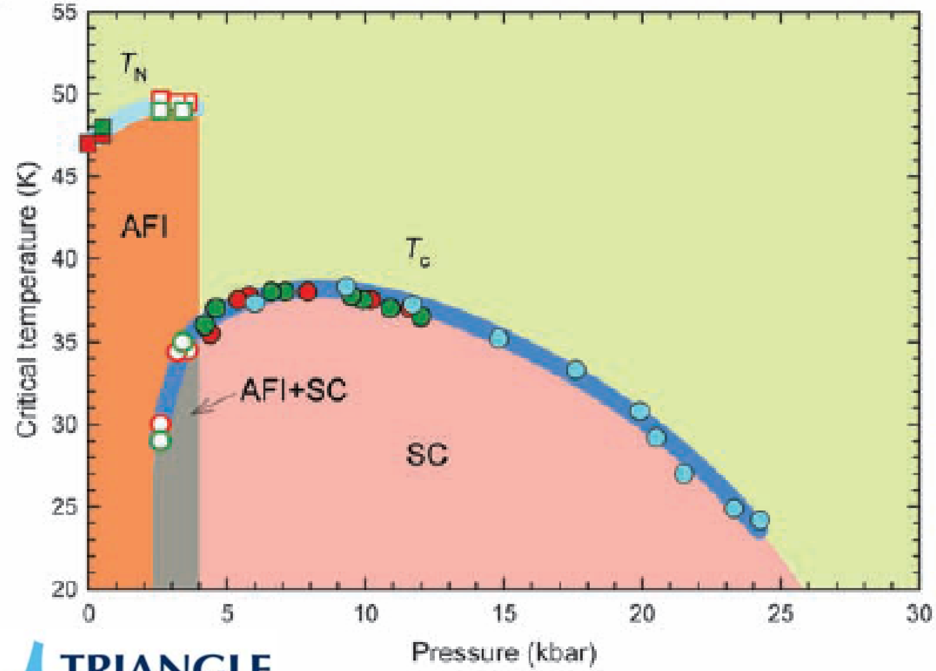
A15 Cs_3C_{60} : magnetic at ambiant pressure



M



T



SC induced by pressure
Proximity of magnetism and SC.

Palstra *et al.* Sol. Stat. Commun. **93** 327 (1995).

A Ganin et al Nature Materials, April 2008

Takabayashi et al , Science 2009

NMR study of the pressure induced Mott transition to superconductivity in the two phases of Cs_3C_{60}

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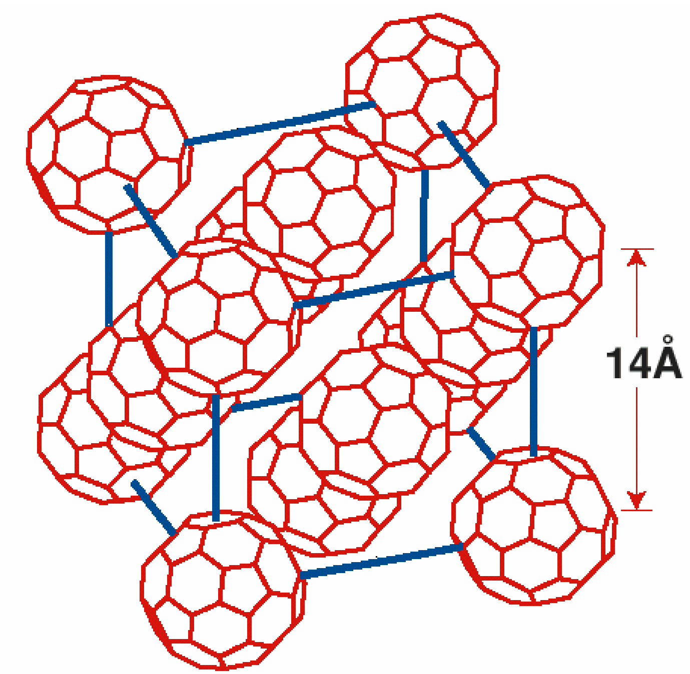
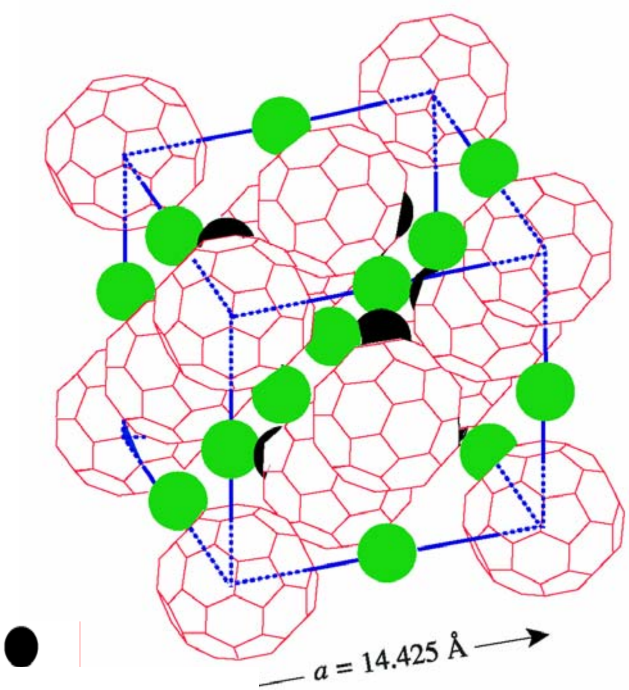
sc and fcc phases of fullerides



Solid C_{60}



tetrahedral A site

octahedral A site

A_3C_{60}
(A= K , Rb, Cs)

High T_c superconductors

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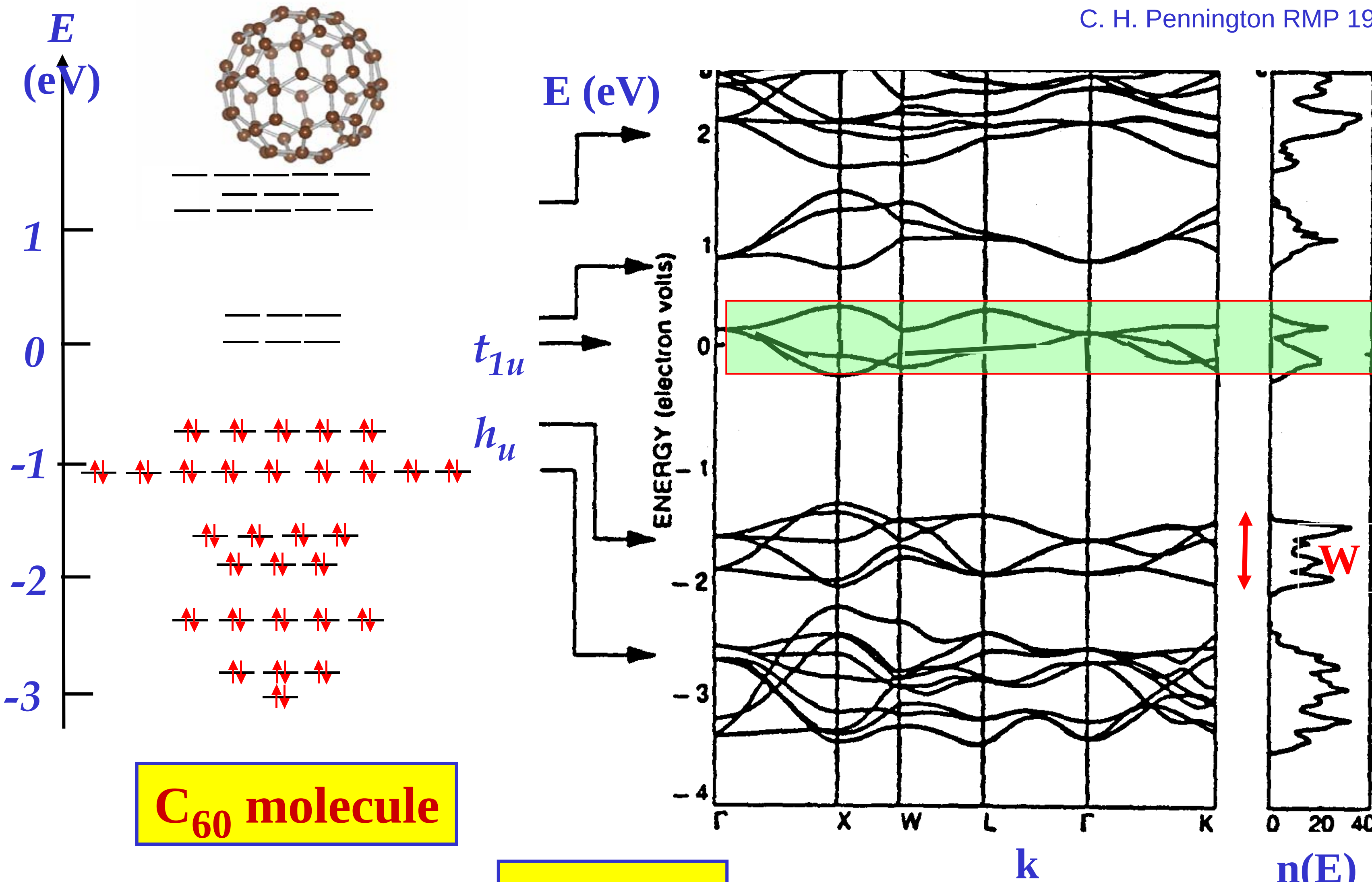
Electronic structure of doped fullerides

C. H. Pennington RMP 1995.

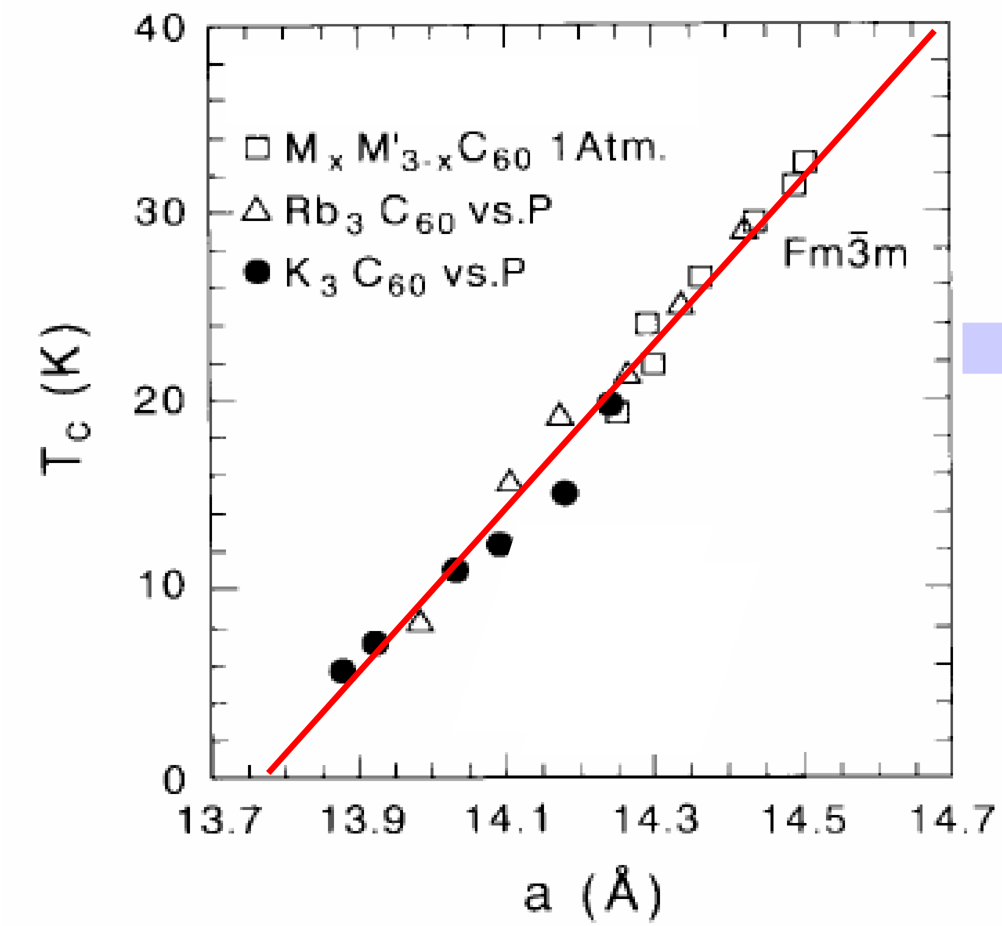


C_{60} molecule

C_{60} solid

W= 0.5 eV

Superconductivity in A_3C_{60}



$$k_B T_c = 1.14 \hbar\omega_D \exp\left[-\frac{1}{V_0\, n(E_F)}\right]$$

T_c depends only on $n(E_F)$

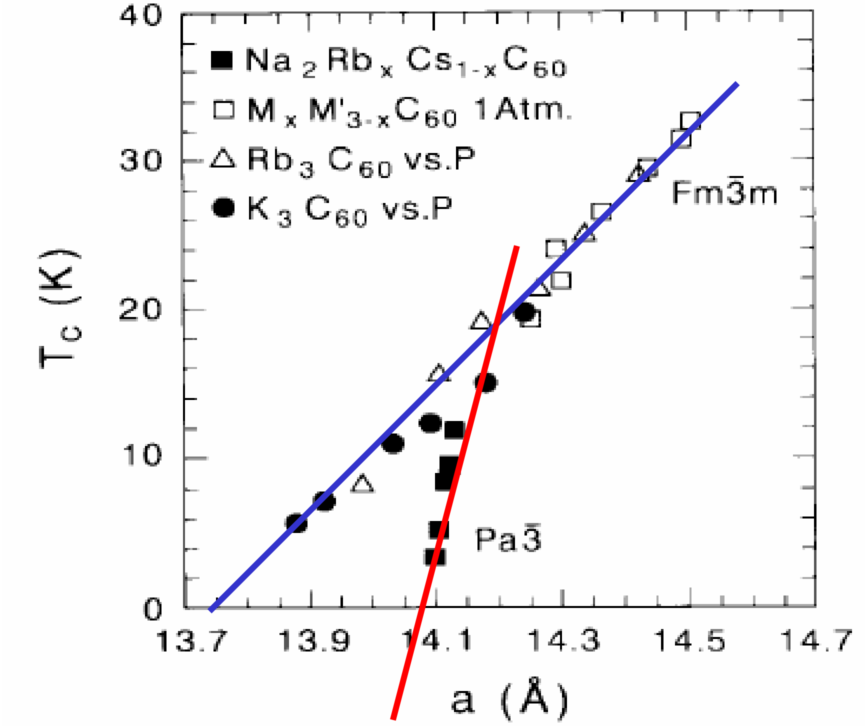
$\hbar\omega_D$ and V_0

are molecular properties

Only C_{60} phonons are involved

The scaling between T_c and the lattice parameter is assumed to support a BCS-like mechanism

Difference between Na_2AC_{60} and the other A_3C_{60}



Yildirim *et al*. MRS proceedings **359** 273 (1995).

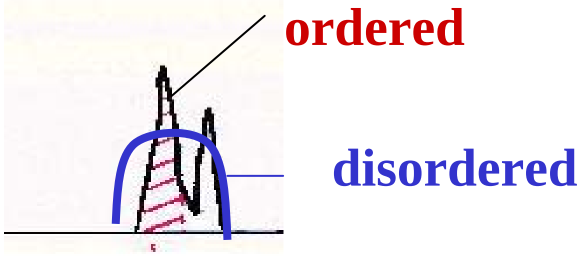
Na_2AC_{60} sc Pa3 structure

other A_3C_{60} are fcc

-Different slope for sc phases ?

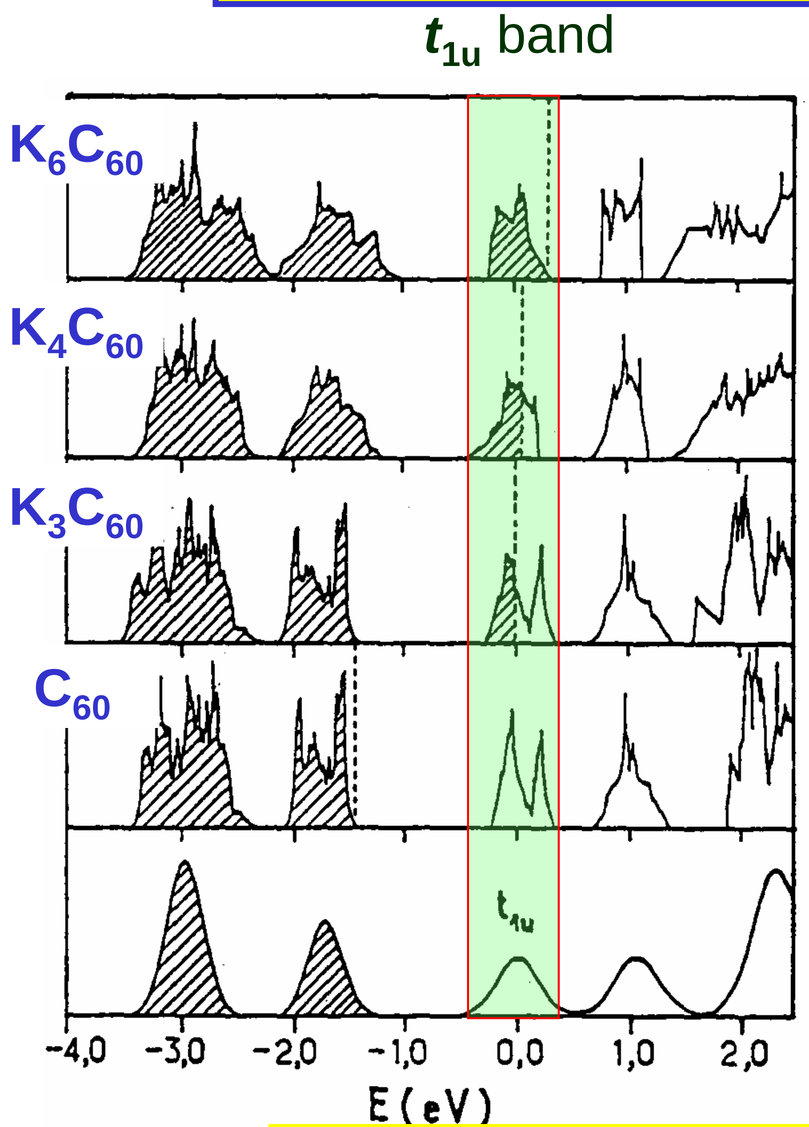
- Different variation of $n(E_F)$?

- Alkali plays a role after all ?



What happens then for other *n* values ?

The anomalous properties of A_nC_{60} compounds



Electronic doping of triply degenerate t_{1u} orbitals

All compounds with n<6 should be metallic.

But…

Crystalline symmetry depends on doping level.

x = 4 : Insulator (bct)
K_4C_{60}, G. Zimmer EPL 1994, V. Brouet PRB 2002.

***x* = 3 : Superconductor** (fcc, cubic)

x = 2 : Insulator (cubic)
Na_2C_{60}, F. Rachdi PRB 1997, V. Brouet PRL 2001, PRB 2002.

x = 1 : Metal-Insulator transition
polymer-Cs_1C_{60}, V. Brouet PRL 1996, B. Simovic Synth. Met. 1999.
cubic-Cs_1C_{60}, V. Brouet PRL 1999.

Strong correlations and Jahn-Teller distortions (JTD)

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H. Alloul, Santa Barbara, 12/02/2010

LPS ORSAY

NMR study of the pressure induced Mott transition to superconductivity in the two phases of Cs_3C_{60}

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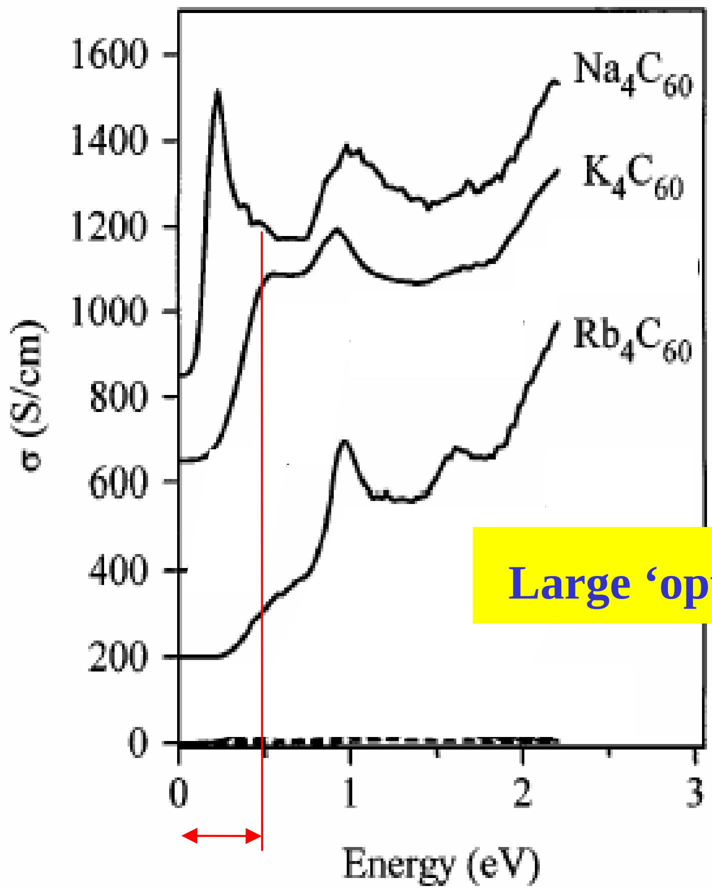
Electronic properties of A_4C_{60} (bct) and Na_2C_{60} (cubic)

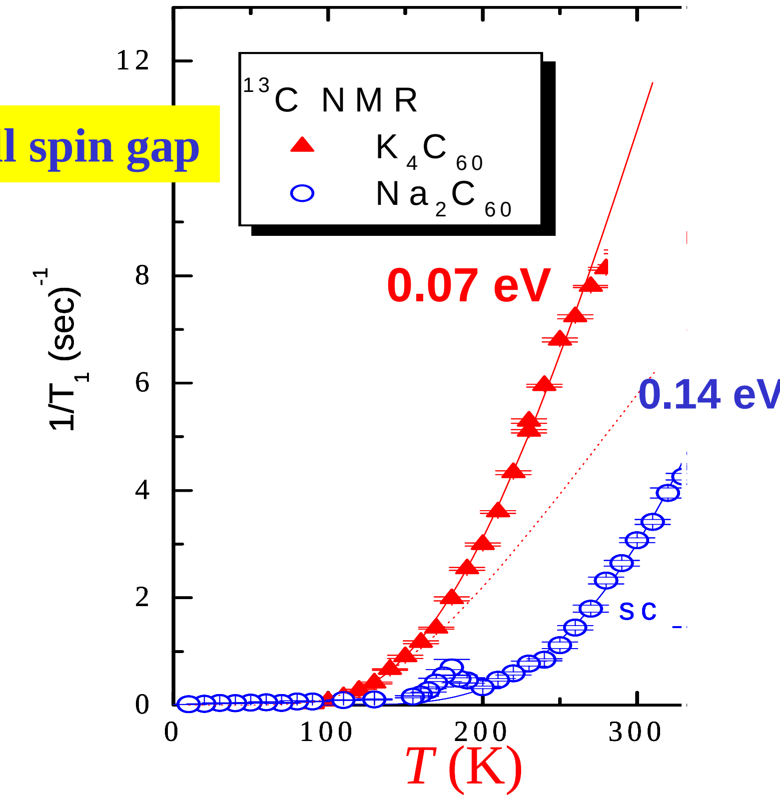
Optical conductivity σ extracted by EELS



0.5eV

M. Knupfer *et al*, PRL **79,** 2714 (1997).

NMR $1/T_1$



K_4C_{60}: G. Zimmer *et al* EPL (1994), V. Brouet *et al* , PRB 66 155122 (2002).Na_2C_{60}, F. Rachdi *et al,* PRB 1997, V. Brouet *et al* PRL 2001. PRB 2002.

Insulating non magnetic ground state

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Stabilization of the singlets

Why singlets? Two electrons on a ball costs an energy ***U***

So there is an energy gain which opposes to *U*

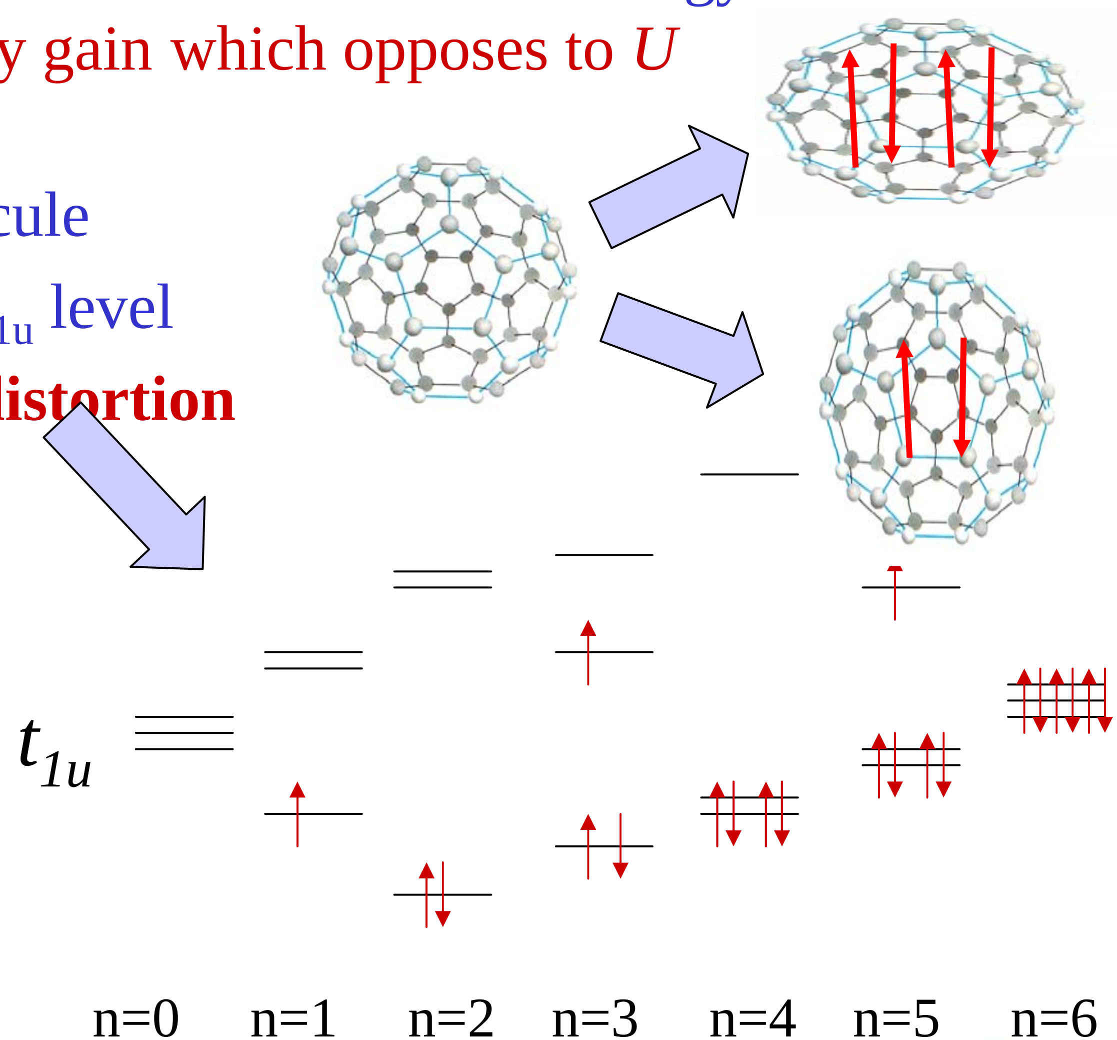
For a charged molecule

The degeneracy of the t_{1u} level

is lifted by a **Jahn Teller distortion**

The Jahn-Teller splitting of the t_{1u} level depends on the C_{60} charge

Manini, Tosatti PRB94



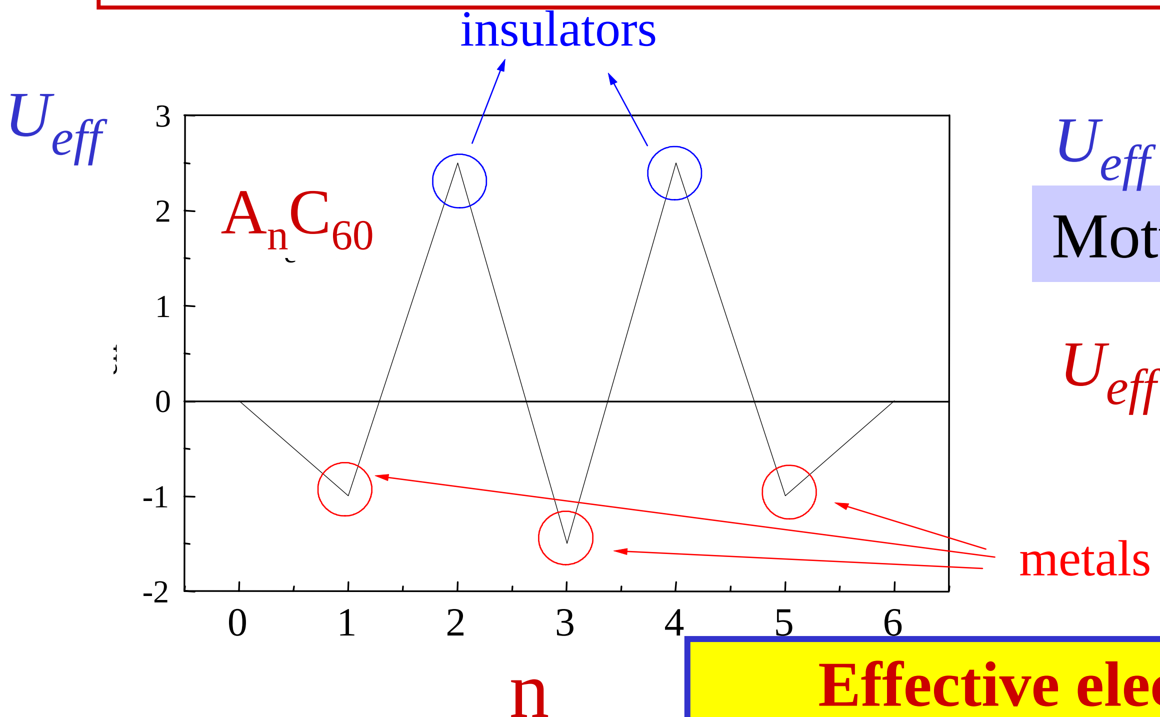
The larger gain per electron is for n=2 or 4

Strong electronic correlations induce the insulating behaviour

In all A_nC_{60} the coulomb repulsion U is large

The Jahn-Teller distortion adds a contribution U_{eff} to U

Yields a S=0 ground state rather than S=1 for n =2 or 4



U_{eff} adds to U for even n

Mott Jahn-Teller insulators

U_{eff} reduces U for odd $\boldsymbol{n}$

M. Héritier, W. Victoroff
O. Gunnarson
M. Fabrizio, E. Tosatti

Effective electronic interactions mediated by Jahn-Teller distortions

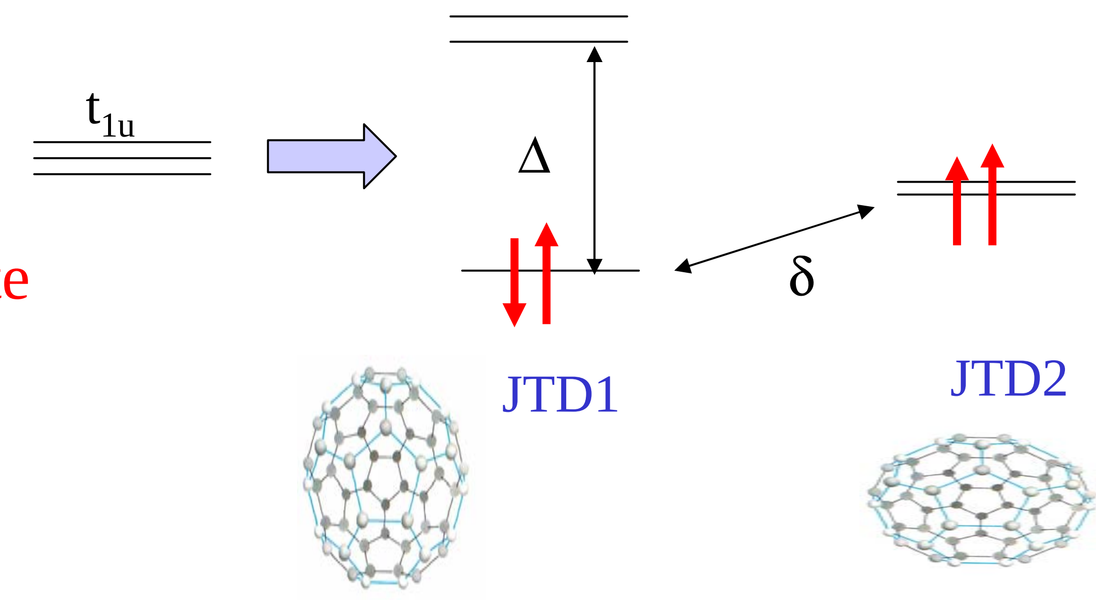
Insulating states of A_4C_{60} and Na_2C_{60}

(Fabrizio, Tosatti PRB97)

Experimental situation

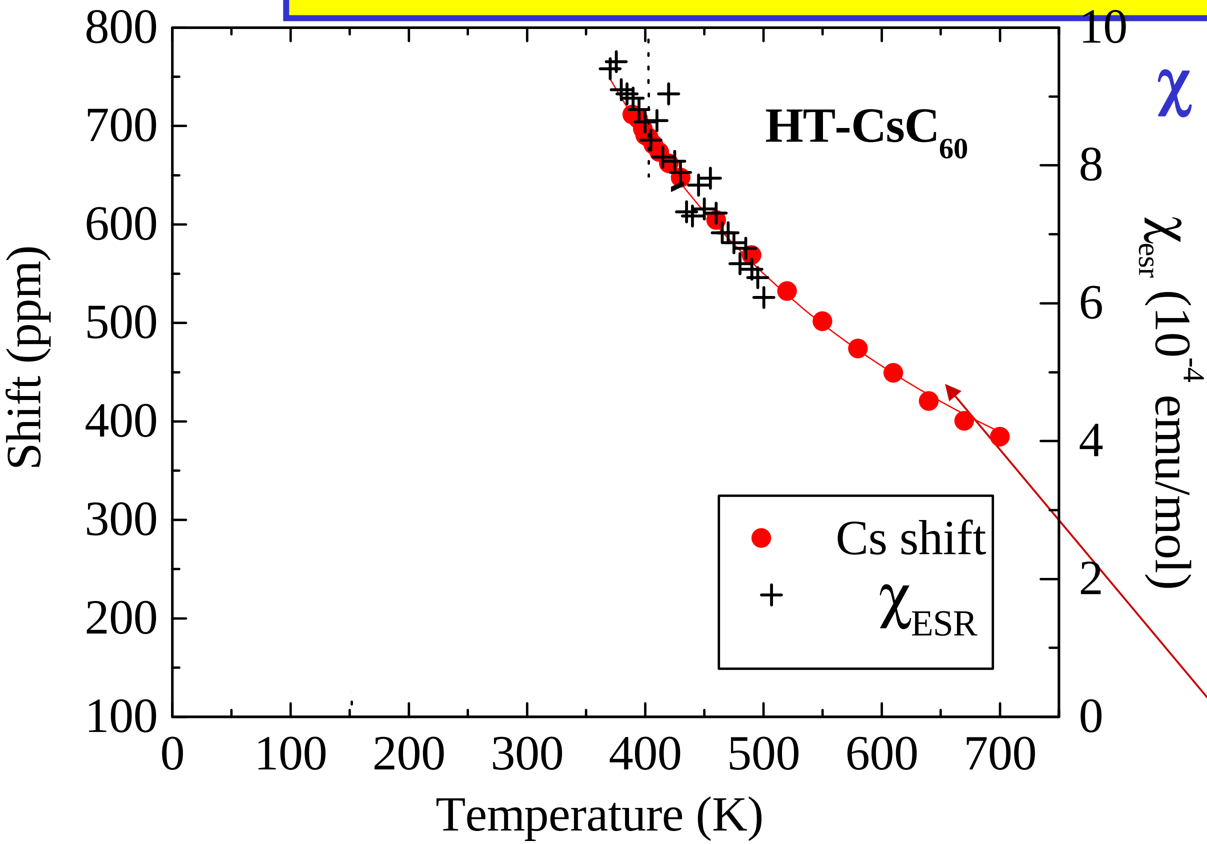
- Non magnetic ground state

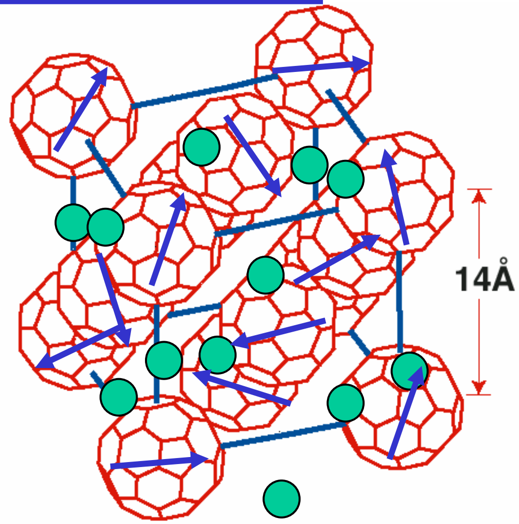
- 2 different gaps



- **Hund's rule is not obeyed** in the fundamental JTD1 state
- **Excited state is triplet** and corresponds to a different JTD2
- NMR detects the **Singlet Triplet excitation δ**

The special case of the cubic phase of CsC_{60}





$$\chi_P = \frac{g^2 \mu_B S(S+1)}{3k_B(T+\theta)}$$

with S=1/2 and θ<5K

Strong correlations for A_3C_{60} as well ?

Both NMR and ESR indicate a paramagnetic Curie law for fcc AC_{60}

So CsC_{60} is a Mott insulator

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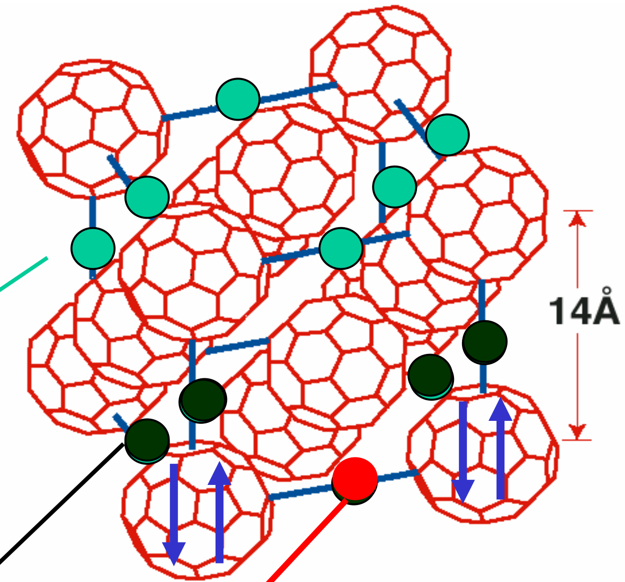
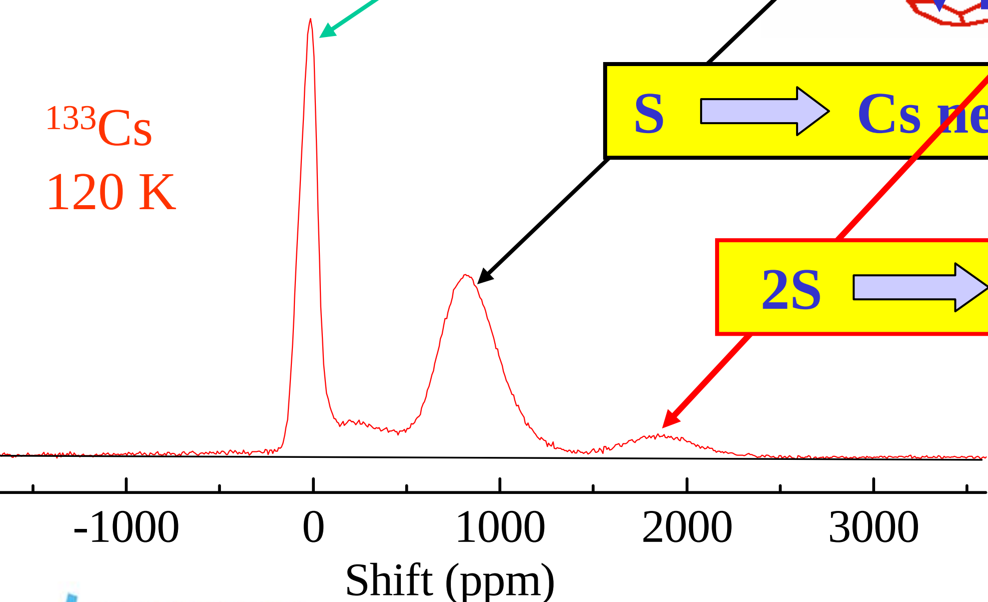
Evidence for charge segregation: singlet states



NS far from singlets^{-}

S ⇨ Cs neighbors of one localized C_{60}^{2-}

2S ⇨ Cs neighbors of two localized C_{60}^{2-}

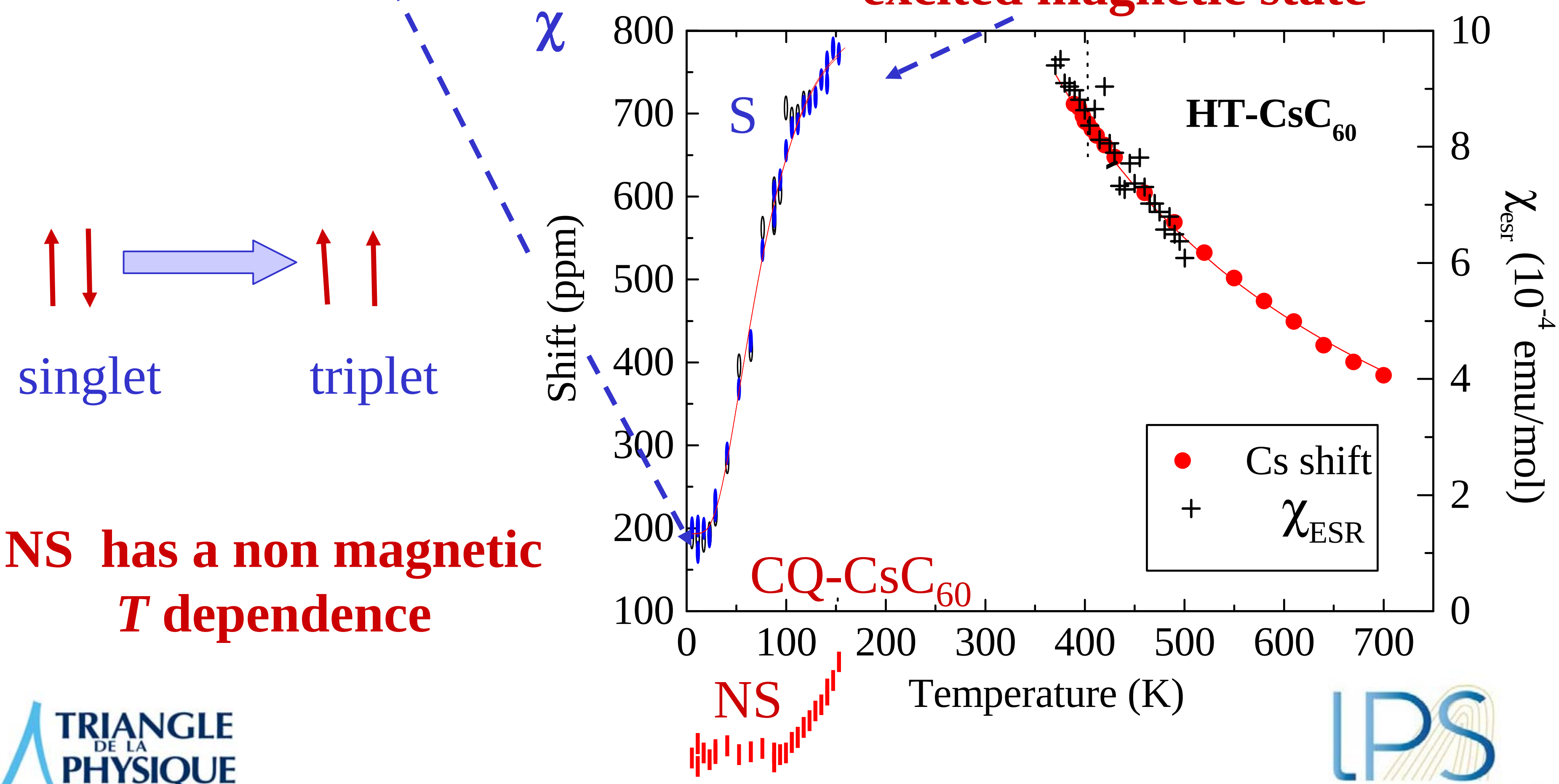


V. Brouet *et al* , PRL 1999.

Comparison of Cs NMR in QC- CsC_{60} and HT- CsC_{60}

Very different from HT cubic phase

S (and 2S) sense a transition from a **local non-magnetic ground state** …… towards an **excited magnetic state**



NS has a non magnetic *T* dependence

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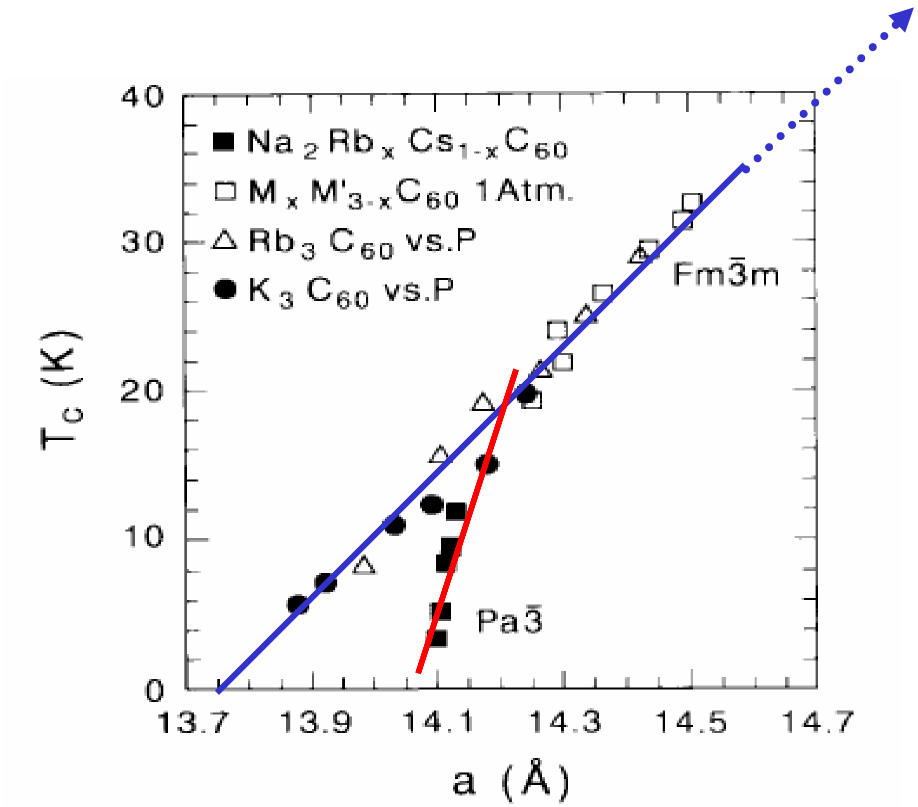
Superconductivity in A_3C_{60}



Higher Tc?

Mott transition?

Alkali Ionic radius

Li
Na
K
Rb
Cs
Fr

Cs_3C_{60}

Expanded Mott state in K_3C_{60} expanded by NH_3 molecules

$(NH_3)_xK_3C_{60}$ a = 14.971 Å (bct)

Rosseinsky Nature **364** 425 (1993).

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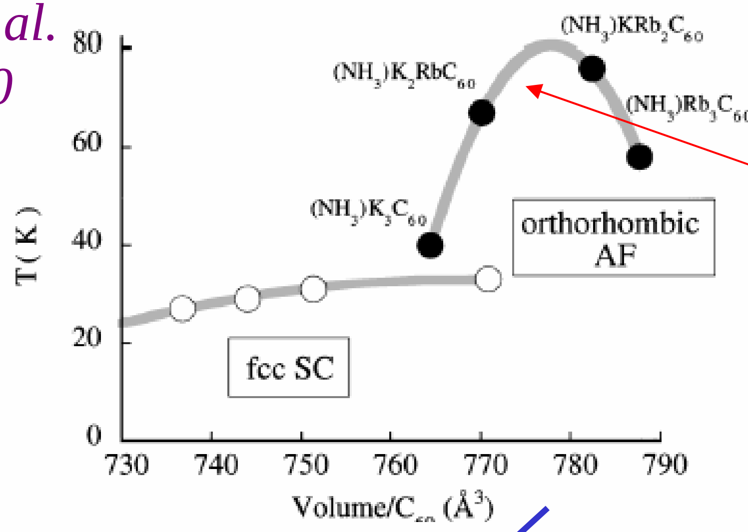
Metal-insulator transition in expanded A_3C_{60}

Iwasa et al. PRL2000

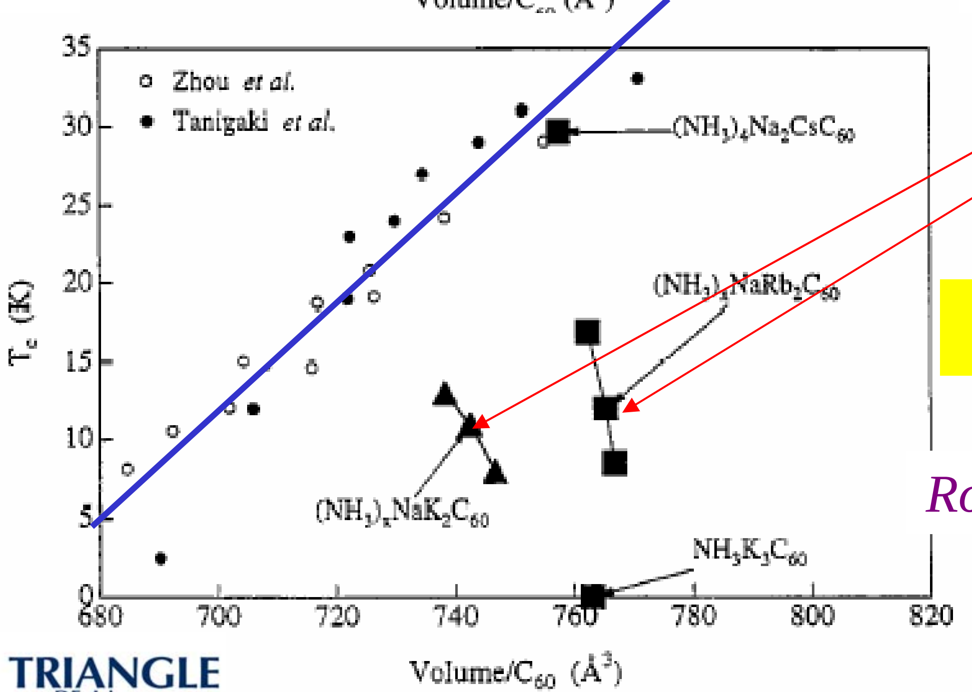


For large lattice spacings:

the system becomes AF.

But:

structural distortions also induce magnetism



Lowering of the degeneracy?

Rosseinsky, Maniwa, Iwasa , Prassides

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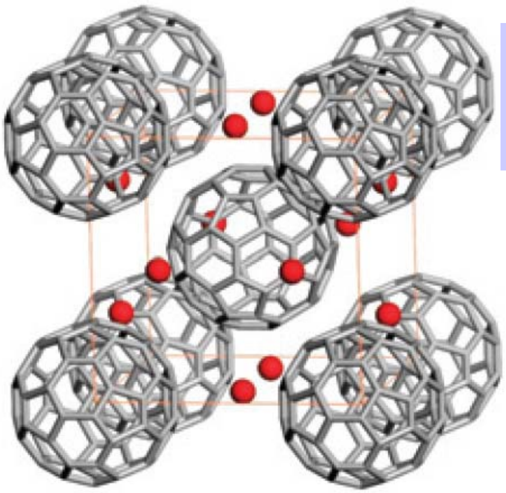
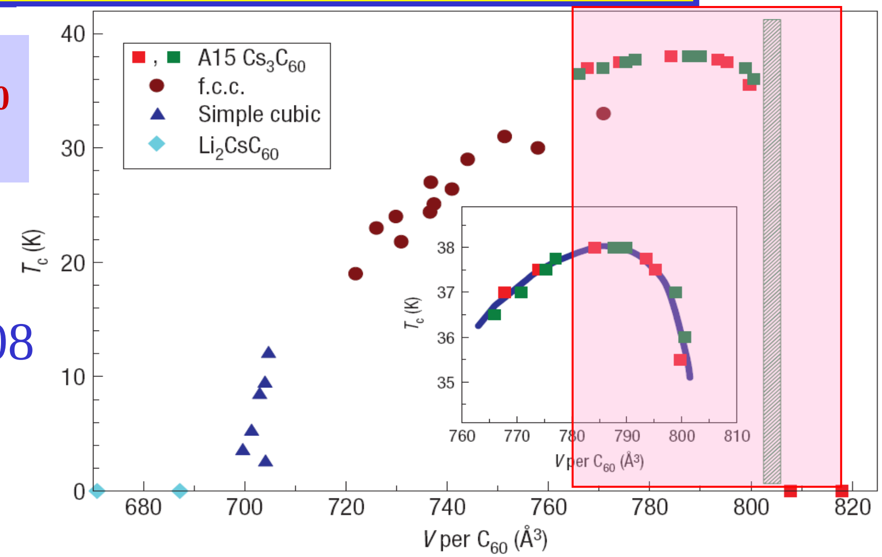
A15 Cs_3C_{60} pressure induced superconductivity

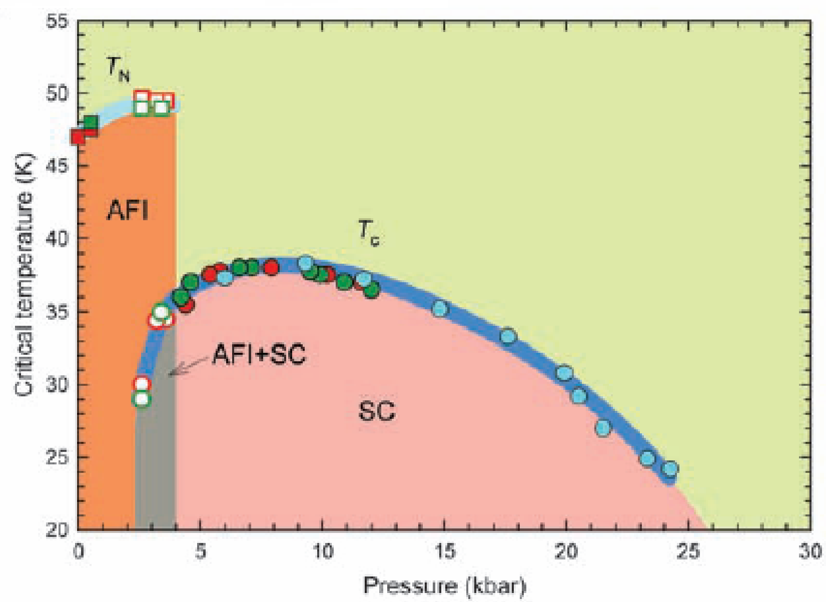


New phase of Cs_3C_{60} with A15 structure

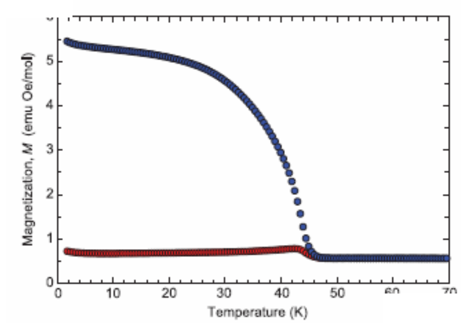
A. Ganin et al
Nature Materials , 2008



M



Takabayashi *et al* , Science 2009

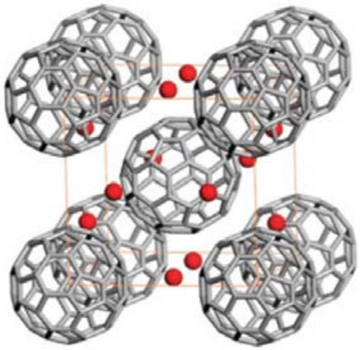


T

Magnetic at ambient pressure

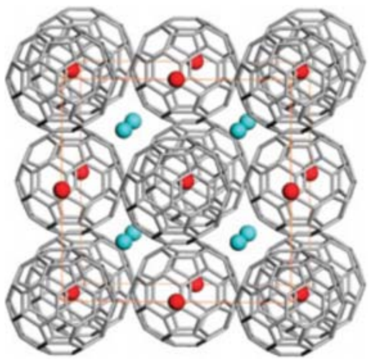
SC and Mott transition

Multiple phases in the samples: ^{133}Cs NMR is very helpful



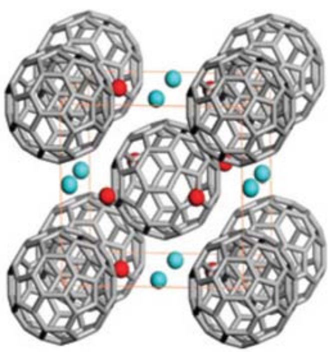
A15 Cs_3C_{60}

A single Cs site
with non cubic
local symmetry



fcc Cs_3C_{60}

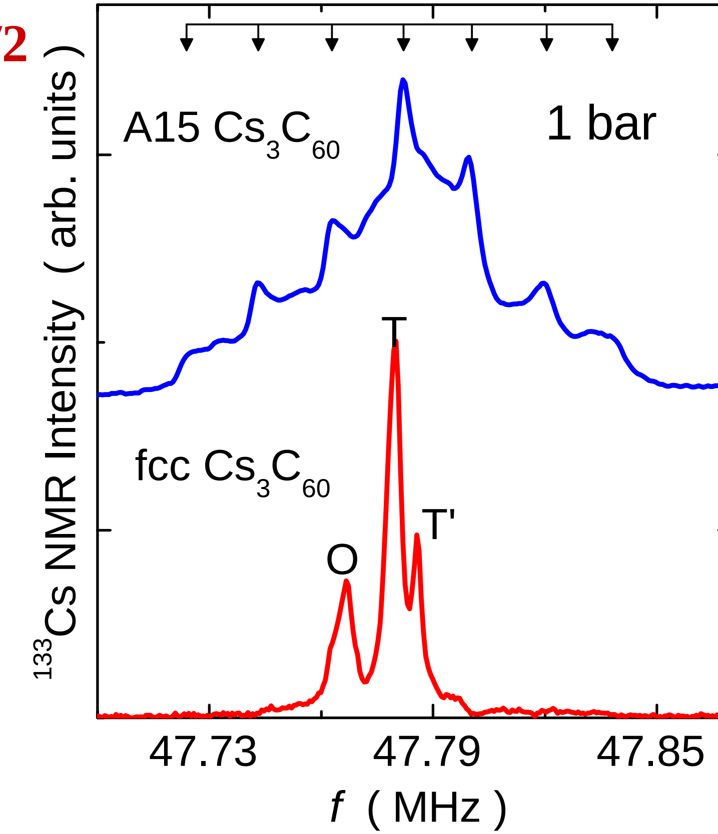
Well known: two Cs sites
O and T = 1 : 2
With cubic local symmetry
T splits at low T
Merohedral disorder of the C_{60}



Cs_4C_{60}

Insulating phase
Eliminated by its very long T_1

^{133}Cs *I*=7/2



A15-rich sample		fcc-ri ch sample	
A15	: **58.4 %**	A15	: 34 %
FCC	: 12 %	FCC	: **55 %**
Cs_4C_{60}	: 29.5 %	Cs_4C_{60}	: 11 %

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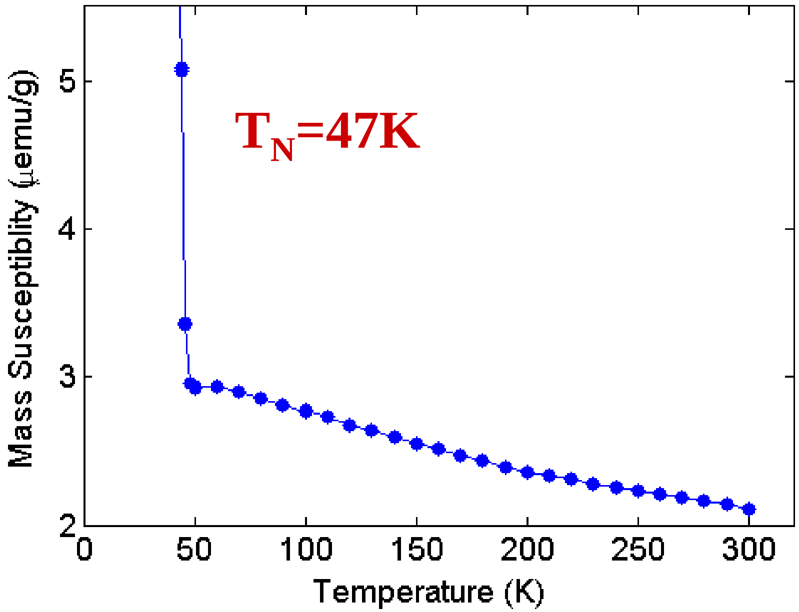
Differences allow selective NMR experiments

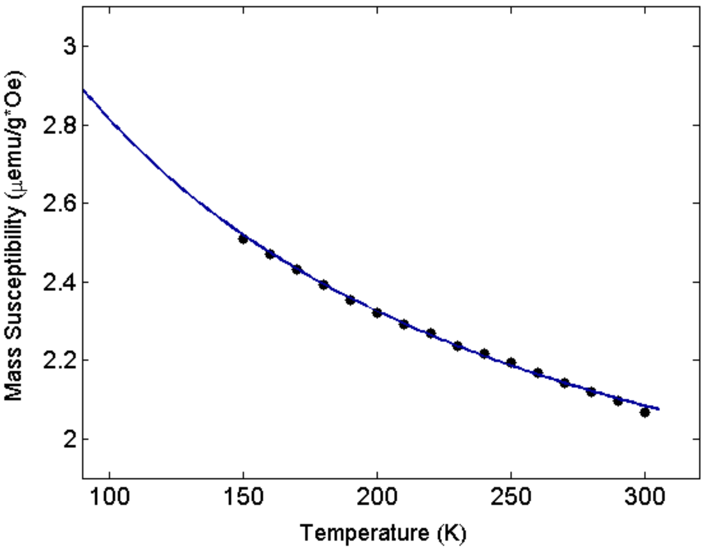
LPS ORSAY

H. Alloul, Santa Barbara, 12/02/2010

Paramagnetic state susceptibility

(30.9% Cs_3C_{60} FCC, 53.6% Cs_3C_{60} A15, 15.5% Cs_4C_{60})





High T Curie-Weiss behavior

Weiss temperature: $\theta \sim 100$ K

Effective moment : $p_{eff} \sim 1.70\ \mu_B$

$$\chi^{-1} = p_{eff}^2/3k_B(T + \theta)$$

Local moment S ~ 1/2 on the C_{60} balls?

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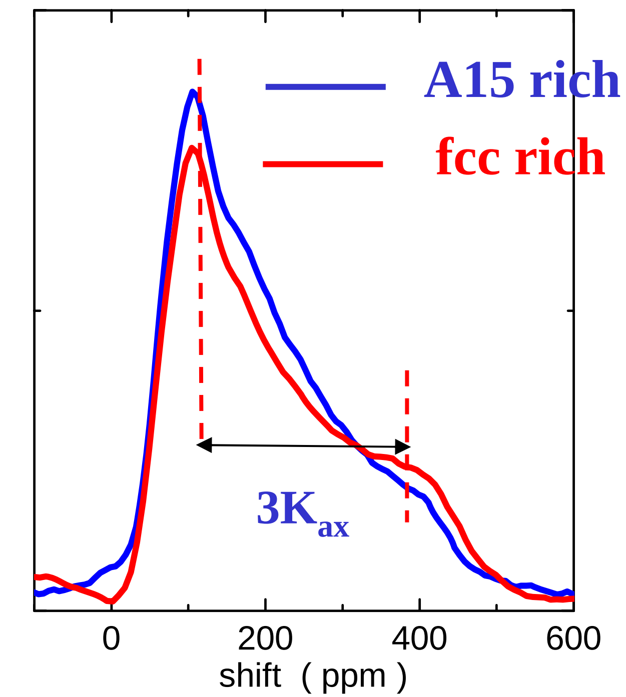
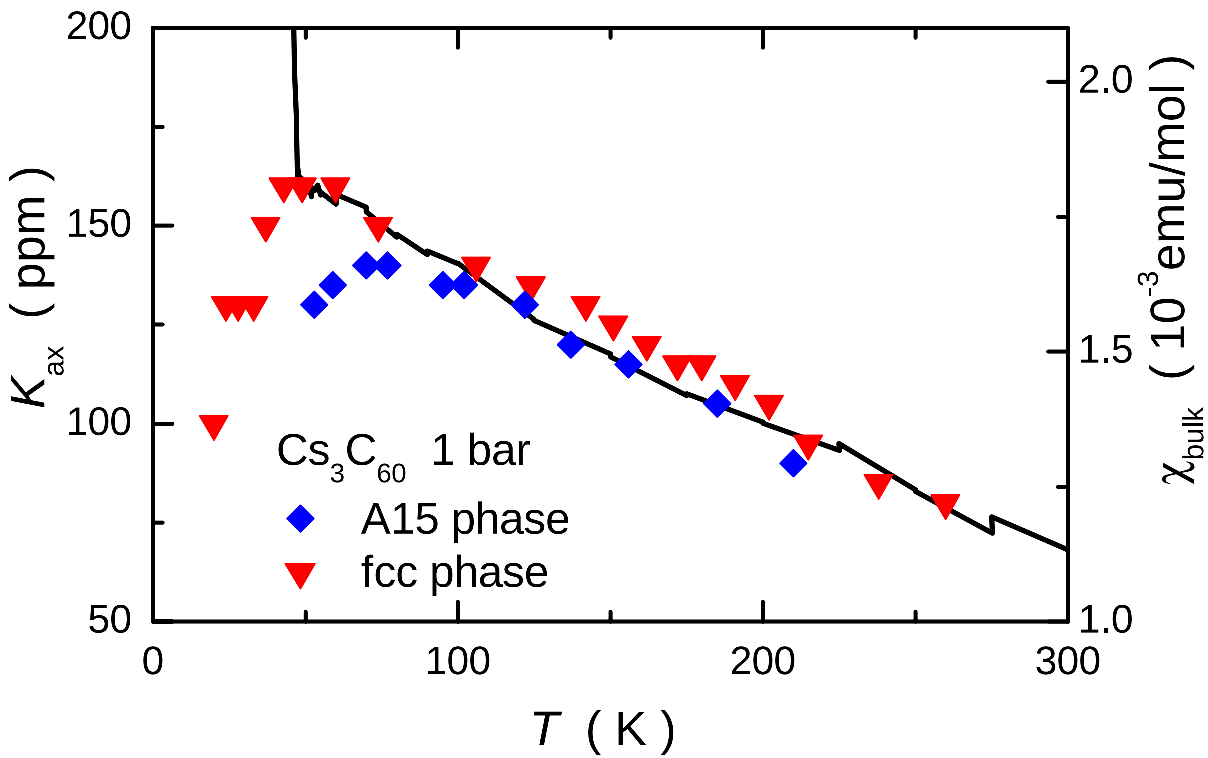
Paramagnetic state susceptibility

^{13}C NMR





Anisotropic hyperfine coupling

$K_{ax} = A \quad \chi_s$

Here A is dipolar A_{exp} ~ 700 Oe/μ_B

Calculated value : A= 640 Oe/μ_B

Local moment is indeed on the C_{60} balls!

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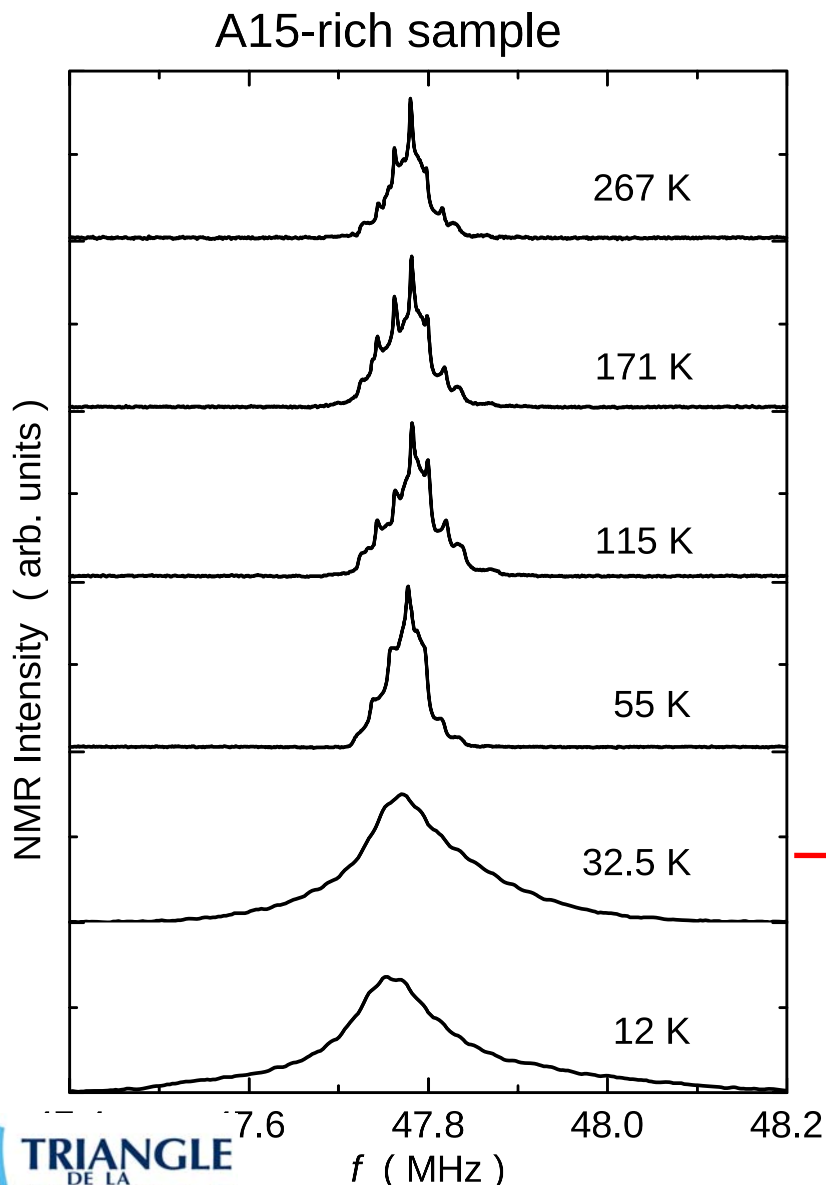
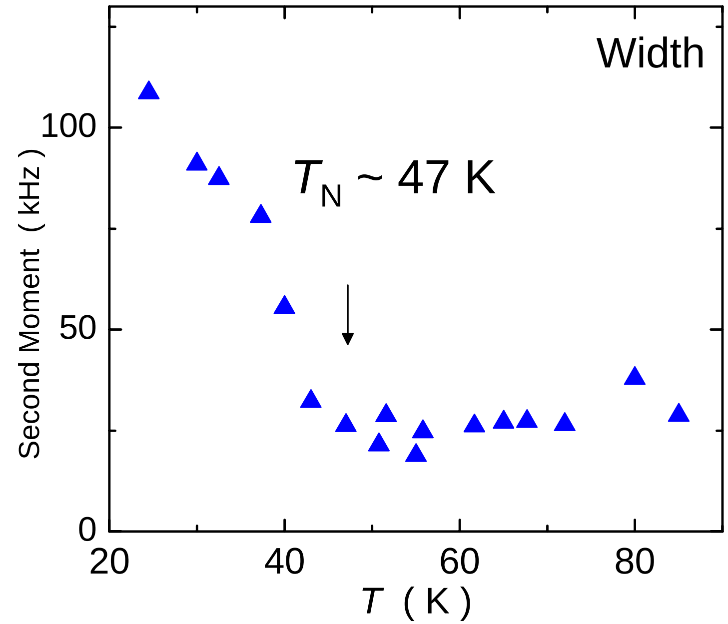
Magnetic transition in A15 phase





Spectral broadening :
static internal magnetic field.

A15 phase shows magnetic order below T_N = 47 K at ambient *p*.

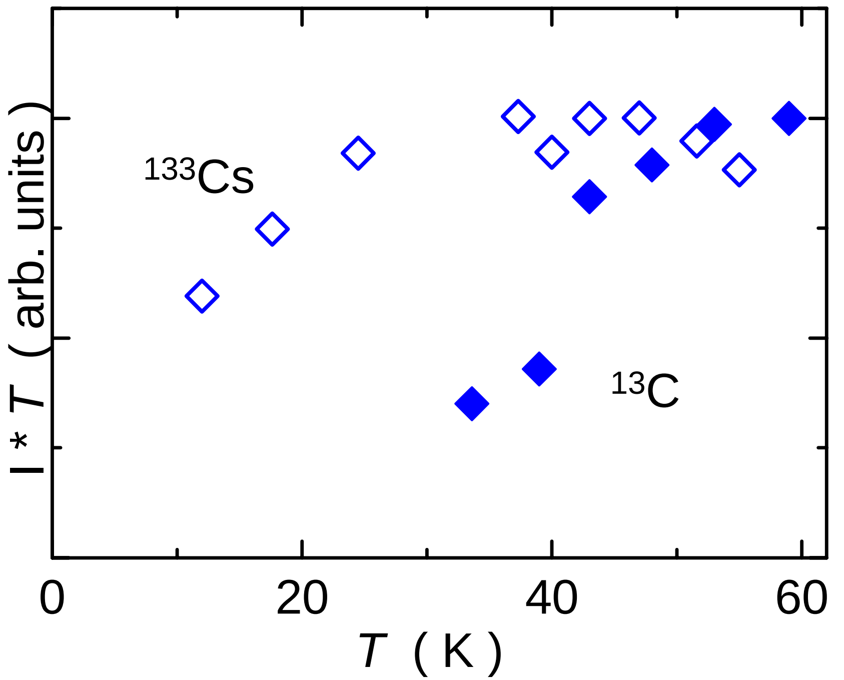
Takabaya

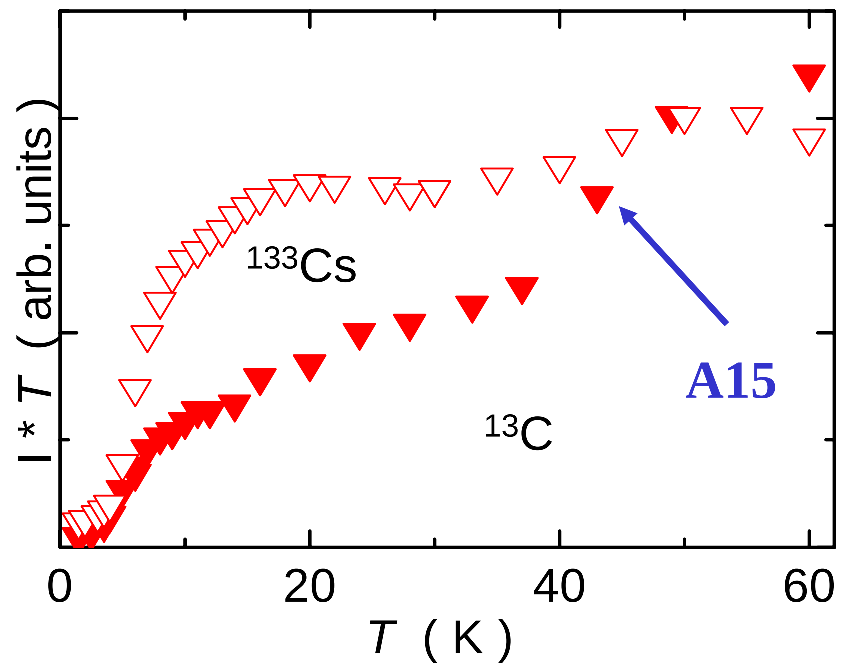
Magnetism and crystal structure

A15 rich

^{13}C and ^{133}Cs NMR intensities

fcc rich





T_N=47K

Gradual freezing
About 10K ?

Geometrical frustration depresses the magnetic ordering (fcc not bipartite)

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Spin dynamics and crystal structure

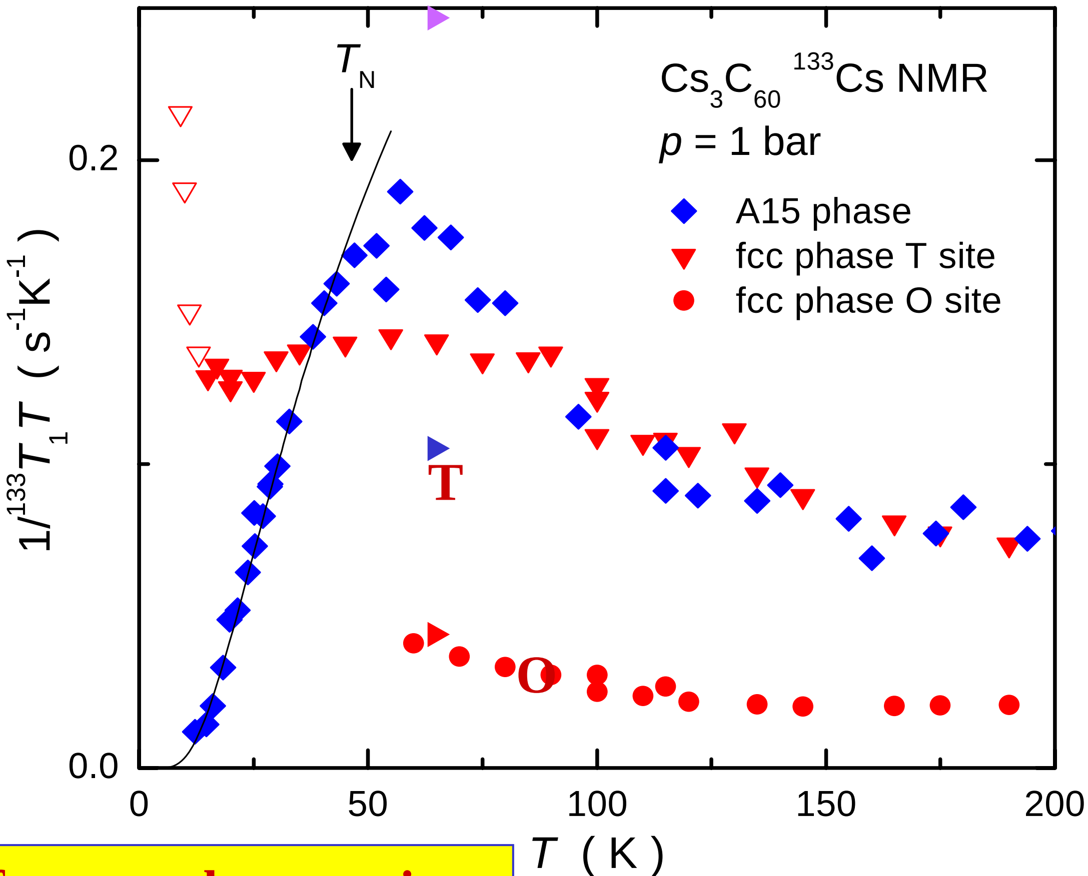
^{133}Cs NMR T_1

Both phases
Enhanced magnetic fluctuations
in the paramagnetic state.

A15 phase
magnetic gap
Δ ~ 50 K

fcc phase
No magnetic gap.
Still some slow dynamics below 10 K.

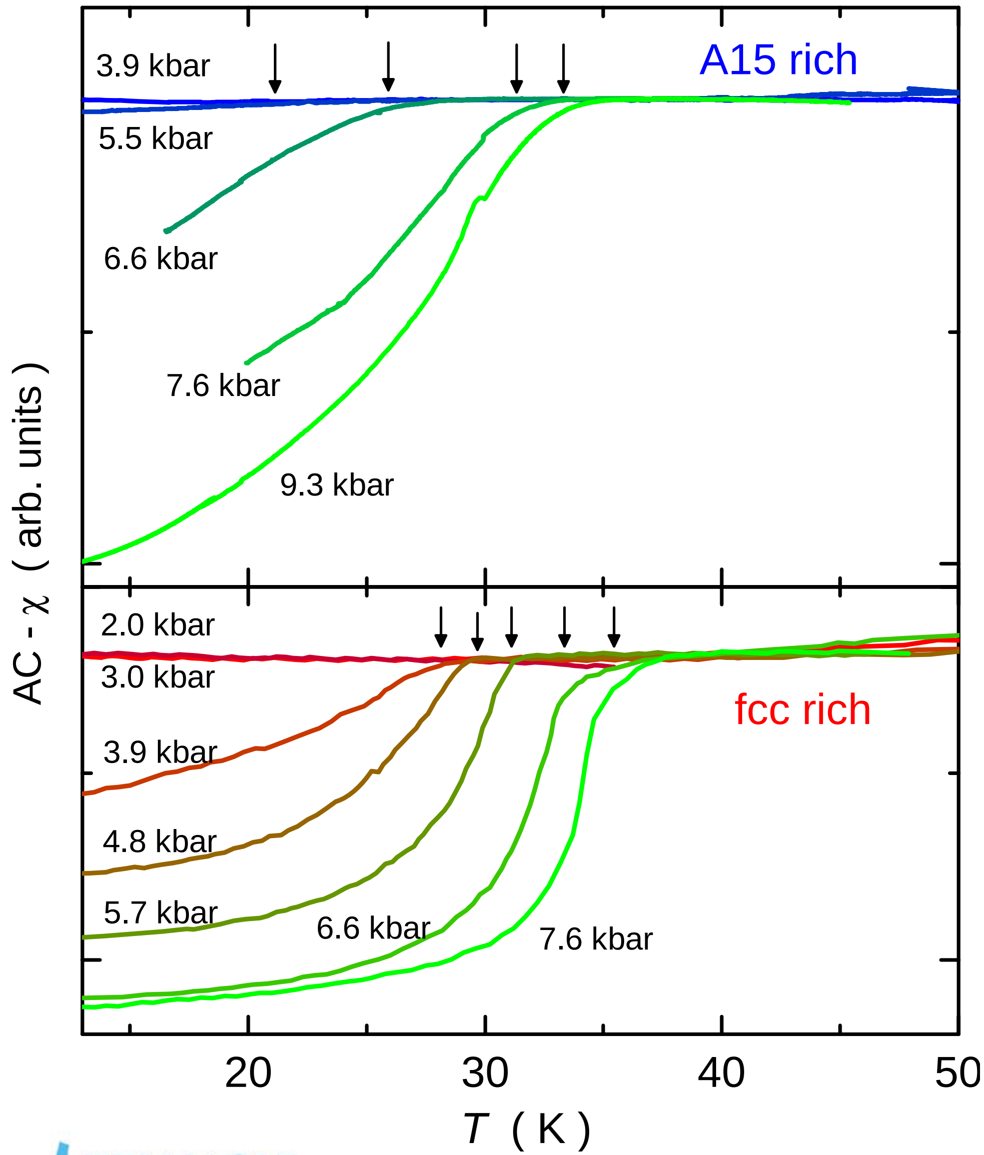
O site has twice smaller coupling constant than that for T site.

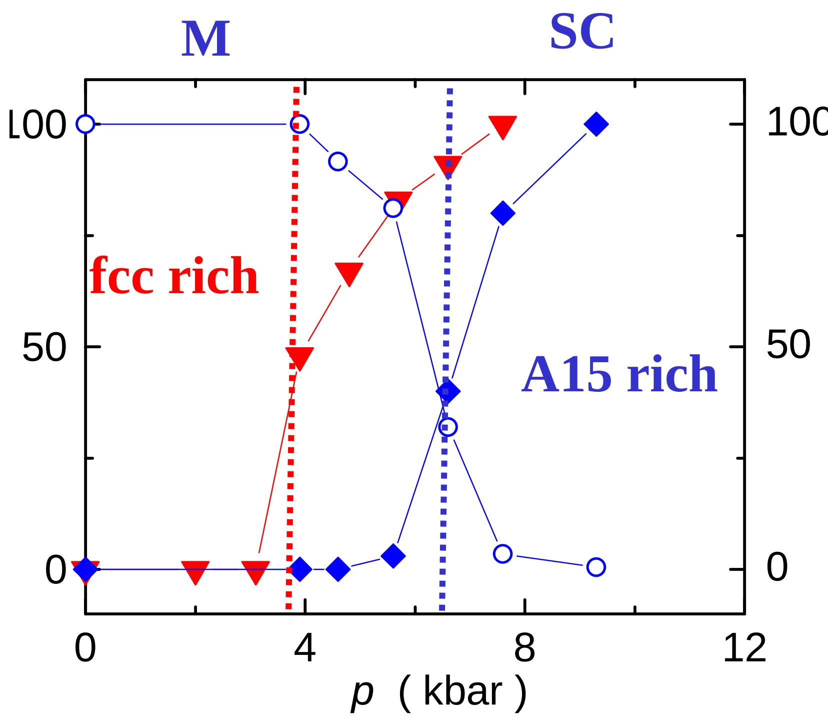


Ordered versus frustrated magnetism (fcc not bipartite)?

H. Alloul, Santa Barbara, 12/02/2010

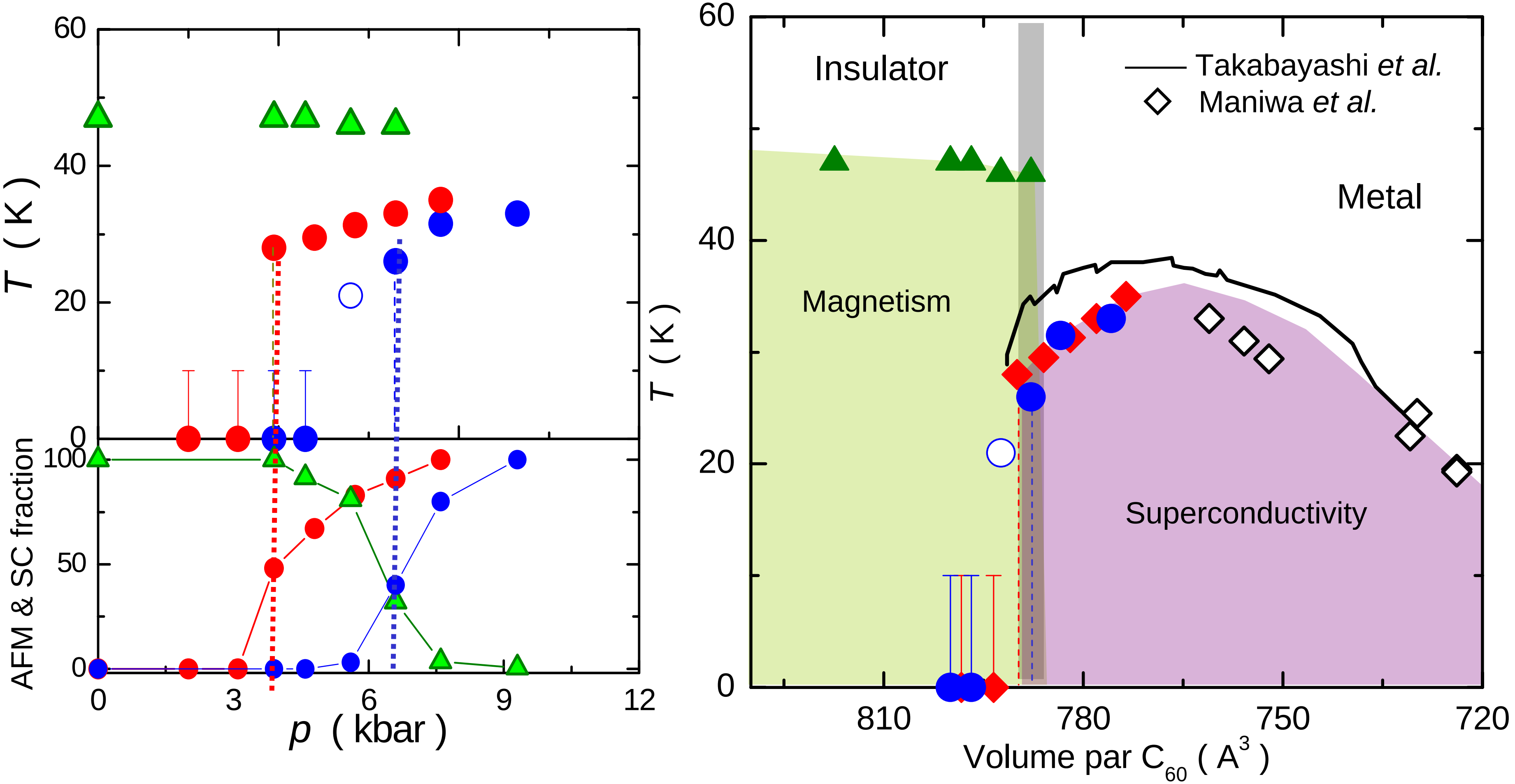
Pressure induced superconductivity





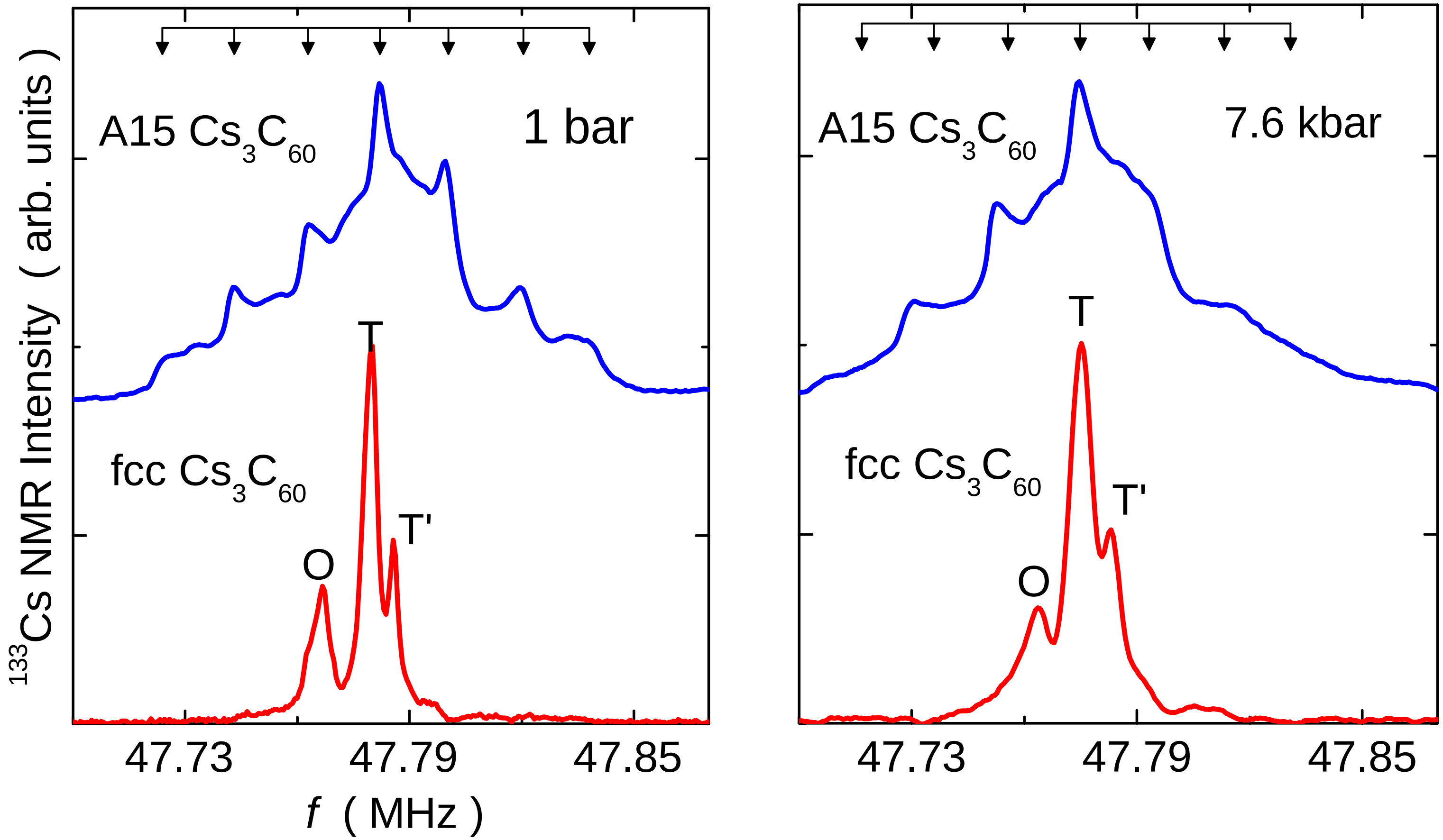
Distinct (*p*,*T*) phase diagrams for the two phases

Phase diagrams



No difference with crystal structure on the SC side

Mott transitions to the metallic state in the A15 phase



p=1bar

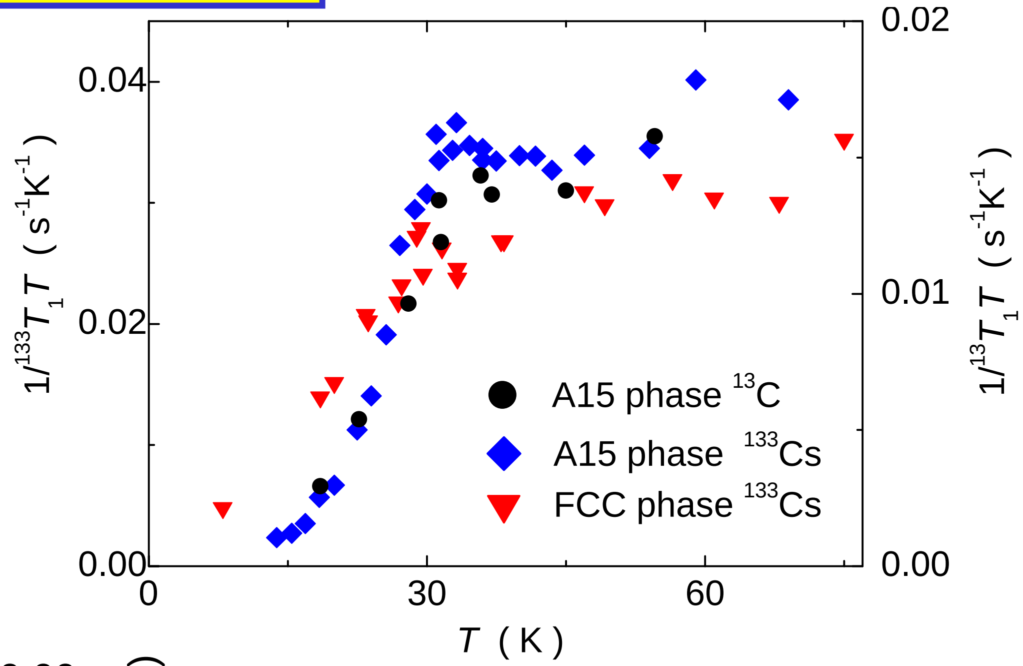
p=7.6 kbar

No change of crystal structures

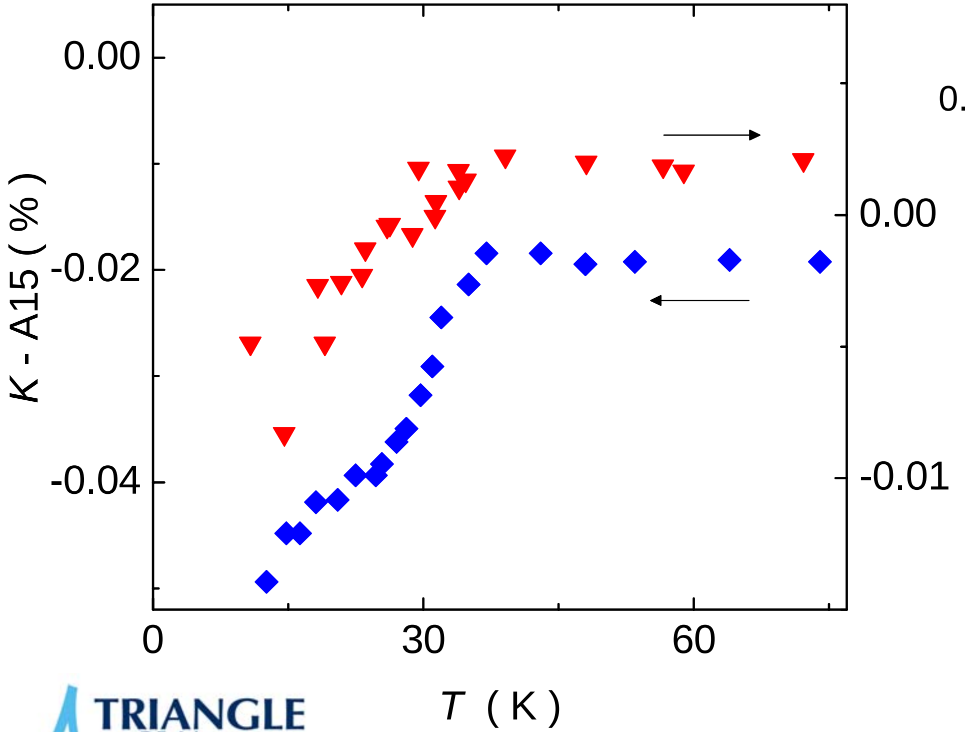
H. Alloul, Santa Barbara, 12/02/2010

Superconductivity

$p >> p_c$



T_1 data



Knight shift (singlet pairing)

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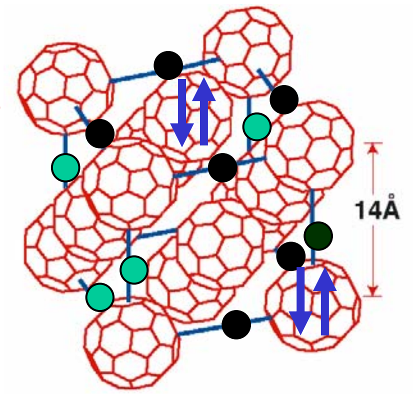
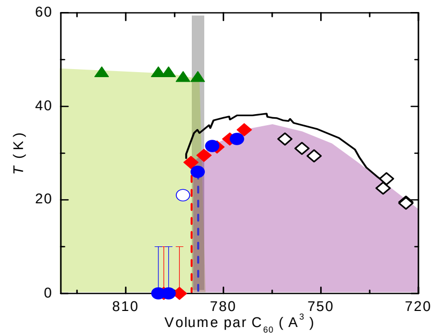
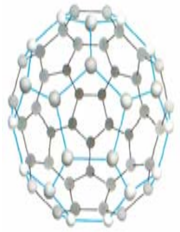
Summary



- Fulleride compounds are peculiar correlated electron systems
- Originalities asociated with their nanostructure
  - Icosaedric symmetry of the soccer ball.
  - Orientational disorder
  - Internal degrees of freedom of the molecule:
  - Phonons, Molecular Jahn-Teller distorsions



- They display many effects driven by correlations
  - Static charge segregation in Cs_1C_{60}
  - High Tc superconductivity near a MIT (Cs_3C_{60})
  - Molecular excitations survive in the solid.
- Those are strongly influenced by molecular Jahn Teller effects
  - Favor singlet formation in Cs_1C_{60}
  - Jahn Teller Mott insulating states in K_4C_{60} or Na_2C_{60}
  - Jahn-Teller effects explain why electronic correlations appear smaller in A_3C_{60} but remain sizable as Cs_3C_{60} is magnetic



- Excellent possibility to study multiorbital Mott transitions

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