

# Emergent phenomena in large interacting communities

Giulio Biroli

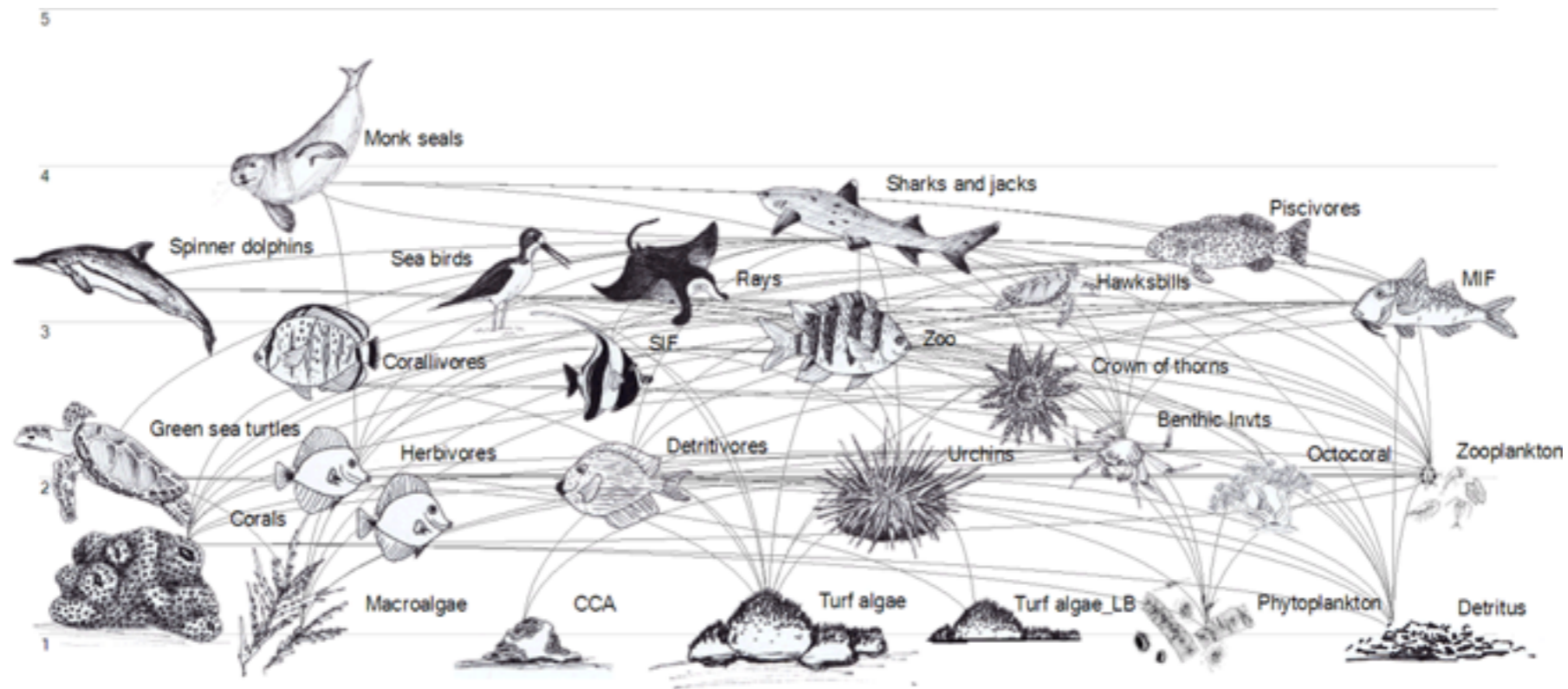
Institute for Theoretical Physics, CEA Saclay, France  
Statistical Physics Lab-ENS Paris

Joint works with G. Bunin, C. Cammarota, V. Ros, F. Roy

# Ecosystems

- Communities formed by individuals belonging to different species.
- Interactions between individuals intra and inter species.
- Competition for resources--Cooperation.
- Abundances of species vary dynamically due to the births and deaths.

# Traditional ecosystem



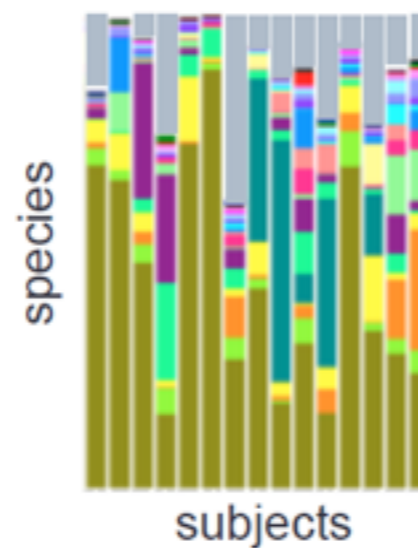
## “Modern” ecosystems



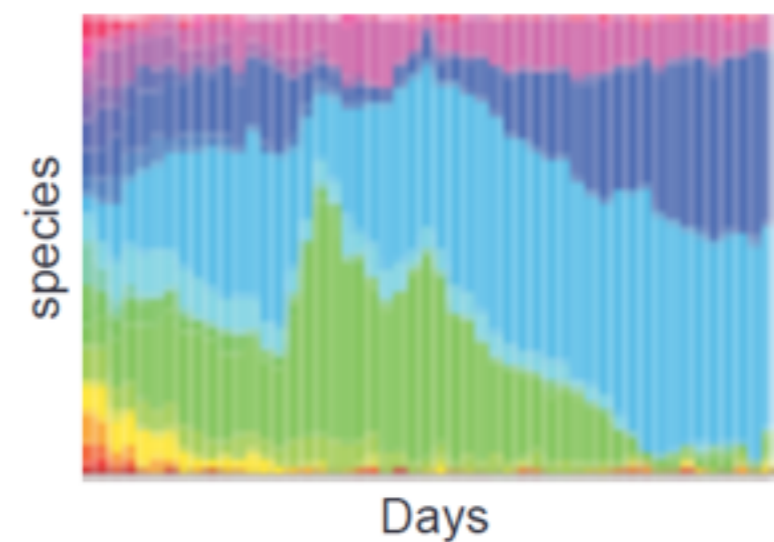
## “Modern” ecosystems



Abundance profiles



Time dependence



Human microbiome project  
hundred-thousand species  
in a given ecosystem

# Questions & Motivations

From a few species to many  $\longrightarrow$  STATISTICAL PHYSICS

Emergent properties of complex ecosystems

- How many equilibria for the same ecosystem?  
In some ecosystems many, Bashan et al. Nature 2016
- Response to perturbations?  
Memory, hysteresis, Dethlefsen, Relman PNAS 2011
- Equilibria or chaotic dynamics?  
Chaos in plankton ecosystem, Beninca et al Nature 2008
- What are the factors determining diversity (number of surviving species)?

# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i (K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i \geq 0$  abundance of species  $i$   
 $S$  is the number of species

Well-mixed population: no-space dependence

Dynamics due to intra-and inter-species interactions

Properties of the community reached dynamically

# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i (K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i$  abundance of species  $i$   
 $S$  is the number of species

$$\frac{dN_i}{dt} = r_i N_i (K_i - N_i)$$

A species alone self-regulates to the abundance  $K_i$



# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i(K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i$  abundance of species  $i$   
 $S$  is the number of species

$$\frac{dN_i}{dt} = N_i \left[ r_i(K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right]$$

↑  
interaction between species

# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i(K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i$  abundance of species  $i$   
 $S$  is the number of species

$$\frac{dN_i}{dt} = N_i \left[ r_i(K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t)$$

↑  
Demographic Noise to  
model fluctuations in  
births and deaths

$$\langle \eta_i(t) \eta_j(t') \rangle = 2\omega^2 \delta(t - t')$$

# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i (K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i$  abundance of species  $i$   
 $S$  is the number of species

$$\frac{dN_i}{dt} = N_i \left[ r_i (K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

Immigration rate



# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i (K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i$  abundance of species  $i$   
 $S$  is the number of species

Large number of species  
( $S \sim 50-100$  is large)

# Lotka-Volterra equations for ecosystems

$$\frac{dN_i}{dt} = N_i \left[ r_i (K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$

$i = 1, \dots, S$

$N_i$  abundance of species  $i$   
 $S$  is the number of species

**Main assumption: complex  $\rightarrow$  random**

(May in ecology & Wigner in physics)

(Determining interactions network: a key inference problem)

$\alpha_{ij}$  Gaussian RVs i.i.d.

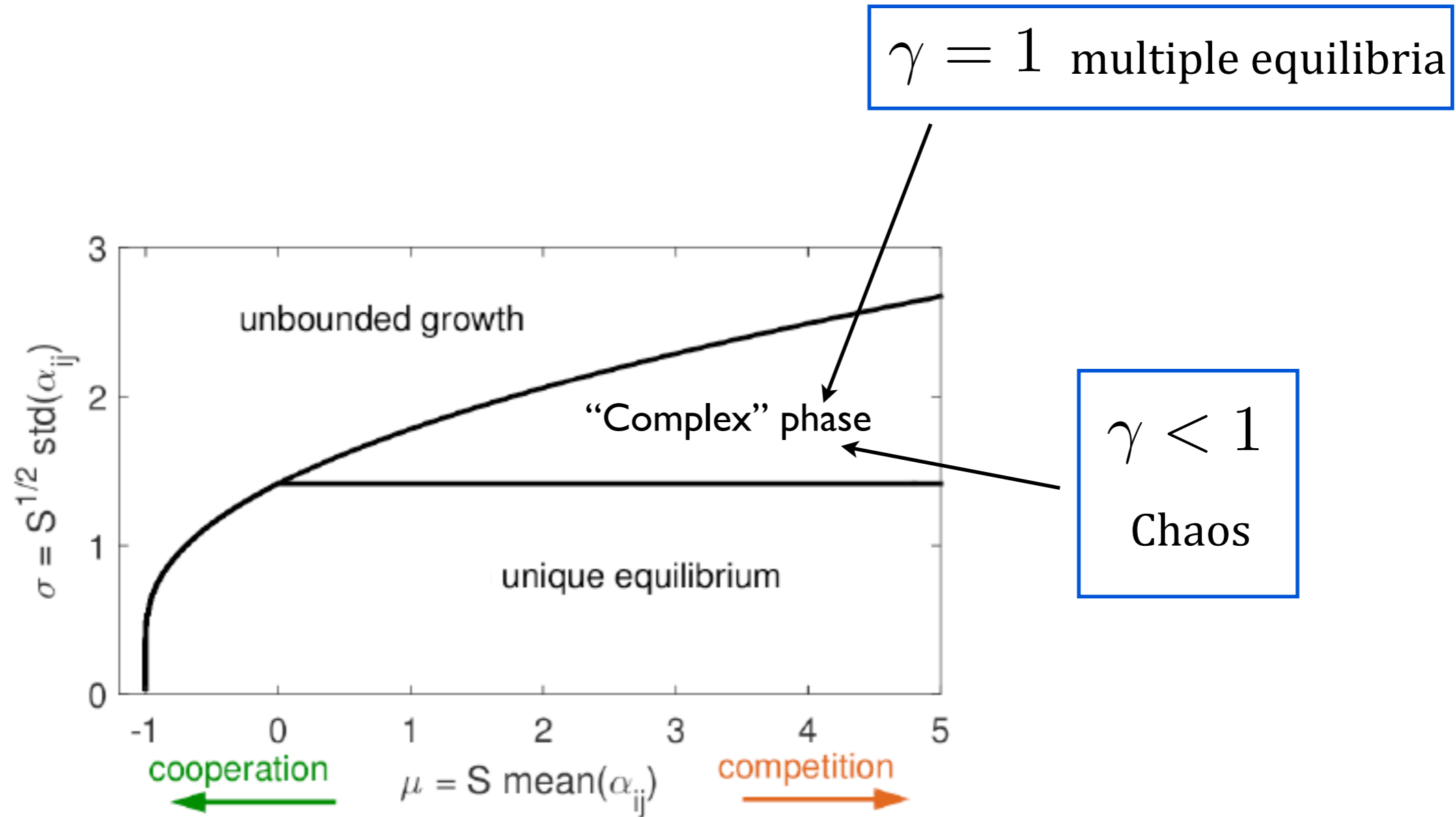
$$\langle \alpha_{ij} \rangle = \frac{\mu}{S} \quad \langle \alpha_{ij}^2 \rangle_c = \frac{\sigma^2}{S}$$

$$\langle \alpha_{ij} \alpha_{ji} \rangle_c = \gamma \langle \alpha_{ij}^2 \rangle_c \quad \gamma = 1 \quad \text{symmetric} \quad -1 \leq \gamma \leq 1$$

Small noise and small immigration rate ( $K_i=1$ )

Representative model, see  
Barbier et al. PNAS to appear

# Ecosystems Phase Transitions



Similar for symmetric and non symmetric interactions (here  $\gamma = 0$ )

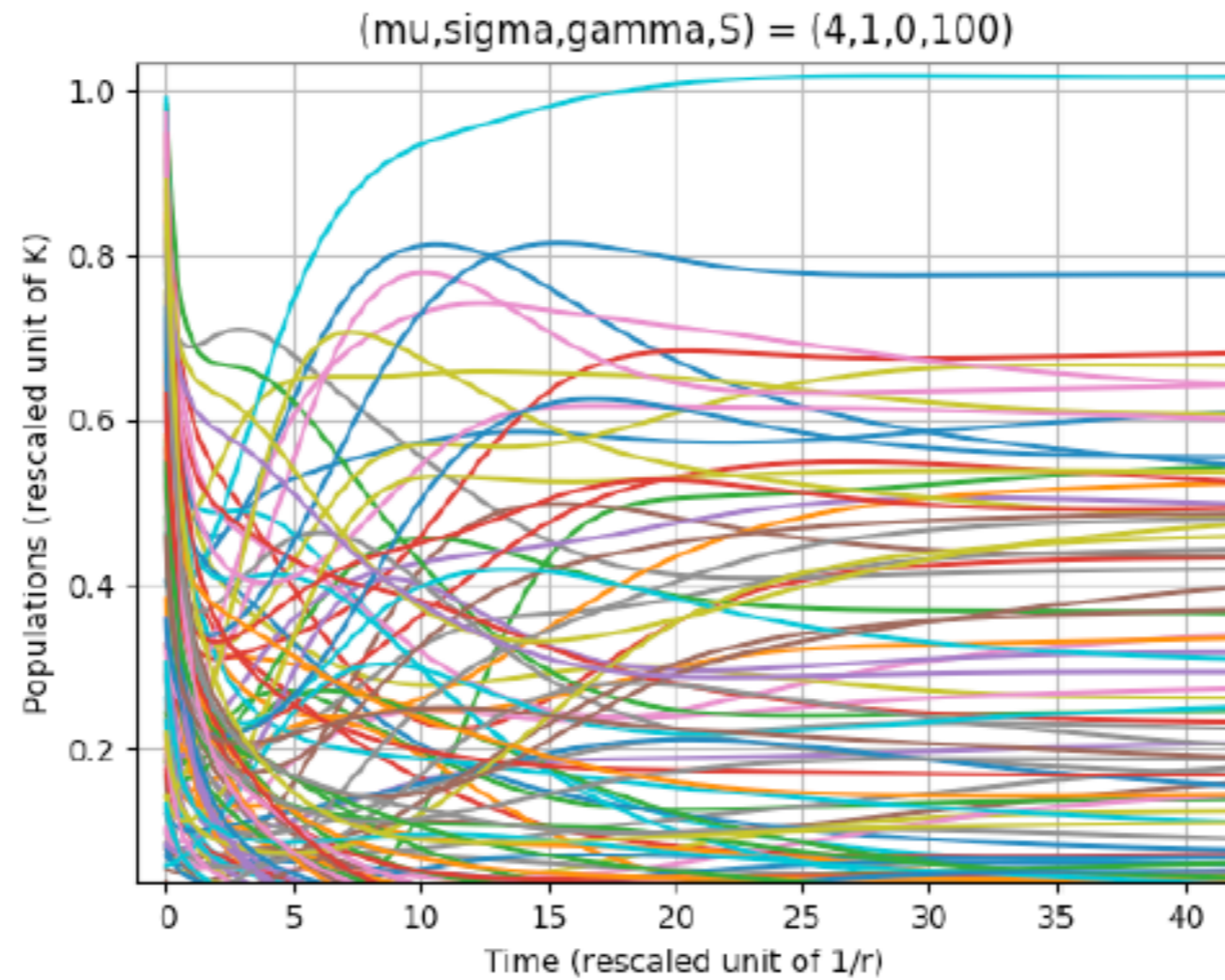
Related works and phase diagrams

Sompolinsky, Crisanti, Sommers '88 ; Diederich, Oppen '89;  
Fisher, Mehta '14; Kessler, Shnerb '15; Bunin '16

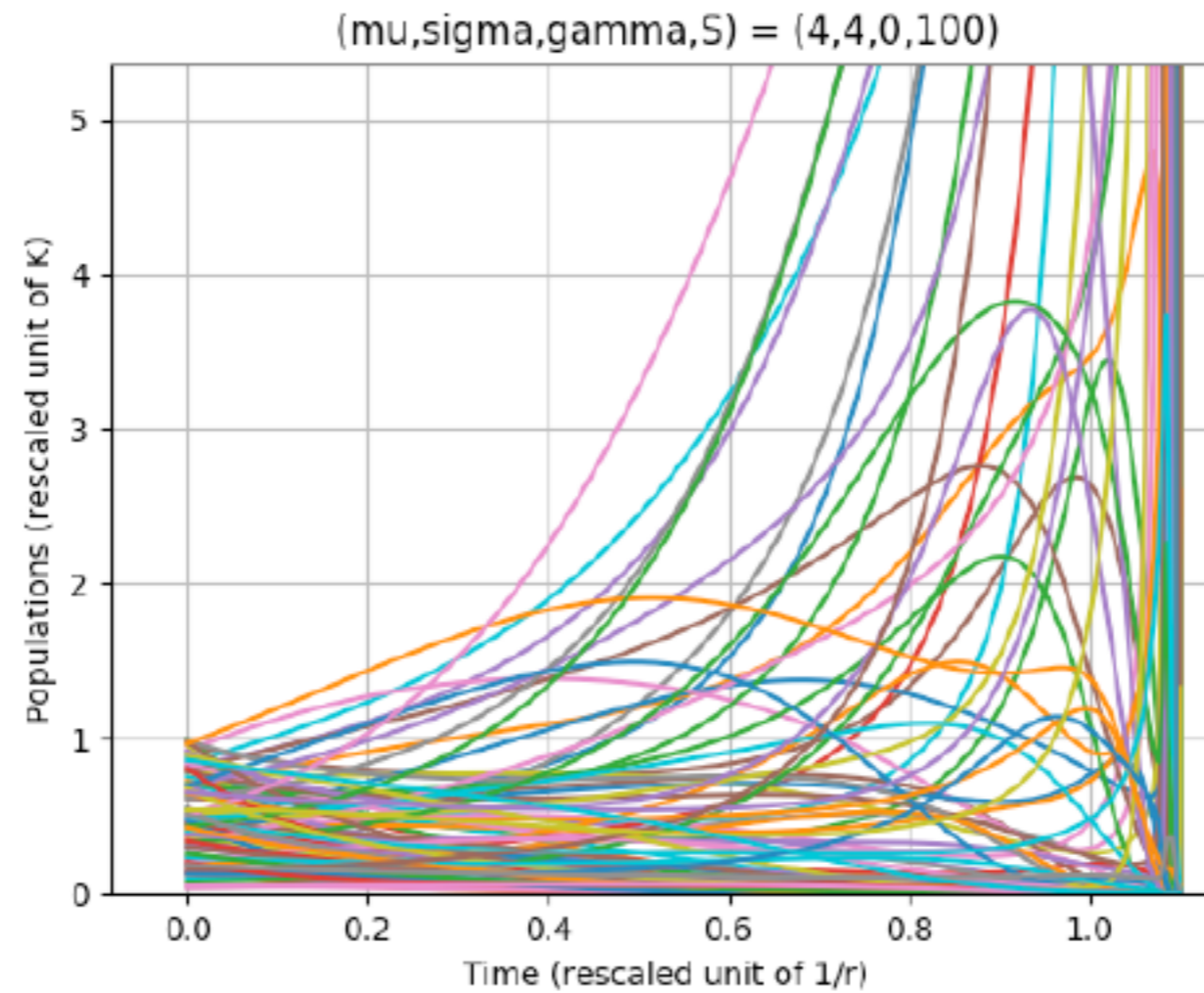
EXACT SOLUTION

G. B., G. Bunin and C. Cammarota arXiv:1710.03606  
and works in progress (F. Roy, V. Ros)

# Unique Equilibrium Phase



# Unbounded Growth Phase





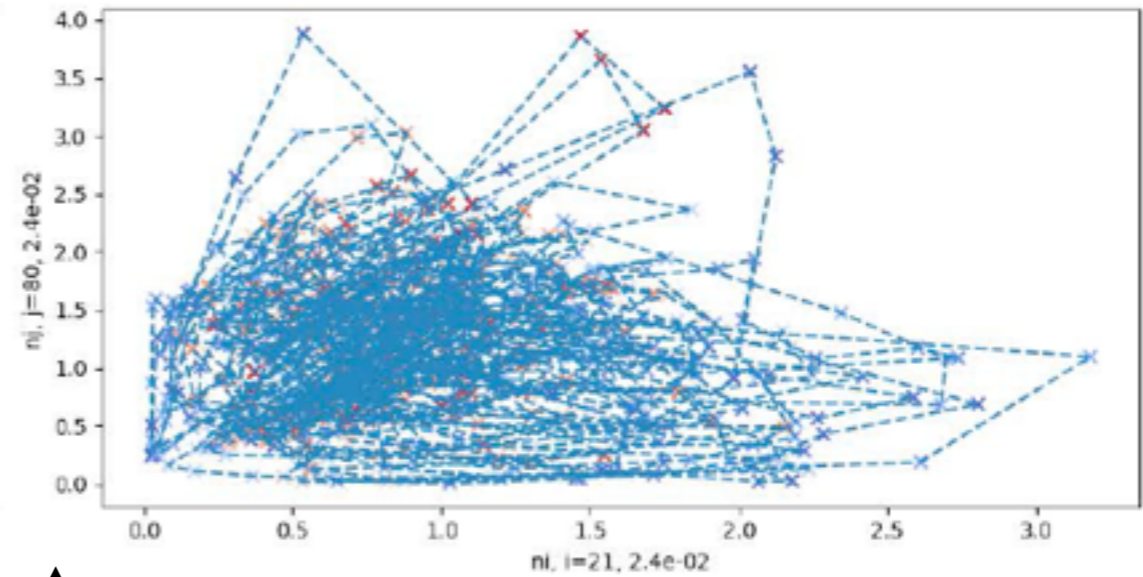
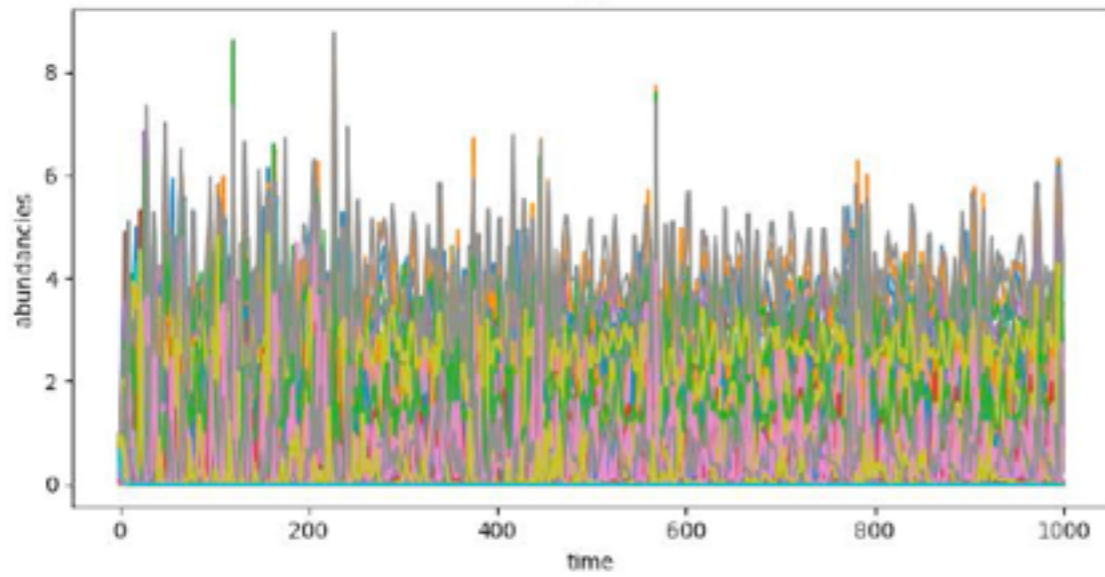
# “Complex” Phase

$\gamma = 1$  symmetric interactions: multiple equilibria

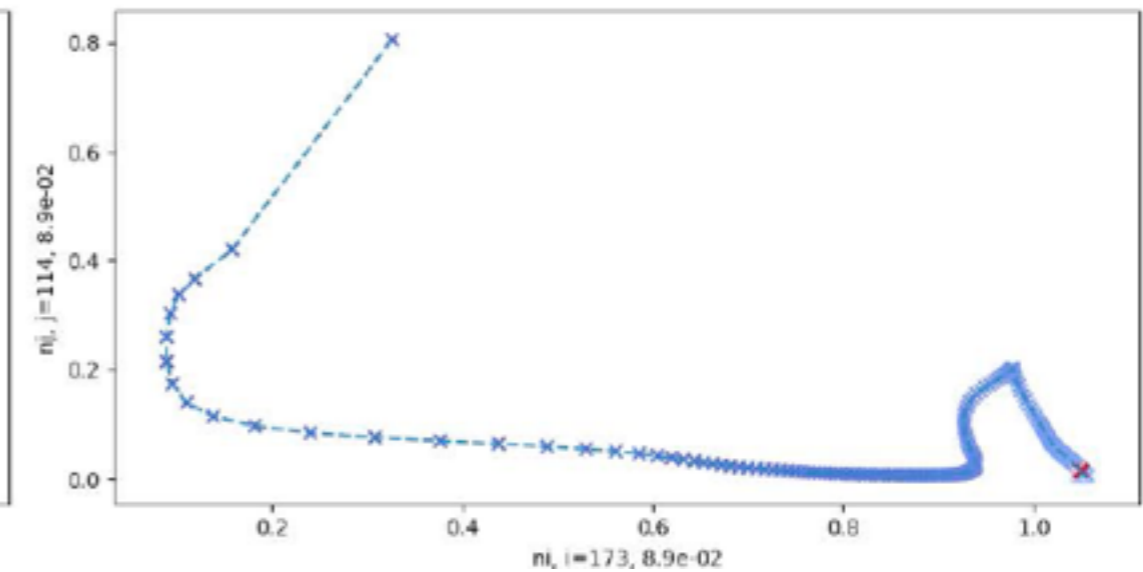
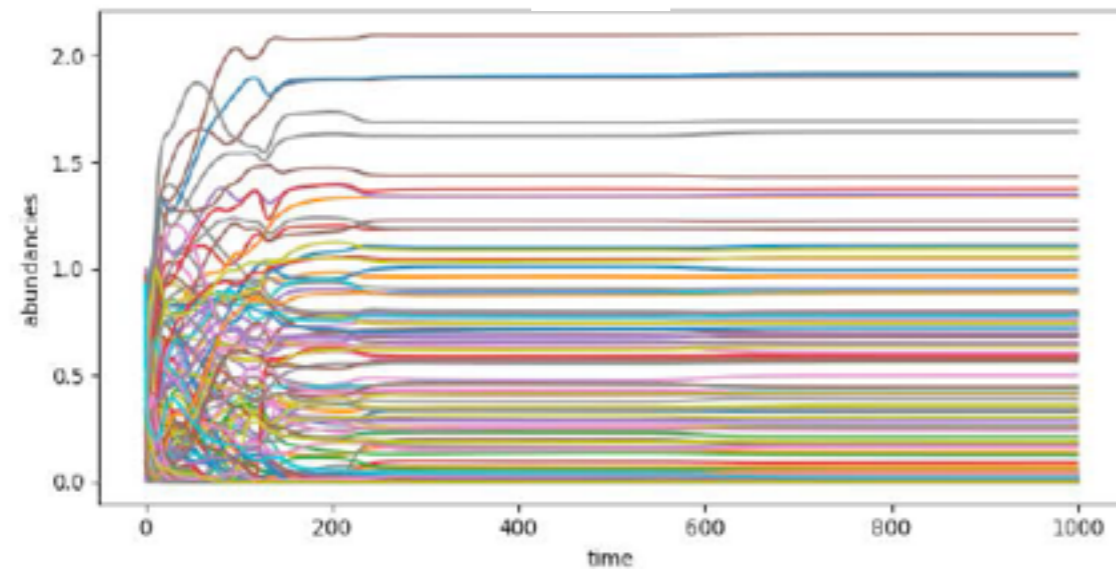
$\gamma < 1$  non-symmetric interactions: chaos

# Transition to Chaos

$$\gamma < 1$$



$\sigma$



# Symmetric interactions

$$\gamma = 1$$

$$\frac{dN_i}{dt} = N_i \left[ r_i(K_i - N_i) - \sum_{j, (j \neq i)} \alpha_{ij} N_j \right] + \sqrt{N_i} \eta_i(t) + \lambda_i$$



$$\frac{dN_i}{dt} = -N_i \partial_{N_i} E(\{N_i\}) + \sqrt{N_i} \eta_i(t)$$

Langevin equation

$$E = \sum_i V_i(N_i) + \frac{1}{2} \sum_{i \neq j} \alpha_{ij} N_i N_j + \sum_i (\omega^2 - \lambda_i) \log N_i$$

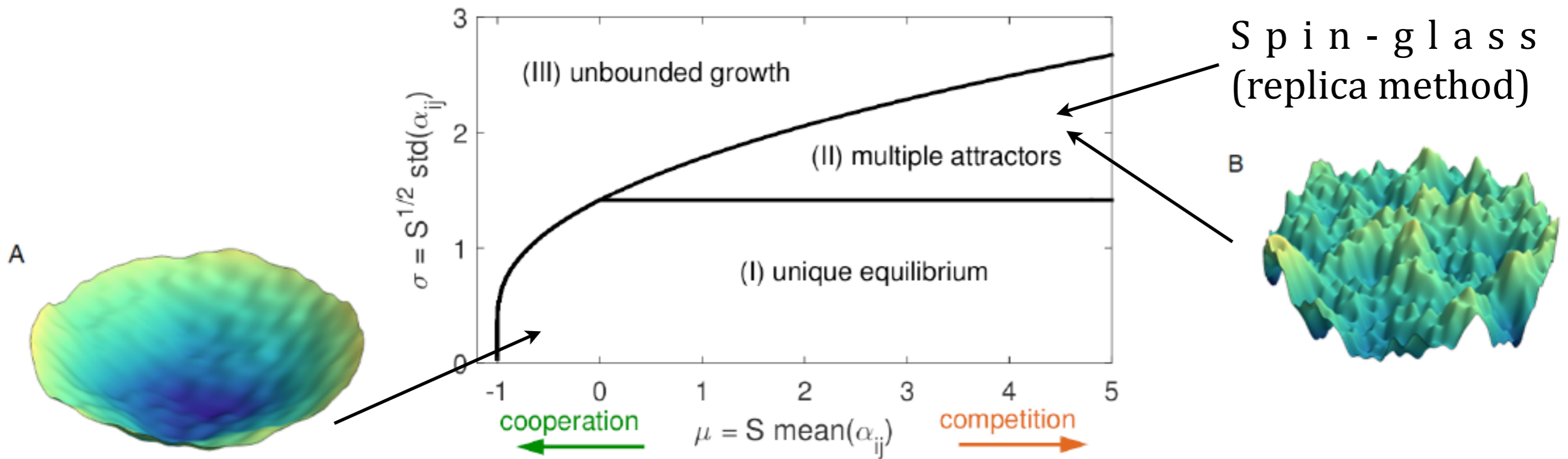
Stochastic dynamics of a disordered system  
(~spin-glass)

small or zero noise  $\longrightarrow$

Low temperature physics

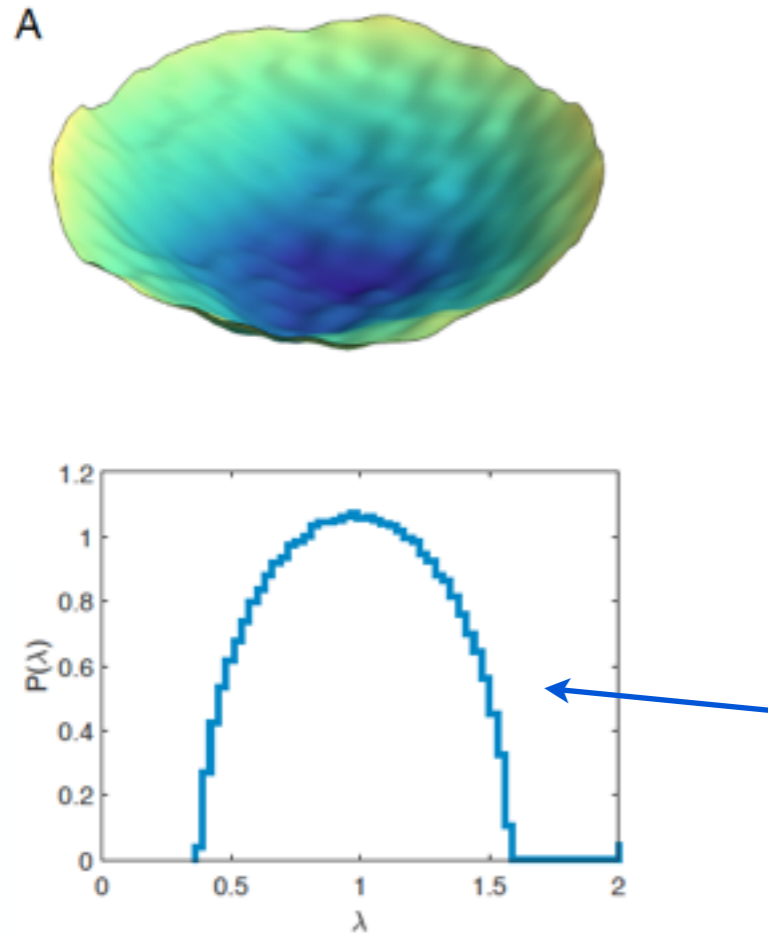
$$T = \omega^2$$

# The phase diagram & the energy landscape



# The Two Phases

One equilibrium

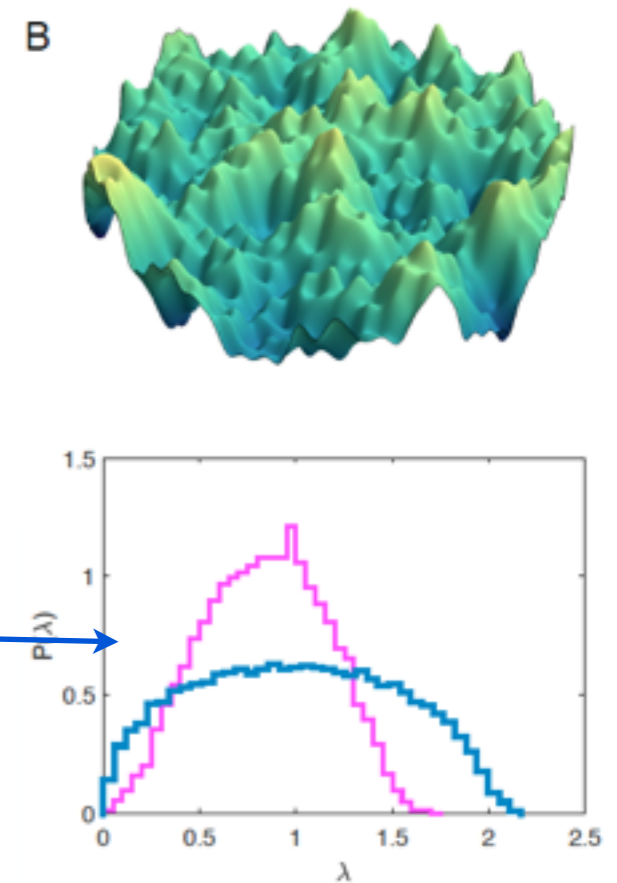


One single  
stable equilibrium

Inverse eigenvalue distribution  
of the stability matrix  $\frac{\partial N_i^*}{\partial K_j}$

Eigenvalue distribution of the  
Hessian

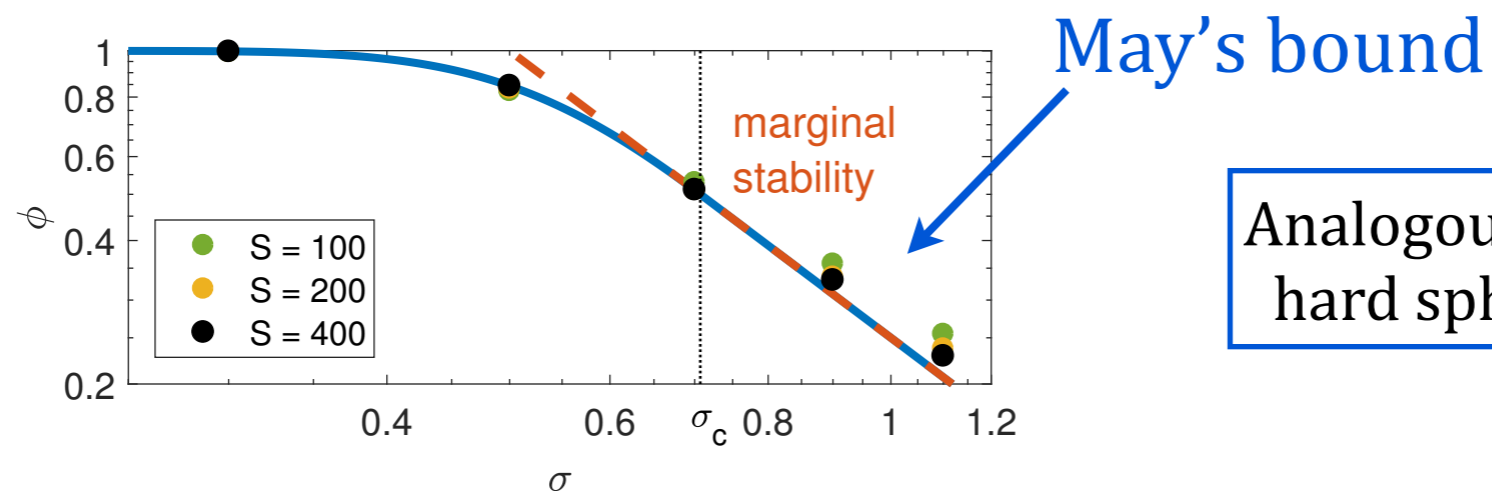
Multiple Equilibria  
(Critical Spin-Glass Phase)



All equilibria  
marginally stable

# Critical Multiple Equilibria Phase

Marginal stability fixes dynamically the diversity

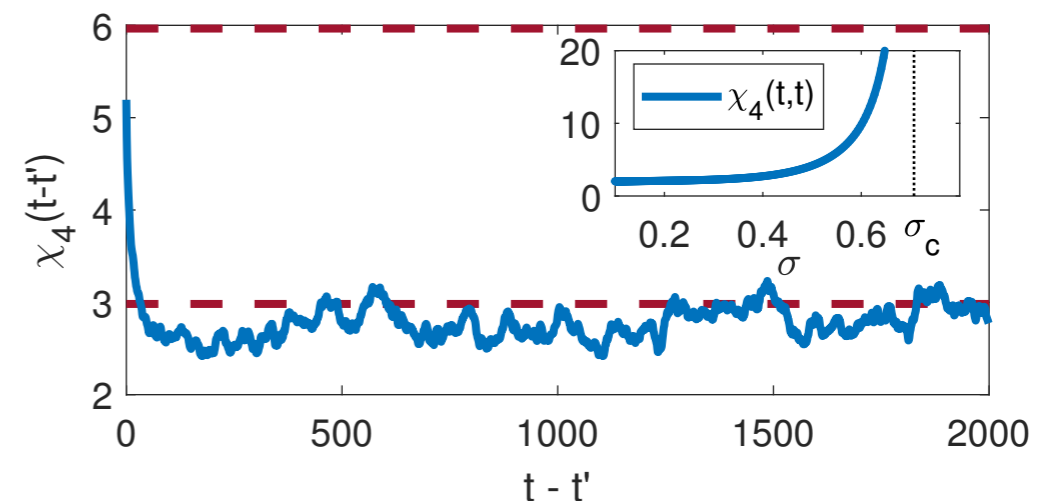


Analogous phenomenon in jamming of hard spheres: isostaticity of packings

- Extreme susceptibility to perturbations (memory only in the one equilibrium phase)

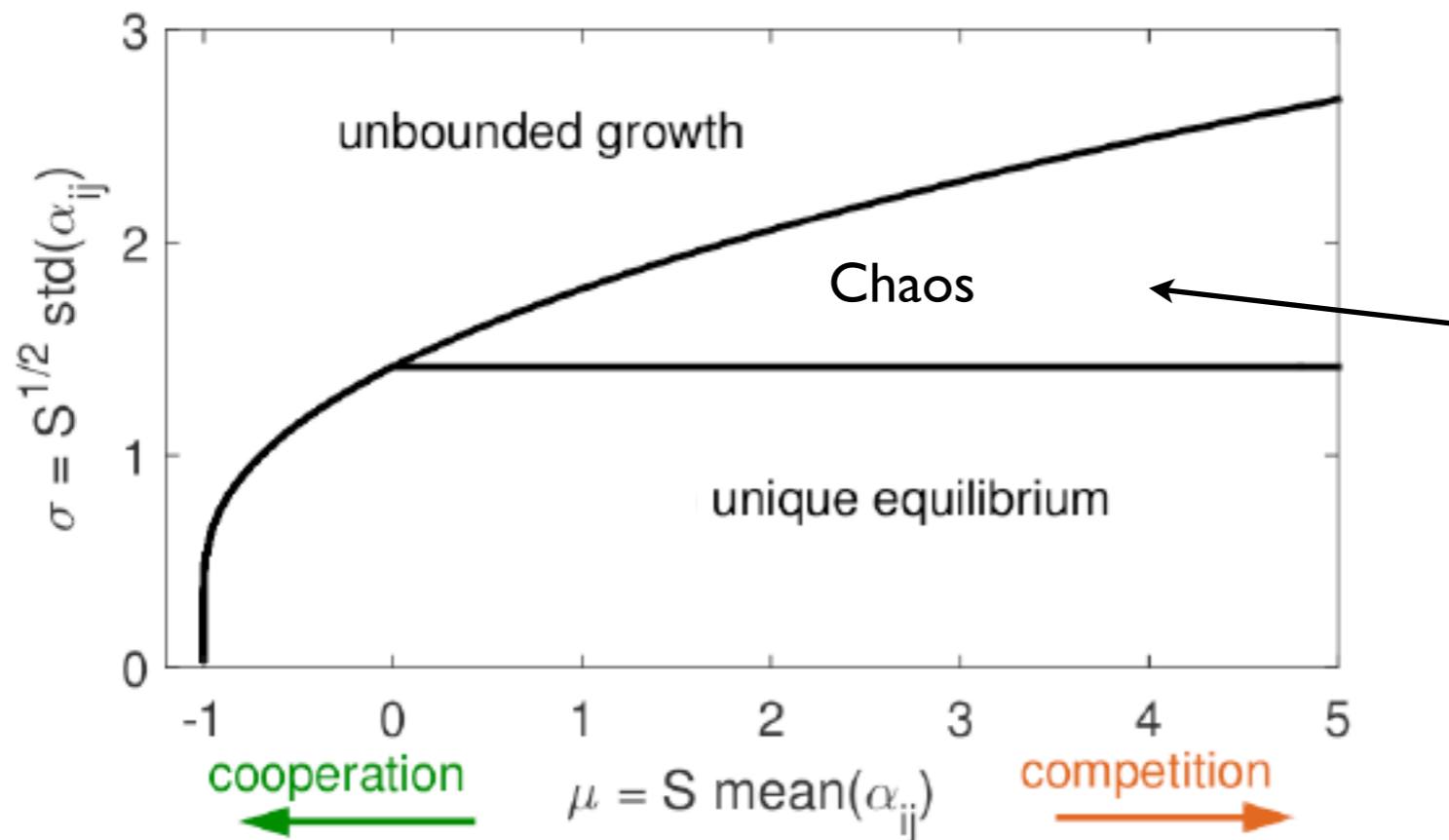
- Large fluctuations-correlations

$$\chi_4(t, t') = \frac{1}{S} \sum_{ij} [\langle \delta N_i(t) \delta N_i(t') \delta N_j(t) \delta N_j(t') \rangle - \langle \delta N_i(t) \delta N_i(t') \rangle \langle \delta N_j(t) \delta N_j(t') \rangle]$$



# Dynamics and Transition to Chaos

$$\gamma < 1$$



All equilibria are unstable  
(Kac-Rice Method)

Chaotic dynamics  
(dynamical mean-field theory)

Ongoing: characterize chaotic dynamics, properties of the transition to chaos

# Emergent phenomena in interacting communities

- Different phases of ecosystems from the exact solution of the Lotka-Volterra model of ecosystems
- An entire region with multiple equilibria poised at the edge of stability:
  - marginal phase, extreme susceptibility to perturbations, large correlations, ...
  - diversity is dynamically fixed by the requirement of being marginal stable: May's bound is saturated
- Chaotic phase where all equilibria are unstable
- Generality beyond the particular model we studied: emergent properties as for phases of matter.
- Many perspectives: Chaotic dynamics, slow dynamics, avalanches, other functional responses, retardation effects, space dependence, ...