

Studies of Charge Order in Single Crystals of LuFe_2O_4 and Fe_2OBO_3

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Work supported by BES Division of Materials Sciences and Engineering

Collaborators

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UMass -- P. Khalifah

Andy Christianson

Julich – R. Hermann, W. Schweika

N.C. State – H. J. Xiang, M.-H. Whangbo

Ames -- J. W. Kim

UK – V. Varadarajan, J. W. Brill



Special thanks to Art Sleight

Manuel Angst

Motivation

- Search for model charge order systems in which electrostatics play an important role
- Charge ordering is an intriguing route to ferroelectricity, and there may be strong coupling between magnetic & ferroelectric order parameters
- ORNL has good tools for these studies – i.e., neutron scattering, electron microscopy, piezoelectric force microscopy

Outline

LuFe_2O_4 – snapshot of work in progress

- Magnetization & thermal properties
- Neutron scattering
- Resonant Ultrasound Spectroscopy
- IR spectroscopy

Fe_2OBO_3 – summarize results presented in 2 papers*

- first growth of single crystals
- first observation of C.O. superstructure
- BVS analysis showing integer iron valence
- discovery of incommensurate phase

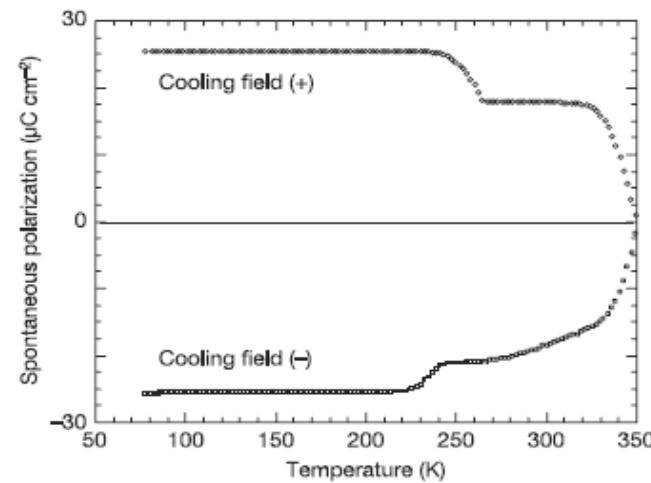
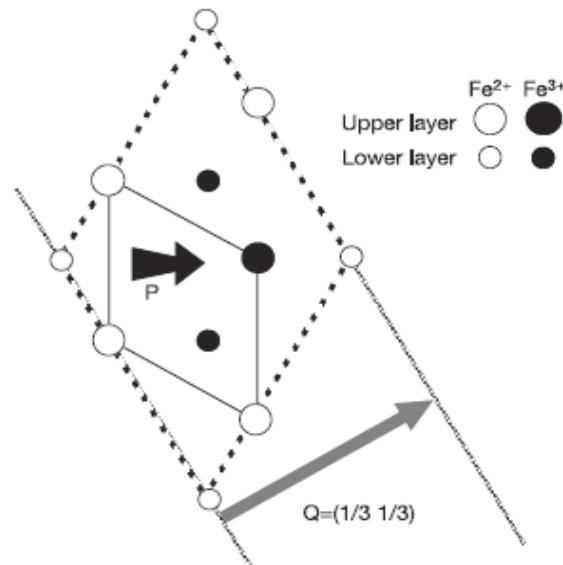
*M. Angst, et al. PRL **99**, 086403 (2007)

M. Angst, et al. cond-mat: 0707.3127

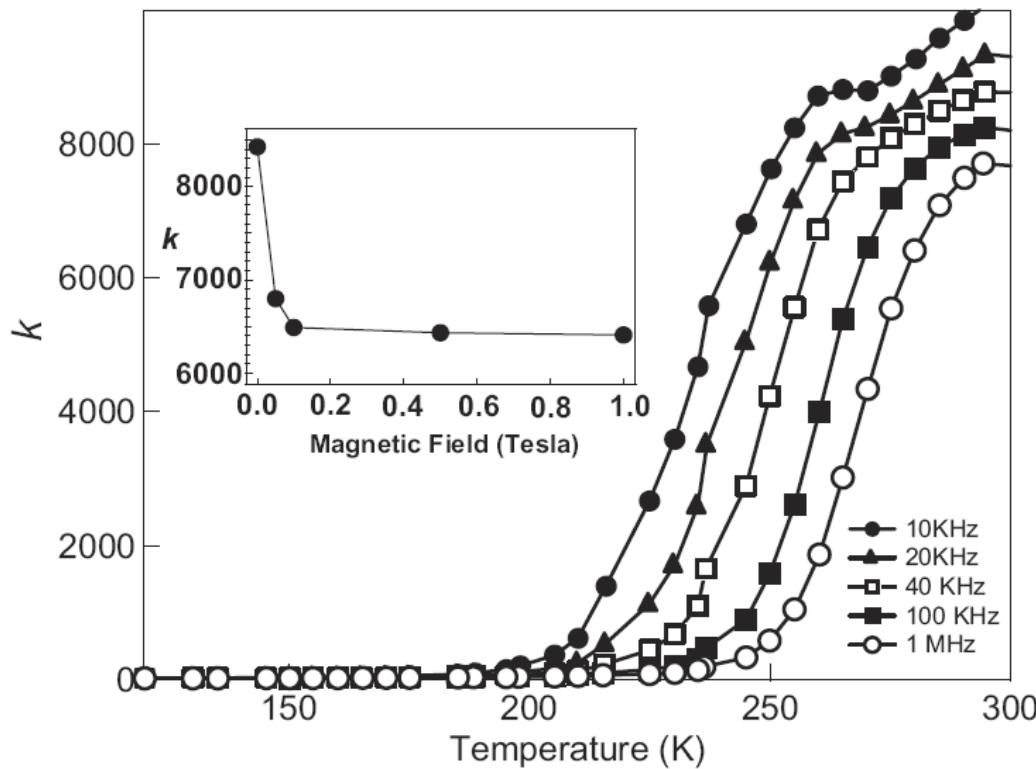
LETTERS

Ferroelectricity from iron valence ordering in the charge-frustrated system LuFe_2O_4

Naoshi Ikeda¹, Hiroyuki Ohsumi¹, Kenji Ohwada², Kenji Ishii², Toshiya Inami², Kazuhisa Kakurai³, Youichi Murakami⁴, Kenji Yoshii², Shigeo Mori⁵, Yoichi Horibe⁵ & Hijiri Kitô⁶



Large Magneto-Dielectric Effect in LuFe₂O₄



M. Subramanian et al., Adv. Mater. **18**, 1737 (2006)

Structure of LuFe₂O₄

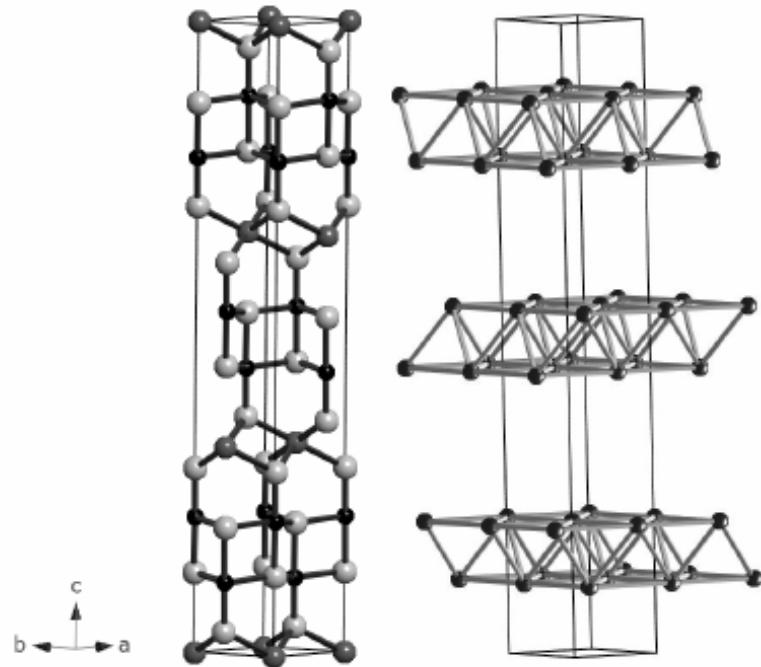
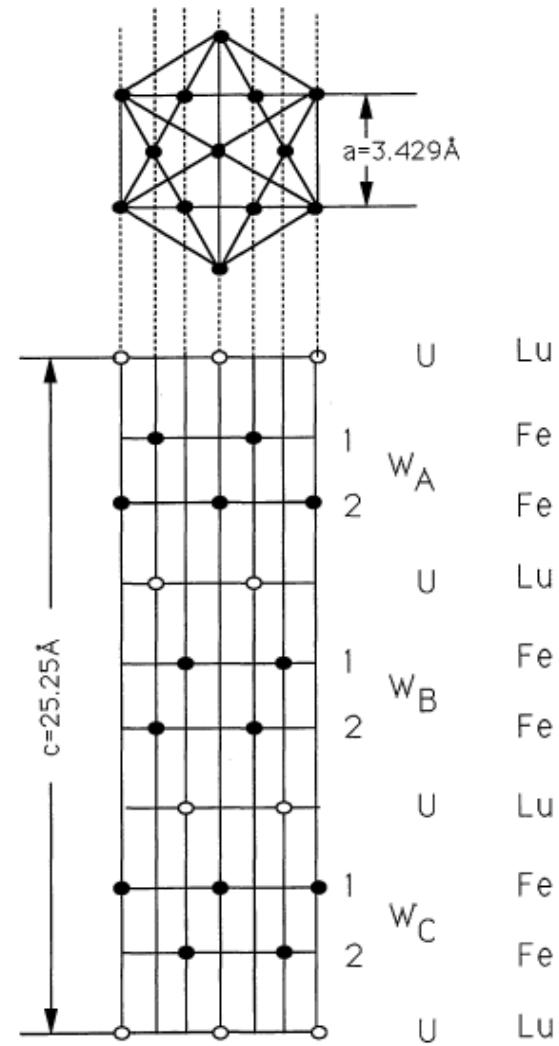


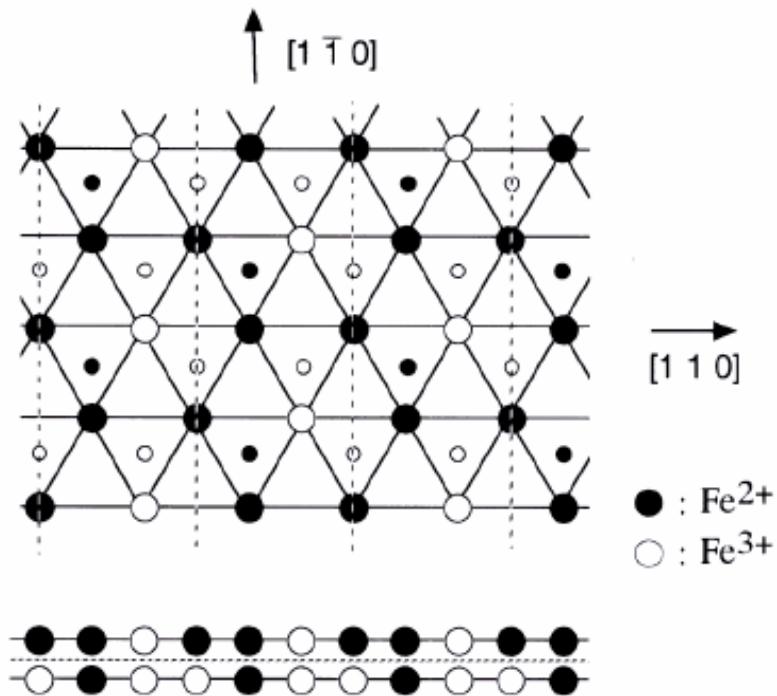
Figure 1. Crystal structure of LuFe₂O₄ showing the layered arrangement of Lu (large dark-grey spheres), Fe (small black spheres), and oxygen (large white spheres) along the c -axis (left). The Fe double layers are shown with a triangular interconnectivity (right). The Fe—Fe distances within a layer (3.44 \AA) are longer than the Fe—Fe distances between the layers (3.156 \AA).

M. Subramanian et al., Adv. Mater. (2006)

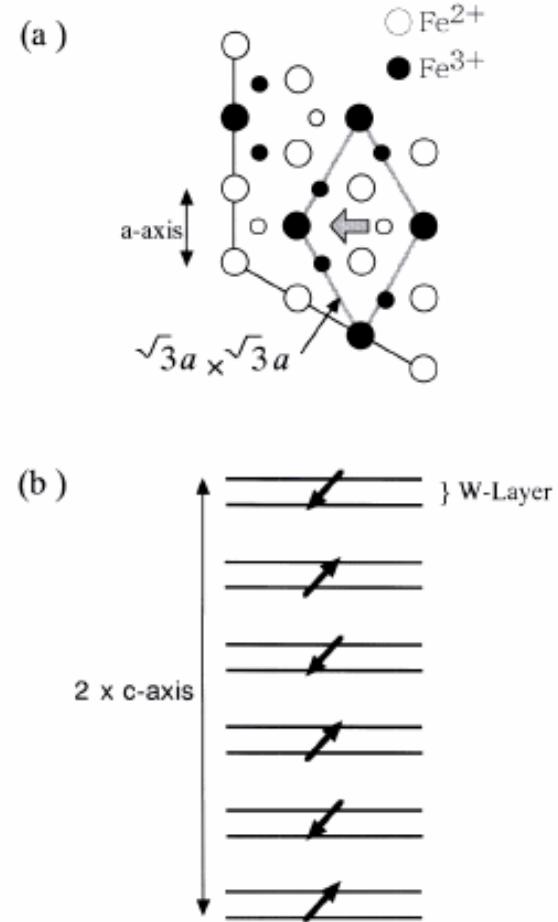


Iida, et al. JPSJ (1993)

Proposed Charge Ordering Pattern (Ikeda)

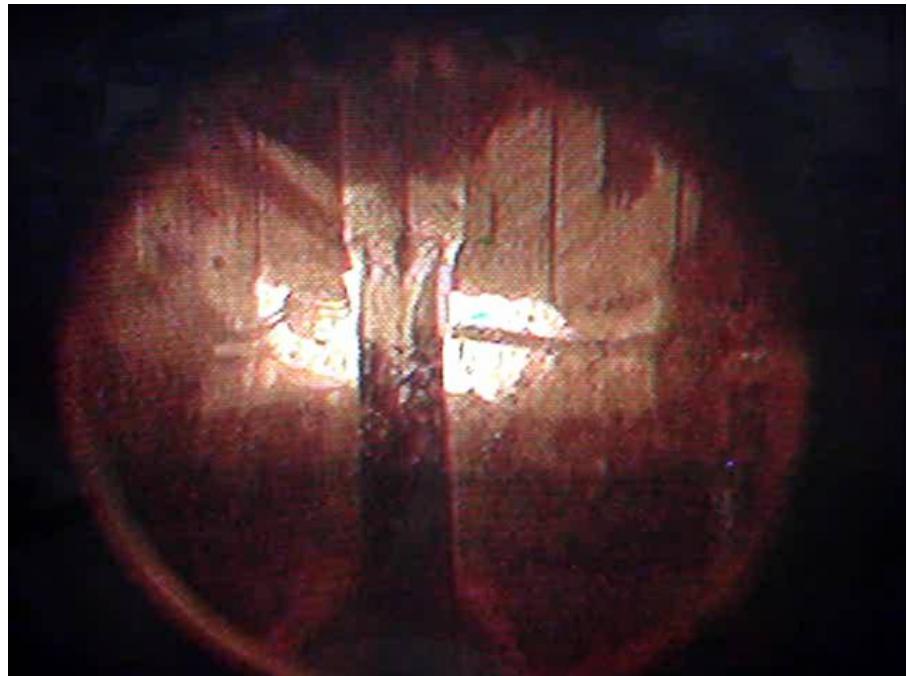
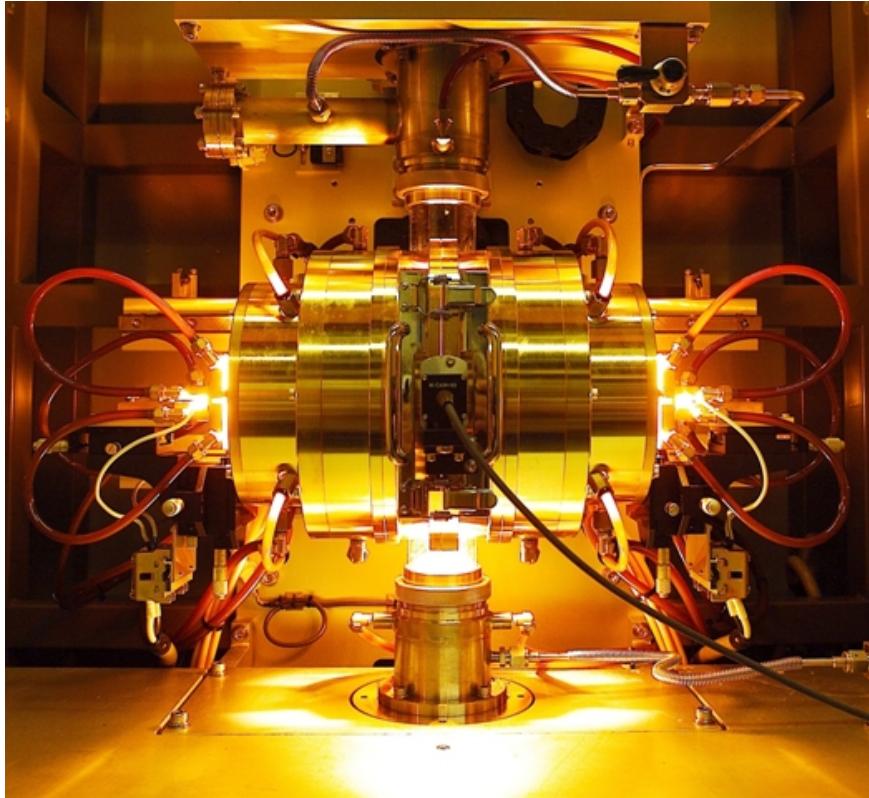


Yamada & Ikeda, JKPS (1998)

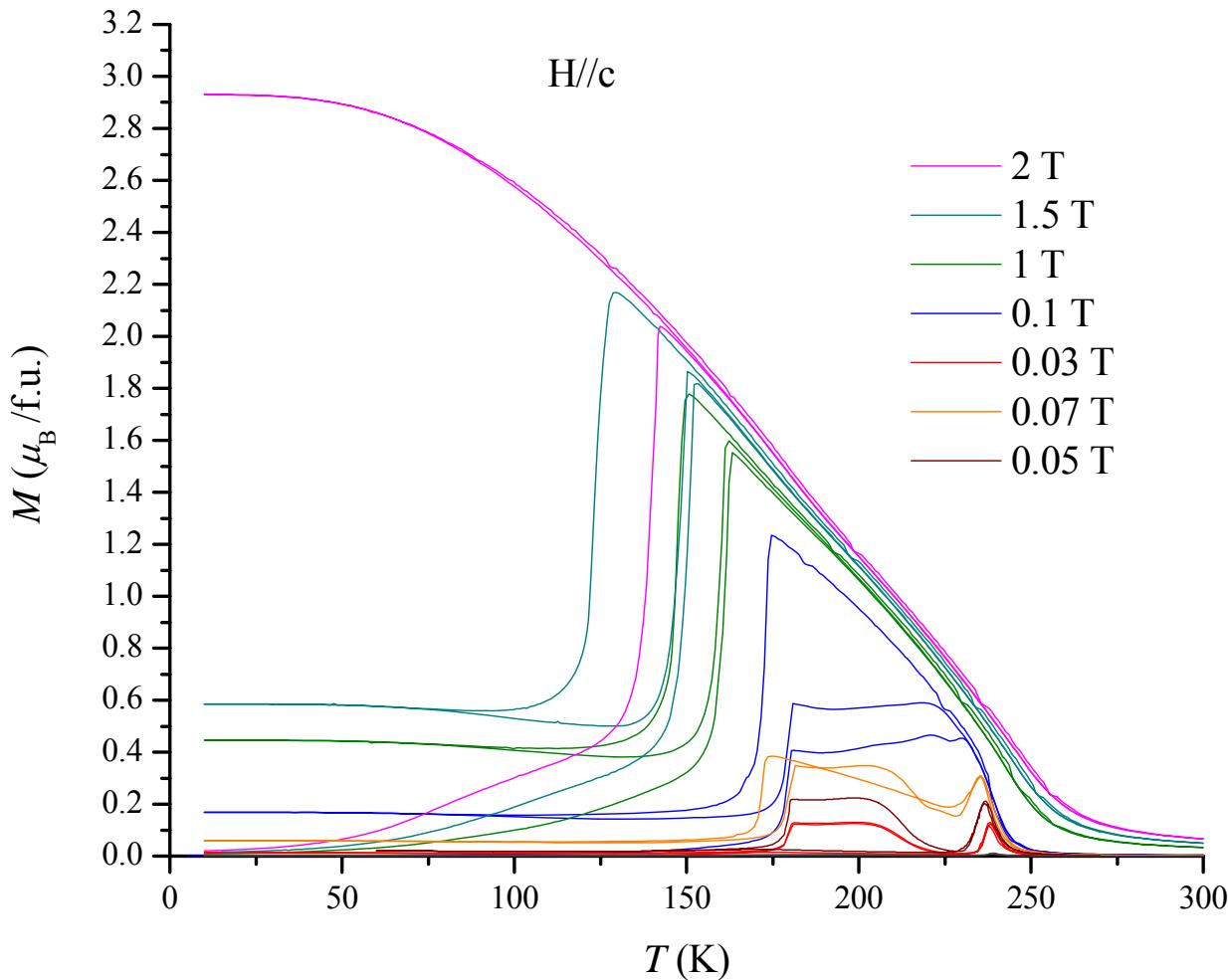


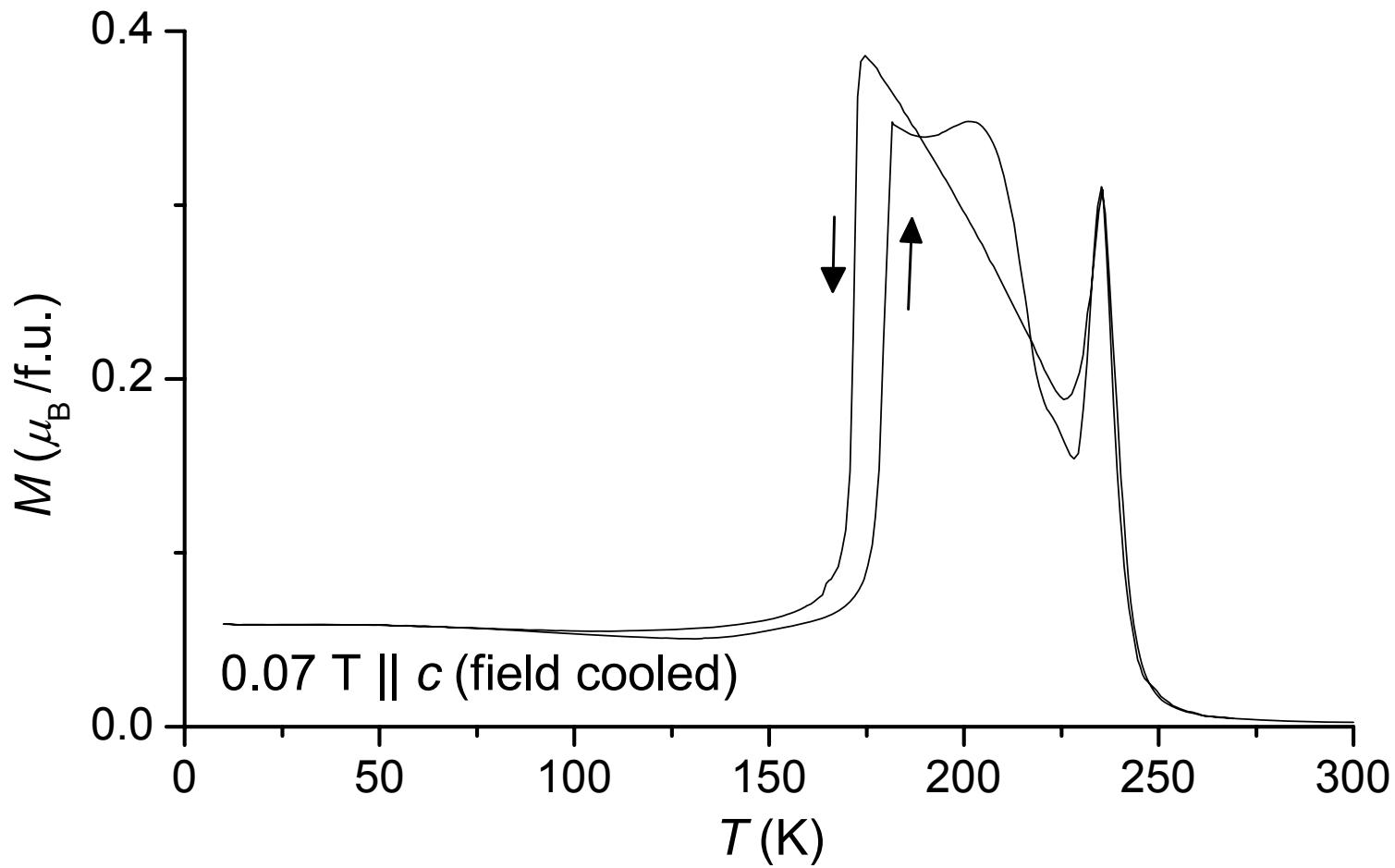
Ikeda, Nohdo, Yamada
JKPS (1998)

LuFe₂O₄ Optical Floating Zone Growth

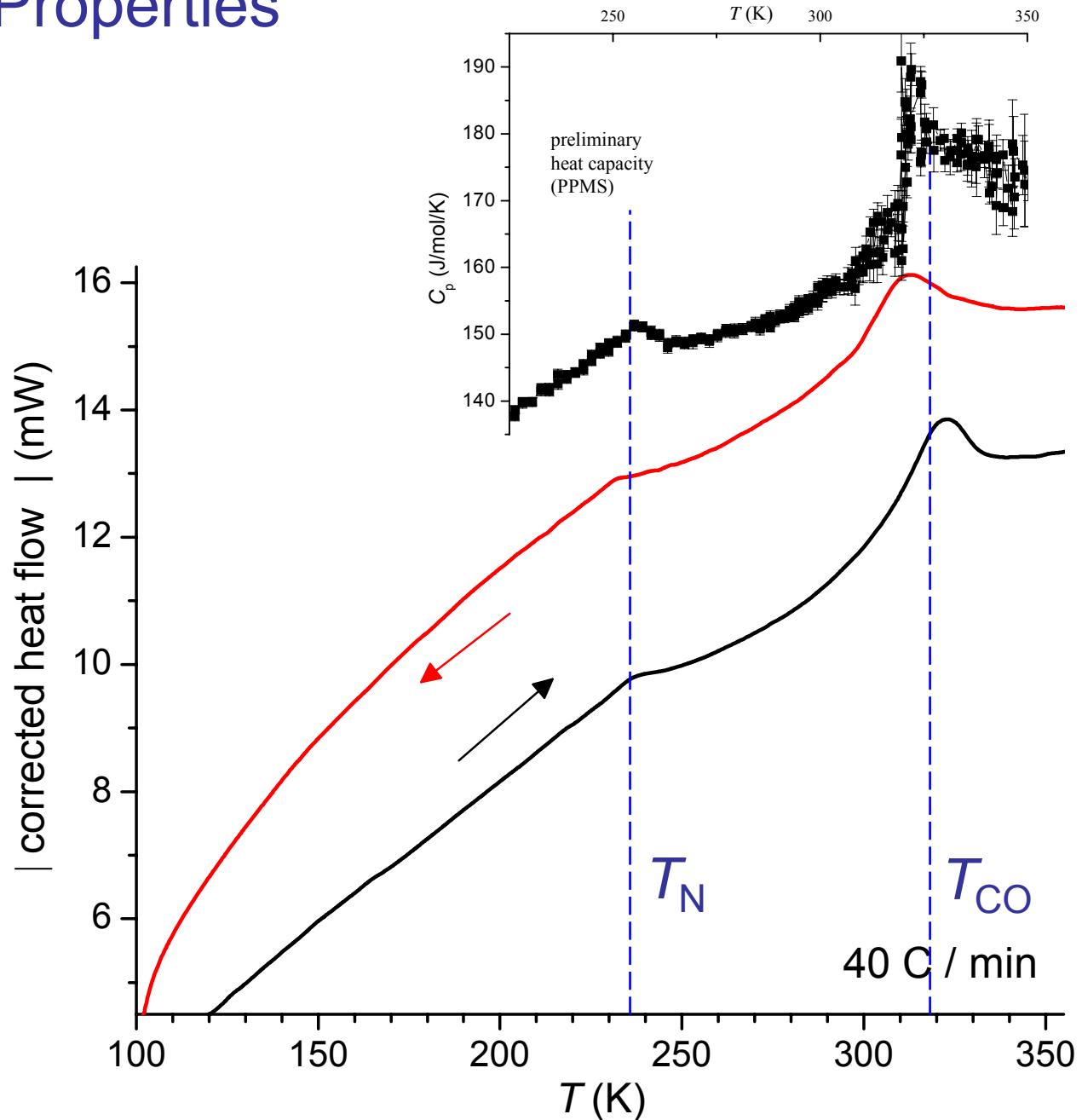


LuFe₂O₄ Magnetic Response





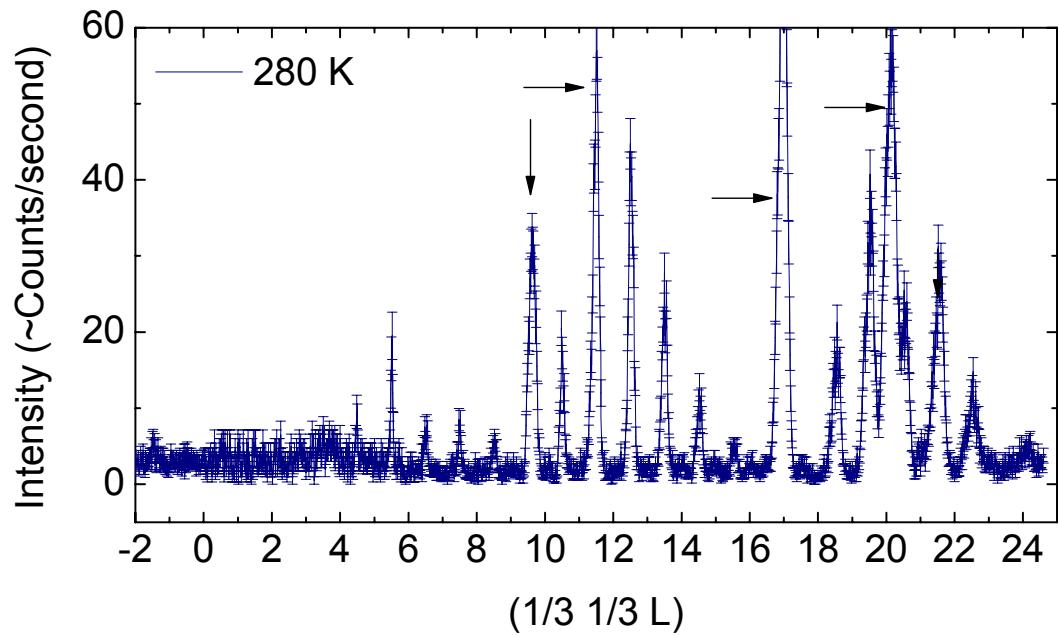
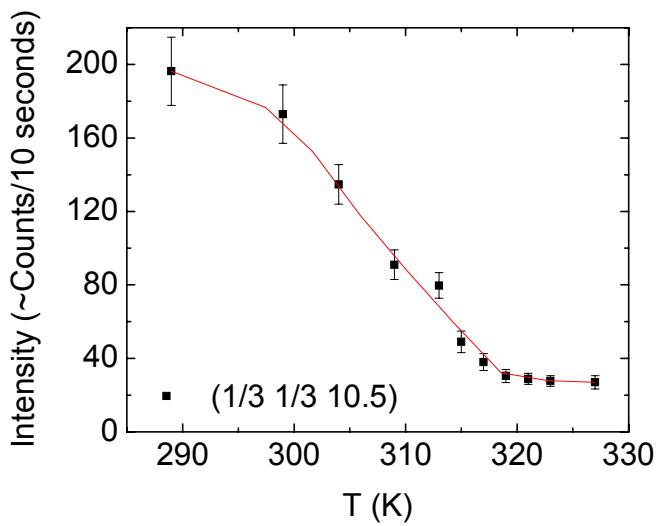
Thermal Properties



3D Charge Order below 320 K

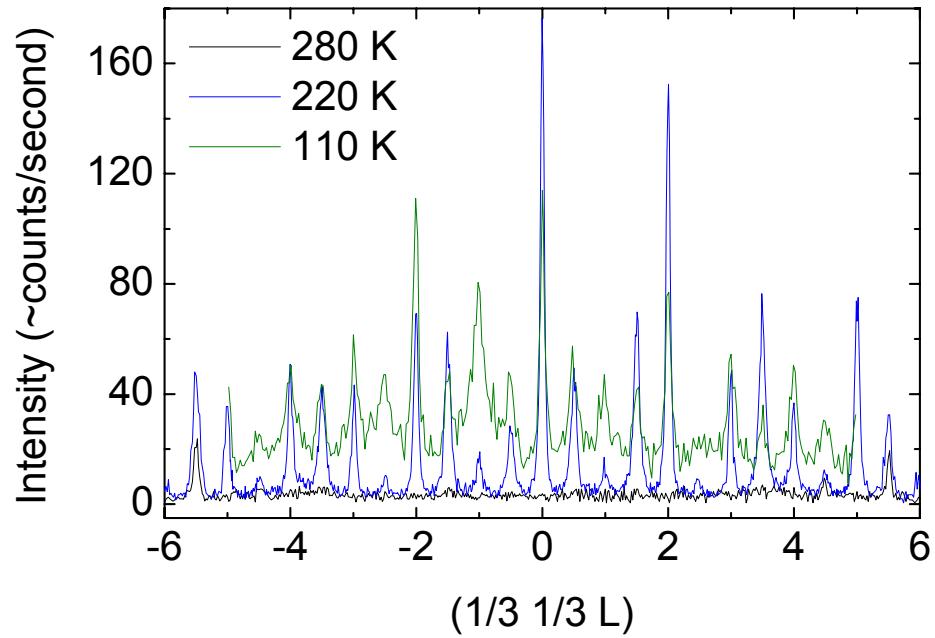
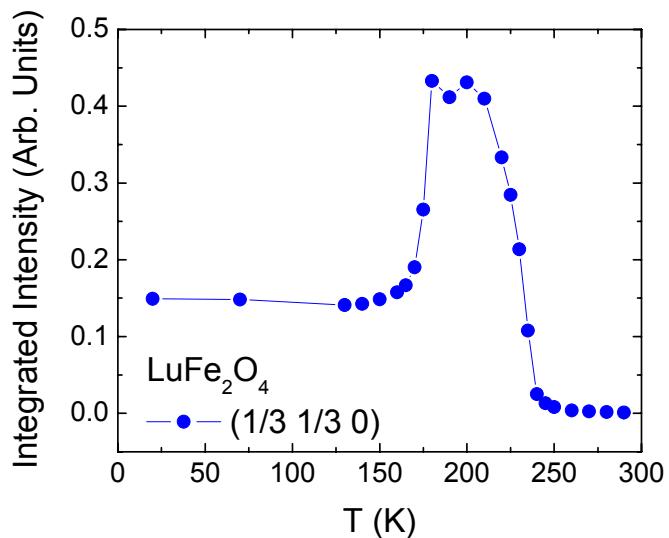


Neutron work by
Andy Christianson
Shull Fellow, ORNL



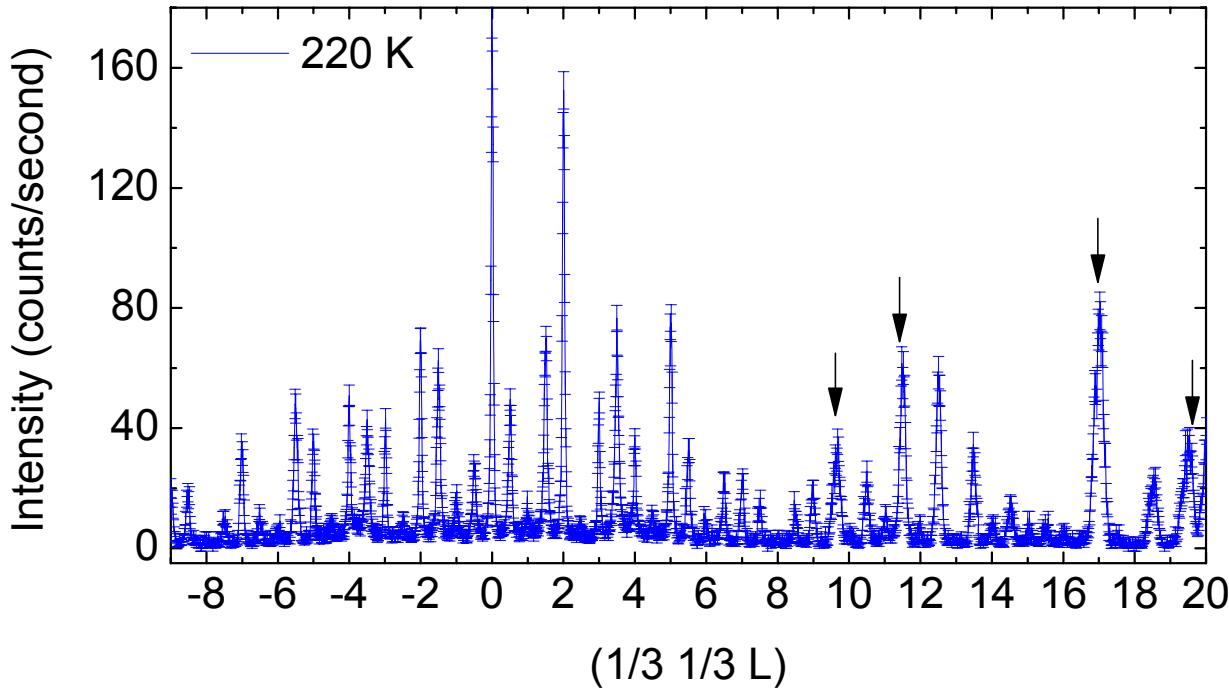
- 3D charge order develops below 320 K and is observed with propagation vector $(1/3 \ 1/3 \ 1/2)$

Magnetic Ordering In LuFe_2O_4



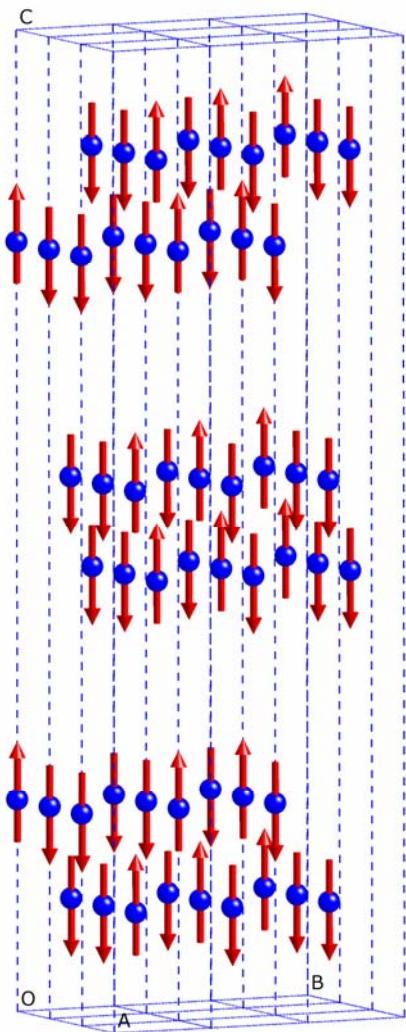
- Two Magnetic transitions-240 K (previously observed) and 180 K (new).
- Both transitions are 3D, however the peaks are not resolution limited (indicating a finite correlation length).
- Intensity occurs on peaks indexed as $(1/3 \ 1/3 \ L)$ where L is either an integer or half integer. The L integer peaks are the fundamental magnetic peaks and the half integer peaks are a consequence of the charge order at 320 K.

Magnetic Structure at 220 K

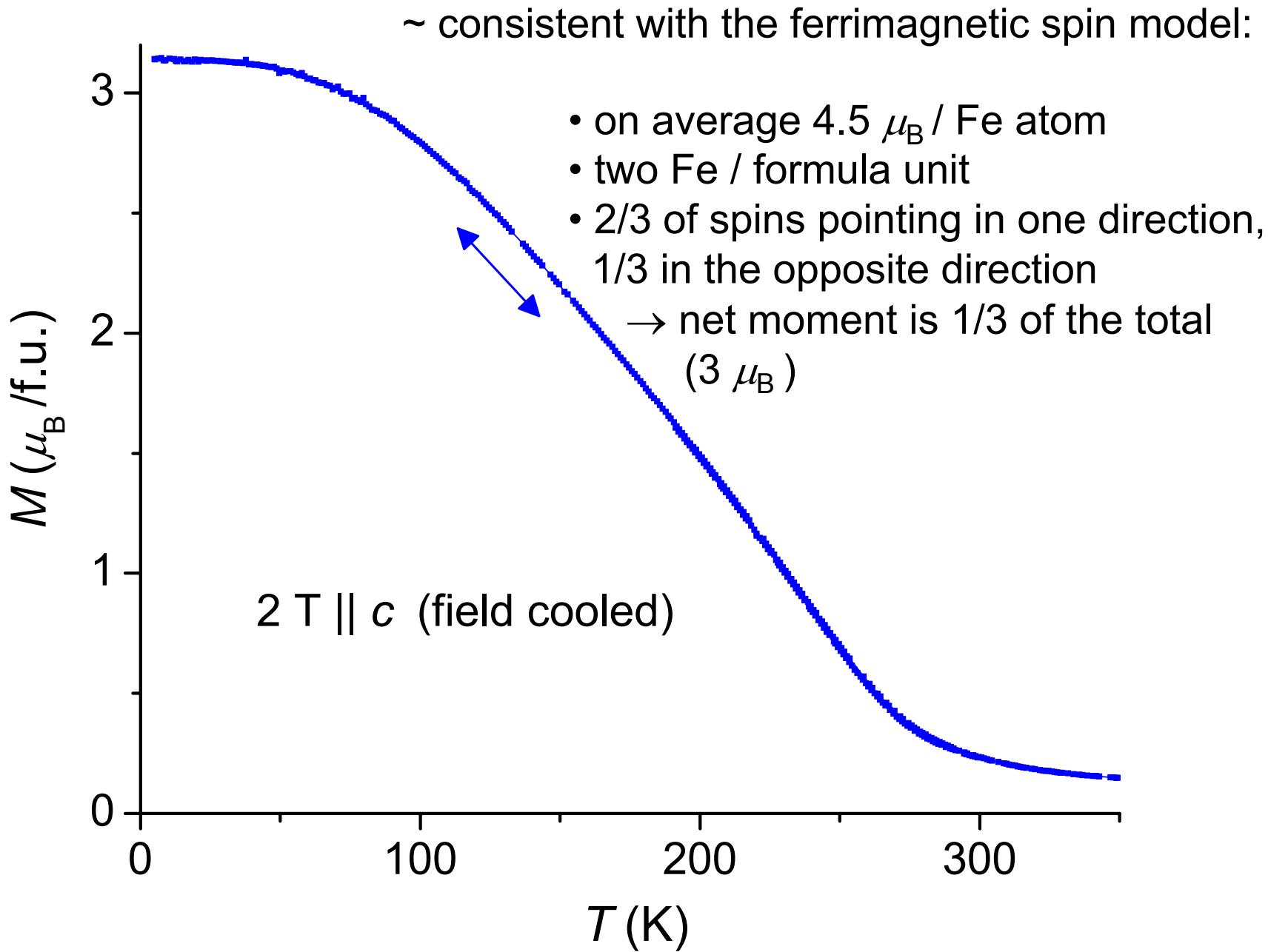


- Peaks at Integer L are purely magnetic while peaks at $\frac{1}{2}$ integer L are related to both the charge order and magnetic order.
- The magnetic contribution to the $L=1/2$ integer peaks arises solely due to the charge order.

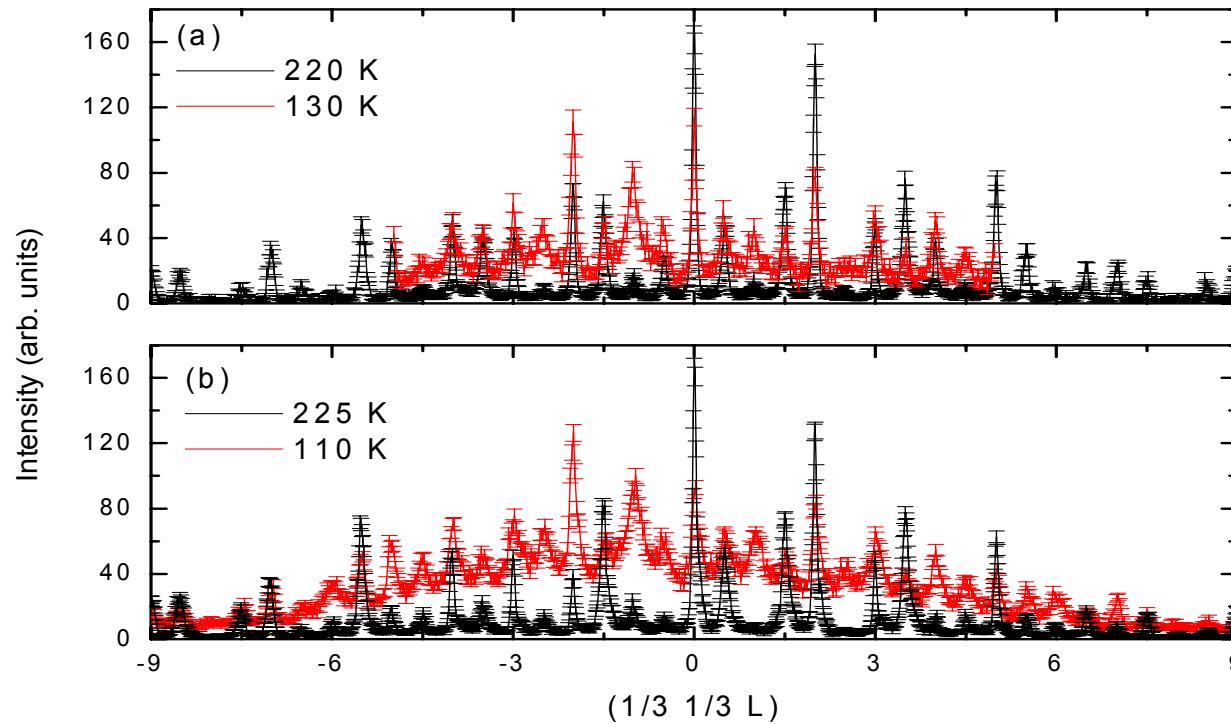
Magnetic Structure of LuFe₂O₄ at T = 220 K



- Spins point along c-axis
- Ordering wavevector (1/3, 1/3, 0)
- 3 symmetry equivalent mag. structures
- 2/3 moments point up, 1/3 down
- Consistent with magnetization

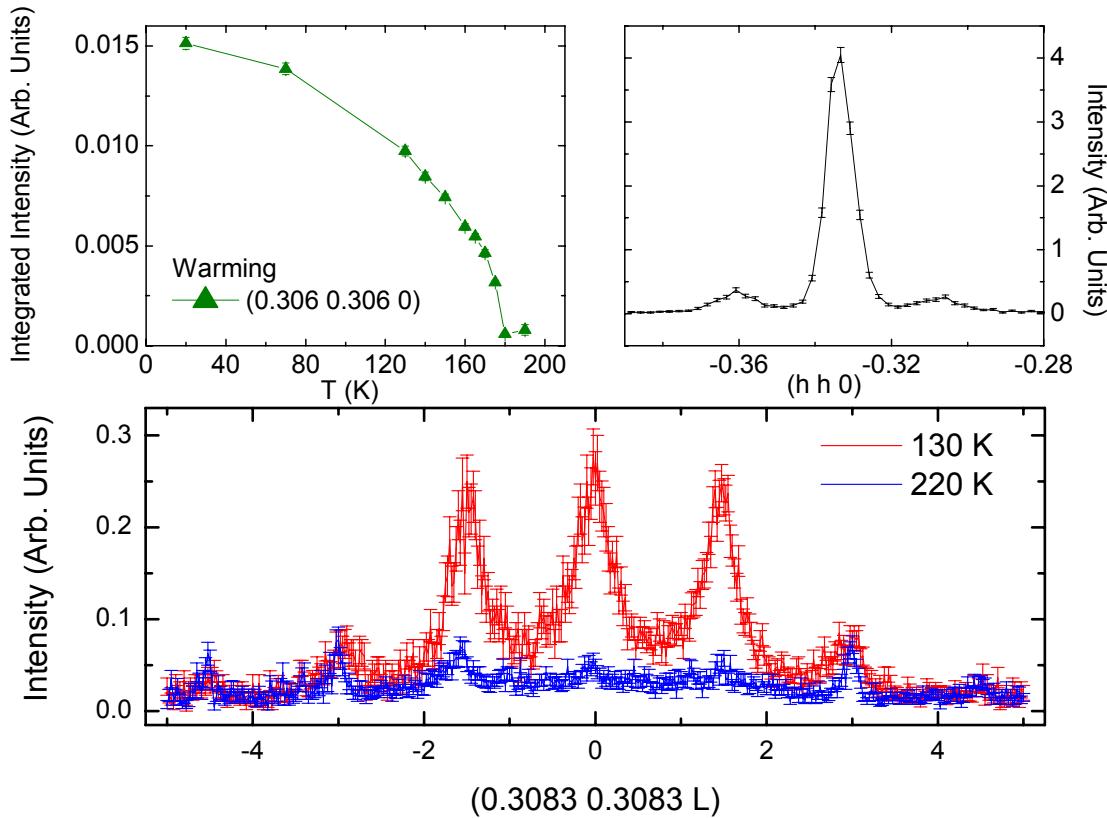


180 K Transition



- Significant broadening occurs in magnetic peaks below 180 K.
- A 2D component to the scattering becomes easily observable below 180 K, however, this appears to be the remnants of unordered parts of the sample as this 2D diffuse scattering is stronger in samples with worse oxygen stoichiometry (comparison above). This is consistent with the old work which shows only 2D scattering.

Low Temperature Satellites

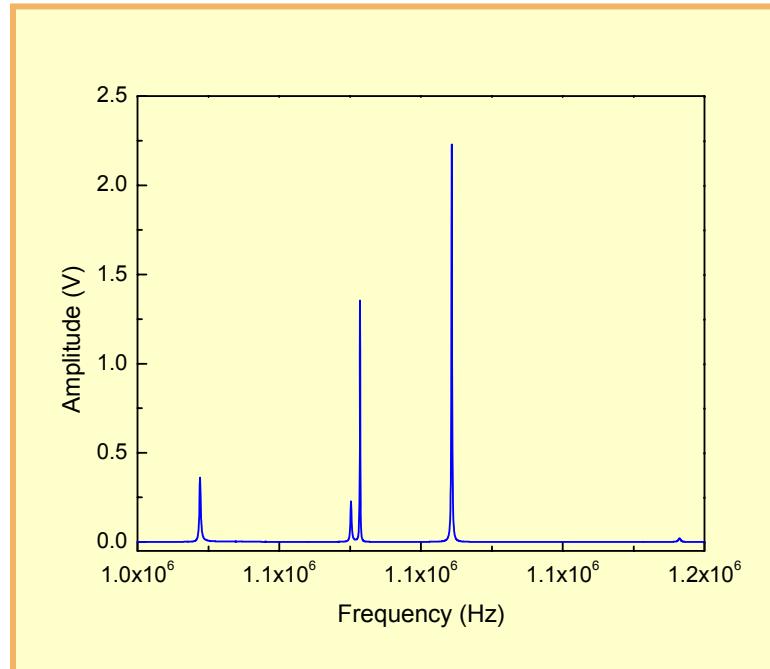
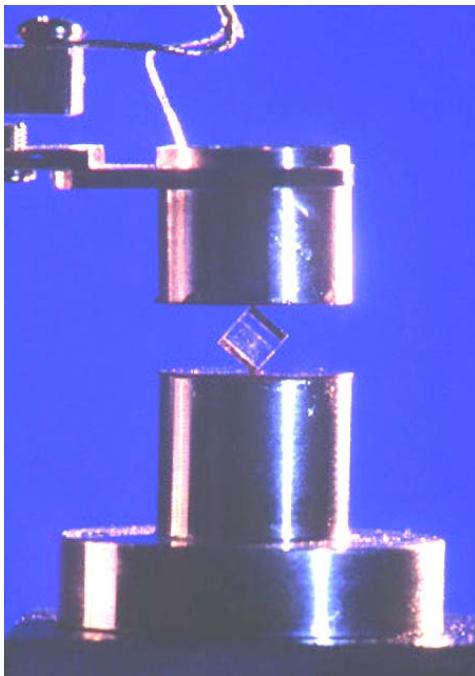


- New satellites appear below 180 K. At present the origin is unknown, however they appear to not be magnetic in origin.



advantages of RUS: **all elastic constants can be obtained in one measurement**
small samples (mm³)

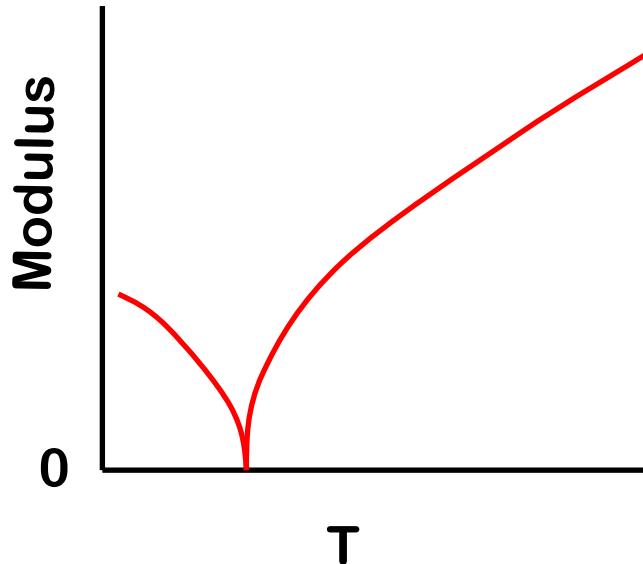
Veerle Keppens



Strain-Order Parameter Coupling at Phase Transitions

Bi-Linear coupling:

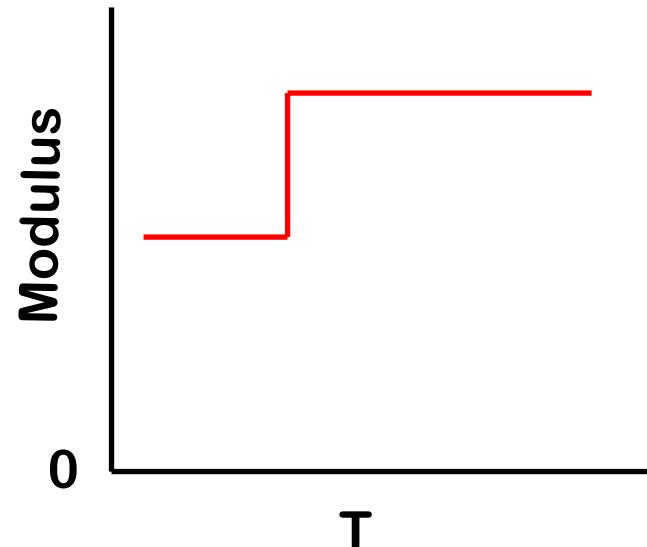
$$F_c = \beta Q \varepsilon$$



$$c = c^0 - \frac{a}{T - \theta}$$

Quadratic coupling:

$$F_c = \beta Q^2 \varepsilon$$



$$c = c^0 \quad T > T_c$$

$$c = c^0 - \frac{2h^2}{\beta} \quad T < T_c$$

Magnetite Displays Bi-Linear Coupling at Verwey Transition

88

T. Goto and B. Lüthi

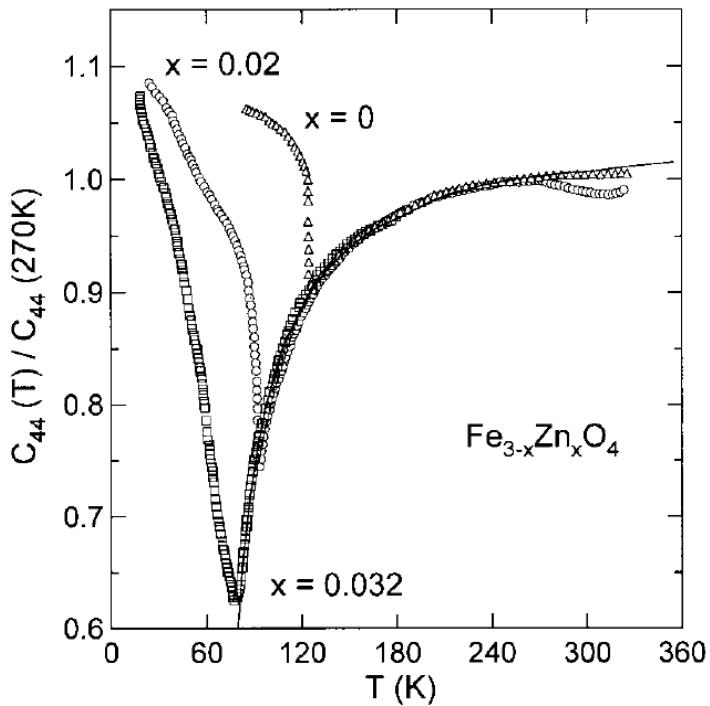
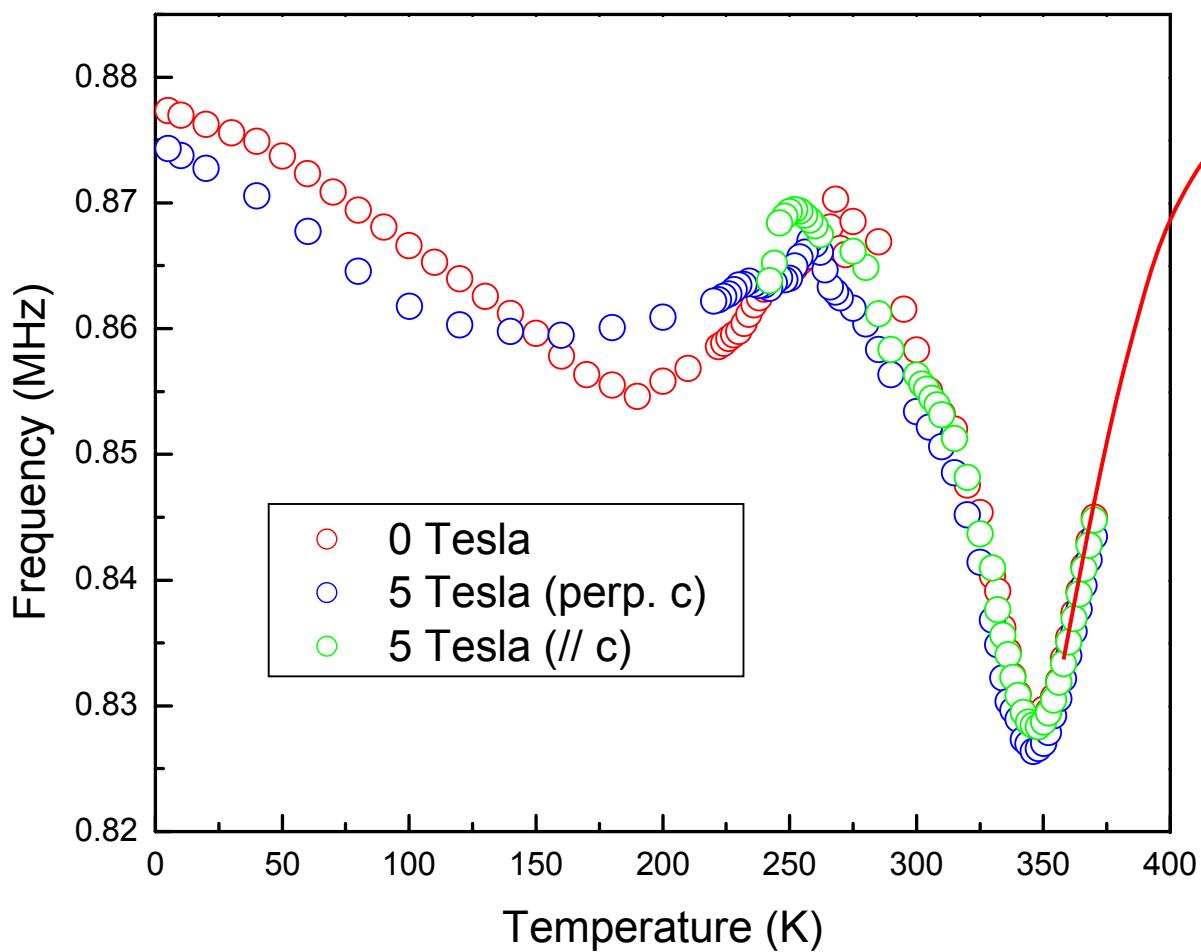


Figure 10. Normalized soft modes $c_{44}(T)/c_{44}(270\text{K})$ for the $\text{Fe}_{3-x}\text{Zn}_x\text{O}_4$ ($x = 0$, 0.02 and 0.032) crystals together with a fit (—) discussed in the text.

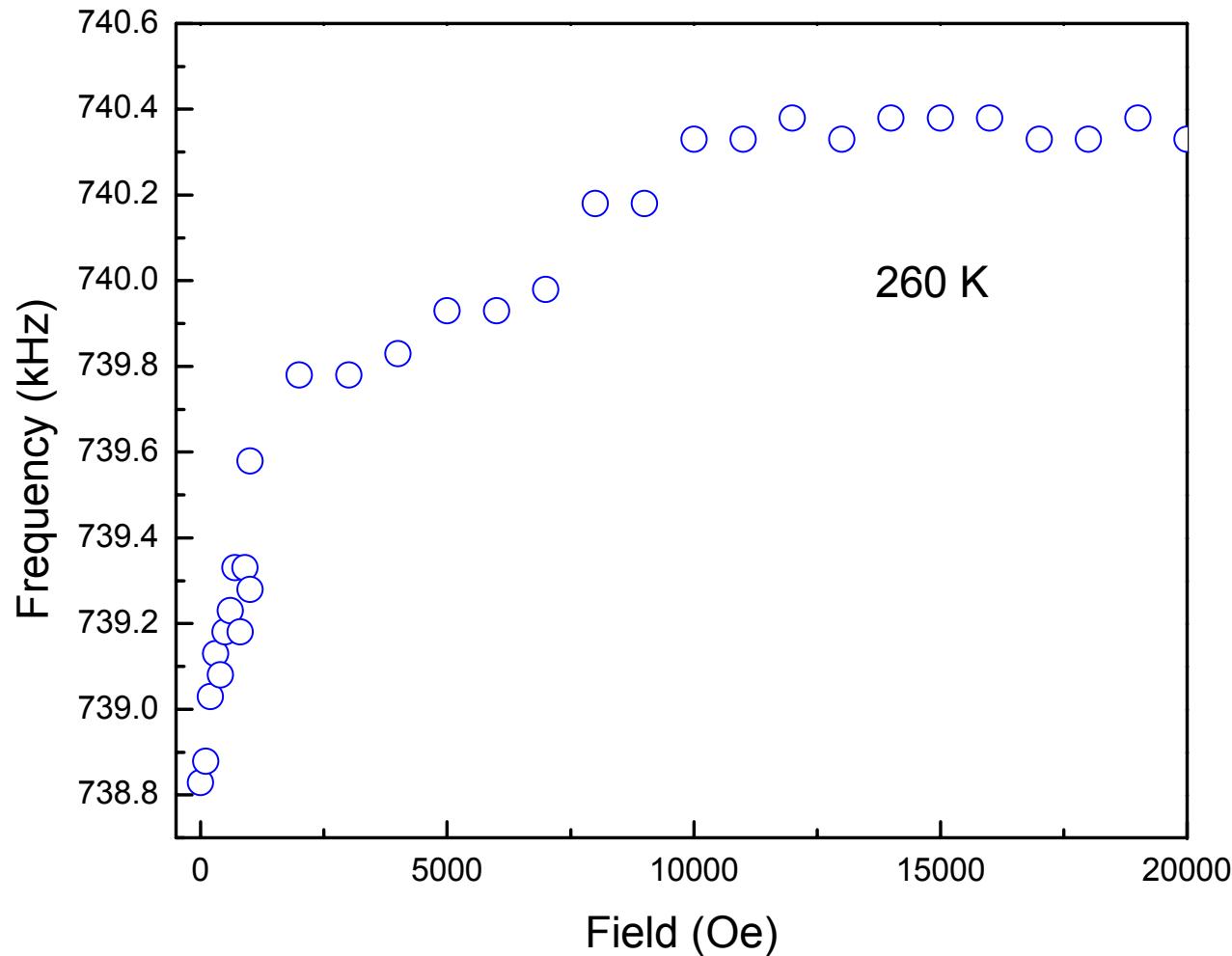
Schwenk et al., Eur. Phys. J. B 13, 491 (2000)

Elastic strain induced by the ultrasonic wave couples to the charge fluctuation modes

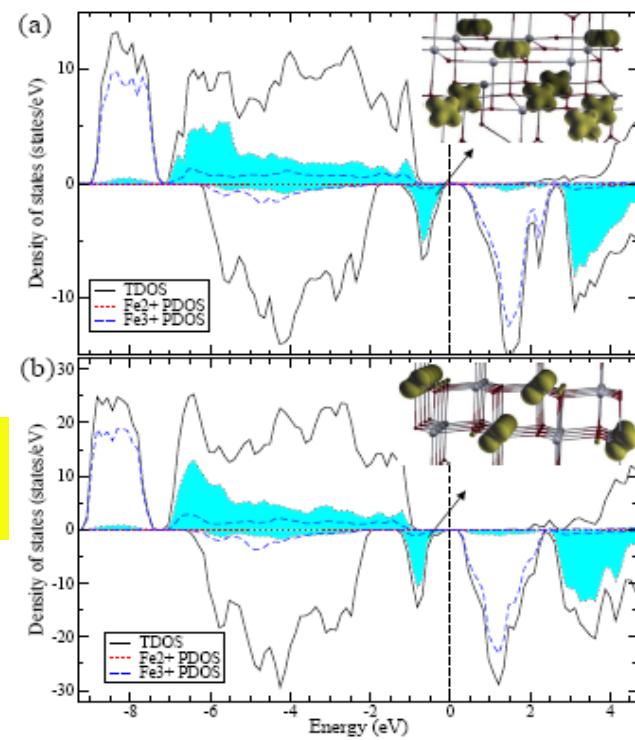
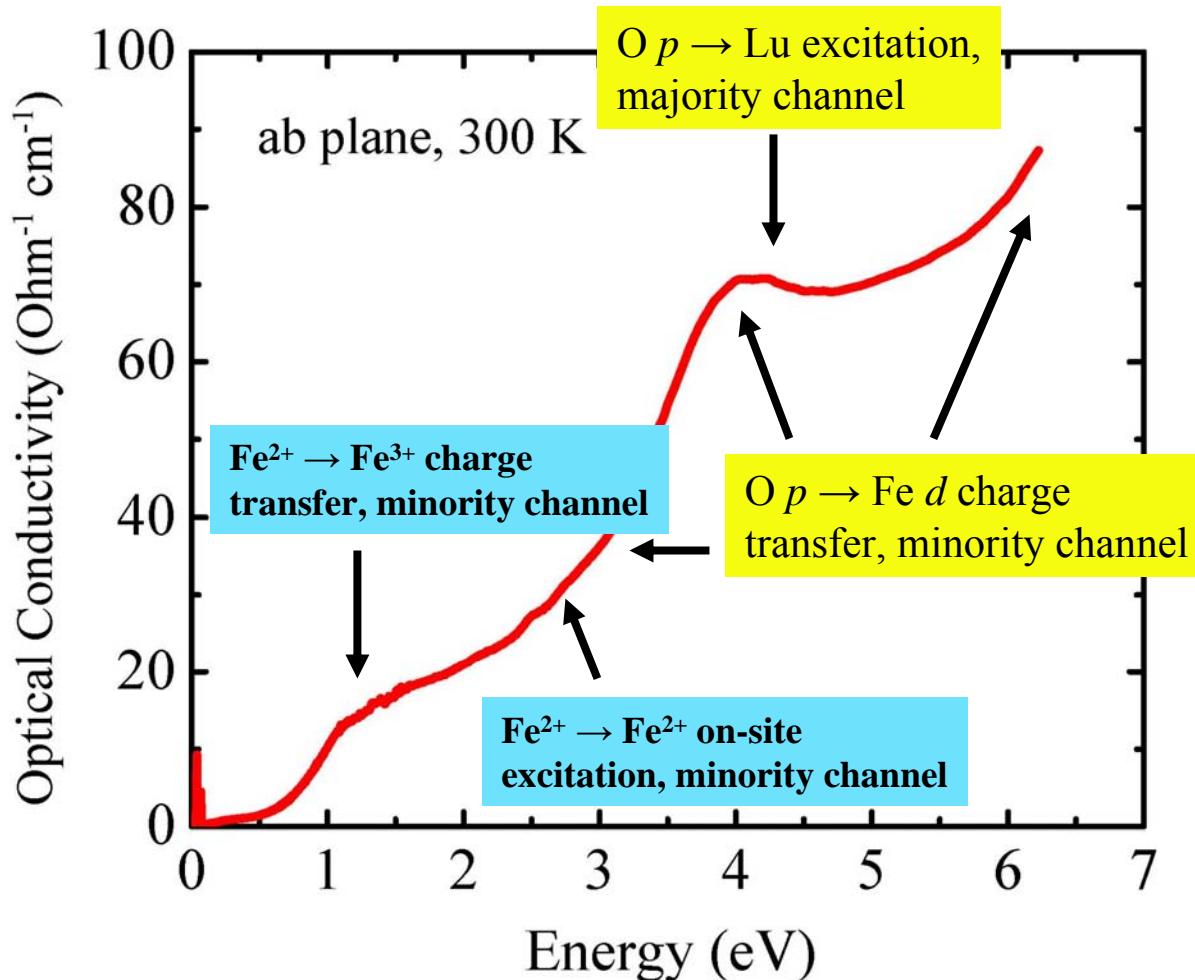
LuFe₂O₄ RUS Measurements in Magnetic Field



RUS Measurements in Magnetic Field



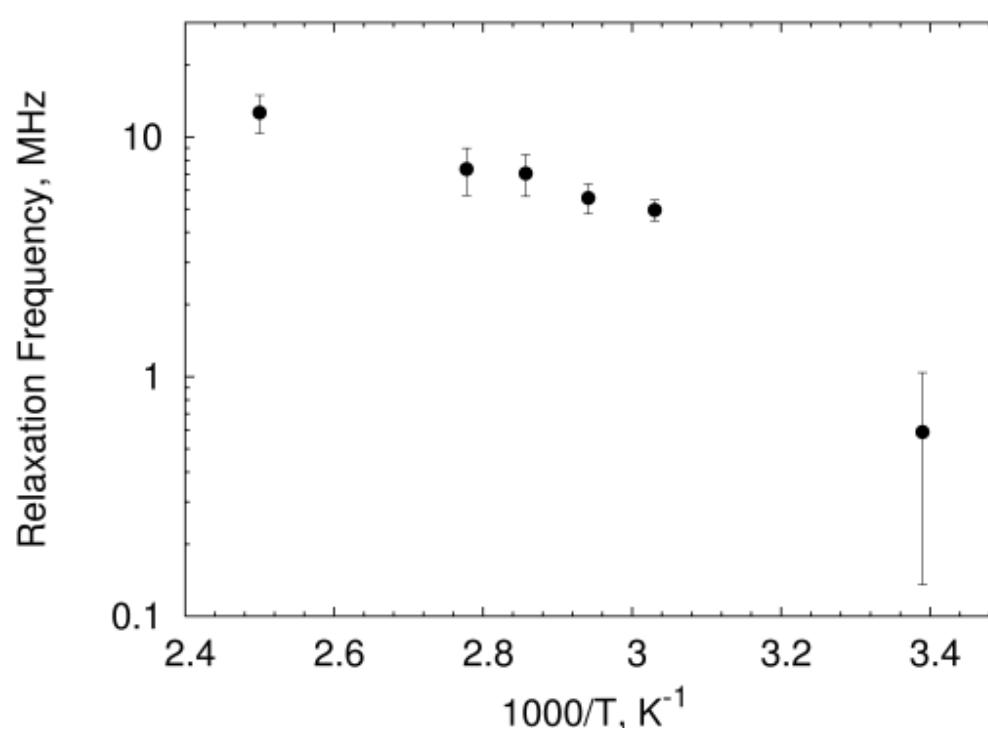
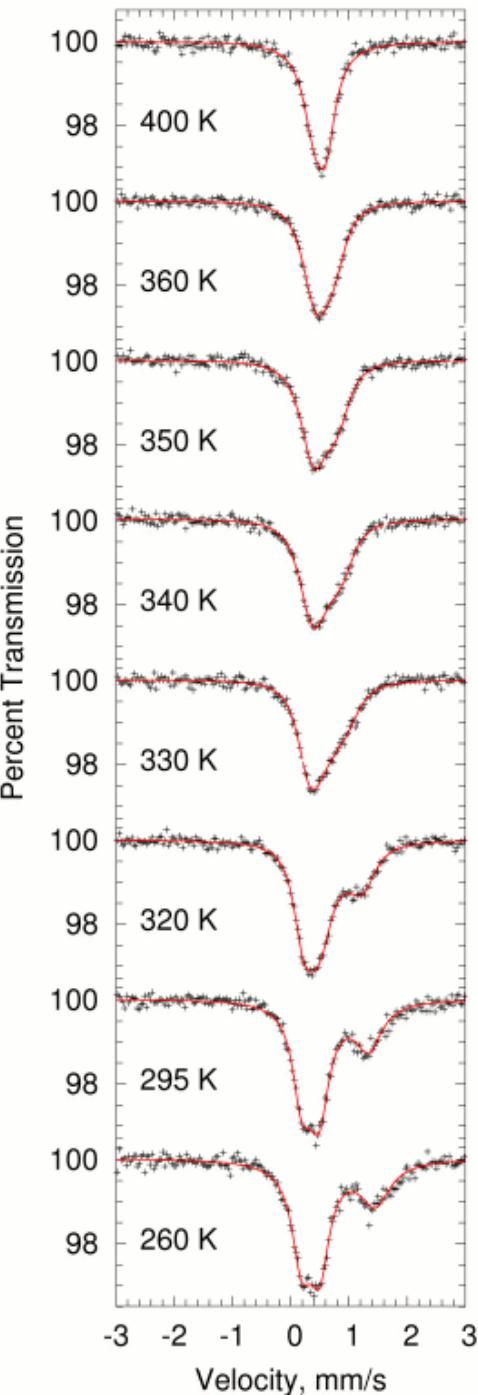
Optical Properties of LuFe₂O₄: Musfeldt Group, University of Tennessee



Xiang and Whangbo, PRL (2007)

- Optical properties are dominated by excitations in the minority spin channel
- $\sqrt{3} \times \sqrt{3}$ and chain structures have similar density of states pictures. Possibility of extracting domain distribution information via temperature dependence of 1.5 eV excitation.
- Magneto-optical work in progress.

Mössbauer Spectroscopy (R. Hermann)



Assuming 50% Fe^{2+} and 50% Fe^{3+} is compatible with the spectra.

Upon heating, the charge order vanishes with an increase in the electron hopping frequency.

Below 260 K the hopping is slower than resolvable, ie than 0.1 MHz.

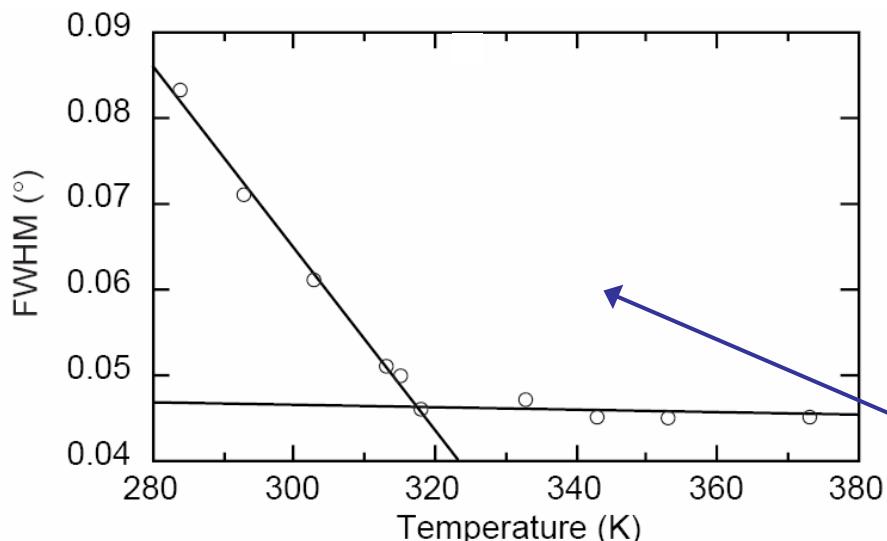
Fe_2OBO_3

Warwickite type $\text{M}^{\text{I}}\text{M}^{\text{II}}\text{OBO}_3$ usually orthorhombic

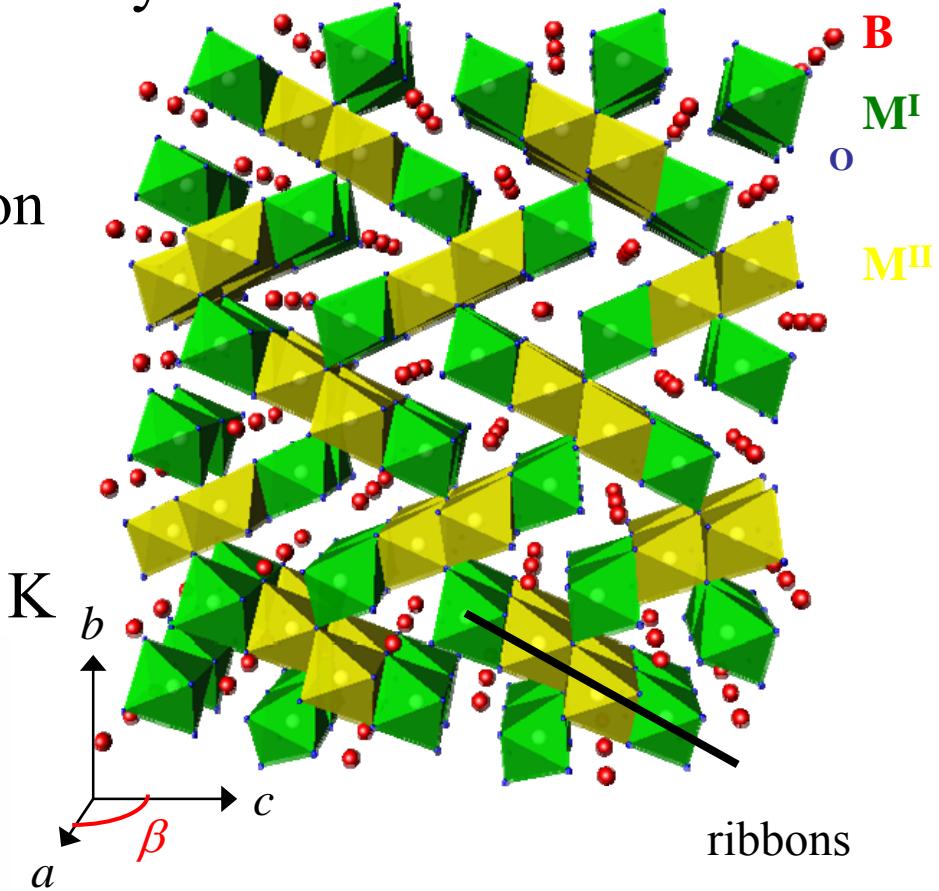
Two structurally distinct metal sites with octahedral coordination
(I, II)



→ monoclinic distortion below 317 K



[Attfield *et al.*, Nature 396, 655 (1998)]

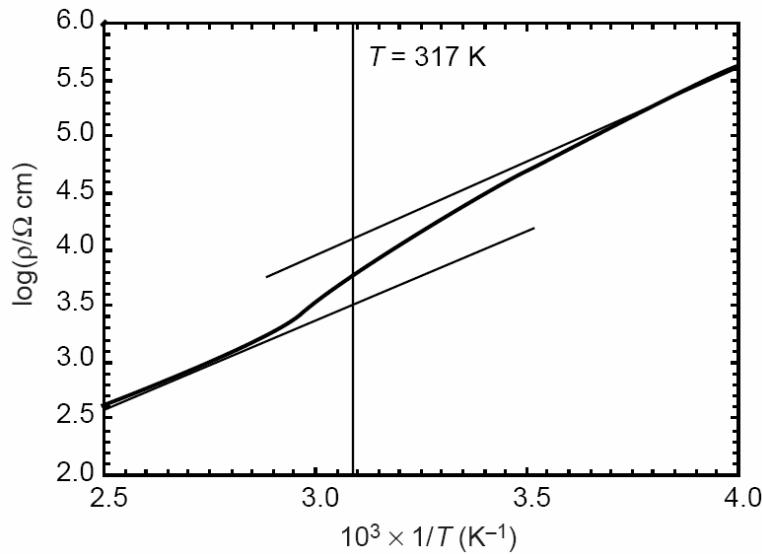


Half-width of the combined
monoclinically split
(102)/(-102) doublet

Fe_2OBO_3

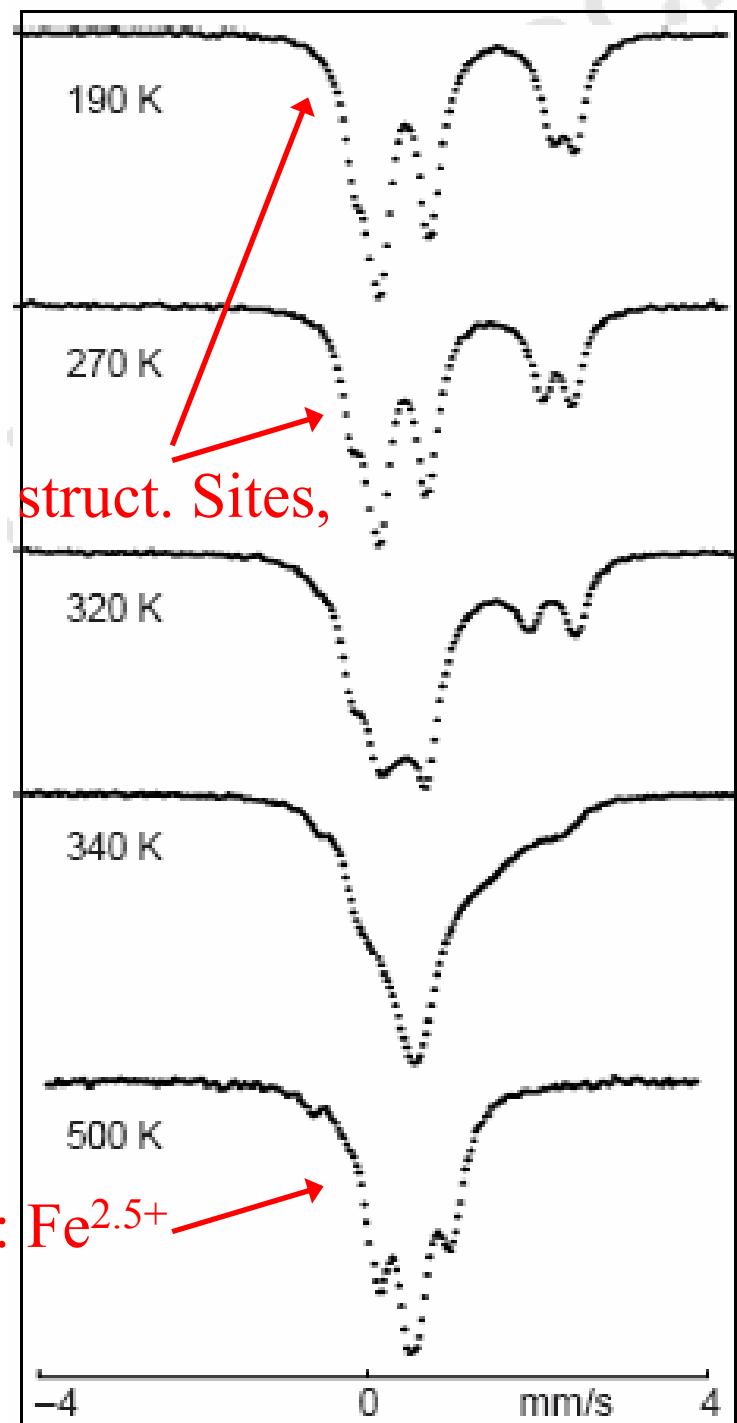
Charge order surmised based on Mössbauer spectroscopy and resistivity:

Four doublets assigned to $\text{Fe}^{2+}/\text{Fe}^{3+}$ on two struct. Sites, no discernible hopping



Rapid hopping on two structural sites: $\text{Fe}^{2.5+}$

[Attfield *et al.*, Nature 396, 655 (1998)]

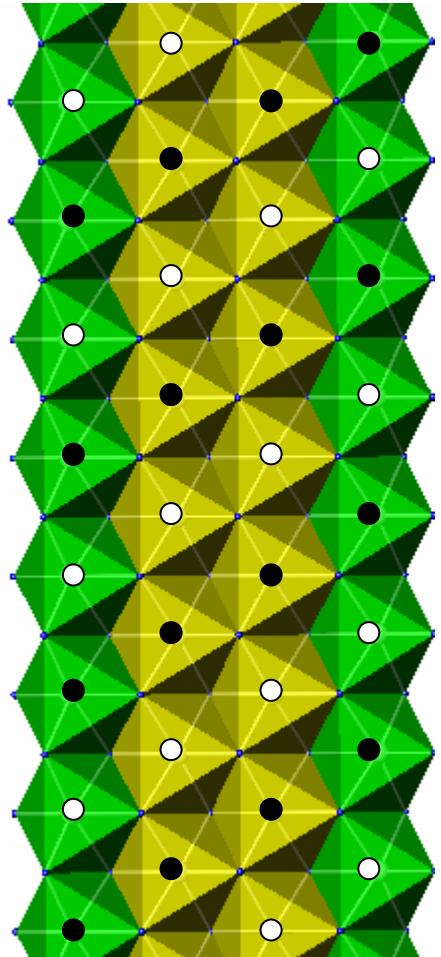
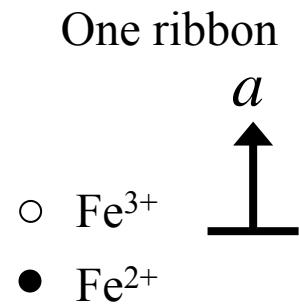


Fe_2OBO_3

Should lead to unit cell doubling in a direction
[weak $(h+0.5 \ k \ l)$ reflections]

not observed (on polycrystalline material)

- No superstructure observed &
- Resistivity feature associated with CO very broad



possible charge order model
[Attfield *et al.*, Nature **396**, 655 (1998)]

→ Single-crystal study highly desirable

Single crystals of Fe_2OBO_3

Try flux growth in Fe – B – O system

Starting point : growing FeBO_3 from

Fe_2O_3 – B_2O_3 – PbO – PbF_2 flux

[following Kotrobova *et al.*,

J. Crystal Growth **71**, 607 (1985)]

Then tune growth parameters to favor crystallisation of Fe_2OBO_3



Fe_2OBO_3 crystallised as **thin needles** ($\parallel a$) of weight **up to ~ 1 mg**

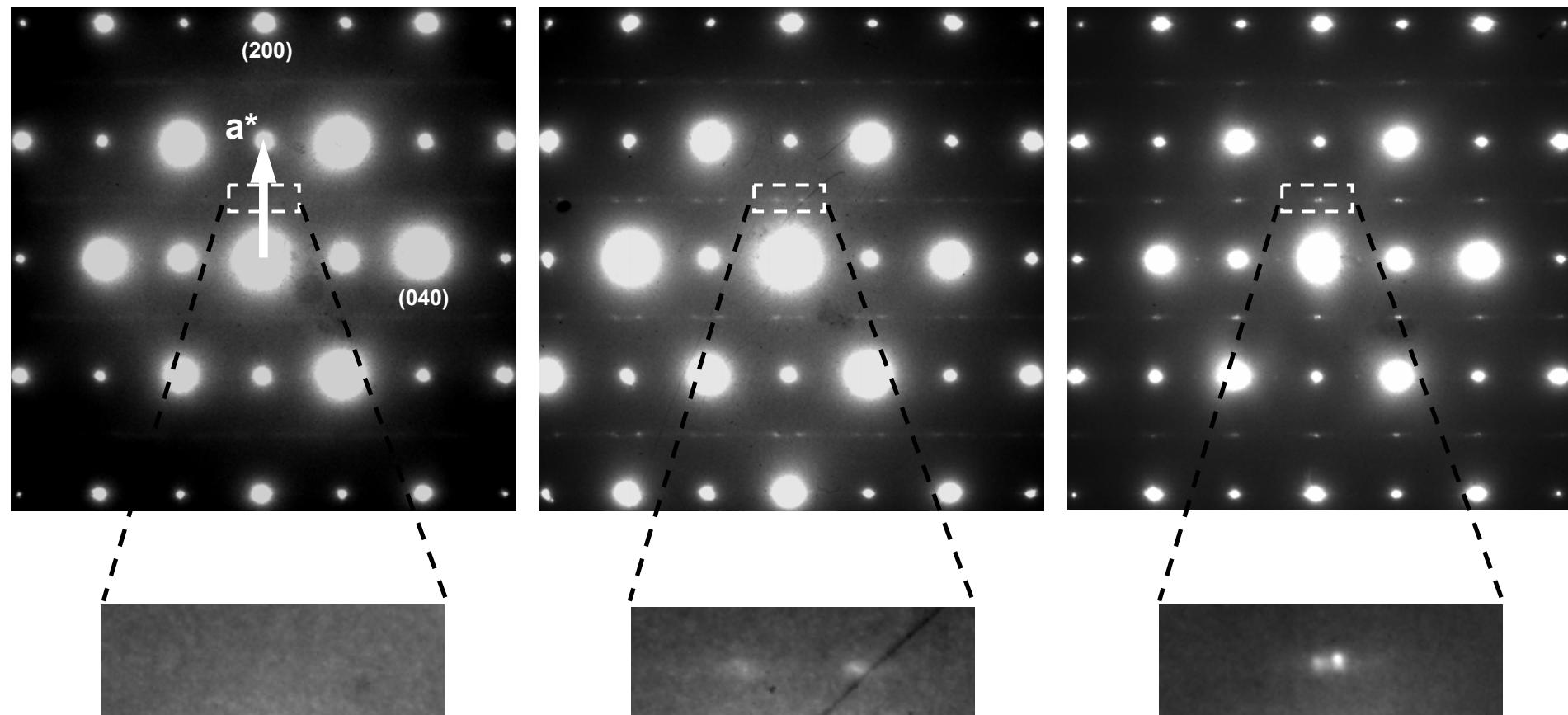
Phase confirmed by powder x-ray diffraction (and various phys. propert.)

Fe_2OBO_3 : electron diffraction ([001] zone)

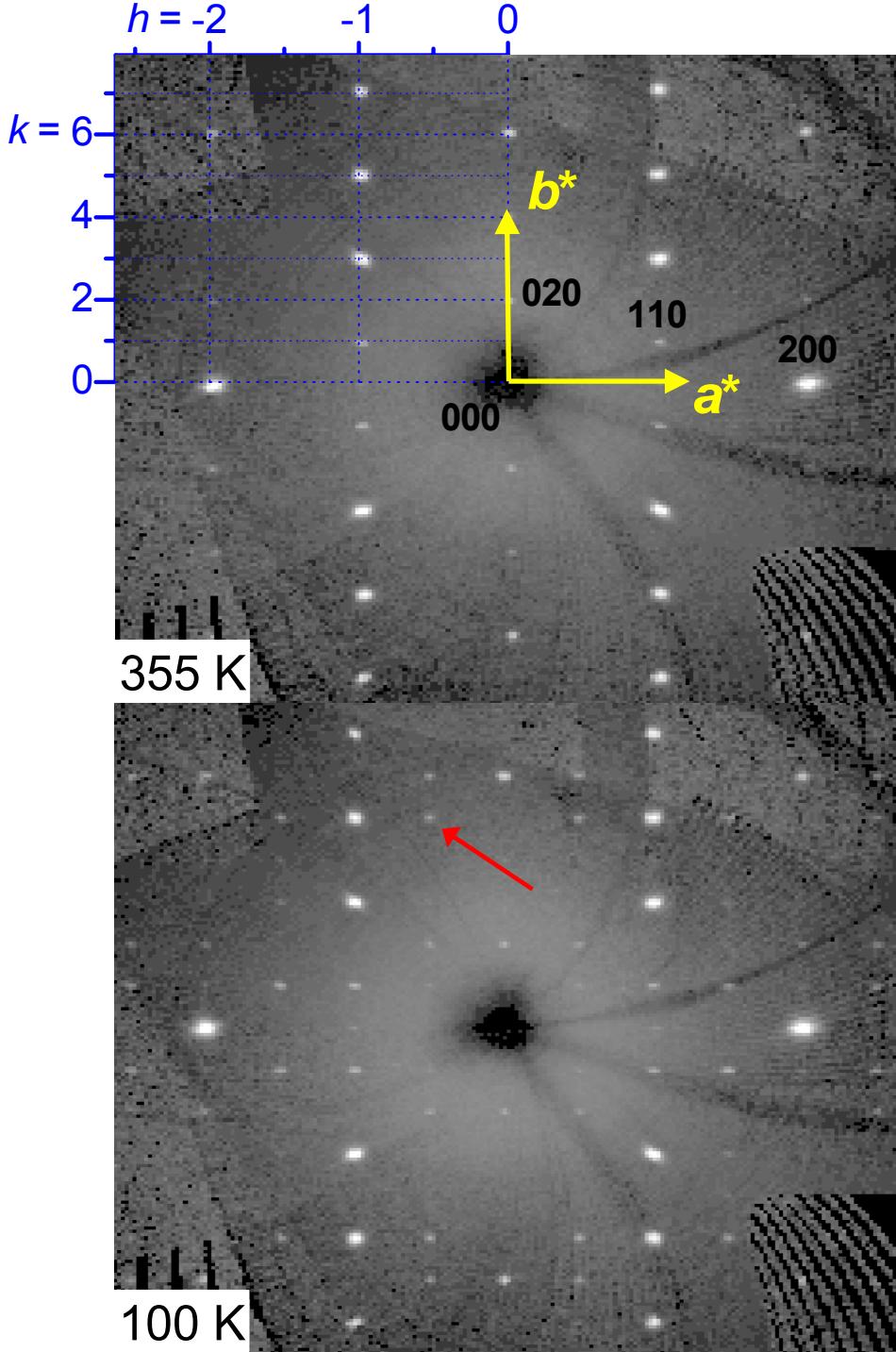
$T = 350 \text{ K}$

$T = 300 \text{ K}$

$T = 117 \text{ K}$

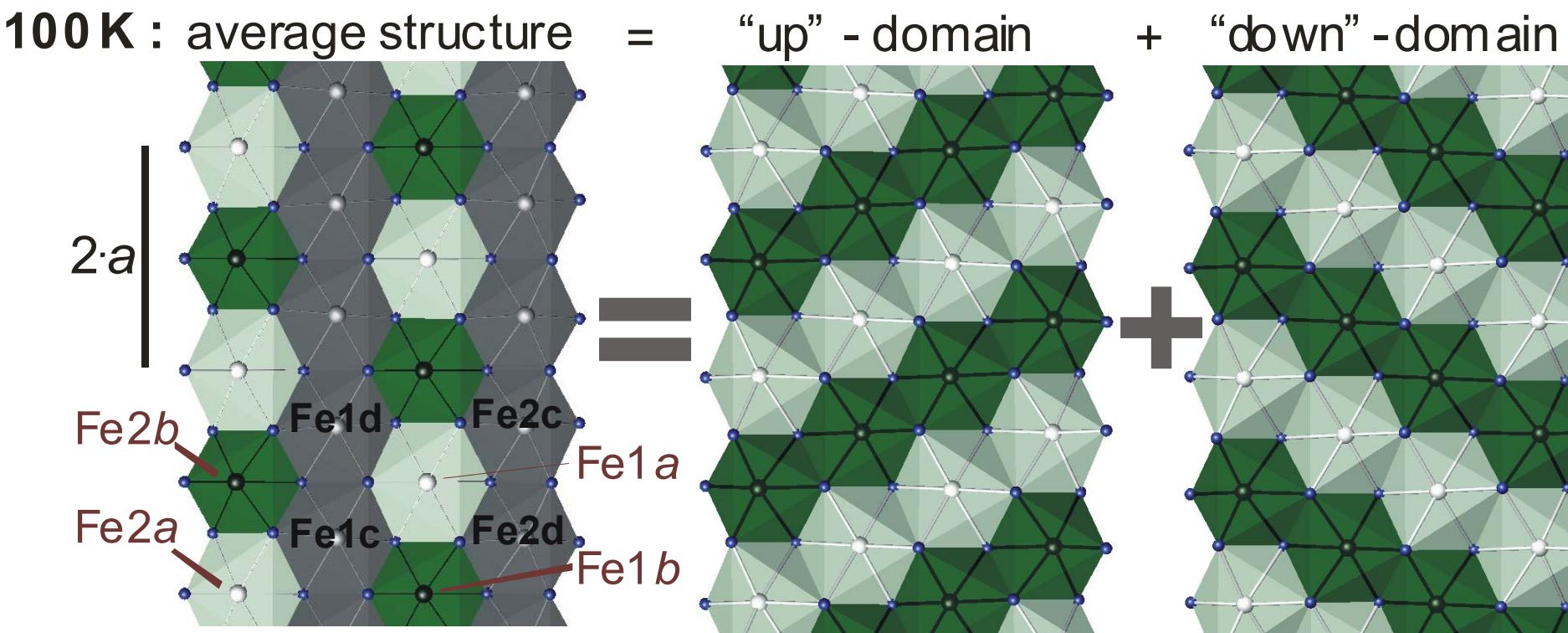


Experiments performed by Jing Tao on Urbana TEM

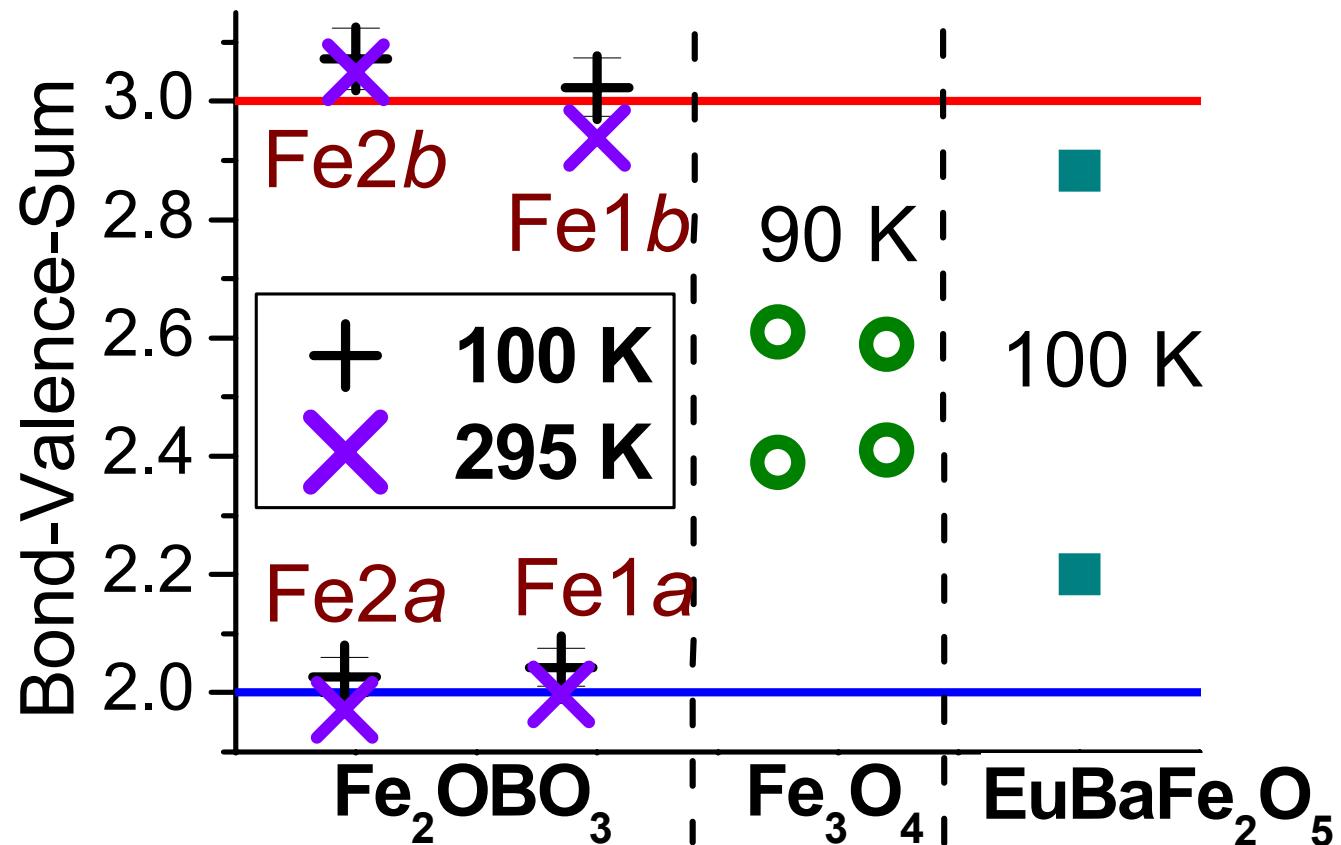


X-ray diffraction:
Doubling of unit cell
along a at low T

Structure refinement



Fe_2OBO_3 seems best example so far of ionic charge order



“...the inductive effect of the antagonistic bond...”

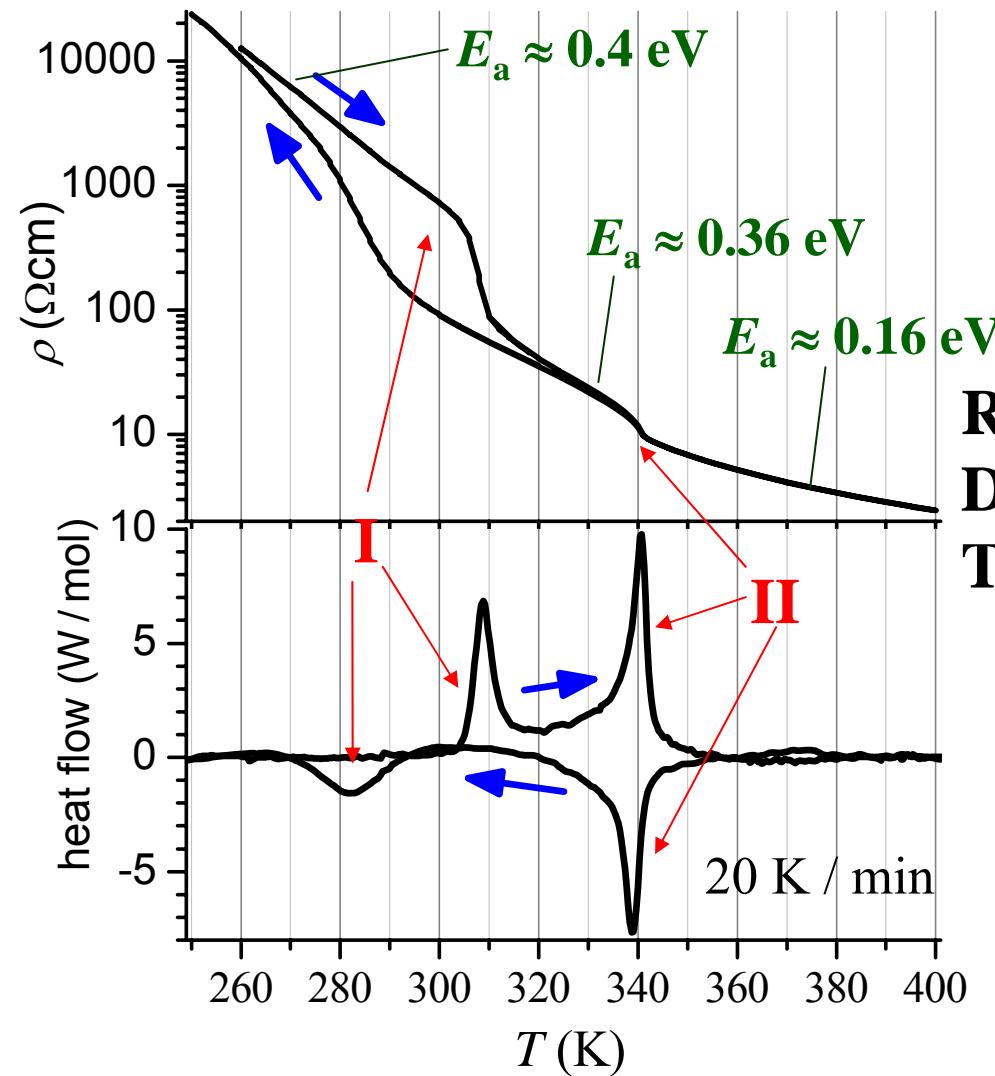
F. Menil, J. Phys. Chem. Solids 46, 763 (1985)

From Art Sleight:

In EuBaFe₂O₅, the impact of the Eu and Ba would be to increase Fe-O covalency relative to Fe₃O₄. In Fe₂OBO₃, the impact of B would be to increase ionicity relative to Fe₃O₄. And I am convinced that it is this increased ionicity that gives basically the ideal values in Fe₂OBO₃.



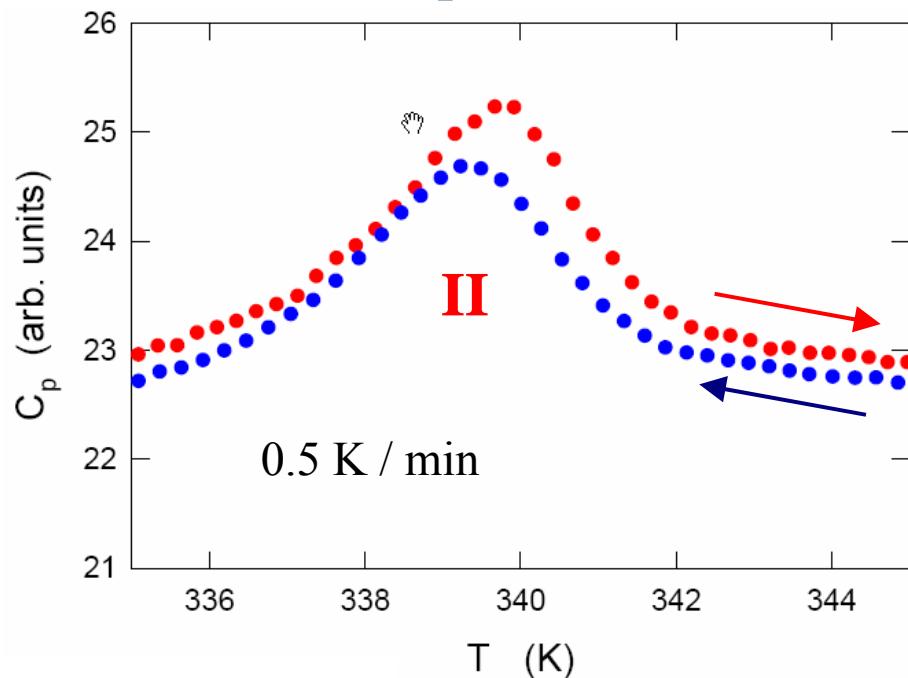
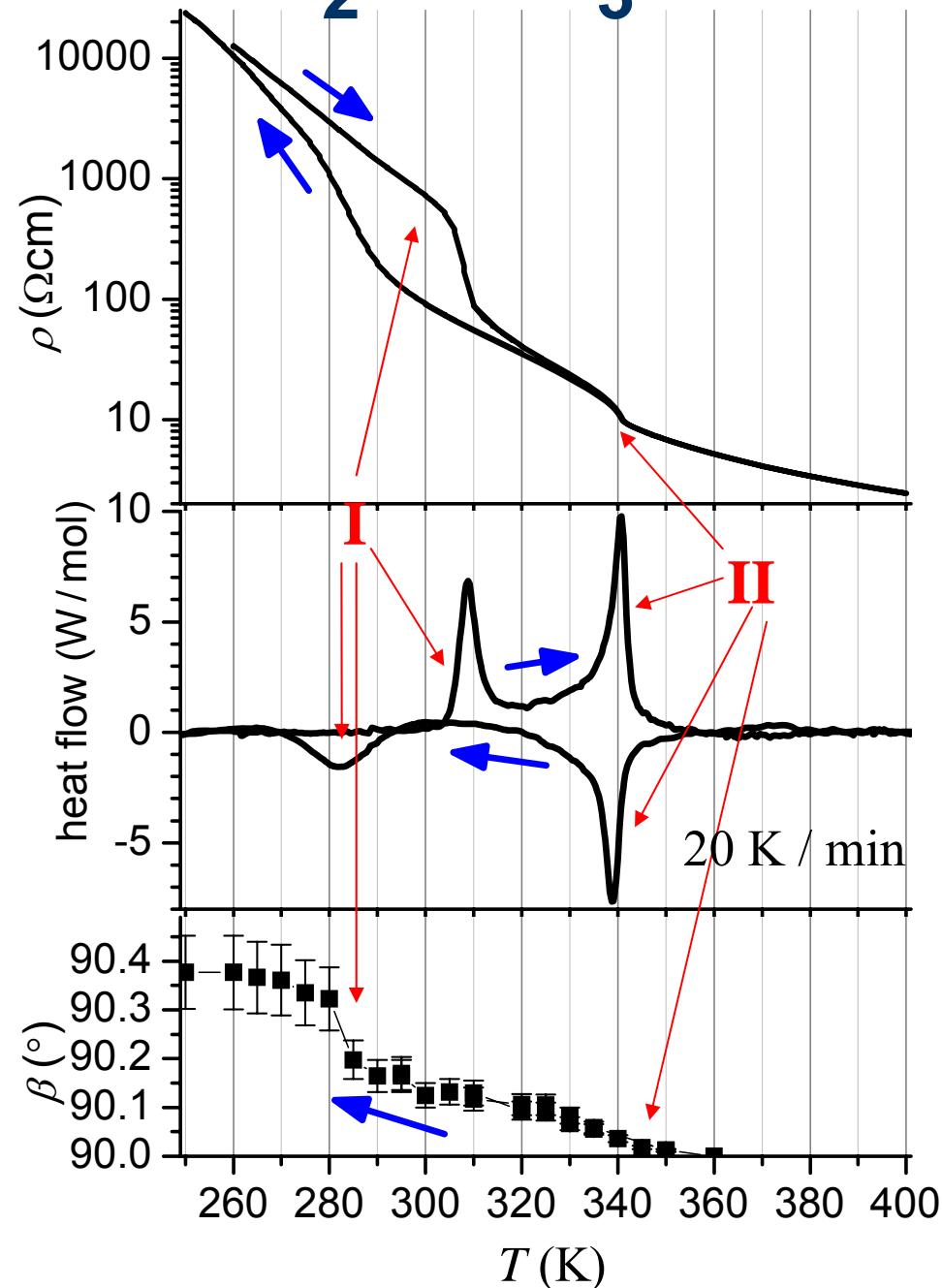
Fe_2OBO_3 : intermediate phase



Activated transport :
 $(\rho \propto e^{E_a/kT})$

Resistivity and
Differential scanning calorimetry :
Two phase transitions

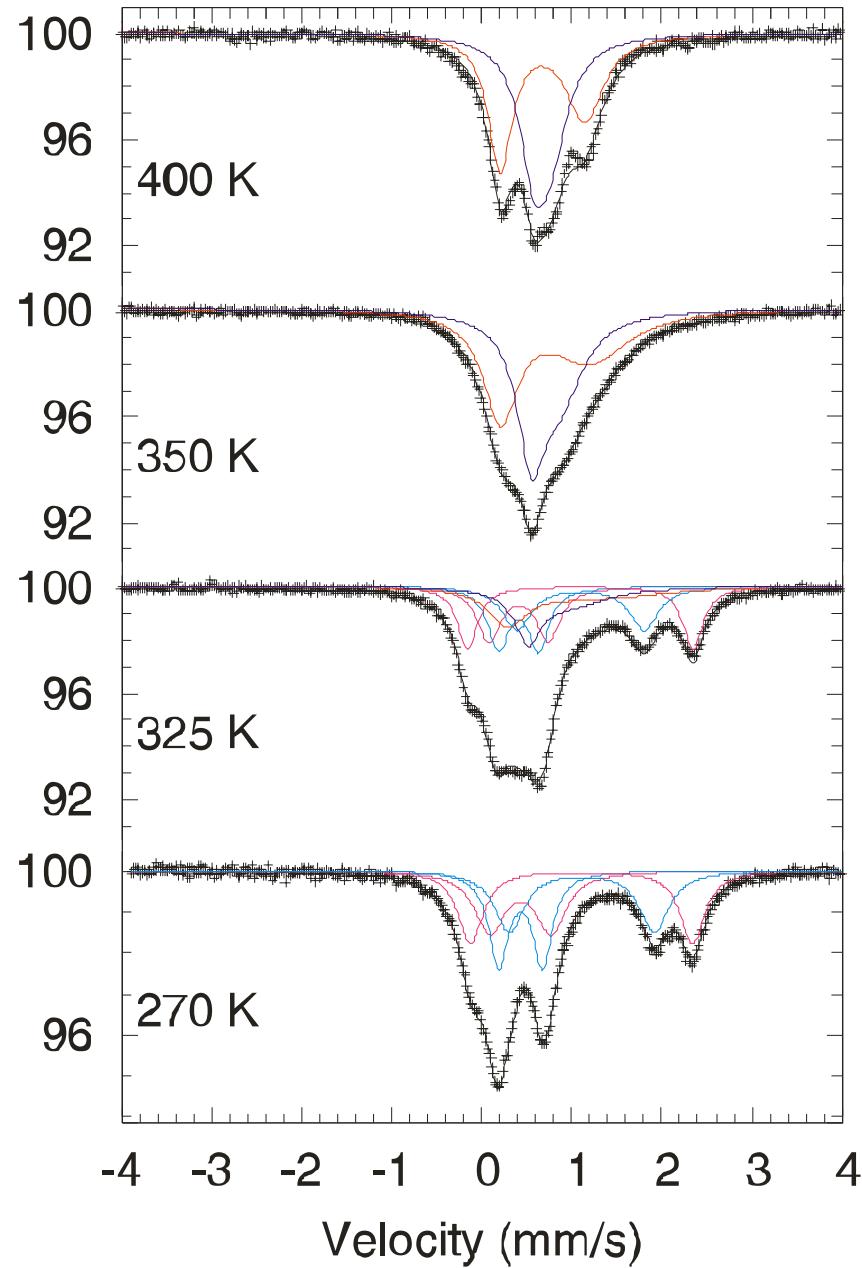
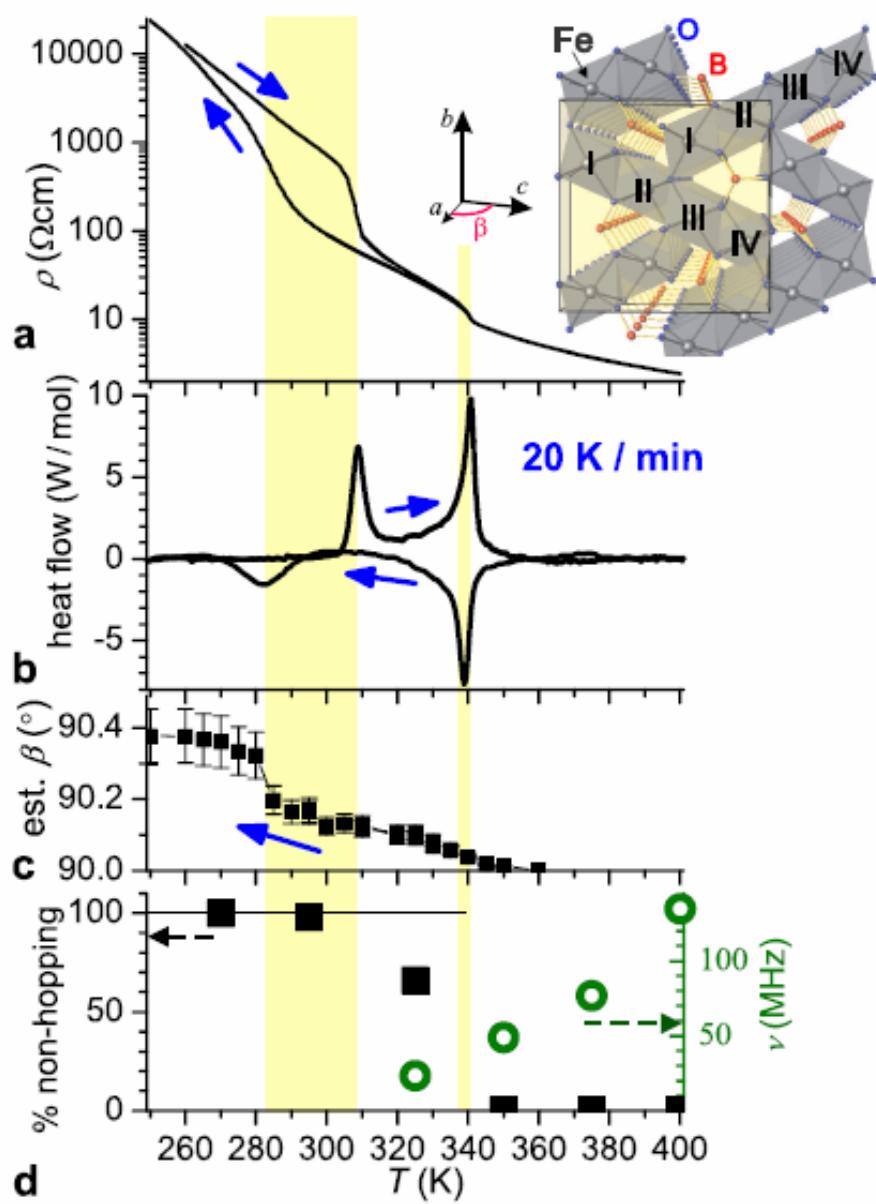
Fe_2OBO_3 : intermediate phase

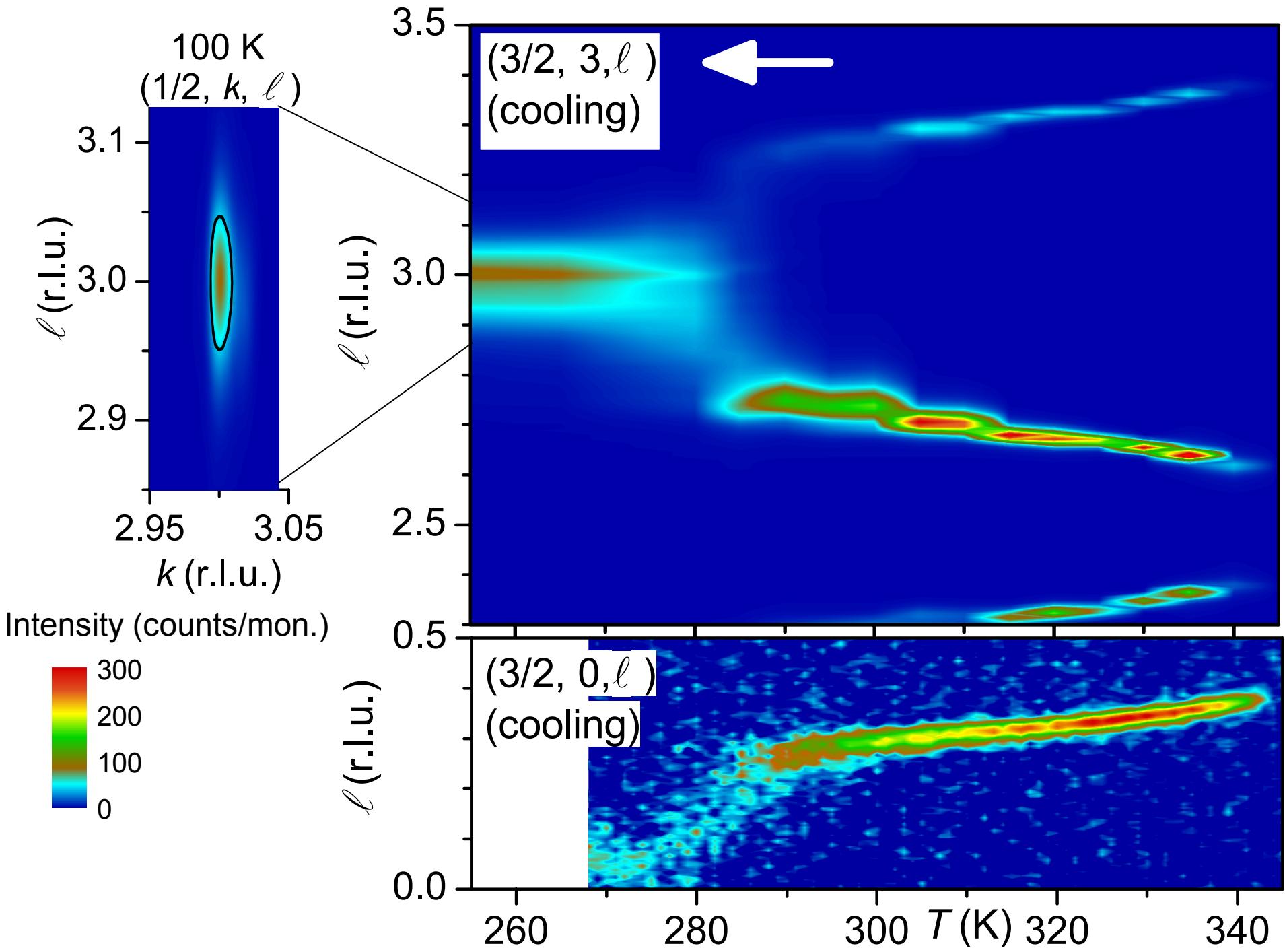


(Viji Varadarajan, Joe Brill, UK)
Hysteresis \rightarrow **II** (weakly) first order
too

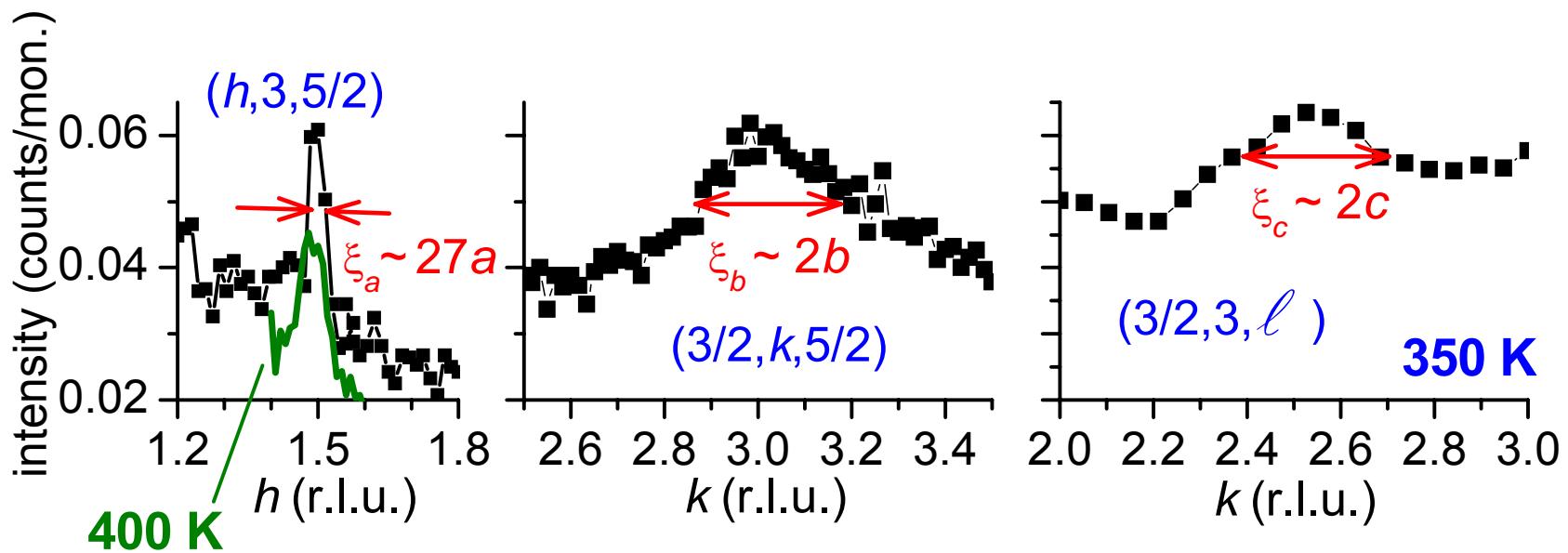
Structurally, **II**
is monoclinic-orthorhombic
transition

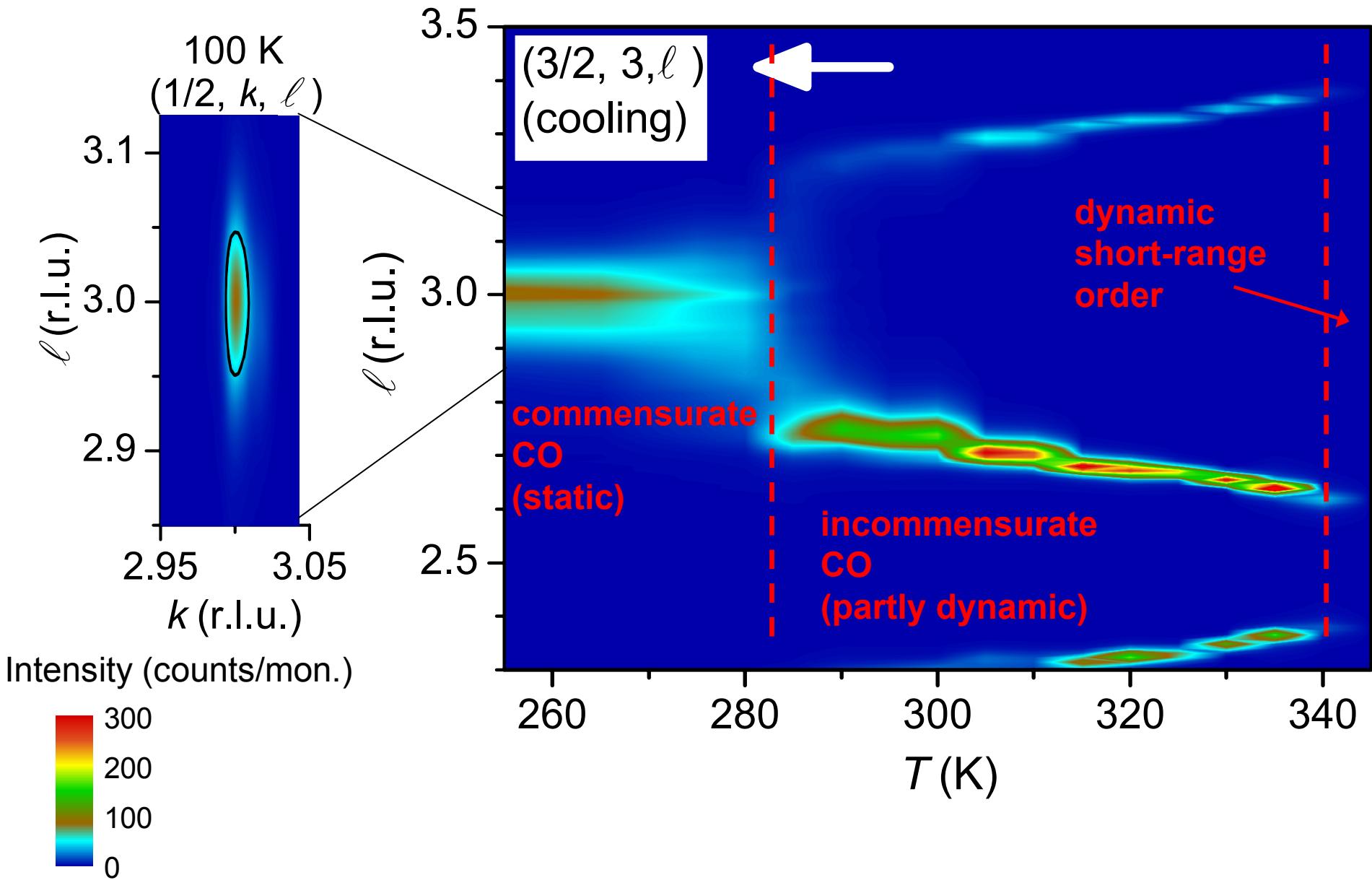
Fe_2OBO_3 : intermediate phase



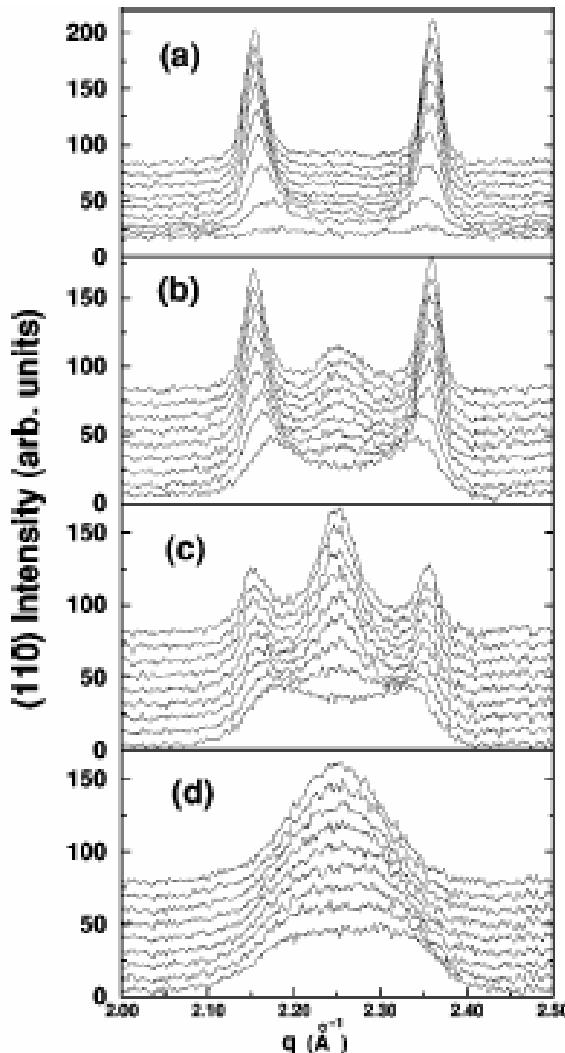


Correlations Above CO Transition





Incommensurate Phase Belongs to Class of Phenomena Known as “Modulated Phases”



Cu-Au 50-50 at. %

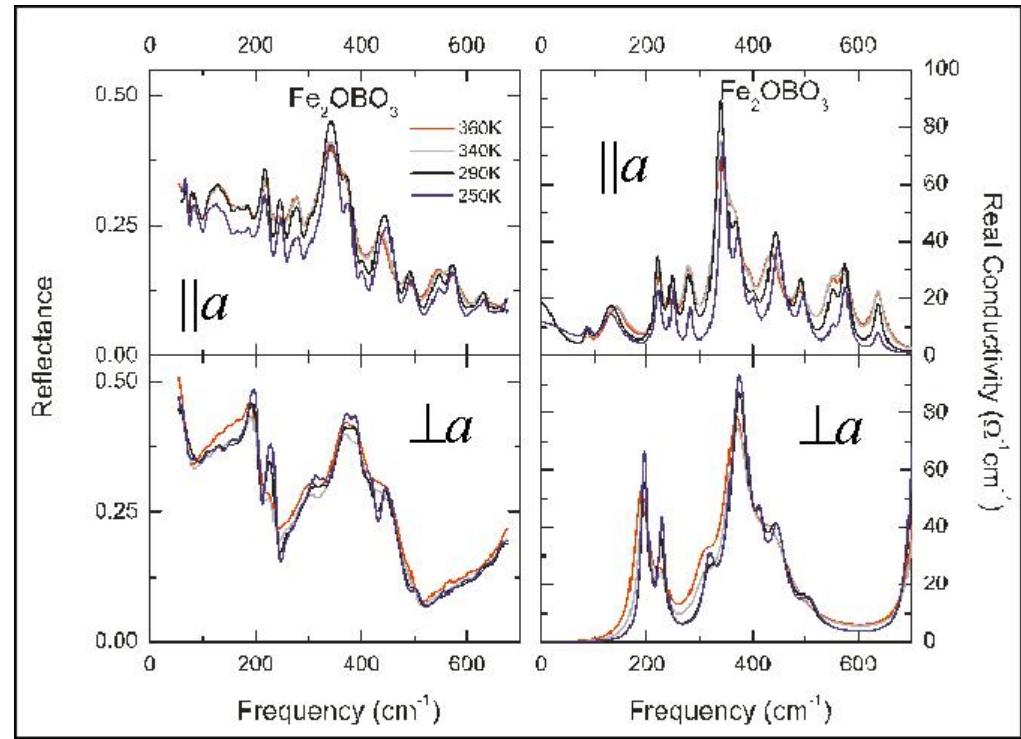
- Above 410 °C stable phase disordered
- Below 385 °C equilibrium phase CuAu I, in which Cu and Au occupy alternating atomic Layers
- Between these 2 transition temperatures CuAu orders into a one dimensional long Period superlattice called CuAu II, which consists of a periodic array of antiphase boundaries with an average modulation wavelength approx. 10 times the size of the underlying CuAu cell

Malis, Ludwig, PRB **60**, 14675 (1999)

“the origin of the modulated phase is still under debate”

Future Work

Optics – Willie Padilla,
Boston College



“Switching” Effects in IV curves -- Ana Akrap, L. Forro, EPFL

Impedance Spectroscopy, search for ferroelectricity

TEM work

Summary

LuFe_2O_4

- Oxygen stoichiometry strongly affects physical properties
- Observation of thermal anomaly at 330 K
- 3D magnetic structure solved
- New phase transition observed at 180 K

Fe_2OBO_3

- Crystal growth of Fe_2OBO_3
- Found C.O. superstructure reflections
- Showed that Fe_2OBO_3 is best example to date of ionic C.O.
- Found the incommensurate phase, assigned it to the same class of phenomena as the CuAu II “modulated phase”