

Modeling the Solid State with Coupled Atomistic/Continuum Methods

Ronald E. Miller Carleton University, Ottawa, Canada

KITP June 7, 2012



Credits

Lead Actors

Collaborator on many things (QC, hot QC, comparisons): Ellad Tadmor, U. Minnesota

Post-Docs: Denis Saraev (PDF, currently at large) Behrouz Shiari (PDF, now at NNIN Michigan)

> Collaborators on CADD: Leo Shilkrot (currently at large) Bill Curtin (Brown and EPFL)

Collaborators on hot QC: Laurent Dupuy, CEA Frederic Legoll, University of Paris W.K. Kim, PDF at U. Minnesota

Supporting Cast

Help with all the comparisons of methods: Mitch Luskin, Univ. of Minnesota Michael Parks, LANL Catalin Picu, RPI Dong Qian, U. Cincinnati

Producers



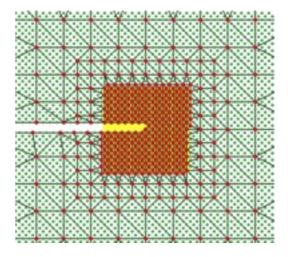


Ontario Premier's Research Excellence Awards (PREA)



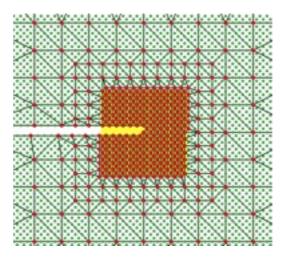


- I. Static partitioned-domain methods for modeling crystalline solids
 - Overview of the essential ingredients of these coupling methods





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- 2. Relative accuracy and efficiency of 14 static methods from the literature

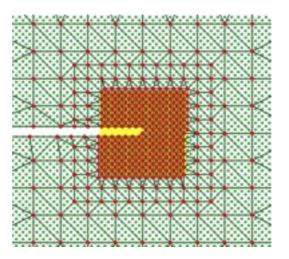


QC-G/FC-10	8D-10	OCICLS-10	ESM-10	COC(13)-E-10	COC(1)-GFC-10	CRCM-10
*	*	*	- 20		-	
90-5/70-20	8D-20	QCICLS-20	854-20	COC(13)-E-20	COC(1)-GPC-20	CACM-20
#	*	#	*	٠	÷	
00-6/10-30	80-30	QCICLS-30	85M-30	000(13)-8-90	COC(1)-GPC-30	CACM-30
15	*	-11	$\frac{4}{r}$			* *

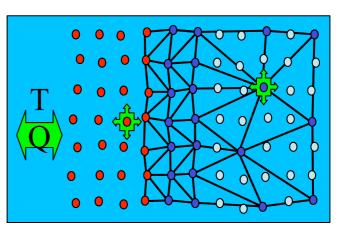


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3. Extending these methods to finite Temperature



00-GFC-10	8D-10	OCICLS-10	85M-10	COC(13)-E-10	COC(1)-OFC-10	CACM-10
*	*	*	•		.	
90-6FC-20	8D-20	QCICLS-20	85M-20	CQC(13)-E-20	COC(1)-GFC-20	CACM-20
#	*	#	*		÷	
00-0/10-30	80-30	QCICLS-90	85M-30	000(13)-8-30	COC(1)-GFC-30	CACM-30
#	*	4	*			2 :2



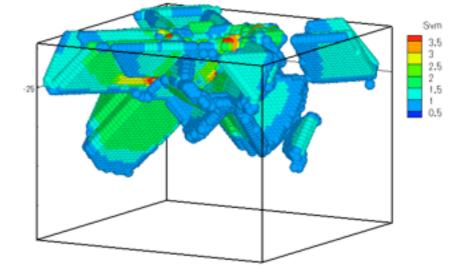
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about 200,000 Ni on Cu at zero K

MD simulations require relatively large systems

Nickel on Copper 0 K 3.5 3 2.5 2 1.5 1.5 1 0.5



Saraev and Miller, Acta Mat., 2006

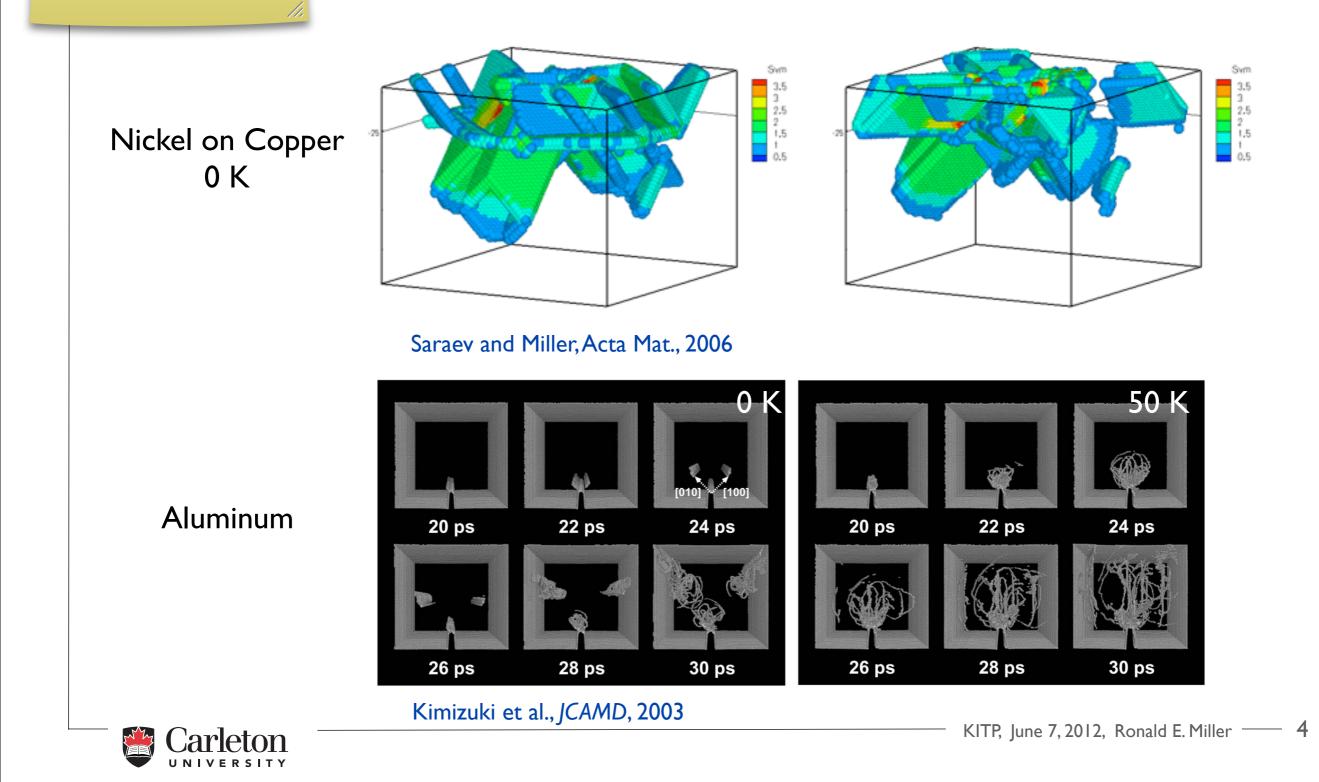


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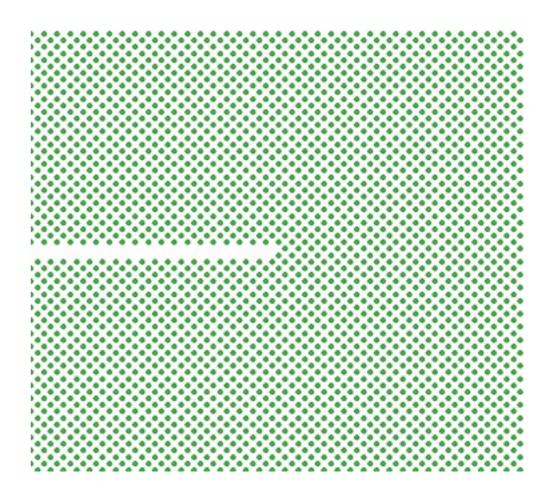
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$$\mathcal{V} = \sum_{\alpha}^{\text{atoms}} E^{\alpha}$$



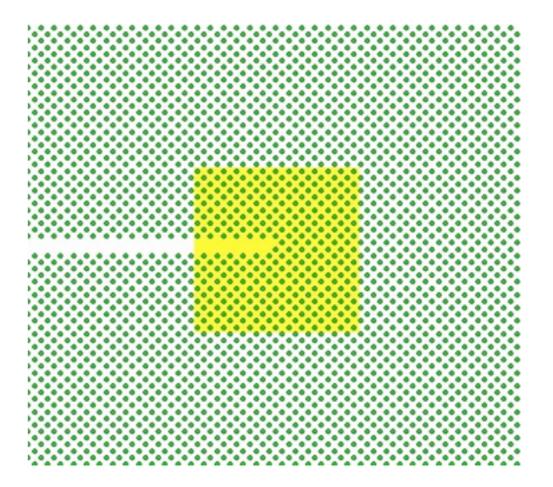


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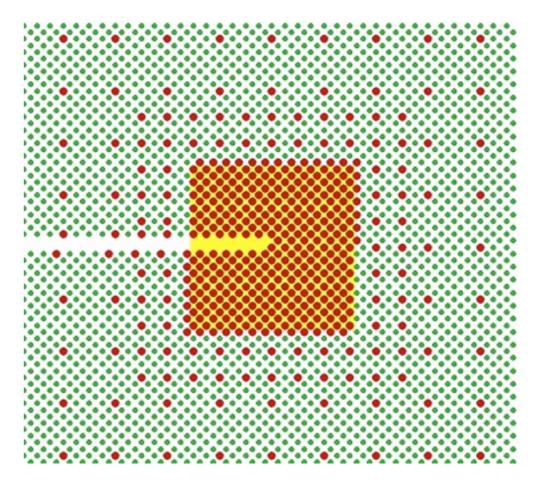
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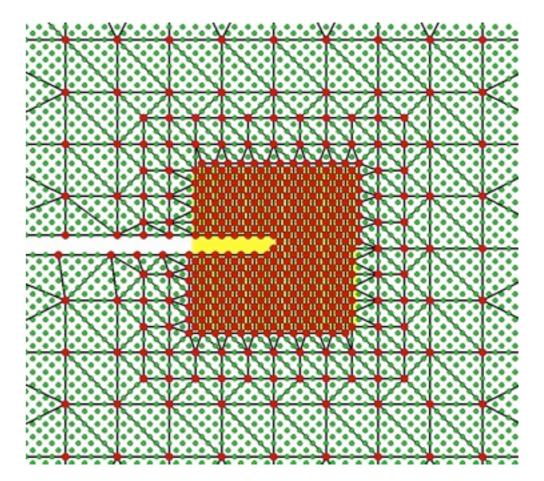
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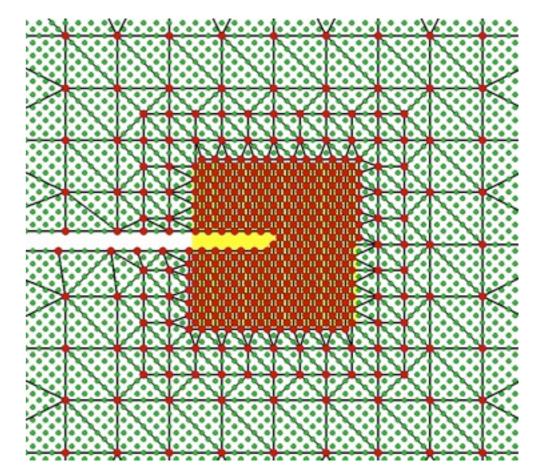
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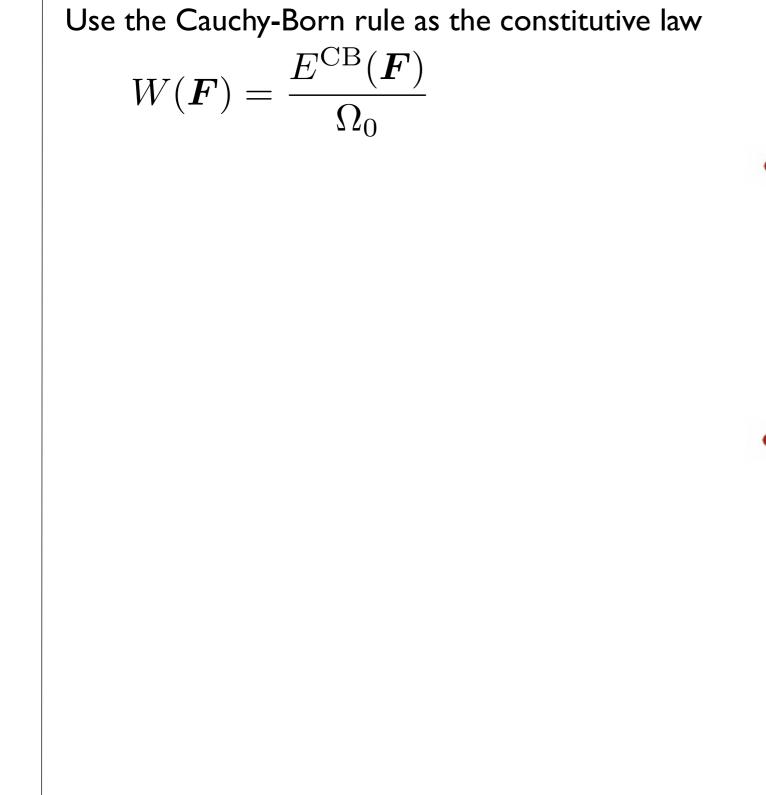
Positions of atoms found by interpolation

$$oldsymbol{u}^lpha = \sum_i^{ ext{nodes}} S_i(oldsymbol{X}^lpha)oldsymbol{u}_i$$

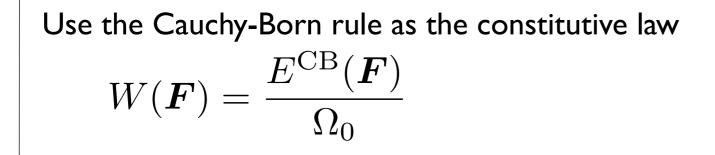




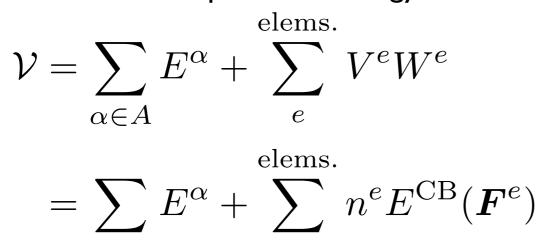
x = FX

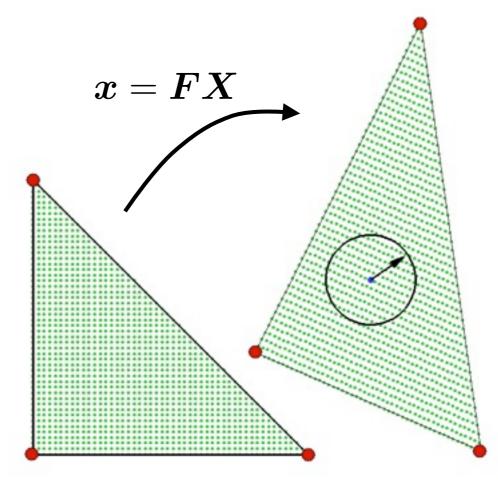






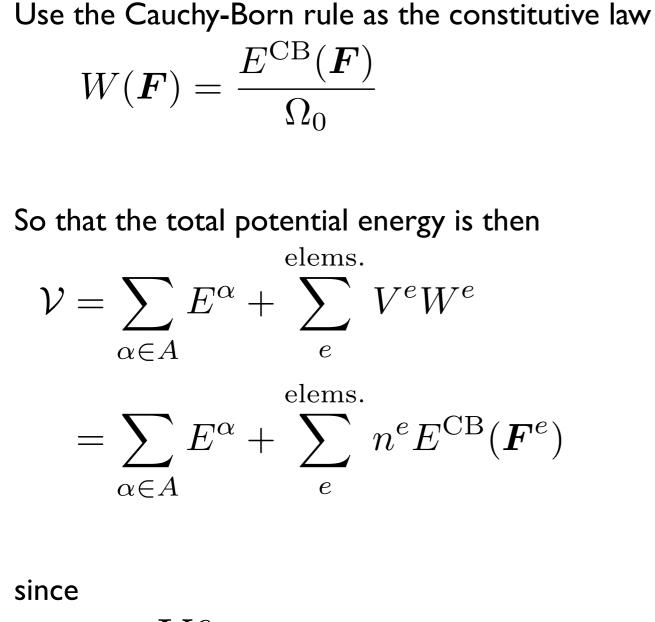
So that the total potential energy is then

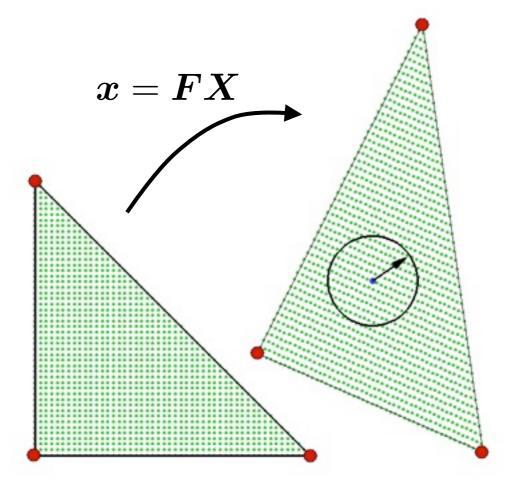






 $\alpha \in A$



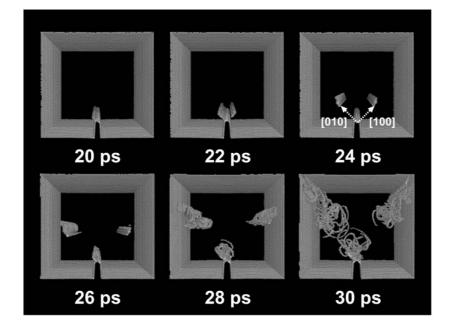


$$\frac{V^e}{\Omega_0} \approx n^e$$



Example: Fracture using the QC Method

Multiscale Goal: To replace this...



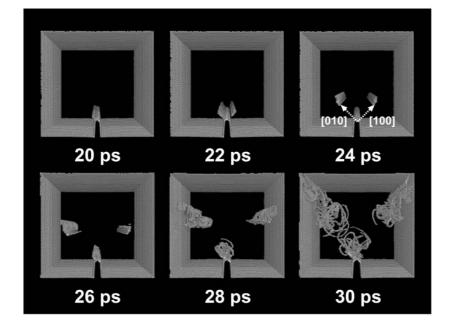
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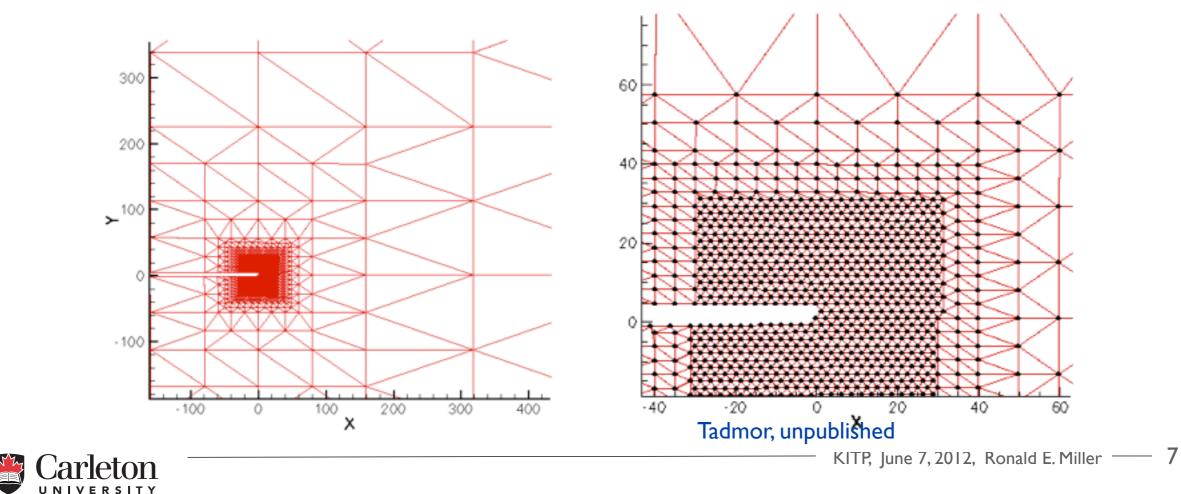
Tadmor, unpublished

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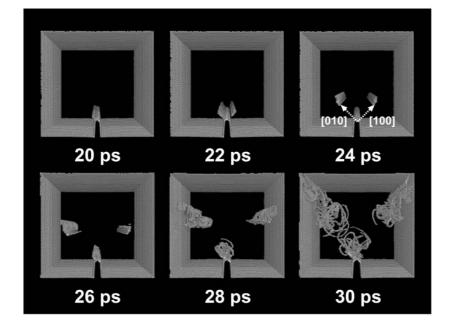


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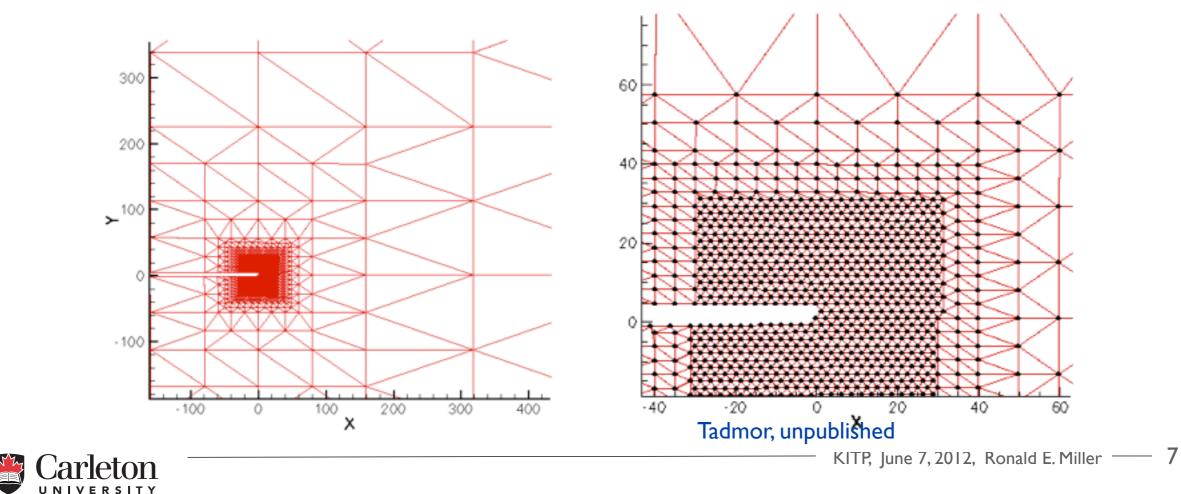


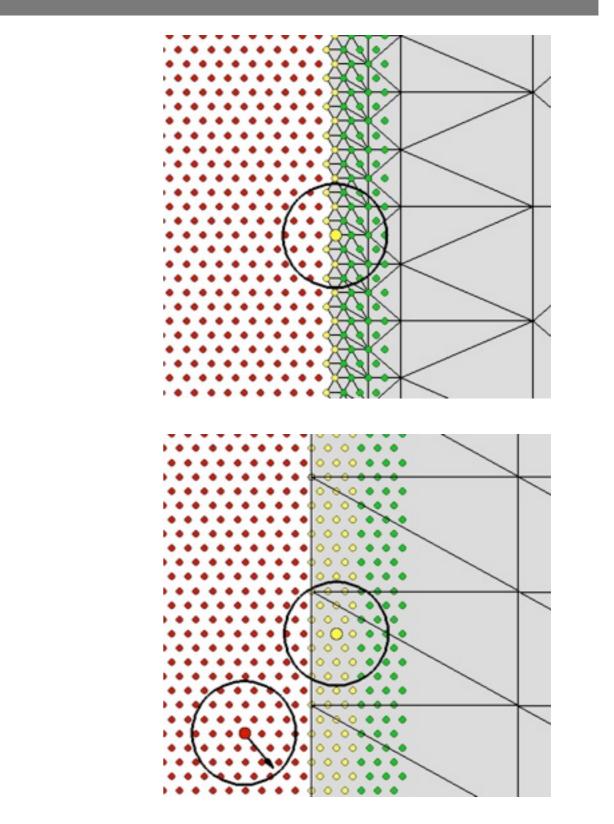
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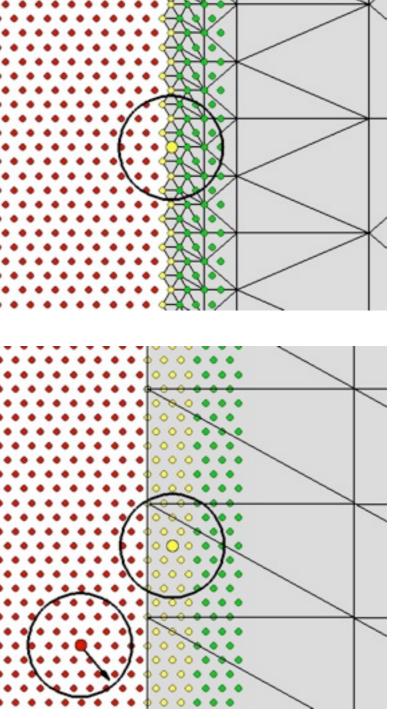
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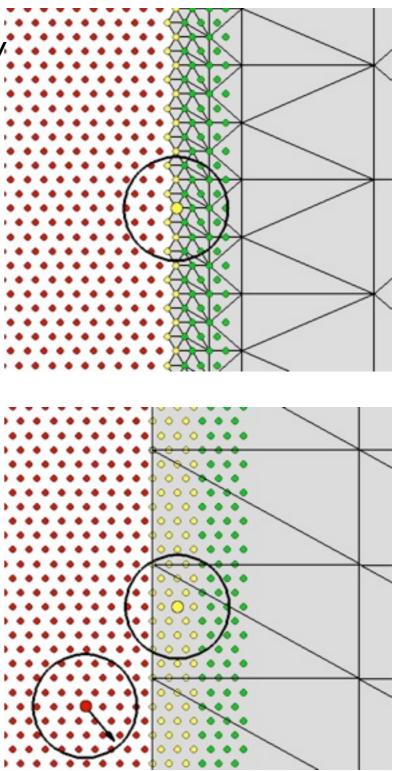
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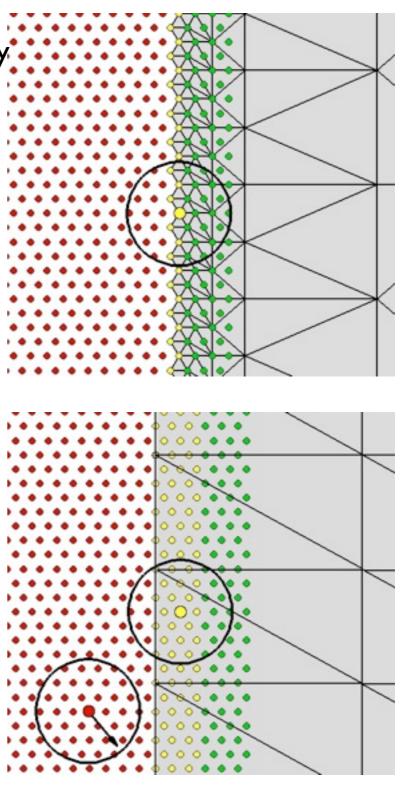
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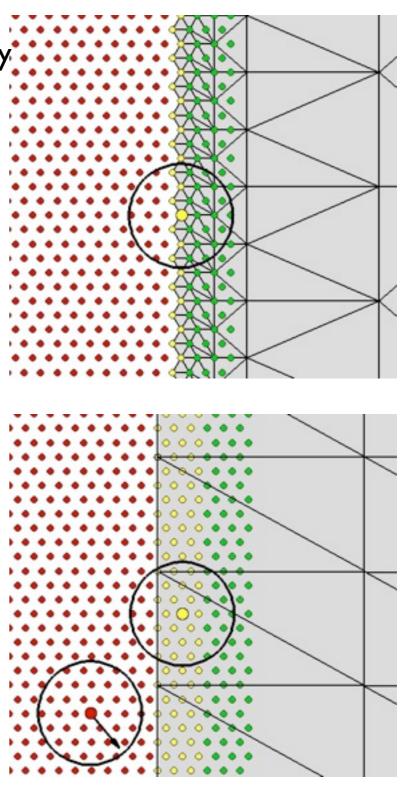
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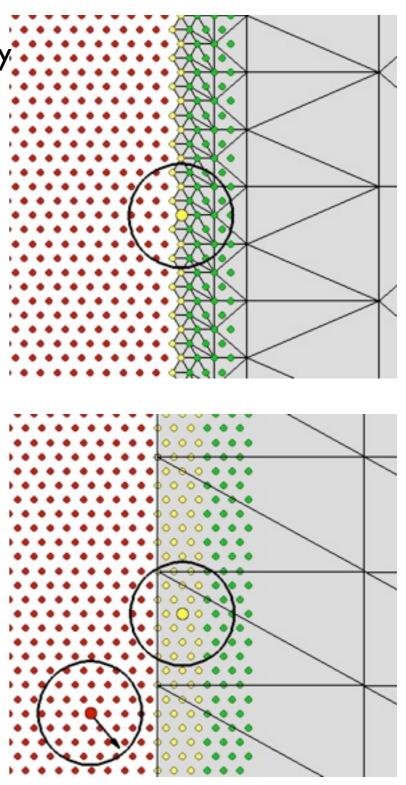
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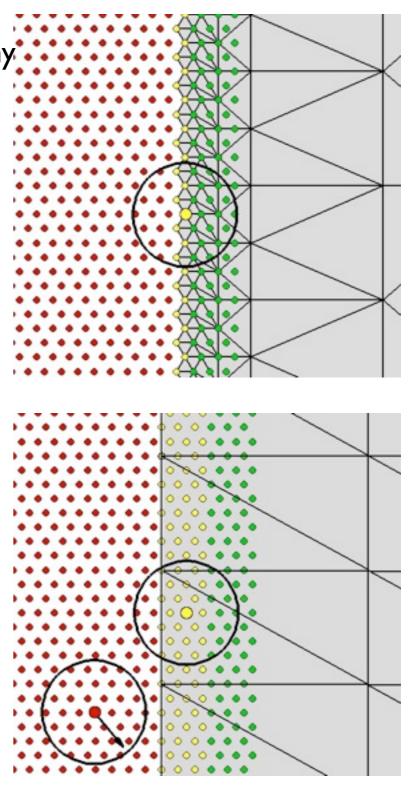
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- 4. The choice of the finite element constitutive law





Comparison of 14 coupled methods (Miller and Tadmor, MSMSE, 2009)

8 Energy-based methods

Key References	Continuum Handshake Model		Coupling Boundary Condition	
[1, 2]	Cauchy-Born	None	Strong Compatibility	
[3]	Linear Elasticity	None	Strong Compatibility	
[4]	Cauchy-Born	Linear mixing of energy	Weak Compatibility (penalty)	
[5, 6]	Cauchy-Born	None	Weak/Stong Mix (least-squares fit)	
[7]	Linear Elasticity	None	Weak Compatibility (average atomic positions)	
[8]	Averaging of atomic clusters	None	Strong Compatibility	
[9]	Cauchy-Born	None	Strong Compatibility	
[8]	Averaging of atomic clusters	None	Strong Compatibility	
	References [1, 2] [3] [4] [5, 6] [7] [8] [9]	ReferencesModel[1, 2]Cauchy-Born[3]Linear Elasticity[4]Cauchy-Born[5, 6]Cauchy-Born[7]Linear Elasticity[8]Averaging of atomic clusters[9]Cauchy-Born[8]Averaging of[8]Averaging of	ReferencesModel[1, 2]Cauchy-BornNone[3]Linear ElasticityNone[4]Cauchy-BornLinear mixing of energy[5, 6]Cauchy-BornNone[7]Linear ElasticityNone[8]Averaging of atomic clustersNone[9]Cauchy-BornNone[8]Averaging of NoneNone[8]Averaging of NoneNone[9]Cauchy-BornNone	



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Acronym	Key References	Continuum Model	Handshake	Coupling Boundary Condition
FEAt	[10]	non-linear, nonlocal elasticity	None	Strong Compatibility
CADD	[11, 12]	Linear Elasticity	None	Strong Compatibility
HSM	[13]	Non-Linear Elasticity	atomic averaging for nodal B.C.	Weak Compatibility (average atomic positions)
AtC	[14, 15, 16, 17]	Linear Elasticity	Linear mixing of stress and atomic force	Strong Compatibility
AtC-GFC	unpublished	Linear Elasticity	Linear mixing of stress and atomic force	Strong Compatibility
CQC(m)-F	[18]	Averaging of atomic clusters	None	Strong Compatibility

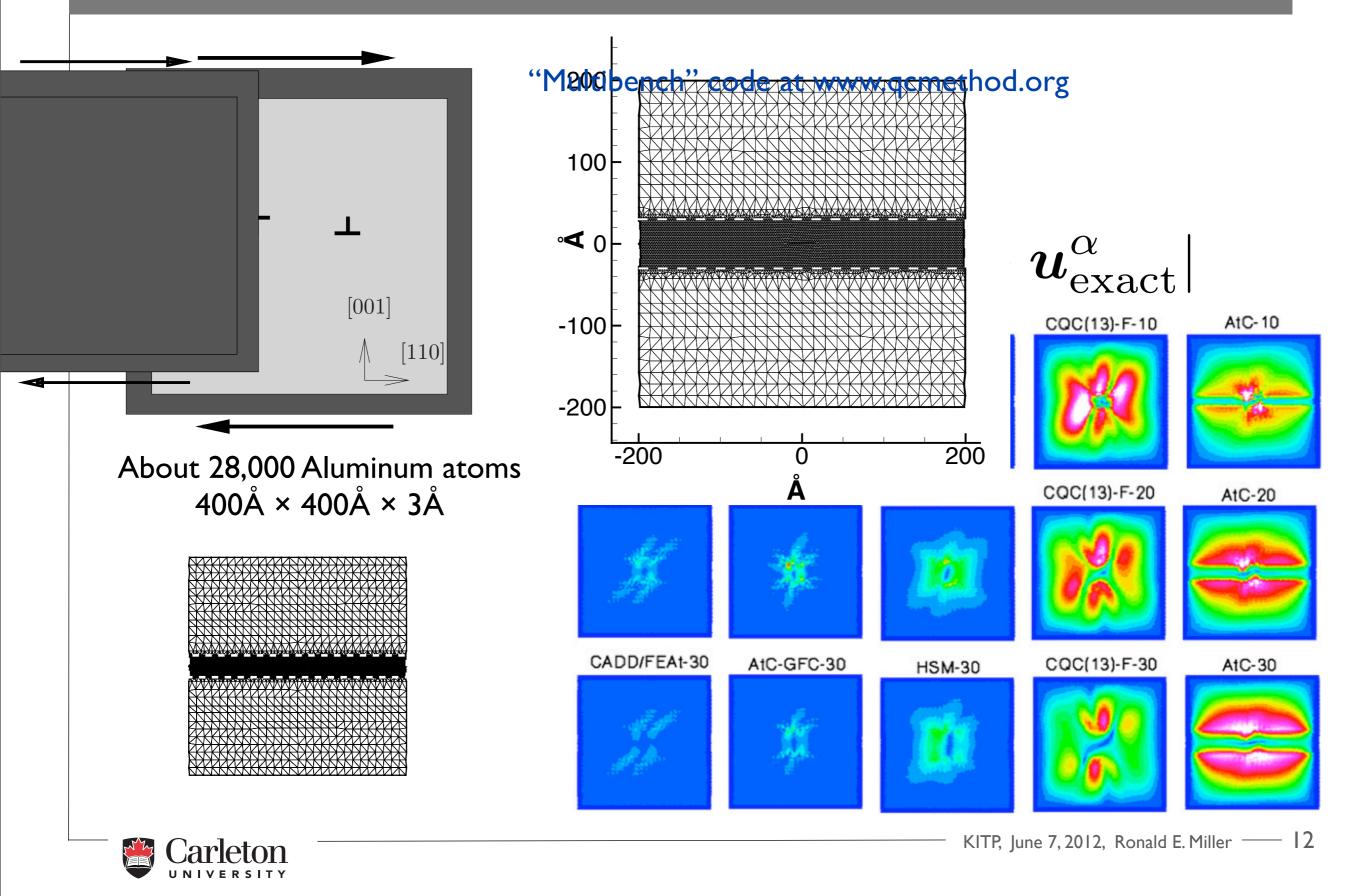


- Tadmor, E. B., Ortiz, M., and Phillips, R. *Philos. Mag. A* 73(6), 1529–1563 (1996).
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- [16] Badia, S., Parks, M., Bochev, P., Gunzburger, M., and Lehoucq, R. Multiscale Model. Simul. 7(1), 381–406 (2008).
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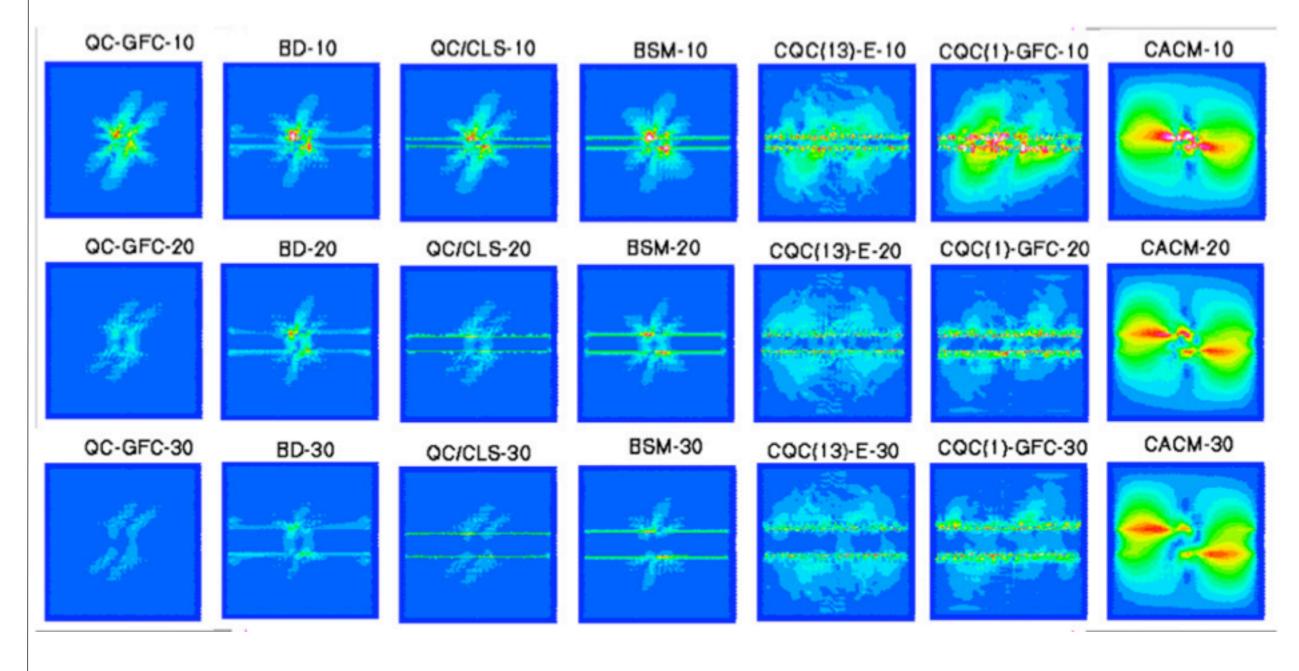


Comparison of Accuracy: Force-based methods

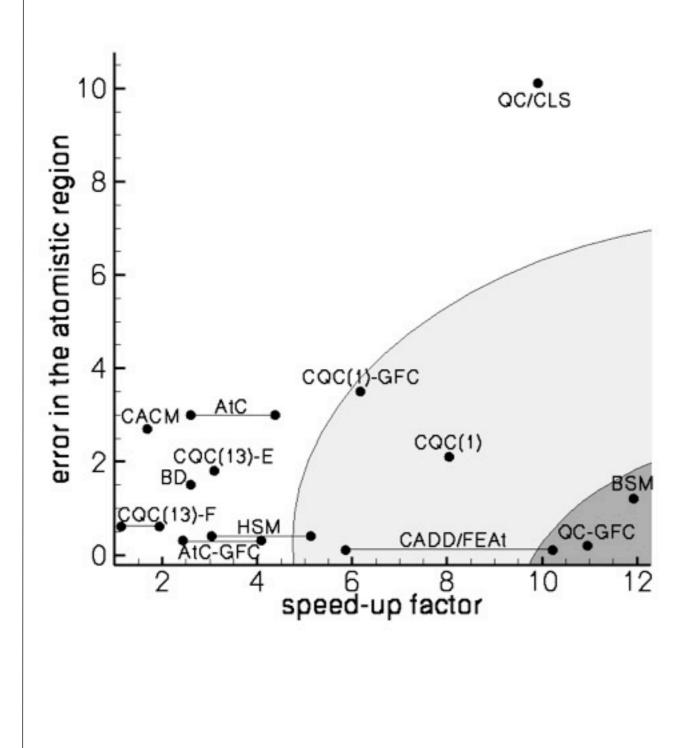


Comparison of Accuracy: Energy-based methods

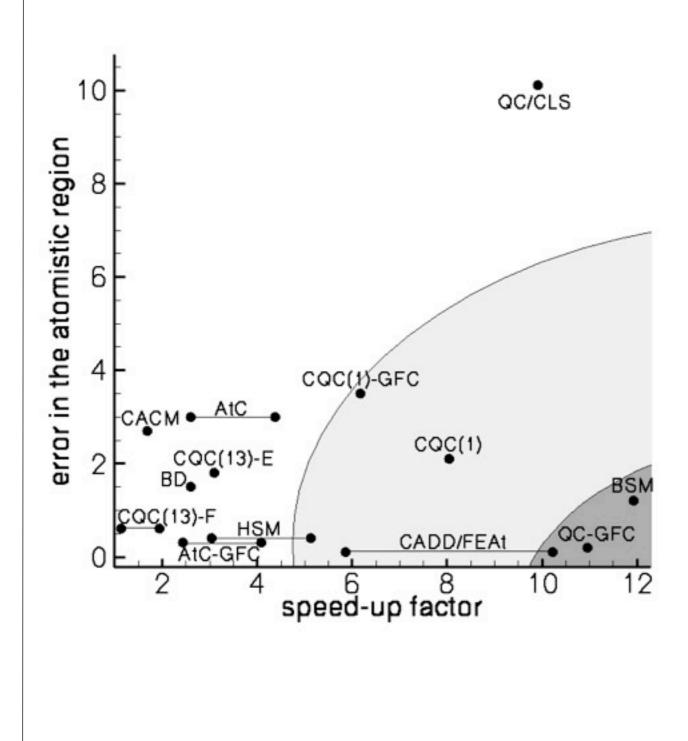
$$\epsilon^{lpha} = | \boldsymbol{u}^{lpha} - \boldsymbol{u}^{lpha}_{ ext{exact}} |$$





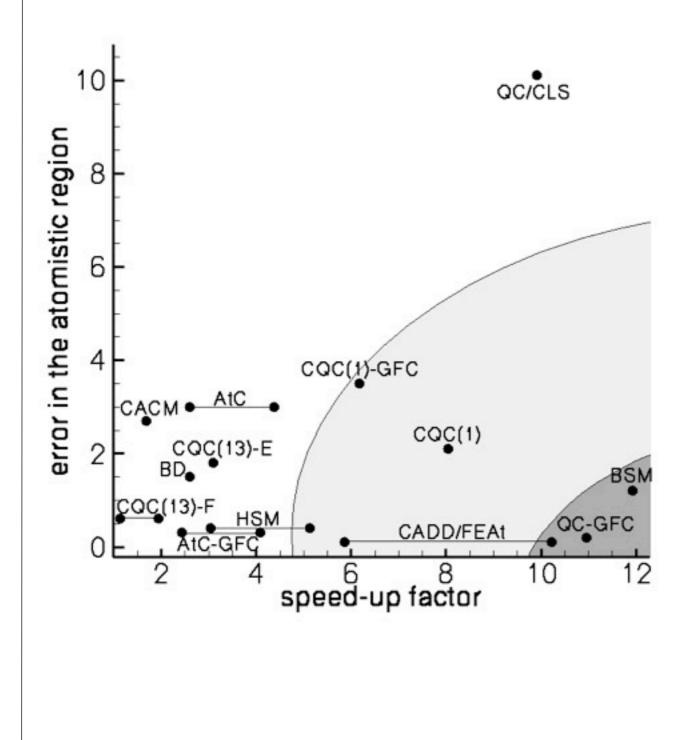






Broad conclusions drawn:

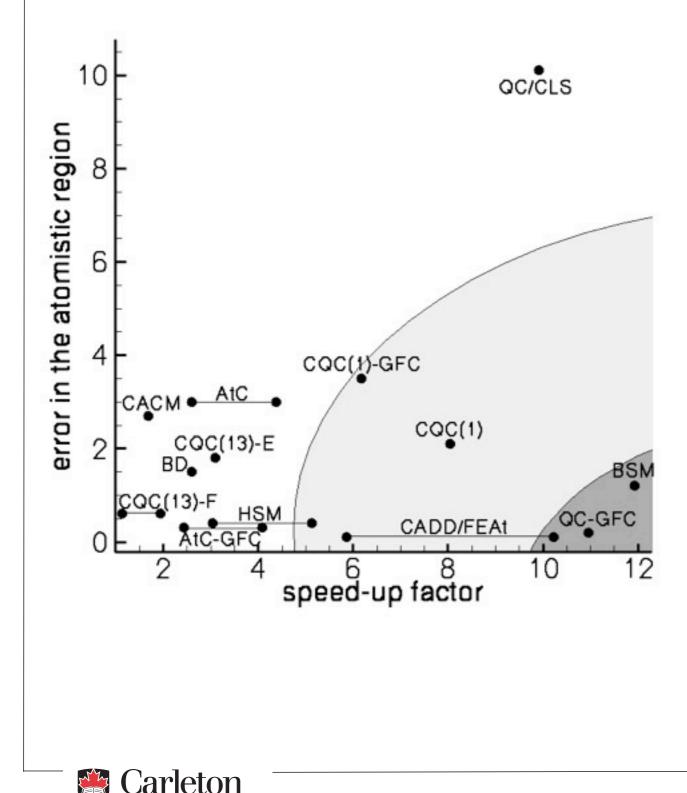




Carleton

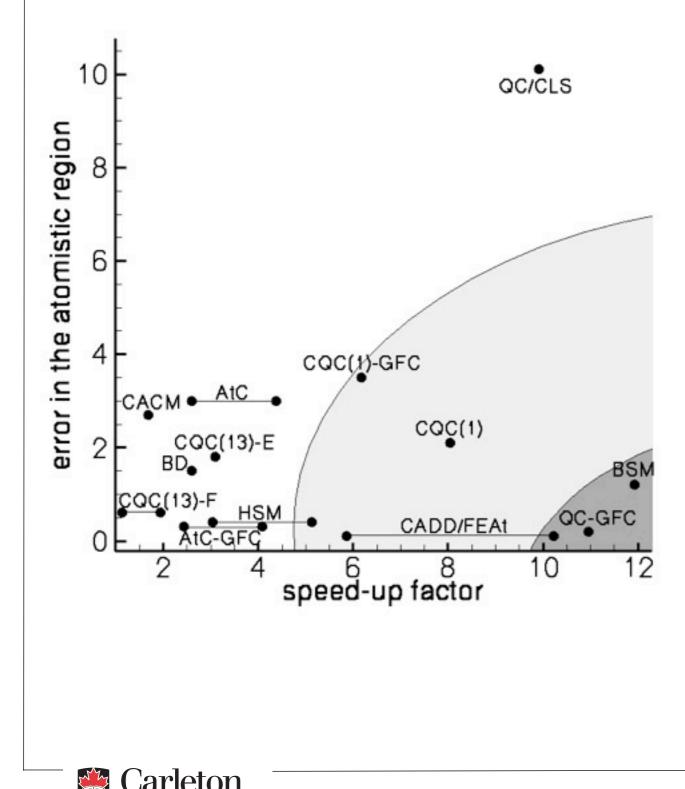
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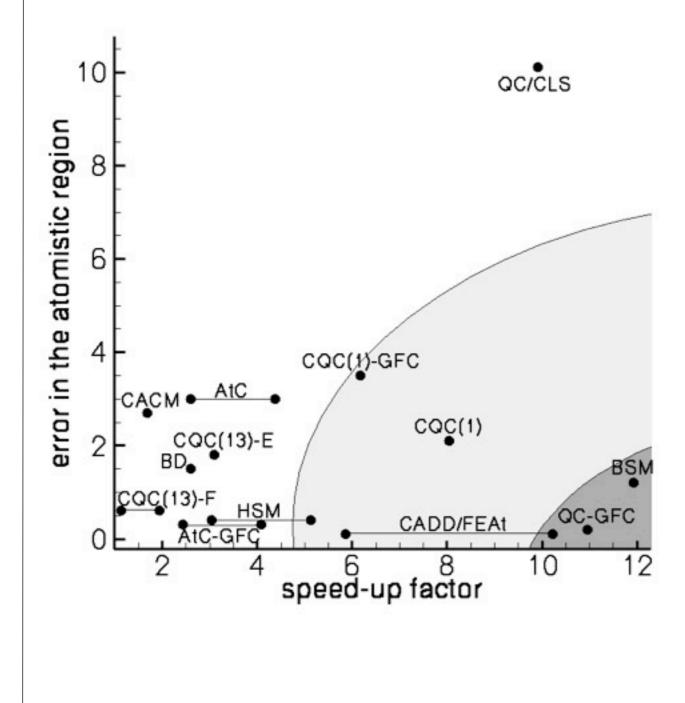
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Comparison of Speed and Accuracy of Coupling Methods

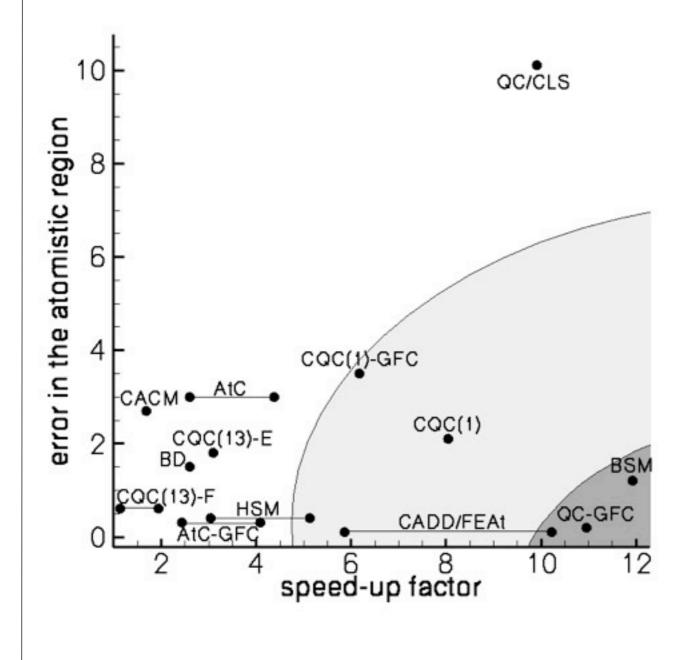


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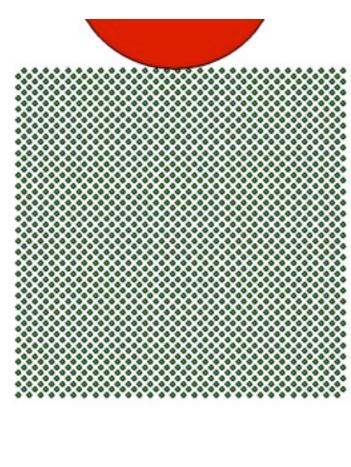
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- Ghost force correction in many methods may be a good idea.



- I. Small atomistic region: thermostats are more intrusive
- 2. How do we correctly account for the missing entropy coarsened degrees of freedom?
- 3. Wave reflections from the atomistic/continuum interface due to:
 - changes in repatom density and material properties

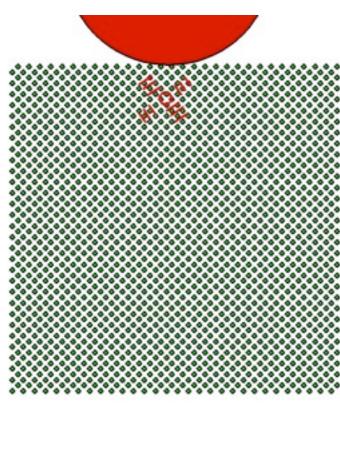


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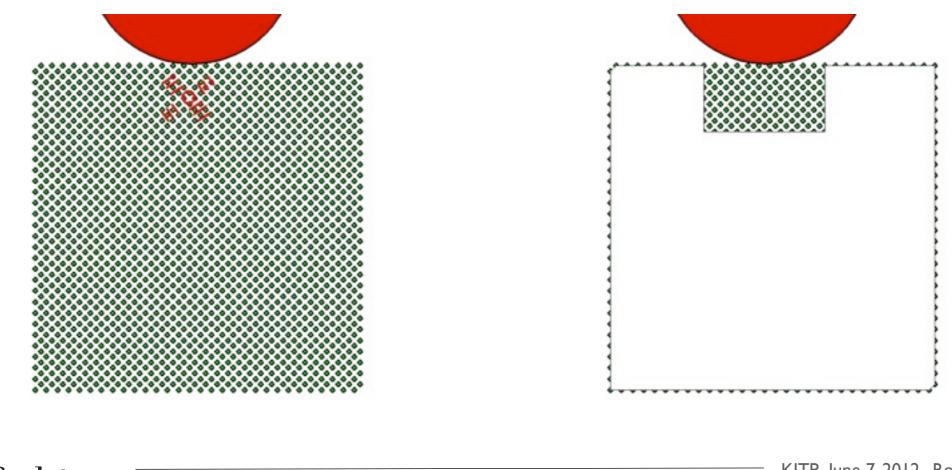




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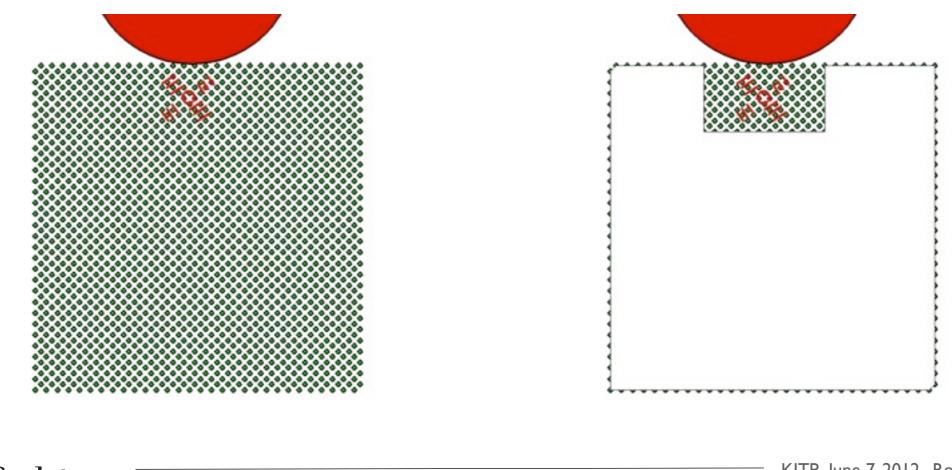


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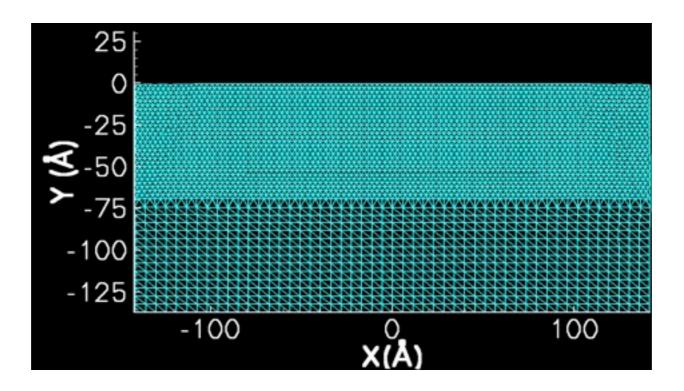


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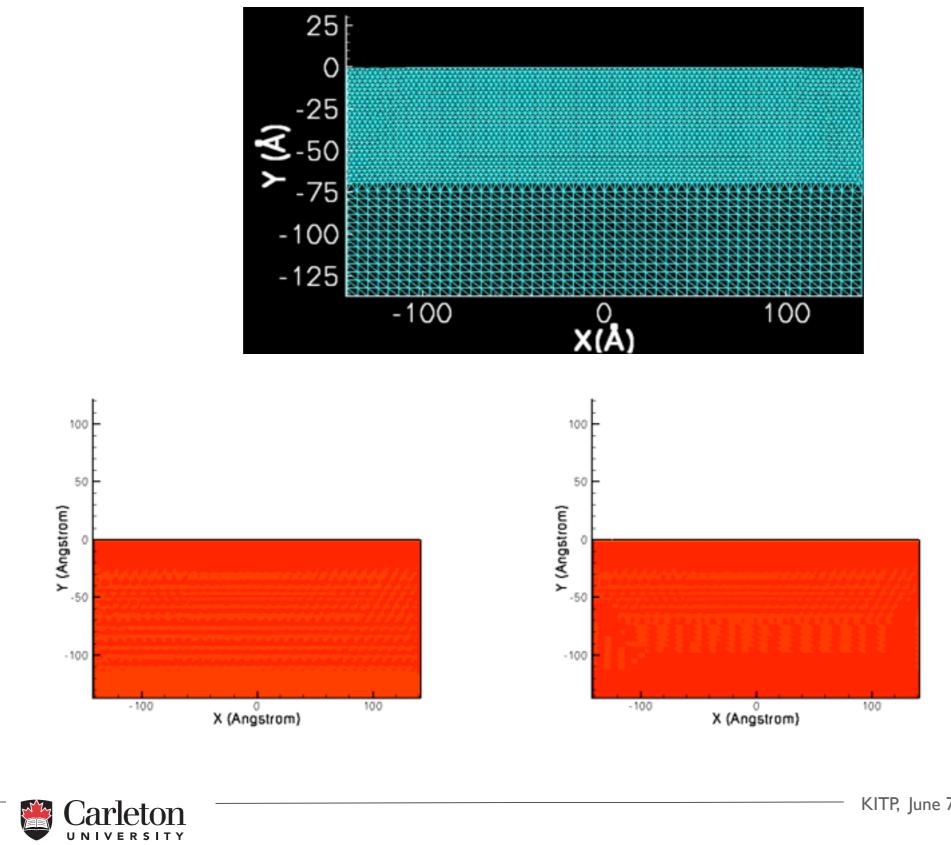


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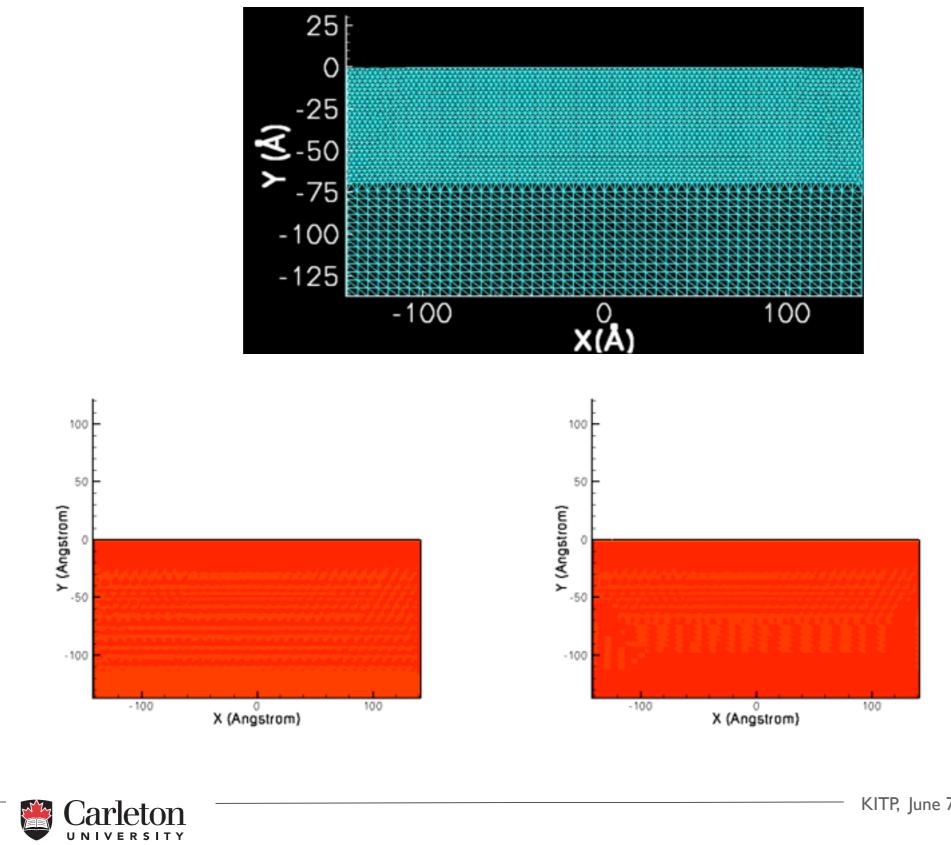




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Approaches to Finite Temperature



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I. Absorbing boundary conditions (ABC)

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- ii) the dynamics in the atomistic region is unfettered by a thermostat (NEMD)
- iii) difficult to impose temperature control
- iv) usually, absorbed heat is simply "lost" instead of transmitted
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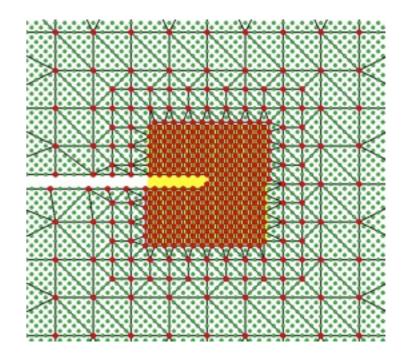
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2. Coarse-Grained Canonical Ensemble

- i) full temperature control
- ii) correct (or at least approximate) treatment of missing entropy
- iii) no direct control of wave reflections
- iv) not suitable for NEMD problems



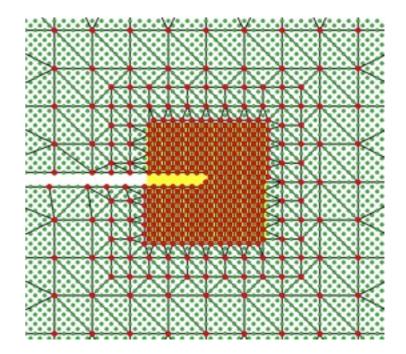
Dupuy et al, PRL, 2005, Tadmor et al, *AMR*, submitted 2012





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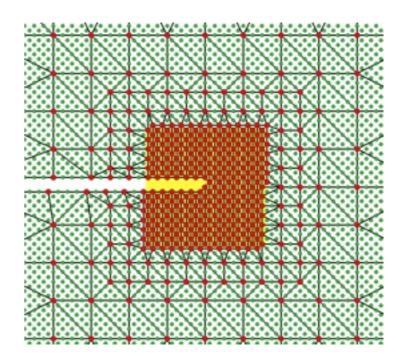




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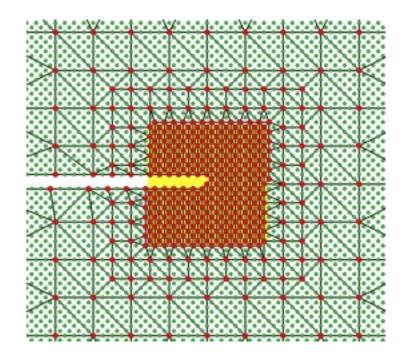


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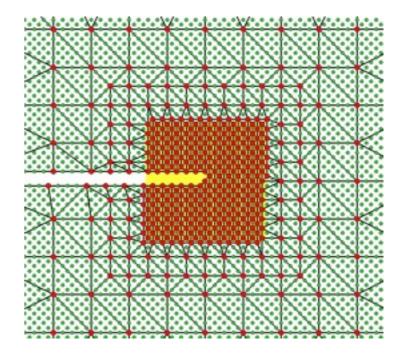
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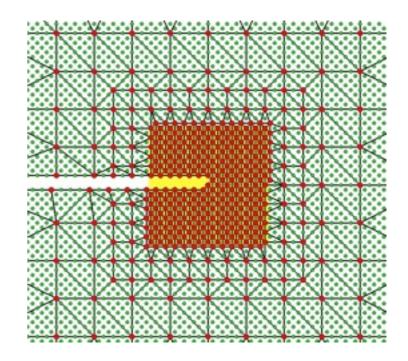
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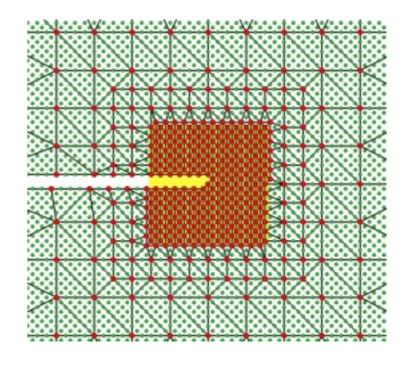
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exact but intractable



Write the potential energy as the sum of two parts:

I. The zero temperature contribution (use Cauchy-Born Rule)

2. A contribution due to the thermal fluctuations of the coarsened atoms about their mean positions (use Local Harmonic Approximation (LeSar et al., PRL, 1989) AND the Cauchy-Born Rule)

$$\mathcal{V} = \sum_{\alpha \in A} E^{\alpha} + \sum_{e}^{\text{elems.}} n^{e} \left[E^{\text{CB}}(\boldsymbol{F}^{e}) + \frac{k_{\text{B}}T}{2} \ln \frac{\left\| \boldsymbol{D}^{\text{CB}}(\boldsymbol{F}^{e}) \right\|}{(2\pi k_{\text{B}}T)^{3}} \right]$$



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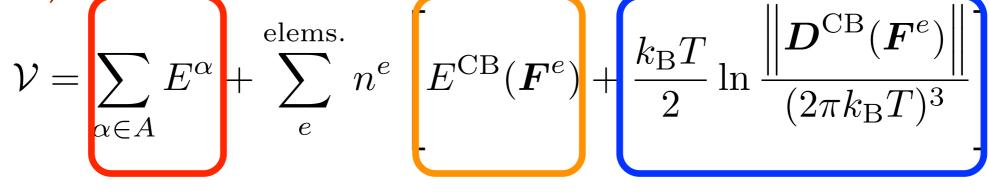
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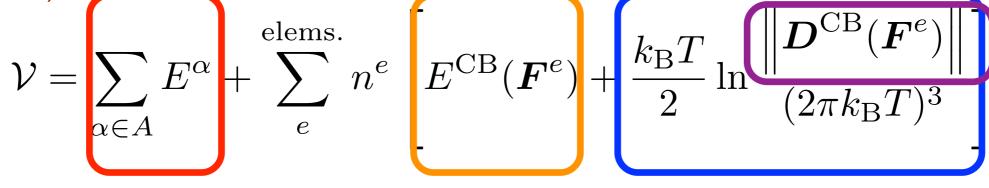
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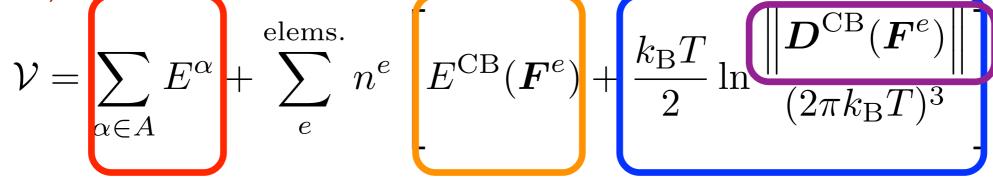
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atomistic part

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This Doesn't Work: Mesh size and shape dependent properties!

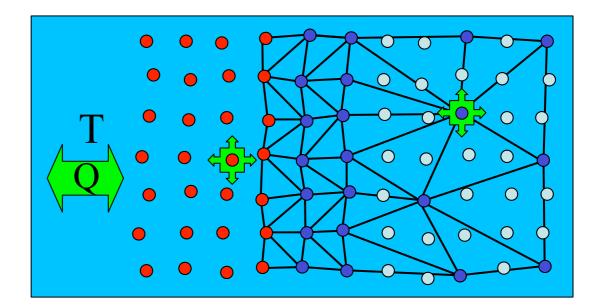


The Culprit: "Mesh entropy"

Problem:

- the motion of the nodes is really the change of the mean positions of the atoms over time
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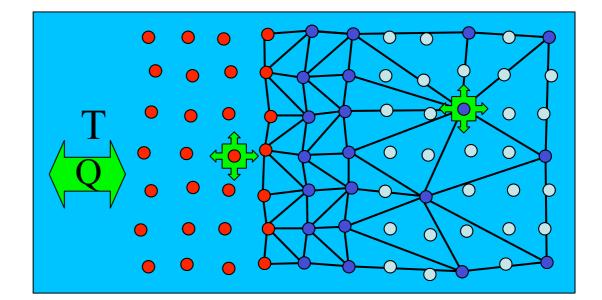
Solution:

• Subtract an approximate mesh entropy correction from the potential energy:

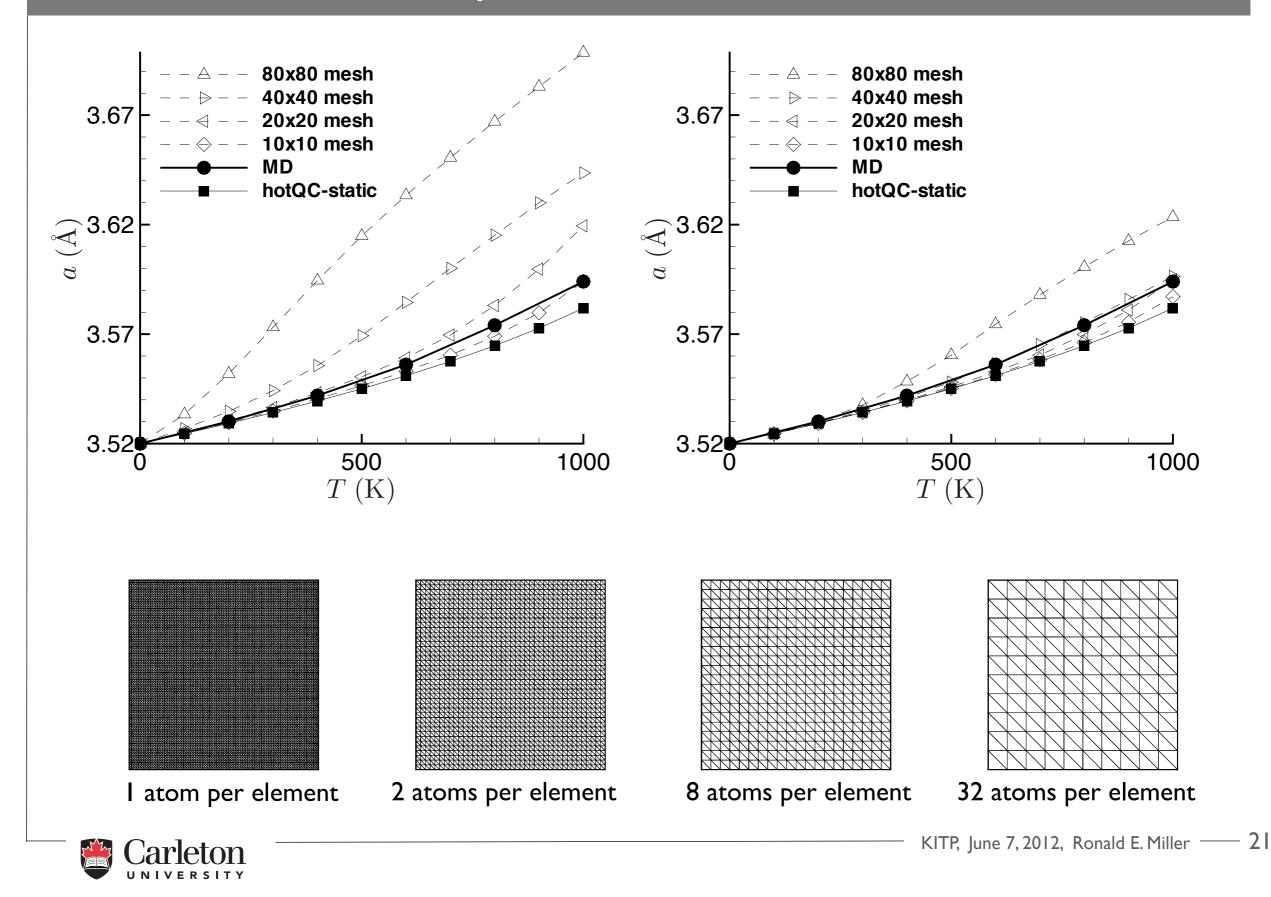
$$\widehat{\mathcal{V}} = \mathcal{V} - TS_{\text{mesh}}$$

$$S_{\text{mesh}} \approx -\frac{k_{\text{B}}}{2} \ln \frac{\det \mathbf{K}}{(2\pi k_{\text{B}}T)^{3N_{\text{loc}}}}$$

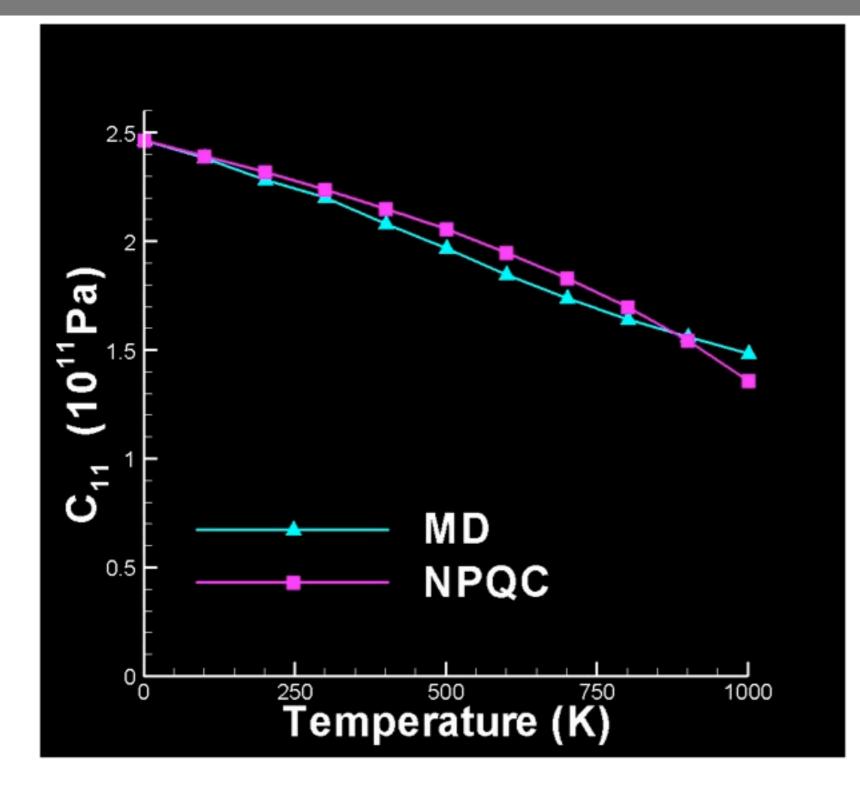




Test: Thermal Expansion of Ni (Angelo et al., MSMSE, 1995)

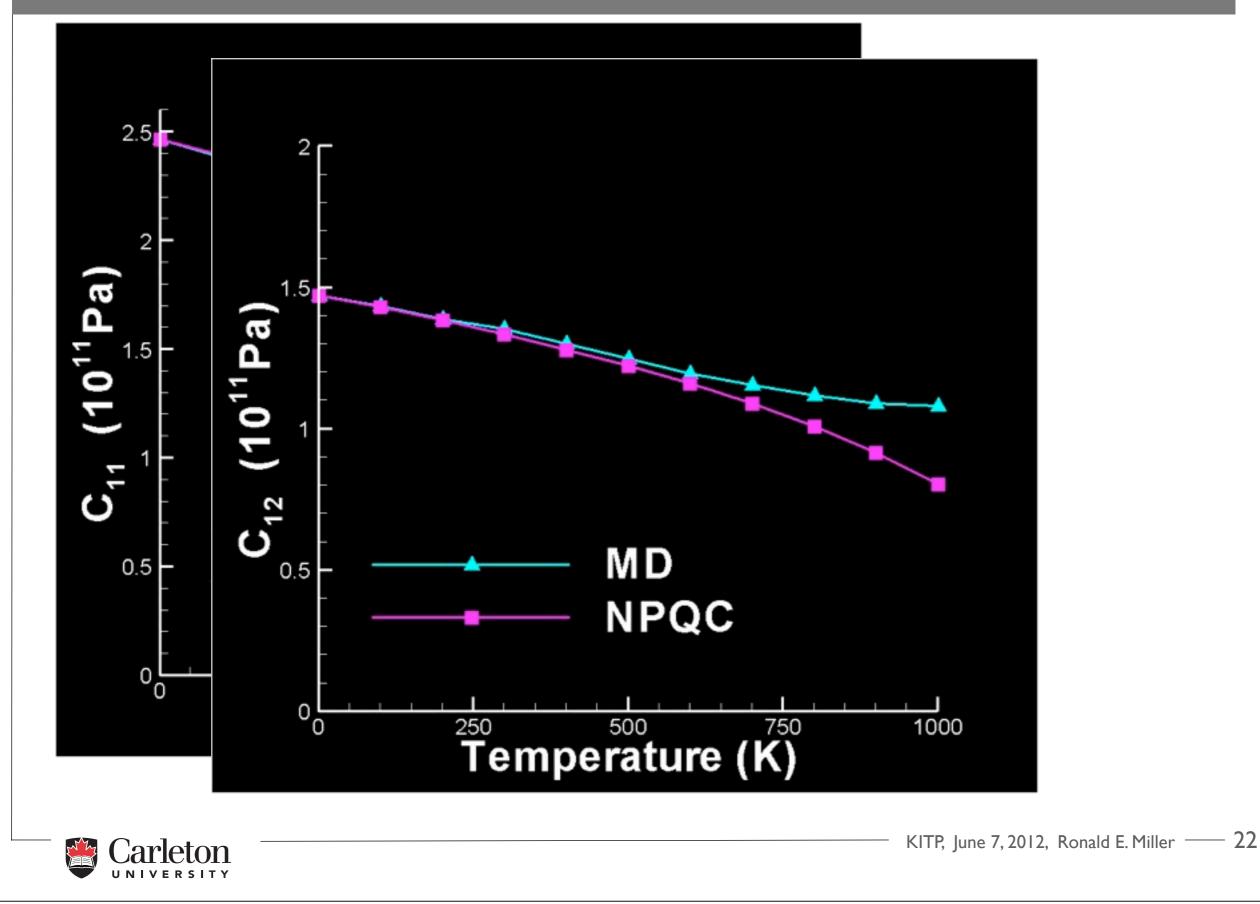


Another Test: Elastic Constants of Ni

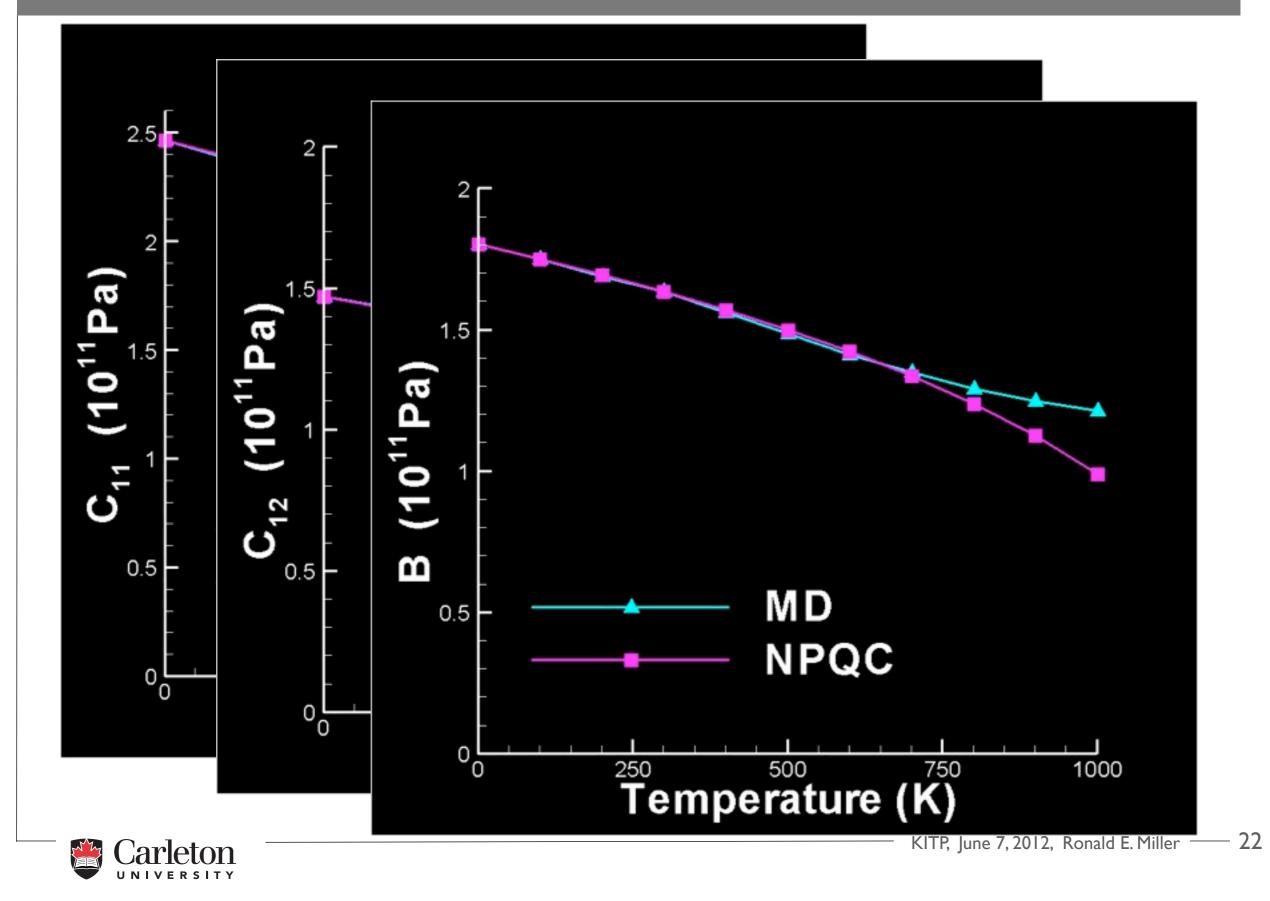


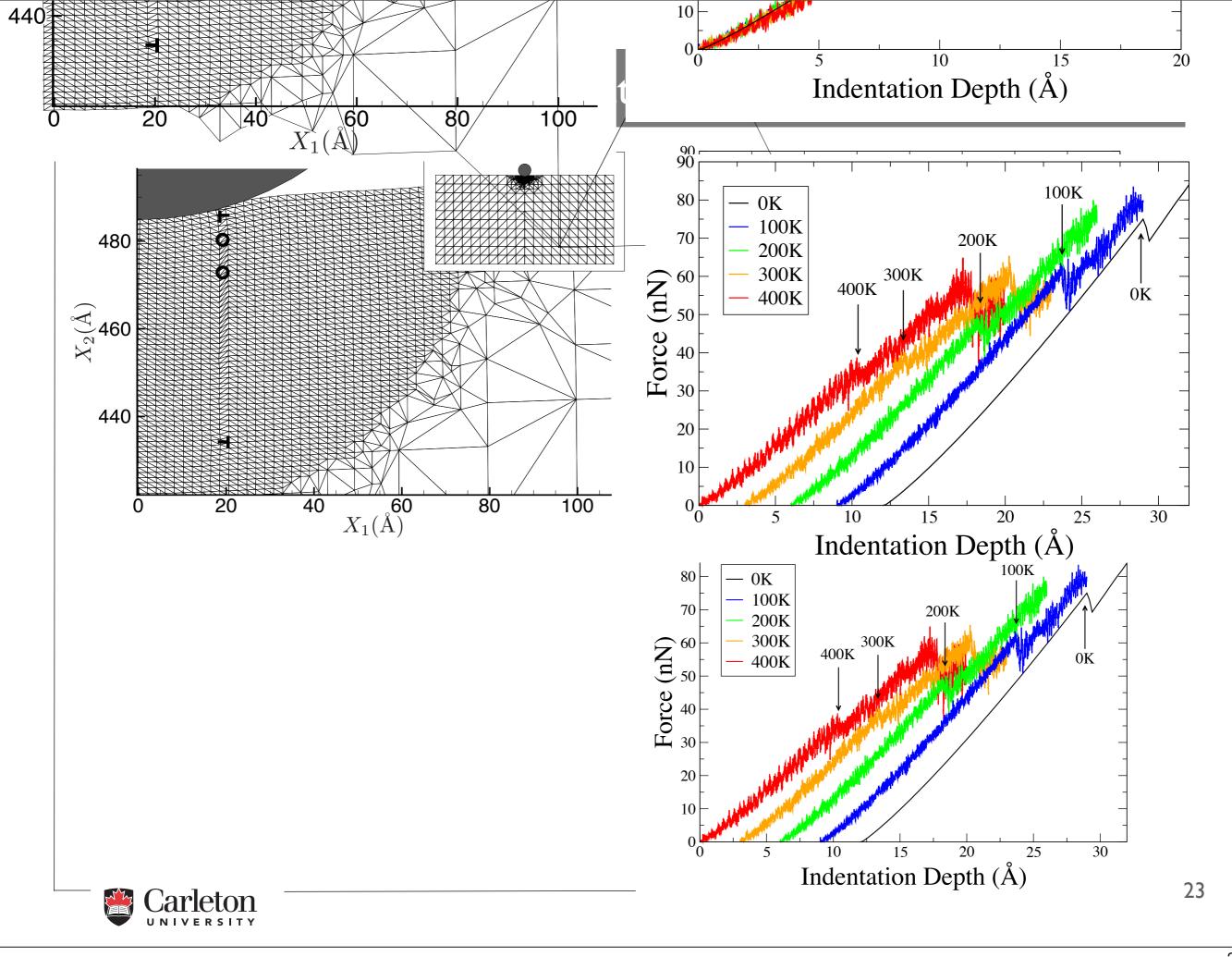


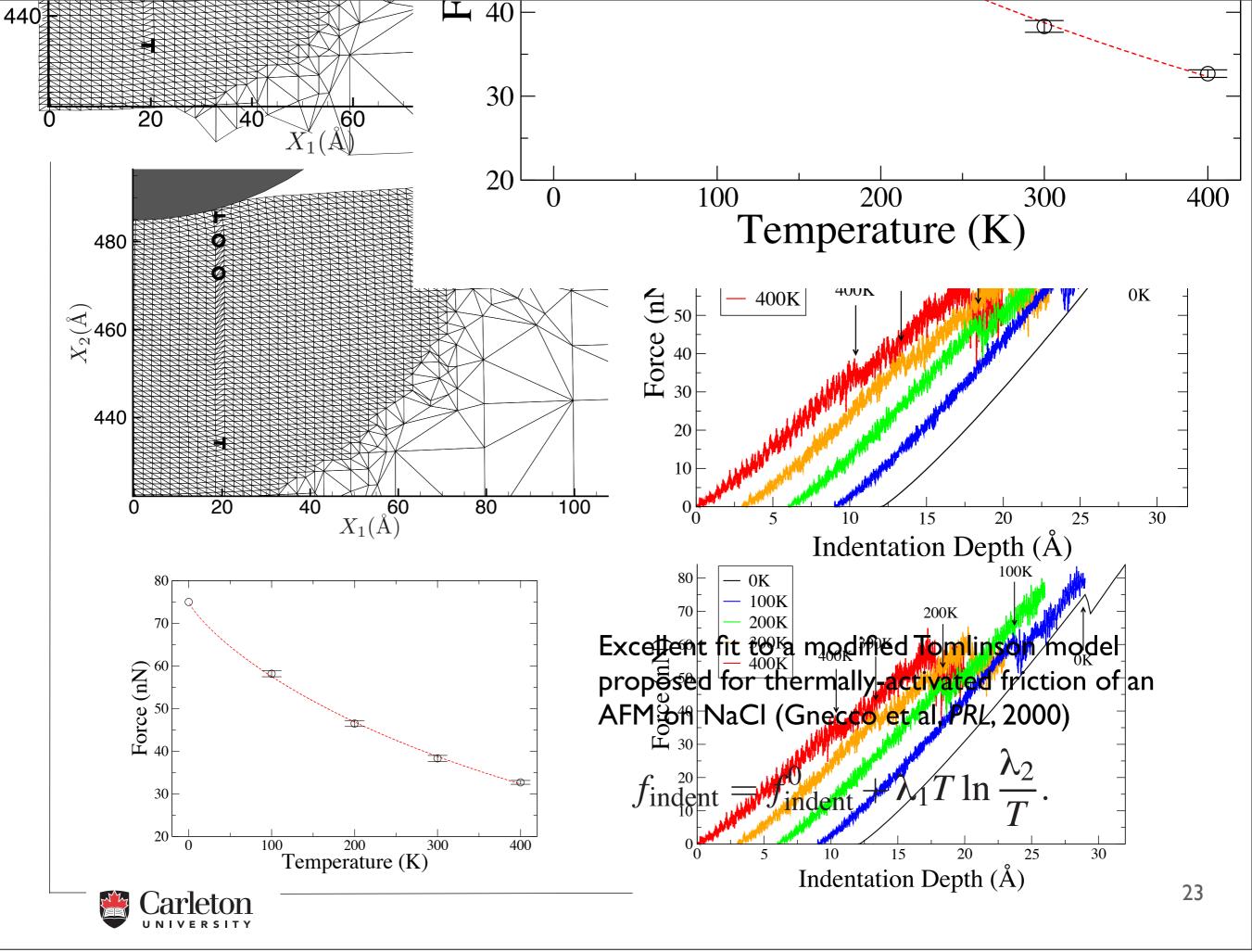
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Credits

Lead Actors

Collaborator on many things (QC, hot QC, comparisons): Ellad Tadmor, U. Minnesota

Post-Docs: Denis Saraev (PDF, currently at large) Behrouz Shiari (PDF, now at NNIN Michigan)

> Collaborators on CADD: Leo Shilkrot (currently at large) Bill Curtin (Brown and EPFL)

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Supporting Cast

Help with all the comparisons of methods: Mitch Luskin, Univ. of Minnesota Michael Parks, LANL Catalin Picu, RPI Dong Qian, U. Cincinnati



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Producers

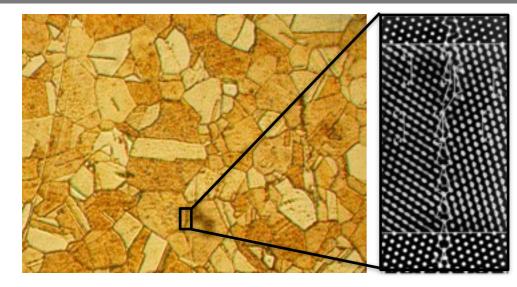




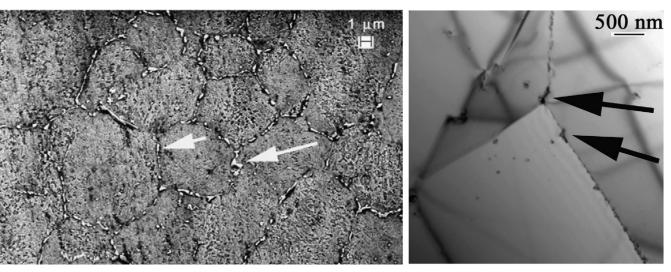
Ontario Premier's Research Excellence Awards (PREA)



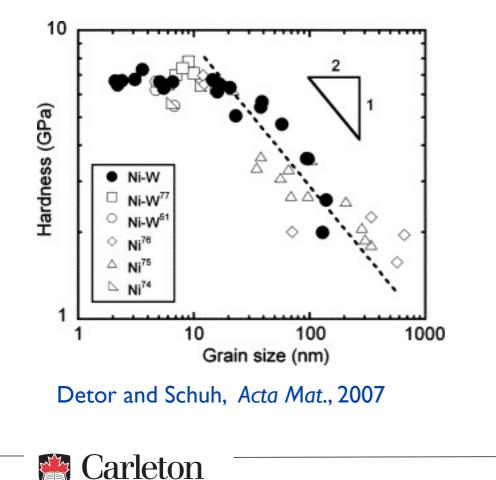
Dislocations, Grain Boundaries and Strength



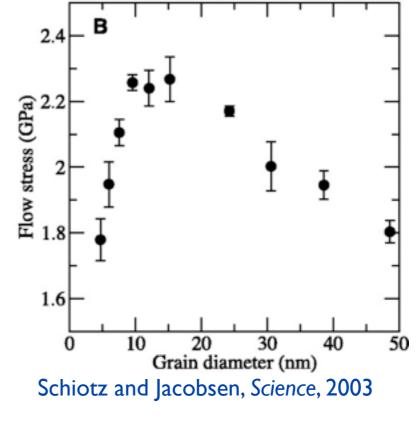
Hu et al., Mater. Chem. Phys., 2002.



Han et al., Mater. Sci. Eng. A, 2009



NIVERSITY



Example 2: The AtC Method

Badia, S., Bochev, P., Lehoucq, R., Parks, M. L., Fish, J., Nuggehally, M., and Gunzburger, M. Int. J. Multiscale Comput. Eng. 5(5), 387–406 (2007). 2 з 8 Handshake Atomistic Continuum I. Introduce blending functions 2. Compute atomic forces 3. Find force on each atom using blending: $\eta^{\alpha,\beta} = \frac{\eta^{\alpha} + \eta^{\beta}}{2}$ $oldsymbol{f}^{lpha}=\sum\eta^{lpha,eta}oldsymbol{f}^{lphaeta}$ $\beta \neq \alpha$ 4. Find nodal forces using complementary blending function: $n_{\rm elem}$ r βCI

$$\mathbf{f}^{I} = -\sum_{e=1} \int_{B_{e}} \Theta(\mathbf{X}) \mathbf{P}(\widetilde{\mathbf{F}}(\mathbf{u})) \frac{\partial S}{\partial \mathbf{X}} dV,$$

5. Impose positions of handshake atoms from interpolated displacements (strong compatibility)

6. Linear elasticity in the continuum



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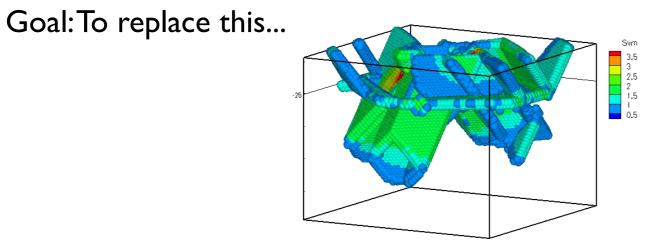
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 - NEMD: non-equilibrium MD (constant energy)



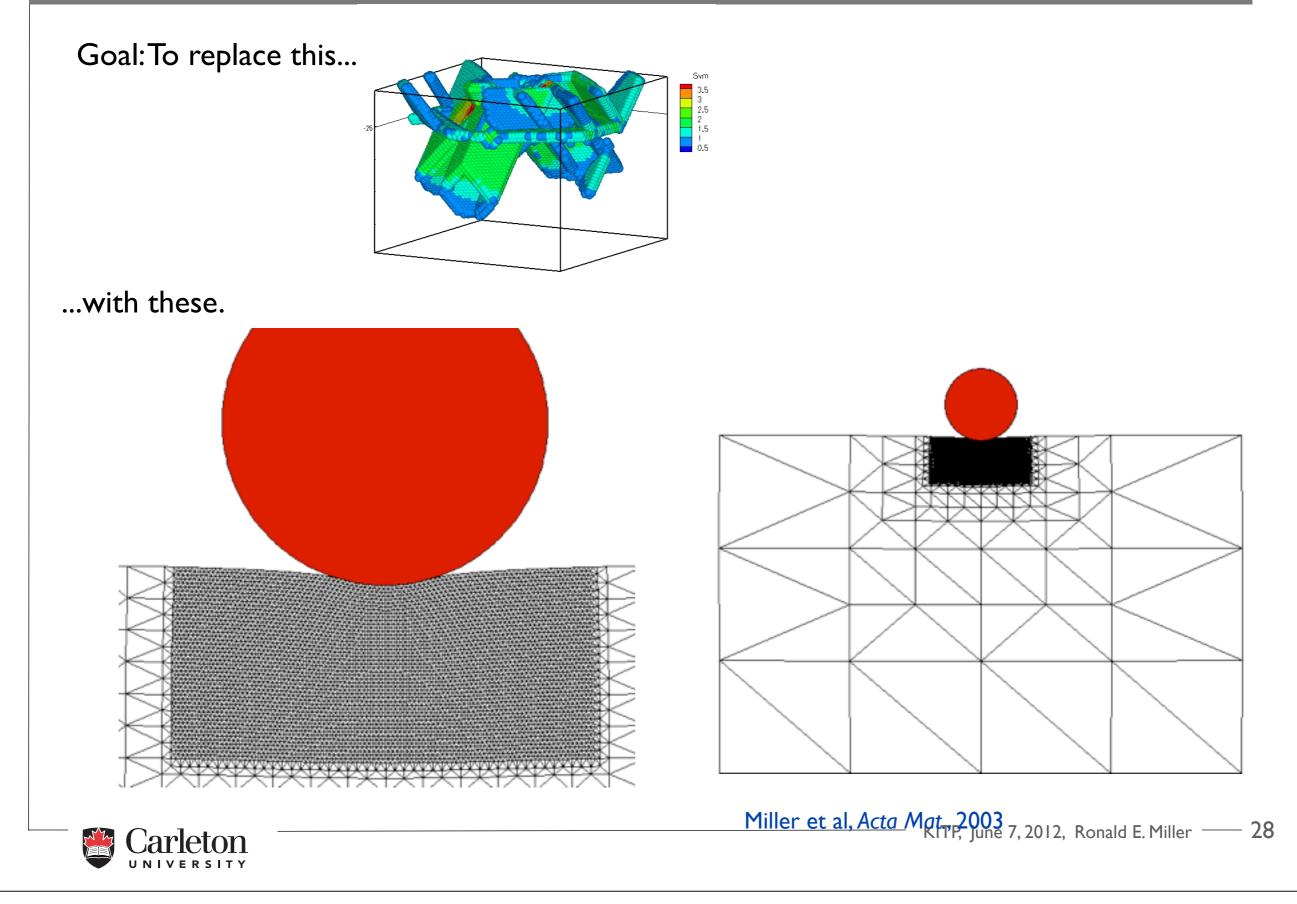
Example: Nano-Indentation using the CADD method



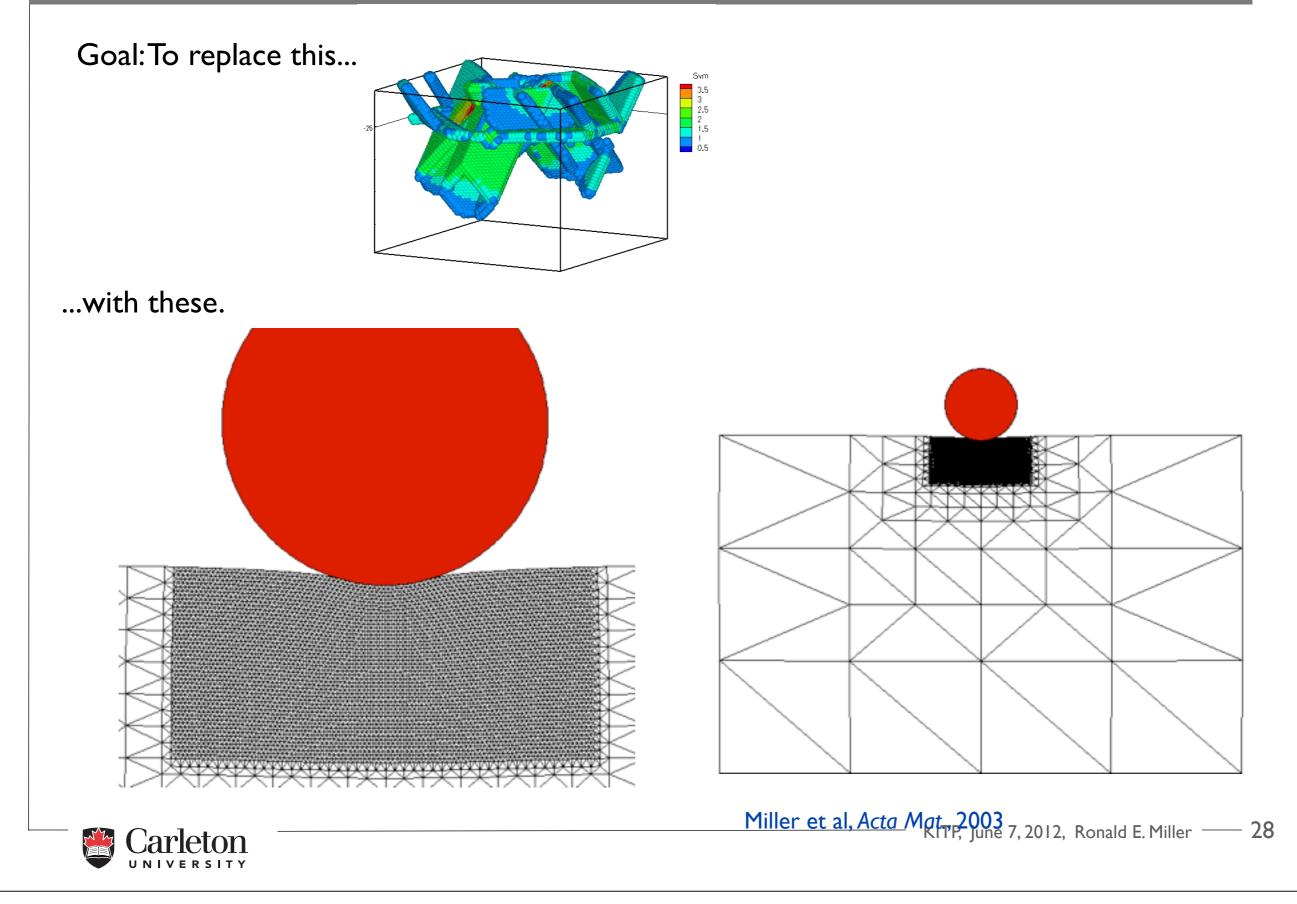
...with these.



Example: Nano-Indentation using the CADD method



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X

Stress-induced dislocation motion

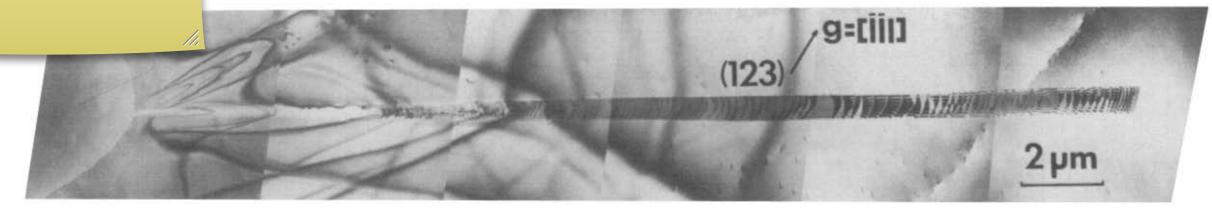
Dislocations move long distances due to high local stresses and low lattice resistance



Ohr is in copper de la fuente is gold.

Stress-induced dislocation motion

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Ohr, Mater. Sci. Eng., 1985

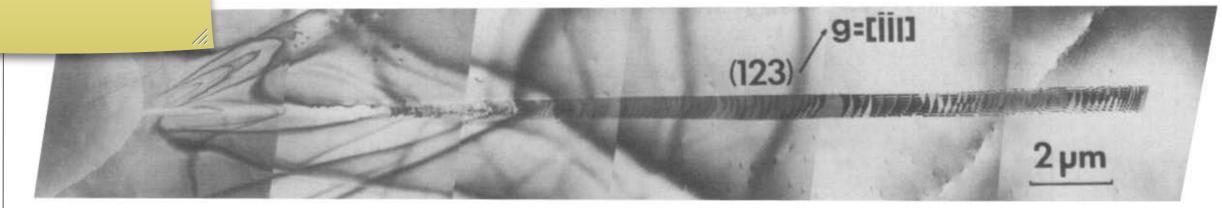
X



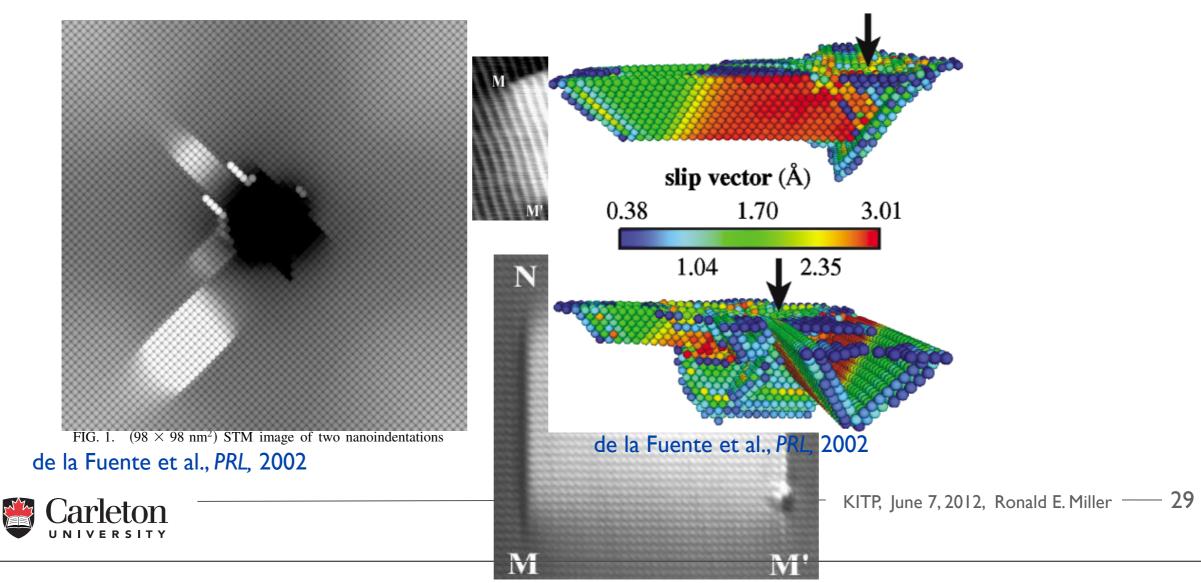
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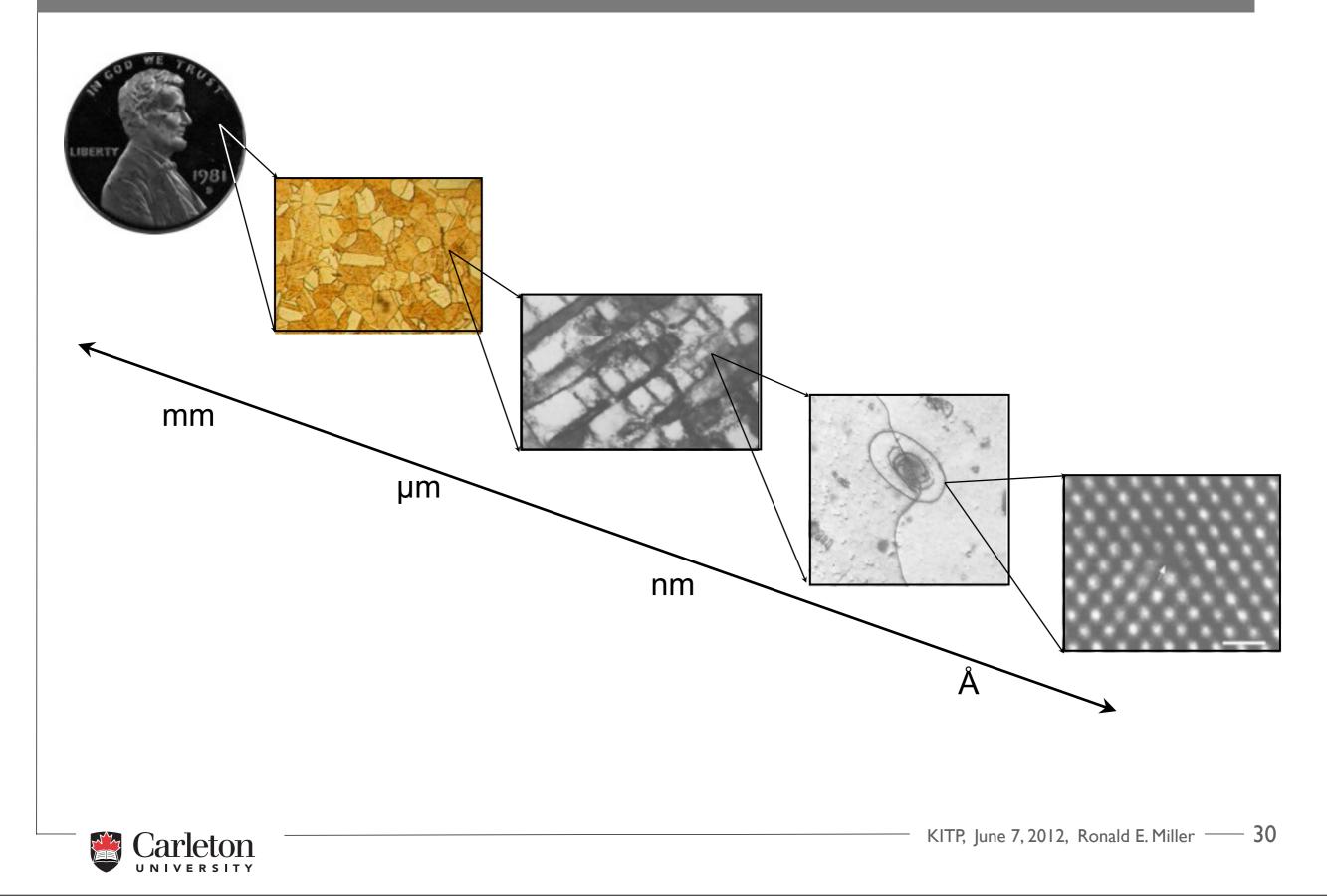
Dislocations move long distances due to high local stresses and low lattice resistance



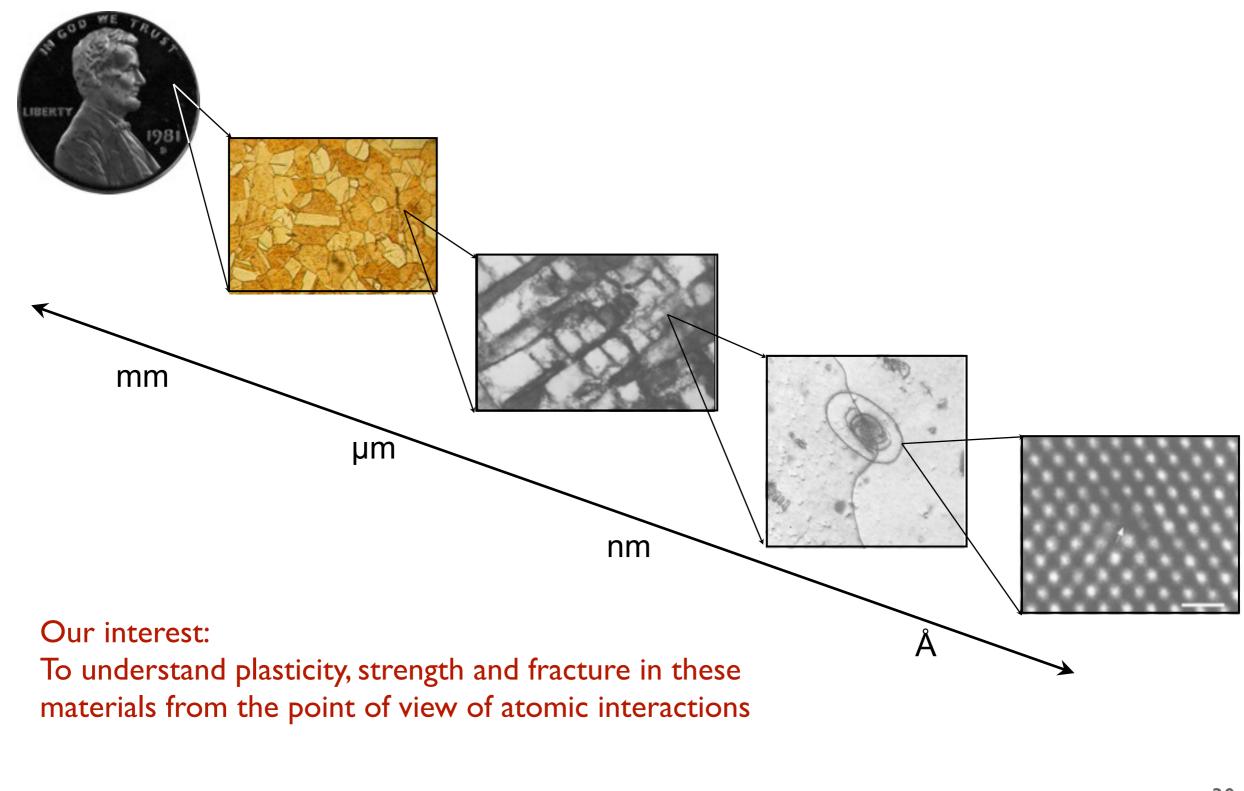
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Length Scales in Crystalline Copper



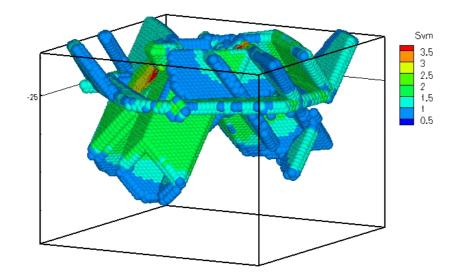
Length Scales in Crystalline Copper



Carleton

Example: Multiscale Nano-Indentation

The goal is to replace this...



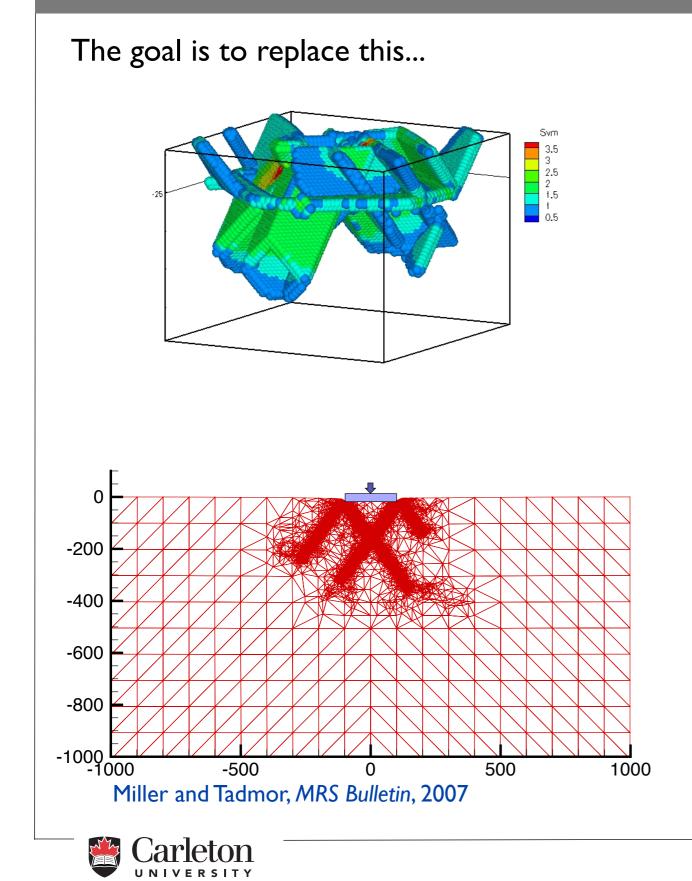
Miller et al., Acta Mat., 2004

Miller and Tadmor, MRS Bulletin, 2007

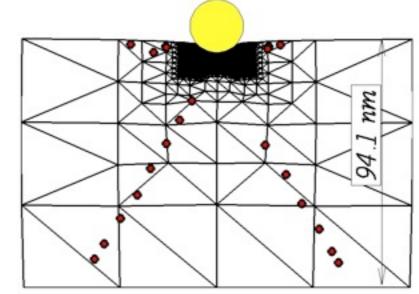


Knap and Ortiz, PRL, 2003

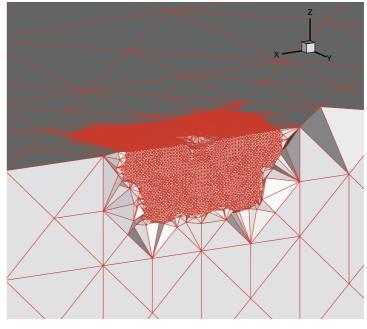
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