

# NSI in oscillation and scattering experiments

Julia Gehrlein

Neutrinos as a portal to new physics and astrophysics

March 2022

# New neutrino interactions (NSI)

- ▶ appear naturally in extensions of SM with new mediators [see Bhupal's talk]
- ▶ focus on NC NSI: general framework [Wolfenstein '78]

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f,P} \epsilon_{\alpha\beta}^f (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{f} \gamma_\mu P f)$$

- ▶ diagonal  $\epsilon_{\alpha\alpha}$  real, off-diagonal  $\epsilon_{\alpha\beta} = |\epsilon_{\alpha\beta}| e^{i\phi_{\alpha\beta}}$
- ▶ signature in **oscillation** and **scattering** experiments

- ▶ framework

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f} \epsilon_{\alpha\beta}^f (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{f} \gamma_\mu P f)$$

- ▶ affect oscillations via **new matter effect**

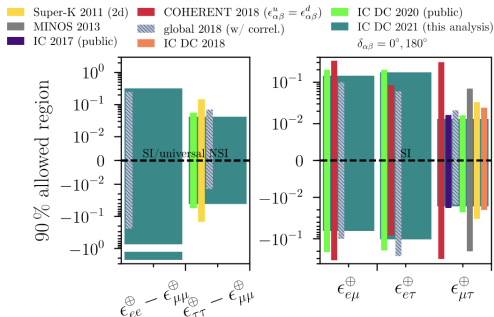
$$H = \frac{1}{2E} \left[ U^\dagger M^2 U + a \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix} \right]$$

matter potential  $a \propto G_F \rho E$

- ▶ can subtract one term on the diagonal  $\Rightarrow$  **8 free** NSI parameter probed by oscillations ( $|\epsilon_{\alpha\beta}|$ ,  $\phi_{\alpha\beta}$ ,  $\epsilon_{ee} - \epsilon_{\mu\mu}$ ,  $\epsilon_{\tau\tau} - \epsilon_{\mu\mu}$ )
- ▶ parameters related to parameters in Hamiltonian by  $\epsilon_{\alpha\beta} = \sum_f \epsilon_{\alpha\beta}^{f,V} \langle N_f(x) / N_e(x) \rangle$

## How to measure $\epsilon$ with oscillations?

- ▶ effect of NSI scales with energy, baseline, and matter density  
 $\Rightarrow$  use high-energy neutrino sources with long baselines  $\rightarrow$   
 atmospheric neutrinos at IceCube



[IceCube '21]

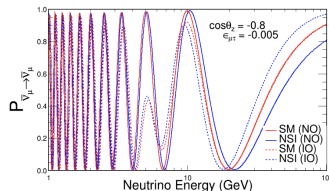
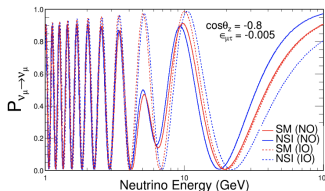
## How to measure $\epsilon$ with oscillations?

►  $\epsilon_{\mu\tau}$ :

Do we need to observe tau neutrinos?

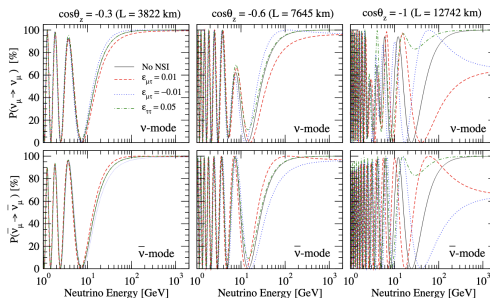
long-baseline and atmospheric oscillations are dominated by  $\nu_3$  which contains  $\nu_\mu$  and  $\nu_\tau$

→ use **muon neutrino disappearance** data from atmospheric neutrinos or long baseline experiments



## How to measure $\epsilon$ with oscillations?

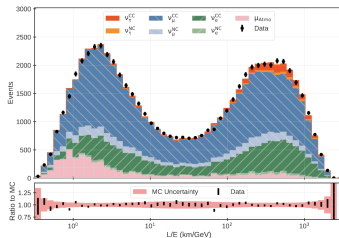
- $\epsilon_{\tau\tau} - \epsilon_{\mu\mu}$ :  
 diagonal NSI parameter  $\rightarrow$  affects atmospheric mass splitting  $\rightarrow$   
 compare results for  $\Delta m_{32}^2$  from **atmospheric muon neutrino disappearance** to measurements of  $\Delta m_{32}^2$  from reactor experiments in vacuum



[Miranda, Nunokawa '15]

## How to measure $\epsilon$ with oscillations?

- ▶  $\epsilon_{e\mu}$ :  
 use **electron appearance from atmospheric muon neutrinos**  
 cascade signature of  $\nu_e$ , low  $\nu_e$  yield  
 use muon appearance from atmospheric  $\nu_e$ : low atmospheric electron component
- ▶  $\epsilon_{e\tau}$ : (no tau neutrinos required)  
 use **electron appearance from atmospheric muon neutrinos**



[IceCube '19]

## How to measure $\epsilon$ with oscillations?

▶  $\epsilon_{ee} - \epsilon_{\mu\mu}$ :

no source of high energy  $\nu_e$ !

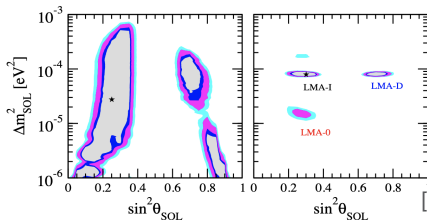
**LMA Dark solution**: transformation

$$\sin \theta_{12} \rightarrow \cos \theta_{12}, \quad \Delta m_{31}^2 \rightarrow -\Delta m_{32}^2, \quad \delta \rightarrow \pi - \delta,$$

$$(\epsilon_{ee} - \epsilon_{\mu\mu}) \rightarrow (\epsilon_{ee} - \epsilon_{\mu\mu}) - 2, \quad (\epsilon_{\tau\tau} - \epsilon_{\mu\mu}) \rightarrow -(\epsilon_{\tau\tau} - \epsilon_{\mu\mu}),$$

$$\epsilon_{\alpha\beta} \rightarrow -\epsilon_{\alpha\beta}^*$$

leaves neutrino evolution invariant [Miranda, Tortola, Valle '04, Coloma, Schwetz '16]

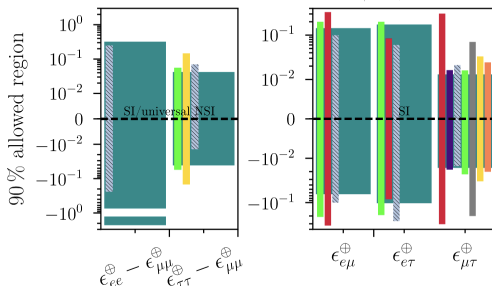


[Miranda, Tortola, Valle '04]



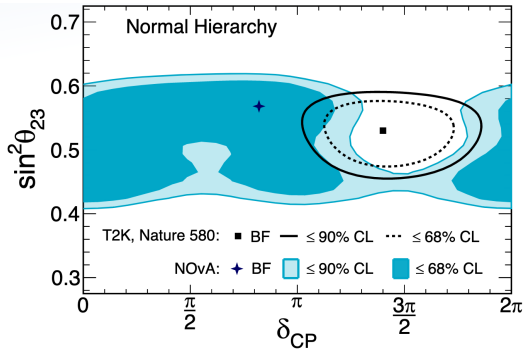
## How to measure $\epsilon$ with oscillations?

- Super-K 2011 (2d)
  - COHERENT 2018 ( $\epsilon_{\alpha\beta}^u = \epsilon_{\alpha\beta}^d$ )
  - IC DC 2020 (public)
  - MINOS 2013
  - global 2018 (w/ correl.)
  - IC DC 2021 (this analysis)
  - IC 2017 (public)
  - IC DC 2018
- $\delta_{\alpha\beta} = 0^\circ, 180^\circ$



[IceCube '21]

## Hint for non-zero $\epsilon$ ?



[Himmel '20]

→ **disagreement** at  $\sim 2\sigma$ !

- ▶ difference between NOvA and T2K are baselines and matter densities → neutrinos at NOvA experience stronger **matter effects**

⇒ Can NSI resolve this tension?

**CP-Violating Neutrino Non-Standard Interactions in Long-Baseline-Accelerator Data**

Peter B. Denton,<sup>1,\*</sup> Julia Gehrlein,<sup>1,†</sup> and Rebekah Pestes<sup>1,2,‡</sup>

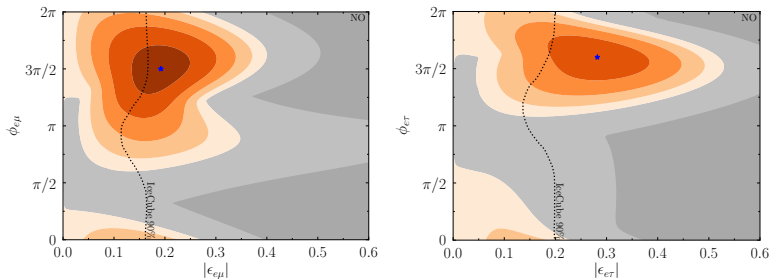
[arxiv: 2008.01110, Phys. Rev. Lett. 126 (2021) no.5, 051801]

- ▶ focus on off-diagonal NSI parameters
- ▶ derived analytical and numerical results

$$|\epsilon_{e\beta}| \approx \frac{s_{12}c_{12}c_{23}\pi\Delta m_{21}^2}{2s_{23}w_\beta} \left| \frac{\sin\delta_{\text{T2K}} - \sin\delta_{\text{NOvA}}}{a_{\text{NOvA}} - a_{\text{T2K}}} \right| \approx \begin{cases} 0.22 & \text{for } \beta = \mu \\ 0.24 & \text{for } \beta = \tau \end{cases},$$

with  $w_\beta = \sin\theta_{23}(\cos\theta_{23})$

- ▶ used disappearance, appearance data from NOvA and T2K, information from vacuum experiments for  $\theta_{13}$ ,  $|\Delta m_{32}^2|$ ,  $\theta_{12}$ ,  $\Delta m_{21}^2$



orange preferred over SM at integer values of  $\Delta\chi^2$ , dark gray disfavored at  $\Delta\chi^2 = 4.61$  [see also Chatterjee, Palazzo '20]

- ▶ complex NSI with  $|\epsilon| \approx 0.2$ ,  $\phi \approx 3\pi/2$ ,  $\delta \approx 3\pi/2$ , NO can fully resolve tension
- ▶ change in MO from NO to IO: improvement of  $\Delta\chi^2 = 2.3$  [Kelly et al '20; Esteban et al '20]

- ▶ framework

$$\mathcal{L}_{NSI} = -2\sqrt{2}G_F \sum_{\alpha,\beta,f} \epsilon_{\alpha\beta}^f (\bar{\nu}_\alpha \gamma^\mu \nu_\beta) (\bar{f} \gamma_\mu P f)$$

→ NSI affect neutrino-SM scattering experiments like CEvNS (coherent elastic neutrino nucleus scattering)

- ▶ NSI effect on weak charge

$$Q_{W\alpha}^2 \propto \left[ N(g_n^V + \epsilon_{\alpha\alpha}^V) \right]^2 + \sum_{\beta \neq \alpha} \left[ \epsilon_{\alpha\beta}^V (Z + N) \right]^2$$

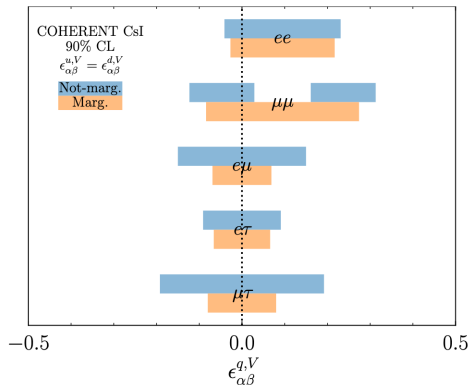
→ no sensitivity to complex NSI phase

- ▶ sensitivity to 5 NSI parameters (not sensitive to  $\epsilon_{\tau\tau}$ )

- ▶ NSI in oscillations:  
forward scattering interactions  $\rightarrow$  constraints **independent of mediator masses** (unless  $m \sim d_{max}^{-1}$  [Wise, Zhang '18])
- ▶ NSI in scattering:  
**dependence on mediator mass**
  - ▶ heavy mediator regime ( $m^2 > q^2 \sim (10 \text{ MeV})^2$ ):  $\epsilon \propto g_X^2/m^2$
  - ▶ light mediator regime  $\epsilon \propto g_X^2/(q^2 + m^2)$

# NSI in scattering: heavy mediators

- ▶ CEvNS detected for the first time in 2016 by COHERENT
- ▶ NSI constraints apply to mediators  $m > 10$  MeV
- ▶ constraints derived in [Denton, JG '20] using Feldman-Cousins framework

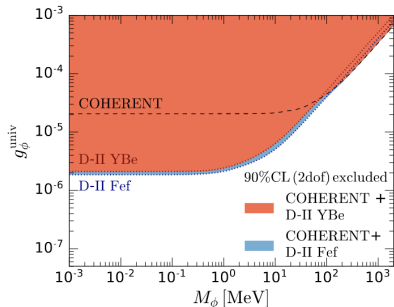




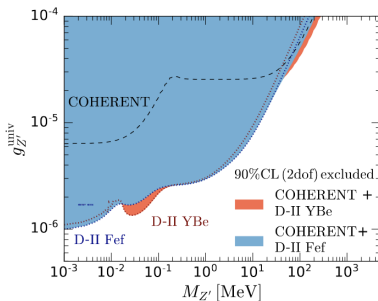
# NSI in scattering: light mediators

use CEvNS at low energy accelerator and reactor CEvNS

scalar mediator



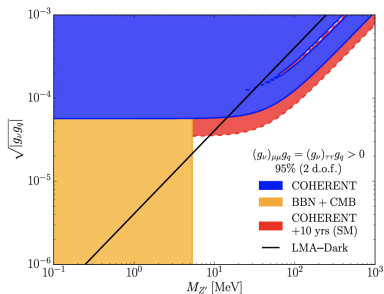
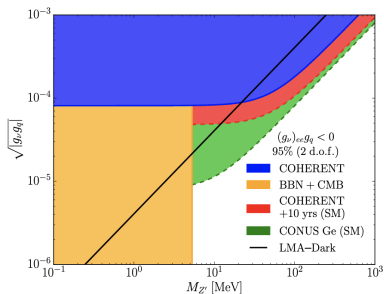
vector mediator



[Coloma et al '22]

# NSI in scattering: light mediators

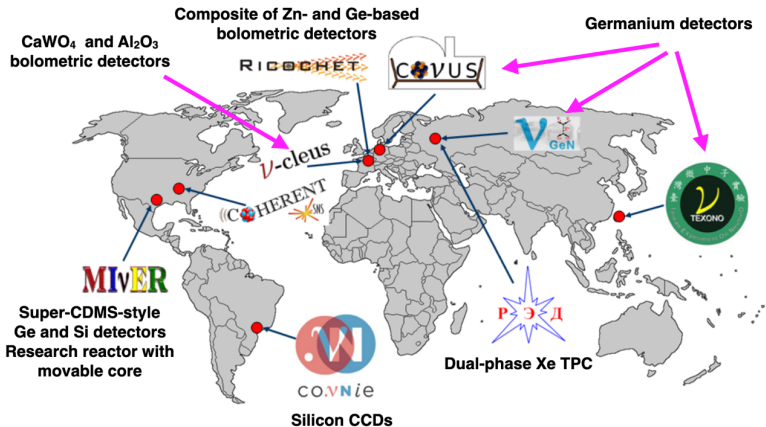
connection to oscillation phenomenology



[Denton, Shoemaker, Farzan '18]

⇒ scattering experiments can probe degeneracies in oscillation parameters in presence of NSI

# NSI in scattering: future



- ▶ NSI can be tested in oscillation experiments and scattering experiments → complementarity between different probes
- ▶ currently weak hint for non-zero NSI from oscillations
- ▶ future atmospheric and long baseline experiments as well as scattering experiments will probe NSI further

# Thank you for your attention!

