

Interpretation of astrophysical neutrino signal, and the multi-messenger context

(a particle physics theorist's perspective)

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KITP conference

Neutrinos: Recent developments

... and future challenges

Nov 3-7, 2014

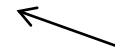
Contents

- > Introduction
- > Simulation of neutrino sources, detector response
- > Interpretations of signal
 - Expected (guaranteed?) contributions
 - Conceptual insights, **UHECR connection**
 - Astrophysical object *speculations*
- > Future challenges:
 - The question of flavor!?
 - Physics case for high-energy extension of IceCube?
- > Summary

KITP conference

*Neutrinos:
Recent
developments*

***... and future
challenges***



Cosmic messengers

Theory
(source
distribution)

Physics of astrophysical
neutrino sources = physics of
cosmic ray sources

Multi-messenger interpretations
must rely on theory (acceleration,
radiation processes, particle escape, geometry, ...)

Theory
(radiation
model)

Large
astrophysical
uncertainties

Theory
(infrared
etc BGs)

Theory
(magnetic
fields, ...)

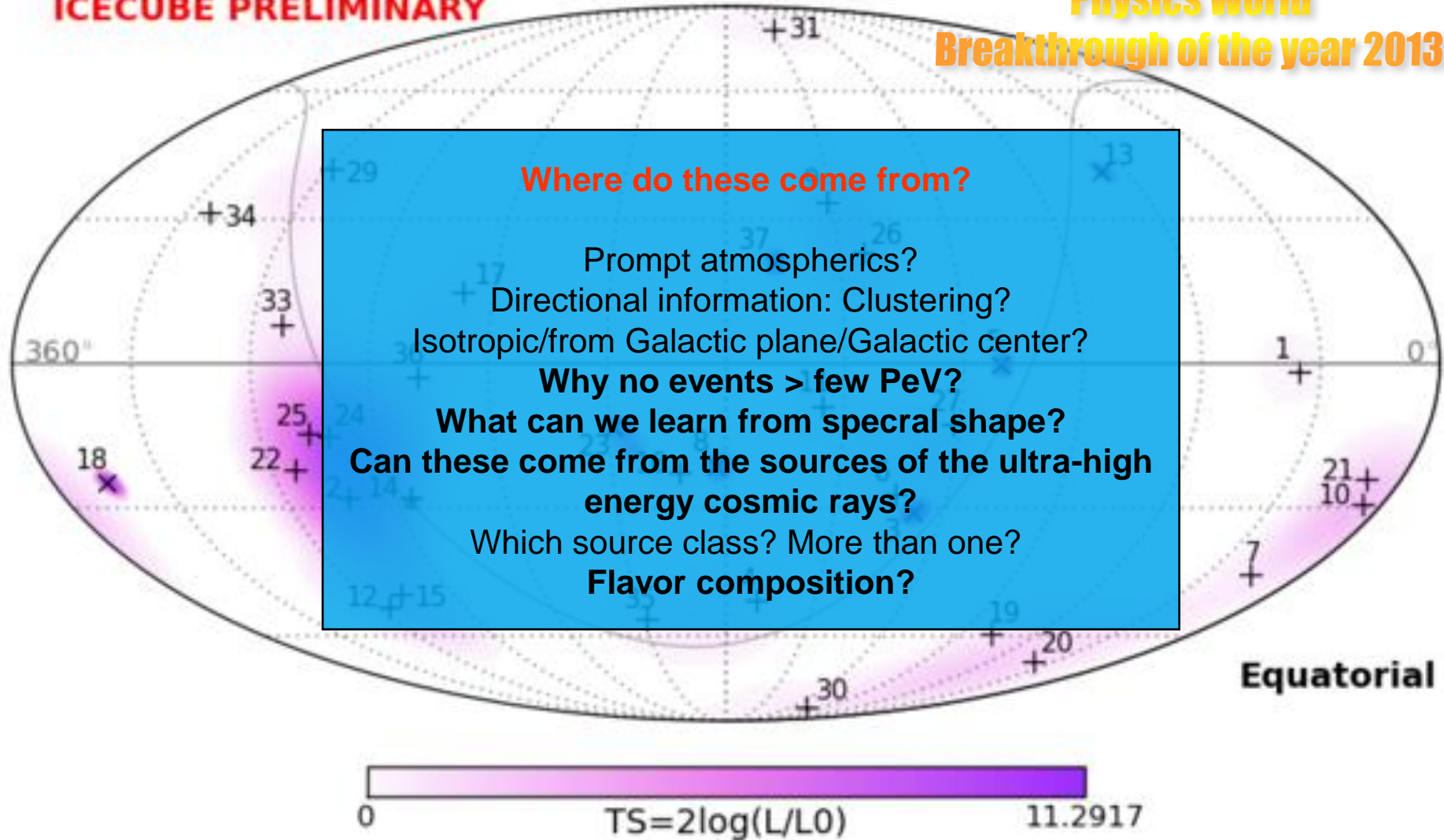
Astrophysical
beam dump



2014: 37 neutrinos in the TeV-PeV range

ICECUBE PRELIMINARY

Physics World
Breakthrough of the year 2013



Science 342 (2013) 1242856; update by Gary Hill @ Neutrino 2014

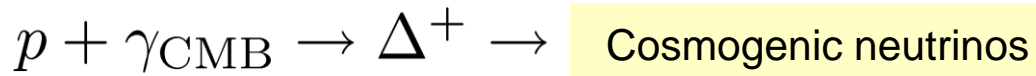
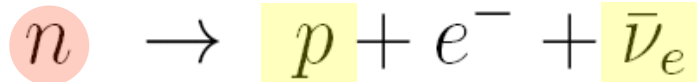


Simulation of neutrino sources

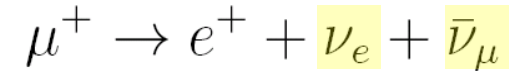
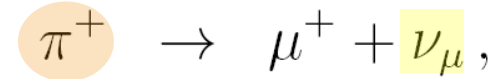


Neutrino and cosmic ray source (illustrative proton-only scenario, $p\gamma$ interactions)

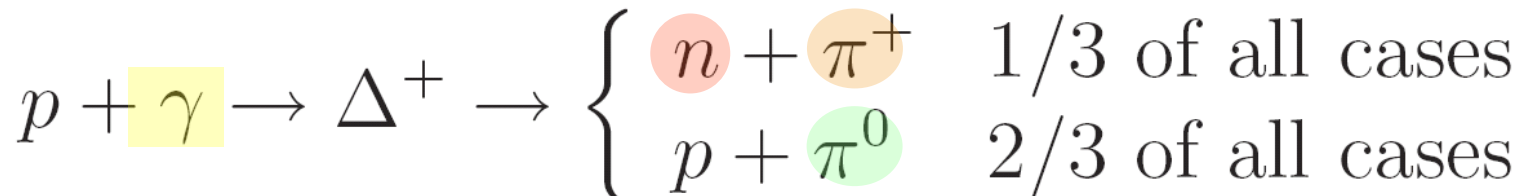
If neutrons can escape:
Source of cosmic rays



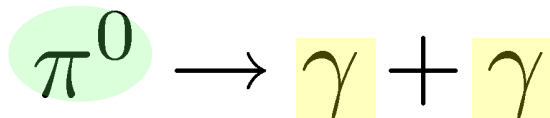
Neutrinos produced in
ratio $(\nu_e:\nu_\mu:\nu_\tau)=(1:2:0)$



Delta resonance approximation:



π^+/π^0 determines ratio between neutrinos and high-E gamma-rays

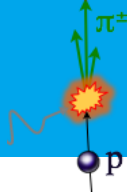


Cosmic messengers

High energetic gamma-rays;
typically cascade down to lower E
Additional constraints!



Source simulation: $p\gamma$ (particle physics)



> $\Delta(1232)$ -resonance approximation:

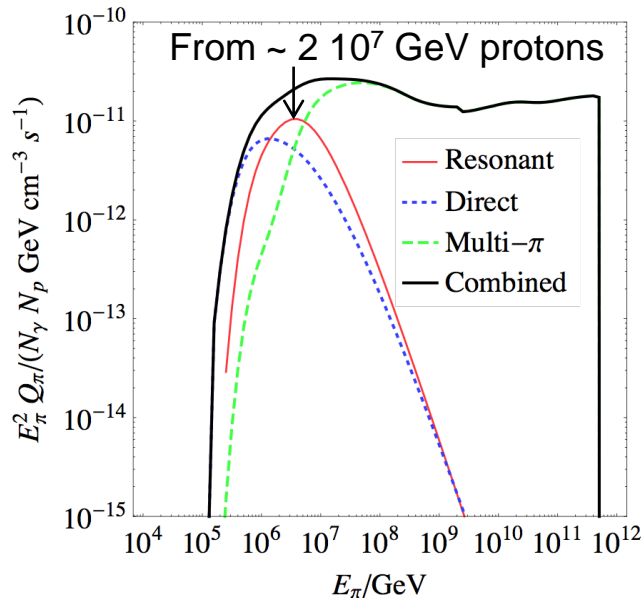
$$p + \gamma \rightarrow \Delta^+ \rightarrow \begin{cases} n + \pi^+ & 1/3 \text{ of all cases} \\ p + \pi^0 & 2/3 \text{ of all cases} \end{cases}$$

Resonance condition: Proton energy [GeV] x Photon energy [GeV] ~ 0.2

> Limitations:

- No π^- production; cannot predict π^+/π^- ratio (Glashow resonance!)
- High energy processes affect spectral shape (X-sec. dependence!)
- Low energy processes (t-channel) enhance charged pion production

> Example: Peak at PeV energies (here pions) from thermal target photons?



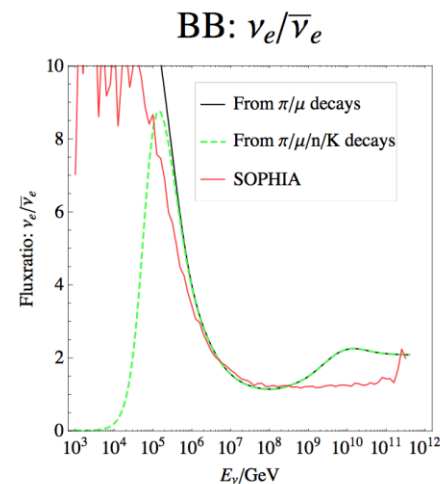
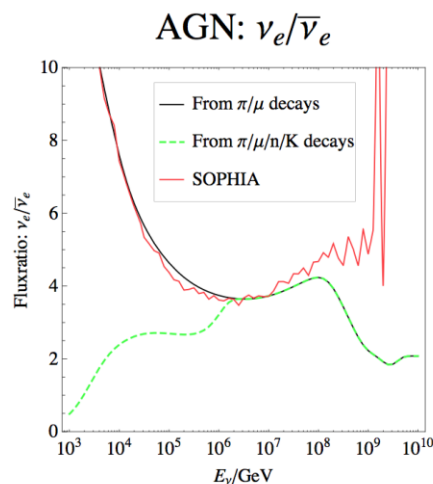
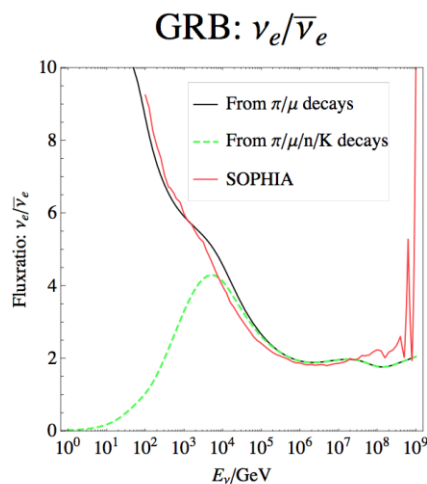
Example: E^{-2} proton spectrum interacting with 10 eV blackbody target photon spectrum
Multi-pion processes exclude this as explanation for observed neutrinos unless $E_{p,max} < 10^8$ GeV!

From: Hümmer et al, *ApJ* 721 (2010) 630
parameterization based on SOPHIA –
Mücke, Rachen, Engel, Protheroe, Stanev, 2000



pp versus p γ interactions: commons and differences

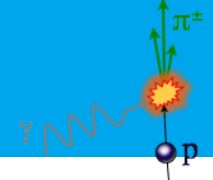
- > In pp and p γ interactions, the secondary pions take about 20% of the proton energy, the neutrinos about 5% (per flavor)
 - PeV neutrinos must come from 20 PeV (pp/p γ) to 1 EeV (Fe-p)
- > The charged to neutral pion ratio is roughly 50-50 (both pp and p γ)
- > The spectral shape of the neutrinos follows the primary nuclei for pp interactions, for p γ it depends on that of the target photons as well
- > There are typically (not necessarily) more electron neutrinos than antineutrinos at the source for p γ interactions (if Glashow resonance used for discrimination, proposed in [Anchordoqui et al, hep-ph/0410003](#); [Barger et al, 1407.3255](#) etc!)



Δ -resonance prediction: infinity!



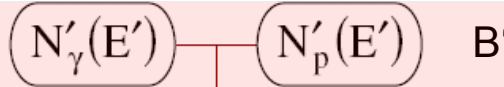
Neutrino production (example: $p\gamma$)



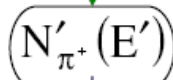
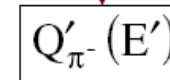
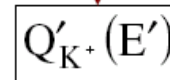
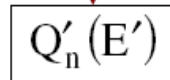
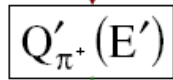
Dashed arrows: kinetic equations include cooling and escape

$Q(E)$ [$\text{GeV}^{-1} \text{cm}^{-3} \text{s}^{-1}$]
per time frame
 $N(E)$ [$\text{GeV}^{-1} \text{cm}^{-3}$]
steady spectrum

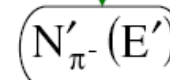
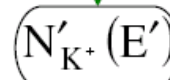
Input \Rightarrow Object-dependent
 \Rightarrow Astrophysics!



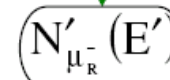
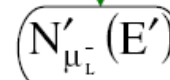
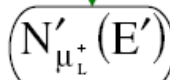
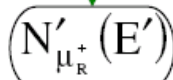
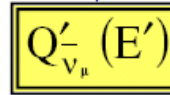
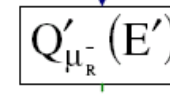
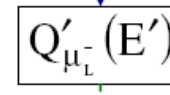
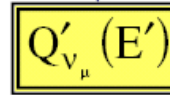
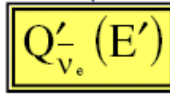
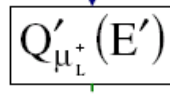
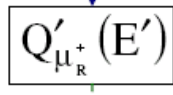
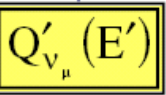
photohadronics



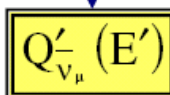
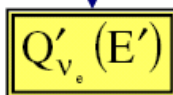
Optically
thin
to neutrons



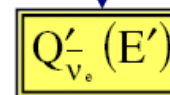
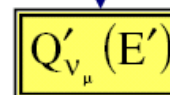
weak decays



weak decays



from:
Baerwald, Hümmer, Winter,
Astropart. Phys. 35 (2012) 508



Kinetic equations (steady state)

- > Treat energy losses/escape in continuous limit:

$$Q(E) = \frac{\partial}{\partial E} (b(E) N(E)) + \frac{N(E)}{t_{\text{esc}}}$$

Injection

Energy losses

Escape

$$b(E) = -E t_{\text{loss}}^{-1}$$

$Q(E,t)$ [$\text{GeV}^{-1} \text{cm}^{-3} \text{s}^{-1}$] injection per time frame (e. g. from acc. zone)

$N(E,t)$ [$\text{GeV}^{-1} \text{cm}^{-3}$] particle spectrum including spectral effects

NB: Need $N(E)$ to compute particle interactions

- > Simple case: No energy losses $b=0$: $N(E) = Q(E) t_{\text{esc}}$

- > Special cases:

- $t_{\text{esc}} \sim R/c$ (free-streaming, aka “leaky box“)
- $t_{\text{esc}} \sim E^{-\alpha}$. Consequence: $N(E) \sim Q_{\text{inj}}(E) E^{-\alpha}$, Escape: $Q_{\text{esc}}(E) = N(E)/t_{\text{esc}} \sim Q_{\text{inj}}$

(Neutrino spectrum from $N(E)$ can have a break which is not present in escaping primaries $Q_{\text{esc}}(E)$)



In the presence of strong B: Secondary cooling

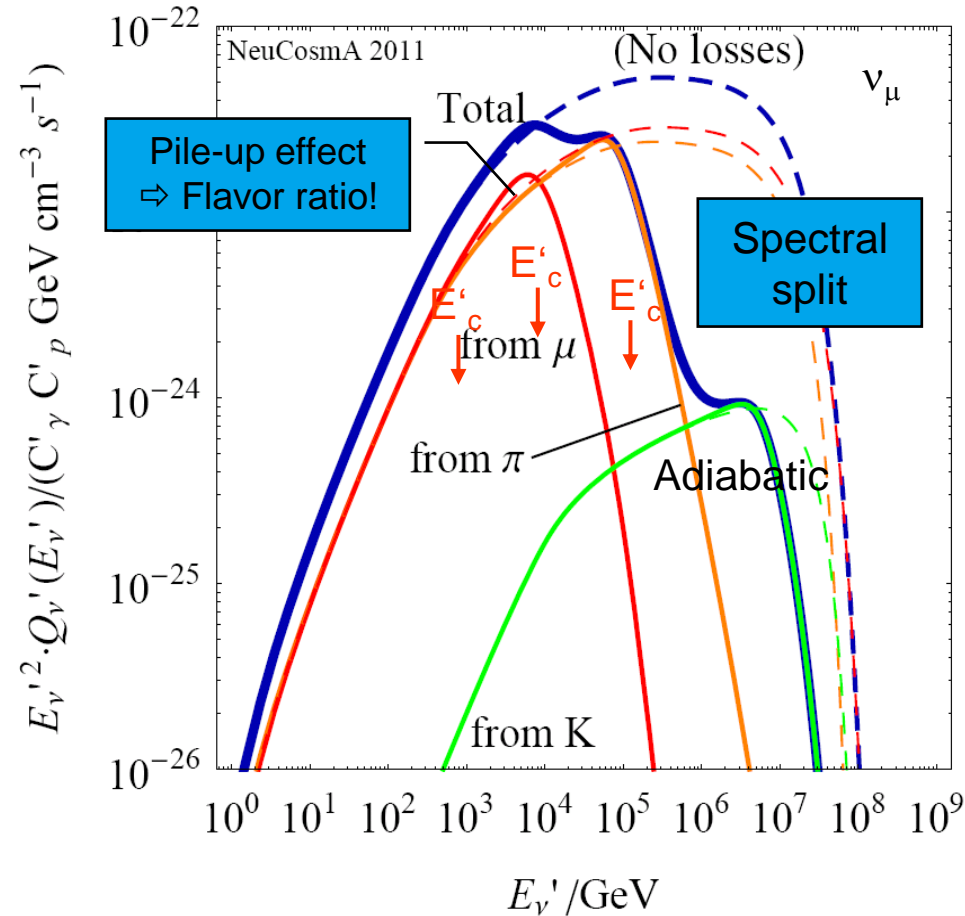
Example: GRB

Secondary spectra (μ , π , K) loss-steepend above critical energy

$$E'_c = \sqrt{\frac{9\pi\epsilon_0 m^5 c^7}{\tau_0 e^4 B'^2}}$$

- E'_c depends on particle physics only (m , τ_0), and \mathbf{B}'
- Leads to characteristic flavor composition and shape
- **Very robust prediction for sources?** [i.e. any additional radiation processes mainly affecting the primaries will not affect the flavor composition]

Decay/cooling: charged μ , π , K



Baerwald, Hümmer, Winter, *Astropart. Phys.* **35** (2012) 508;
 also: Kashti, Waxman, 2005; Lipari et al, 2007

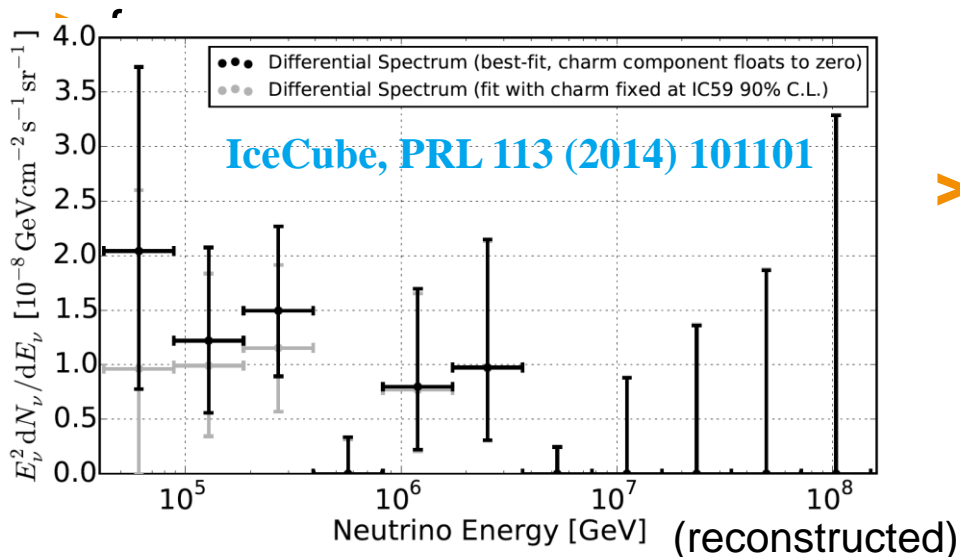
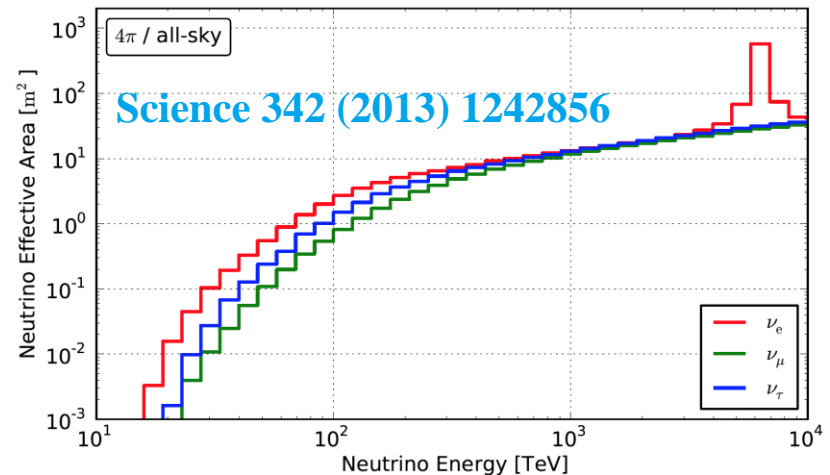


Neutrino propagation and detector response

- > Neutrino propagation: Flavor mixing $P_{\alpha\beta} = \sum_{i=1}^3 |U_{\alpha i}|^2 |U_{\beta i}|^2$
- > Event rate computation: $N = \int_{E_{\min}}^{E_{\max}} \phi_{\nu}(E) A_{\text{eff}}(E) t_{\text{obs}} [d\Omega] dE$

(A_{eff} typically includes analysis cuts;
 ϕ in units of $\text{cm}^{-2} \text{s}^{-1} \text{GeV}^{-1} [\text{sr}^{-1}]$)

- > A complication for interpretations:
 neutrino energy reconstruction
 (incident energy \rightarrow deposited
 energy \rightarrow reconstructed energy)



- > Reconstructed distribution does not exhibit statistically significant gap, cutoff seems to be more evident

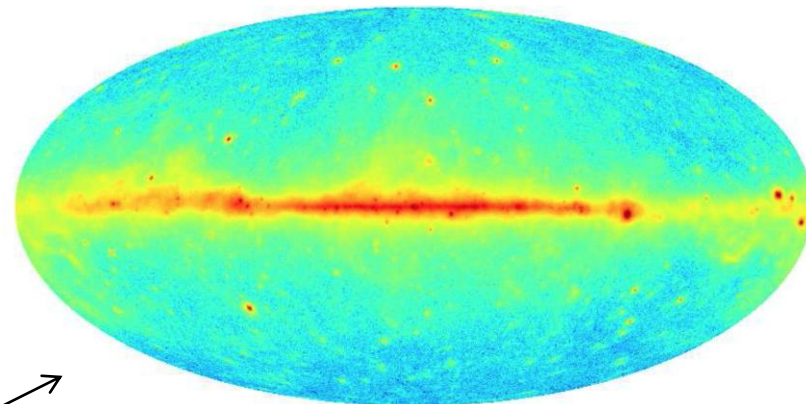


On the signal interpretation



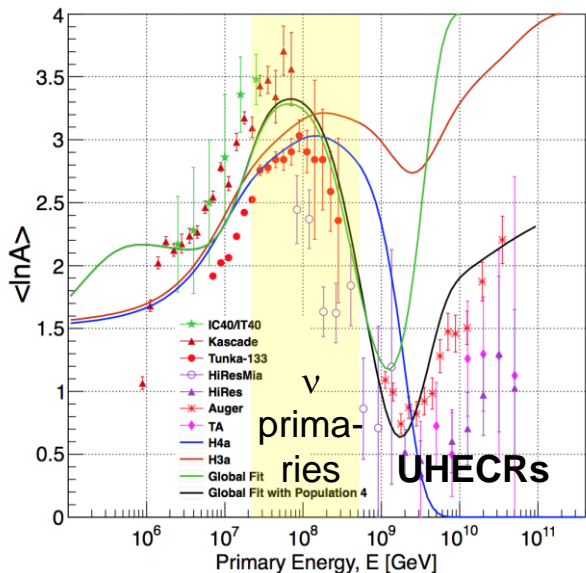
Galactic “guaranteed” contribution?

- Cosmic rays interact with hydrogen in our Galaxy
- Cosmic ray density from local observations; extension of production region can be inferred from diffuse gamma-ray observations (very narrow around Galactic plane)

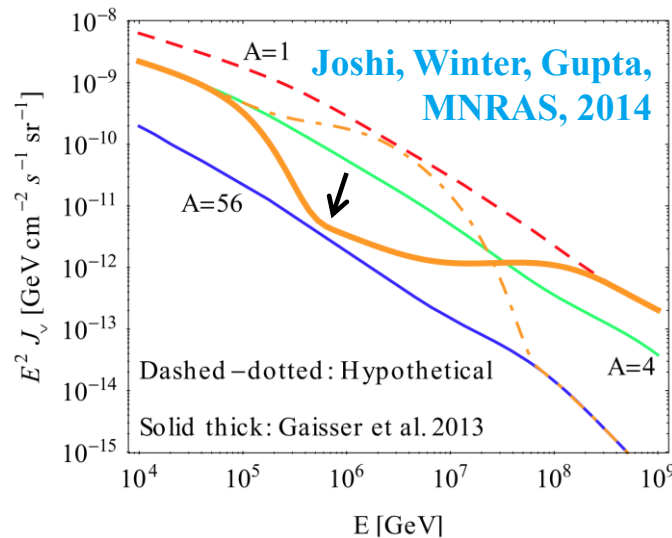


Fermi-LAT, ApJ 750 (2012) 3
(see also arXiv:1410.3696)

- Complication: the CR composition changes non-trivially in relevant range:



Gaisser, Stanev, Tilav, 2013



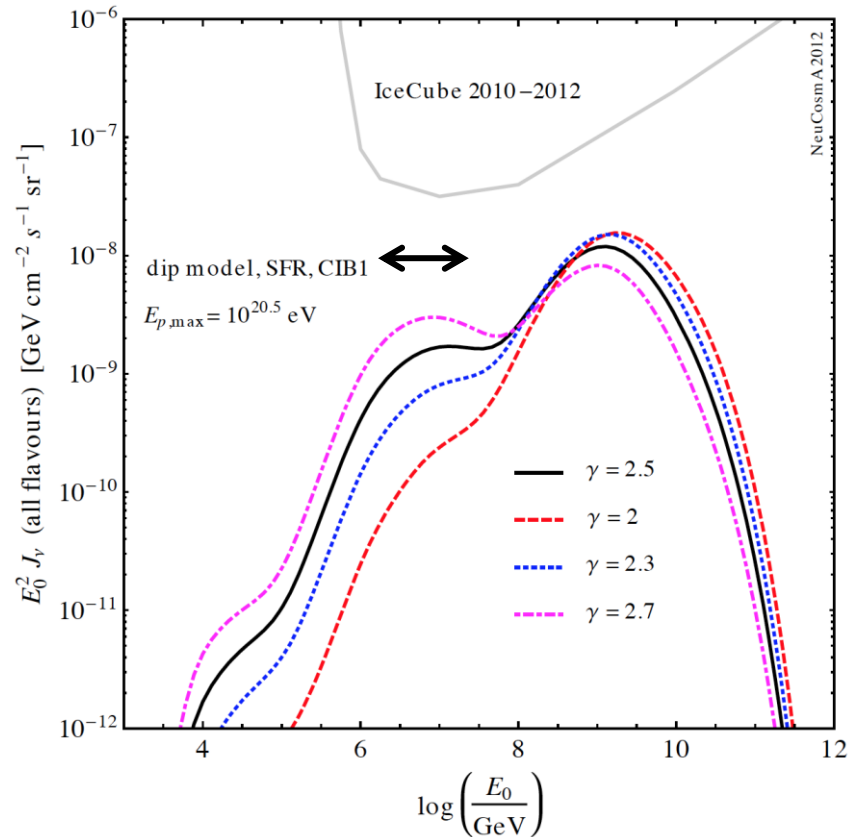
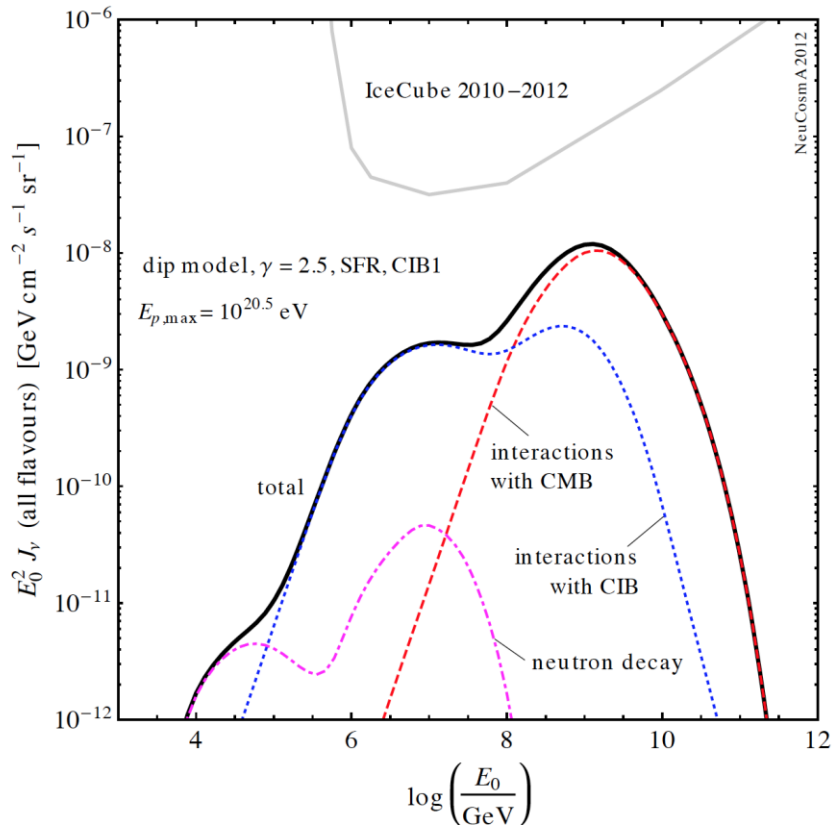
(all-sky averaged prediction)

O(0.1-1) event plausible to satisfy n_H , composition and γ -ray constraints



Cosmogenic (from CR propagation) origin?

- PeV neutrinos from cosmic infrared background interactions of UHECRs (here: protons); depends somewhat on model (high-z evolution?)
- Even if protons, soft spectra etc, difficult to reach required flux



Figures: Bustamante, Evoli, Sigl, WW, in prep. See also Roulet, Sigl, van Vliet, Mollerach, JCAP 1301 (2013) 028



Conceptual insights? Neutrinos from Ap interactions

... in cosmologically distr. sources (star formation rate evolution)

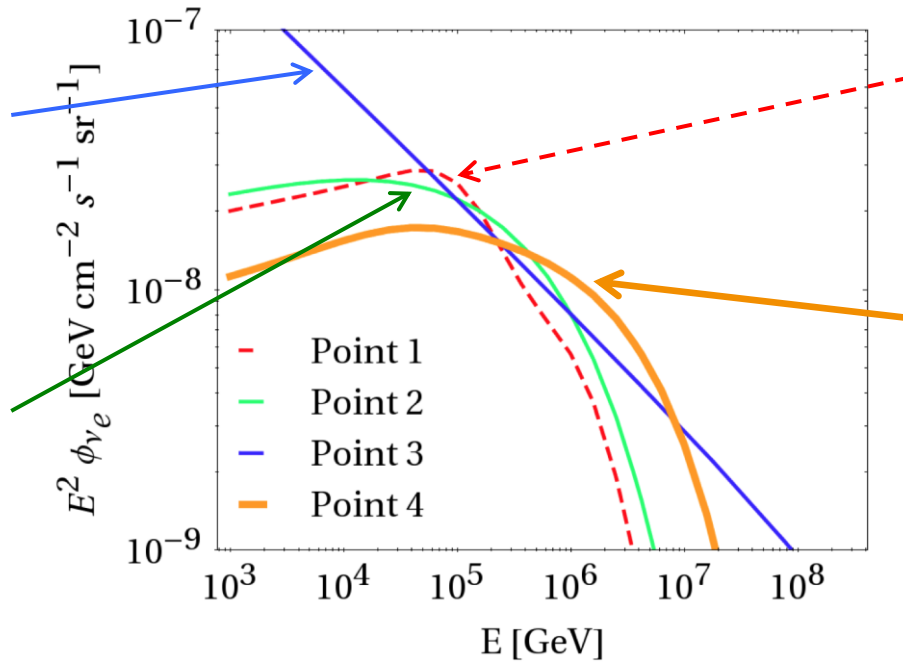
➤ Simplest model, as the neutrino spectrum follows that of the primaries

➤ Possible fits to data:

Parameter	Description	Unit
α	Spectral index of primary nuclei	none
E_{\max}	Maximal energy	GeV
B	Magnetic field	Gauss (G)
A	Mass number	none

Protons, $\alpha=2.5$
[Problem:
Fermi diffuse
 γ -ray bound
Murase, Ahlers,
Lacki, PRD 2013]

Protons
 $\alpha=2$
 $E_{\max}=10^{7.5}$ GeV



Protons
 $\alpha=2$
 $B \sim 10^4$ G
(magnetic field effects on
sec. pions, muons, kaons)

Nuclei
 $\alpha=2, E_{\max}=10^{10.1}$ GeV
Composition at source

$$A(E) = \max \left(1, 56 \times \left(\frac{E}{E_{\max}} \right)^\beta \right)$$

with $\beta=0.4$



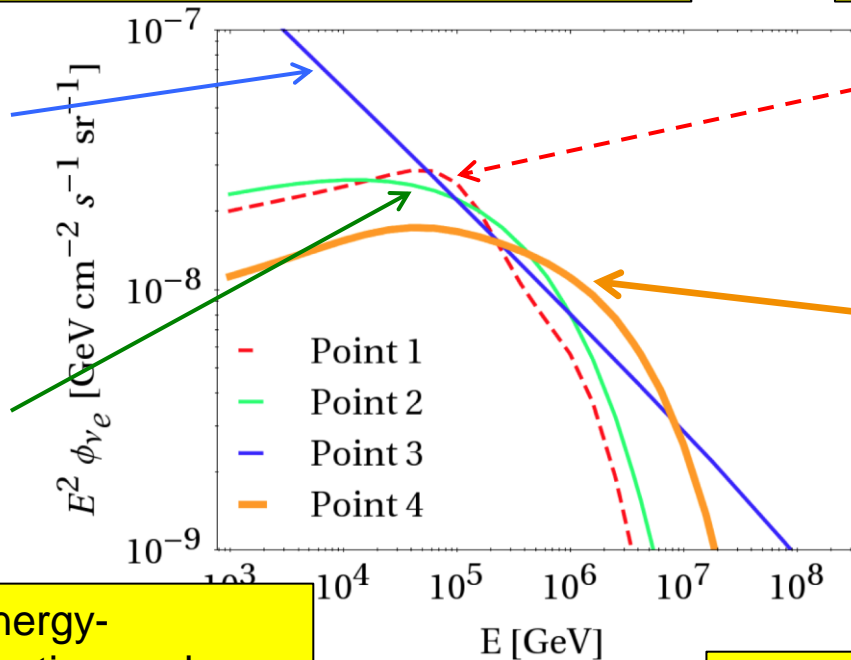
Connection to ultra-high energy cosmic rays (UHECR)?

Yes, but: Energy input per decade very different in neutrino-relevant and UHECR energy ranges (Energetics seem to favor $\alpha \sim 2$ – Waxman/Bahcall!)
 see e. g. Katz et al, 1311.0287 for generic discussion;
 for GRBs specifically: extremely large baryonic loadings implied: Baerwald et al, *Astropart. Phys.* 62 (2015) 66

Yes, but: Synchrotron losses limit maximal proton energies as well. Need large Doppler factors (e. g. GRBs)

Protons, $\alpha=2.5$
 [Problem: Fermi diffuse γ -ray bound
 Murase, Ahlers, Lacki, PRD 2013]

Protons
 $\alpha=2$
 $E_{\max}=10^{7.5}$ GeV



Protons
 $\alpha=2$
 $B \sim 10^4$ G

Nuclei
 $\alpha=2$, $E_{\max}=10^{10.1}$ GeV
 Composition at source

$$A(E) = \max \left(1, 56 \times \left(\frac{E}{E_{\max}} \right)^\beta \right)$$

with $\beta=0.4$

Yes, but: Need energy-dependent escape timescale leading to break/cutoff within source (diff. from ejection!)
 see e.g. Liu et al, PRD, 2014;
 arXiv:1310.1263 or starburst galaxies

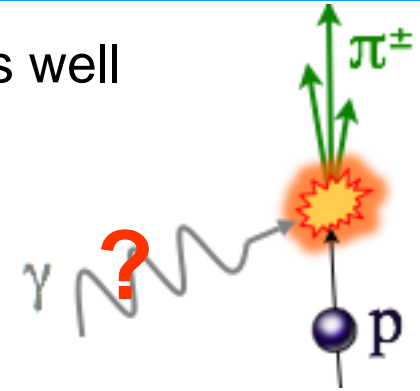
WW, arXiv:1407.7536
 (PRD, in print)

Yes, but: $A(E)$ change somewhat too shallow to match observation; difference source-observation from propagation?



Neutrinos from $p\gamma$ interactions

- > More freedom, as spectral shape depends on photons as well
 - Strategies to address the large parameter space?
- > Target photon field typically:
 - Put in by hand (derived from observed spectrum)
 - Thermal target photon field (e. g. accretion disk)
 - From interplay of radiation processes (synchrotron radiation, Bremsstrahlung, inverse Compton, pair production ...).
- > One of simplest self-consistent cases:
From synchrotron radiation of co-accelerated electrons (AGN-like model)
- > Requires few model parameters, mainly



Parameter	Units	Description	Typical values used
R	km (kilometers)	Size of acceleration region	10^1 km ... 10^{21} km
B	G (Gauss)	Magnetic field strength	10^{-9} G ... 10^{15} G
α	1	Universal injection index	1.5 ... 4

Model by Hümmer, Maltoni, Winter, Yaguna, *Astropart. Phys.* **34** (2010) 205



Parameter space: Hillas plot?

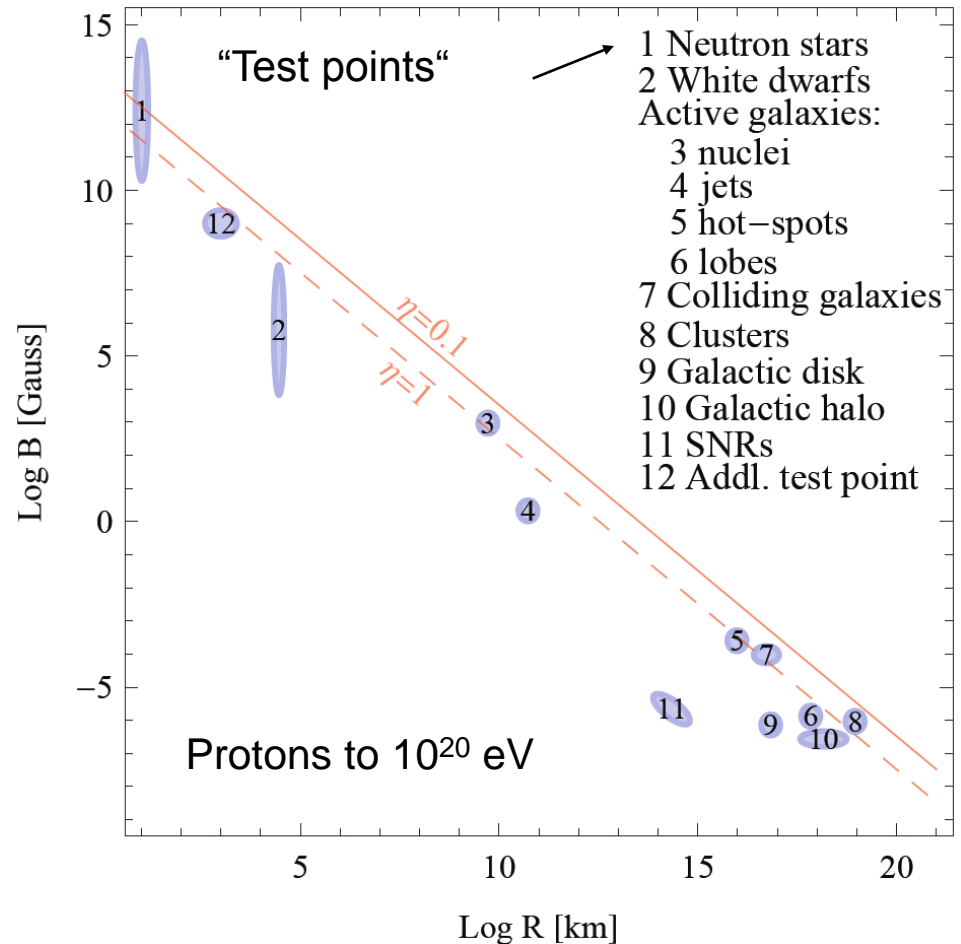
- > Model-independent (necessary) condition for acceleration of cosmic rays:

$$E_{\max} \sim \eta Z e B R$$

(Larmor-Radius < size of source;
 η : acceleration efficiency)
Particles confined to within
accelerator!

- > Caveat: condition relaxed if source heavily Lorentz-boosted (e.g. GRBs)
- > Test points in these figures not to be taken too seriously (large astrophysical uncertainties)

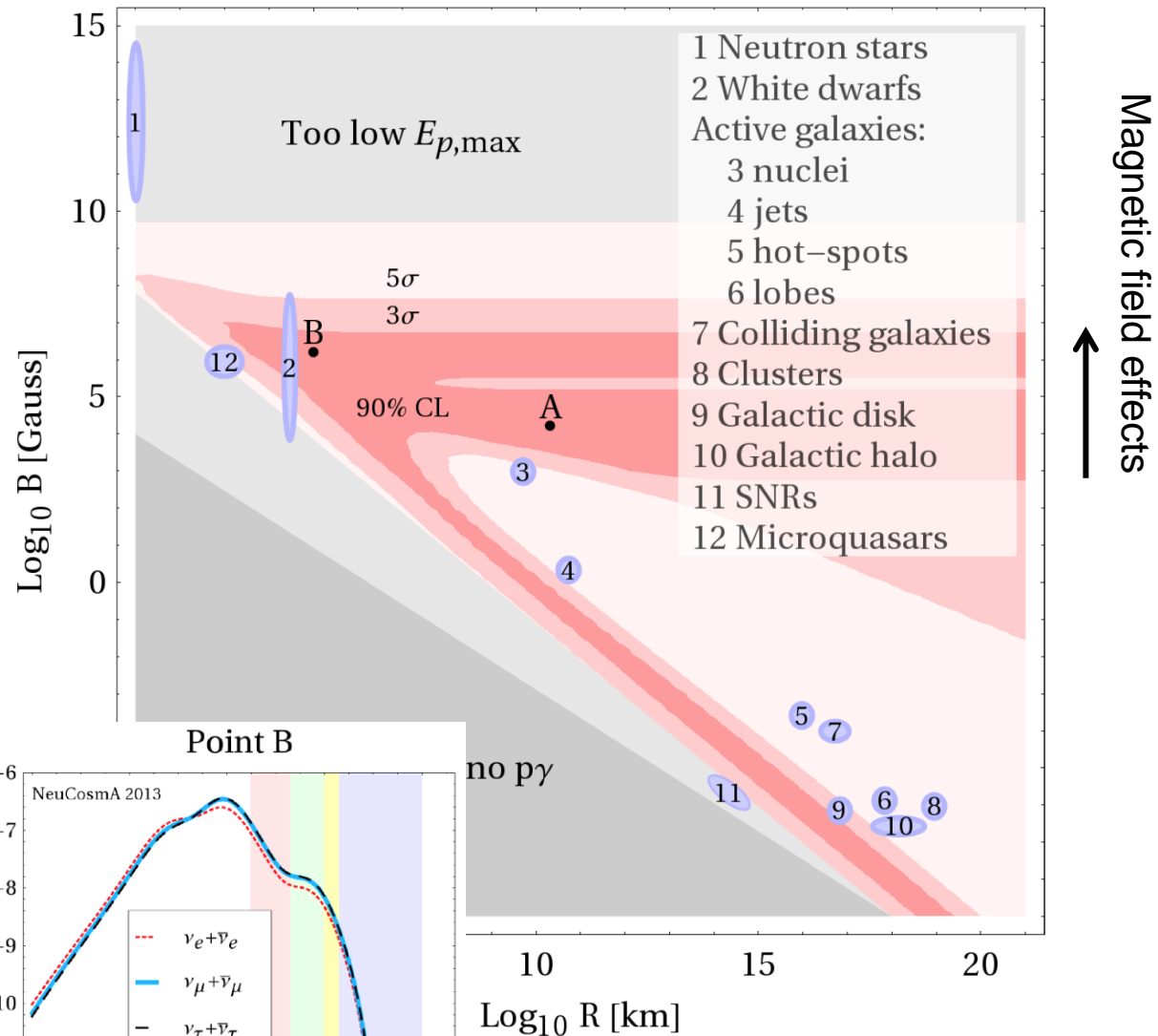
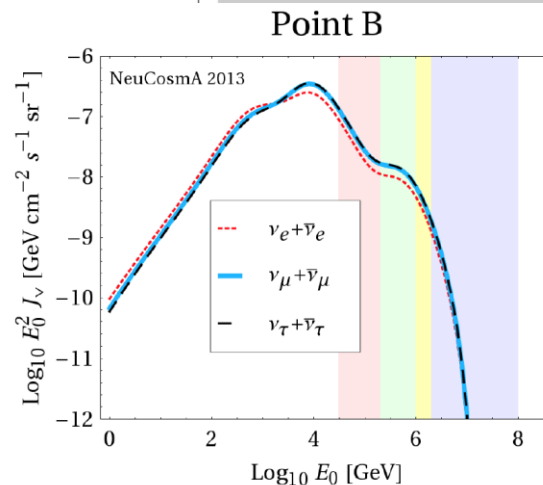
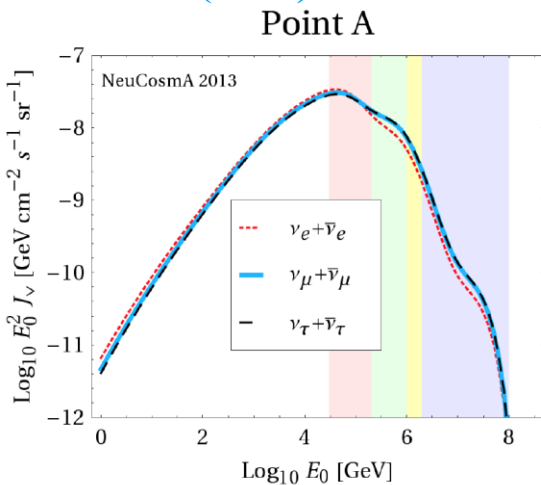
Hillas 1984; version adopted from M. Boratav



Parameter space constraints (SFR evolution, $p\gamma$ model)

- Current data start to constrain parameter space of radiation models
- Again, regions preferred where either maximal proton energy or magnetic field effects lead to cutoff

WW, arXiv:1307.2793,
PRD88 (2013) 083007



So, where do the neutrinos come from? (my personal bias)

- > We have no clue. There are, however, plenty of candidate classes
 - Galactic sources
 - + testable by directional correlations (some workarounds in literature ...)
 - no evidence, so far; perhaps only a few events
 - Gamma-ray bursts
 - + may have enough power, cutoff “natural” from magnetic field effects
 - needs to be a “non-standard” population, as strong bounds on gamma-ray detected GRBs from stacking; low luminosity GRBs?
 - Active Galactic Nuclei
 - + wide playground, may have enough power; there is hope for some directional/flaring correlations
 - wide playground; not easy to accommodate a cutoff (unless from primaries)
 - Starburst galaxies, hypernova remnants, ... (diffuse flux contributions)
 - + good to blame for if not better solution found; cutoff/break expected if energy-dependent escape timescale
 - direct evidence difficult (not easily resolvable)
 - From interactions with photon backgrounds (cosmogenic or nearby populations)
 - + kind of expected fluxes (at some level)
 - models which produce IceCube flux require tweaking

(see [Anchordoqui et al, arXiv:1312.6587](#) for a review)
- > There are some recent proposals for generic tests, e. g. limit source density from clustering of events
(see e. g. [Ahlers, Halzen, arXiv:1406.2160](#))



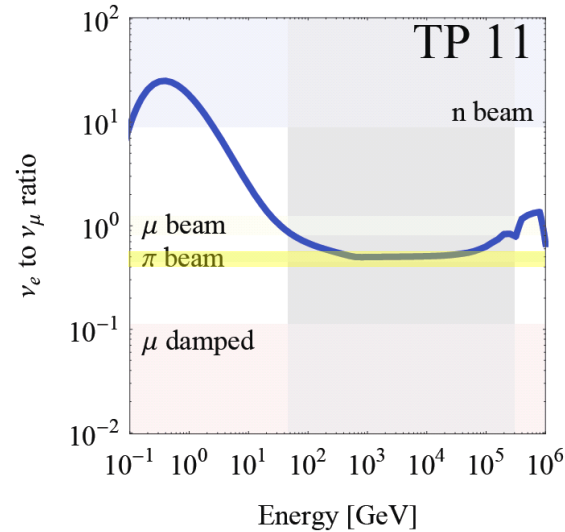
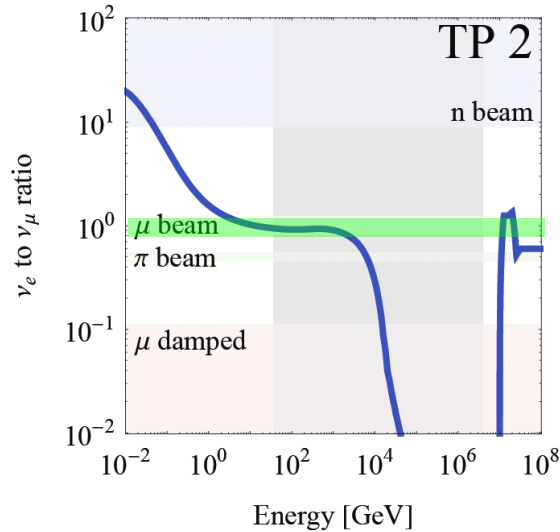
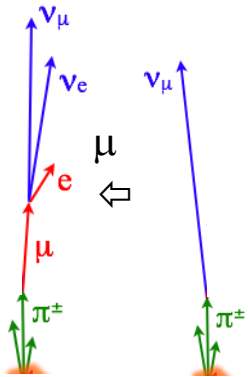
Future challenges



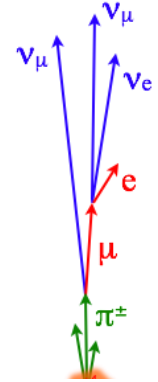
The question of flavor:

Flavor composition is energy-dependent! (earlier $p\gamma$ model)

Muon beam
- muon damped

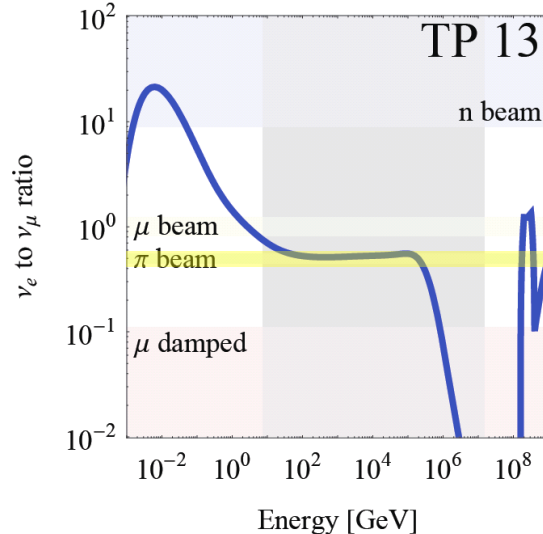
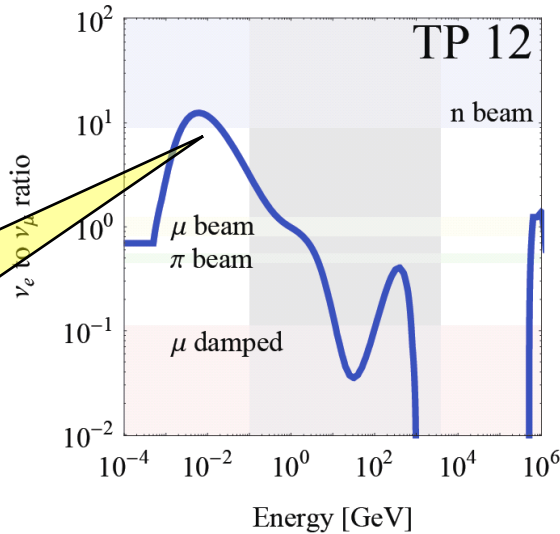


Pion beam
 $(\nu_e:\nu_\mu:\nu_\tau)=(1:2:0)$

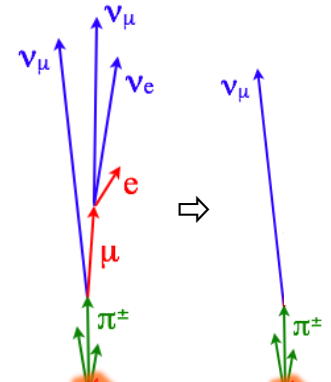


Undefined
(mixed source)

Typically
n beam
for low E
(from $p\gamma$)



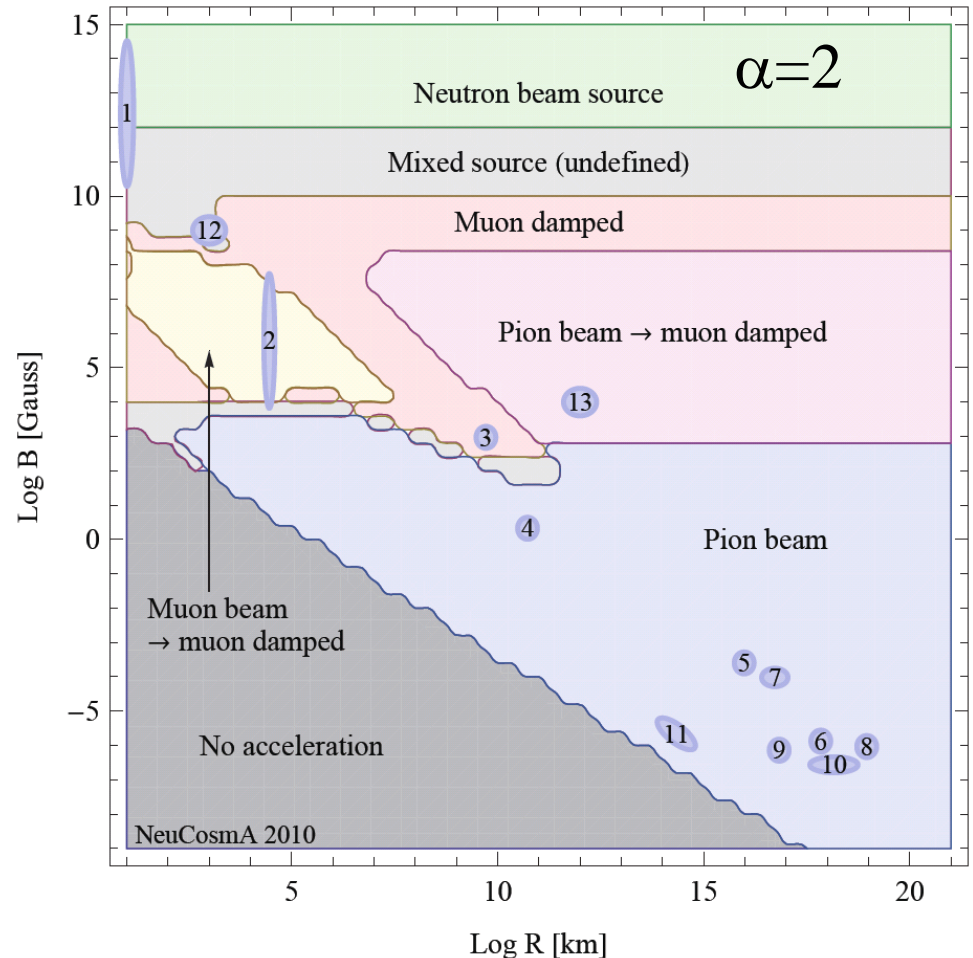
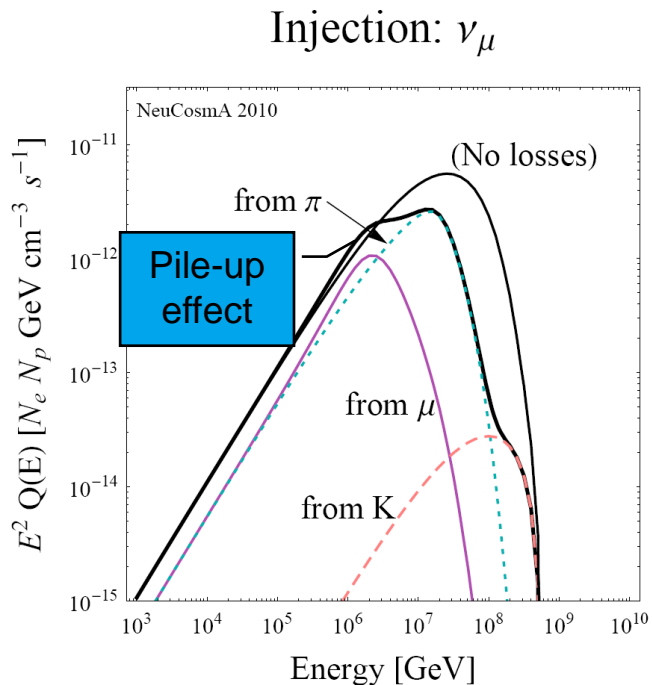
Pion beam
⇒ muon damped



Flavor composition ... can be predicted

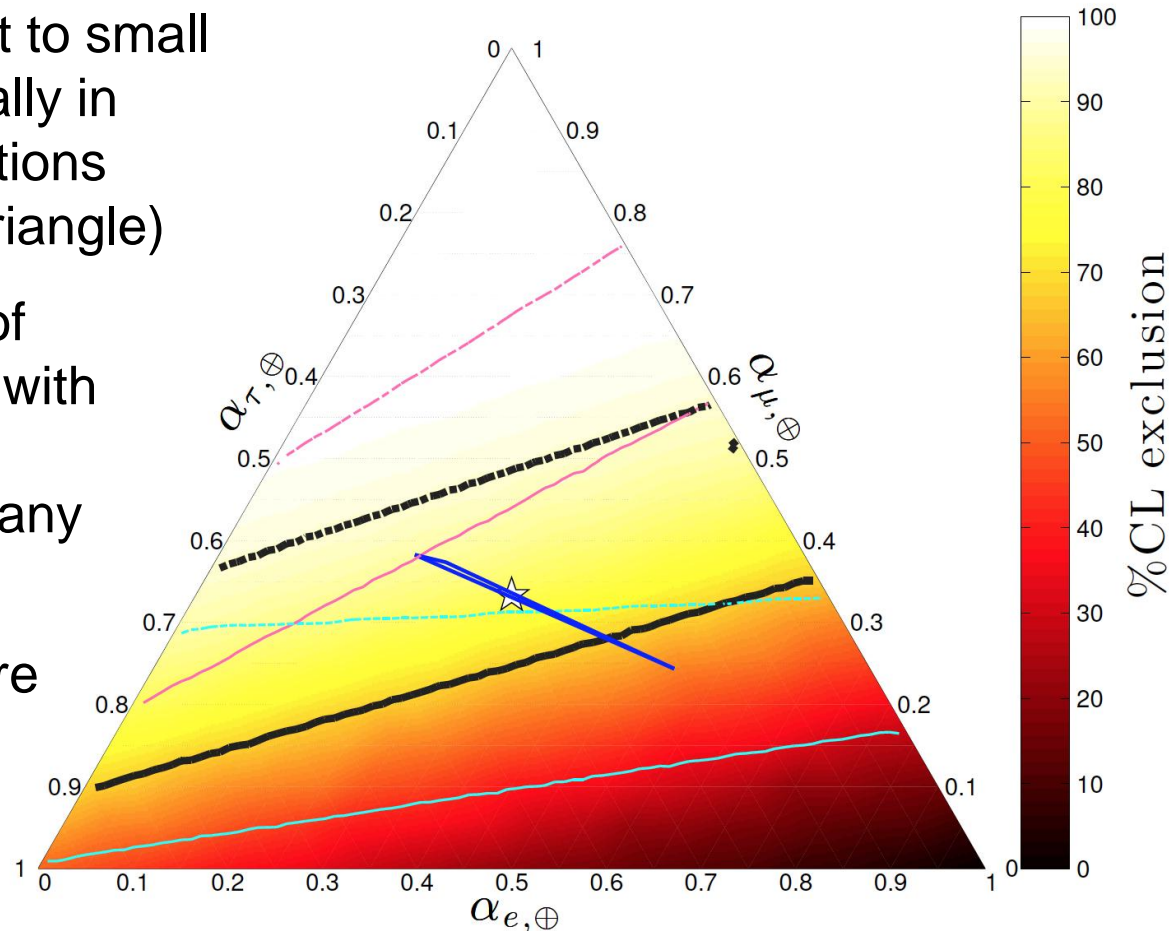
- Pion beam good assumption for sources on galactic scales
- Muon beam sources if muon pile-up:

Hümmer, Maltoni, Winter, Yaguna,
Astropart. Phys. 34 (2010) 205



Flavor composition ... can be measured!

- Statistical significance yet too small to be meaningful, especially in region of standard predictions after flavor mixing (blue triangle)
- There is a slight tension of astrophysical predictions with data, as the background prediction contains too many muon tracks
- Future challenge: measure flavor composition with precision meaningful in blue triangle!



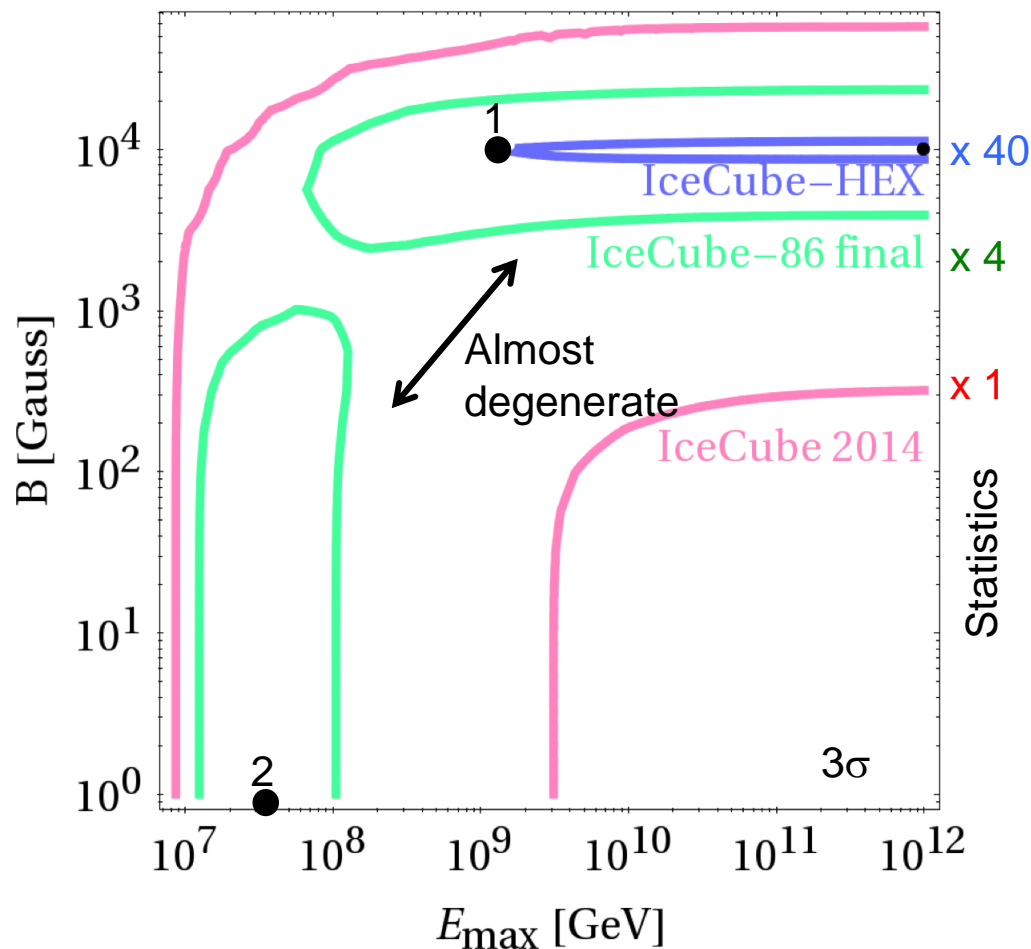
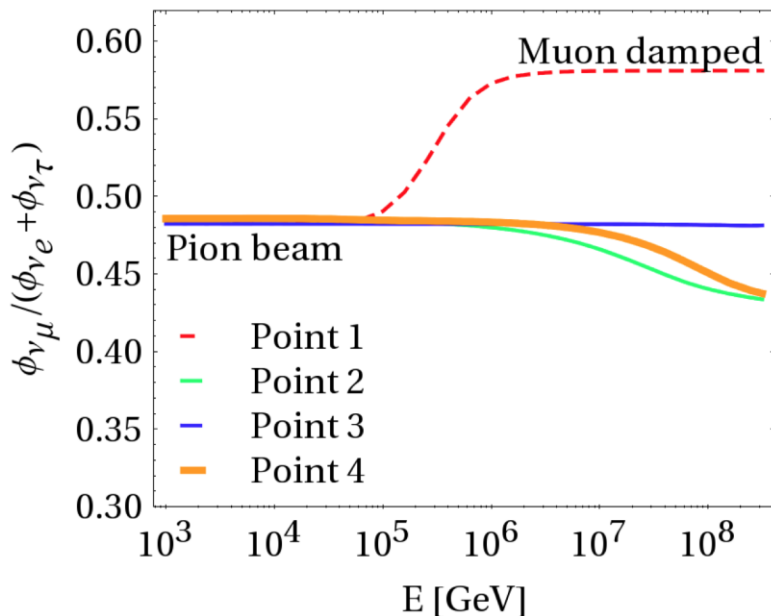
Discussion by Sergio Palomares-Ruiz on Oct. 2, 2014 at KITP; watch podcast at <http://online.kitp.ucsb.edu/online/neutrinos14/palomaresruiz/>

Mena, Palomares-Ruiz, Vincent,
PRL 113 (2014) 091103; see also Fu, Ho,
Weiler, 1411.1174 for a new related work



Physics case for high-E extension? Flavor! (Ap model)

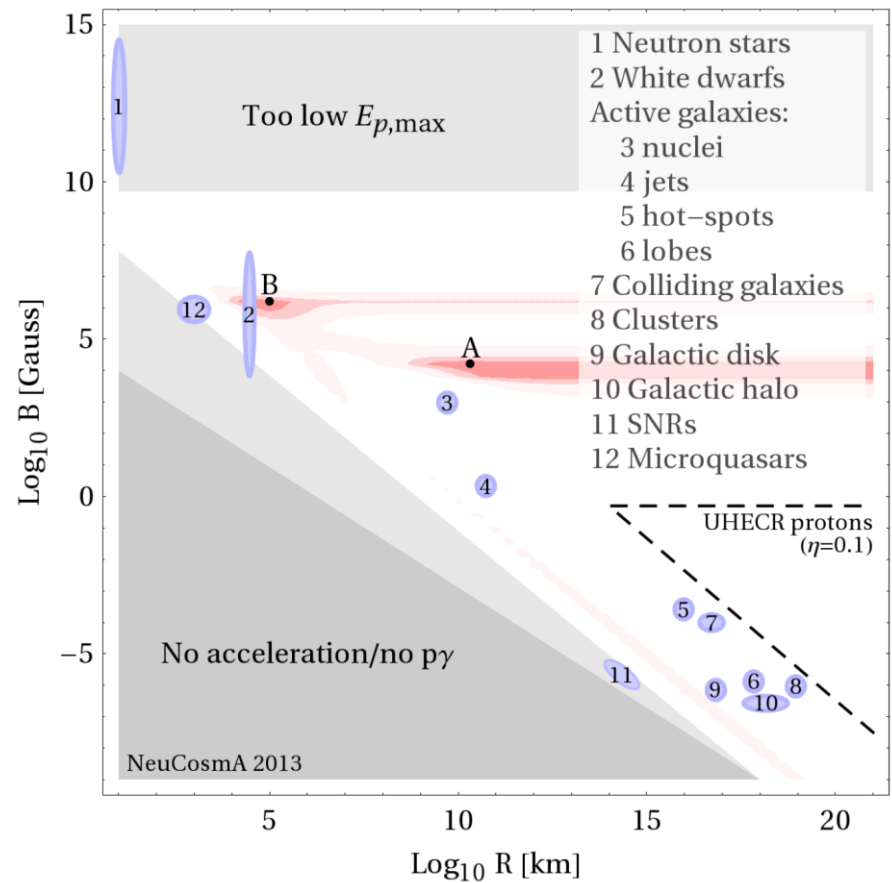
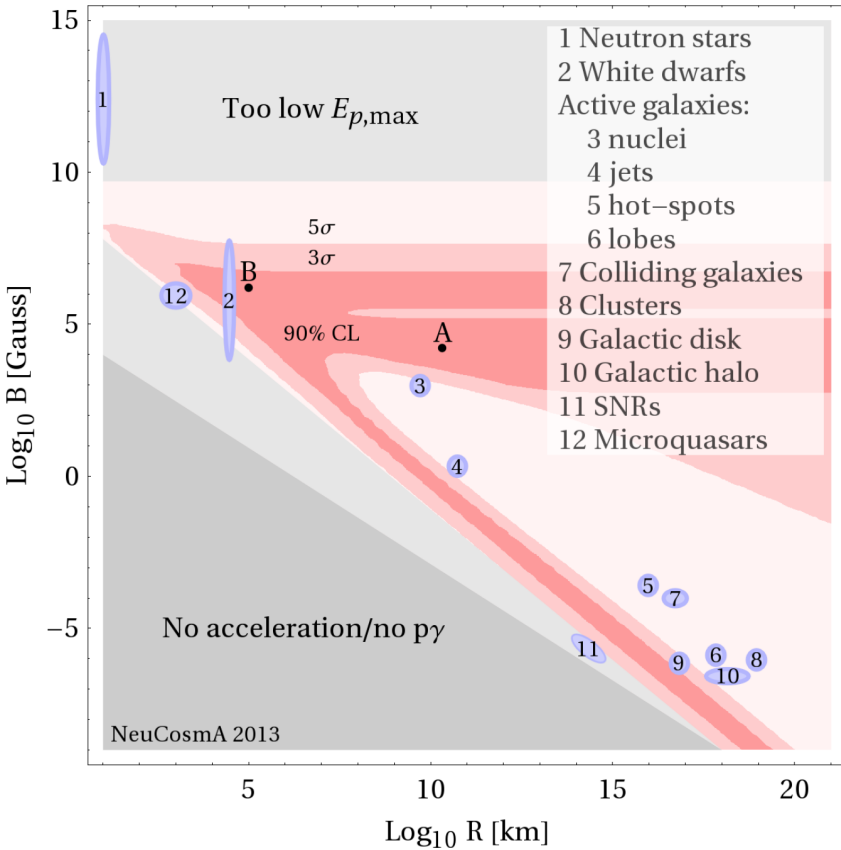
- Cutoff from magnetic field effects degenerate with maximal energy cutoff in terms of spectral shape
- Flavor composition different (see below)
- Need sufficient statistics in cutoff region (about 1 PeV)



WW, arXiv:1407.7536, PRD, in print



Physics case for high-E extension? Precision! ($p\gamma$ model)



WW, arXiv:1307.2793, PRD88 (2013) 083007



Summary and conclusions

- > Cosmic neutrinos are interesting, as they point towards the sources of the cosmic rays
- > Next major qualitative step: resolve the sources of the neutrinos; perhaps neutrino astronomy can do better than cosmic ray observations?
- > If not, we have to rely on conceptual arguments, which require the statistics of an high-energy extension:
 - Spectral shape
 - Flavor composition
 - Anisotropies
 - Multiplets, ...
- > Several constraints come from gamma-ray and cosmic ray observations; *theory* needed to draw a self-consistent multi-messenger picture
- > Connection to UHECRs requires extrapolation over several orders of magnitude in energy, and composition; **source modeling in presence of heavier nuclei one of future key issues for CR-neutrino theory**

