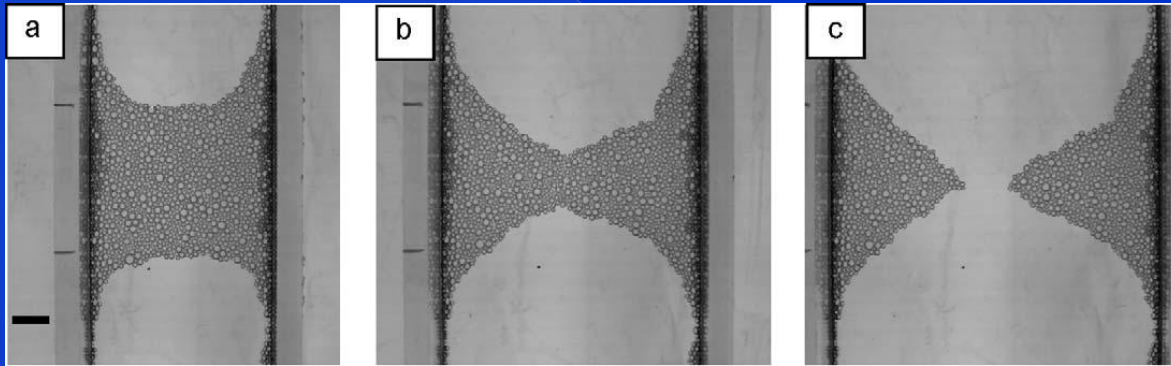
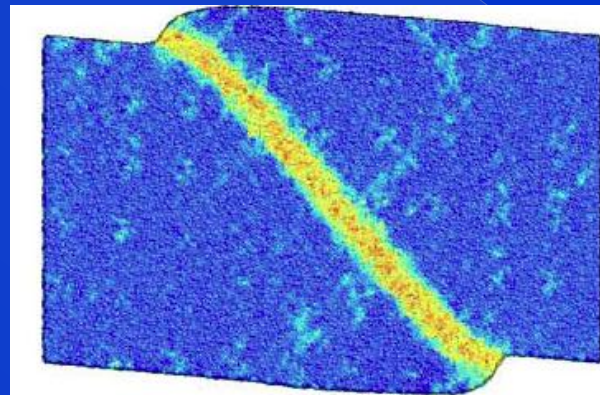


Continuum physics of deforming glasses: From oscillatory shear to fracture

Eran Bouchbinder
Weizmann Institute of Science



Arciniaga, Kuo and Dennin
Colloids and Surfaces A (2011)



Cao, Cheng and Ma
Acta Materialia (2009)



Lowhaphandu and Lewandowski
Scripta Materialia (1998)

Work with:

Jim Langer (UCSB)

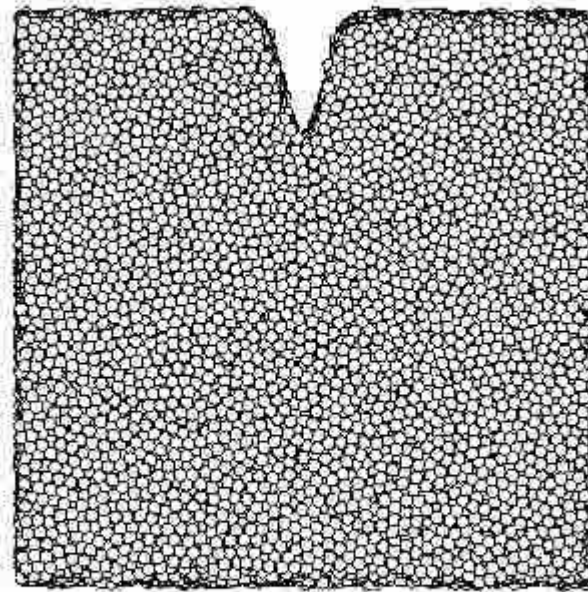
Chris Rycroft (Harvard)

Nathan Perchikov (Weizmann)

The challenge (as we see it)

To develop a
focus on no
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cking instabilities,

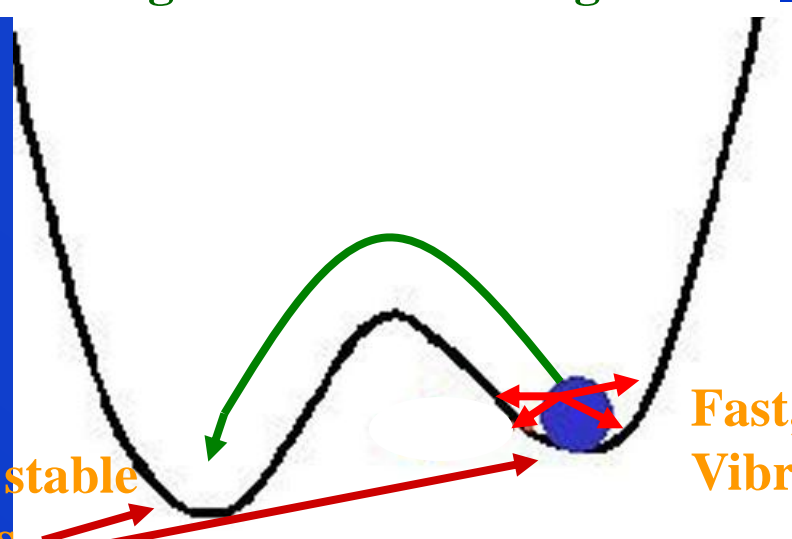


N. Bailey et al. PRB 69, 144205 (2004)
Simulation of Cu-Mg Metallic Glass

Theoretical approach: Nonequilibrium thermodynamics

Basic idea 1: Separable Configurational + Kinetic/Vibrational Subsystems

Slow, Non-Equilibrated,
Configurational rearrangements



Total internal energy:

$$U_{total} \approx U_C + U_K$$

Total entropy:

$$S_{total} \approx S_C + S_K$$

Mechanically stable
configurations

Fast, Equilibrated,
Vibrational motion

Weak coupling between these two subsystems, Timescales separation, Quasi-ergodicity due to external driving forces

Basic idea 2: The degrees of freedom relevant to irreversible deformation can be described by non-equilibrium coarse-grained internal variables

$$U_{total} \approx U_C(S_C, E_{ij}, \Lambda, m_{ij}) + U_K(S_K, E_{ij})$$

Define two different temperatures:

$$\chi = \left(\frac{\partial U_C}{\partial S_C} \right) \quad T = \left(\frac{\partial U_K}{\partial S_K} \right)$$

“Effective” temperature, non-equilibrium degrees of freedom

Ordinary, equilibrium temperature

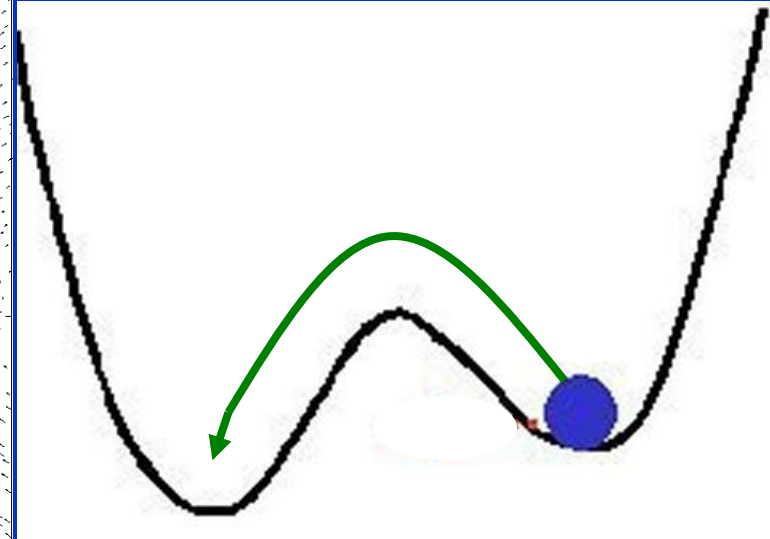
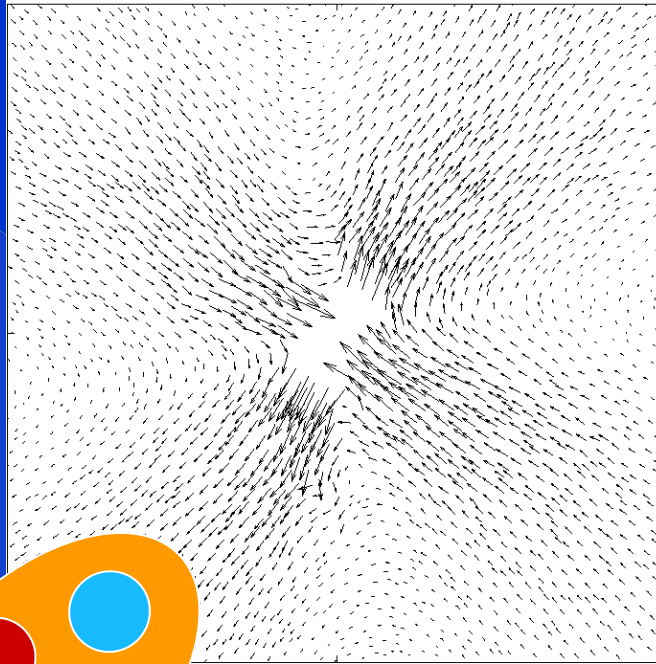
Early ideas in the glass/granular materials community:

Edwards, Cugliandolo, Kurchan, Coniglio, Barrat, Berthier and more

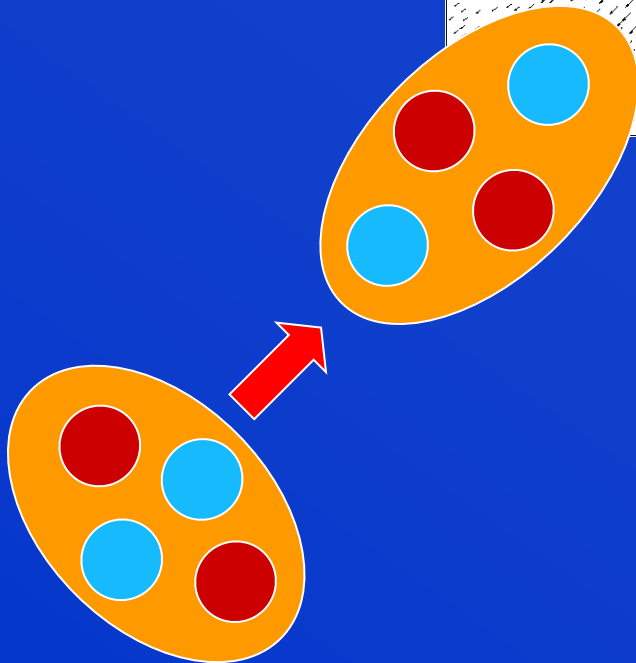
χ

is a thermodynamic temperature, e.g. it appears in equations of state, it controls the probability of configurational fluctuations etc.

Theoretical approach: Constitutive law



Potential Energy Landscape



Shear-Transformation-Zones (STZ)

$$\Lambda$$

Probability to observe an STZ
(average normalized density)

$$m_{ij}$$

Normalized orientational bias of STZ

ML Falk & JS Langer, Physical Review E 57, 7192 (1998)

EB, JS Langer & I Procaccia, Physical Review E 75, 036107 (2007)

EB & JS Langer, Physical Review E 80, 031133 (2009)

ML Falk & JS Langer, Annu. Rev. Condens. Matter Phys. 2, 353 (2011)

K Kamrin & EB, J. Mech. Phys. Solids, In press (2014)

The Equations (dimensionless, simple shear, low T)

$$\dot{\gamma}^{pl}(s, T, \Lambda, m) = \Lambda e^{-1/T} \cosh\left(\frac{\Omega s}{T}\right) \left[\tanh\left(\frac{\Omega s}{T}\right) - m \right]$$

STZ density thermal activation dynamics

STZ density change for v

$$\Lambda = e^{-1/\chi}$$

$$\dot{m} = \frac{2s\dot{\gamma}^{pl}}{\Lambda} (1 - ms)$$

Generalized Boltzmann factor controls STZ density

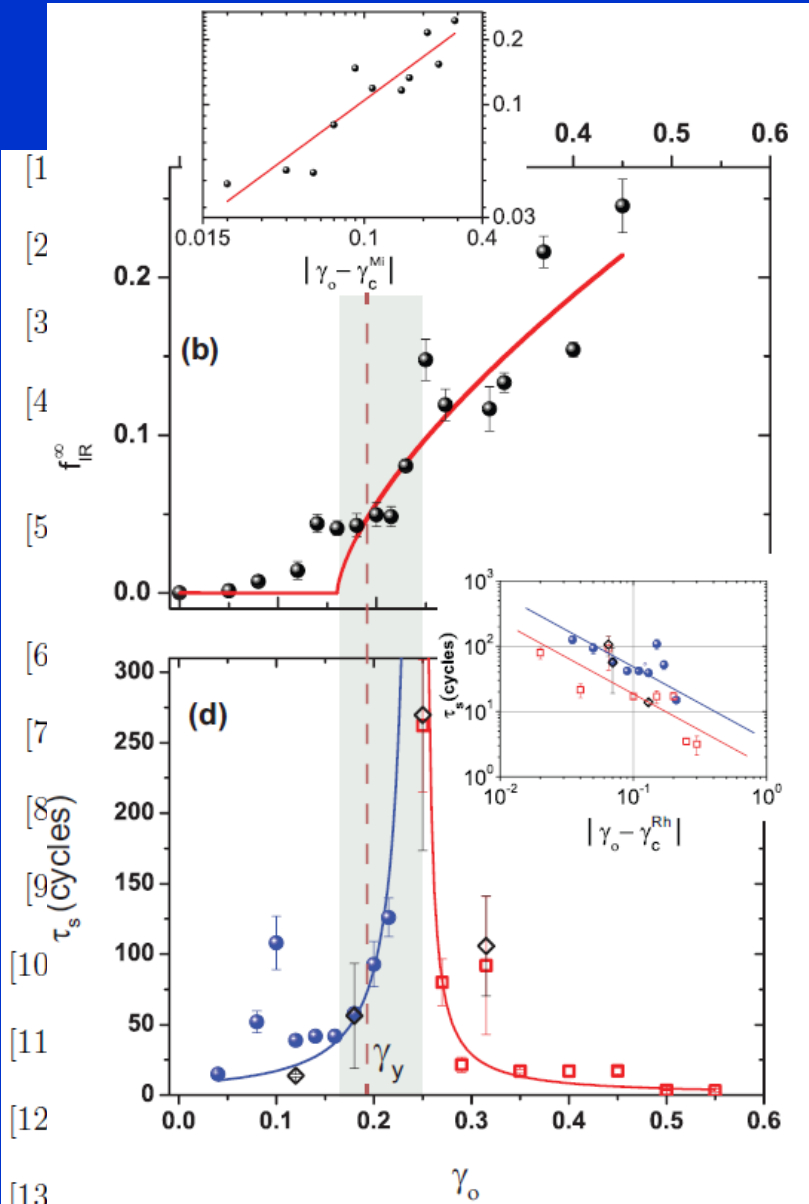
Orientational order dynamics

$$c_0 \dot{\chi} = 2s\dot{\gamma}^{pl} (\chi_\infty - \chi) + \kappa_c \nabla^2 \chi$$

Elasto-plastic rates decomposition

$$\dot{\gamma} = \dot{\gamma}^{el} + \dot{\gamma}^{pl} \quad \implies \quad \dot{s} = 2\mu [\dot{\gamma} - \dot{\gamma}^{pl}]$$

Application I: Variable-amplitude oscillatory shear



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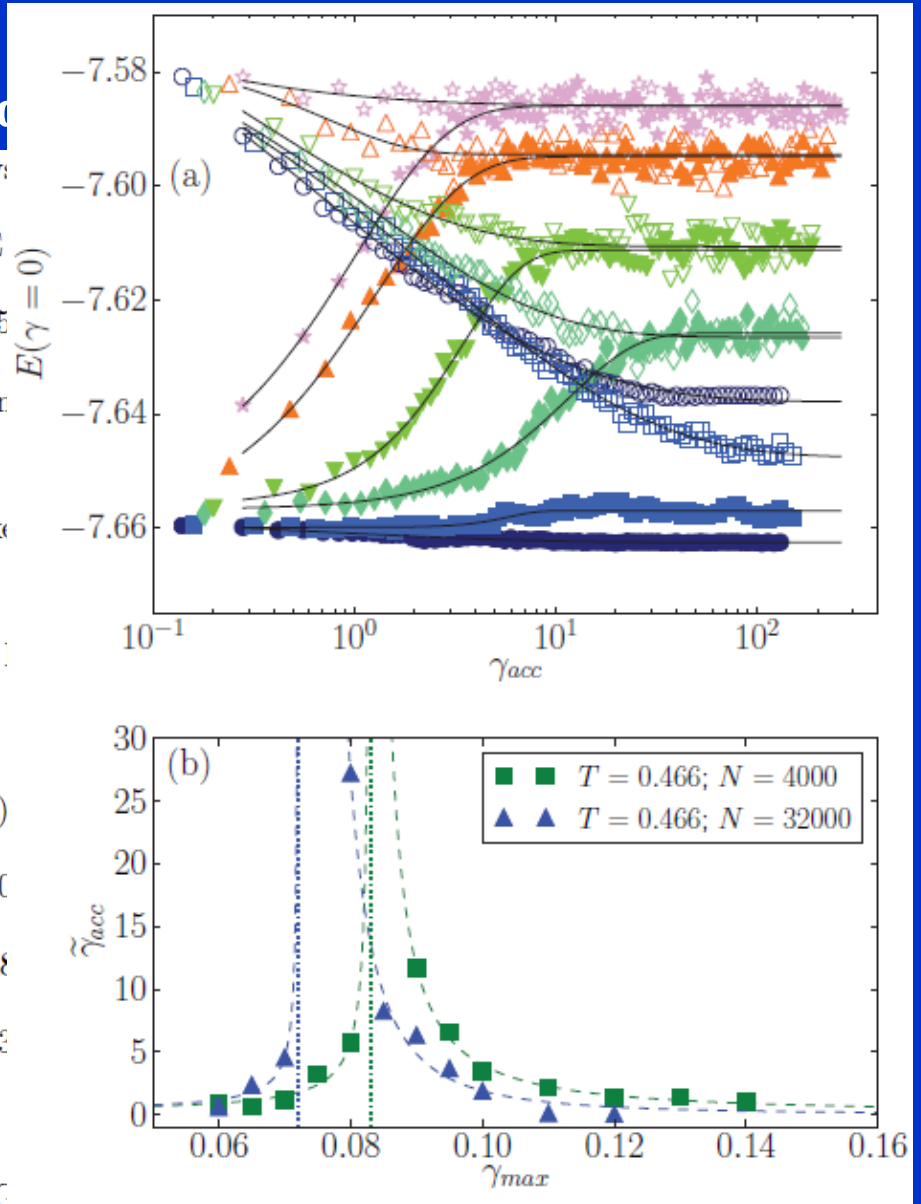
(2013)

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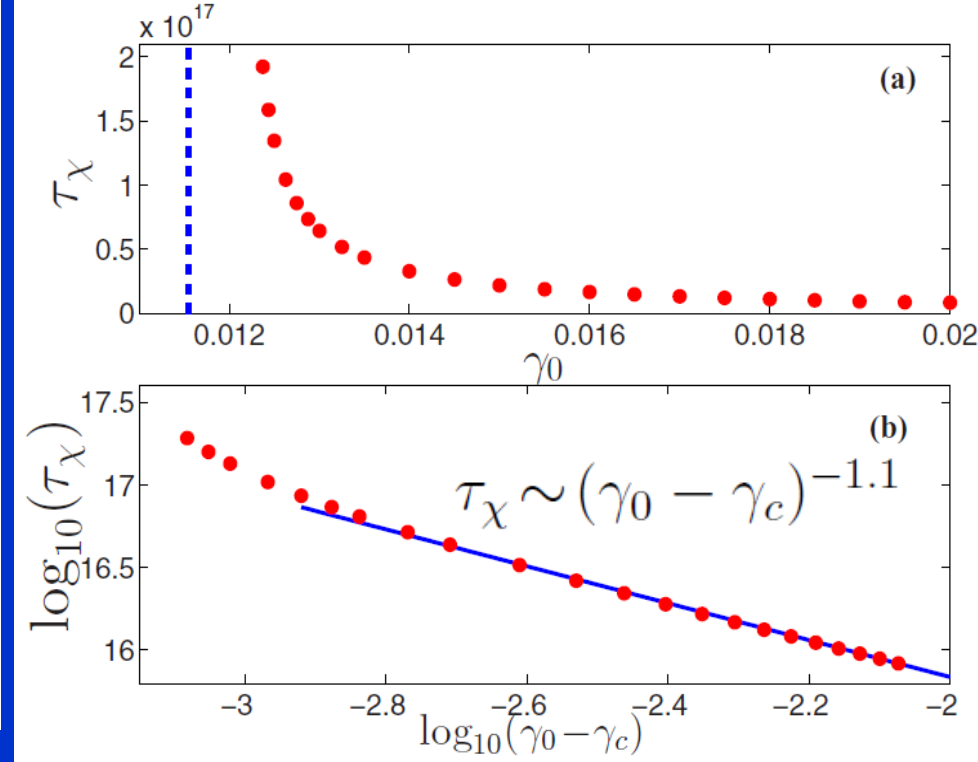
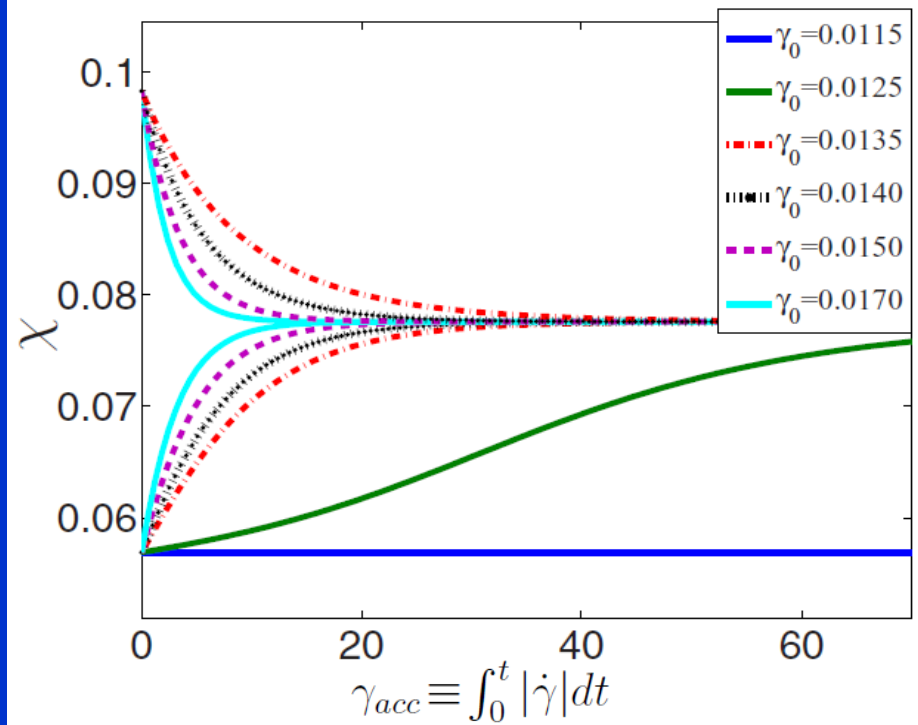
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Variable-amplitude oscillatory shear: Results



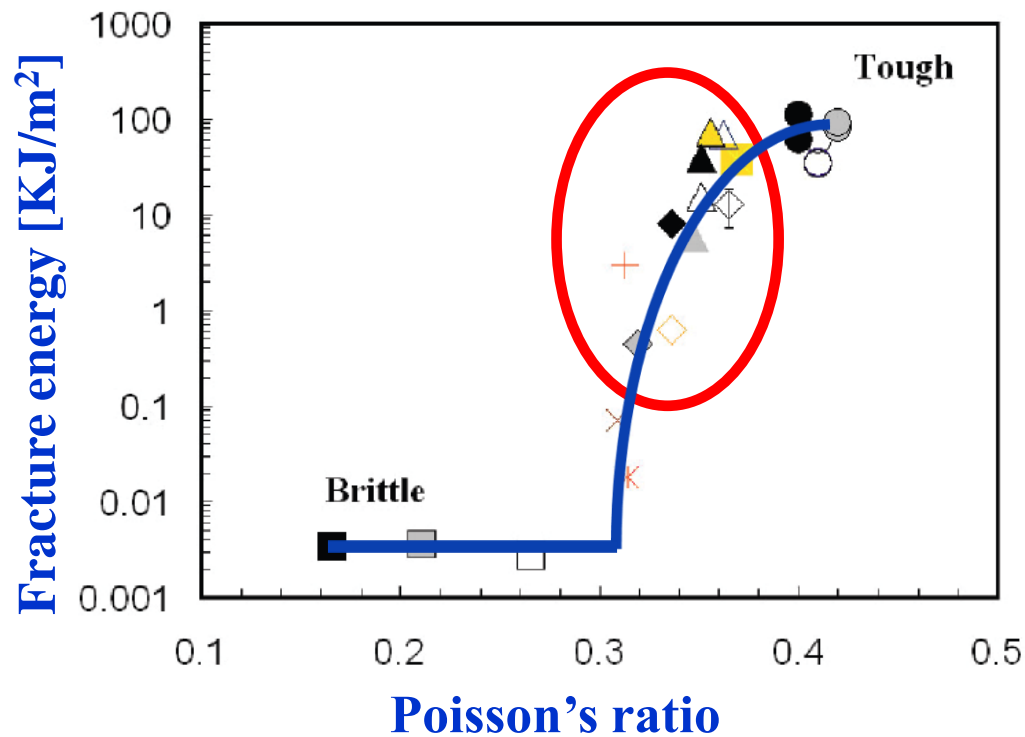
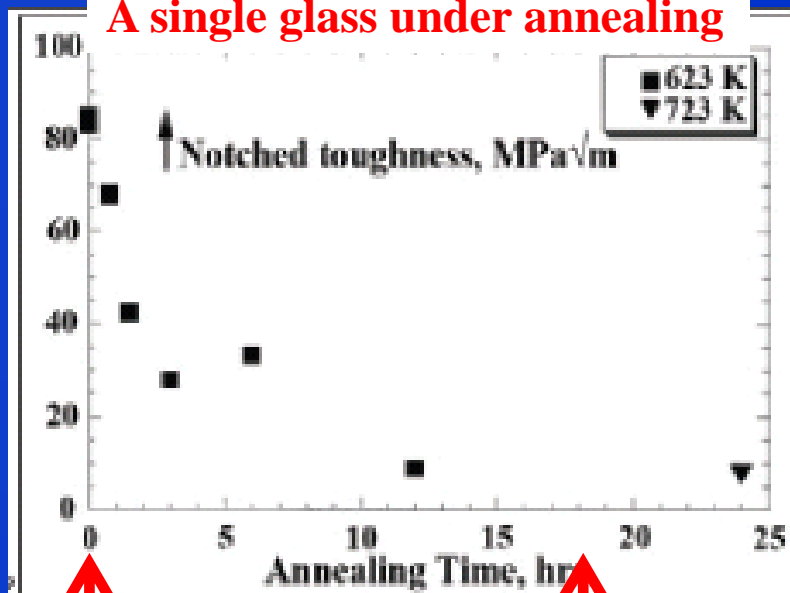
$$c_0 \dot{\chi} = 2s \dot{\gamma}^{pl} (\chi_\infty - \chi)$$

$$A \equiv \lim_{t \rightarrow \infty} \int_0^t s \dot{\gamma} dt$$

$$c_0 \dot{\chi} \simeq 2 \dot{\gamma}^{pl} = e^{-1/\chi} e^{-1/T} \cosh\left(\frac{\Omega s}{T}\right) \left[\tanh\left(\frac{\Omega s}{T}\right) - m \right]^{-1}$$

Application II: Ductile-to-brittle transition

A single glass under annealing



High initial χ

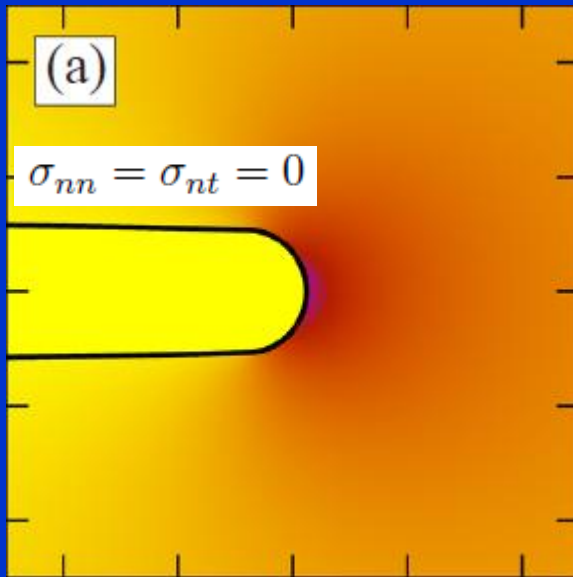
χ

Low initial χ

χ

$$\chi > T$$

Notch Fracture Toughness



Serious computational challenges were involved, especially in achieving realistic loading rates

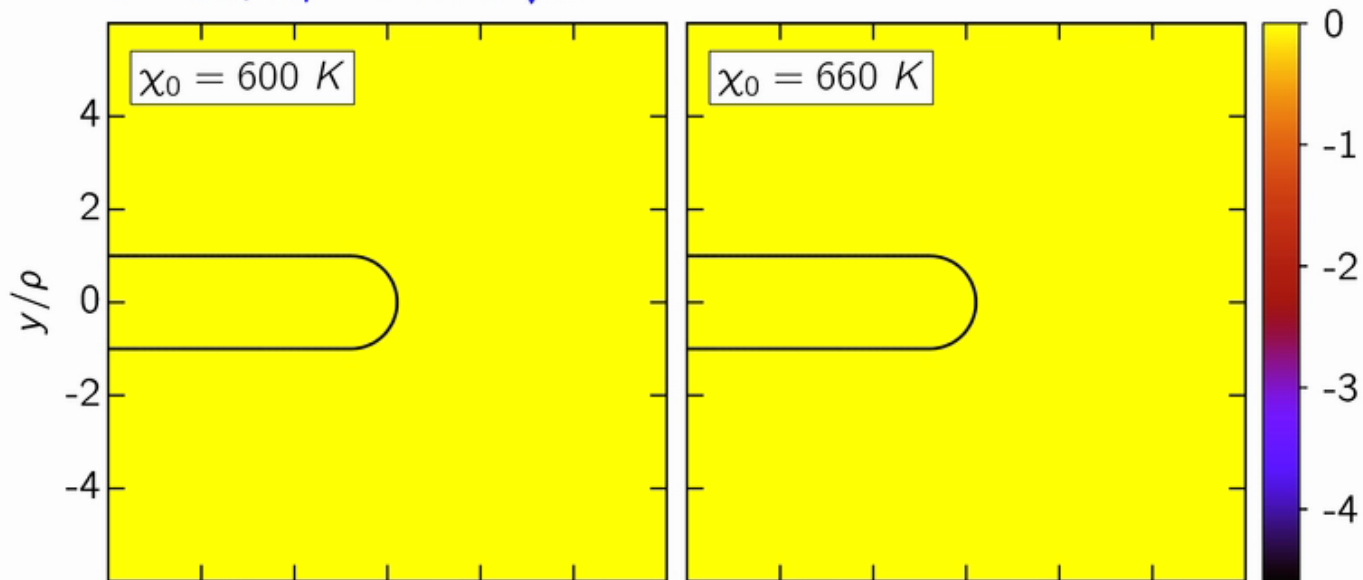
An Eulerian projection method for quasi-static elastoplasticity
 CH Rycroft, Y. Sui & EB, arxiv 1409.2173 (2014)

$Zr_{41.2}Ti_{13.8}Cu_{12.5}Ni_{10}Be_{22.5}$ bulk metallic glass (Vitreloy 1)

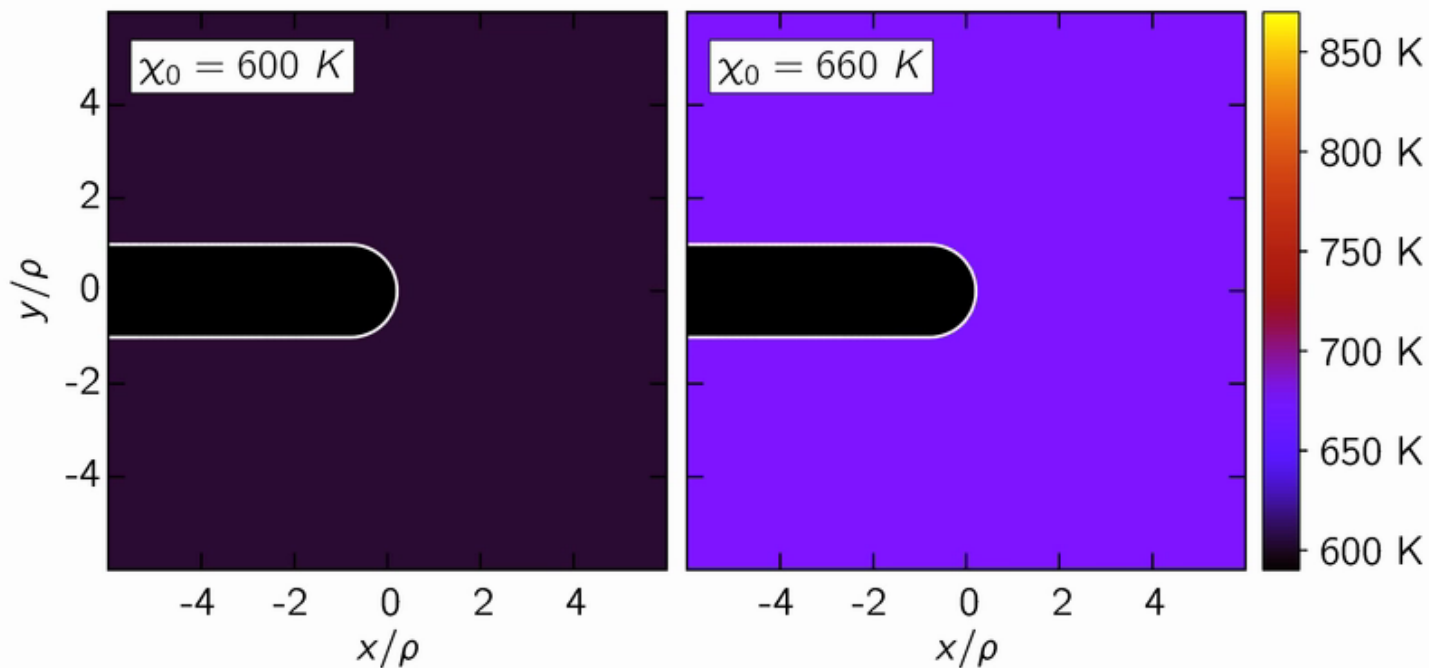
Parameters determined from independent sources	
$\tau_0 \sim 10^{-13}$ sec	Molecular vibration timescale
$s_y \sim 1$ GPa	Shear yield stress
$K, \mu \sim 100$ GPa	Elastic moduli
$e_z \sim 1$ eV	STZ formation energy
$\chi_\infty \sim 1.4T_g$	Steady state of χ
$T \sim 0.5T_g$	Test temperature

Results

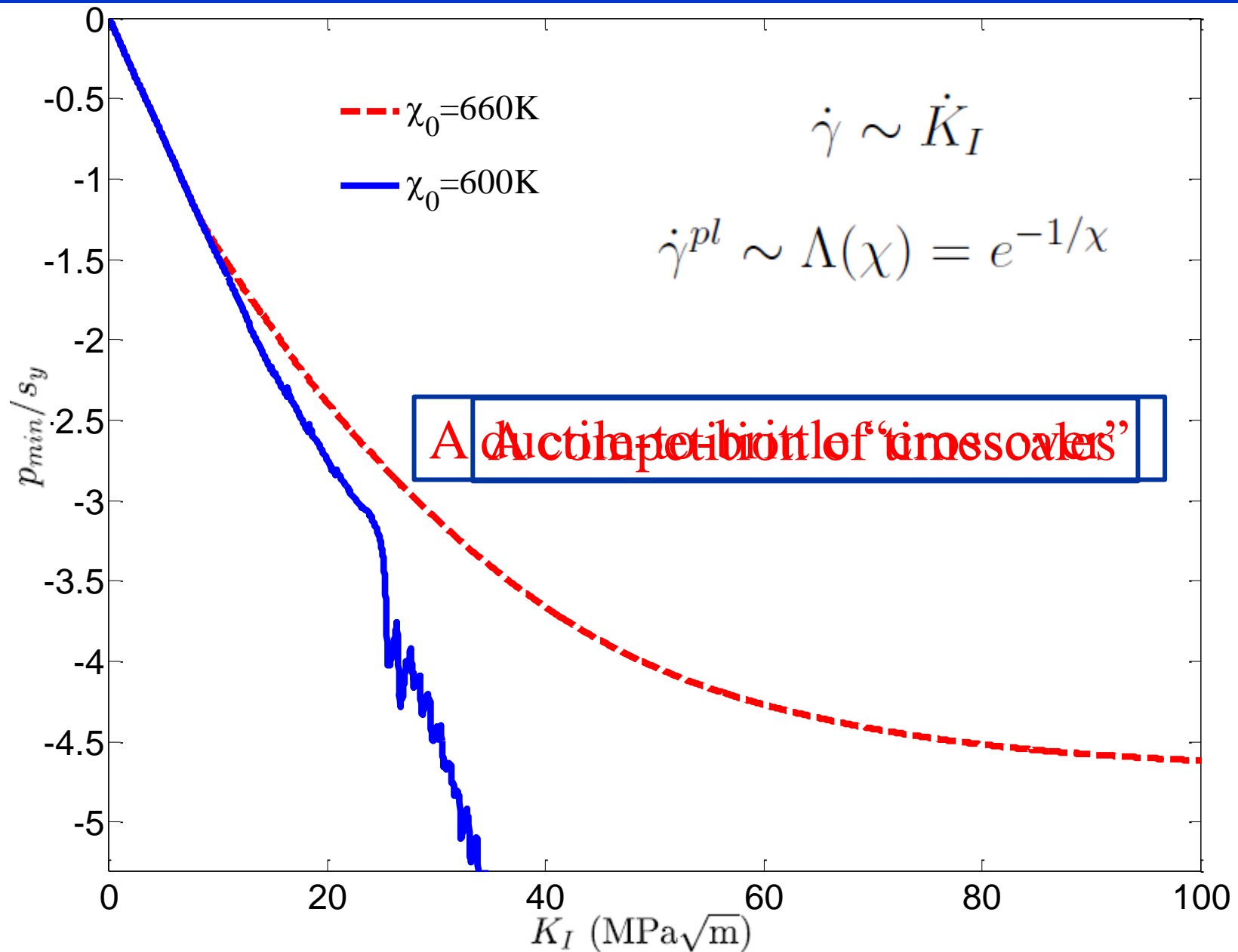
$t = 0 \text{ s}, K_I = 0 \text{ MPa } \sqrt{\text{m}}$



$t = 0 \text{ s}, K_I = 0 \text{ MPa } \sqrt{\text{m}}$

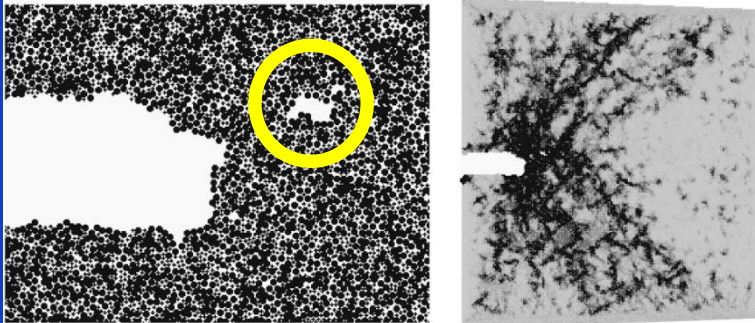


Results (cont'd)



Initiation criterion: A cavitation instability

MD simulation

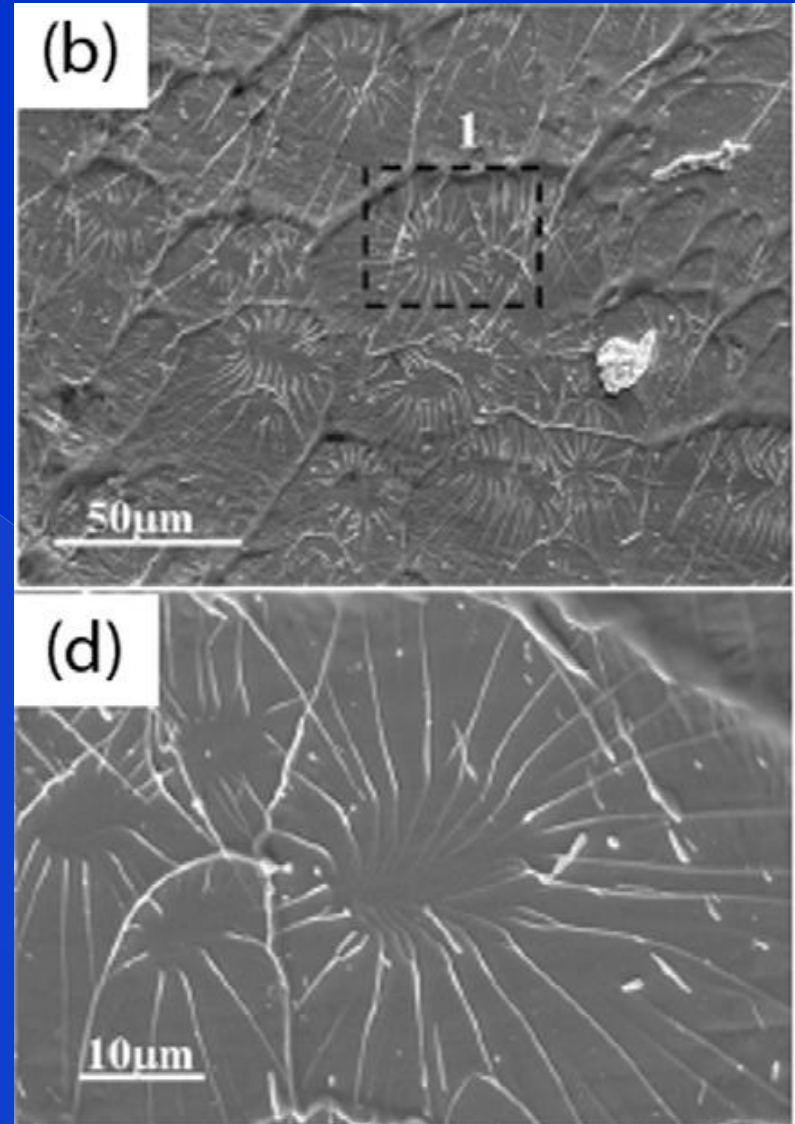


M. L. Falk, PRB 60, 7062 (1999)

$$\sigma_c = \frac{2}{\sqrt{3}} \left[1 + \log \left(\frac{2E}{3\sqrt{3}s_y} \right) \right] s_y$$

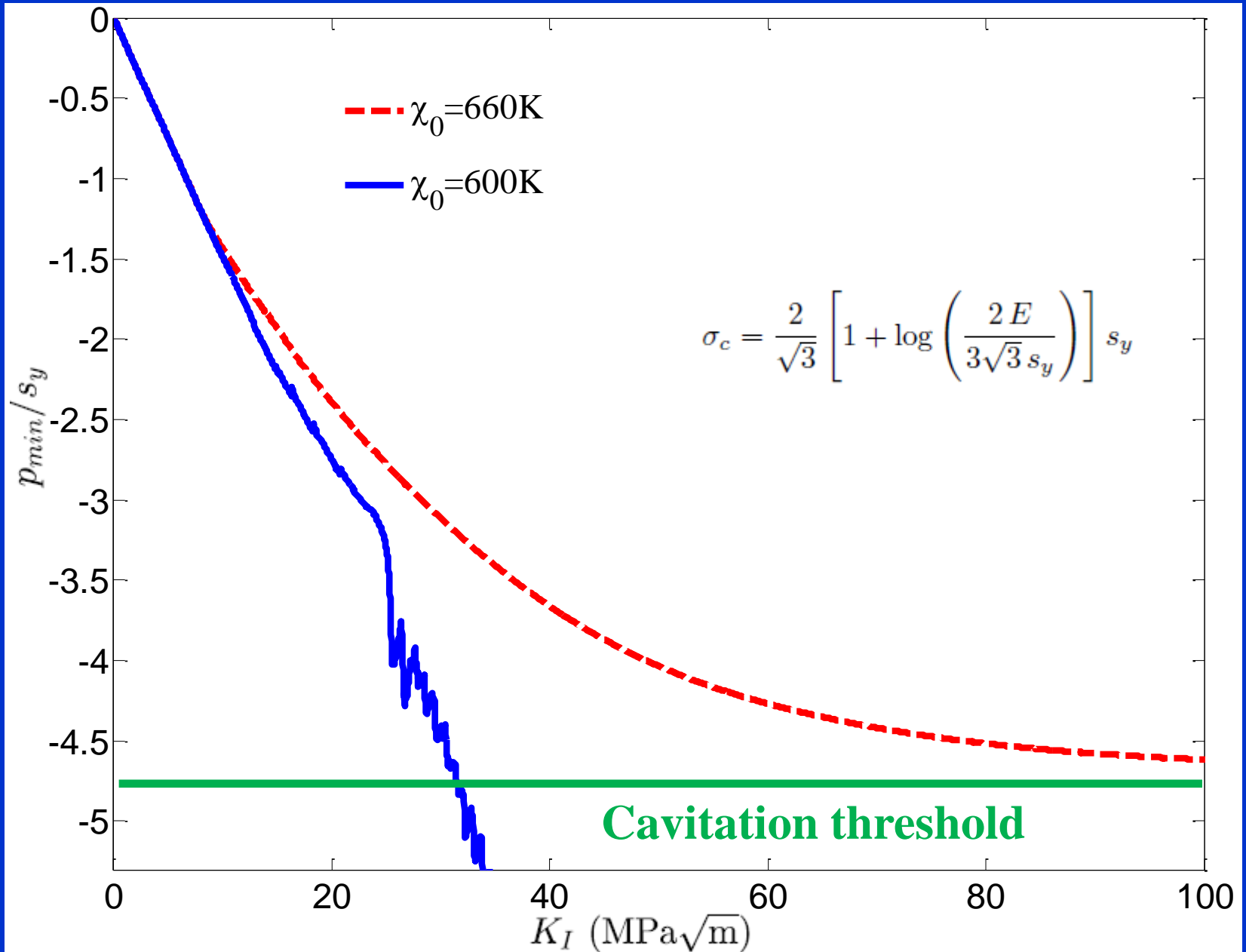
Huang et al., JMPS 39, 223 (1991)

Experiment



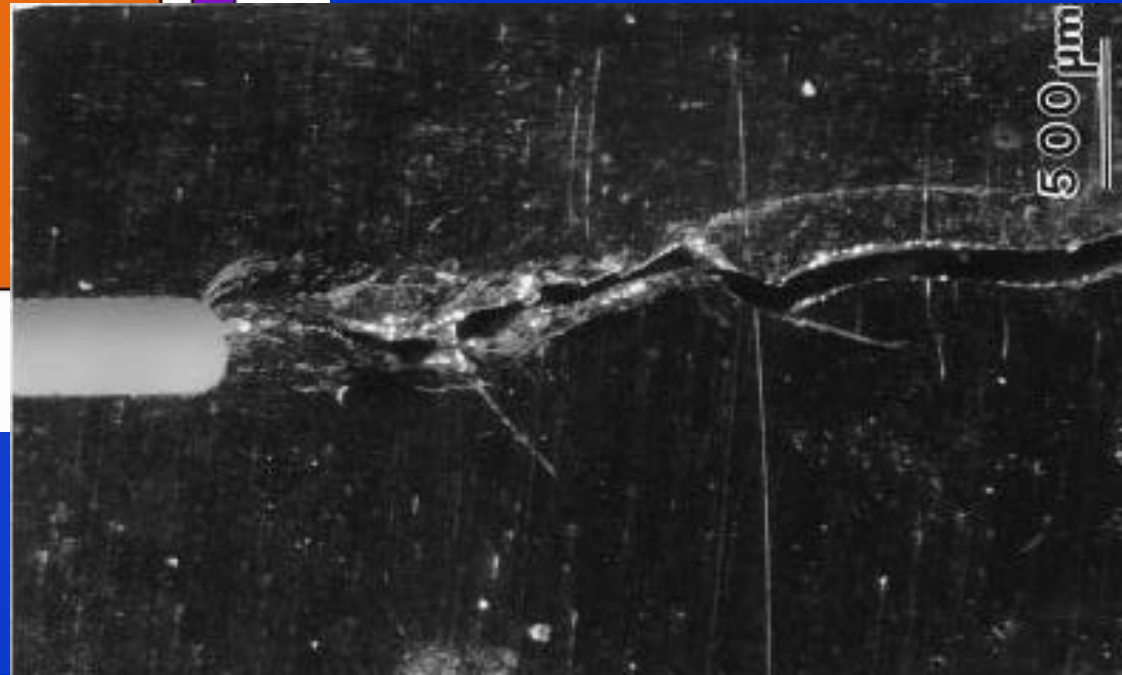
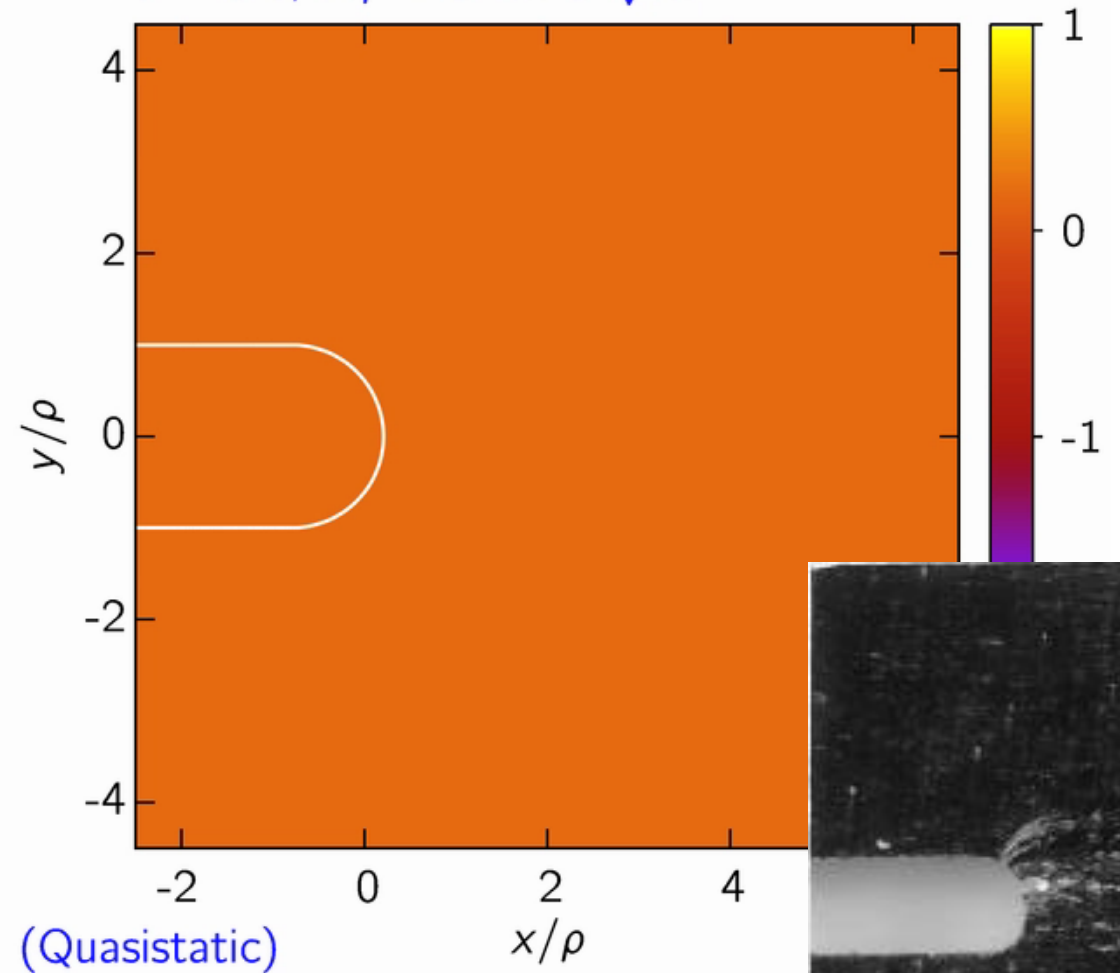
Jiang et al., Philosophical Magazine 2008, 88 407

Results (cont'd)

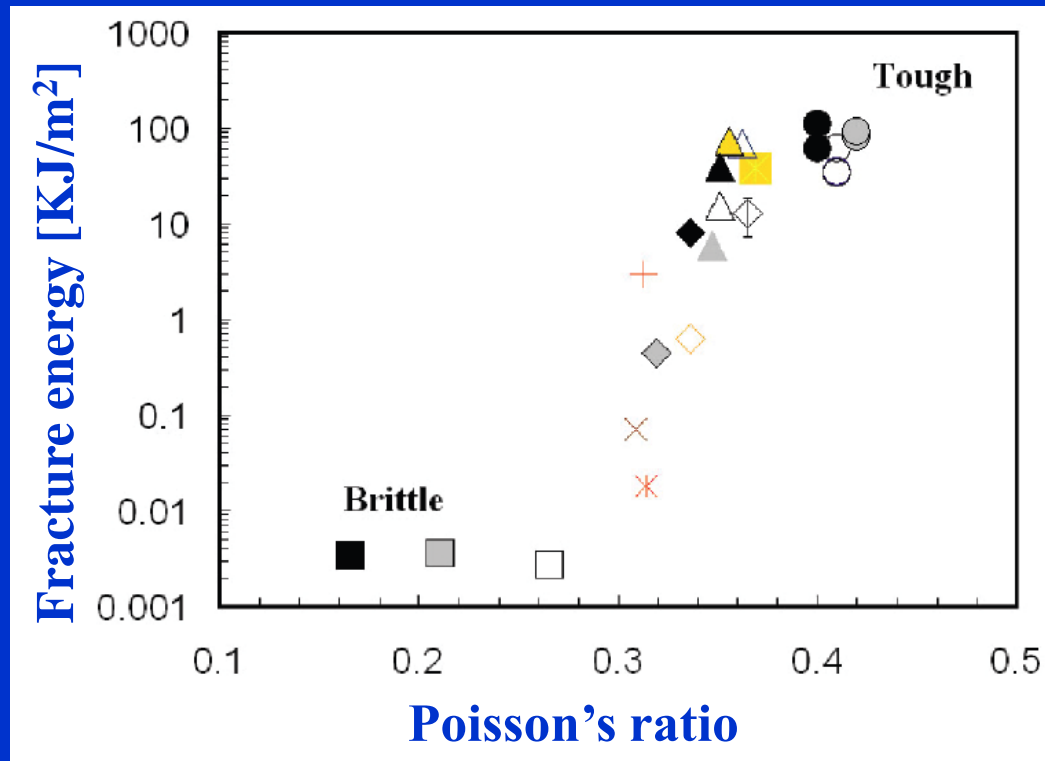


Does local cavitation lead to global failure?

$t = 0 \text{ s}, K_I = 0 \text{ MPa} \sqrt{\text{m}}$



Final comment



$$\nu(\chi)$$

$$\dot{\gamma}^{pl} \sim \Lambda(\chi) = e^{-1/\chi}$$

Summary

The macroscopic theory of glassy deformation based on Shear-Transformation-Zones (STZ) has been discussed.

The theory incorporates coarse-grained internal state variables and characterizes the structural state of the deforming glass by a temperature that may differ from the bath temperature.

We discussed two recent applications of the theory:

- The theory predicts salient features of the variable-amplitude oscillatory shear response of amorphous materials, as observed in recent experiments and computer simulations.**
- The theory offers an explanation for an observed annealing-induced ductile-to-brittle transition and may open the way for a better understanding of the toughness of metallic glasses.**