Implications from GW170817 and I-Love-Q relations for relativistic hybrid stars

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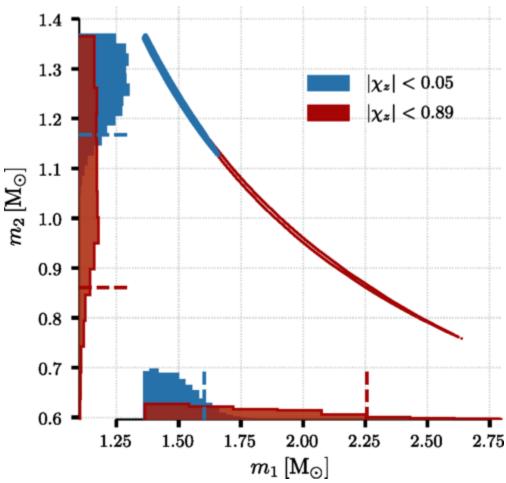
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Collaborators: K. Yagi, D. Alvarez-Castillo, D. Blaschke, A. Sedrakian

GW170817 nature of the progenitor: binary neutron star?

Mass - mass plane

$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{(m_1 + m_2)^{1/5}}$$

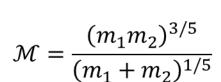


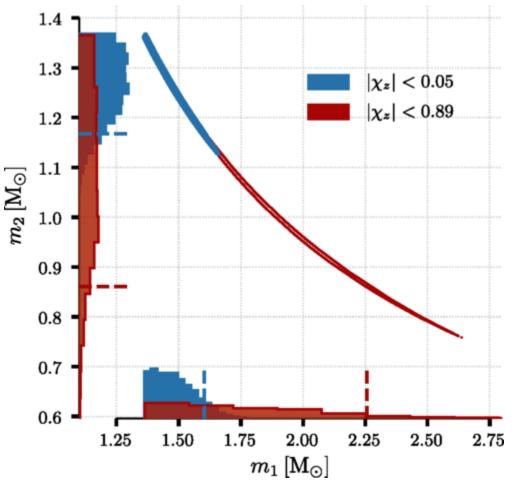
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Could be a long-lived hypermassive NS or a BH.

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GW data do not rule out black hole – neutron star or other exotic compact objects

LIGO/Virgo GW detection of GW170817

- Waveform model
 - Point mass (black hole binary) Post-Newtonian (PN) that include spin-spin effects
 - Finite size effects post-1PN $[O(\frac{v}{c})^2]$ parametrized through the tidal deformabilities $\Lambda_i = \left(\frac{2}{3}\right) k_{2,i} \left[\left(\frac{c^2}{G}\right) \left(\frac{R_i}{m_i}\right)\right]^5$, $k_{2,i}$ = tidal Love numbers

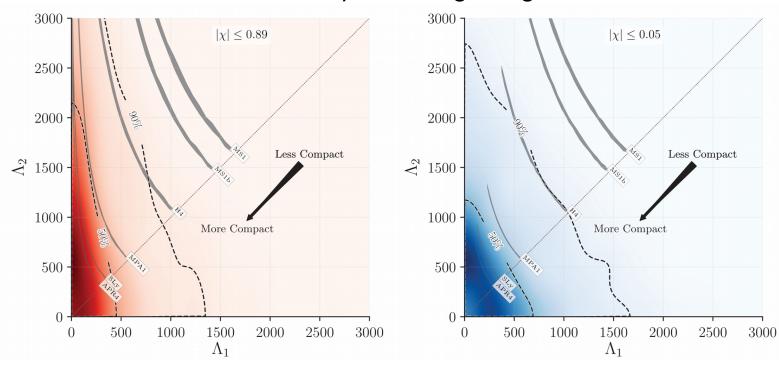
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- For an external gravitational potential (tidal field) $\mathcal{E}_{ij} = \frac{\partial^2 \Phi_{ext}}{\partial x^i \partial x^j}$, there is a quadrupole moment indunced on a star $Q_{ij} = \int d^3x \; \delta \rho(x) \left(x_i x_j \frac{1}{3} r^2 \delta_{ij} \right)$ with $Q_{ij} = -\frac{2R^5}{3G} k_2 \mathcal{E}_{ij}$
- For neutron stars: $k_2 \sim 0.05 0.15$, $\Lambda \sim 150 2000$; For black holes: $k_2 = 0$, $\Lambda = 0$

GW170817: implications for nuclear physics

• Tidal deformability: assuming a single EOS describes both neutron stars

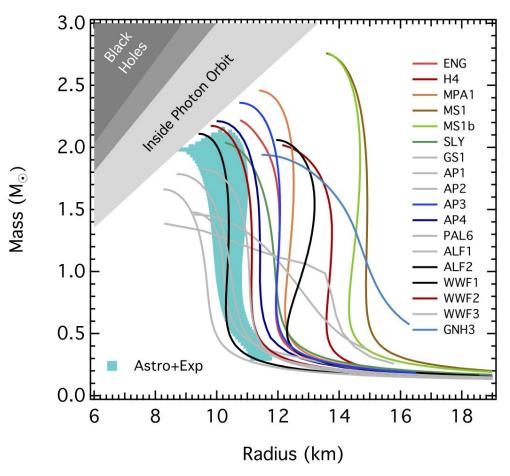


$$\Lambda_{\rm i} = \left(\frac{2}{3}\right) {\rm k}_{2,\rm i} \left[\left(\frac{{\rm c}^2}{\rm G}\right) \left(\frac{{\rm R}_{\rm i}}{{\rm m}_{\rm i}}\right) \right]^5$$
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Data favor more compact stars

GW170817: implications for nuclear physics



Data favor more compact stars

- For sufficiently high densities quark deconfinement can take place
- In 1968 Ulrich Gerlach considered equations of state which allow a first-order transition between the hadronic phase and quark matter phase.
- For a sufficiently large "jump" in energy density over which the pressure remains constant at the phase transition a third family of compact objects can emerge
- Stars with a quark core surrounded by a hadronic shell with a first-order phase transition in between are called hybrid compact stars

There exist two prescriptions for matching the low-density nucleonic EoS to the quark matter EoS; which one is realized in nature depends on the surface tension between nuclear and quark matter (e.g. Glendenning 1997)

If the tension between these phases is low, a mixed phase of quark and nucleonic matter is formed in-between purely nuclear and quark matter phases.

If the tension is high, a sharp transition boundary is energetically favorable → third family of compact object (usually referred to as the "third family")

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Consider two sets of EOSs

Low-density regime

Set I: density functional theory with density-dependent couplings applied to hadronic matter Colucci and Sedrakian 2013

Set II: either stiffest EOS of Hebeler, Lattimer, Pethick, and Schwenk (2013) or the density-dependent relativistic meanfield EoS DD2-p30 Alvarez-Castillo, Ayriyan, Benic, Blaschke, Grigorian, and Typel (2016)

Consider two sets of EOSs

High-density regime

Set I: Constant sound speed parametrization Alford, Han, and Prakash 2013

Set II: Piecewise polytropic representation (Read et al. 2009, Hebeler et al. 2013 etc.)

The parameterizations chosen can be reproduced by a relativistic density functional approach

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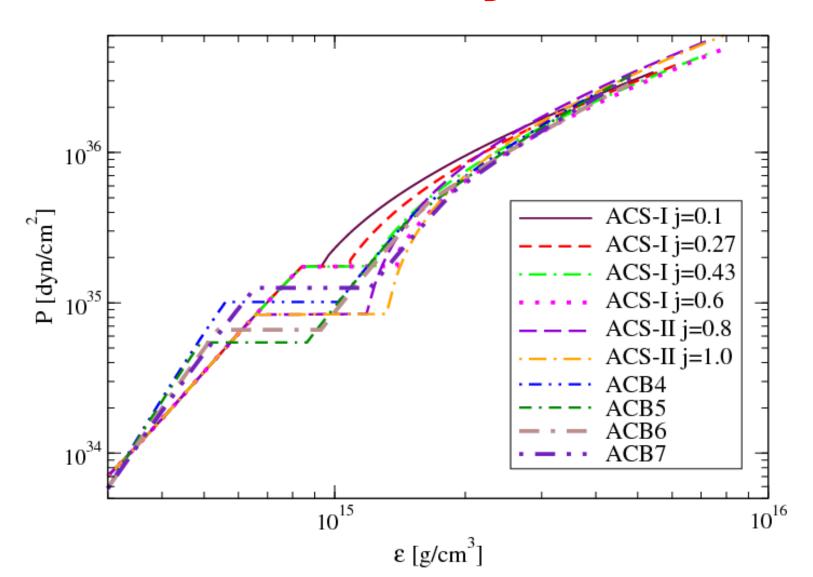
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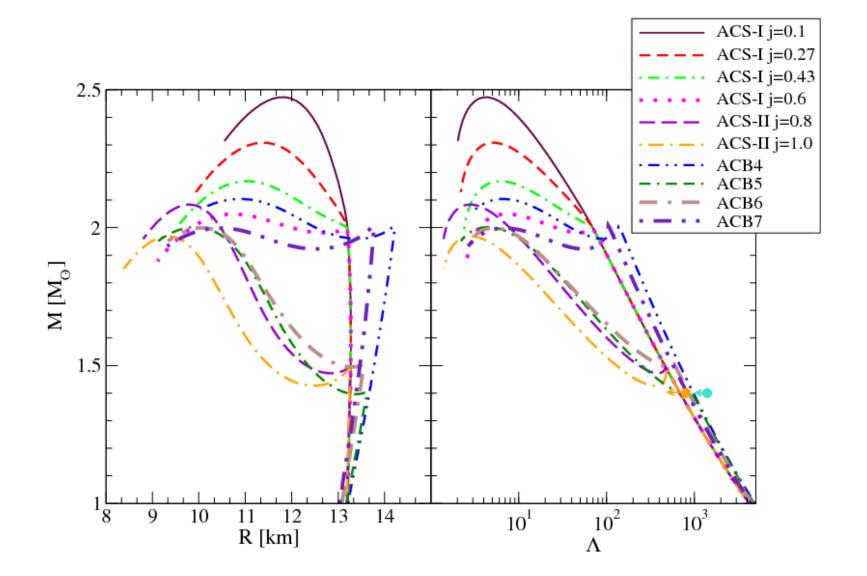
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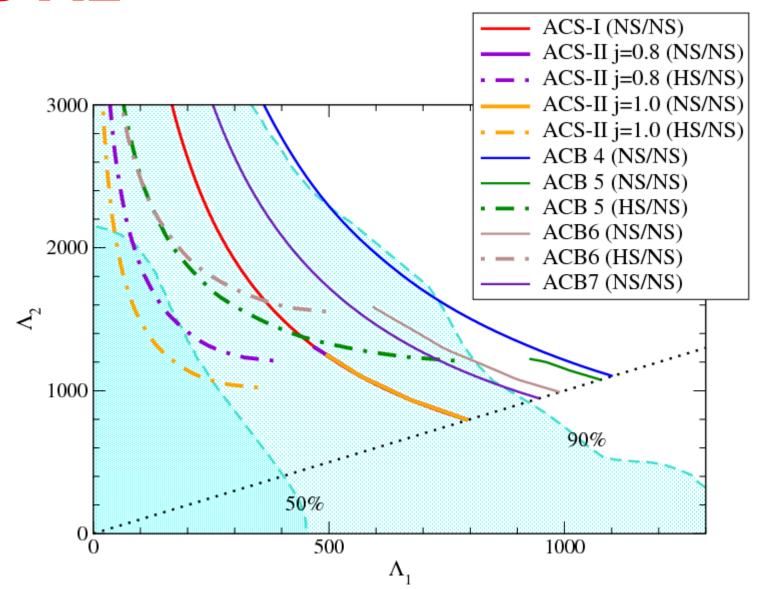
Pressure vs density



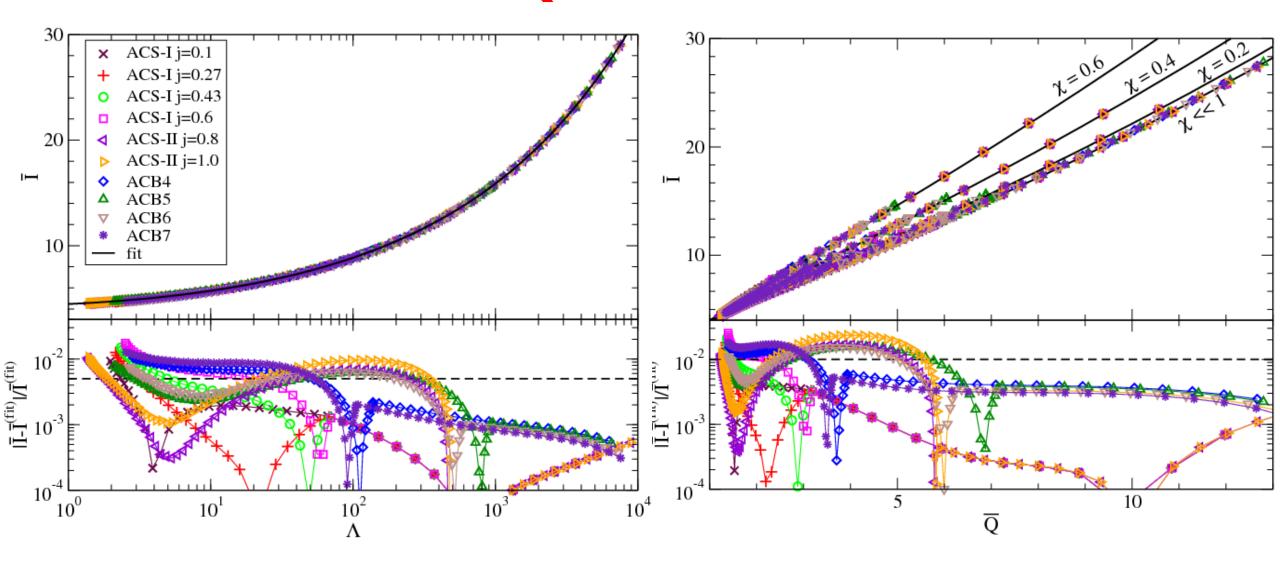
Mass - Radius & Mass - A



$\Lambda 1 \text{ vs } \Lambda 2$



I-Love and I-Q



GW170817 and **GRB 170817A** implications for fundamental physics: nuclear physics

- What could the merger remnant and hence the sGRB engine be?
 - NSNS merges -> Remnant collapse to BH -> BH + disk -> jet
 - NSNS merges -> Remnant is a massive, ms mangetar + disk -> jet
- Observations cannot cleanly distinguish these scenarios, but constraints on the nuclear equation of state (EOS) can be placed from either scenario

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- Observations cannot cleanly distinguish these scenarios, but constraints on the nuclear equation of state (EOS) can be placed from either scenario
- NS matter cannot support an arbitrary amount of baryonic mass
- For any given EOS there is a maximum baryonic mass ${\cal M}_{B,max}$ that can be supported
- If $m_{1B}+m_{2B}-m_{\infty}>M_{B,max}$, the remnant will ultimately collapse and form a BH
- If $m_{1B} + m_{2B} m_{\infty} < M_{B,max}$, the remnant will not collapse