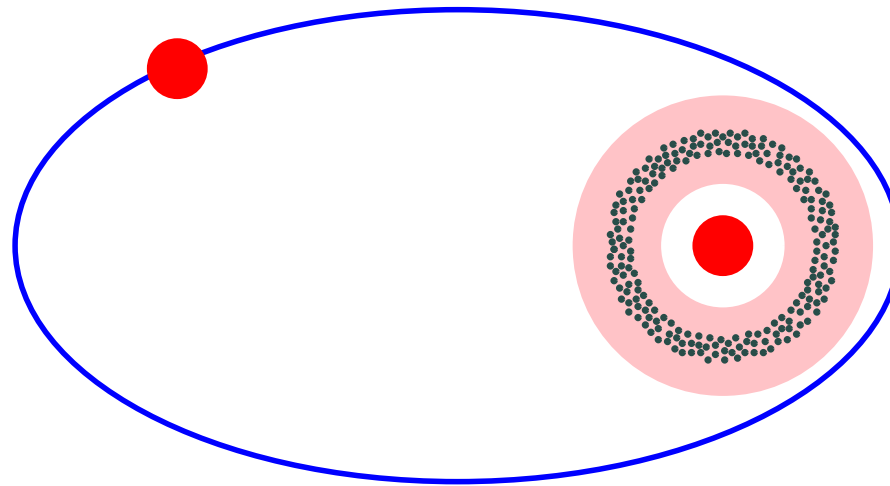


Planetesimal Dynamics in Binary Stellar Systems



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A 100% Unix-Based Presentation

Motivation

Observation

Binaries: more common than single stars

Disks: common among young binaries

Planets: At least 5 extrasolar planets in binaries

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Standard Scenario of Planet Formation

Single Star + Protoplanetary Disk \Rightarrow Planetary System

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Standard Scenario of Planet Formation

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A Theory for Planet Formation in Binaries

Binary + Protoplanetary Disk \Rightarrow Planetary System

Low-Mass PMS Binaries

Frequency

$f_{\text{binary}} \gtrsim 0.5$ for systems with $M_1 + M_2 \lesssim 3M_{\odot}$

Properties

- Mass ratio of secondary to primary $0 \lesssim M_2/M_1 \lesssim 1$
- Orbital characteristics:
 - $1\text{day} \lesssim P \lesssim 10^8\text{days}$ with $\langle P \rangle \simeq 10^5\text{days}$
($10^{-2}\text{AU} \lesssim a \lesssim 10^4\text{AU}$ with $\langle a \rangle \simeq 10^2\text{AU}$)
 - $0 \lesssim e \lesssim 1$ and e_{max} increases with P
- Disk frequency $f_{\text{disk}} \simeq 0.5$
 - circumstellar disk for $a \gtrsim$ a few AU
 - circumbinary disk for $a \lesssim$ a few AU

Previous Studies

Planetesimal Dynamics

Heppenheimer (1978)

- perturbation by secondary and gas drag
- 2-D, no self-gravity

Marzari & Scholl (2000)

- perturbation by secondary, gas drag, and collision
- 2-D, no self-gravity

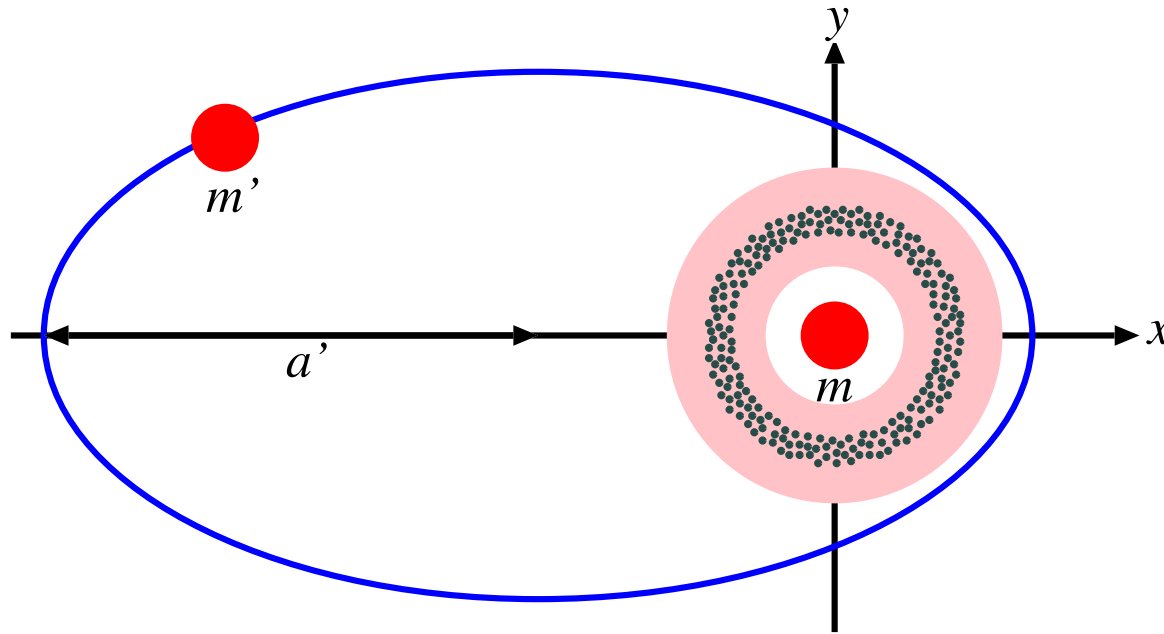
Planetesimal Accretion

Barbieri, Marzari, & Scholl (2002)

Quintana et al. (2002)

- late accretion stage from protoplanets to planets

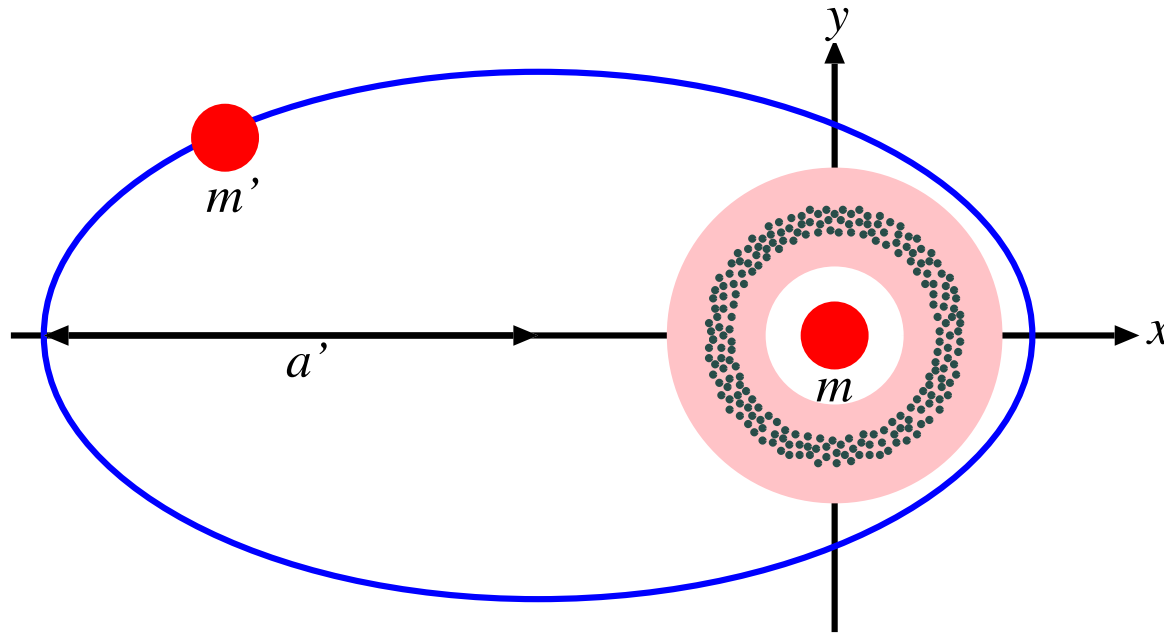
Equation of Motion for Planetesimals



$$\frac{d\mathbf{v}_i}{dt} = \underbrace{-Gm \frac{\mathbf{x}_i}{|\mathbf{x}_i|^3}}_{\text{gravity of primary}} + \underbrace{Gm' \frac{\mathbf{x}' - \mathbf{x}_i}{|\mathbf{x}' - \mathbf{x}_i|^3}}_{\text{gravity of secondary}} + \underbrace{\sum_{j \neq i}^N Gm_j \frac{\mathbf{x}_j - \mathbf{x}_i}{|\mathbf{x}_j - \mathbf{x}_i|^3}}_{\text{mutual gravity}}$$

$$+ \underbrace{\mathbf{f}_{\text{gas}}}_{\text{gas drag}} + \underbrace{\mathbf{f}_{\text{col}}}_{\text{collision effect}} + \underbrace{-Gm' \frac{\mathbf{x}'}{|\mathbf{x}'|^3}}_{\text{indirect term}}$$

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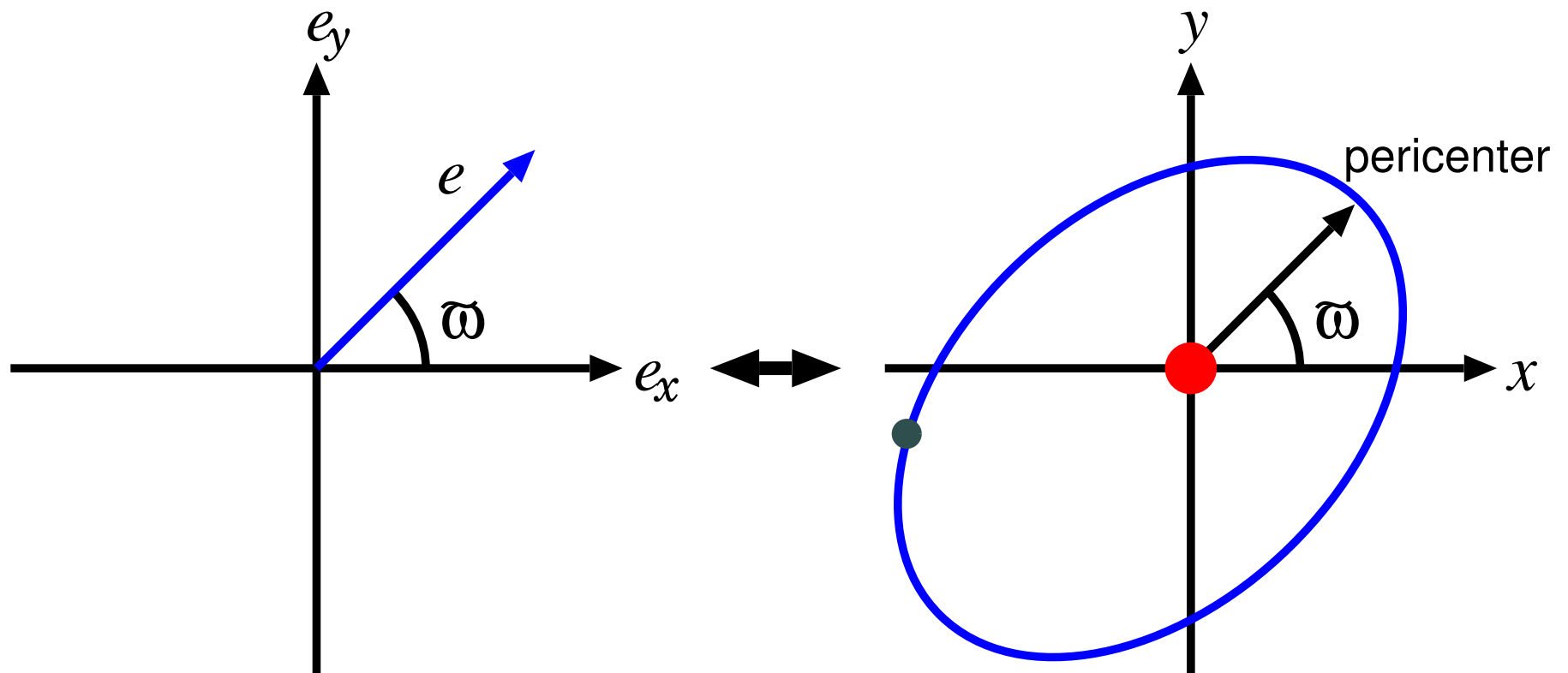
$$\underbrace{-Gm' \frac{\mathbf{x}'}{|\mathbf{x}'|^3}}_{\text{indirect term}}$$

Eccentricity Vector

$$\mathbf{e} = (e_x, e_y) = (e \cos \varpi, e \sin \varpi) = (k, h)$$

e : eccentricity

ϖ : longitude of pericenter



Relative Velocity and Eccentricity

Relative Velocity

$$|\mathbf{v}_{ij}| = |\mathbf{v}_j - \mathbf{v}_i| \simeq |\mathbf{e}_j - \mathbf{e}_i|v_K = |\mathbf{e}_{ij}|v_K \simeq \sigma_e v_K$$

Eccentricity Dispersion

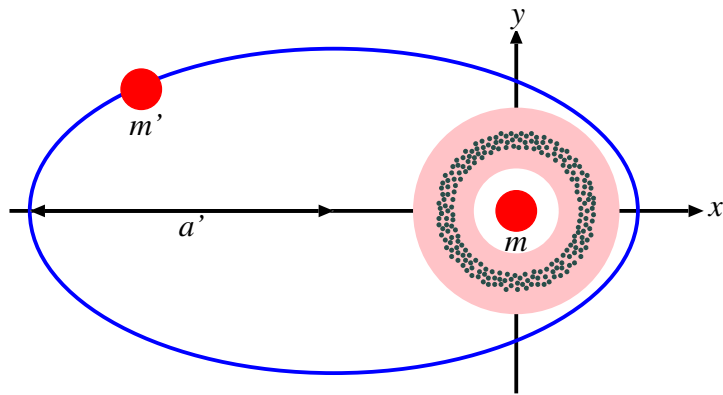
$$\sigma_e = \sqrt{\langle |\mathbf{e}_i|^2 \rangle - |\langle \mathbf{e}_i \rangle|^2}$$

controls

- the growth mode
- the growth time scale

Numerical Experiments

Model



| | |
|--------------|-----------------------------------|
| primary | m |
| secondary | $m', a', e', i' = 0, \varpi', n'$ |
| planetesimal | $m_p, a, e, i, \varpi, \Omega, n$ |

$$m = m' = M_{\odot}, a' = 25[\text{AU}], e' = 0.5$$

(α Cen model)

Initial Conditions

$$N = 1000, m_p = 10^{24}\text{g}$$

$$a = 1\text{AU}(\Delta a = 0.1\text{AU}), \sigma_e = 2\sigma_i$$

$$(\text{minimum-mass disk model: } \Sigma = 10\text{gcm}^{-2})$$

Method of Calculation

4th-order Hermite integrator with GRAPE-6

Disturbing Function of Secondary

Assumptions

$$a/a' \ll 1, m_p \ll m + m', e, i \ll 1$$

Disturbing Function

$$R \simeq \frac{m'}{m + m'} n'^2 a^2 \left[\frac{1}{4(1 - e'^2)^{3/2}} \left(1 + \frac{3}{2} e'^2 \right) - \frac{3}{2(1 - e'^2)^{3/2}} \sin^2 \frac{i}{2} - \frac{15}{16} \frac{a}{a'} \frac{e'}{(1 - e'^2)^{5/2}} e \cos(\varpi - \varpi') \right]$$

Lagrange's Planetary Equation

New Variables

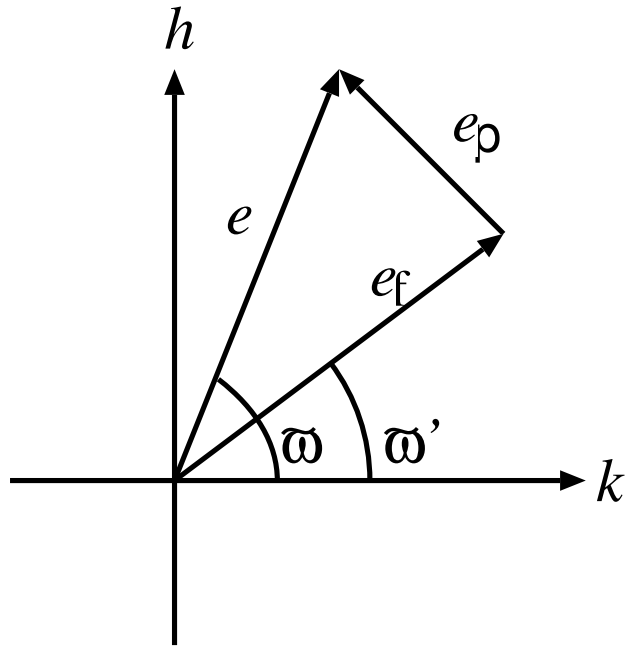
$$h = e \sin \varpi, \quad k = e \cos \varpi$$
$$p = \sin \frac{i}{2} \sin \Omega, \quad q = \sin \frac{i}{2} \cos \Omega$$

Linearized Planetary Equation

$$\begin{cases} \frac{dh}{dt} = \frac{1}{na^2} \frac{\partial R}{\partial k} = Ak - B \cos \varpi' \\ \frac{dk}{dt} = -\frac{1}{na^2} \frac{\partial R}{\partial h} = -Ah - B \sin \varpi' \end{cases} \quad \begin{cases} \frac{dp}{dt} = \frac{1}{4na^2} \frac{\partial R}{\partial q} = -Aq \\ \frac{dq}{dt} = -\frac{1}{4na^2} \frac{\partial R}{\partial p} = Ap \end{cases}$$

$$A = \frac{3}{4} \frac{m'}{m + m'} \frac{n'^2}{n} \frac{1}{(1 - e'^2)^{3/2}}$$
$$B = \frac{15}{16} \frac{m'}{m + m'} \frac{n'^2}{n} \frac{a}{a'} \frac{e'}{(1 - e'^2)^{5/2}}$$

Solution to Planetary Equation



e_p proper eccentricity
 e_f forced eccentricity
 i_p proper inclination
 $i_f = 0$ forced inclination

$$\left\{ \begin{array}{l} h = e_p \sin(At + \varpi_p) + e_f \sin \varpi' \\ k = e_p \cos(At + \varpi_p) + e_f \cos \varpi' \end{array} \right. \quad \left\{ \begin{array}{l} p = \sin \frac{i_p}{2} \sin(-At + \Omega_p) \\ q = \sin \frac{i_p}{2} \cos(-At + \Omega_p) \end{array} \right.$$

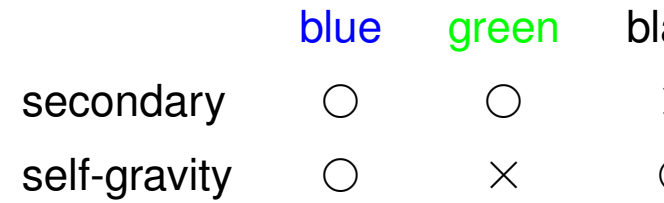
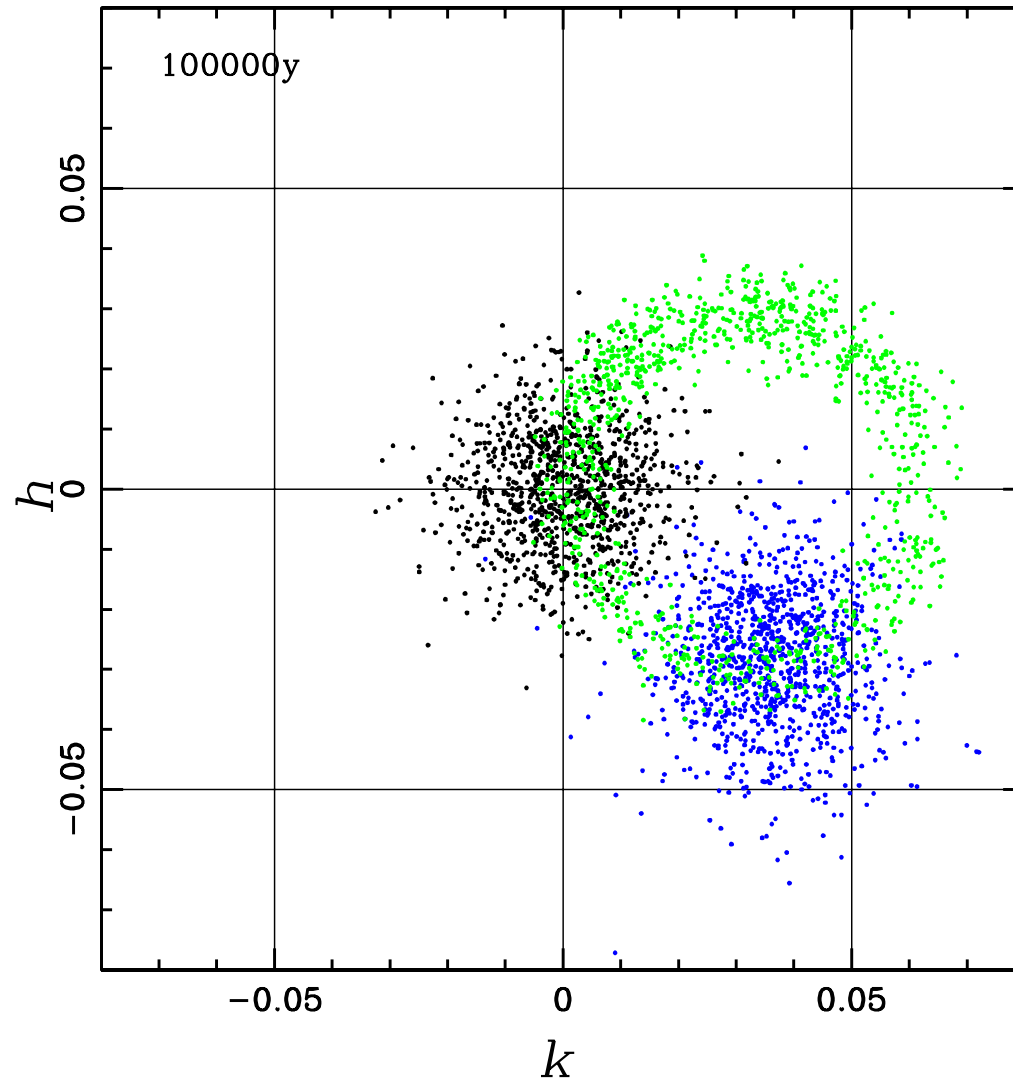
$$e_f = \frac{B}{A} = \frac{5}{4} \frac{a}{a'} \frac{e'}{(1 - e'^2)}$$

Animations

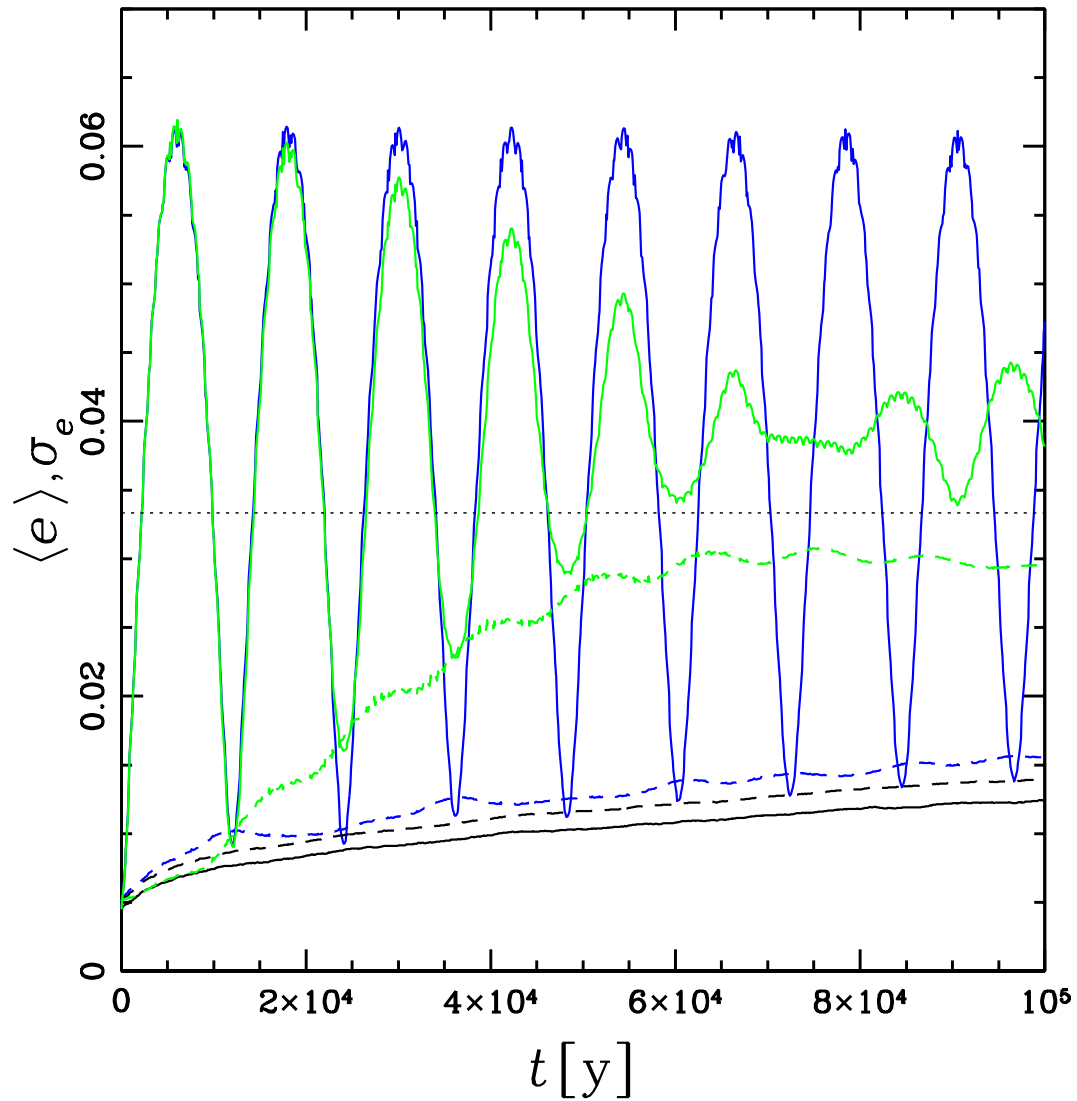
No Self-Gravity Case

Self-Gravity Case

Evolution of Eccentricity

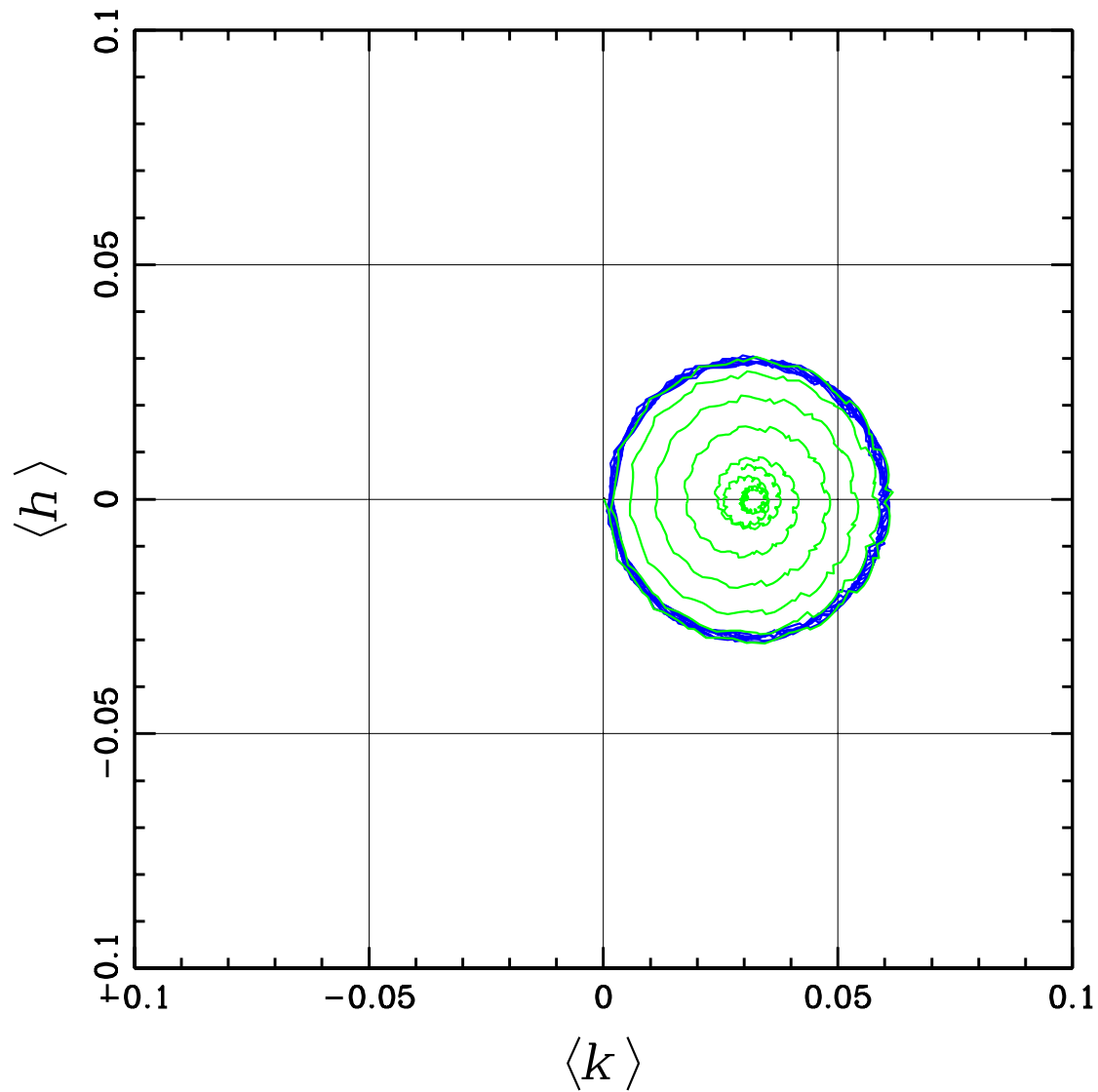


Evolution of Eccentricity



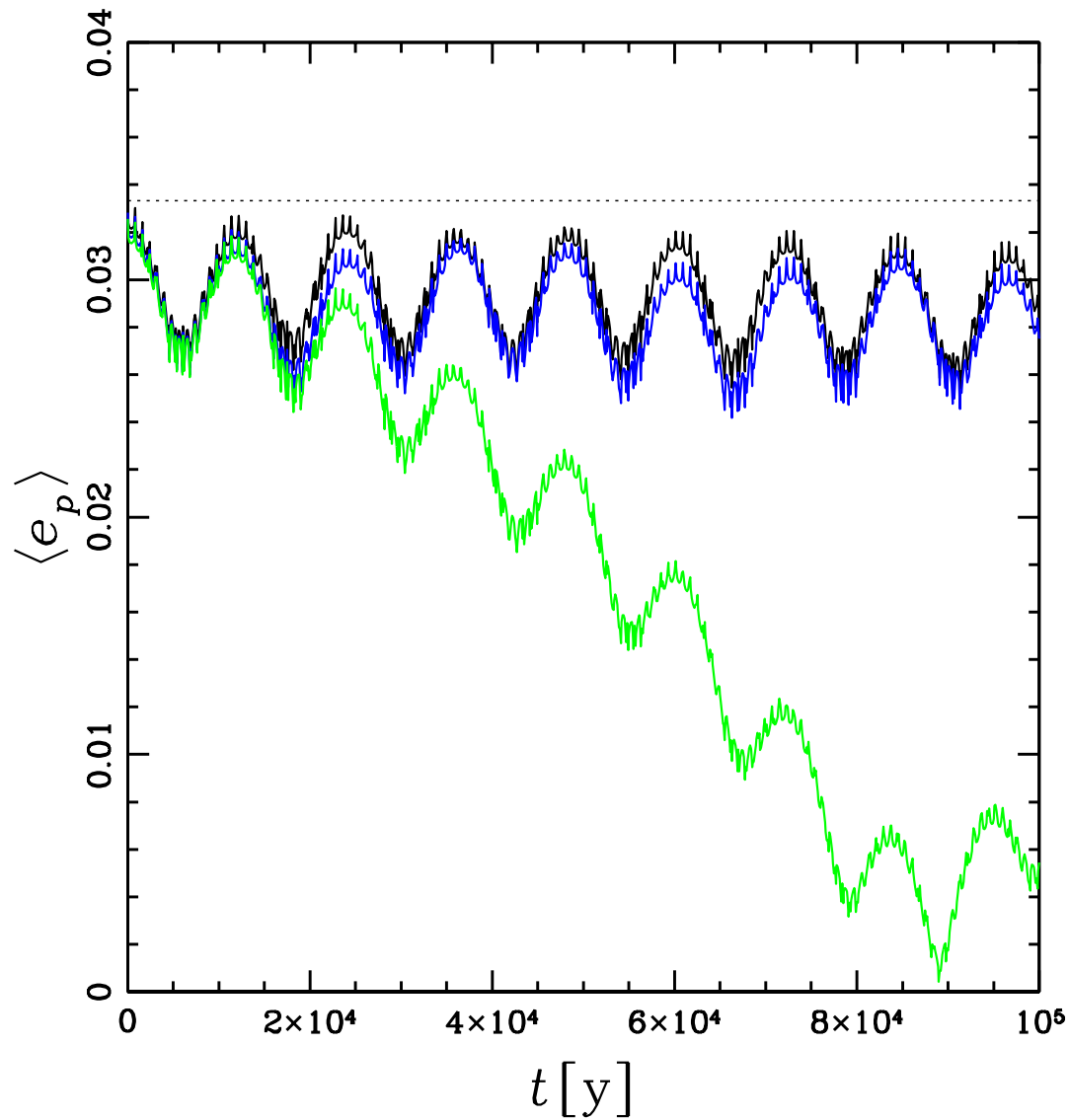
| | | | |
|--------------|---------------------|-------|-------|
| | blue | green | black |
| secondary | ○ | ○ | ○ |
| self-gravity | ○ | × | ○ |
| solid: | $\langle e \rangle$ | | |
| dashed: | σ_e | | |

Evolution of Eccentricity



| | | | |
|--------------|------|-------|----|
| | blue | green | bl |
| secondary | ○ | ○ | ○ |
| self-gravity | ○ | × | ○ |

Evolution of Eccentricity



| | | | |
|------------|-------|------|-------|
| σ_e | black | blue | green |
| | 0.005 | 0.02 | 0.03 |

Time Scales

Time Scale of Shear

$$T_{\text{shear}} \simeq \frac{2\pi}{\frac{dA}{da} \Delta a} = \frac{16\pi}{9} G^{-1/2} m'^{-1} a'^3 (1 - e'^2)^{3/2} m^{1/2} a^{-1/2} \Delta a^{-1}$$

Time Scale of Two-Body Relaxation

$$T_{\text{relax}} \simeq \frac{\sigma^3}{\sqrt{2}\pi G^2 n_m m_p^2 \ln \Lambda} \simeq \frac{m^{3/2} \sigma_e^4}{\sqrt{2}\pi G^{1/2} \Sigma m_p a^{1/2} \ln \Lambda}$$

n_m number density

Σ surface mass density

Condition for Orbital Alignment

Preliminary

Mean Orbital Separation

$$\Delta a = \frac{1}{2\pi a n_s} = \frac{m_p}{2\pi a \Sigma}$$

Condition

$$T_{\text{relax}} \ll T_{\text{shear}}$$

↓

$$\sigma_e < \sigma_e^{\text{crit}} \propto m^{-1/4} m'^{-1/4} a'^{3/4} (1 - e'^2)^{3/8} \Sigma^{1/2} a^{1/4}$$

Summary

Towards a Theory for Planet Formation in Binaries



Planetesimal Dynamics

Planetesimal orbits perturbed by
secondary and self-gravity



Orbital alignment of planetesimals



Small relative velocity between planetesimals
(same relative velocity as the single star case)



Same growth mode and time scale
as the single star case

Things to Do

Near Future (While at KITP?)

- To perform more simulations with other parameters
- To derive the condition for orbital alignment
- To include gas drag

Next Step

- To include accretion

Application

- Satellite-ring interaction
- Eccentric ring