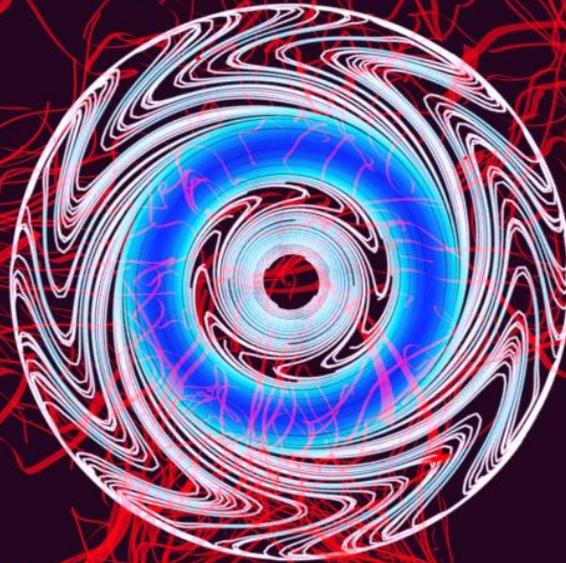


# Beyond chaos: the continuing enigma of turbulence



**Nigel Goldenfeld**

University of Illinois at Urbana-Champaign

Original work with Hong-Yan Shih and Tsung-Lin Hsieh, grudgingly partially supported by NSF



Kavli Institute for  
Theoretical Physics

University of California, Santa Barbara

# What can theoretical physics tell us about turbulence?

# Nothing ... according to Feynman

- There is a physical problem that is common to many fields, that is very old, and that has not been solved. **It is not the problem of finding new fundamental particles, but something left over from a long time ago—over a hundred years.** Nobody in physics has really been able to analyze it mathematically satisfactorily in spite of its importance to the sister sciences. It is the analysis of *circulating or turbulent fluids*.

# Nothing ... according to Feynman

- There is a physical problem that is common to many fields, that is very old, and that has not been solved. **It is not the problem of finding new fundamental particles, but something left over from a long time ago—over a hundred years.** Nobody in physics has really been able to analyze it mathematically satisfactorily in spite of its importance to the sister sciences. It is the analysis of *circulating or turbulent fluids*.
- The simplest form of the problem is to take a pipe that is very long and push water through it at high speed. **We ask: to push a given amount of water through that pipe, how much pressure is needed? No one can analyze it from first principles and the properties of water.** If the water flows very slowly, or if we use a thick goo like honey, then we can do it nicely.

# Richard Feynman



**That is the  
central problem  
which we ought  
to solve some  
day, and we  
have not.**

**What can theoretical physics tell  
us about anything?**

# **What can theoretical physics tell us about anything?**

And actually ... what is theoretical physics?

# Nobel Prize 2016: condensed matter theory



David Thouless

Mike Kosterlitz

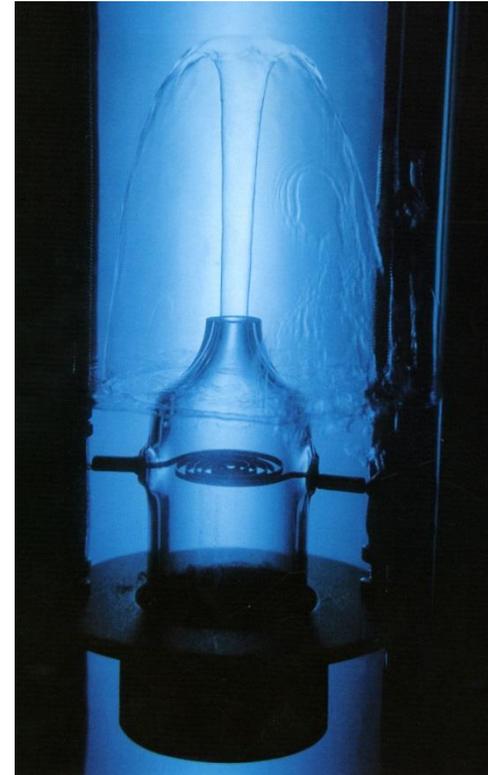
Duncan Haldane

# Superfluids



Honey – normal fluid  
with viscosity at  
room temperature

<http://www.honeyassociation.com/>



Superfluid helium –  
no viscosity at one  
degree above  
absolute zero

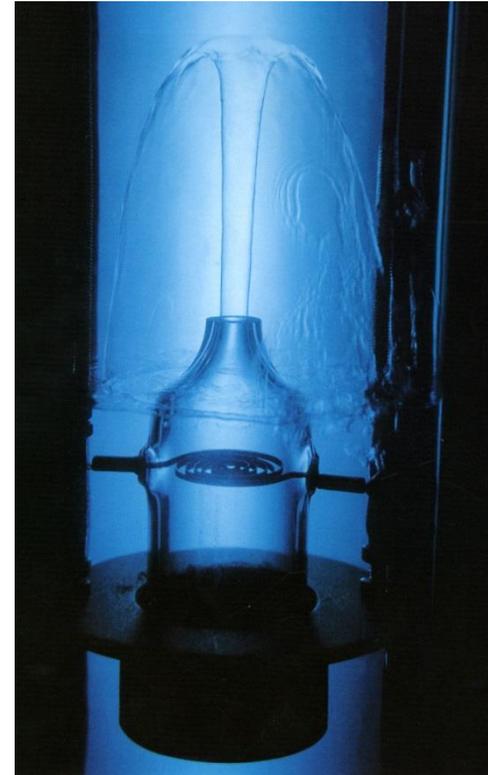
<http://pitp.physics.ubc.ca/archives/CWSS/showcase/topics/He4-fountain.jpg>

# Superfluids



Classical fluid:  
velocity can “swirl”  
as slow as you like

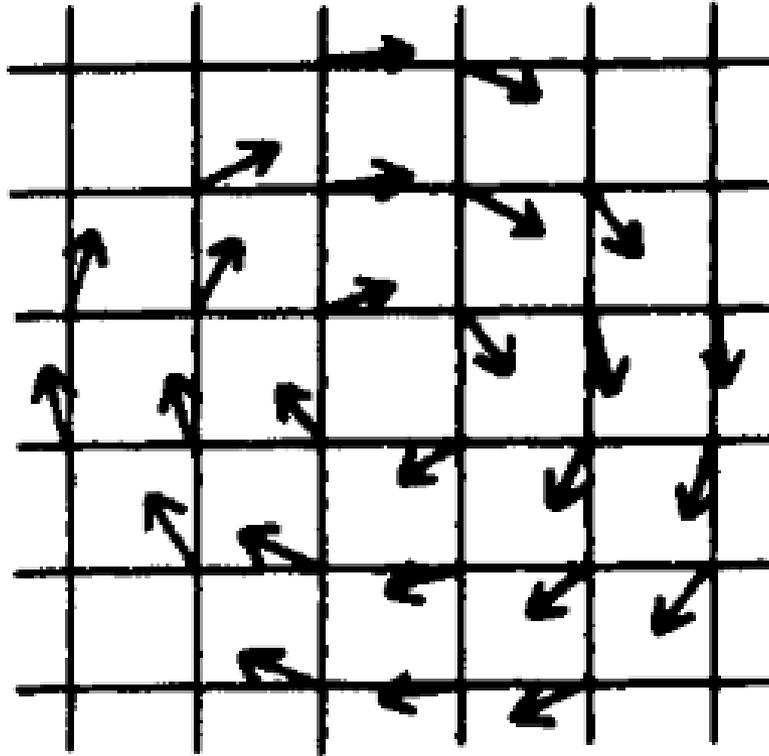
<http://www.honeyassociation.com/>



Quantum fluid: velocity  
cannot “swirl” slower  
than a certain amount

<http://pitp.physics.ubc.ca/archives/CWSS/showcase/topics/He4-fountain.jpg>

# Arrows on a plane



Arrows represent the flow velocity of a quantum fluid.

# Arrows on a plane – predict superfluid film phase transitions

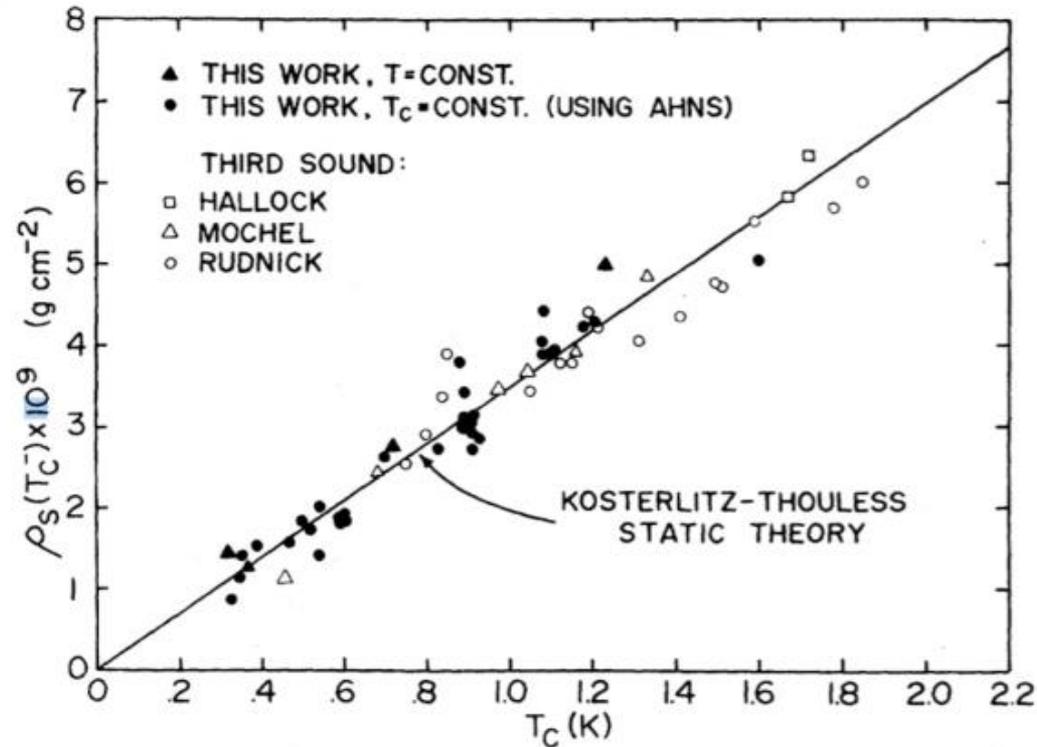


Figure 4: At the Kosterlitz-Thouless transition the superfluid density and critical temperature are predicted to have a linear relation depending only on fundamental constants  $\rho(T_c) = T_c \frac{2}{\pi} \frac{m^2 k_B}{h^2}$ . (Figure from Ref. [10].)

# Arrows on a plane – predict superfluid film phase transitions

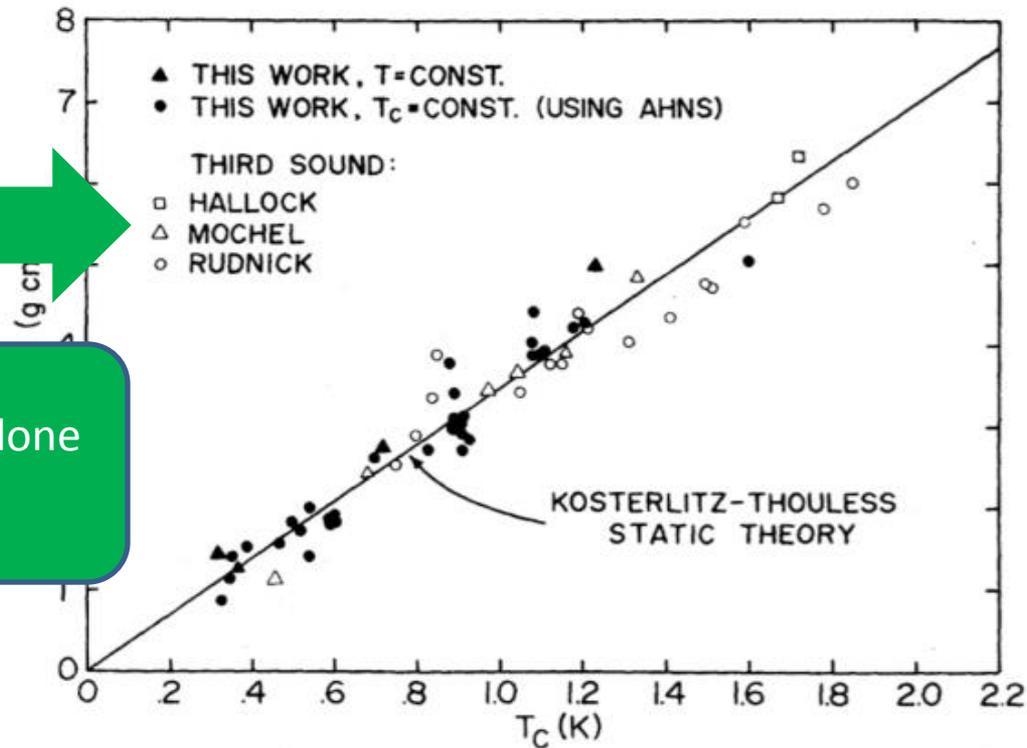


Figure 4: At the Kosterlitz-Thouless transition the superfluid density and critical temperature are predicted to have a linear relation depending only on fundamental constants  $\rho(T_c) = T_c \frac{2}{\pi} \frac{m^2 k_B}{h^2}$ . (Figure from Ref. [10].)

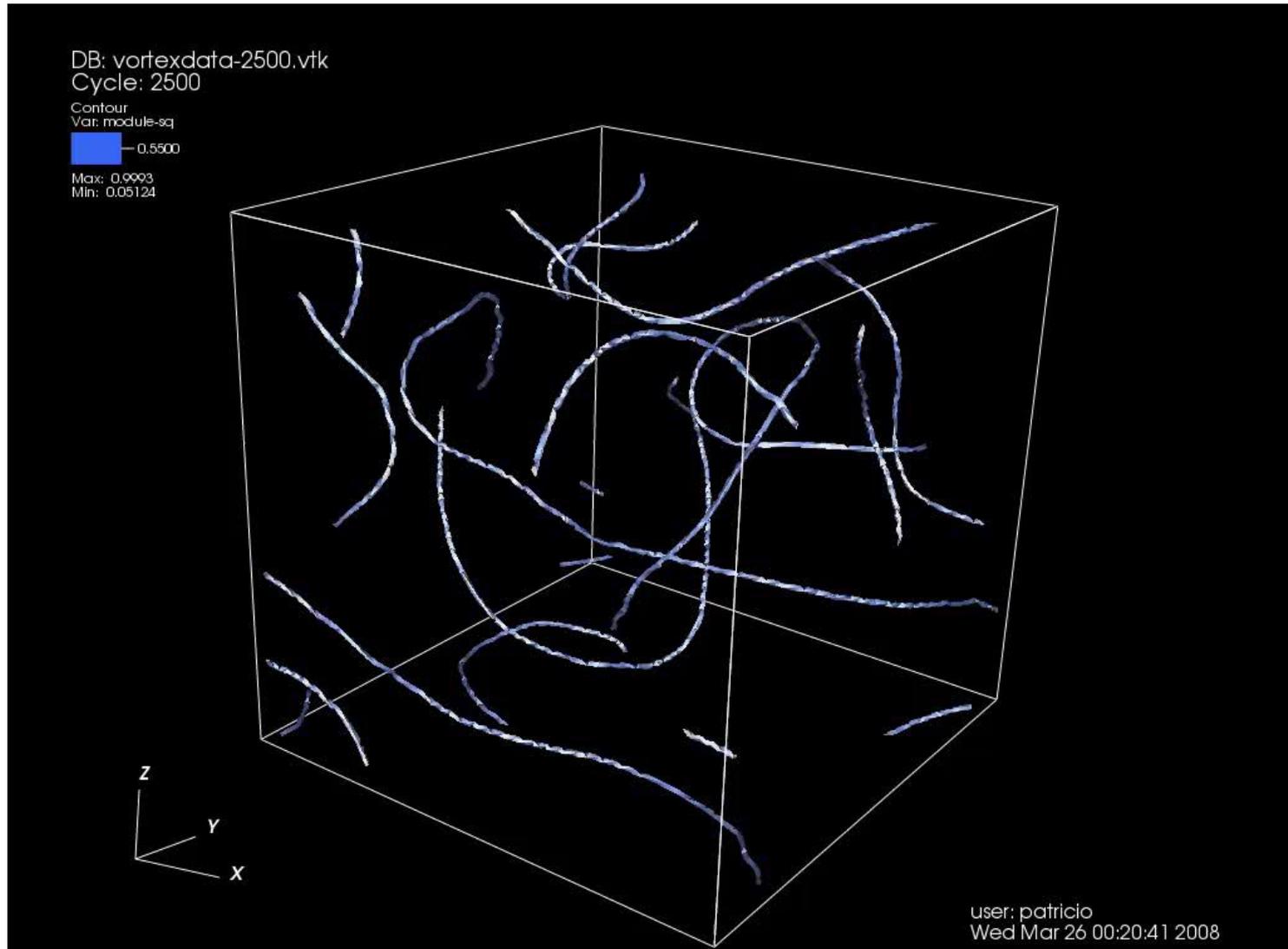
# Amazing and surprising fact

- Superfluids in two dimensions are pretty complicated
- Kosterlitz and Thouless' Nobel Prize work predicted how helium goes from behaving like a normal liquid to a superfluid just by ...
- **Looking at the simplest way that arrows on a plane could be described energetically**

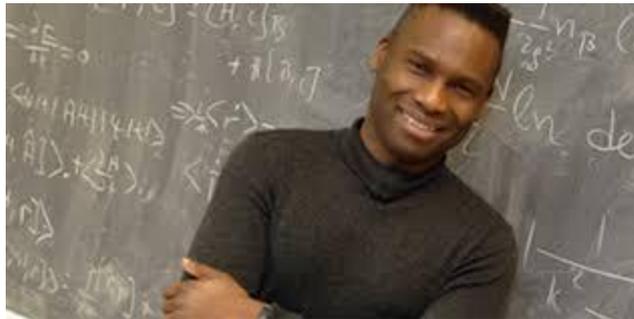
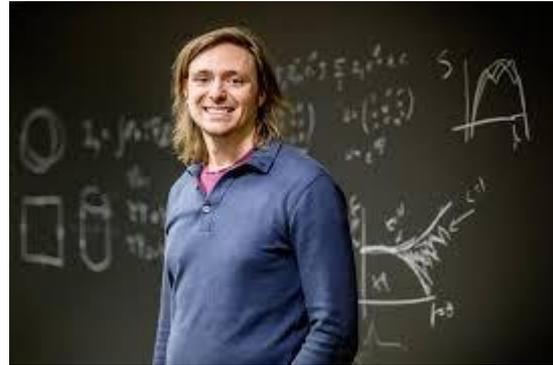
# Amazing and surprising fact

- Superfluous complexity is not needed to describe the world like
  - Kostant's prediction of a normal distribution of the number of particles in a system like
- The moral of the story is that simple models can predict very complex phenomena. In fact everything we know about the physical world comes from looking for the simplest mathematical models (that have the right symmetry and topology).
- **Looking at the simplest way that arrows on a plane could be described energetically**

# Superfluid turbulence in 3D



# Theoretical condensed matter physicists at UIUC



# Is this theoretical physics?

$(\xi_1, \xi_2) \sim \dots$

$P(t) = \dots$   
 Finite size scaling  
 $\sum_1 \sim (P - P_c)^{-\beta}$   
 $\sum_2 \sim (P - P_c)^{-\gamma}$   
 Dynamic scaling  $\xi \sim \xi_c^z$   
 $\xi_c \sim (P_c - P)^{-\nu}$   
 $\nu = \frac{1}{\beta - \gamma}$

①  $P < P_c$

②  $P > P_c$

$P(t, r) = A F\left(\frac{t-r}{S_1(r)}\right)$   
 $\int_{t_0}^{\infty} dt F\left(\frac{t-r}{S_1(r)}\right) = 1$   
 $\frac{d}{dt} F\left(\frac{t-r}{S_1(r)}\right) = -\frac{1}{S_1(r)}$   
 $\int_{t_0}^{\infty} dt \frac{d}{dt} F\left(\frac{t-r}{S_1(r)}\right) = -\frac{1}{S_1(r)} \int_{t_0}^{\infty} dt = -\frac{1}{S_1(r)} \lim_{t \rightarrow \infty} F\left(\frac{t-r}{S_1(r)}\right) + \frac{1}{S_1(r)} F\left(\frac{t_0-r}{S_1(r)}\right)$   
 $0 = -\frac{1}{S_1(r)} + \frac{1}{S_1(r)} F\left(\frac{t_0-r}{S_1(r)}\right)$   
 $F\left(\frac{t_0-r}{S_1(r)}\right) = 1$

Survival probability  
 $P_s(t, r) = \int_{t_0}^{\infty} P(t, r) dt = \int_{t_0}^{\infty} A F\left(\frac{t-r}{S_1(r)}\right) dt$   
 $\tau = \frac{t_0 - r}{S_1(r)}$   
 $F = P(F)$   
 $P_s(t, r) = \int_{t_0}^{\infty} dt F\left(\frac{t-r}{S_1(r)}\right) = \int_{\tau}^{\infty} d\tau F(\tau) = \int_{\tau}^{\infty} F(\tau) d\tau$   
 $\frac{d}{d\tau} F(\tau) = -F(\tau)$   
 $F(\tau) = e^{-\tau}$   
 $P_s(t, r) = \int_{\tau}^{\infty} e^{-\tau} d\tau = e^{-\tau} = e^{-\frac{t_0 - r}{S_1(r)}}$   
 $P_s(t, r) = e^{-\frac{t_0 - r}{S_1(r)}}$   
 $P_s(t, r) = e^{-\frac{t_0 - r}{S_1(r)}}$   
 $P_s(t, r) = e^{-\frac{t_0 - r}{S_1(r)}}$

③  $P(L, r)$

$\times 5000$   $\times 5000$

$t_0 - t_0$   $P_1$   $P_2$

$t_0$   $t_0$

$t_0 < t_0$  and time  $t_0$

Survival probability  
 $P_s(t, r) = \frac{\# \text{ of molecules that survived to } t}{\# \text{ of molecules}}$   
 $P_s(t, r) = \frac{P_r(t_0 > t_0)}{\dots}$   
 $P_s(t, r) = \int_{t_0}^{\infty} P(L, r) dt = 1 - e^{-\dots}$

Occupied sites  $P > P_c$

Holes  $P > P_c$

$L_0(t) > L_1(t) > L_2(t)$

$P_0(L_0) = \frac{1}{\beta} e^{-z} e^{-z}$   
 $\int_{-\infty}^{\infty} P_0(L_0) dL_0 = e^{-e^{-z}}$   
 $z = L_0 - \mu$   
 $P_r(L_0 > t)$

Basic claims:  
 time of puff survival against splitting with separation  $\geq \delta$  is  $\sim L_0$

What does Biot measure?  
 $\log(1 - P) \approx -P$  (if  $P$  small)  
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 $\log(1 - P) \approx -P$  (if  $P$  small)

We need the constraint that the puff splitting events are well-defined and determine the large scale trajectory of a puff. We also want daughter puffs to be viable puffs and not instantly decaying.

$P_s \approx e^{-(t-t_0)/\tau} = P_r(t_0 > t) = 1 - e^{-e^{-\dots}}$   
 $\Rightarrow -\frac{(t-t_0)}{\tau} = \log(1 - e^{-e^{-\dots}}) = -e^{-\dots}$   
 Set  $t_0 = \mu$ ,  $t - t_0 = a$   
 $\frac{a}{\tau} = e^{-e^{-a/P(R_0)}}$

# Is this theoretical physics?

It's my blackboard,  
so probably!

$(\xi_1, \xi_2) \sim \dots$

$P(t) = \dots$   
Finite size scaling

①  $\sum_{L=1}^{\infty} (P-P_c)^{-1/4}$   
 $\sum_{L=1}^{\infty} (P-P_c)^{-1/4}$   
Dynamic scaling  $\tau \sim \xi_c^2$   
 $\xi_c \sim (\xi_1^2)^{1/2}$   
 $\sim (1-P_c)^{-1/2}$   
 $\nu_{11} = \nu_{12}$

②  $P \rightarrow P_c$   
 $P(t, P) = A F\left(\frac{t_0}{\xi_1(P)}\right)$   
Survival probability  
 $P_s(t, P) = \int_{t_0}^{\infty} P(t, P) dt = \int_{t_0}^{\infty} \frac{1}{\xi_1(P)} dt$   
 $\tau = t_0$   
 $F(x) = \dots$   
 $P_s(t, P) = \int_{t_0}^{\infty} dt F(t) = \int_{t_0}^{\infty} F(t) dt$   
 $(1-F)^{1/4} = \dots$   
 $\times (P_c - P) + C(P) = 1 - P_c^2 = (P_c - P)^{2/4}$   
 $P_c < P$

Occupied sites  $P > P_c$

Holes  $P > P_c$

$P_0(L_0) = \frac{1}{\beta} e^{-z L_0}$   
 $z = L_0^{-1}$   
 $\int_0^L P_0(L_0) dL_0 = e^{-e^{-z}}$   
 $P_r(L_0 > t)$

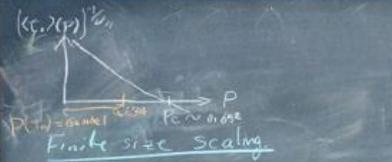
and that the puff splitting is preferred to filter out small scale fluctuations so we have many void events that and determine the large scale trajectory. I also want daughter puffs to be readily extinctly decaying.

$P_r(L_0 > t) = 1 - e^{-e^{-t}}$   
 $\log(1 - e^{-e^{-t}}) = -e^{-t}$   
 $t_0 = a$   
 $\frac{1}{\tau} = e^{-a/P(P_c)}$

$P_s(t, P) = \frac{\# \text{ of molecules that survived to } t}{\# \text{ of molecules}}$   
 $P_s(t, P) = P_r(L_0 > t)$   
 $P_s(t, P) = \int_{t_0}^{\infty} P(L(t)) dt = 1 - e^{-e^{-t}}$

# Is this theoretical physics?

But it's in the  
Institute for  
Genomic Biology,  
so who knows?



①  $\sum_{t=1}^{\infty} P(t) \sim (P - P_c)^{-1}$   $P < P_c$   
 $\sum_{t=1}^{\infty} t P(t) \sim (P - P_c)^{-2}$   
Dynamic scaling  $z \sim \xi_c^2$   
 $\xi_c \sim (P - P_c)^{-1/z}$   
 $\sim (P - P_c)^{-2}$

②  $P \rightarrow P_c$   
 $v_{11} = v_{22}$

$P(t, t_0) = A F\left(\frac{t - t_0}{\xi(t_0)}\right)$

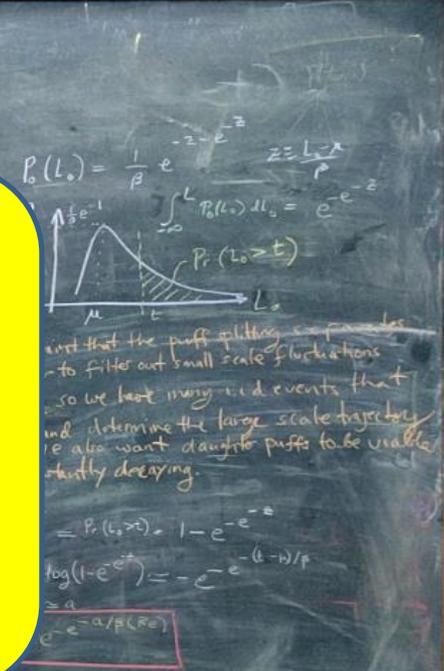
Survival probability

$P_s(t, t_0) = \int_{t_0}^{\infty} P(t, t_0) dt = \int_{t_0}^{\infty} \frac{1}{\xi(t_0)} F\left(\frac{t - t_0}{\xi(t_0)}\right) dt$

$\tau = \frac{t_0}{v_{11}}$   
 $F = P_s$

$P_s(t, t_0) = \int_{t_0}^{\infty} dt F(t) = \int_{t_0}^{\infty} F(t) dt$   
 $(1 - F) = F(t_0) = 1 - \frac{t_0}{\xi(t_0)} F(0) = 1 - \frac{t_0 F(0)}{\xi(t_0)}$   
 $\times (P_c - P) + C(A) = 1 - P_c \xi_c^2 (P - P_c)^{-2} = 1 - \frac{t_0 F(0)}{\xi(t_0)}$

Occupied sites  $P > P_c$   
Holes  $P < P_c$   
③  $P_s(L, t)$



and that the puff splitting is a process to filter out small scale fluctuations so we have many red events that and determine the large scale trajectory. I also want daughter puffs to be readily extinctly decaying.  
 $P_s(t_0, \infty) = 1 - e^{-e^{-t}}$   
 $\log(1 - e^{-e^{-t}}) = -e^{-t}$   
 $\approx -e^{-t}$   
 $e^{-e^{-t}} = e^{-a/P(t_0)}$

# Is this theoretical physics?

③

$P(L, t)$

$x 5000$

$t_{00} - t_0$

$t_0(P_1)$

$t_0(P_2)$

$t_{00}$

$t_0$

Survival probability  $P_S(t, T)$

# of molecules that survived to  $t$

# of molecules

①

$P(L, t)$

$x 5000$

$t_{00} - t_0$

$t_0(P_1)$

$t_0(P_2)$

$t_{00}$

$t_0$

Survival probability  $P_S(t, T)$

# of molecules that survived to  $t$

# of molecules

②

$P \rightarrow P_c$

$P(t, P) = A F\left(\frac{t-t_0}{S_1(P)}\right)$

Survival probability

$P_S(t, T) = \int_{t_0}^T P(t, P) dP = \int_{t_0}^T \frac{1}{S_1(P)} F\left(\frac{t-t_0}{S_1(P)}\right) dP$

$T = t_0$

$F = P(T)$

$P(t, T) = \int_{t_0}^T dP F(P) = \int_{t_0}^T F(P) dP$

$(1-F)^{1/\nu} = F(t, T)^{1/\nu}$

$\times (P_1 - 1) + C(A) = 1 - P_2^{1/\nu} = (P_1 - P_c)^{1/\nu}$

④ Occupied sites  $P > P_c$

Holes  $P < P_c$

$L_0(t) > L_1(t) > L_2(t)$

Basic claims:

time of puff survival against splitting with separation  $\geq \delta$  is  $\sim L_0$

at low Biot number?

$\log(1-P) \approx -P$  prob. puff does not split up to time  $t$

$\log(1-P) \approx -P$

$\log(1-P) \approx -P$

We need the constraint that the puff splitting is a power law  $\delta x \geq \delta$  in order to filter out small scale fluctuations and  $\delta x \leq L_{puff}$  so we have many red events that are well-defined and determine the large scale trajectory of a puff. We also want daughter puffs to be viable puffs and not instantly decaying.

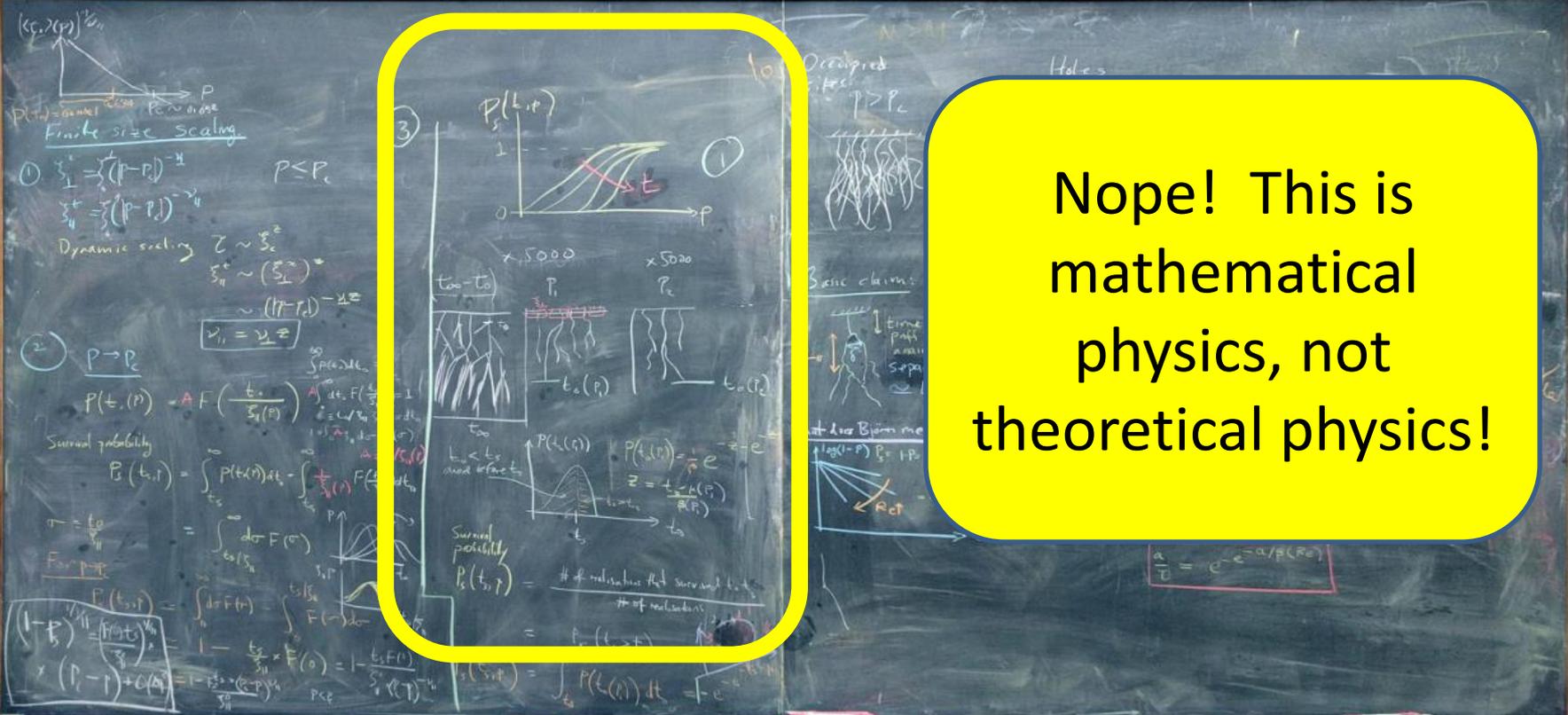
$P_S \approx e^{-(t-t_0)/\tau} = P_c(t_0, T) = 1 - e^{-e^{-\frac{t-t_0}{\tau}}}$

$\Rightarrow -\frac{(t-t_0)}{\tau} = \log(1 - e^{-e^{-\frac{t-t_0}{\tau}}}) = -e^{-\frac{(t-t_0)}{\tau}}$

Set  $t_0 = \mu$ ,  $t - t_0 = a$

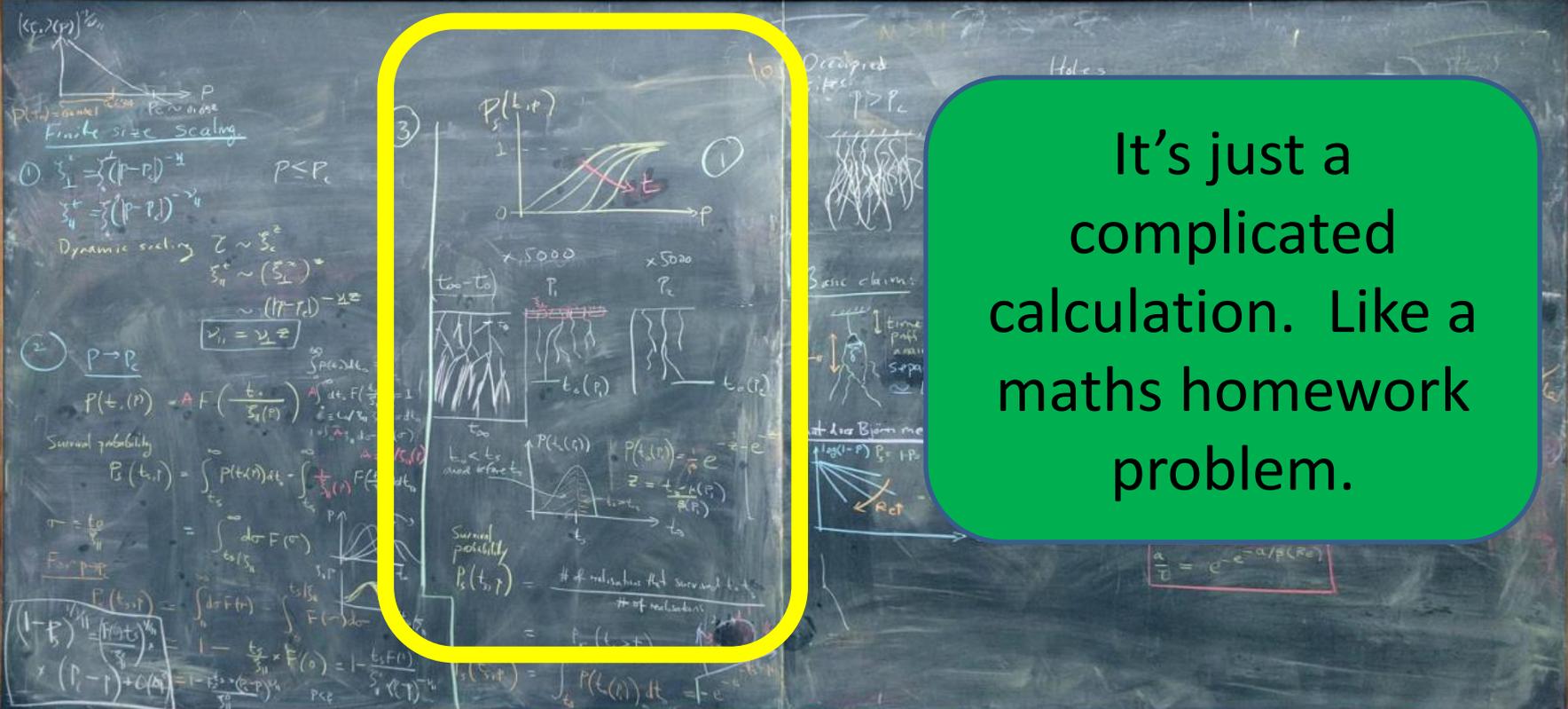
$\frac{a}{\tau} = e^{-e^{-a/\mu}}$

# Is this theoretical physics?



Nope! This is mathematical physics, not theoretical physics!

# Is this theoretical physics?



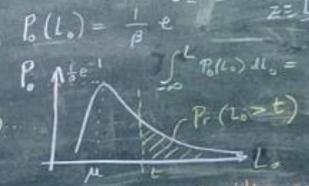
It's just a complicated calculation. Like a maths homework problem.



# Is this theoretical physics?

How about this bit?

$\frac{dL}{dt} = -\frac{L}{\tau}$   
 $P > P_c$



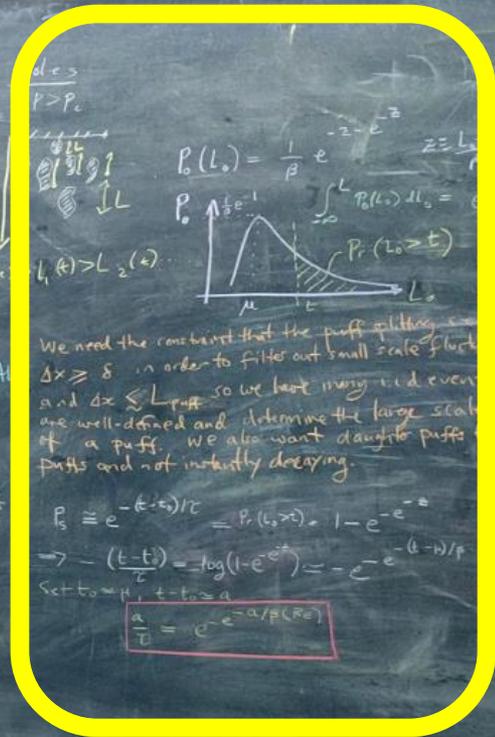
$P_0(L_0) = \frac{1}{\beta} e^{-z/\beta}$   
 $P_0(L_0) = \frac{1}{\beta} e^{-L_0/\beta}$   
 $\int_{L_0}^{\infty} P_0(L) dL = e^{-L_0/\beta}$   
 $P_r(L_0 > t) = e^{-L_0/\beta}$

We need the constraint that the puff splitting is well-defined and determining the large scale behavior of a puff. We also want daughter puffs to be well-defined and not instantly decaying.

$P_s \equiv e^{-(t-t_0)/\tau} = P_r(L_0 > t) = 1 - e^{-L_0/\beta}$   
 $\Rightarrow -\frac{(t-t_0)}{\tau} = \log(1 - e^{-L_0/\beta}) = -e^{-L_0/\beta}$   
Set  $t_0 = \mu$ ,  $t - t_0 = a$   
 $\frac{a}{\tau} = e^{-a/\beta(P_0)}$

# Is this theoretical physics?

Yup! It's a guess about how something behaves -- actually how fluids become turbulent!



# Is this theoretical physics?

Mathematical physics

Theoretical physics

Handwritten notes on a chalkboard, including:

- Left side:**
  - Graph of  $(\xi_1, \xi_2) \sim \omega$  vs  $P$ .
  - Equations:  $\sum_1 \sim (P-P_c)^{-1/2}$ ,  $\sum_2 \sim (P-P_c)^{-1/4}$ .
  - Dynamic scaling:  $\xi \sim \xi_c^z$ ,  $\xi_1 \sim (\xi_c^z)^z$ .
  - Equation:  $\nu_1 = \nu_2 = z$ .
  - Equation:  $P(t, P) = A F\left(\frac{t-t_0}{S_1(P)}\right)$ .
  - Survival probability:  $P_S(t, T) = \int_{t_0}^T P(t, P) dP$ .
  - Equation:  $P_S(t, T) = \int_{t_0}^T dF(P) = F(T) - F(t_0)$ .
  - Equation:  $P_S(t, T) = 1 - \frac{t-t_0}{T-t_0} F(0) = 1 - \frac{t-t_0}{T-t_0} F(P_c)$ .
- Middle (highlighted):**
  - Graph of  $P(L, t, P)$  vs  $P$ .
  - Diagram of branching structures with labels  $t_0$  and  $t$ .
  - Equation:  $P(L, t, P) = \frac{1}{P} e^{-z} e^{-zL}$ .
  - Equation:  $z = t - t_0 F(P)$ .
  - Survival probability:  $P_S(t, T) = \frac{\text{\# of realizations that survived to } t}{\text{\# of realizations}}$ .
- Right side (highlighted):**
  - Equation:  $P_S(L_0) = \frac{1}{P} e^{-z}$ .
  - Equation:  $P_S(L_0) = \int_{L_0}^{\infty} P(L, t, P) dL$ .
  - Equation:  $P_S(L_0) = 1 - e^{-z}$ .
  - Equation:  $z = t - t_0 F(P)$ .
  - Equation:  $z = -\log(1 - e^{-z}) = -\log(1 - e^{-z})$ .
  - Equation:  $\frac{a}{T} = e^{-a/P(P_c)}$ .
- Other notes:**
  - Occupied sites  $P > P_c$ .
  - Basic claims: time of puff survival against splitting with separation  $\geq \delta$  is  $\sim L_0$ .
  - at low Biot number?  $\log(1-P) \approx -P$  prob. puff does not split up to time  $t$ .
  - Equation:  $\log(1-P) \approx -P$ .



**Why do we need to guess in  
turbulence?**

**Why do we need to guess in  
turbulence?**

$$\mathbf{F} = m \mathbf{a}$$

# Why do we need to guess in turbulence?

$$F = m a$$

Force

mass

acceleration

# Why do we need to guess in turbulence?

$$a = F/m$$

acceleration

Force

mass

# Why do we need to guess in turbulence?

$$a = F/m$$

The force in fluid dynamics is usually gravity, pressure, viscosity ...

# Why do we need to guess in turbulence?

$$a = F/m$$

The force in fluid dynamics is usually gravity, pressure, viscosity ... **but it could be magnetic fields in stars or a plasma fusion experiment**

...

# Why do we need to guess in turbulence?

$$a = F/m$$

The force in fluid dynamics is usually gravity, pressure, viscosity ... **but it could be magnetic fields in stars or a plasma fusion experiment ... or even come from quarks and gluons in a particle physics experiment**

**If we know acceleration, we can  
calculate speed**

**Acceleration means “how much your speed changes every second”**

# Mid-term exam

**Question:** A car accelerates from rest at 9 meters per second per second for three seconds.

Please tell me:

# Mid-term exam

**Question:** A car accelerates from rest at 9 meters per second per second for three seconds.

Please tell me: **What is the brand of the car?**

# Mid-term exam

**Question:** A car accelerates from rest at 9 meters per second per second for three seconds.

Please tell me: **What is the brand of the car?**

**Answer:** After 1 second it is going 9 meters/second. After 2 seconds it is going 18 meters/second. After 3 seconds it is going 27 meters/second.

# Mid-term exam

**Question:** A car accelerates from rest at 9 meters per second per second for three seconds.

Please tell me: **What is the brand of the car?**

**Answer:** After 1 second it is going 9 meters/second. After 2 seconds it is going 18 meters/second. After 3 seconds it is going 27 meters/second.

27 meters/second = 60 miles per hour

**0-60 mph in 3 seconds. Must be a Porsche!**

# Acceleration of a fluid

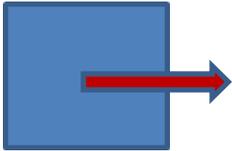
When you try to apply

$$\mathbf{a} = \mathbf{F}/m$$

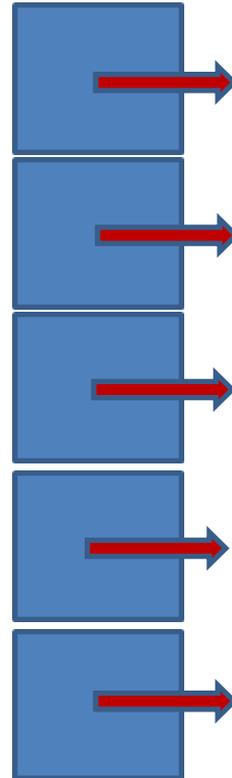
to every point in a fluid, not just a single particle, it gets very complex.

# Acceleration of a fluid

**PARTICLE**

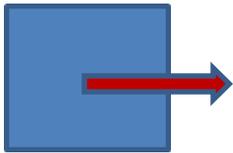


**FLUID**



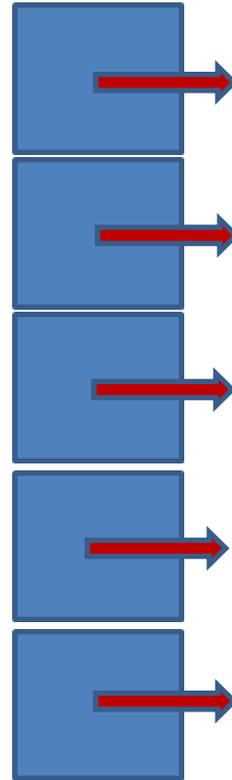
# Acceleration of a fluid

**PARTICLE**



Predictable

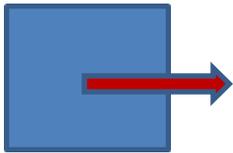
**FLUID**



Predictable

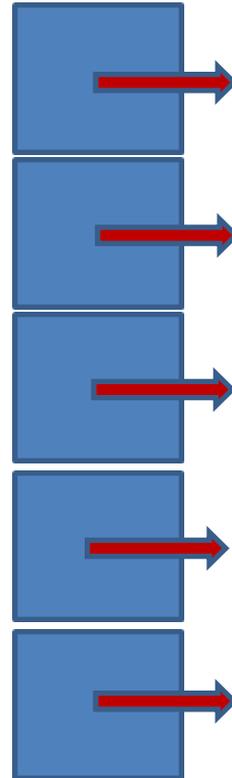
# Acceleration of a fluid

**PARTICLE**



Predictable

**FLUID**

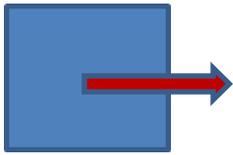


Predictable

Laminar

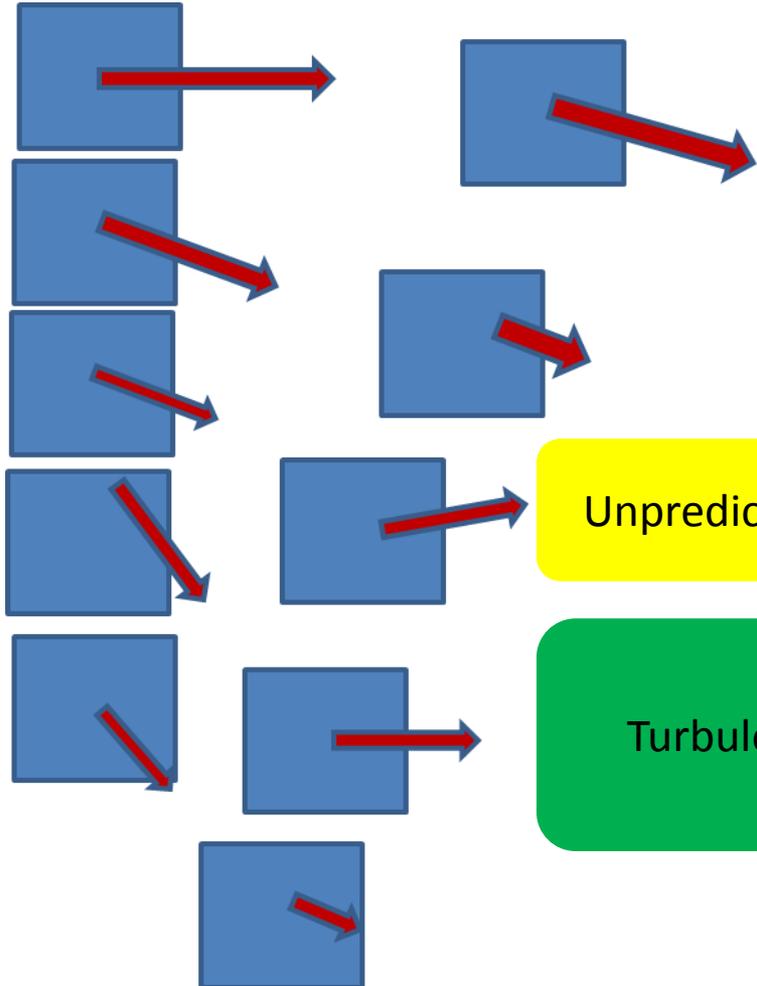
# Acceleration of a fluid

**PARTICLE**



Predictable

**FLUID**

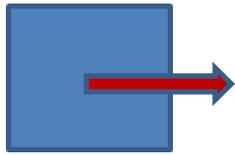


Unpredictable

Turbulent

# Chaos vs. Turbulence

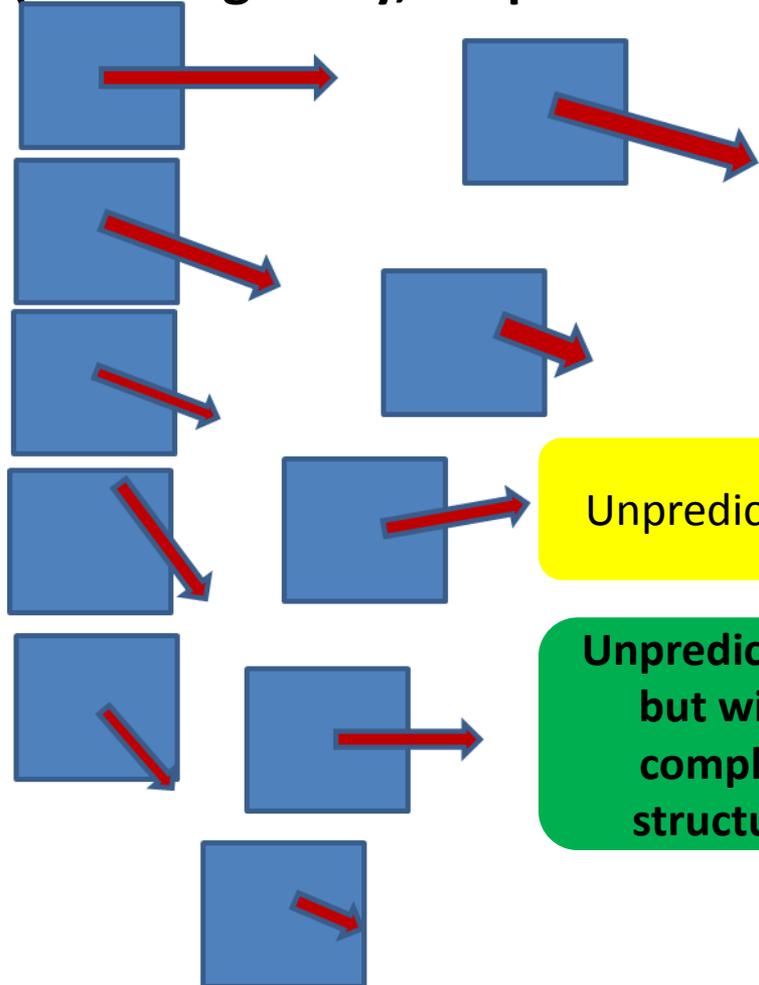
Chaos – one particle  
(behaving badly, in time)



Unpredictable

Unpredictable  
but with  
“simple”  
structure

Turbulence – many particles  
(behaving badly, in space & time)

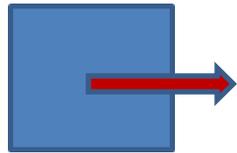


Unpredictable

Unpredictable  
but with  
complex  
structure

# Acceleration of a fluid

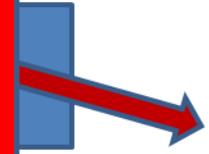
**PARTICLE**



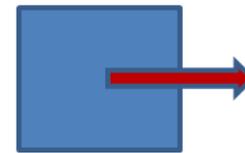
Predictable

Velocity of fluid blob depends on where it was, and where it is going to, and how fast it is going.

So the velocity governs its own behavior in some self-referential way.



Unpredictable



Turbulent

**What is turbulence?**

# Turbulence is stochastic and wildly fluctuating

Soap film experiment



M. A. Rutgers, X-l. Wu, and W. I. Goldburg.  
"The Onset of 2-D Grid Generated Turbulence in Flowing Soap Films,"  
Phys. Fluids 8, S7, (Sep. 1996).

# Turbulence generates structure at many scales

Soap film experiment



M. A. Rutgers, X-l. Wu, and W. I. Goldburg.  
"The Onset of 2-D Grid Generated Turbulence in Flowing Soap Films,"  
Phys. Fluids 8, S7, (Sep. 1996).

# The big guess in turbulence



A.N. Kolmogorov

- Eddies spin off other eddies without losing energy.
  - Does not involve friction!
  - Hypothesis due to Richardson, Kolmogorov, ...
- Turbulence is a scale-invariant cascade of energy.

# Scale-invariant cascade

## Biology

**Swift & de Morgan poem**

Great fleas have little fleas

Upon their backs to bite 'em,

And little fleas have lesser  
fleas,

And so *ad infinitum*.

# Scale-invariant cascade

## Biology

### Swift & de Morgan poem

Great fleas have little fleas

Upon their backs to bite 'em,

And little fleas have lesser  
fleas,

And so *ad infinitum*.

## Turbulence

### Richardson poem

Big whorls have little whorls

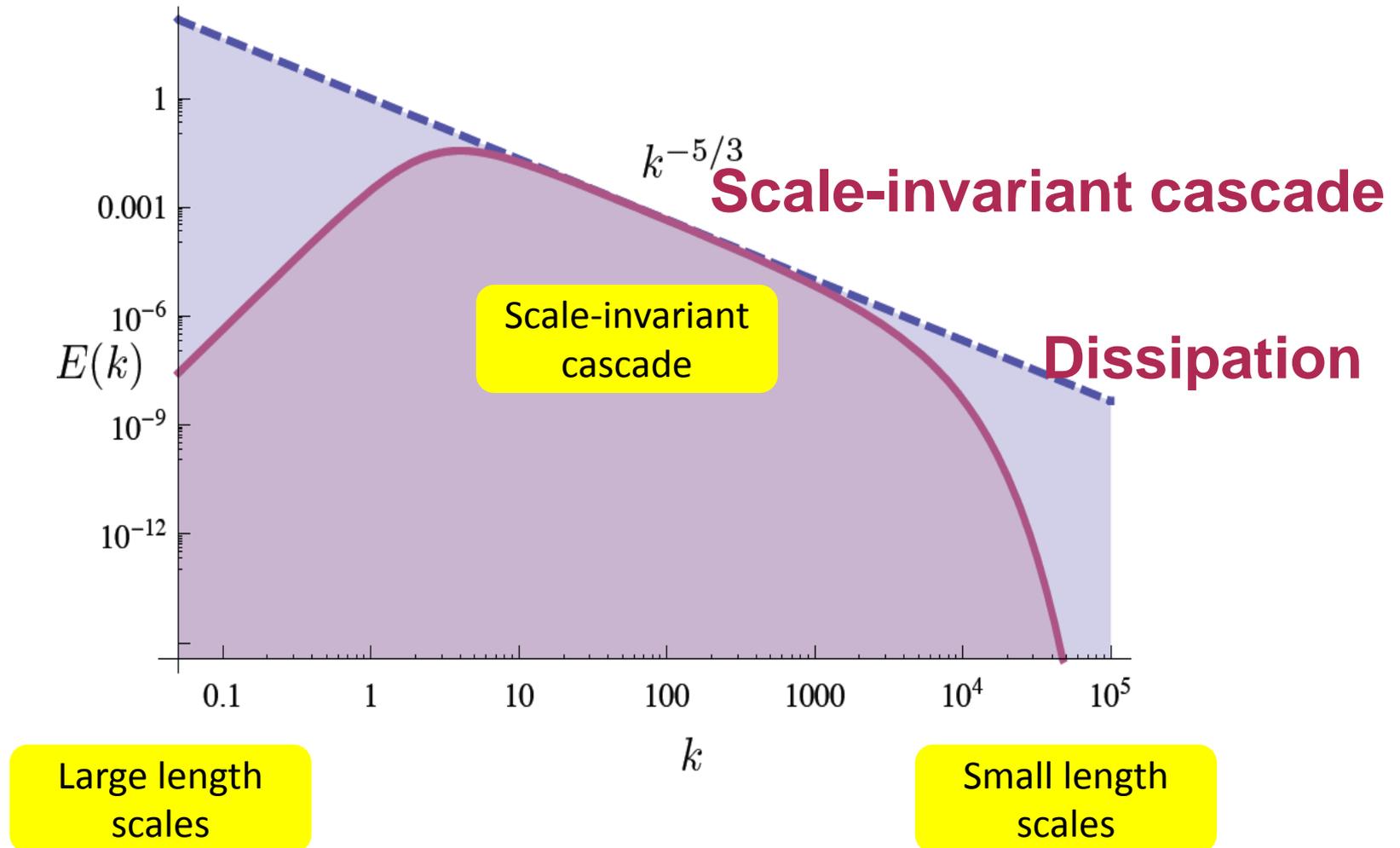
That feed on their velocity,

And little whorls have lesser  
whorls

And so on to viscosity.

# The big guess in turbulence

Energy input



# Turbulent cascades



**3D forward cascade**



**2D inverse cascade**

**Energy flows to small scales**

**Energy flows to large scales**

# Turbulent cascades



**3D forward cascade**

$$E(k) \propto \epsilon^{2/3} k^{-5/3}$$



**2D inverse cascade**

$$E(k) \propto \epsilon^{2/3} k^{-5/3}$$

# Turbulent cascades



**2D forward cascade**



**2D inverse cascade**

**Vorticity flows to small scales**

**Energy flows to large scales**

# Turbulent cascades



**2D forward cascade**

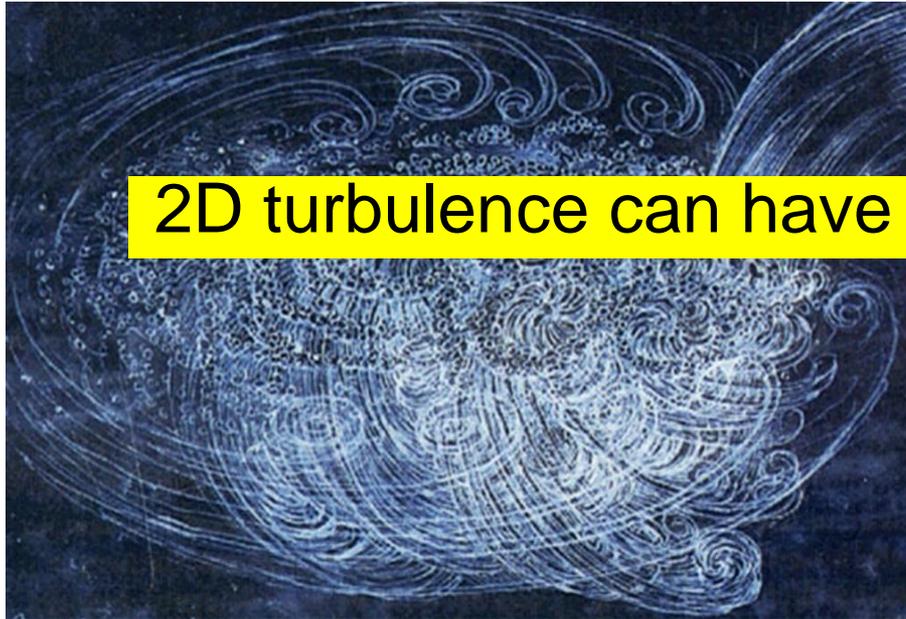
$$E(k) \propto \lambda^{2/3} k^{-3}$$



**2D inverse cascade**

$$E(k) \propto \epsilon^{2/3} k^{-5/3}$$

# Turbulent cascades



2D turbulence can have two separate cascades



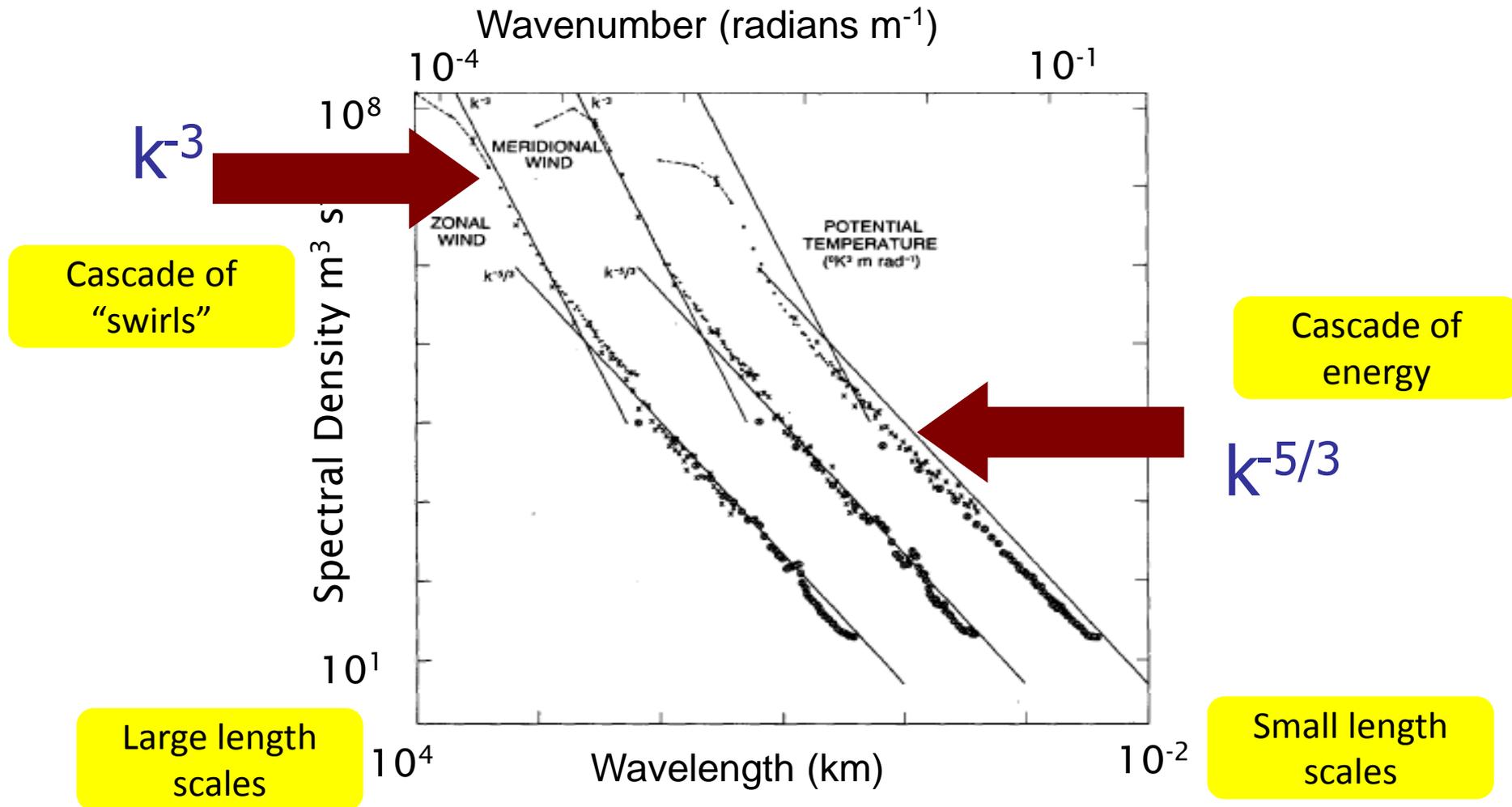
**2D forward cascade**

$$E(k) \propto \lambda^{2/3} k^{-3}$$

**2D inverse cascade**

$$E(k) \propto \epsilon^{2/3} k^{-5/3}$$

# Scale-invariant cascades in the atmosphere



G. D. Nastrom and K. S. Gage, "A Climatology of Atmospheric Wavenumber Spectra of Wind and Temperature Observed by Commercial Aircraft", Jour. Atmos. Sci. vol 42, 1985 p953

# The story so far ...

**The flow tells the flow how to change the flow**

**We can guess the outcome: out of complexity comes  
a simple picture of scale-invariant cascades**

**... sounds nice**

**So is it solved? What's new?**

# Reynolds & Turbulence



- Reynolds (1883) first to study turbulence systematically

Sees “Flashes” of turbulence

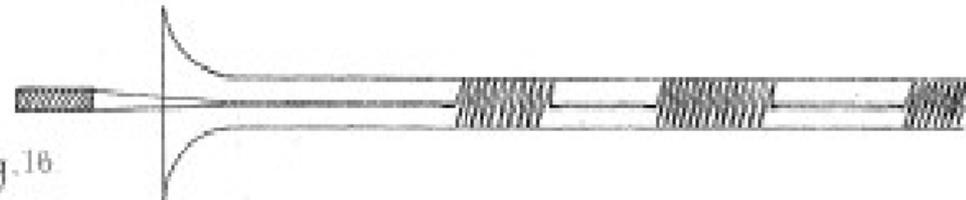
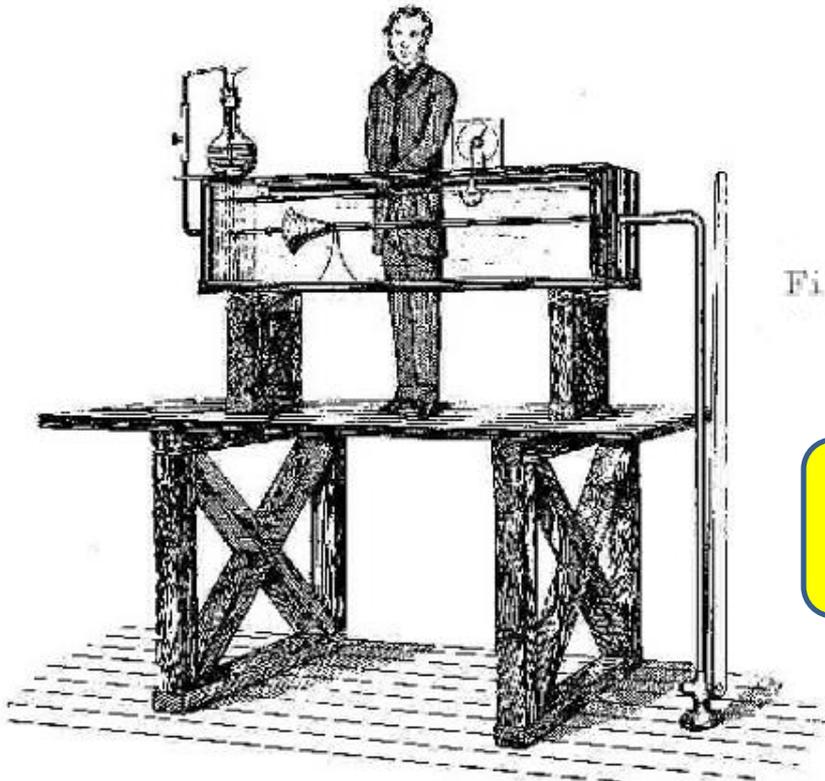


Fig. 16

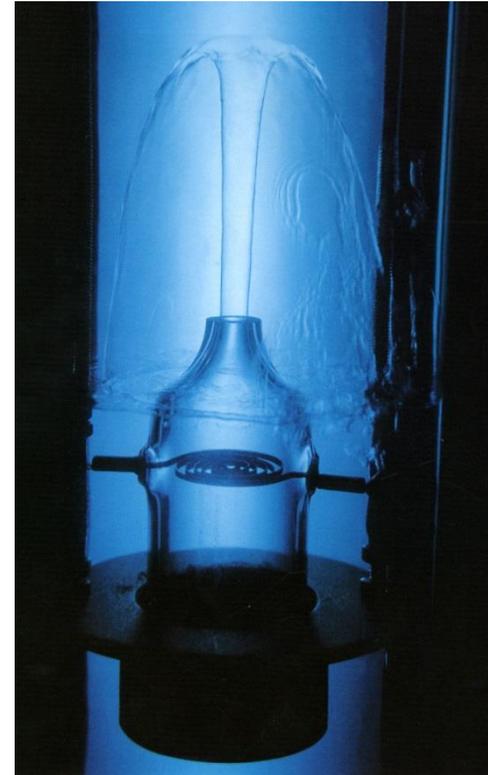
Introduces “Reynolds number”  
 $Re = \text{flow speed} * \text{size} / (\text{viscosity} / \text{density})$

# Superfluids



Classical fluid:  
velocity can “swirl”  
as slow as you like

<http://www.honeyassociation.com/>



Quantum fluid: velocity  
cannot “swirl” slower  
than a certain amount

<http://pitp.physics.ubc.ca/archives/CWSS/showcase/topics/He4-fountain.jpg>

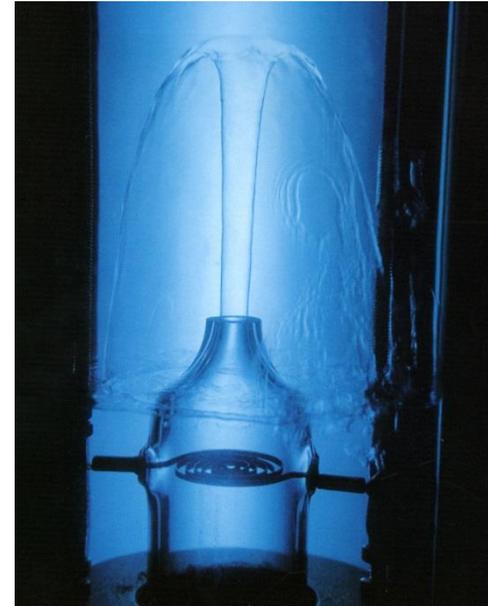
# Superfluids



**LAMINAR:**

Re much less than 1

Classical fluid:  
velocity can “swirl”  
as slow as you like



**TURBULENT:**

Re  $\sim$  1 million

Quantum fluid: velocity  
cannot “swirl” slower  
than a certain amount

Scale-invariant cascade in  
fully-developed turbulence?



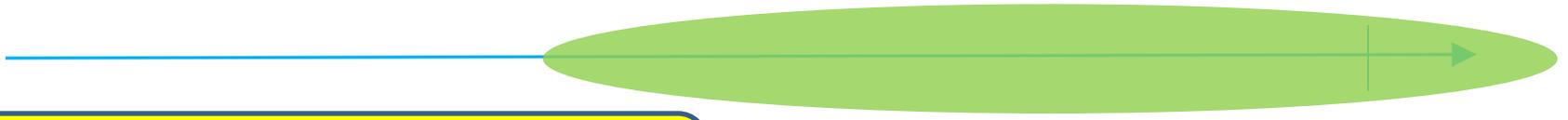
Re  $\sim$  1800 - 2400

Re  $\sim$  infinity

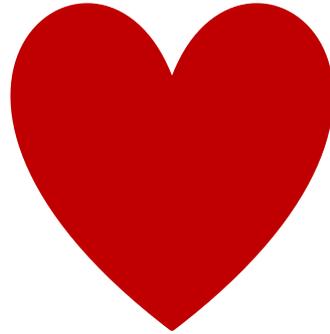
Scale-invariant cascade in fully-developed turbulence?

Re ~ 1800 - 2400

Re ~ infinity



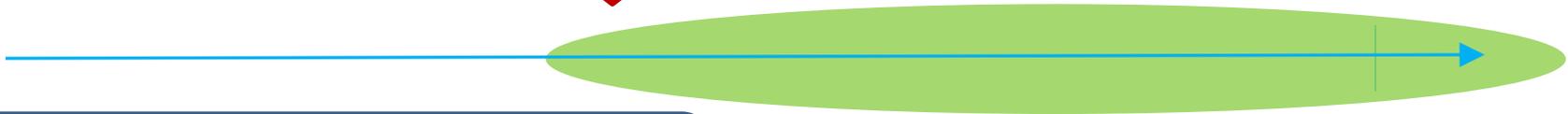
You are here



Scale-invariant cascade in fully-developed turbulence?

Re ~ 1800 - 2400

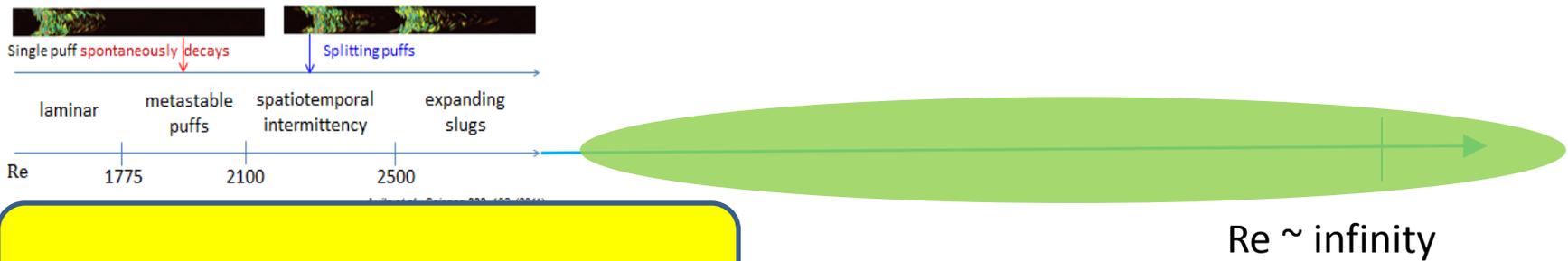
Re ~ infinity



Scale-invariance, but no cascade at laminar-turbulence transition

Scale-invariant cascade in fully-developed turbulence?

### Phase diagram of pipe flow



Re ~ 1800 - 2400

Re ~ infinity

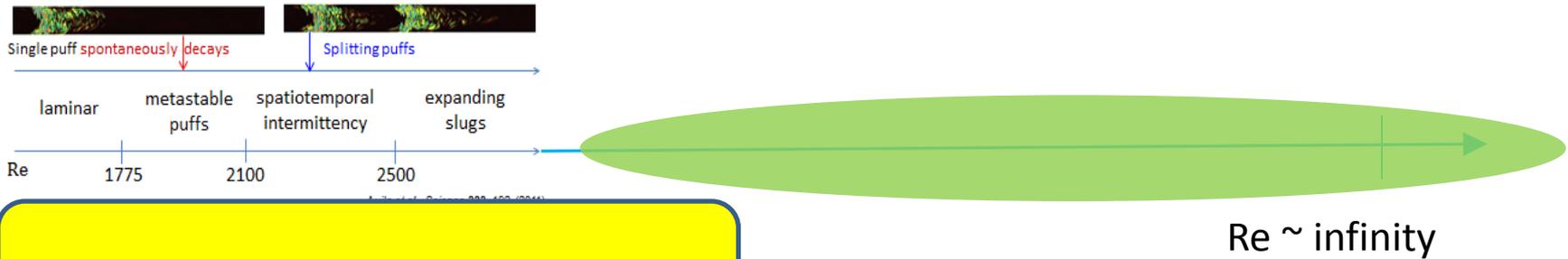
**Q. Is there universal scaling behaviour in fully-developed turbulence?**

**Q. How do fluids become turbulent as you increase their flow speed?**

Scale-invariance, but no cascade at laminar-turbulence transition

Scale-invariant cascade in fully-developed turbulence?

### Phase diagram of pipe flow



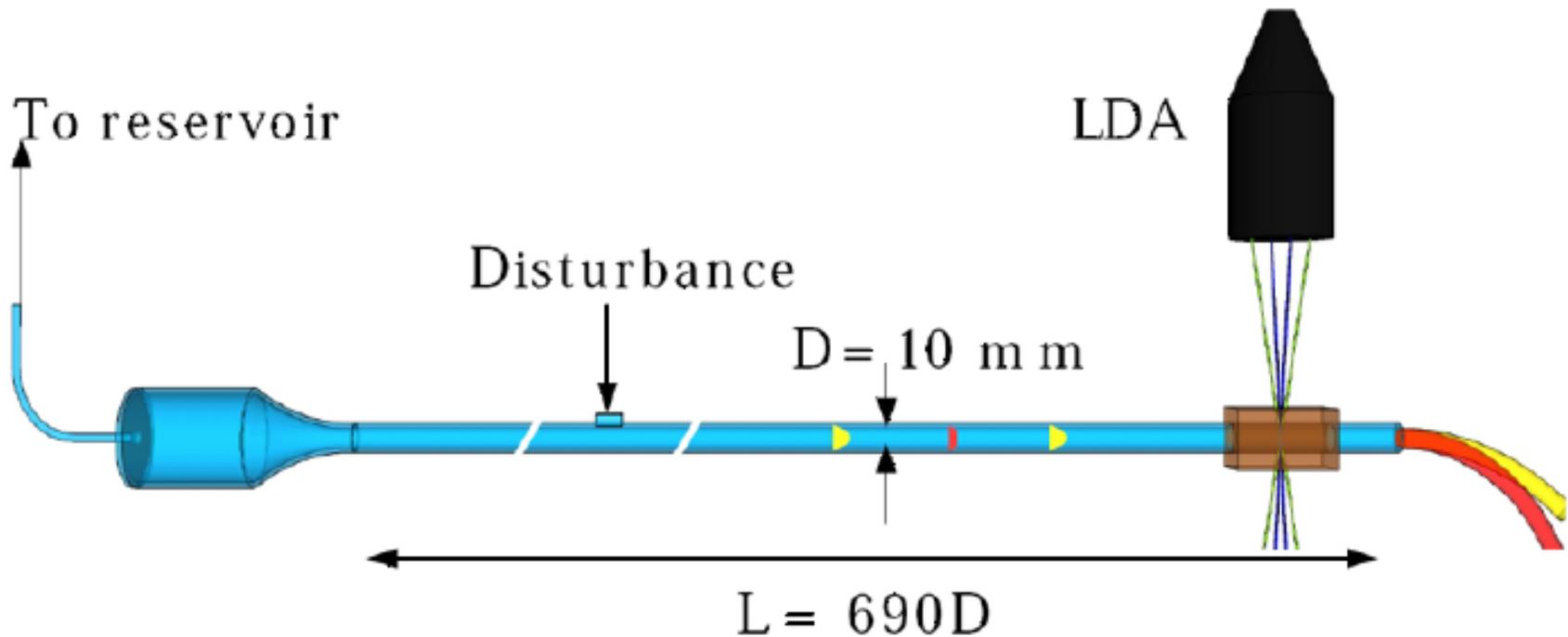
Re ~ 1800 - 2400

Q. Is there universal scaling behaviour in fully-developed turbulence?

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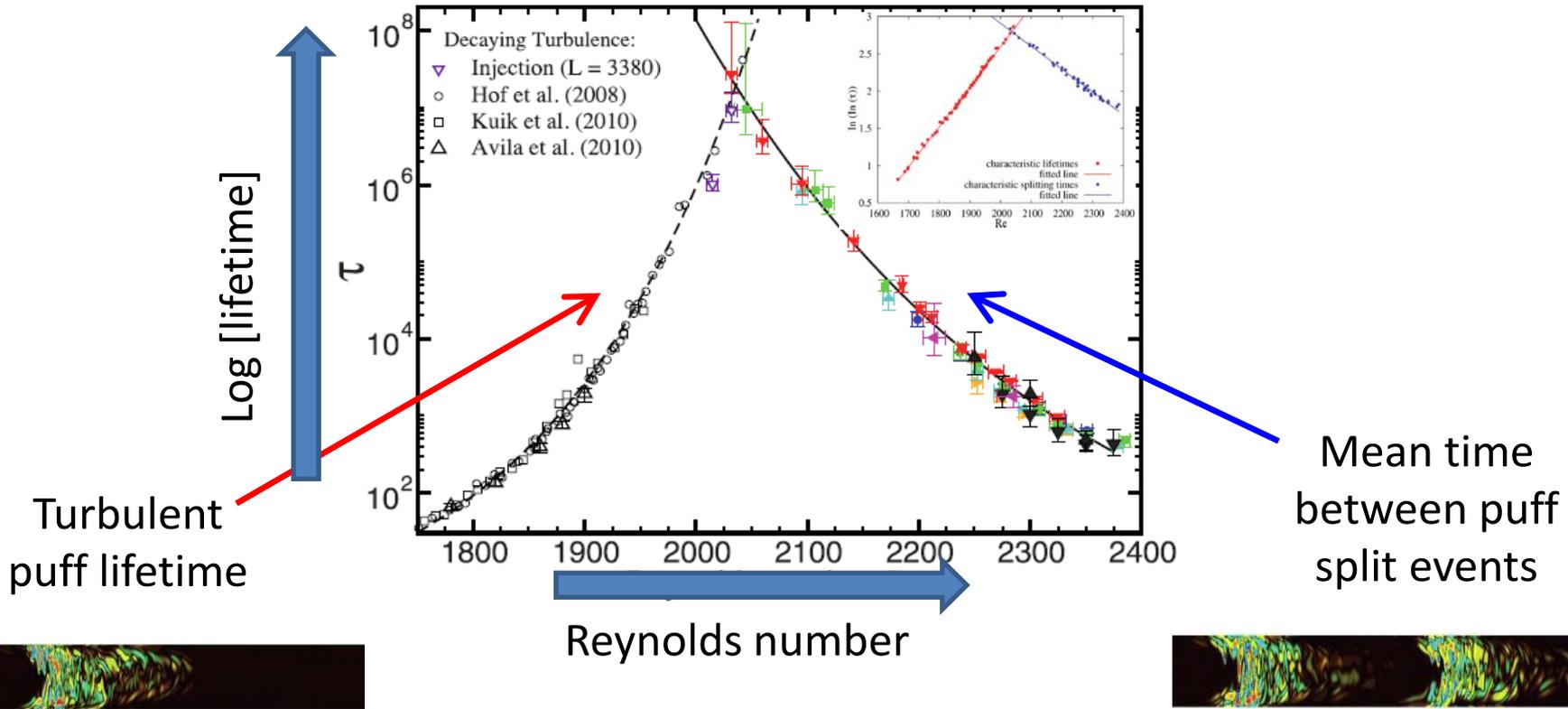
# Precision measurement of turbulent transition

Q: will a puff survive to the end of the pipe?



Many repetitions  $\rightarrow$  survival probability  $P(\text{Re}, t)$

# Fluid in a pipe near onset of turbulence



# Predator-prey ecosystem near extinction



# Predator-prey ecosystem near extinction



Prey birth rate too small  
→ Predator eats all prey and starves and goes extinct

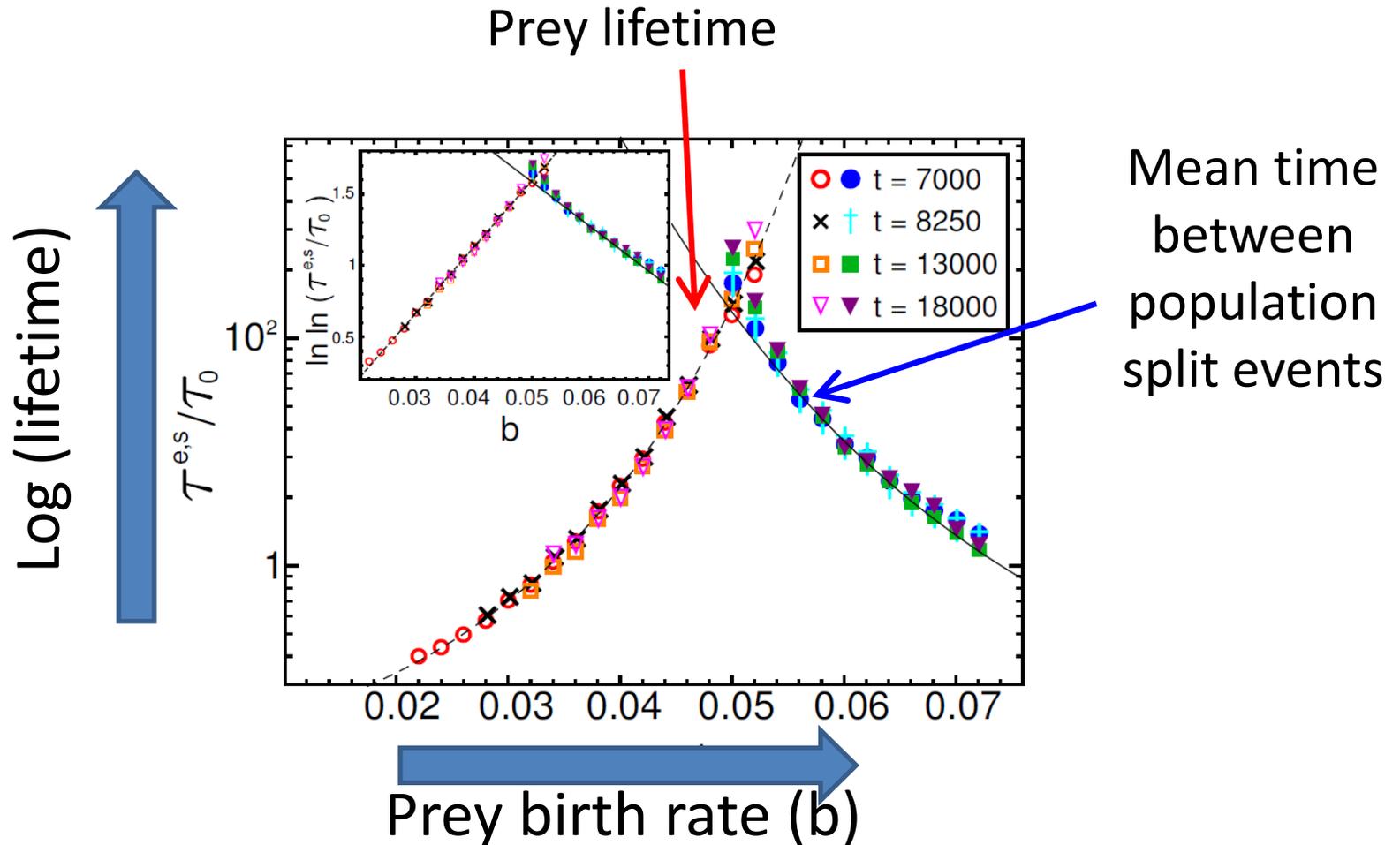
Extinction

Prey birth rate high enough  
→ Predator population supported by prey: coexistence

Coexistence

Birth rate

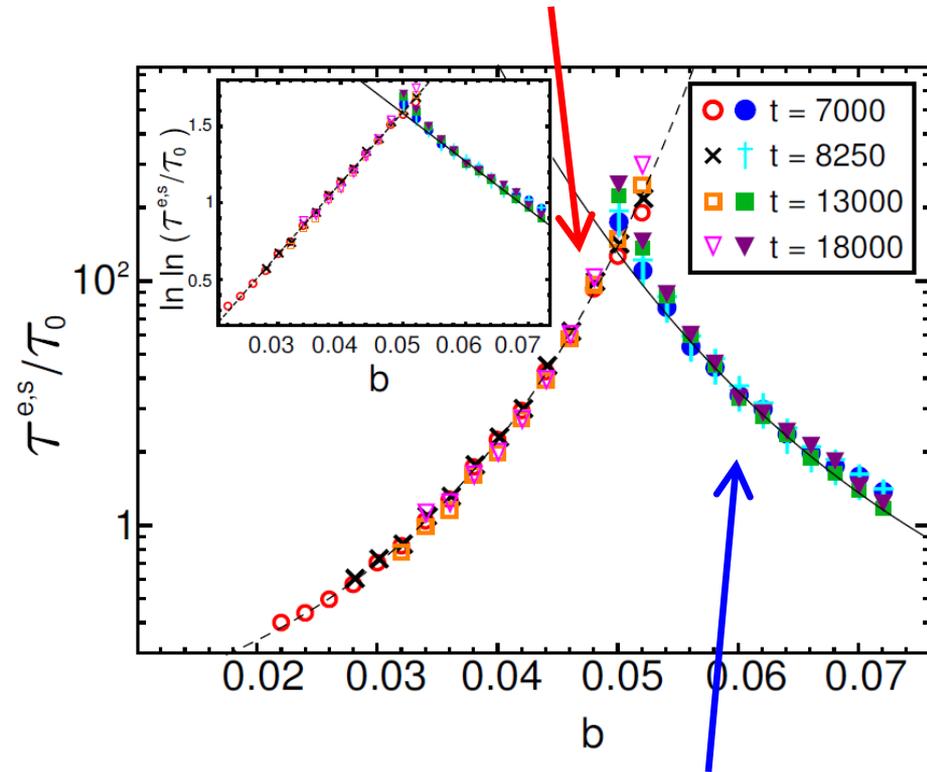
# Predator-prey ecosystem near extinction



Super-exponential scaling:  $\frac{\tau}{\tau_0} \sim \exp(\exp b)$

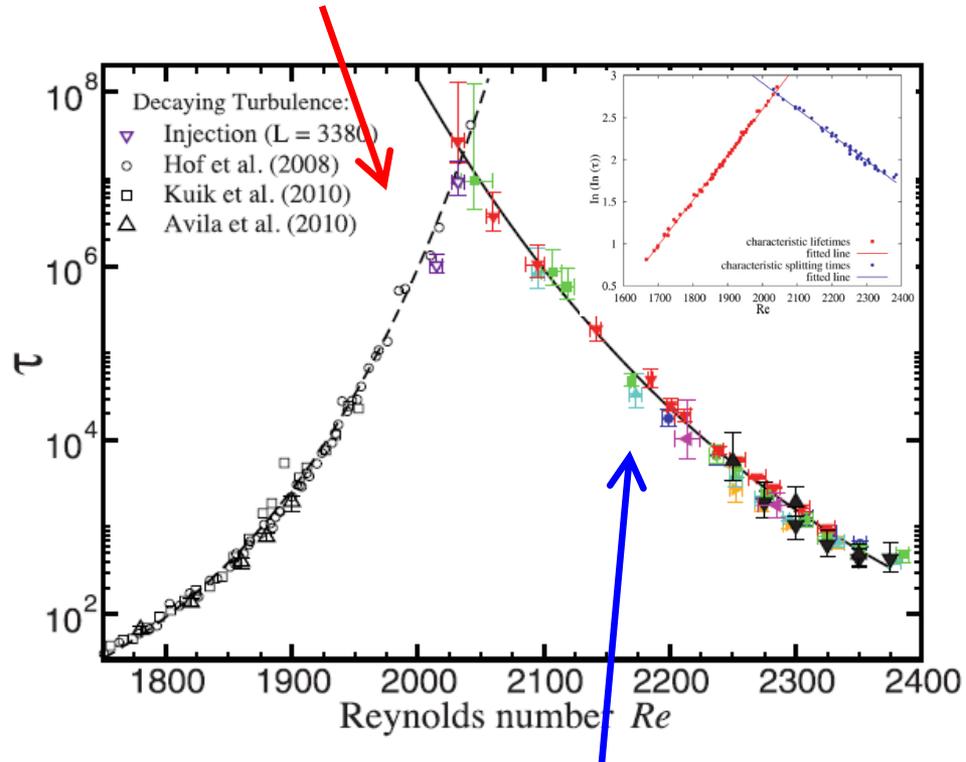
# Predator-prey vs. transitional turbulence

Prey lifetime



Mean time between population split events

Turbulent puff lifetime

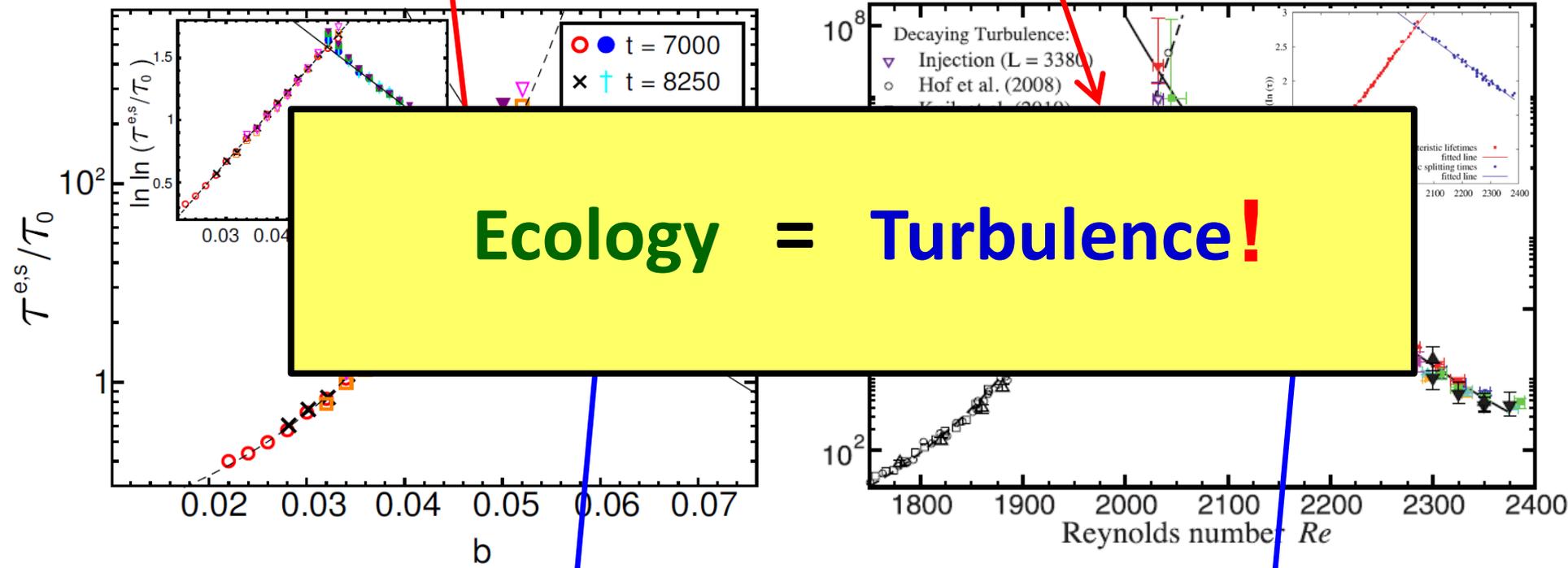


Mean time between puff split events

# Predator-prey vs. transitional turbulence

Prey lifetime

Turbulent puff lifetime



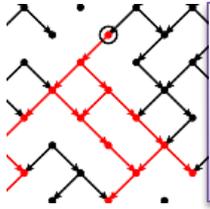
**Ecology = Turbulence!**

Mean time between population split events

Mean time between puff split events

# Turbulence transition – highly connected!

(Boffetta and Ecke, 2012)

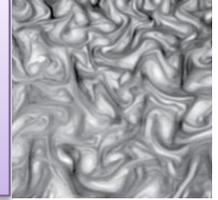


Directed Percolation

(Wikimedia Commons)

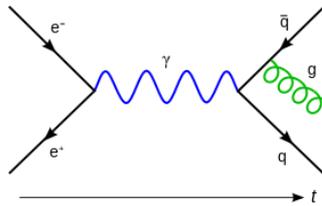
Reggeon field theory  
(Janssen, 1981)

(Classical) Turbulence



Direct Numerical Simulations  
of Navier-Stokes

High energy physics

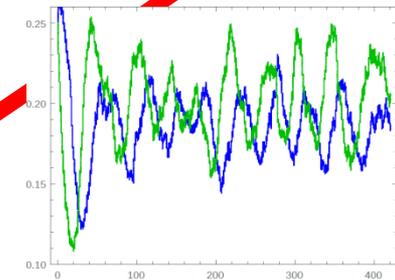


(Wikimedia Commons)

Two-fluid model

Extinction transition  
(Mobilia et al., 2007)

Ecology

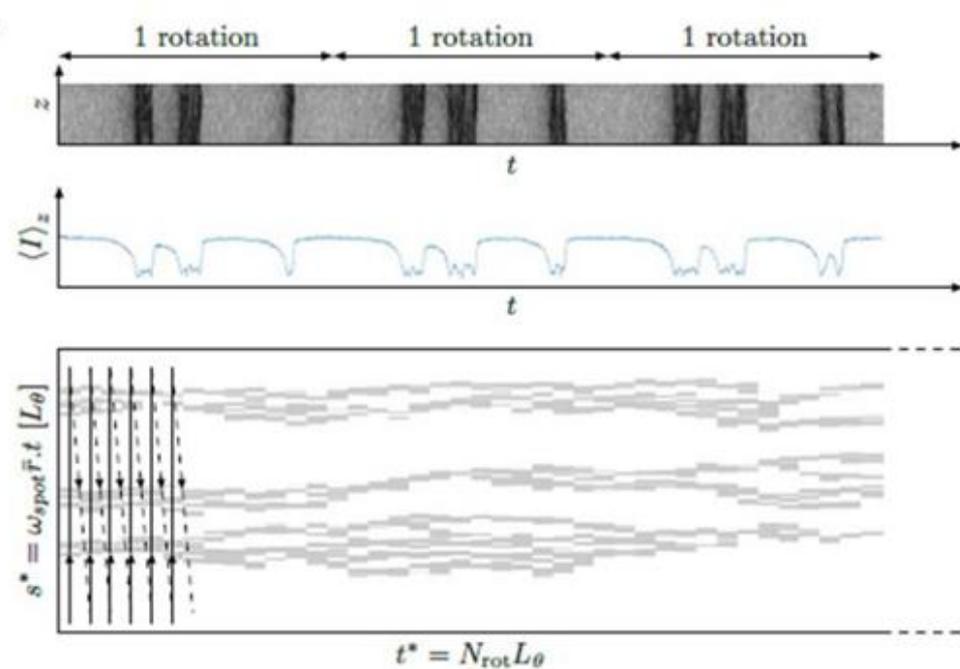
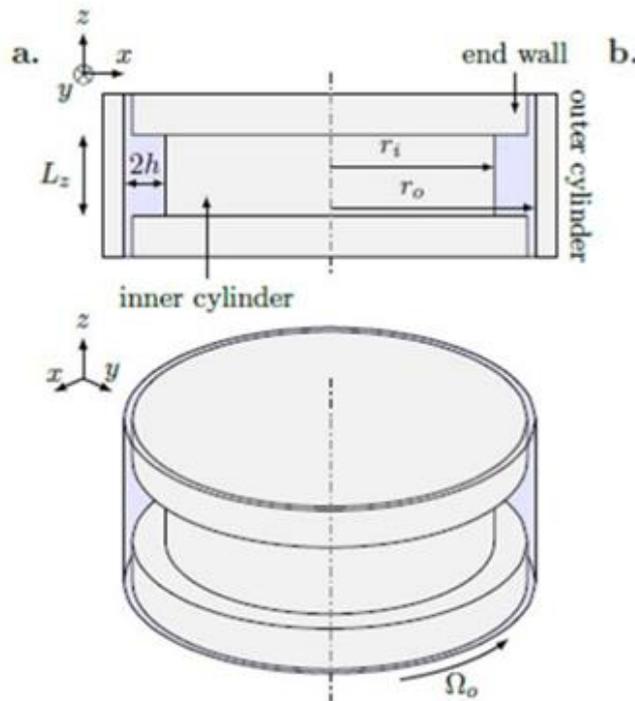


(Pearson Education, Inc., 2009)

# Turbulence and “directed percolation”

Fluid between concentric cylinders, outer one rotating

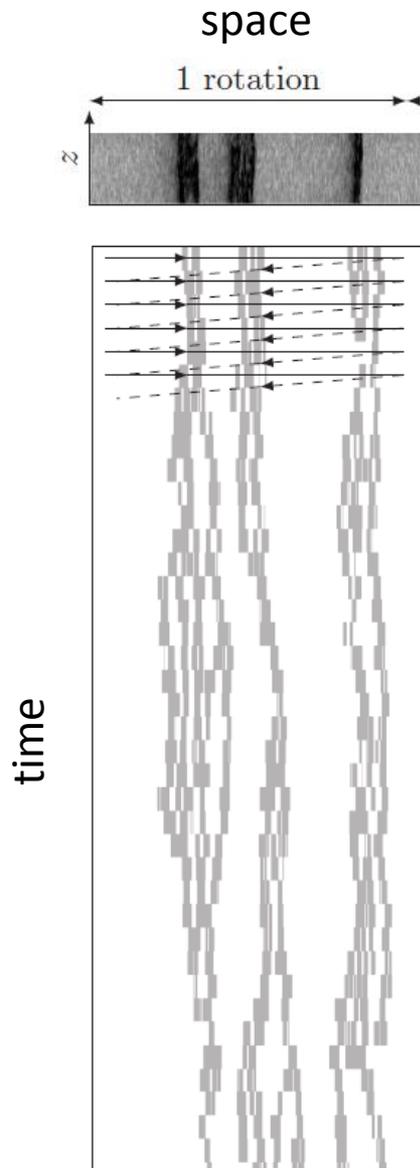
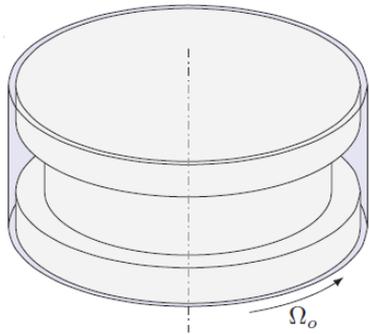
Turbulent patches



Position of turbulent patches changes in time

# Turbulence and “directed percolation”

**Couette**



space



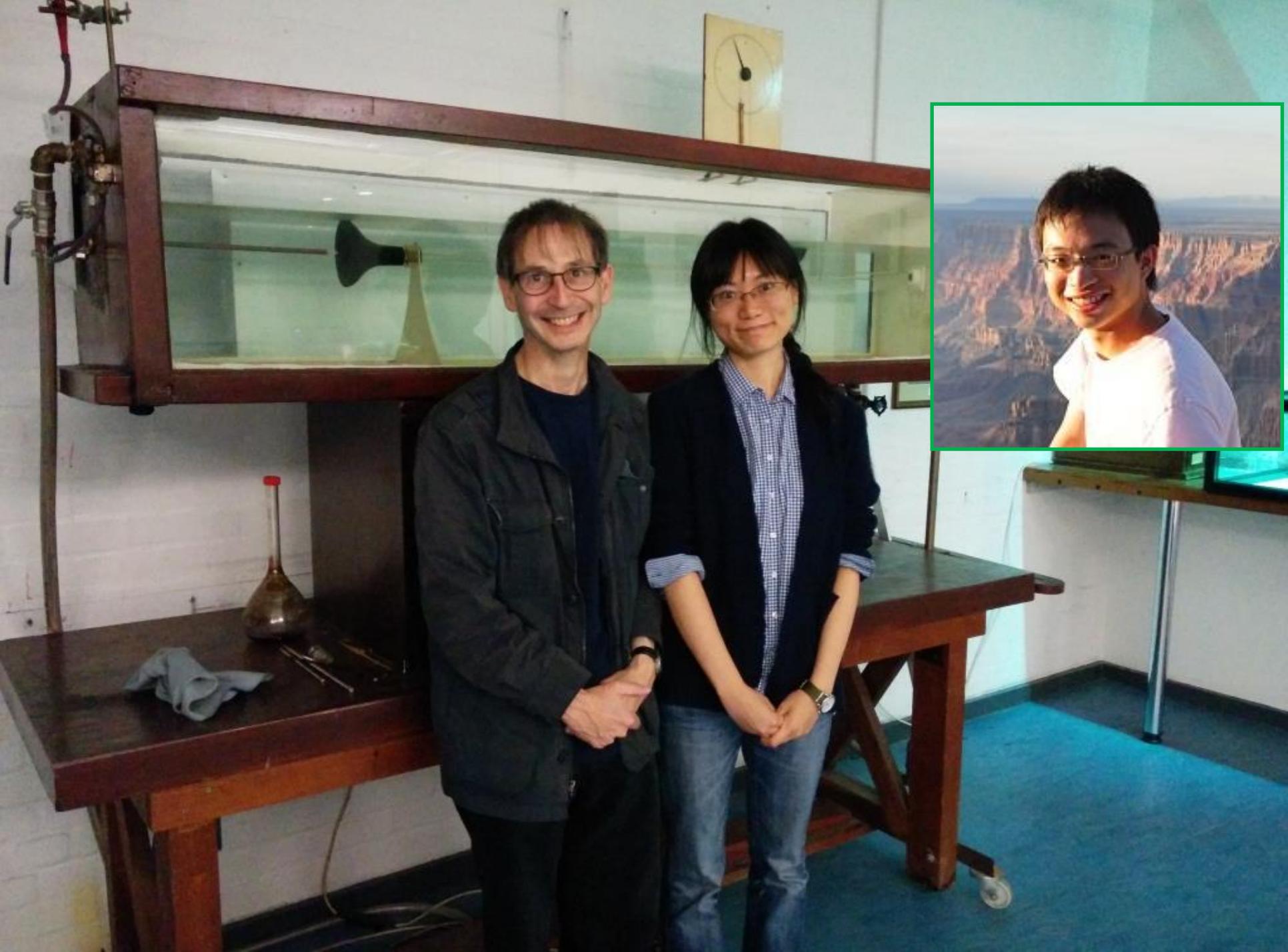
**Ecology**



time

# What did you learn today?

- Turbulence is an unpredictable complex flow with structure at a wide range of length scales
- In your body, the flow of blood is at the edge of turbulence, and crossing that transition can be fatal
- Through great ingenuity by many people we now can measure the transition very accurately and quantify its behavior.
- The transition to turbulence is mathematically like the extinction transition in an ecosystem, directed percolation and even the scattering of high energy elementary particles.



# Take-home messages

- Theoretical physics starts with a simple guess, because the exact equations are too hard to solve
  - Compare the consequences of that guess with experiment
  - And the experiments suggest new guesses etc...
- Simple models --- but not too simple --- can describe very complicated phenomena
  - Physics is successful because we only work on simple problems
  - We do not really know to what extent this approach can be used in more complex sciences like biology
- The problem of turbulence is not solved --- but we understand a lot more than we did in Feynman's day
  - Thanks to brilliant and controlled experiments
  - Healthy interactions between experimenters, theorists and computational physicists

**Turbulence is a life force. It is opportunity.  
Let's love turbulence and use it for change.**

**Lucky Numbers 34, 15, 28, 4, 19, 20**

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