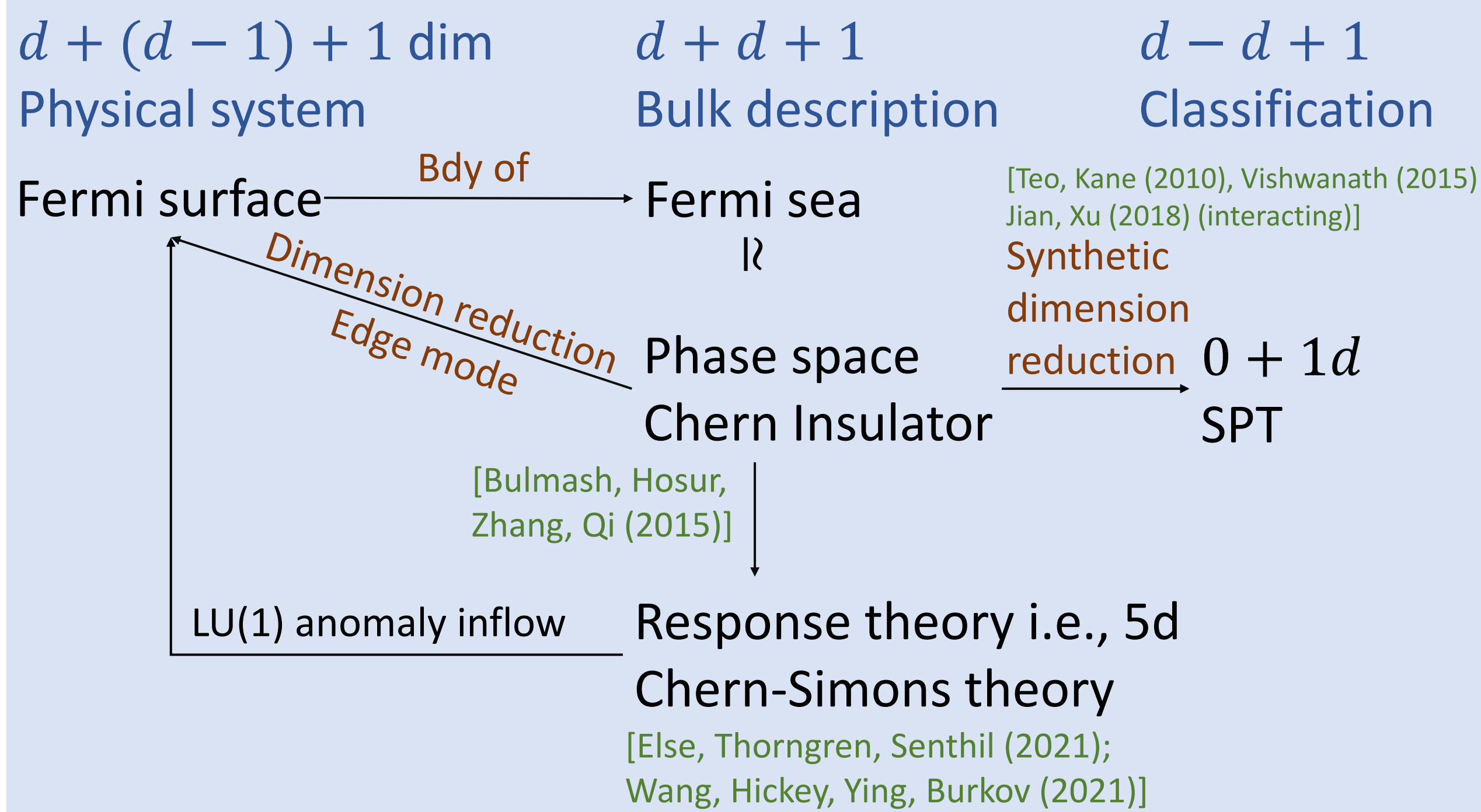


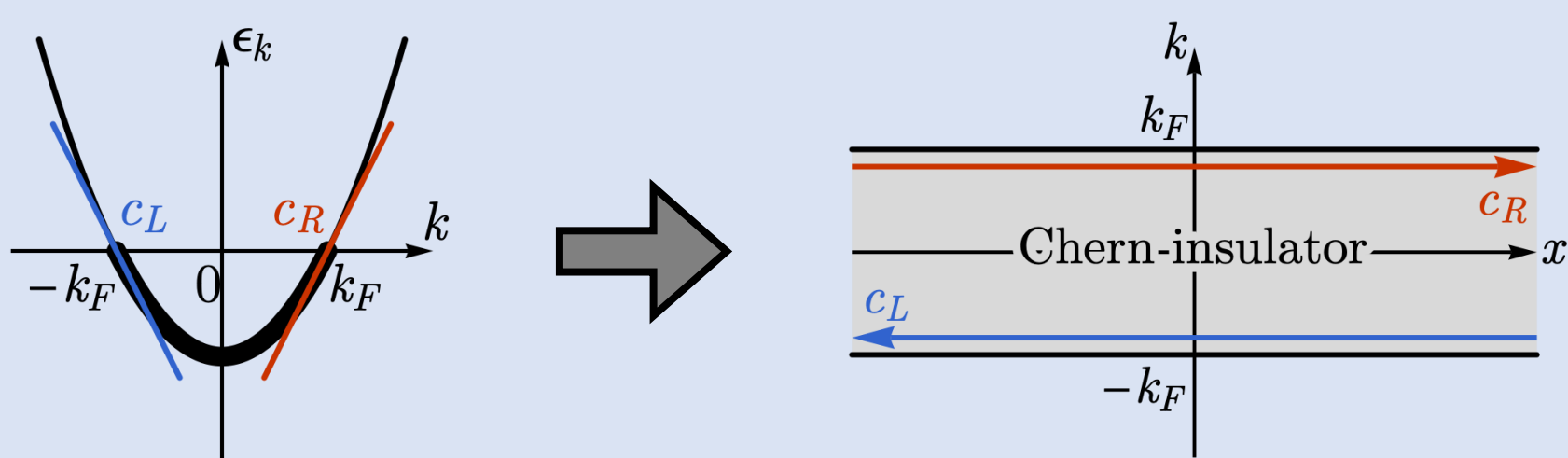
Fermi Surface Anomaly and Symmetric Mass Generation

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Outline



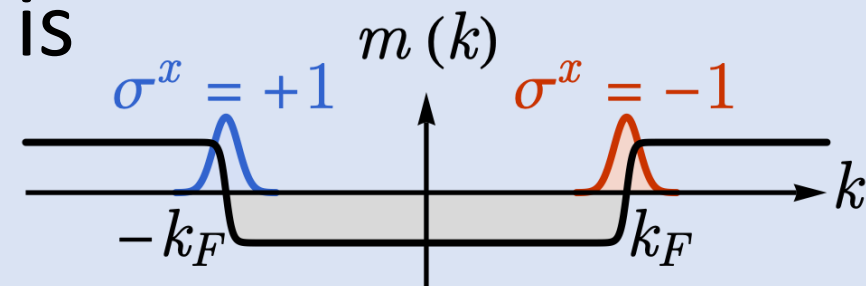
Bulk of Fermi surface \simeq Phase space Chern insulator



Given the phase space Chern insulator,

$$H = \int dx dk \psi^\dagger (i\partial_x \sigma^x + i\partial_k \sigma^y + m(x, k) \sigma^z) \psi$$

And the mass profile is



There are localized modes when mass changes sign.

$$H_L = \int dx \psi^\dagger (+i\partial_x) \psi, \quad H_R = \int dx \psi^\dagger (-i\partial_x) \psi$$

Higher dimensional phase space Chern insulator

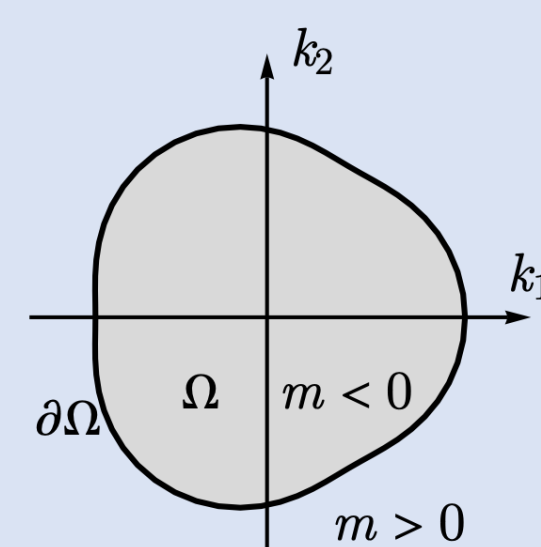
In higher dimension,

$$H_{blk} = \int d^d \mathbf{x} d^d \mathbf{k} \psi^\dagger (i\partial_x \cdot \Gamma^x + i\partial_k \Gamma^k + m(\mathbf{k}) \Gamma^0) \psi$$

Γ 's are from complex Clifford algebra $\mathcal{C}\ell_{2d+1}$.

The mass profile is

$$m(\mathbf{k}) = \begin{cases} \leq 0, \mathbf{k} \in \Omega \\ > 0, \mathbf{k} \notin \Omega \end{cases}$$



To restore the commutator $[x, k] = i$, we use,

$$x = i\partial_k \Leftrightarrow k = -i\partial_x$$

The boundary theory is

$$H_{bdy} = \int_{\partial\Omega} d\mathbf{k}_F \int d(\mathbf{n} \cdot \mathbf{x}) \psi_{\mathbf{k}_F}^\dagger i(\mathbf{n} \cdot \partial_x) \psi_{\mathbf{k}_F}$$

Bulk U(1) symmetry \rightarrow boundary emergent loop group symmetry

$$L_{\partial\Omega} U(1): \psi_{\mathbf{k}_F} \rightarrow e^{i\phi(\mathbf{k}_F)} \psi_{\mathbf{k}_F}, \quad \forall \mathbf{k}_F \in \partial\Omega$$

Response theory

Like the ordinary Chern insulator, the theory can minimally couple to the background gauge field, tracking the global symmetry,

$$H = \int dx dk \psi^\dagger (iD_x \sigma^x + iD_k \sigma^y + m(k) \sigma^z) \psi$$

where $D_i = i\partial_i - A$. After integrating out the fermion,

$$S = \frac{1}{4\pi} \int_{(t,x,k)} \frac{1 - \text{sgn } m(k)}{2} A \wedge dA$$

The Chern-Simons term exists only in the Chern insulator.

The gauge connection $A = A_0 dt + A_x dx + A_k dk$ has a uniform background curvature through the (x, k) plane to ensure the commutator,

$$F = dA = dx \wedge dk$$

And we can reproduce the Luttinger theorem,

$$v = \frac{\delta S}{\delta A_0} = \int_{-k_F}^{k_F} \frac{F_{xk}}{2\pi} = \frac{\text{vol } \Omega}{2\pi}$$

For codimension-1 Fermi surfaces

Fermi surface d	Fermi sea d	Phase spacetime	Bulk CS theory
0	1	1+1+1=3	$A \wedge dA$
1	2	2+2+1=5	$A \wedge dA \wedge dA$
2	3	3+3+1=7	$A \wedge dA \wedge dA \wedge dA$

with background flux through every (x_i, k_i) -plane.

Classification – 0+1d SPT

Classification of $(d + d + 1)D$ phase space Chern insulator = classification of $(0 + 1)D$ topological insulator.

- Start with the bulk phase space Chern insulator,

$$H_{blk} = \int d^d \mathbf{x} d^d \mathbf{k} \psi^\dagger (i\partial_x \cdot \Gamma^x + i\partial_k \Gamma^k + m \Gamma^0) \psi$$

- Using the operator relation, $i\partial_k = x$,

$$H_{blk} = \int d^d \mathbf{x} \psi^\dagger (i\partial_x \cdot \Gamma^x + x \cdot \Gamma^k + m \Gamma^0) \psi$$

- After projecting out every pair of x_i, k_i , $H_{blk} = m \psi^\dagger \psi$.
- The bulk state is trivial if the gapless fermion modes at $m = 0$ can be symmetrically gapped by interaction.
- For \mathbb{Z}_4 -symmetric fermions, trivialization can be achieved by $g(\psi_1 \psi_2 \psi_3 \psi_4 + h.c.)$ at multiplicity 4 $\Rightarrow \mathbb{Z}_4$ classified!

Conclusion

Classification of Fermi surface anomaly

Codimension- p Fermi surface anomaly of symmetry group G is classified by G symmetric interacting fermionic SPT states in p -dimensional spacetime.

Fermi surface symmetric mass generation can happen when the Fermi surface anomaly vanishes.

$$v = \sum_{\alpha} k_{\alpha} \frac{\text{vol } \Omega_{\alpha}}{(2\pi)^d} = 0 \pmod{1}$$

- Potential applications to the pseudo-gap physics
- Spectral signatures: Green's function zeros on the Fermi surface
- Functional renormalization group of FS SMG models
- Matching Fermi surface anomaly with topological order

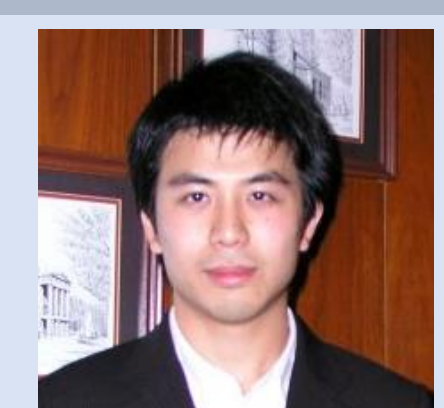
Collaborators and References



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