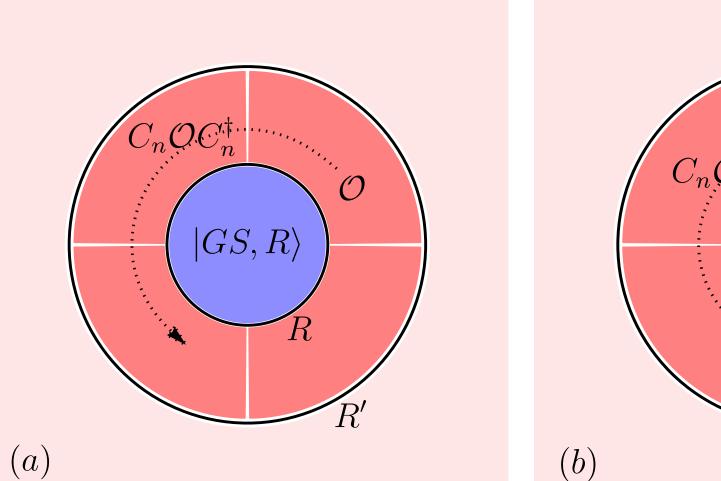
Many-Body Real Space Invariants for Interacting 2D Topology

Jonah Herzog-Arbeitman¹, Andrei Bernevig^{1,2,3}, Zhi-da Song⁴ Princeton University¹, Donostia International Physics Center², IKERBASQUE³, Peking University⁴

1. Many-Body Real Space Invariants

- Topological indices are global in momentum space \rightarrow local in real space, act as "Noether charges" \bullet
- RSIs as quantum numbers: define symmetry operators on OBC's s.t. spectrum independent of OBC's
- Space group operators give local angular momentum and local charge as RSIs (n even, spin-less): \bullet $e^{i\frac{\pi}{n}\hat{N}}C_n |GS\rangle = e^{i\frac{\pi}{n}\Delta_1} |GS\rangle, \ C_n^2 |GS\rangle = e^{i\frac{2\pi}{n/2}\Delta_2} |GS\rangle, \quad \Delta_1, \Delta_2 \in \mathbb{Z}_{2n} \times \mathbb{Z}_{n/2}$
- Proof makes use of a "Many-body atomic limit" of arbitrarily strong interactions, but onsite-local: \bullet







(b)

$$GS, R' \rangle = \prod_{i=0}^{n-1} \mathcal{O}_i |GS, R\rangle, \quad \mathcal{O}_i = \prod_{\mathbf{R} + \mathbf{r} \in \mathcal{D}} C_n^i \mathcal{O}_{\mathbf{R}, \mathbf{r}} C_n^{\dagger i}$$

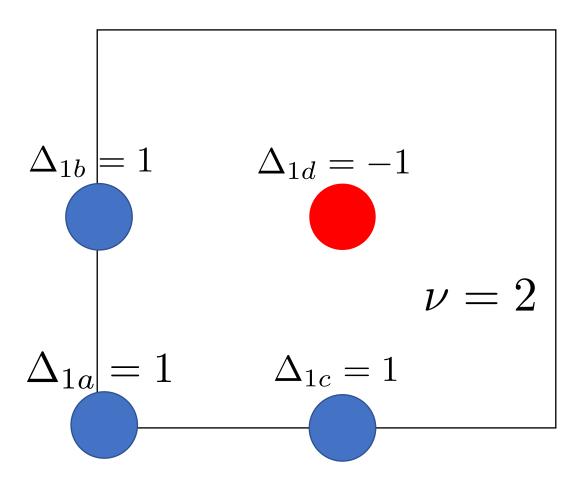
2. Many-Body Fragile Topology

- Definition: Many-body Fragile Topology iff adiabatically disconnected from many-body atomic state at fixed particle number, but connected at higher particle number. \bullet
- Interplay between U(1) charge and space group G. Topological indices are obstructions in the form of inequalities.
- RSIs form an abelian group. In a fragile phase, there exists n such that $\Delta = \Delta_{\nu+n} \Delta_n$ defined by coupling trivial states. \bullet
- Intuition: nonzero RSIs (implying local angular momentum or charge) require local operators to carry them. Formally,

 $\nu \geq \sum N_{\mathbf{x}} \geq \sum |\Delta_{\mathbf{x}}|$ holds in all atomic limit states. Violating this inequality presents an obstruction. x = 1a, 1b, 1c, 1dx = 1a, 1b, 1c, 1d

3. Weak Coupling Limit, Reduction of Single-Particle Topology, and Model Examples

Understand RSIs at weak-coupling by acting on product states \bullet



 $C_4 \mathcal{O} C_4^{\dagger}$

 $C_2 \mathcal{O} C_2^{\dagger}$

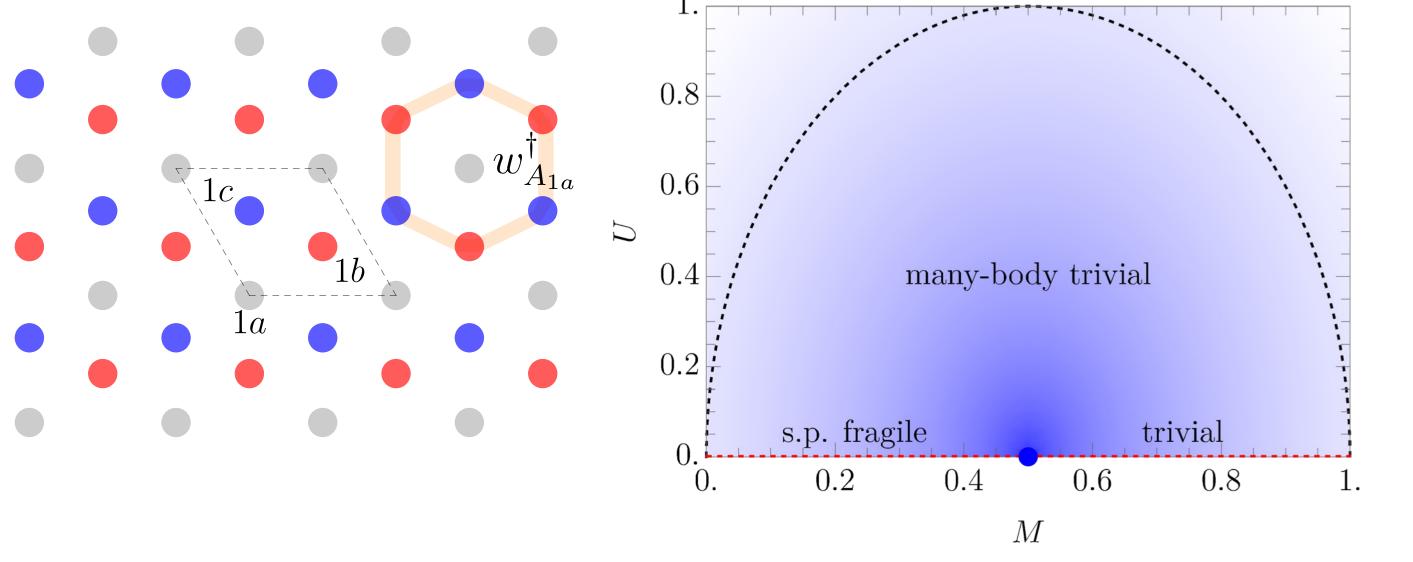
(e)

 $e^{i\frac{\pi}{2}\hat{N}}C_2 |GS\rangle = e^{i\frac{\pi}{2}(m(A) + m(B)) + i\pi m(B)} |GS\rangle$ $\Delta = m(A) - m(B) \mod 4$

- Interpretation from single-particle RSI $\Delta = \delta$ $\mod 4$. \bullet
- Momentum space computation at weak-coupling and reduction of \bullet single-particle affine-monoid structure.
- Some single-particle fragile phases are many-body fragile. Some are not. Example:

4. Many-Body Stable Topology and Interaction-Enabled Topological Phases

- Proof of many-body RSIs in atomic and fragile states relied on a gap for open boundary conditions. No such gap in stable topological phases!
- Proposal: define <u>global</u> many-body RSIs as the quantum numbers of the RSI operators on PBCs (gapped groundstate).
- We prove (1) all global RSIs (defined above) vanish in atomic and fragile states, and (2) Chern insulators have non-zero global RSIs
- Classification suggests strongly correlated topological states that do not exist at weak coupling. Ex: Chern insulator with nonzero momentum.
- 4. Topological Response Theory: RSIs as Wen-Zee Coefficients



- Topological insulators are gapped states with no low-energy electronic excitations. Their response to external probes is described by TQFTs. \bullet
- Recent work identified Wen-Zee terms that describe the response to geometric/curvature perturbations $d\omega$ (defects in spatial symmetries). \bullet





 \mathcal{O}

 $C_4^{\dagger} \mathcal{O} C_4$

 $\mathcal{O}_{\mathbf{x}}$