

Moiré atoms, Wigner molecules, and emergent Kagome lattice

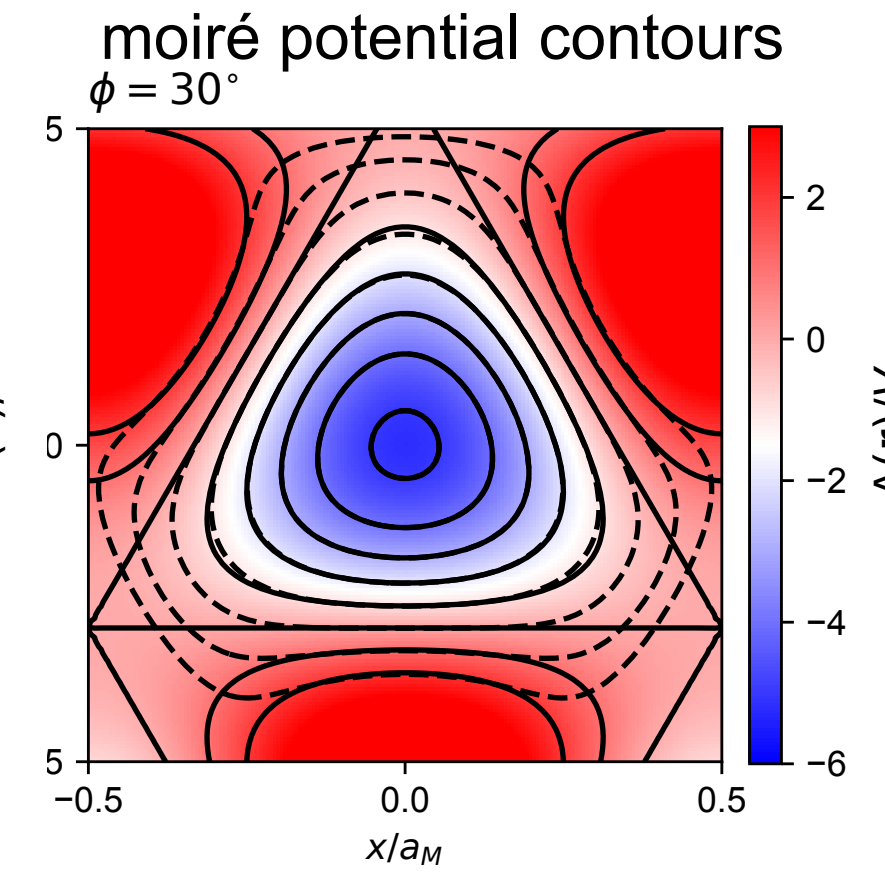
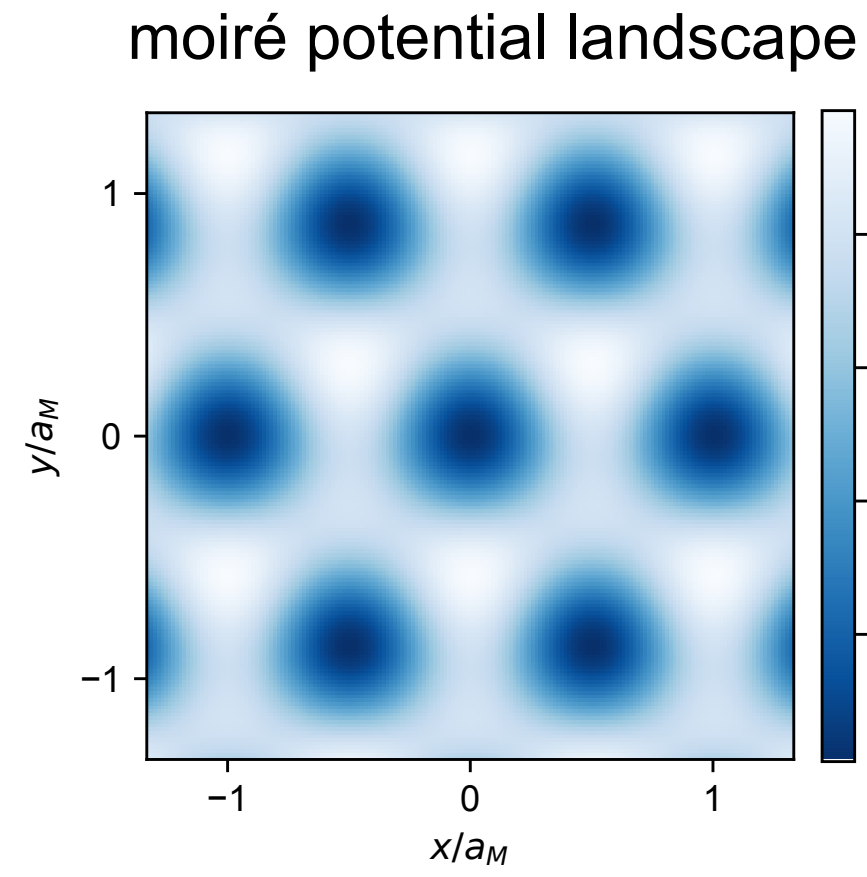
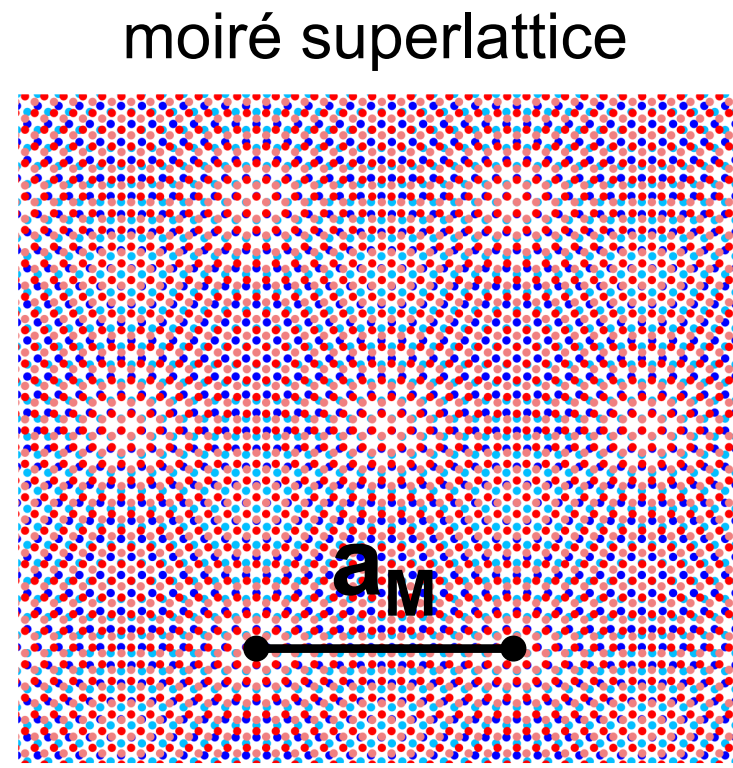
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Q: What happens in semiconductor moiré superlattices at higher filling factors ≥ 1 per potential minimum where interaction and moiré potential energies are frustrated?

Model: 2DEG + smooth periodic potential

$$H = \sum_i \frac{p_i^2}{2m} + \Delta(r_i) + \sum_{i \neq j} \frac{e^2}{\epsilon |r_i - r_j|}$$

$$\Delta(r) = -2V \sum_{i=1,3,5} \cos(g_i \cdot r + \phi)$$



“moiré atom” = local minimum of moiré potential
= parabolic quantum dot + trigonal moiré crystal field

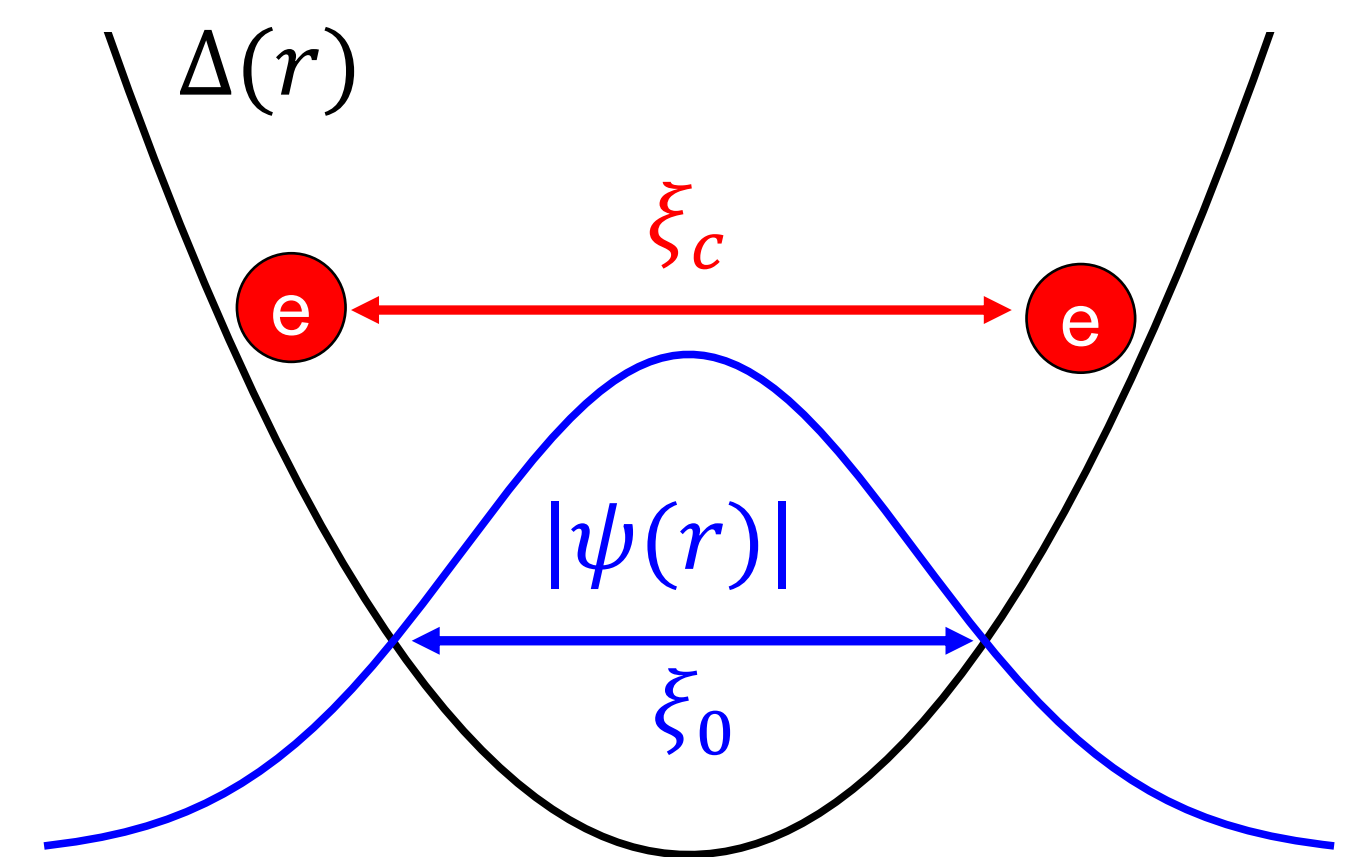
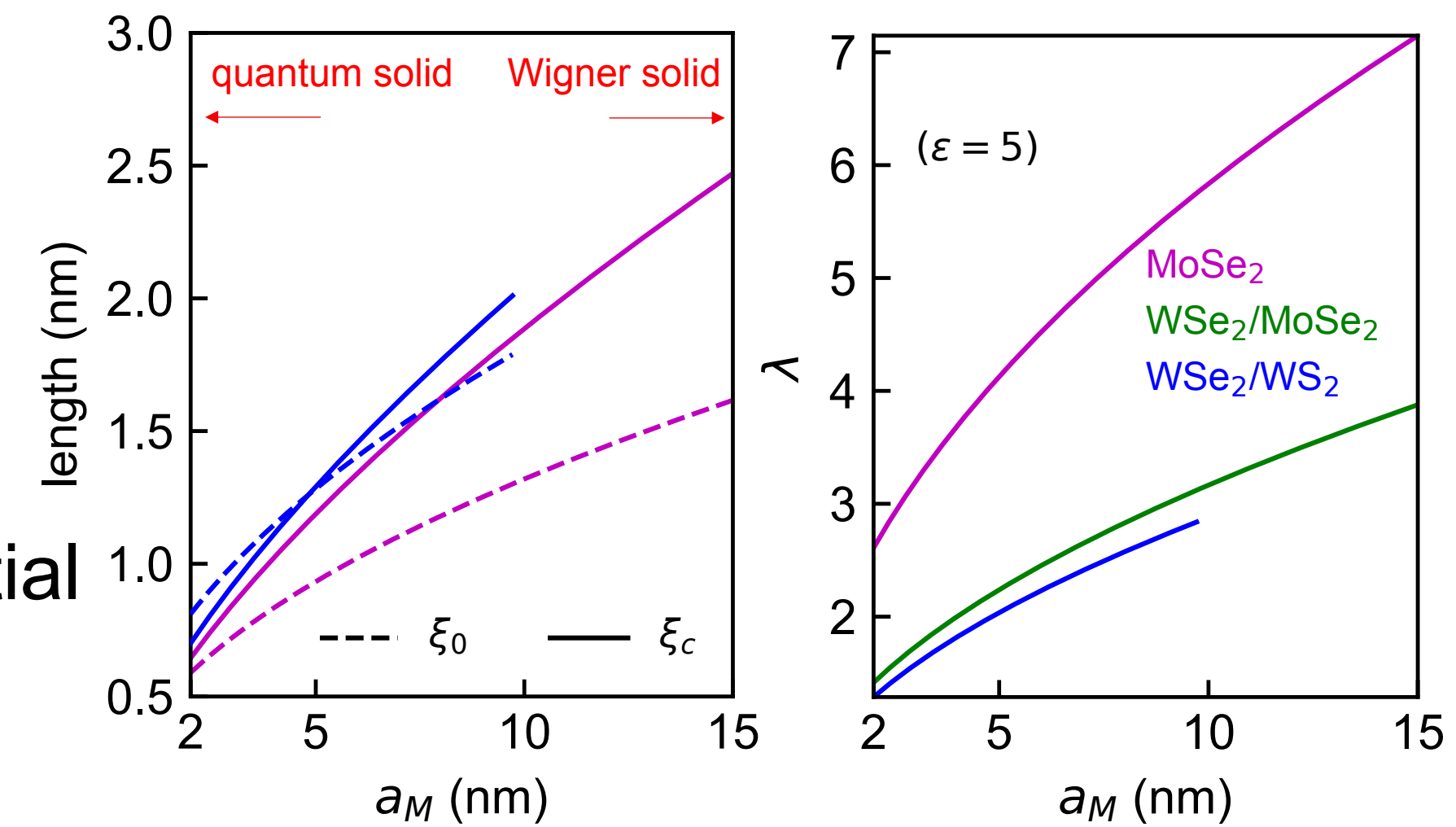
$$\Delta(r) = \frac{1}{2}kr^2 + c_3 \sin(3\theta)r^3 + \dots$$

3 key length scales:

- ξ_0 : radius of one-electron **quantum** oscillator wavefunction
- ξ_c : radius of two **classical** point charges in parabolic potential with Coulomb repulsion
- a_M : distance between moiré atoms

Atomic coupling constant: $\lambda = \frac{\text{coulomb energy}}{\text{bandgap}} = \frac{e^2/(\epsilon\xi_0)}{\hbar\omega} = (\xi_c/\xi_0)^3$

$$\xi_0 = \left[\frac{\hbar^2}{mk} \right]^{1/4} \propto a_M^{1/2} \quad \xi_c = \left[\frac{e^2}{2\epsilon k} \right]^{1/3} \propto a_M^{2/3}$$



Hierarchy of length scales $\xi_0 \ll \xi_c \ll a_M$ realized in limit $a_M \rightarrow \infty$

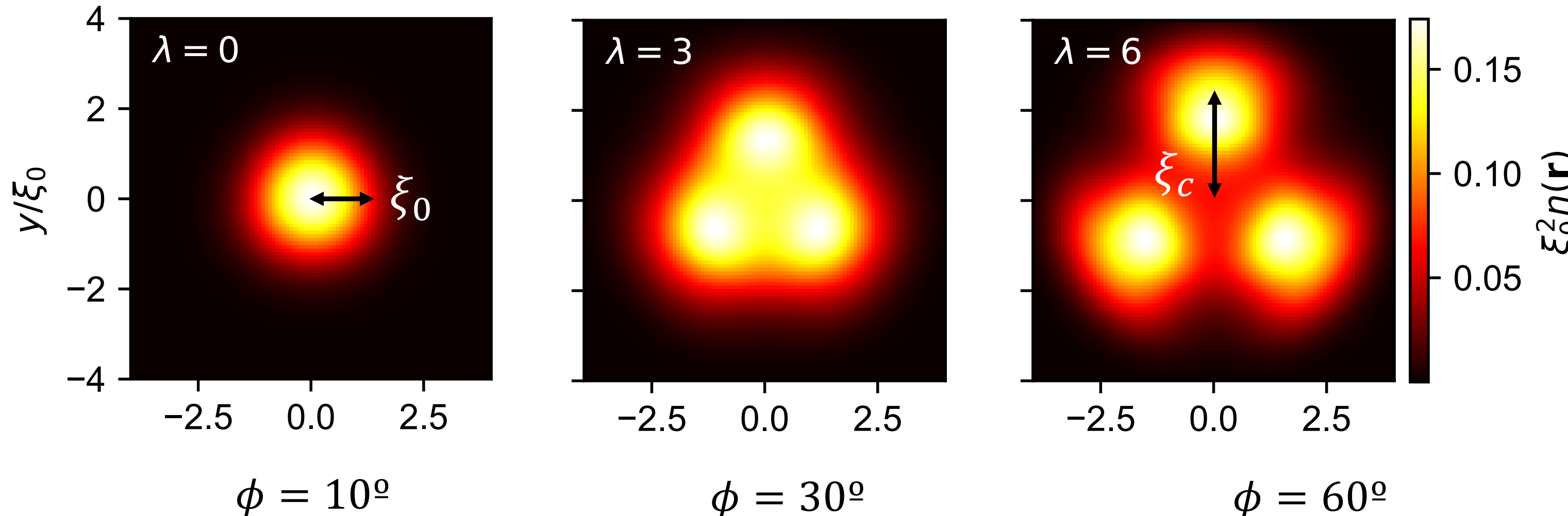
Atomic limit: $\xi_0, \xi_c \ll a_M$

Classical limit: $\xi_0 \ll \xi_c, a_M$

Exact diagonalization of isolated moiré atom

- Evolution from “Schrödinger atom” to “Wigner molecule”
- Wigner molecules are observable with STM

$\xi_c < \xi_0 \longrightarrow$ increasing interaction strength $\longrightarrow \xi_0 < \xi_c$



Moiré solids ($\max(\xi_0, \xi_c) \sim a_M$)

- Phase factor ϕ shapes moiré potential landscape
- Self-consistent Hartree-Fock solution for $\phi = 60^\circ$ has Kagome lattice charge density and quasiparticle bands, whereas moiré potential minima form honeycomb lattice!
- Competition between interaction and moiré potential energy causes electrons to localize to *saddle points* of moiré potential

