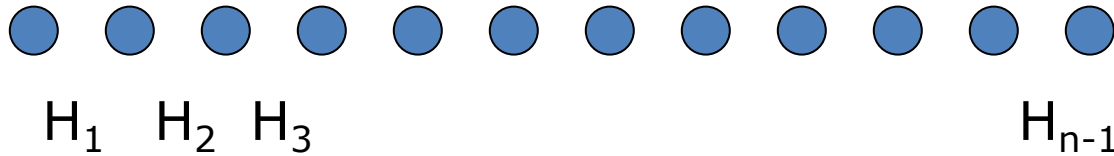


Rigorous RG: An efficient algorithm for low energy states of 1D quantum systems

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Joint work with Itai Arad, Zeph Landau
and Thomas Vidick

Local Hamiltonian in 1D



Assume each particle is a d -level system (qudit) with nearest neighbor interactions

Each term H_i is a $d^2 \times d^2$ Hermitian matrix

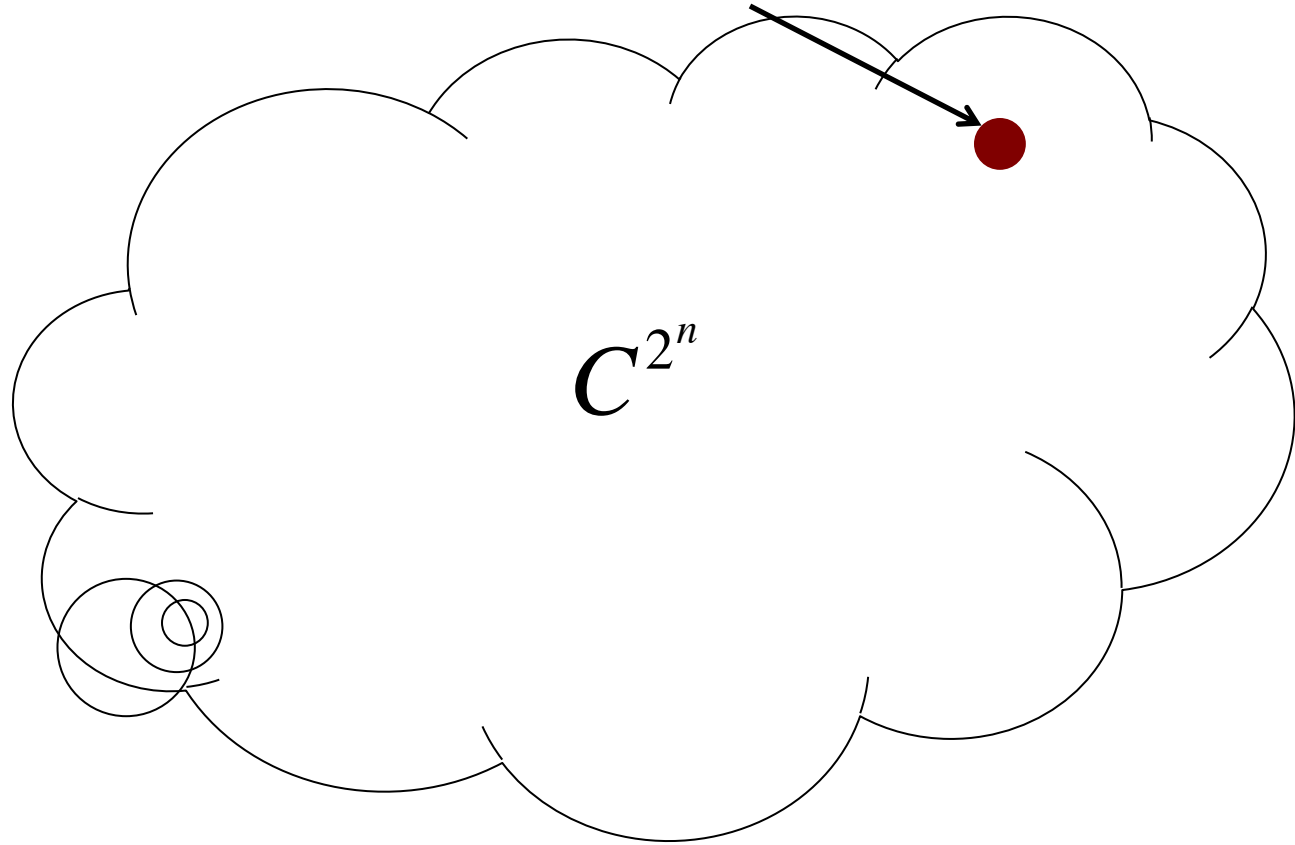
$$H = H_1 + \dots + H_m$$

- Do low energy states of H have succinct classical description?
- Is there an efficient classical algorithm to find such a description?

Physically relevant corner of Hilbert space



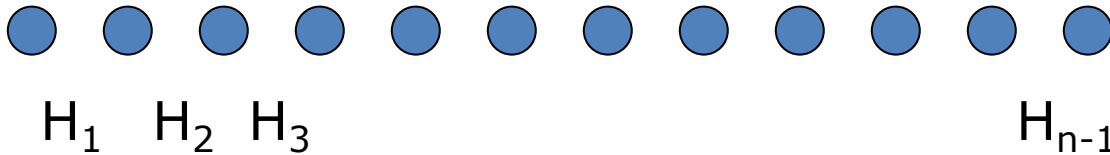
n particles



CMP: low energy states are special

- DMRG (Density Matrix Renormalization Group) [White '92] has been remarkably successful in practice for 1D quantum systems. Quickly outputs compact representation of ground/low energy state.
- Doesn't always work. Artificial hard examples known [Eisert '06]

Local Hamiltonian in 1D



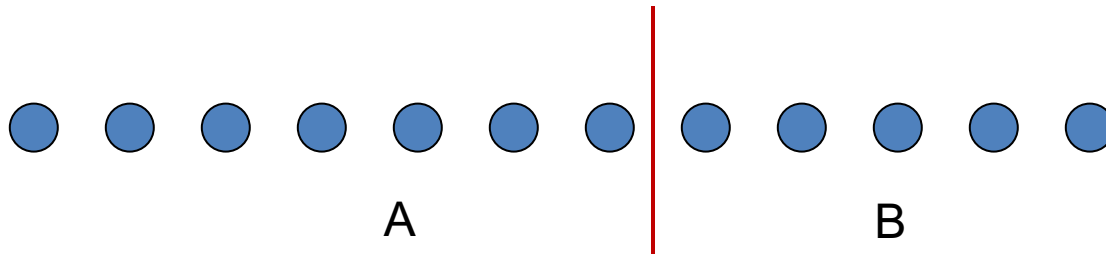
- Do low energy states of H have succinct classical description?

[Hastings '06] (Close approximations to) Ground states of gapped 1D Hamiltonians ($E_1 = E_0 + c$) have compact description as an MPS with $\text{poly}(n)$ bond dimension.

- Is there an efficient classical algorithm to find such a description?

[Landau, V, Vidick '14] Polynomial time algorithm to find ground states of gapped 1D Hamiltonians.

Area Law/Succinct Description for Ground State



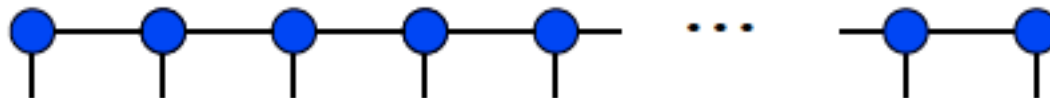
$$|\psi\rangle = \sum c_i |a_i\rangle \otimes |b_i\rangle, \quad c_i \geq 0$$

Schmidt decomposition: $\{|a_i\rangle\}$, $\{|b_i\rangle\}$ orthonormal sets

Schmidt rank = number of non-zero terms

If $|\psi\rangle$ has low Schmidt rank then if $|a_i\rangle$ and $|b_i\rangle$ vectors compactly representable then so is $|\psi\rangle$

If $|\psi\rangle$ has low Schmidt rank across all cuts then $|\psi\rangle$ has a compact representation (in the form of a tensor network called an MPS).



Two classes of Hamiltonians:

1. Gapless Hamiltonians with a low density of low-energy states
 $r = \text{poly}(n)$ dimensional space of eigenstates with energy in the range $[E_0, E_0 + c]$

2. Degenerate gapped ground space:

H has smallest eigenvalue E_0 with associated eigenspace of dimension $r = \text{poly}(n)$, and $E_1 = E_0 + c$

Constant degeneracy: [Huang; Chubb, Flammia 2014]

Do Low Energy States have compact descriptions?

- Will consider the case of a degenerate ground space of polynomial dimension
- All proofs bounding entanglement across cuts show that for every cut there is some ground state with low complexity.
- Want to show: there is a ground state with low complexity across all cuts.
Or better still, all ground states have low complexity across all cuts.

Results:

- Polynomial time algorithm for computing basis for $\text{poly}(n)$ -degenerate ground space for 1D gapped Hamiltonian.

Implies area law + succinct MPS description for $\text{poly}(n)$ -degenerate ground spaces.

- $n^{O(\log n)}$ algorithm for computing MPS descriptions of low energy states of 1D systems in gapless case ($\text{poly}(n)$ eigenstates of energy $\epsilon_0 + c$).

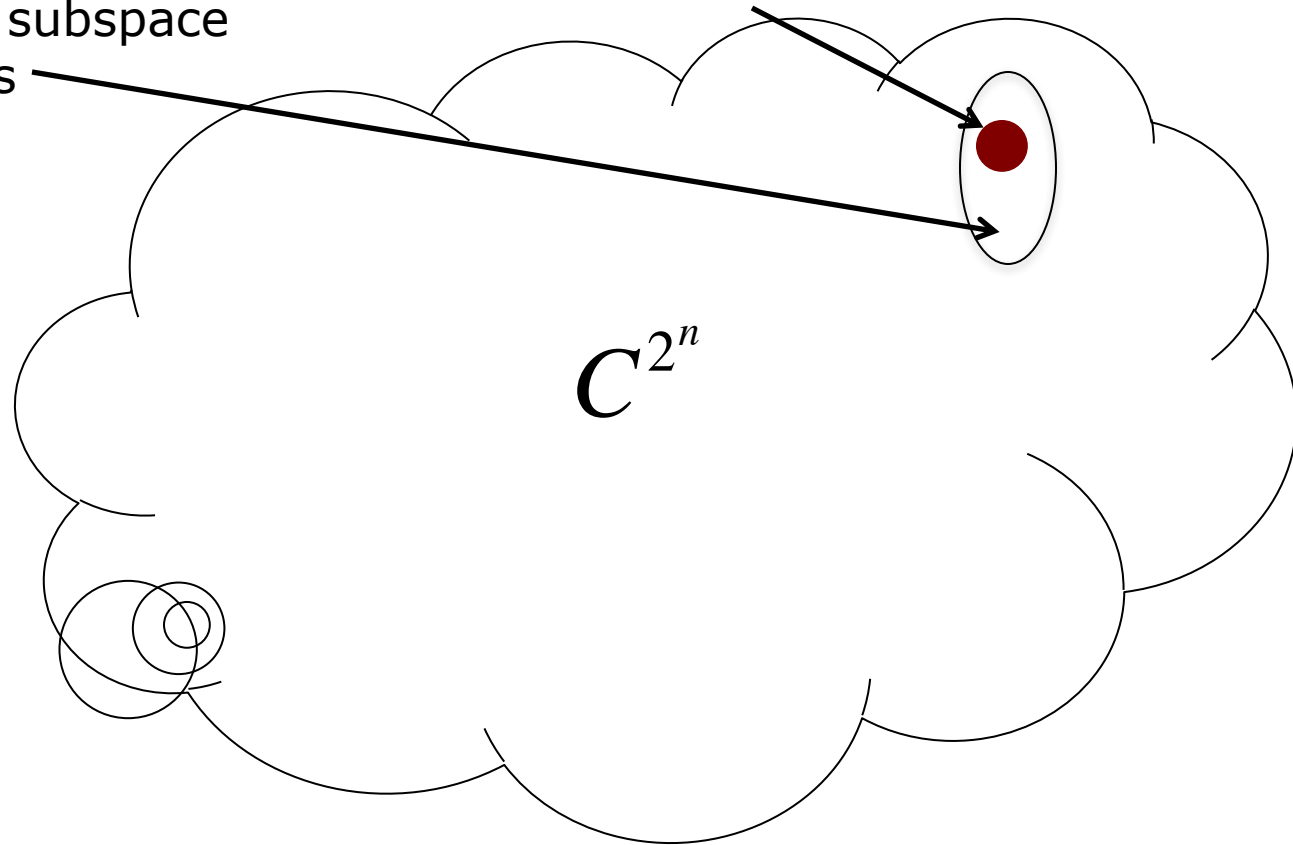
Implies area law (up to log correction) + succinct
MPS description.

What we would like!

Quickly identify a small subspace
in which the solution is
guaranteed to lie

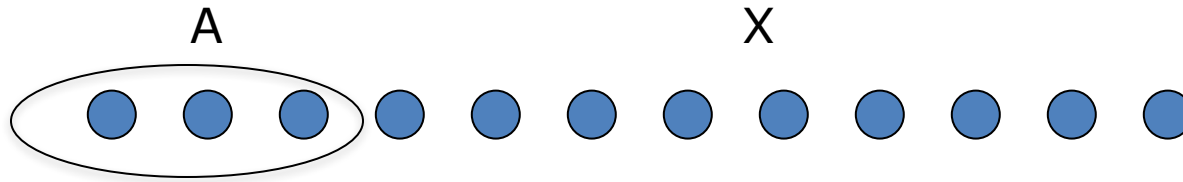
Subspace T of low energy states

.....
 $H = H_1 + H_2 + \dots H_{n-1}$



Once this is done, can compute eigenvectors quickly!

Local Approach



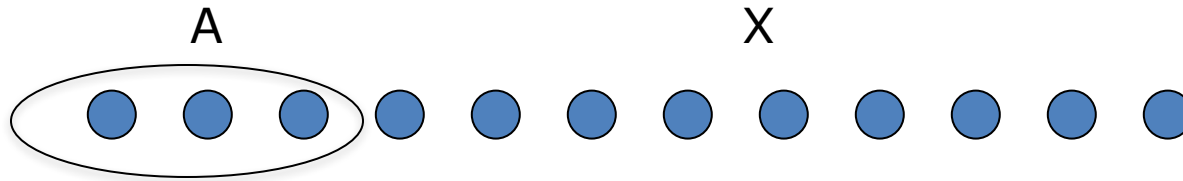
Suppose we are trying to find unique ground state $|\psi\rangle$

What does a partial solution on the first few particles look like?

$$\text{If } |\psi\rangle = \sum c_i |a_i\rangle \otimes |b_i\rangle, \quad c_i \geq 0$$

then partial solution looks like $S = \text{span}\{|a_i\rangle\}$

Local Approach



T = Target subspace of low energy states

Ideally: Identify a subspace $S \subseteq T$ such that $T \subseteq S \otimes X$

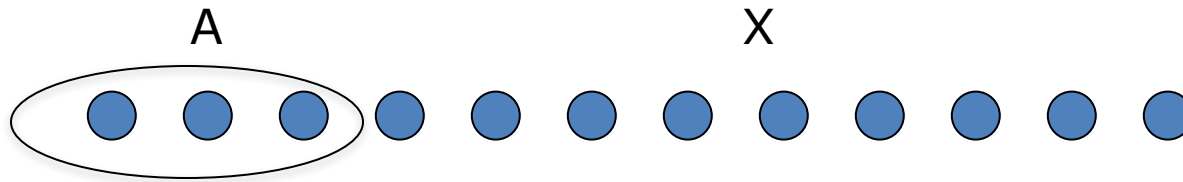
Definition: We will say that S is a δ -viable set if

$$P_T P_{S \otimes X} P_T \geq (1 - \delta) P_T$$

δ = error. $(1 - \delta)$ = overlap

Want: $\dim(S)$ small, δ small.

Local Approach



Definition: We will say that S is a δ -viable set if

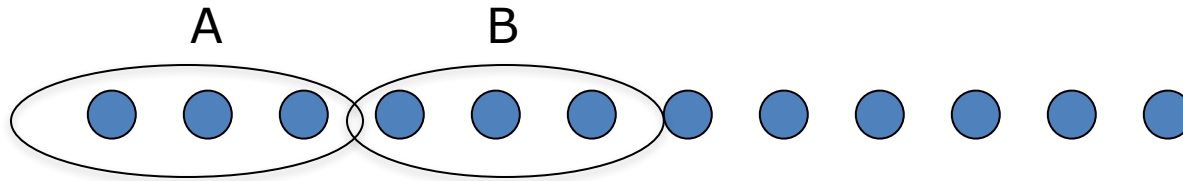
$$P_T P_{S \otimes X} P_T \geq (1 - \delta) P_T$$

Interpretation: A unit ball of $S \otimes X$ when projected onto T contains a ball of radius $(1 - \delta)$

Viability sets

- Think of S interchangeably as subspace or as a compactly represented basis for the subspace.
- $\delta = \text{error}$. $(1 - \delta) = \text{overlap}$
- No efficient test for whether a set is viable or to estimate δ

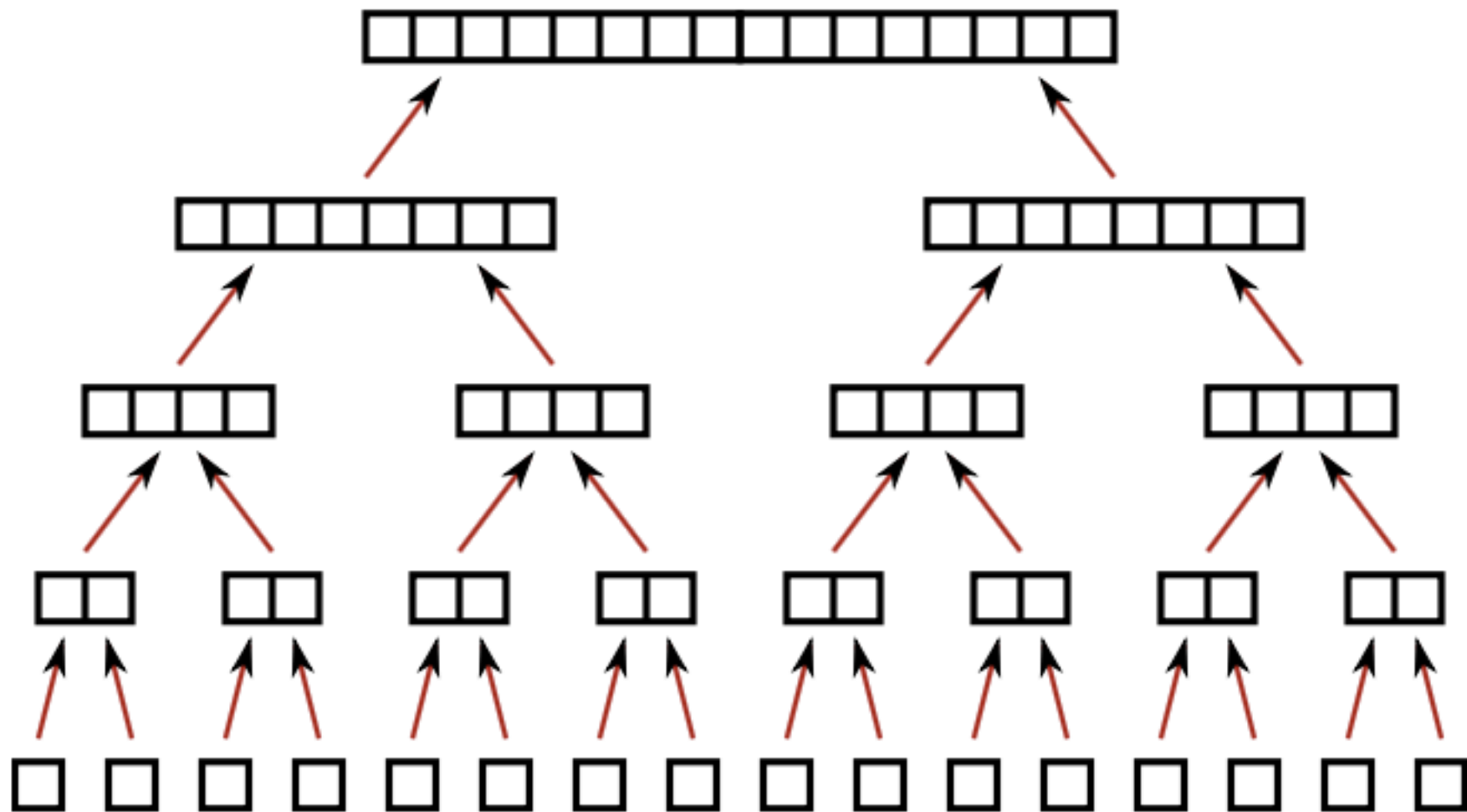
Algorithm design primitives for viable sets:



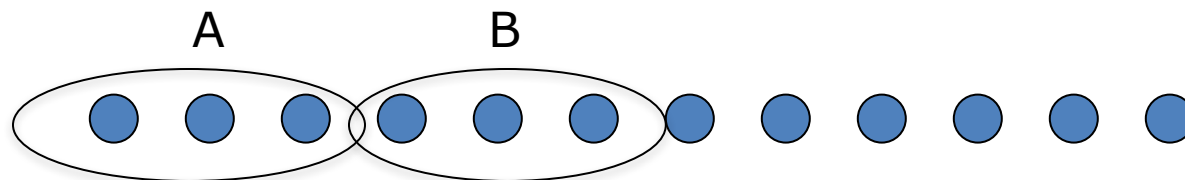
1. Tensoring

If S_1 is δ_1 -viable and S_2 is δ_2 -viable then
 $S_1 \otimes S_2$ is $(\delta_1 + \delta_2)$ -viable

$$\text{Dim}(S_1 \otimes S_2) = \text{Dim}(S_1) \text{Dim}(S_2)$$

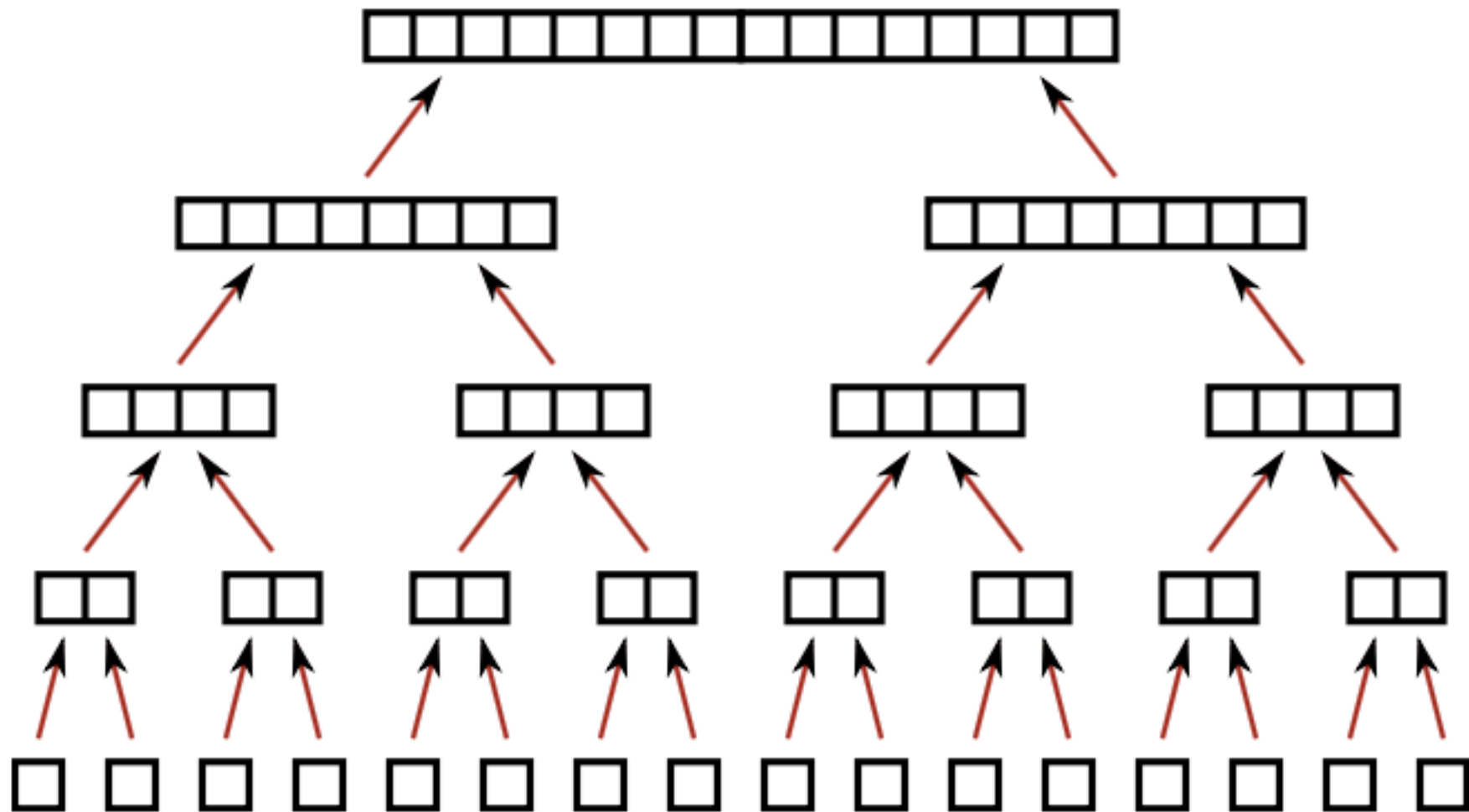


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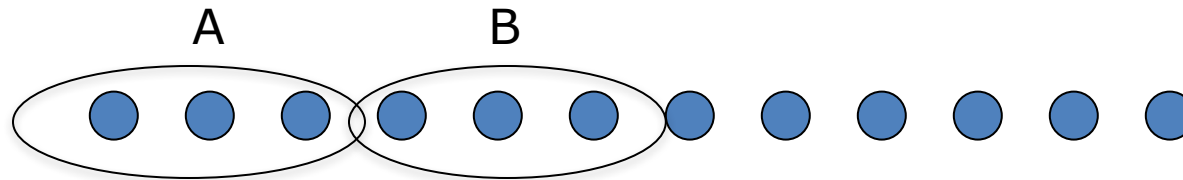


2. Random projection

If S is δ -viable i.e. $\text{overlap} = 1 - \delta$, and $\text{Dim}(S) = s$,
Let R be a random subspace of S of dimension r .
Then R has $\text{overlap} \sim (1 - \delta) r/s$



Algorithm design primitives for viable sets:

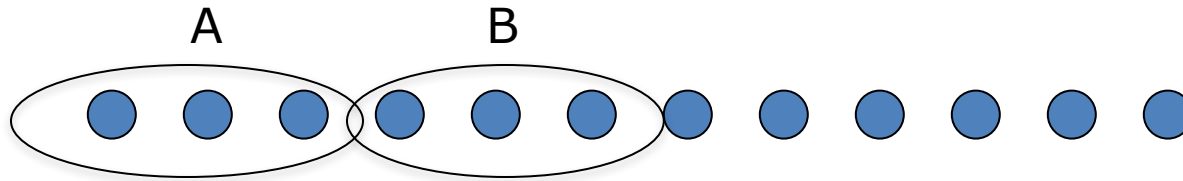


3. Error Reduction

If S is δ -viable and $\text{Dim}(S) = s$. Then error reduction yields S' which is $\Delta / (1 - \delta)^2$ -viable and with $\text{Dim}(S') = D^2s$.

Error reduction carried out by applying D - Δ AGSP, which satisfies $D^{16}\Delta \ll 1$

Merge Process

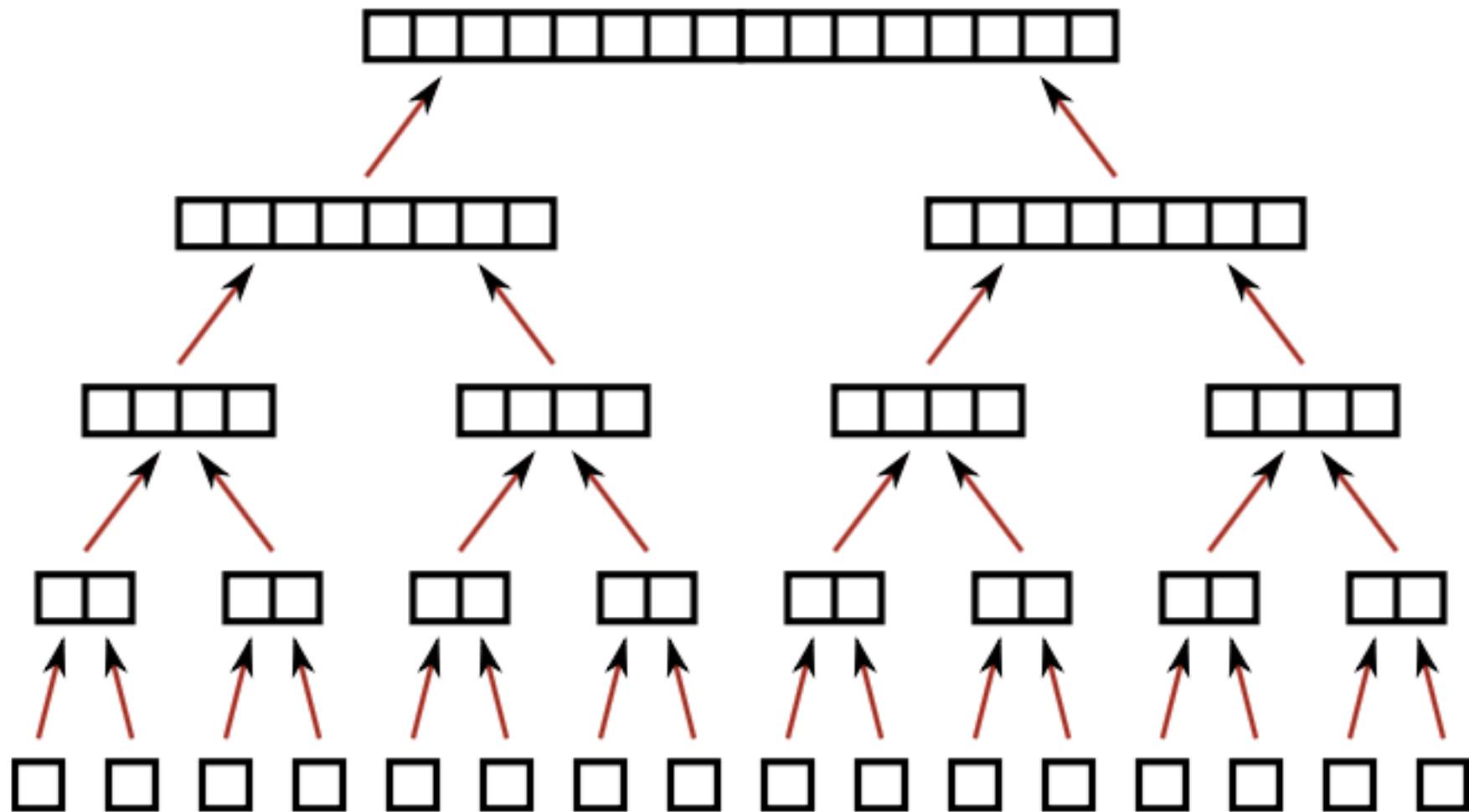


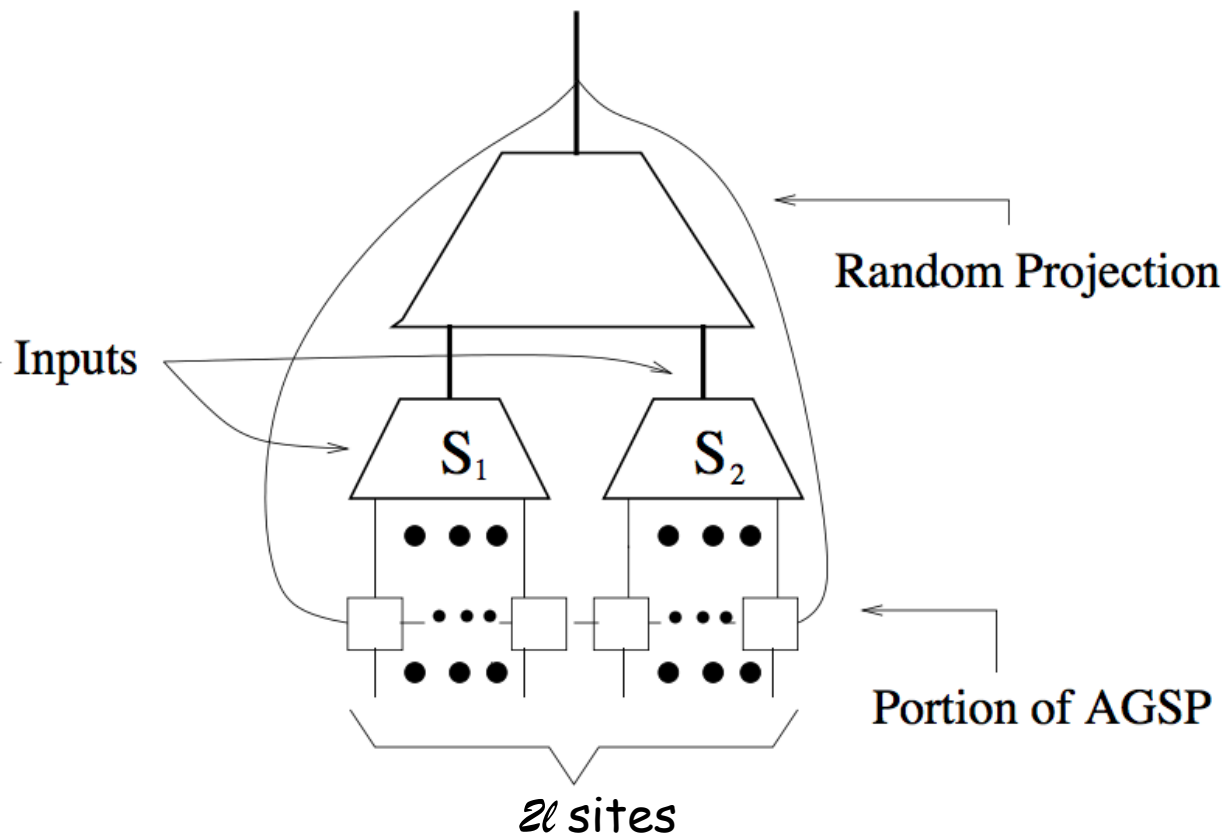
Step 1: Tensor S_1, S_2 δ -viable, each of $\text{Dim} = s$,
to get $S_1 \otimes S_2$ 2δ -viable and of $\text{Dim} = s^2$

Step 2: Random projection to $\text{Dim} = r$
to get $\sim(1 - r/s^2)$ -viable

Step 3: Error reduction to $\text{Dim} = s = D^2 r$ to get
 $\Delta / (1 - \delta)^2$ -viable

$$\Delta s^4 / r^2 \ll s^4 / r^2 \quad r^8 / s^8 = r^6 / s^4 \ll 1$$





Results:

- Polynomial time (n^{1/c^2}) algorithm for computing basis for $r = \text{poly}(n)$ -degenerate ground space for 1D gapped Hamiltonian.

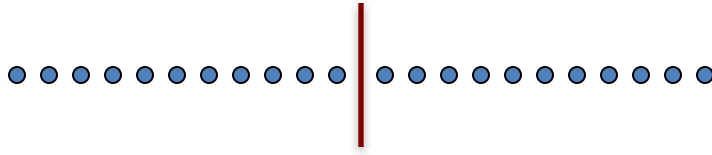
Implies area law: $O(\log r + \frac{\log^3 d}{c})$

- $n^{O(\log n)}$ algorithm for computing MPS descriptions of low energy states of 1D systems in gapless case ($r = \text{poly}(n)$ eigenstates of energy $\varepsilon_0 + c$).

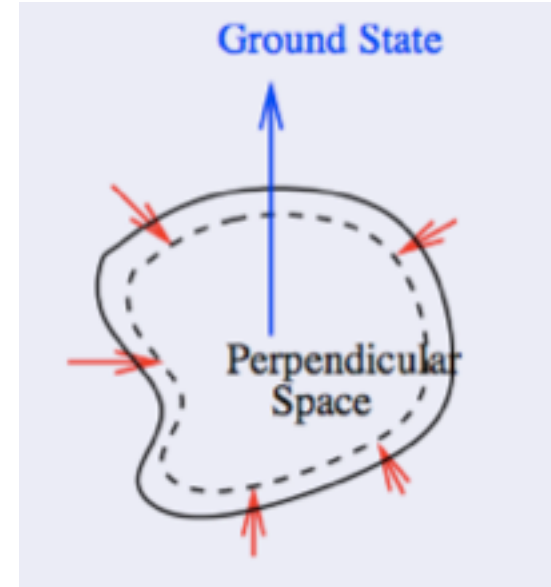
Implies area law (up to log correction):

$$O(\log r + \frac{\log^3 d}{c}) \log n$$

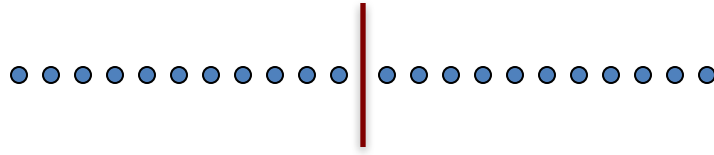
AGSP: Approximate Ground State Projector



An AGSP is an operator K that is not “too complex” and approximately projects onto the ground state:

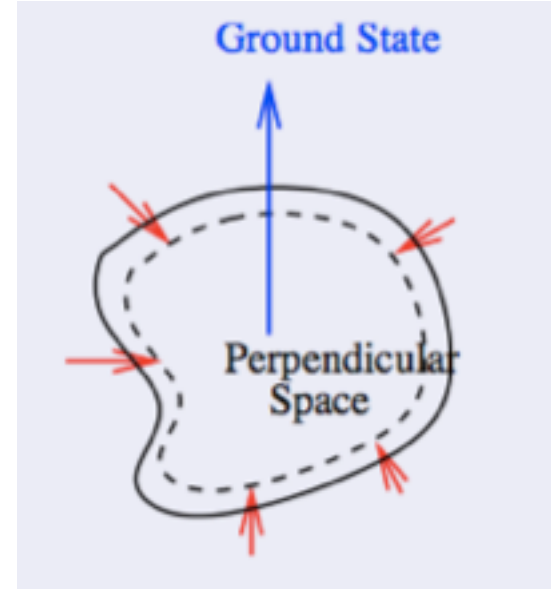


AGSP: Approximate Ground State Projector

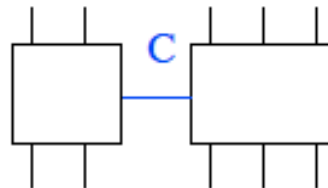


An AGSP is an operator K that is not “too complex” and approximately projects onto the ground state:

- $K|GS\rangle = |GS\rangle$
- Shrinks orthogonal space by $\Delta < 1$
- Has low entanglement rank D : $D\Delta \ll 1$

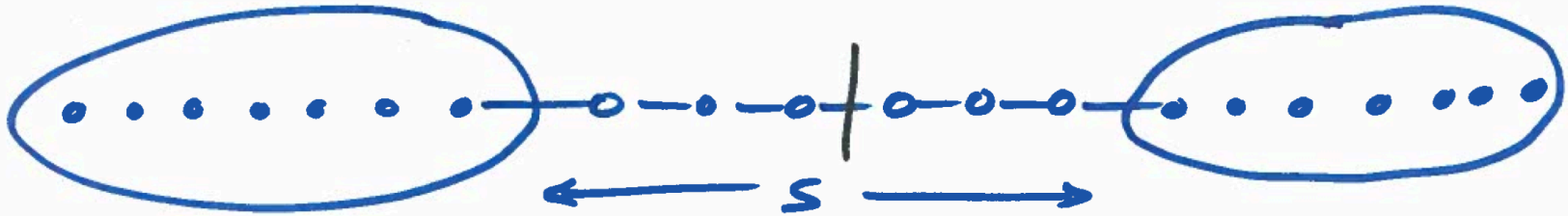


An operator on $\mathcal{H}_1 \otimes \mathcal{H}_2$ of the form $\sum_1^C A_i \otimes B_i$ will be said to have **entanglement rank C** .



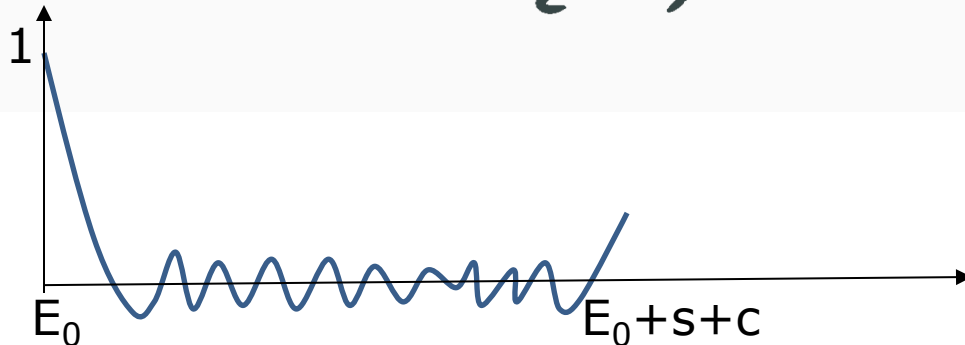
AGSP Construction

[Arad, Kitaev, Landau, V '13]



$$H = H_L + H_1 + H_2 + \dots + H_s + H_R$$

Theorem : $s = O\left(\frac{\log^2 d}{\epsilon}\right)$, then $S_{1D} = O\left(\frac{\log^3 d}{\epsilon}\right)$



$K = C_1(H')$, where C_1 is a scaled Chebyshev polynomial

Must modify AGSP construction to ensure:

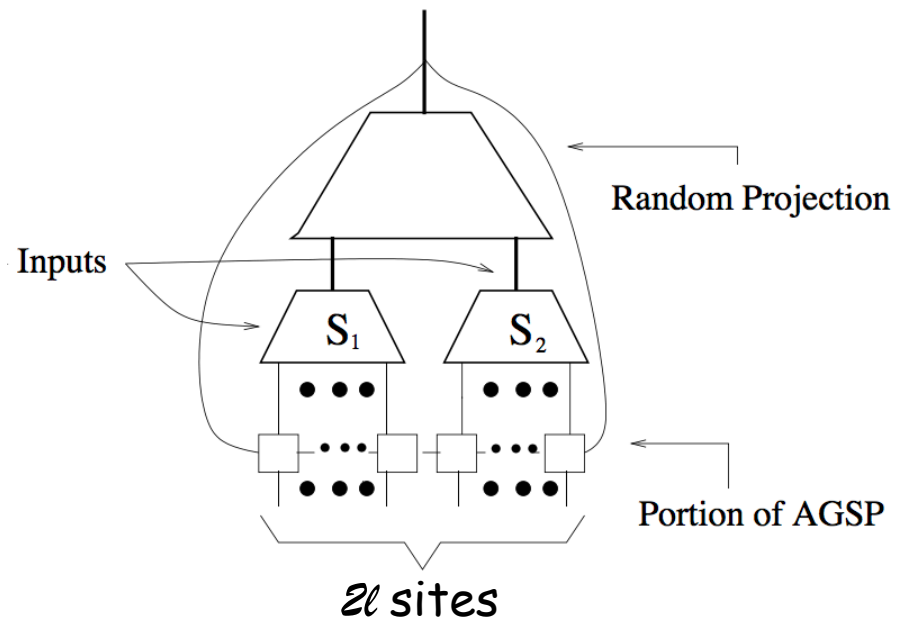
1. Bond dimension away from special cuts is $O(\text{poly}(n))$
2. Computationally efficient

Discussion

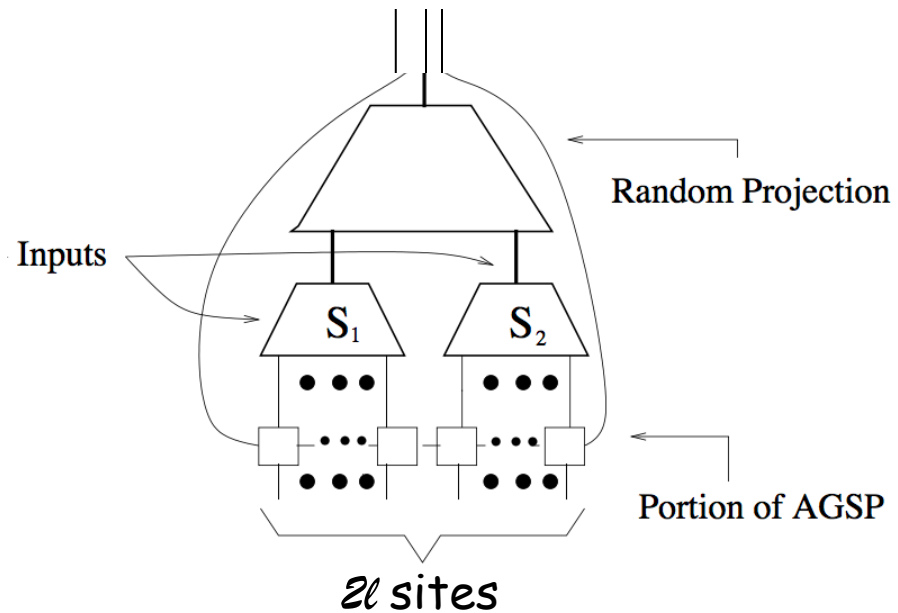
Implementation of heuristic version by [Roberts, Vidick, Motrunich]:

- Uses heuristic version of AGSP, trotter expansion of $e^{-kH/t}$
- Random sampling replaced by taking bottom s minimal eigenvectors of restricted Hamiltonian
- 5-10 times slower than DMRG
- RRG appears to do better when the ground state degeneracy is large or when reconstructing low energy space (Bravyi-Gosset model, XY model with random couplings)

- For frustration-free Hamiltonians with unique ground state, the algorithm works in $\sim O(nM(n))$ time. If a conjecture about bond trimming is correct, then the running time can be reduced to $\sim O(n)$.
- Renormalization



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- Renormalization



- Random projection step acts on a bond in the tensor network.
- Consider a toy model for quantum states where we think of an underlying physically manifested tensor network. Suppose the bonds of the tensor network are subjected to a noise process, described by random projections. What is the computational complexity of this model? BQP? BPP?
i.e. does it support fault-tolerance?