



# Hierarchy of fluctuation theorems and experimental test of the differential fluctuation theorem

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# Outline

## Background

- Length of time's arrow
- Stochastic thermodynamics and fluctuation theorems (FTs)
- Motivation

## Hierarchy of fluctuation theorems

- Microscopic reversibility (MR)
- Differential fluctuation theorem (DFT)
- A generalized Jarzynski equality (GJE) for arbitrary initial states

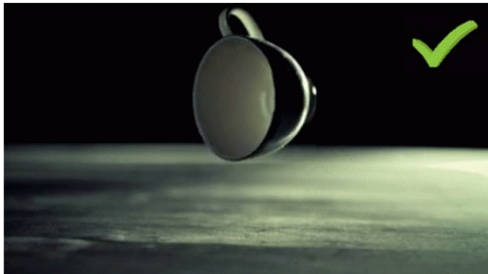
## Experimental test

- Setup and data
- Generate arbitrary initial states

## Summary

# Length of time's arrow

## Macroscopic



- Thermal fluctuations are negligible
- Deterministic
- Newton's equation



- $W \geq \Delta F$
- $\Delta S \geq 0$

## Microscopic



- Thermal fluctuations are dominant
- Stochastic
- Langevin equation
- Fokker-Planck equation



- $\langle W \rangle \geq \Delta F$
- $\langle \Delta S \rangle \geq 0$

# Stochastic thermodynamics and Fluctuation theorems (FTs)

Thermodynamic quantities based on stochastic trajectories:

**Work** <sup>[1,2]</sup>:  $W[\Gamma(t)] = \int_0^\tau \frac{\partial H}{\partial \lambda} \dot{\lambda} dt$

**Heat** <sup>[1,2]</sup>:  $Q[\Gamma(t)] = \int_0^\tau (-\gamma v_t + \xi_t) \circ dx_t = \int_0^\tau \frac{\partial H}{\partial \Gamma} \dot{\Gamma} dt$

**Entropy production** <sup>[4]</sup>:

$$S_{sys}(t) = -\ln \rho(\Gamma_t, t)$$

$$\Delta S_{tot} = \Delta S_{sys} - \frac{Q}{T}$$

$$\Gamma_t := (x_t, p_t)$$

• **Jarzynski equality (JE)** <sup>[1]</sup>:

$$\langle e^{-\beta(W-\Delta F)} \rangle = 1$$

• **Crooks fluctuation theorem (CFT)** <sup>[3]</sup>:

$$\frac{P_R(-W)}{P_F(W)} = e^{-\beta(W-\Delta F)}$$

• **Hummer-Szabo relation (HSR)** <sup>[5]</sup>:

$$\langle \delta(\tilde{\Gamma} - \Gamma_\tau) e^{-\beta W} \rangle = e^{-\beta U_\tau(\tilde{\Gamma})} / Z_0$$

• **FT of entropy production** <sup>[4]</sup>:

$$\frac{P_R(-\Delta S_{tot})}{P_F(\Delta S_{tot})} = e^{-\Delta S_{tot}/k_B}$$

$$\langle e^{-\Delta S_{tot}} \rangle = 1$$

• **Hatano-Sasa relation** <sup>[6]</sup>:

$$\langle e^{-Y} \rangle = 1, Y := \int_0^\tau \frac{\partial \phi}{\partial \alpha} (x(t), \alpha(t)) \dot{\alpha} dt$$

• **Sagawa-Ueda relation** <sup>[7]</sup>:

$$\langle e^{-\beta(W-\Delta F)-I} \rangle = 1$$

.....

# What is the origin of fluctuation theorems?

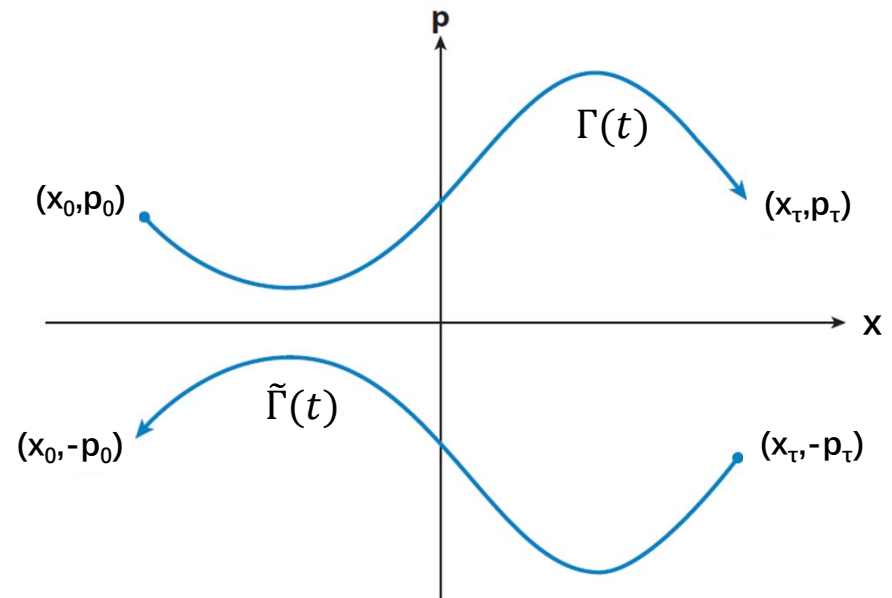
Microscopic reversibility [3,4,8,22]:

$$\ln \frac{p(\Gamma(t)|\Gamma_0)}{\tilde{p}(\tilde{\Gamma}(t)|\tilde{\Gamma}_0)} = -\beta Q[\Gamma(t)]$$

$$\tilde{\lambda}_t := \lambda_{\tau-t}, \tilde{\Gamma}_t := (x_{\tau-t}, -p_{\tau-t})$$

**Impossible** to test it in experiment!

How about it if we do a little bit coarse-graining?



[3] G. E. Crooks, Phys. Rev. E 60, 2721 (1999).

[4] U. Seifert, Phys. Rev. Lett. 95, 040602 (2005).

[8] G. E. Crooks, *Journal of Statistical Physics* 90.5-6 (1998).

[22] C. Jarzynski, Annu. Rev. Condens. Matter Phys. 2, 329 (2011).

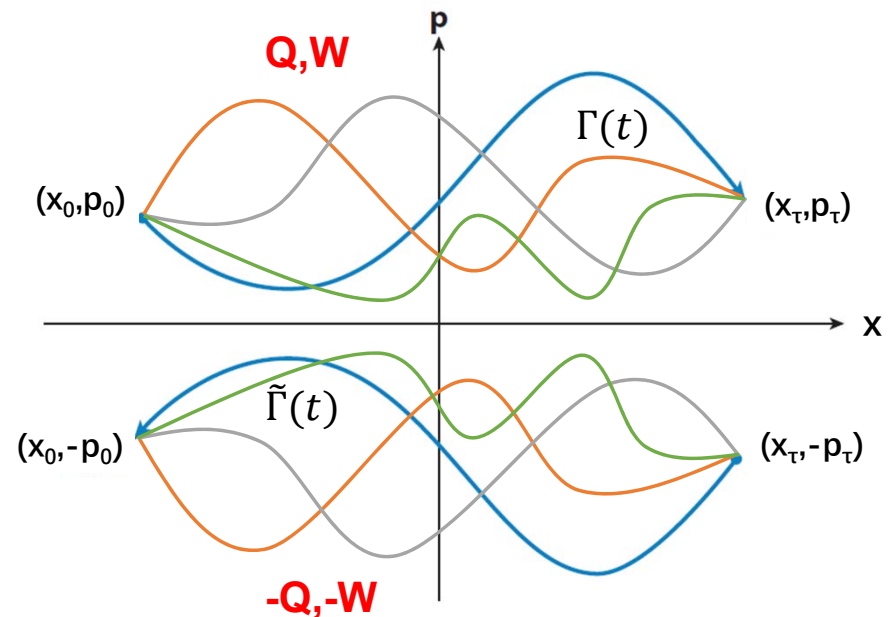
# Differential fluctuation theorem (DFT)

- **DFT version 1** (2000, Jarzynski) <sup>[9]</sup>:

$$\frac{P_R(-Q, \Gamma_0^\dagger | \Gamma_\tau^\dagger)}{P_F(Q, \Gamma_\tau | \Gamma_0)} = e^{\beta Q}$$

- **DFT version 2** (2008, Karplus) <sup>[10]</sup>:

$$\frac{P_R(-W, \Gamma_\tau^\dagger \rightarrow \Gamma_0^\dagger)}{P_F(W, \Gamma_0 \rightarrow \Gamma_\tau)} = e^{-\beta(W - \Delta F)}$$



The most detailed fluctuation theorem that can be tested **experimentally**.

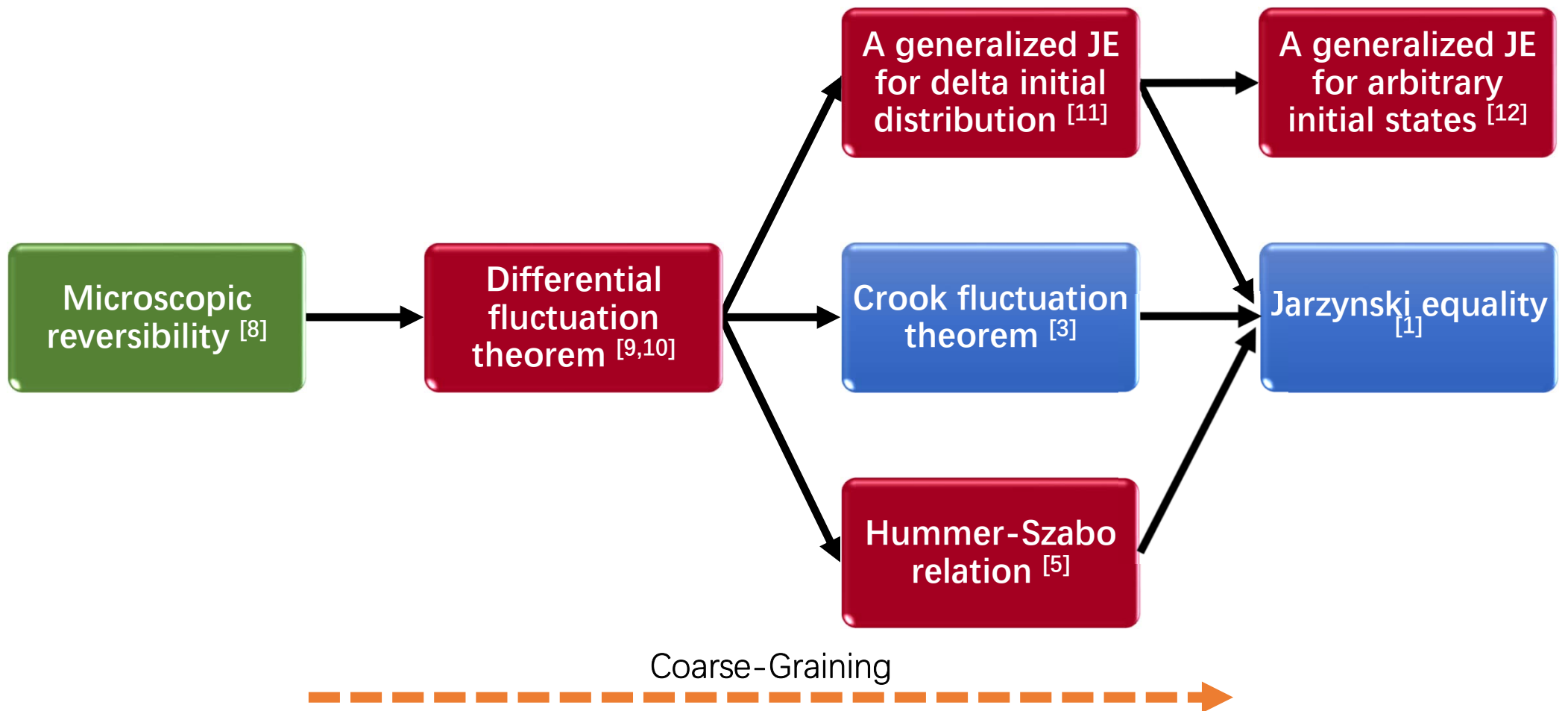
[9] C. Jarzynski, J. Stat. Phys. 98, 77 (2000).

[10] P. Maragakis, et al. J. Phys. Chem. B 112, 6168 (2008).

[11] R. Kawai, et al. Phys. Rev. Lett. 98, 080602 (2007).

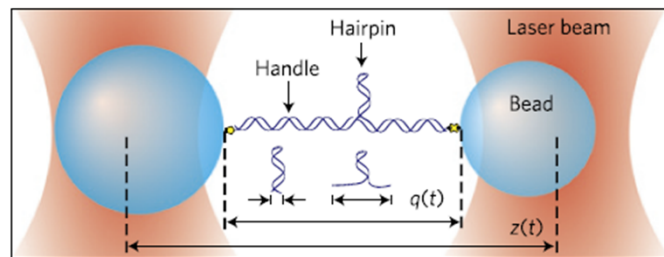
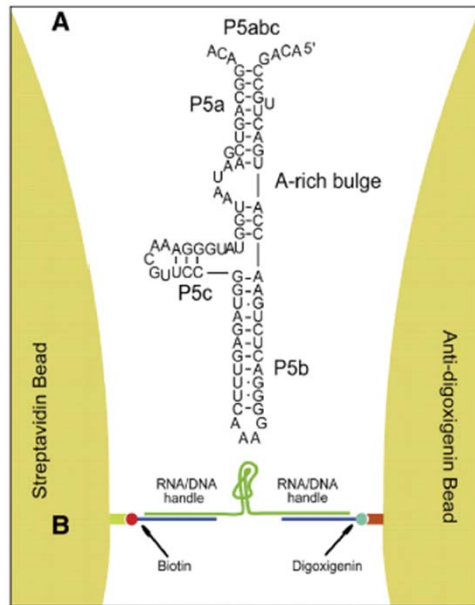
[12] Z. Gong and H. T. Quan, Phys. Rev. E 92, 012131 (2015).

# Hierarchy of fluctuation theorems

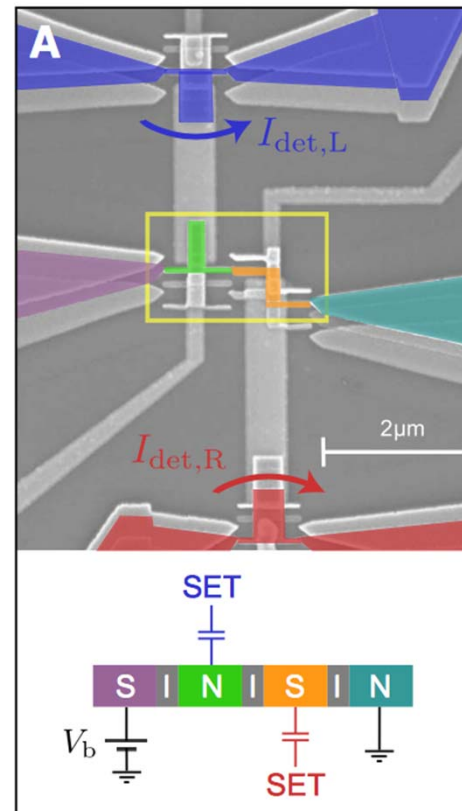


# Previous experiments

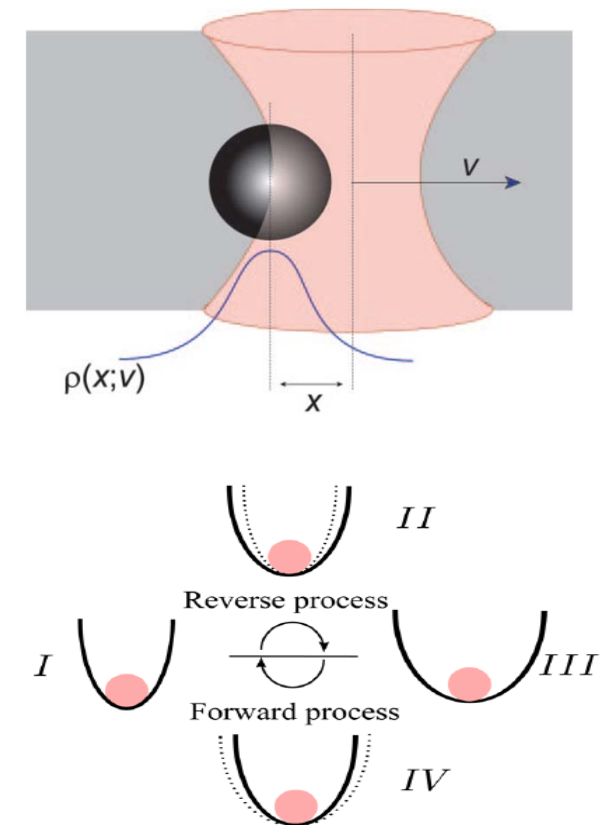
Stretching RNA molecules [13-15]



Electronic circuit [24,25]



Brownian particle trapped in water [16,17]

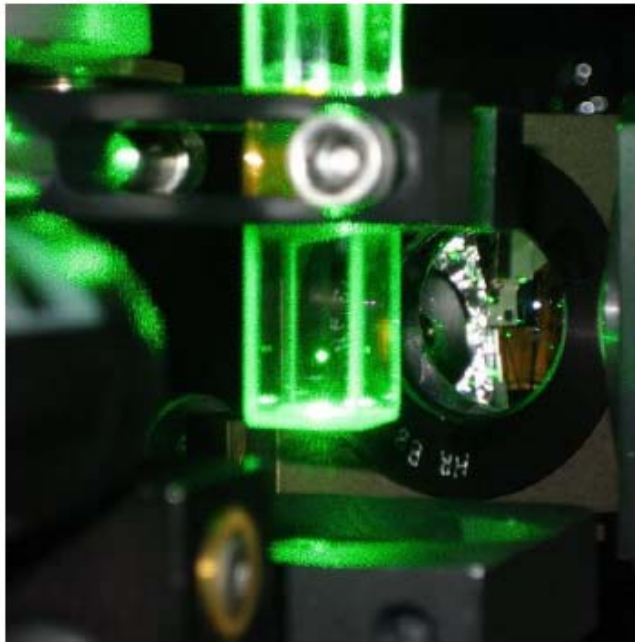


- [13] J. Liphardt, et al. *Science* 296, 1832 (2002).
- [14] D. Collin, et al. *Nature (London)* 437, 231 (2005).
- [15] A. N. Gupta, et al. *Nature Physics* 7,631 (2011).
- [16] E. Trepagnier, et al. *Proc. Natl. Acad. Sci.* 101, 15038 (2004).
- [17] D. Y. Lee, et al. *Phys. Rev. Lett.* 114, 060603 (2015).
- [24] J. P. Pekola, *Nat. Phys.* 11, 118 (2015).
- [25] S. Singh, et al. *arXiv:1712.01693* (2017).



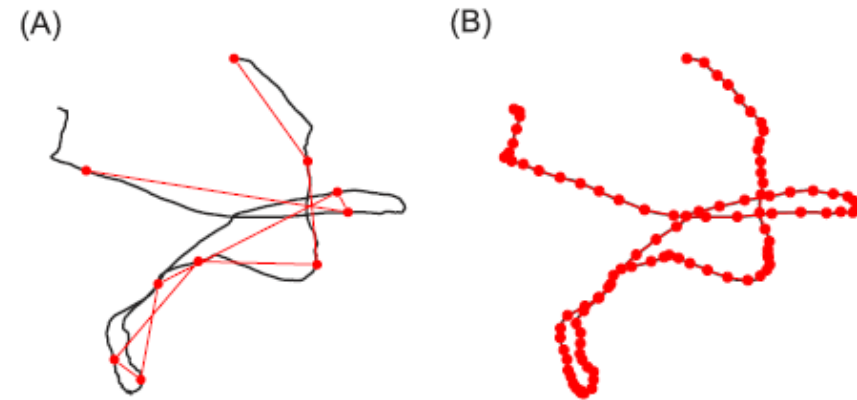
# Our experiment

Silica nanosphere levitated in **AIR** using optical tweezers [18-20]

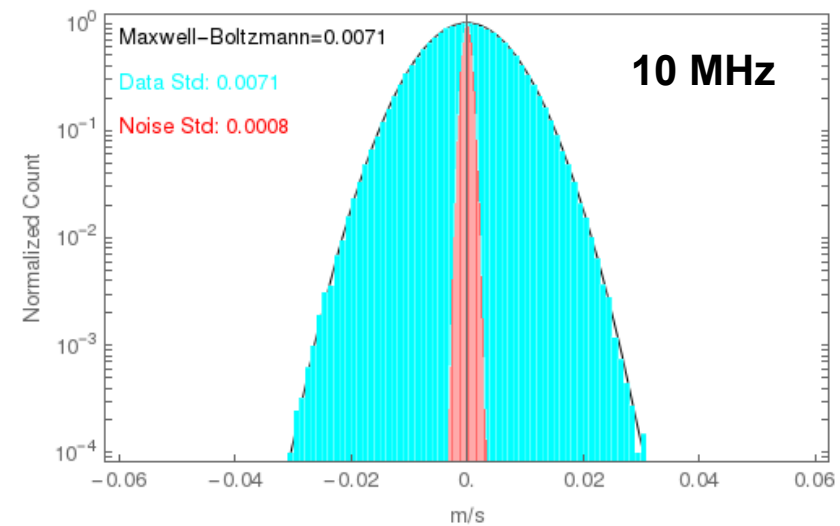


We can study the thermodynamics of the **phase space**.

- [18] T. Li, et al. Science 328, 1673 (2010).
- [19] S. Kheifets et al. Science 343, 1493 (2014) .
- [20] T. M. Hoang, et al. Nat. Commun. 7, 12250 (2016).



Measurement of the instantaneous velocity



# Our experiment

- In air, room temperature (296K)

$$H(x, v, t) = \frac{1}{2}kx^2 - f(t)x + \frac{1}{2}mv^2$$

$$f(t): f_{\text{off}} \rightarrow f_{\text{on}}$$

$$\Delta F = -\frac{f_{\text{on}}^2 - f_{\text{off}}^2}{2k}, \quad W = \int_0^\tau \frac{\partial H}{\partial f} \dot{f}(t) dt$$

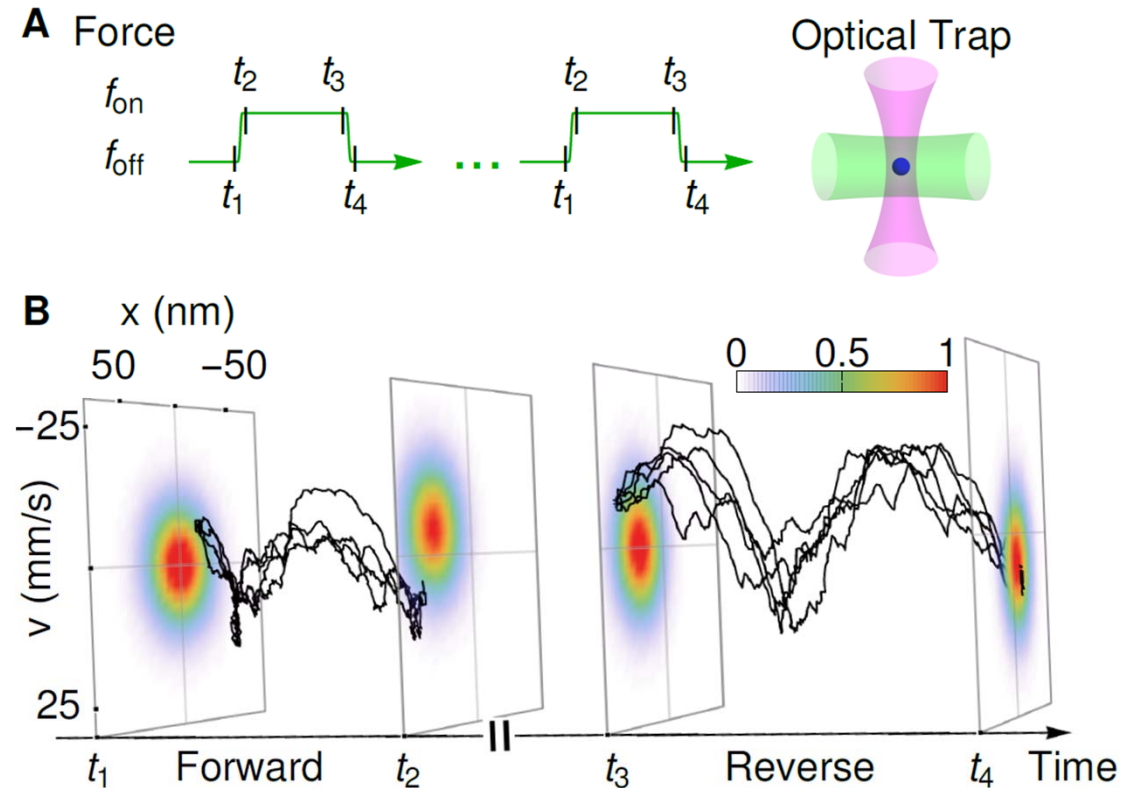
- Challenges of our experiment:

➤ Very large statistics

> **one million** cycles,  $500 \mu\text{s}$  /cycle

$\sim 10^{10}$  data points in the phase space with 10 MHz acquisition rate

➤ Track individual trajectories in the phase space (instantaneous velocity measurement)



# Test the differential fluctuation theorem

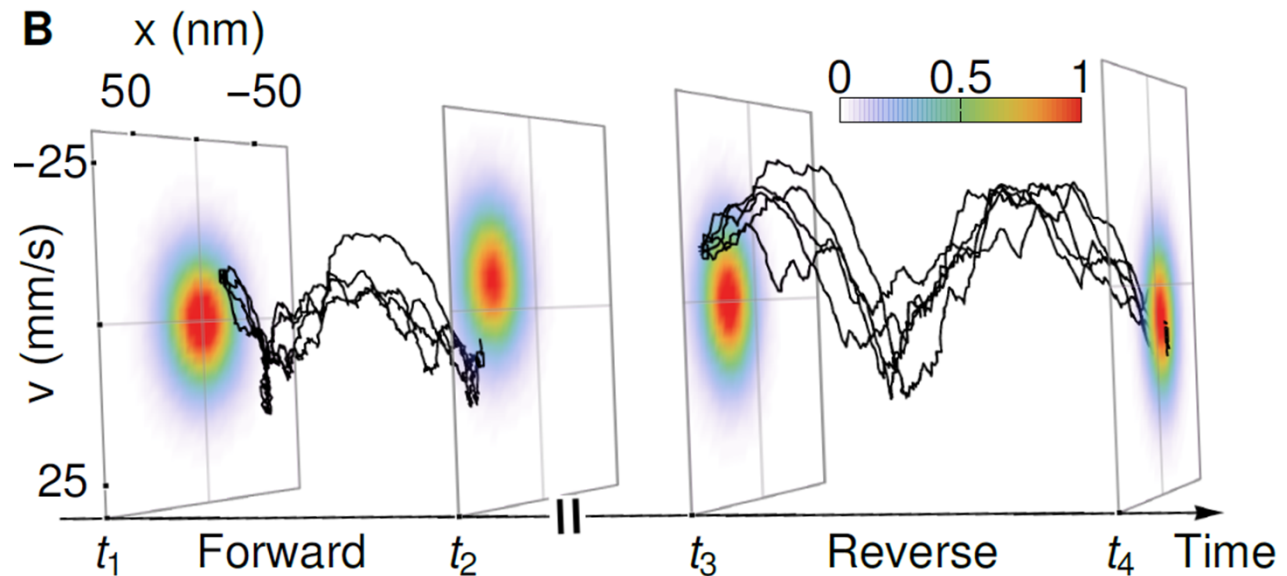
- Underdamped regime (50 Torr)

$$\frac{P_R(-W, x_2 \rightarrow x_1)}{P_F(W, x_1 \rightarrow x_2)} = e^{-\beta(W - \Delta F)}$$

$$\frac{P_R(-W, -v_2 \rightarrow -v_1)}{P_F(W, v_1 \rightarrow v_2)} = e^{-\beta(W - \Delta F)}$$

$$(x, v) := \left(x \pm \frac{\sigma_x}{11}, v \pm \frac{\sigma_v}{11}\right)$$

121 combinations of  $\{x_1, x_2\}$  or  $\{v_1, v_2\}$



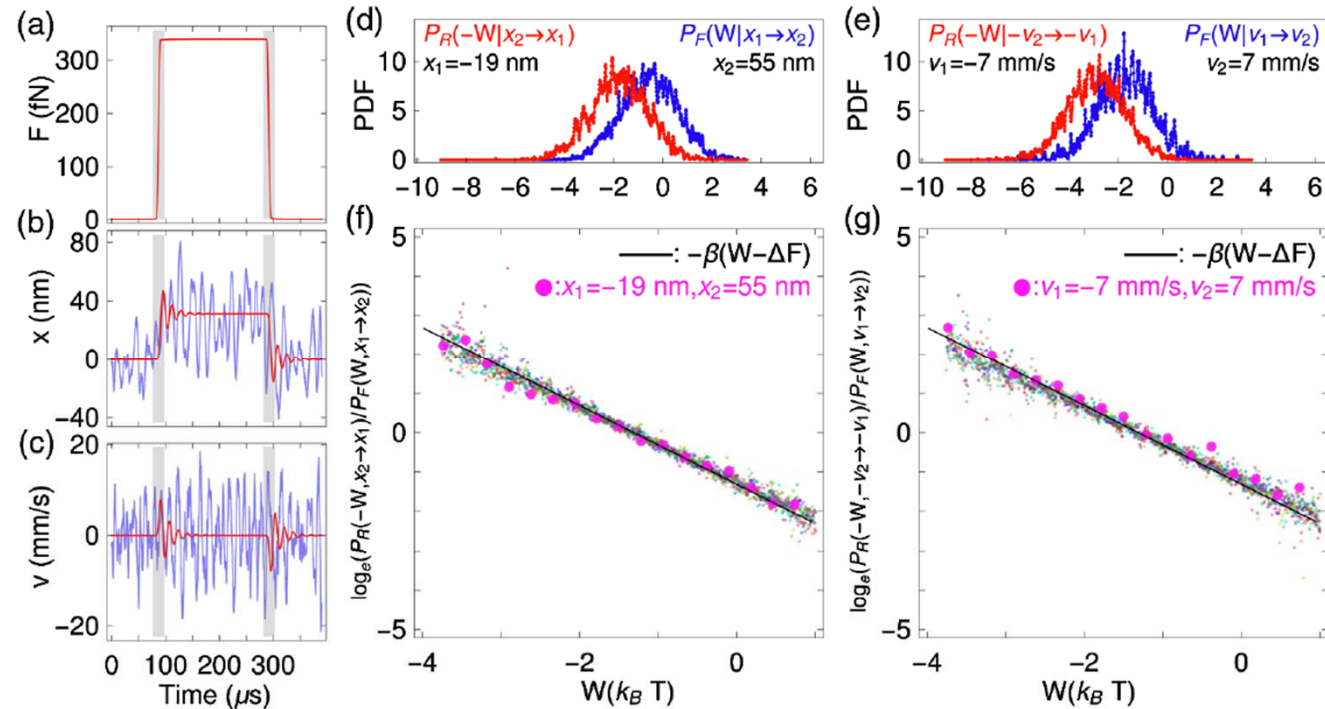
# Test the differential fluctuation theorem

- Underdamped regime (50 Torr)

$$\frac{P_R(-W, x_2 \rightarrow x_1)}{P_F(W, x_1 \rightarrow x_2)} = e^{-\beta(W - \Delta F)}$$

$$\frac{P_R(-W, -v_2 \rightarrow -v_1)}{P_F(W, v_1 \rightarrow v_2)} = e^{-\beta(W - \Delta F)}$$

- $f_{off} = 0, f_{on} = 340 \text{ fN}$
- Ramp time  $\sim 4.6 \mu\text{s}$
- Particle size  $r = 209 \pm 9 \text{ nm}$
- $\Omega = 60.4 \pm 0.3 (2\pi \cdot \text{kHz})$
- $\Delta F = -1.3 k_B T$



# Test the GJE for delta distribution and the HSR

- GJE for delta distribution:

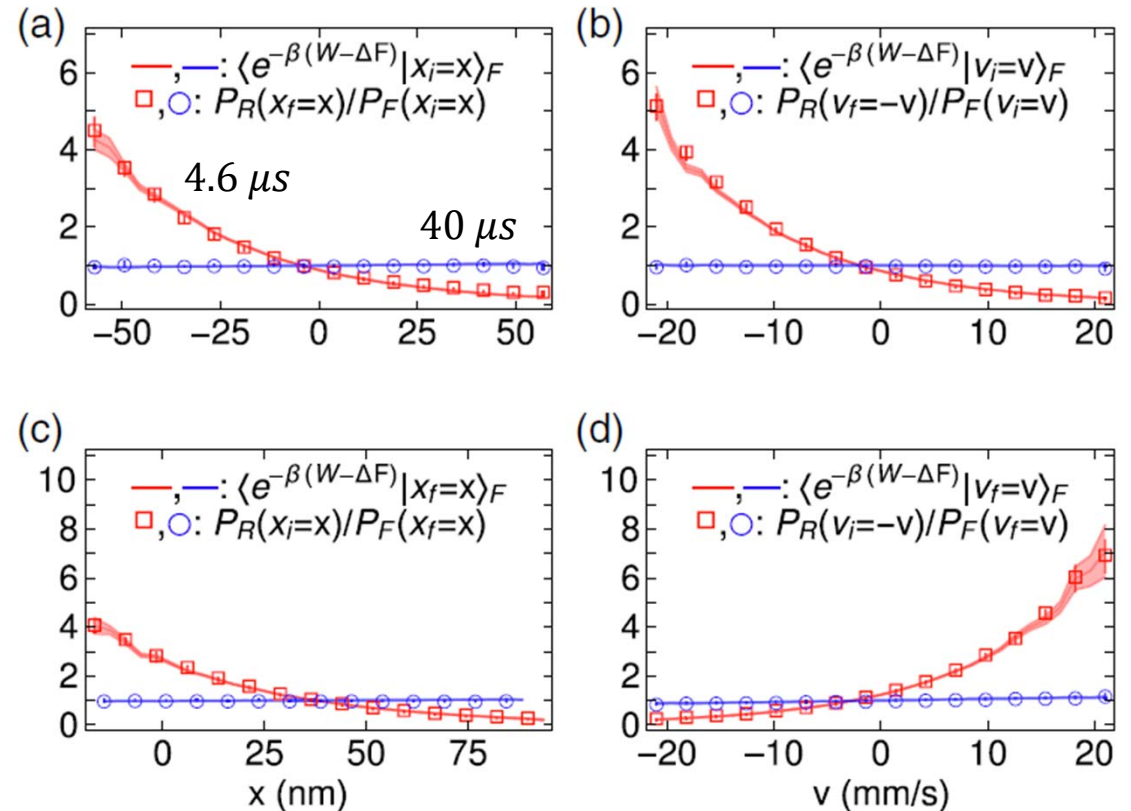
$$\langle e^{-\beta(W-\Delta F)} |x_i = x\rangle_F = \frac{P_R(\tilde{x}_f=x)}{P_F^{eq}(x_i=x)}$$

$$\langle e^{-\beta(W-\Delta F)} |v_i = v\rangle_F = \frac{P_R(\tilde{v}_f=-v)}{P_F^{eq}(v_i=v)}$$

- HSR:

$$\langle e^{-\beta(W-\Delta F)} |x_f = x\rangle_F = \frac{P_R^{eq}(\tilde{x}_i=x)}{P_F(x_f=x)}$$

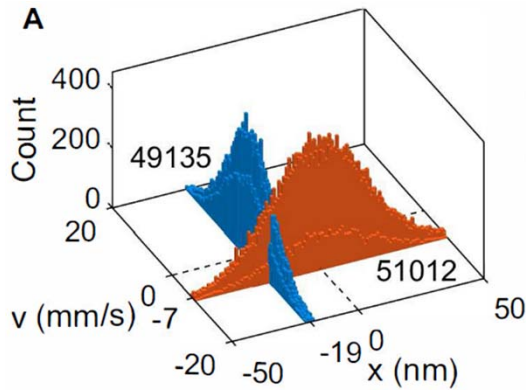
$$\langle e^{-\beta(W-\Delta F)} |v_f = v\rangle_F = \frac{P_R^{eq}(\tilde{v}_i=-v)}{P_F(v_f=v)}$$



# Test the GJE for arbitrary initial states

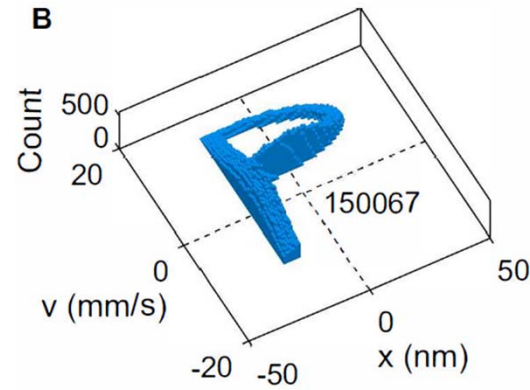
- GJE for arbitrary initial states

$$\langle e^{-\beta(W-\Delta F)} \rangle_{P_{ini}(x_i, v_i)} = \int \frac{P_R(\tilde{x}_f = x, \tilde{v}_f = -v)}{P_F^{eq}(x_i = x, v_i = v)} P_{ini}(x_i = x, v_i = v) dx dv$$



**Delta  
distribution**

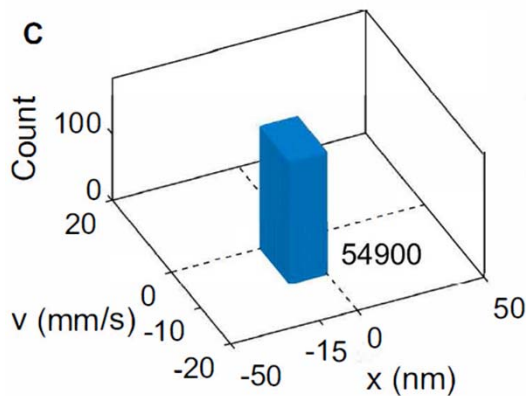
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**P-shaped state**

lhs:  $0.92 \pm 0.02$

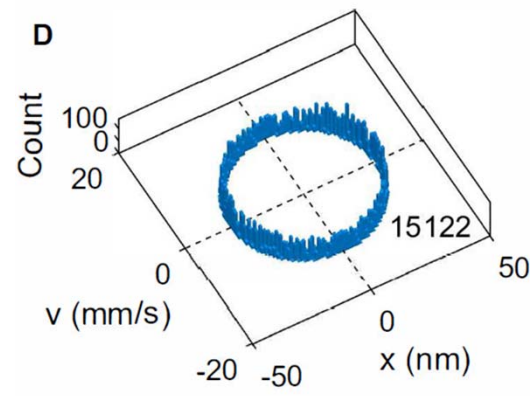
rhs: 0.90



**Uniform  
distribution**

lhs:  $1.42 \pm 0.03$

rhs: 1.42



**Microcanonical  
ensemble**

lhs:  $1.08 \pm 0.02$

rhs: 1.07

# Data in the overdamped regime

- Overdamped regime (760 Torr)

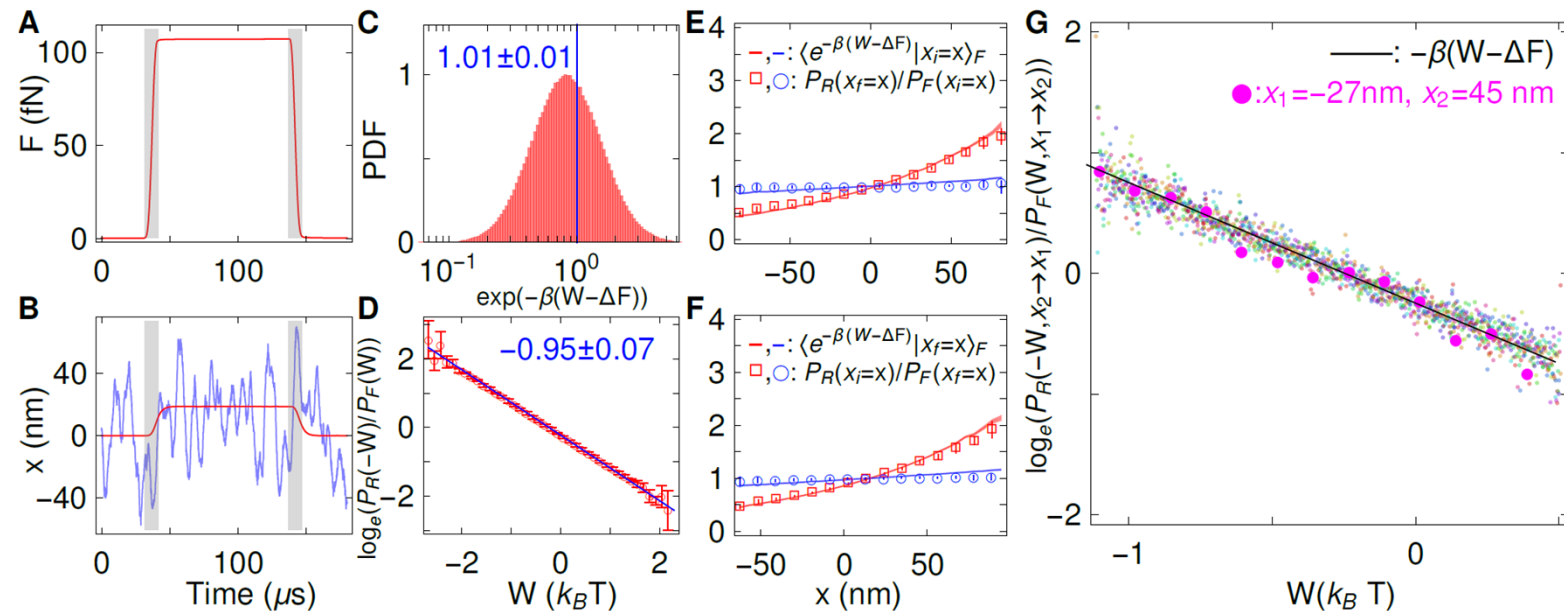
➤  $f_{off} = 0, f_{on} = 107 fN$

➤ Ramp time  $\sim 4.8 \mu s$

➤ Particle size  $r = 145 \pm 5 \text{ nm}$

➤  $\Omega = 76 \pm 3 (2\pi \cdot \text{kHz})$

➤  $\Delta F = -0.24 k_B T$



- [12] Z. Gong and H. T. Quan, Phys. Rev. E 92, 012131 (2015).  
[21] E. H. Feng and G. E. Crooks, Phys. Rev. Lett. 101, 090602 (2008).  
[23] T. Li et al. *Nature Physics* 7, 527 (2011).

# Summary

- Test the DFT and the GJE in both underdamped and overdamped regimes
  - The most detailed fluctuation theorem that can be tested in experiment
  - DFT can unify most of the FTs
  - Length of time's arrow <sup>[21]</sup>
  - Unprecedentedly detailed level
- Technique to generate arbitrary initial states
  - Post-selection of trajectories
- Outlook: extension to quantum regime <sup>[12,23]</sup>



# Collaborators

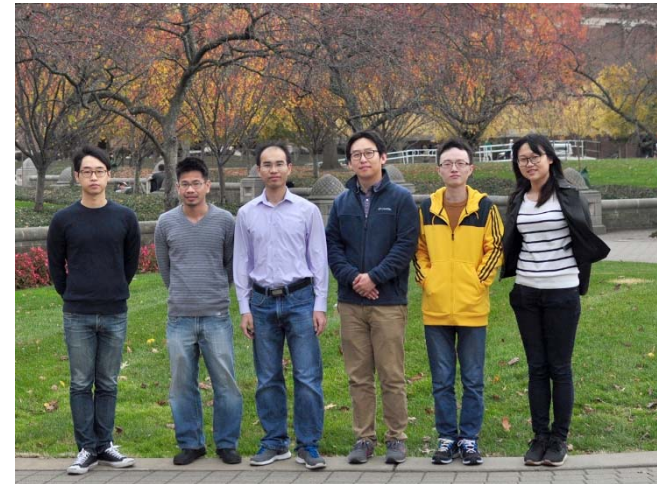
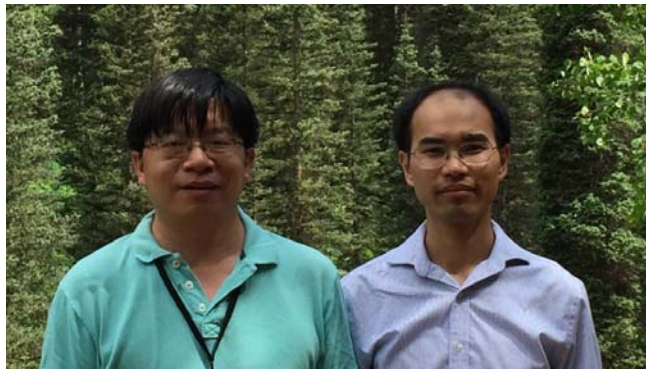
Haitao Quan group (Peking University):

Haitao Quan, Rui Pan



Tongcang Li group (Purdue University):

Tongcang Li, Thai M. Hoang, Jonghoon Ahn, Jaehoon Bang



# References

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- [25] S. Singh, et al. *arXiv:1712.01693* (2017).

**Thanks for your attention!**



# A generalized JE (GJE) for arbitrary initial states

- $e^{-\beta(W-\Delta F)} P_F(W, \Gamma_0 \rightarrow \Gamma_\tau) = P_R(-W, \Gamma_\tau^\dagger \rightarrow \Gamma_0^\dagger)$

$$\iint d\Gamma_0 d\Gamma_\tau: e^{-\beta(W-\Delta F)} P_F(W) = P_R(-W) \quad \text{(CFT)}$$

$$\iint d\Gamma_0 dW: \langle e^{-\beta(W-\Delta F)} | \Gamma(\tau) = \Gamma_\tau \rangle_F = \frac{P_R^{eq}(\tilde{\Gamma}(0)=\Gamma_\tau^\dagger)}{P_F(\Gamma(\tau)=\Gamma_\tau)} \quad \text{(HSR)}$$

A generalized JE (GJE) for delta initial distribution <sup>[11]</sup>:

$$\iint d\Gamma_\tau dW: \langle e^{-\beta(W-\Delta F)} | \Gamma(0) = \Gamma_0 \rangle_F = \frac{P_R(\tilde{\Gamma}(\tau)=\Gamma_0^\dagger)}{P_F^{eq}(\Gamma(0)=\Gamma_0)}$$

A generalized JE (GJE) for arbitrary initial states <sup>[12]</sup>:

$$\langle e^{-\beta(W-\Delta F)} \rangle_{P_{ini}(\Gamma(0)=\Gamma_0)} = \int \frac{P_R(\tilde{\Gamma}(\tau)=\Gamma_0^\dagger)}{P_F^{eq}(\Gamma(0)=\Gamma_0)} P_{ini}(\Gamma(0) = \Gamma_0) d\Gamma_0$$