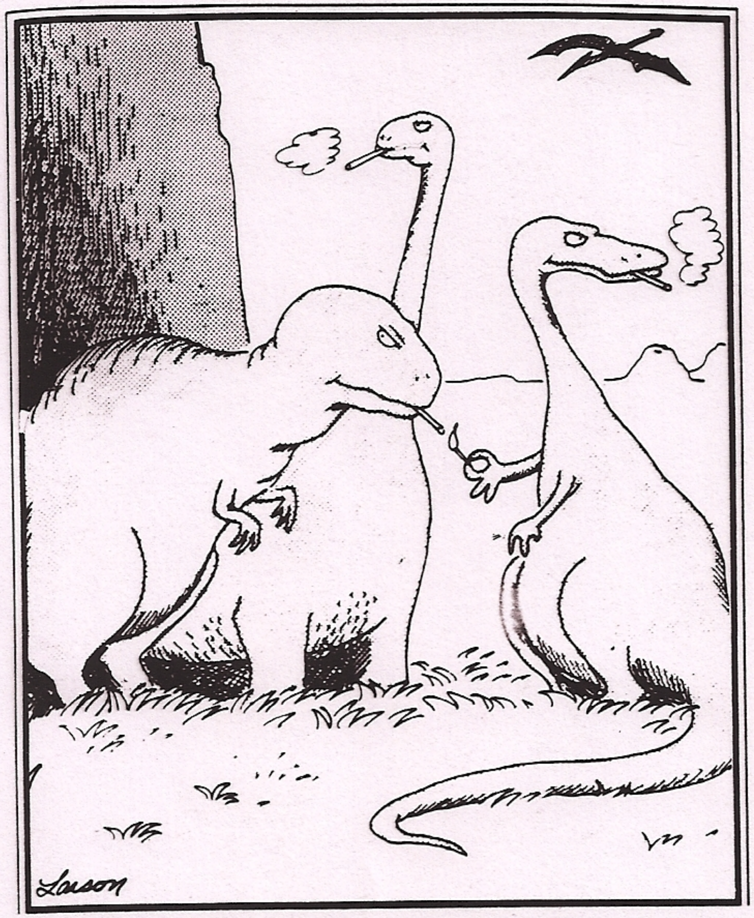
PT-symmetric quantum theory, nonlinear eigenvalue problems, and the Painlevé transcendents



The real reason dinosaurs became extinct

Carl M. Bender
Washington University

"Resurgence in gauge and string theory"
***KITP,* 2017**

The idea of *PT*-symmetric quantum theory:

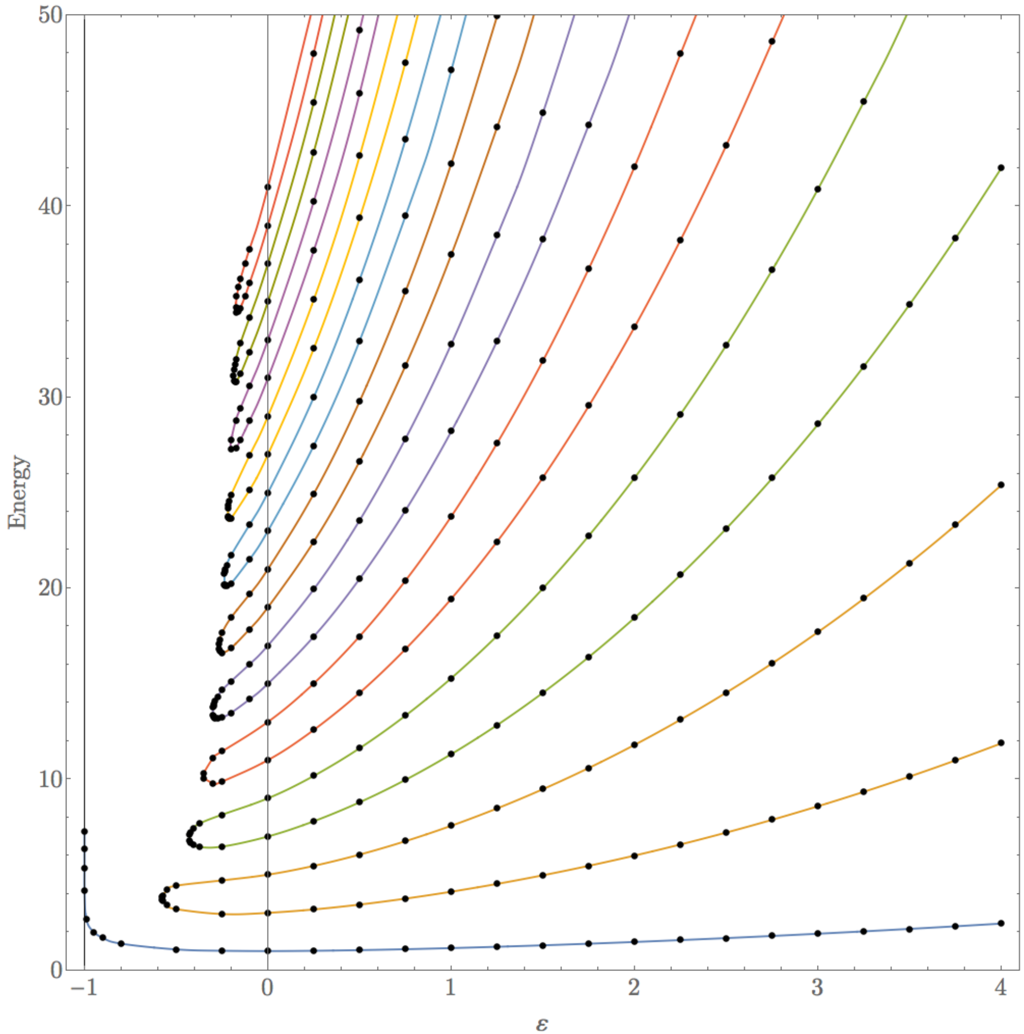
Replace the mathematical condition of Hermiticity by the *weaker* and *physical* condition of *PT* symmetry, where

P = *parity*, ***T*** = *time reversal*

(*Physical* because *P* and *T* are elements of the Lorentz group.)

A class of ***PT***-symmetric Hamiltonians:

$$H = p^2 + x^2(ix)^\varepsilon \quad (\varepsilon \text{ real})$$



Look! The energies are real, positive, and discrete for ε > 0 (!!)

P: $x \rightarrow -x$, $p \rightarrow -p$

T: $x \rightarrow x$, $p \rightarrow -p$, $i \rightarrow -i$

CMB and S. Boettcher
Physical Review Letters **80**, 5243 (1998)

Examples of *PT*-symmetric Hamiltonians

cubic: ε = 1

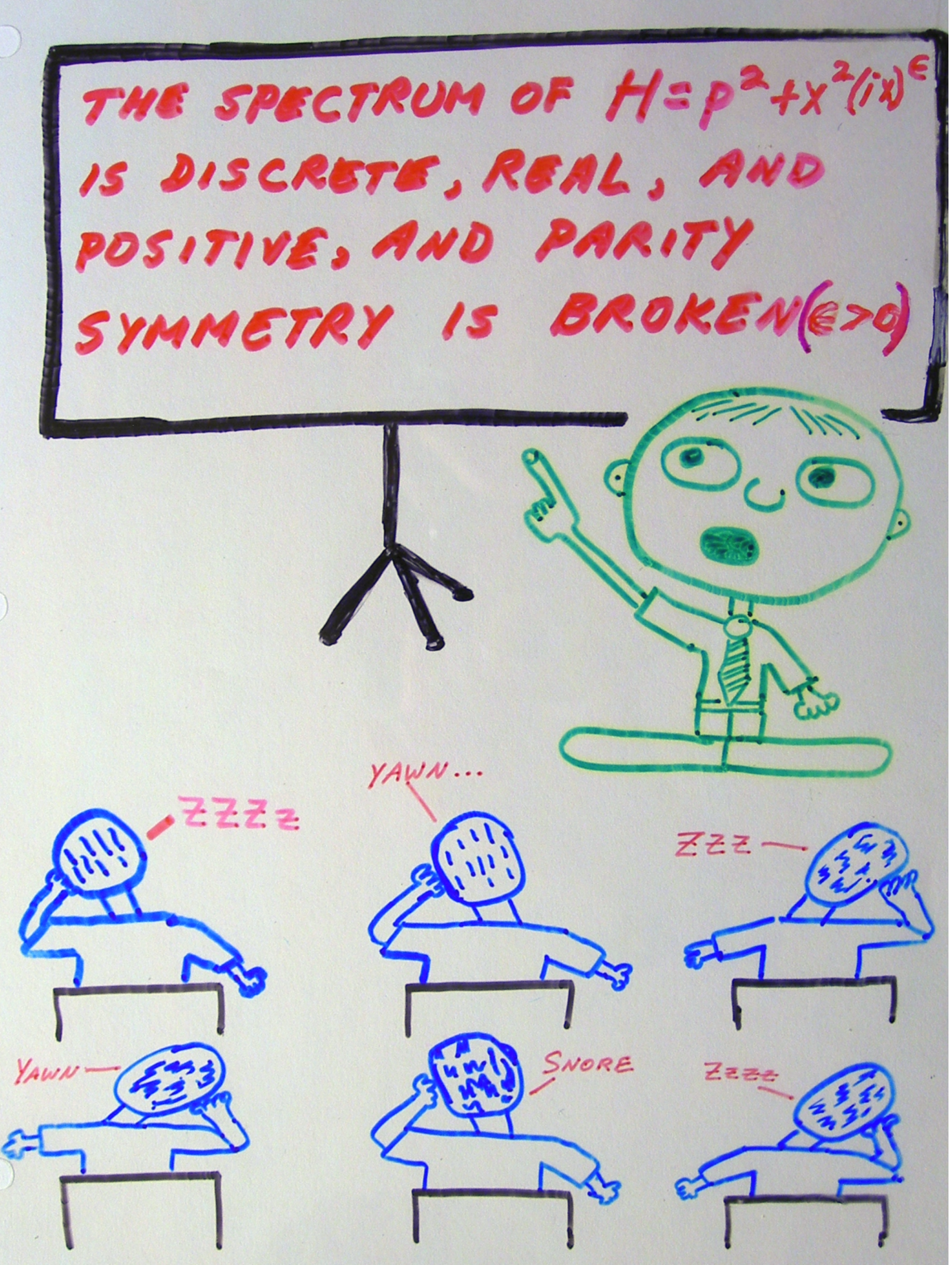
$$H = p^2 + ix^3$$

These Hamiltonians have ***PT* symmetry!**

quartic: ε = 2

$$H = p^2 - x^4$$

← ***An upside-down potential with real positive eigenvalues!***

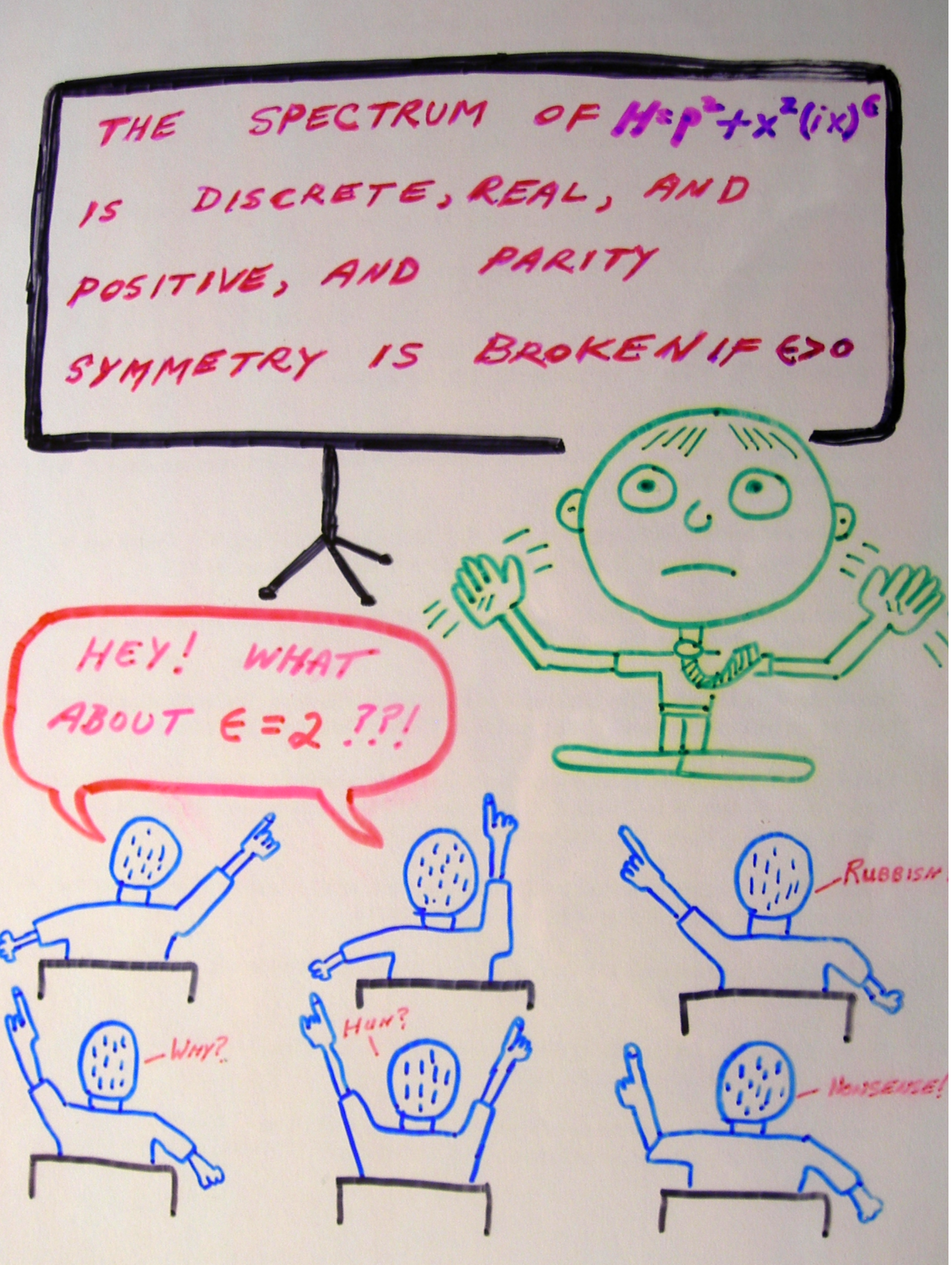


Proof of real eigenvalues:

"ODE/IM Correspondence"
P. Dorey, C. Dunning, and R. Tateo,
J. Phys. A **40**, R205 (2007)

PT **symmetry** ***controls instabilities***

Physical systems that you might *think* are unstable become *stable* in the complex domain…



Upside-down potential with real positive eigenvalues?!

$$V(x) = -x^4$$

Z. Ahmed, CMB, and M. V. Berry,
J. Phys. A: Math. Gen. **38**, L627 (2005)
[arXiv: quant-ph/0508117]

CMB, D. C. Brody, J.-H. Chen, H. F. Jones, K. A. Milton, and M. C. Ogilvie,
Phys. Rev. D **74**, 025016 (2006)
[arXiv: hep-th/0605066]

Stability of the Higgs vacuum:

"*PT*-symmetric interpretation of unstable effective potentials"
CMB, D. W. Hook, N. E. Mavromatos, and S. Sarkar
Journal of Physics A **49**, 45LT01 (2016) [arXiv: hep-th/1506.01970]

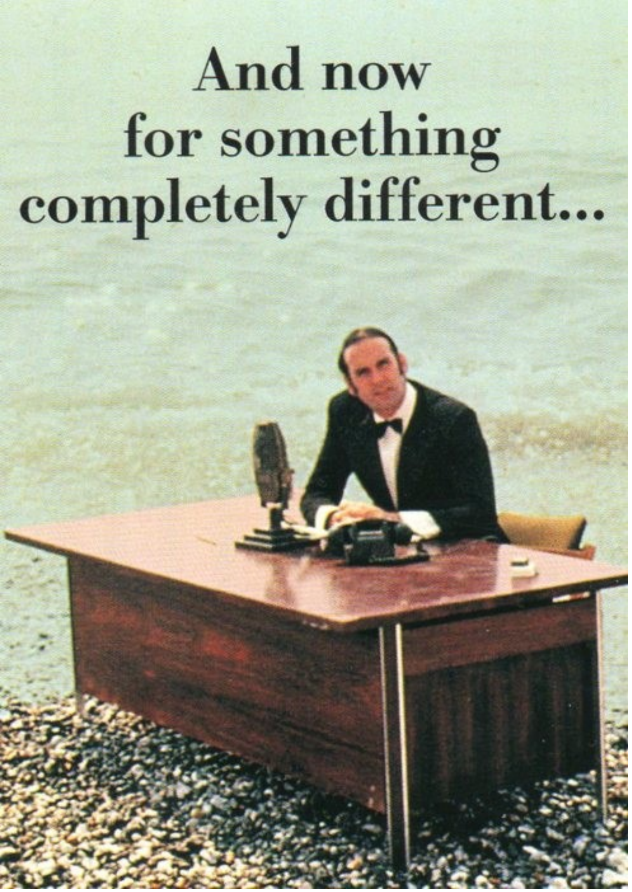
Stability of the double-scaling limit in QM and QFT:

"*PT*-symmetric Interpretation of double-scaling"
CMB, M. Moshe, and S. Sarkar
Journal of Physics A **46**, 102002 (2013) [arXiv: hep-th/1206.4943]

"Double-scaling limit of the O(*N*)-symmetric anharmonic oscillator"
CMB and S. Sarkar
Journal of Physics A **46**, 442001 (2013) [arXiv: hep-th/1307.4348]

Liouville quantum field theory:

"Infinite class of *PT*-symmetric theories from one timelike Liouville Lagrangian"
CMB, D. H. Hook, N. E. Mavromatos, and S. Sarkar
Physical Review Letters **113**, 231605 (2014) [arXiv: hep-th/1408.2432]



Instabilities associated with *nonlinear* eigenvalue problems...

CMB, A. Fring, Q. Wang, and J. Komijani

Linear eigenvalue problems...

$$-\psi''(x) + V(x)\psi(x) = E\psi(x) \qquad \psi(\pm\infty) = 0$$

For *linear* problems *WKB* gives a good approximation for large eigenvalues

$$\int_{x_1}^{x_2} dx\sqrt{E_n - V(x)} \sim (n+1/2)\pi \quad (n\to\infty)$$

*n*th energy level grows like a *constant* times a power of *n*

Example 1: harmonic oscillator

$$V(x) = x^2$$

$$E_n \sim n \quad (n\to\infty)$$

Example 2: anharmonic oscillator

$$V(x) = x^4$$

$$E_n \sim Bn^{4/3} \quad (n\to\infty)$$

$$B = \left[\frac{3\Gamma(3/4)\sqrt{\pi}}{\Gamma(1/4)}\right]^{4/3}$$

WKB **works for** ***PT*****-symmetric Hamiltonians as well:**

$$H = p^2 + x^2(ix)^\varepsilon \qquad (\varepsilon \text{ real})$$

$$E_n \sim \left[\frac{\Gamma\left(\frac{3}{2}+\frac{1}{\varepsilon+2}\right)\sqrt{\pi}\,n}{\sin\left(\frac{\pi}{\varepsilon+2}\right)\Gamma\left(1+\frac{1}{\varepsilon+2}\right)}\right]^{\frac{2\varepsilon+4}{\varepsilon+4}} \qquad (n \to \infty)$$

Asymptotics beyond all orders

Leading asymptotic behavior of solutions to

$$-\psi''(x) + V(x)\psi(x) = E\psi(x)$$

for large positive *x*:

$$\psi(x) \sim C[V(x) - E]^{-1/4} \exp\left[\int^x ds \sqrt{V(s) - E}\right] \quad (x \to \infty)$$

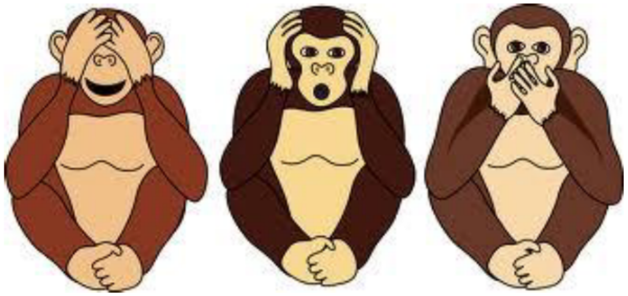
NOTE: There is only *ONE* arbitrary constant.

Second arbitrary constant is invisible with Poincaré asymptotics because it is contained in the *subdominant* solution:

$$\psi(x) \sim D[V(x) - E]^{-1/4} \exp\left[-\int^x ds \sqrt{V(s) - E}\right] \quad (x \to \infty)$$

***Physical* solution is *Unstable* under small changes in *E*.**

Eigenfunctions: 3 characteristic properties



(1) Oscillatory in *classically allowed* region (*n*th eigenfunction has *n* nodes)

(2) Monotone decay in *classically forbidden* region

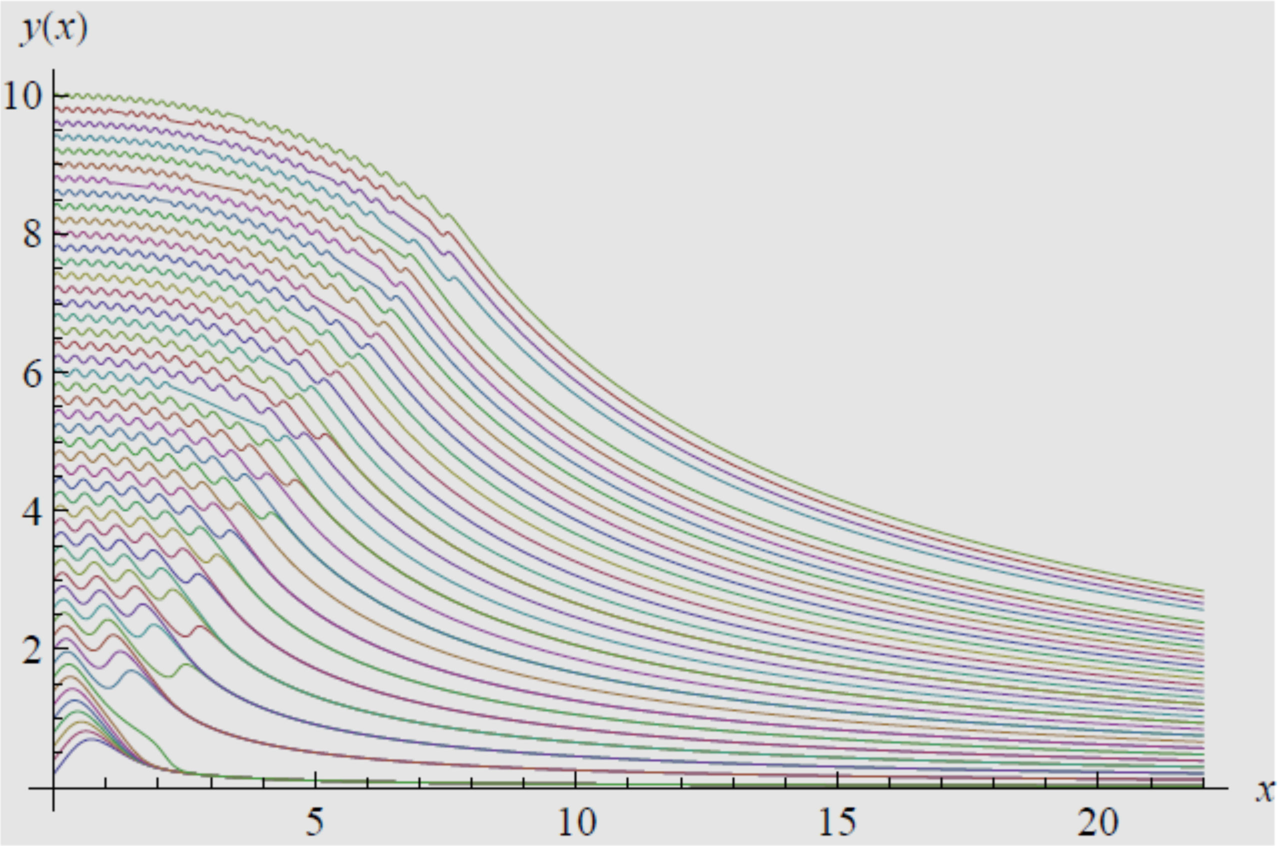
(3) Transition at the boundary (*turning point*)

Toy nonlinear eigenvalue problem

$$y'(x) = \cos[\pi x y(x)], \quad y(0) = a$$

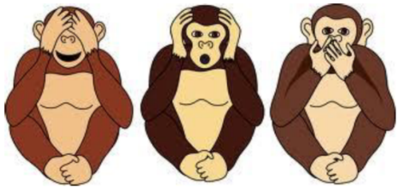
Some references:

[1] C. M. Bender and S. A. Orszag, *Advanced Mathematical Methods for Scientists and Engineers* (McGraw Hill, New York, 1978), chap. 4.
[2] C. M. Bender, D. W. Hook, P. N. Meisinger, and Q. Wang, Phys. Rev. Lett. **104**, 061601 (2010).
[3] C. M. Bender, D. W. Hook, P. N. Meisinger, and Q. Wang, Ann. Phys. **325**, 2332-2362 (2010).
[4] J. Gair, N. Yunes, and C. M. Bender, J. Math. Phys. **53**, 032503 (2012).



Solutions for 50 initial conditions

Note: (1) oscillation (2) monotone decay (3) transition

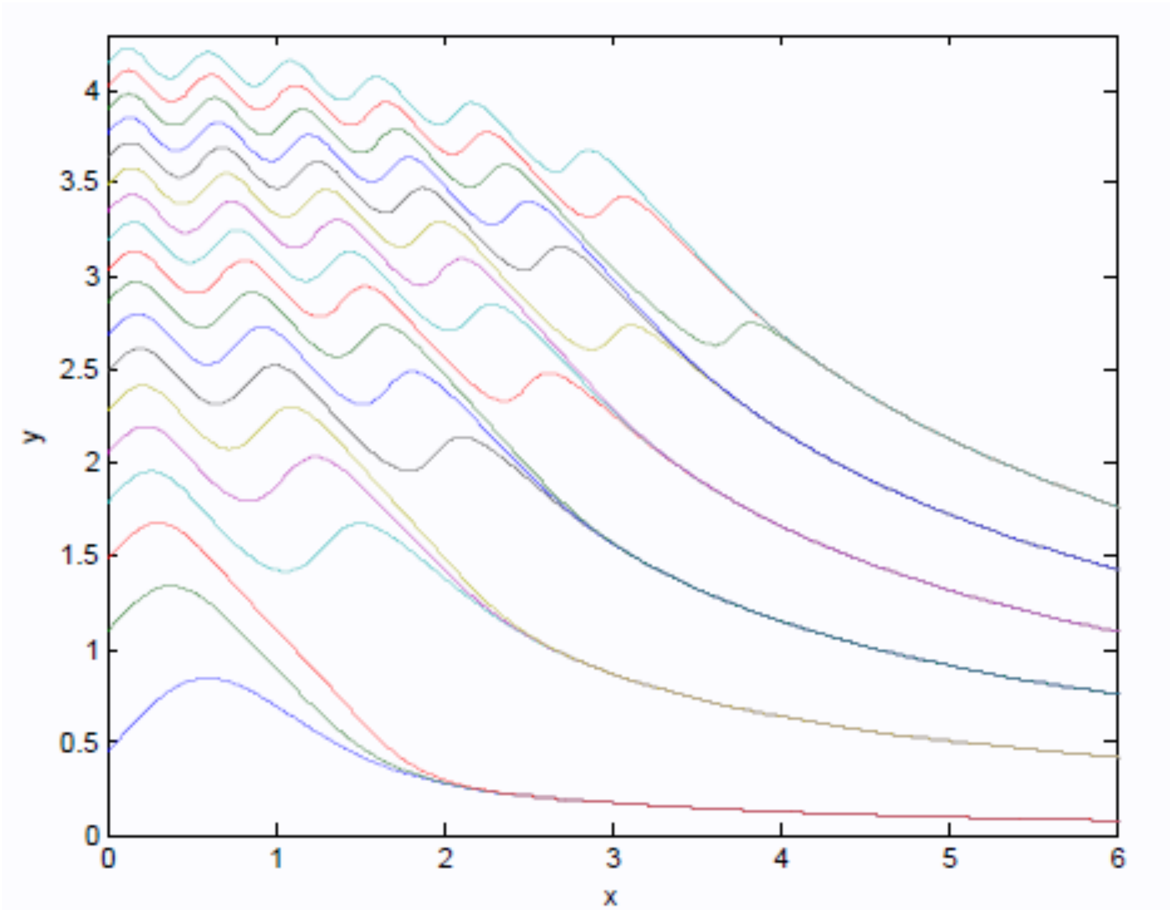


Asymptotic behavior for large x

Solution behaves like: $$y(x) \sim \frac{m + 1/2}{x}$$

***m* = 0, 1, 2, 3, ... is an integer**

There's a ***big*** problem here...



m **= 10**
m **= 8**
m **= 6**
m **= 4**
m **= 2**
m **= 0**

Where are the odd-*m* solutions?!?

Furthermore, no arbitrary constant appears in the asymptotic behavior!!

Where is the arbitrary constant?!?

Is it in higher order?

Higher-order asymptotic behavior for large *x* still contains <u>no arbitrary constant</u>!

$$y(x) \sim \frac{m+1/2}{x} + \sum_{k=1}^{\infty} \frac{c_k}{x^{2k+1}} \quad (x \to \infty)$$

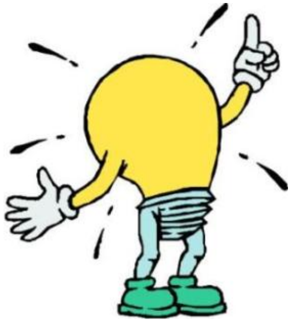
$$\begin{aligned}
c_1 &= \frac{(-1)^m}{\pi}(m+1/2), \\
c_2 &= \frac{3}{\pi^2}(m+1/2), \\
c_3 &= (-1)^m \left[\frac{(m+1/2)^3}{6\pi} + \frac{15(m+1/2)}{\pi^3}\right], \\
c_4 &= \frac{8(m+1/2)^3}{3\pi^2} + \frac{105(m+1/2)}{\pi^4}, \\
c_5 &= (-1)^m \left[\frac{3(m+1/2)^5}{40\pi} + \frac{36(m+1/2)^3}{\pi^3} + \frac{945(m+1/2)}{\pi^5}\right], \\
c_6 &= \frac{38(m+1/2)^5}{15\pi^2} + \frac{498(m+1/2)^3}{\pi^4} + \frac{10395(m+1/2)}{\pi^6}.
\end{aligned}$$

Asymptotics beyond all orders

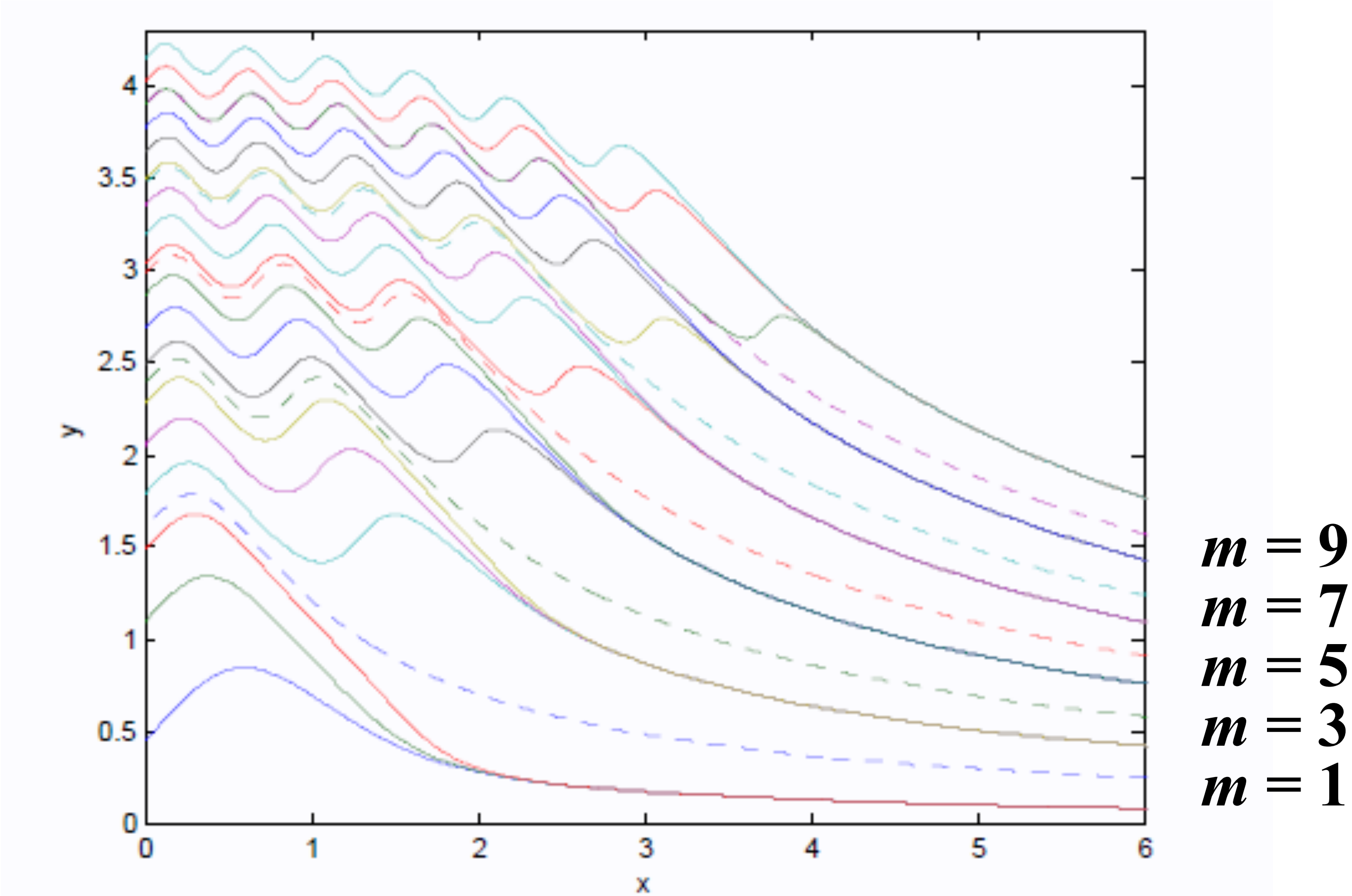
Difference of two solutions in one bundle: $Y(x) \equiv y_1(x) - y_2(x)$

$$\begin{aligned} Y'(x) &= \cos[\pi x y_1(x)] - \cos[\pi x y_2(x)] \\ &= -2\sin\left[\tfrac{1}{2}\pi x y_1(x) + \tfrac{1}{2}\pi x y_2(x)\right]\sin\left[\tfrac{1}{2}\pi x y_1(x) - \tfrac{1}{2}\pi x y_2(x)\right] \\ &\sim -2\sin\left[\pi\left(m+\tfrac{1}{2}\right)\right]\sin\left[\tfrac{1}{2}\pi x Y(x)\right] \quad (x\to\infty) \\ &\sim -(-1)^m \pi x Y(x) \quad (x\to\infty). \end{aligned}$$

$$Y(x) \sim K \exp\left[-(-1)^m \pi x^2\right] \quad (x\to\infty)$$



Aha! *K* is the invisible arbitrary constant!
Odd-*m* solutions are *unstable*;
even-*m* solutions are *stable*.



$$y(0) = a \in \{1.6026, 2.3884, 2.9767, 3.4675, 3.8975, 4.2847, \ldots\}$$

Eigenvalues **correspond to odd-*m* initial values.**
Eigenfunctions **are (*unstable*) *separatrices*, which begin at eigenvalues.**

We calculated up to m=500,001

Let $m = 2n - 1$

For large n the nth eigenvalue grows like the *square root* of n times a constant A, and we used Richardson extrapolation to show that

A= 1.7817974363...

and then we <u>guessed</u> A.

Result:

$$a_n \sim A\sqrt{n} \quad (n \to \infty)$$

$$A = 2^{5/6}$$

This is a rather nontrivial problem...

Analytic calculation of the constant A

Construct moments of $z(t)$:

$$A_{n,k}(t) \equiv \int_0^t ds \cos[n\lambda s z(s)] \frac{s^{k+1}}{[z(s)]^k}$$

Moments are associated with a semi-infinite linear one-dimensional random walk in which random walkers become static as they reach n=1

$$2\alpha_{1,k} + \alpha_{2,k-1} = 0, \qquad 2\alpha_{n,k} + \alpha_{n-1,k-1} + \alpha_{n+1,k-1} = 0 \quad (n \geq 3).$$

$$2\alpha_{2,k} + \alpha_{3,k-1} = 0,$$

Solve the random walk problem exactly and get $A = 2^{5/6}$

CMB, A. Fring, and J. Komijani
J. Phys. A: Math. Theor. **47**, 235204 (2014)
[arXiv: math-ph/1401.6161]

Possible connection with the *power series constant P*???

(Remember the numerical constant $A = 1.7818$)

W. K. Hayman, *Research Problems in Function theory*
[Athlone Press (University of London), London, 1967]

J. Clunie and P. Erdös, Proc. Roy. Irish Acad. **65**, 113 (1967).
J. D. Buckholtz, Michigan Math. J. **15**, 481 (1968).

$$1 \le P \le 2$$

$$\sqrt{2} \le P \le 2$$

$$1.7 \le P \le 12^{1/4}$$

$$1.7818 \le P \le 1.82$$

Three ***nontrivial*** **second-order nonlinear eigenvalue problems**

separatrix

Painlevé equations

Paul Painlevé (1863-1933)

Six Painlevé equations known as Painlevé I – VI

Only spontaneous singularities are poles

Painlevé I $$\frac{d^2y}{dt^2} = 6y^2 + t$$

Painlevé II $$\frac{d^2y}{dt^2} = 2y^3 + ty + \alpha$$

Painlevé III $$ty\frac{d^2y}{dt^2} = t\left(\frac{dy}{dt}\right)^2 - y\frac{dy}{dt} + \delta t + \beta y + \alpha y^3 + \gamma t y^4$$

Painlevé IV $$y\frac{d^2y}{dt^2} = \tfrac{1}{2}\left(\frac{dy}{dt}\right)^2 + \beta + 2(t^2 - \alpha)y^2 + 4ty^3 + \tfrac{3}{2}y^4$$

Painlevé V $$\frac{d^2y}{dt^2} = \left(\frac{1}{2y} + \frac{1}{y-1}\right)\left(\frac{dy}{dt}\right)^2 - \frac{1}{t}\frac{dy}{dt} + \frac{(y-1)^2}{t^2}\left(\alpha y + \frac{\beta}{y}\right) + \gamma\frac{y}{t} + \delta\frac{y(y+1)}{y-1}$$

Painlevé VI $$\frac{d^2y}{dt^2} = \tfrac{1}{2}\left(\frac{1}{y} + \frac{1}{y-1} + \frac{1}{y-t}\right)\left(\frac{dy}{dt}\right)^2 - \left(\frac{1}{t} + \frac{1}{t-1} + \frac{1}{y-t}\right)\frac{dy}{dt} + \frac{y(y-1)(y-t)}{t^2(t-1)^2}\left(\alpha + \beta\frac{t}{y^2} + \gamma\frac{t-1}{(y-1)^2} + \delta\frac{t(t-1)}{(y-t)^2}\right)$$

(1) First Painlevé transcendent

$$y''(t) = 6[y(t)]^2 + t, \qquad y(0) = b,\ y'(0) = c$$

Solution *y*(*x*) must *choose* between two possible asymptotic behaviors as *x* gets large and negative:

$$+\sqrt{-t/6} \text{ or } -\sqrt{-t/6}.$$

Example of a *difficult* choice ...

Two possible asymptotic behaviors

Lower square-root branch is *stable*:

$$y(x) \sim -\sqrt{-x} + c(-x)^{-1/8} \cos\left[\tfrac{4}{5}\sqrt{2}(-x)^{5/4} + d\right] \quad (x \to -\infty)$$

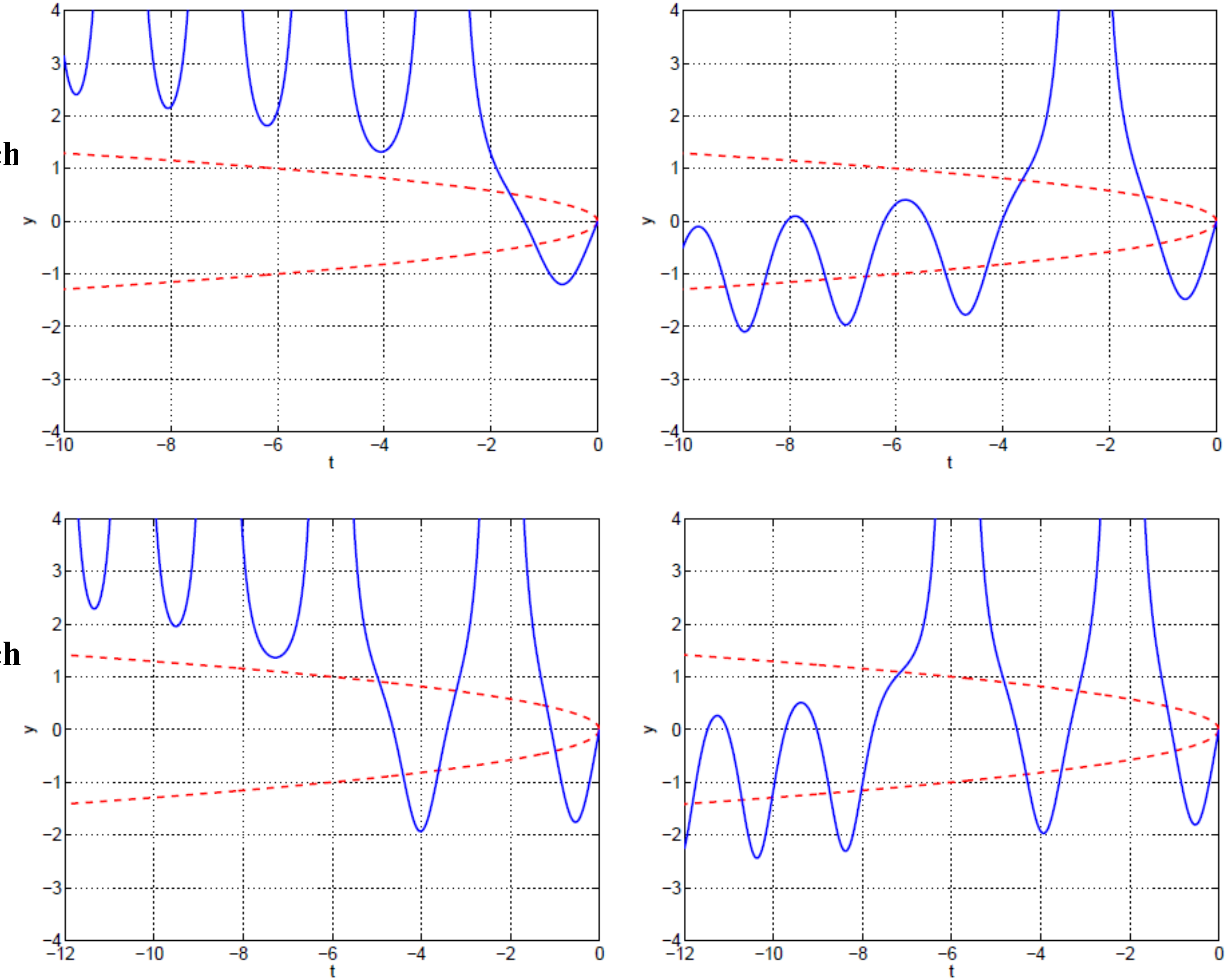
Upper square-root branch is *unstable*:

$$y(x) \sim \sqrt{-x} + c_{\pm}(-x)^{-1/8} \exp\left[\pm\tfrac{4}{5}\sqrt{2}(-x)^{5/4}\right] \quad (x \to -\infty)$$

Two possible kinds of solutions (NOT eigenfunctions):

Unstable branch

Stable branch



Unstable branch

Stable branch

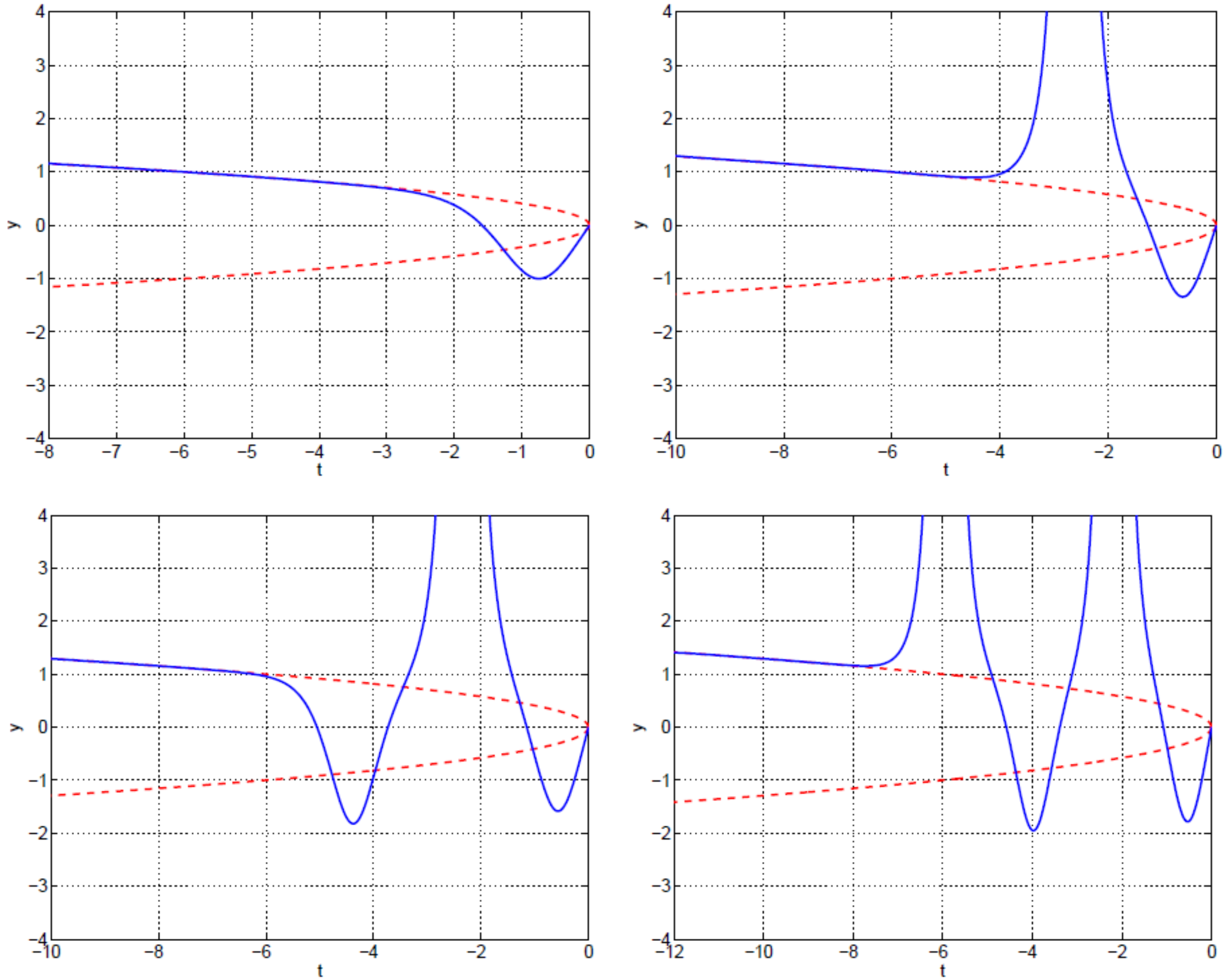
First four separatrix (eigenfunction) solutions:

Unstable

Stable

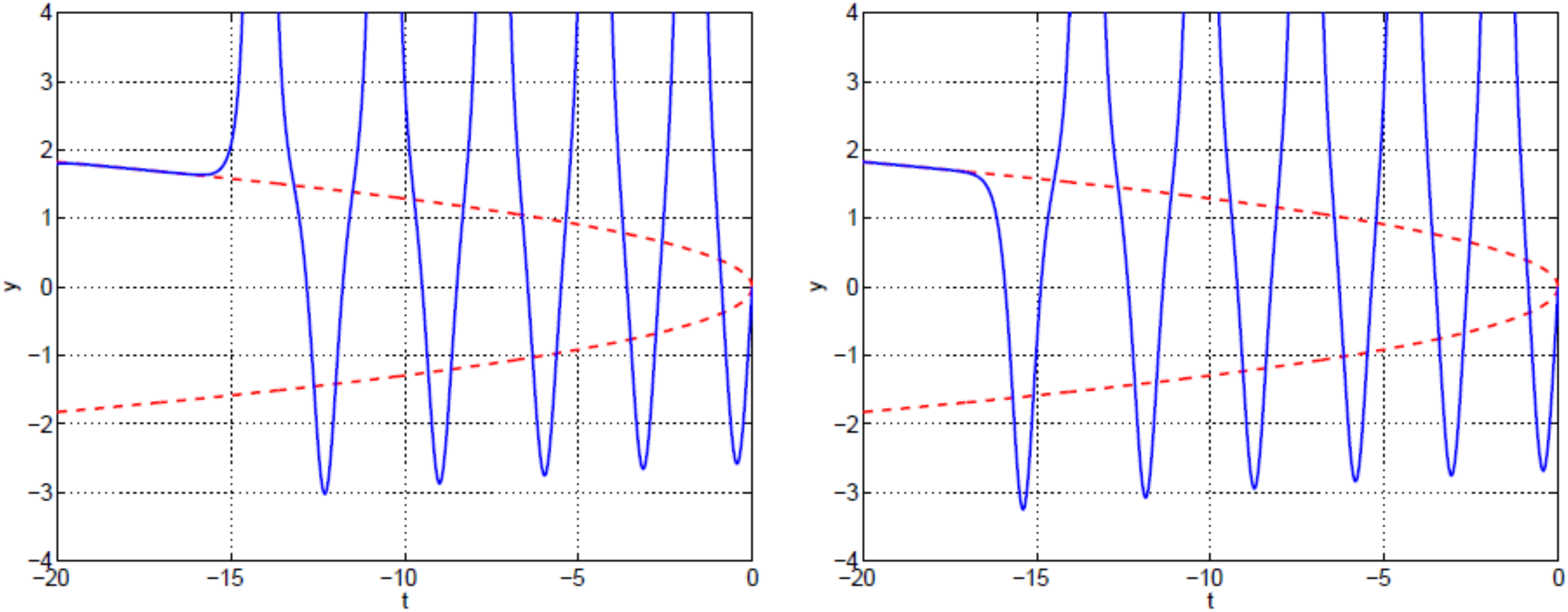
Unstable

Stable



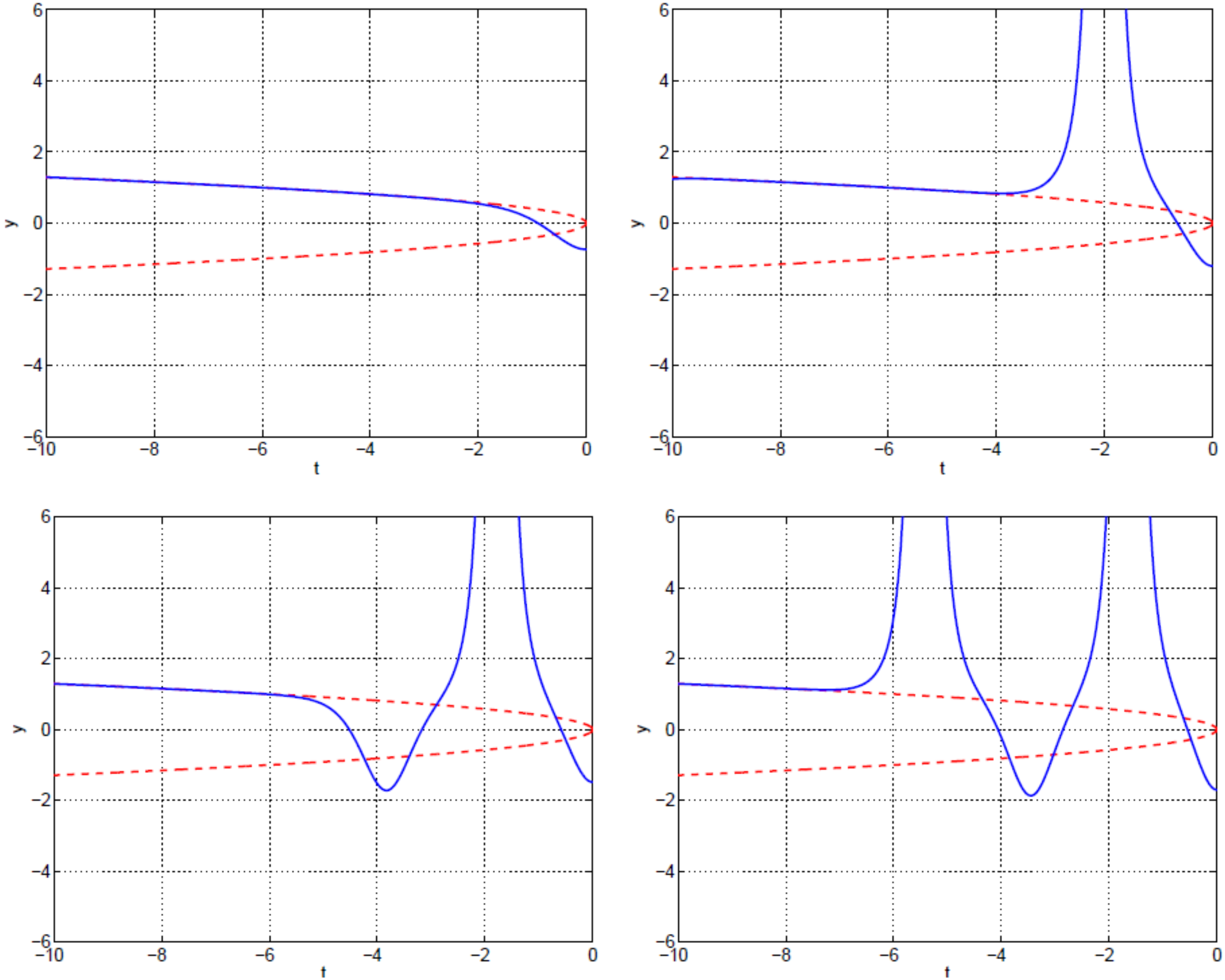
***Initial slope is the eigenvalue*, initial value $y(0) = 0$**

Tenth and eleventh separatrix (eigenfunction) solutions:



***Initial slope* is the eigenvalue, initial value *y*(0) = 0**

First four separatrix solutions with 0 initial slope:



Numerical calculation of eigenvalues

(*nonlinear* semiclassical large-*n* limit)

$$y'(0) = b_n \qquad y(0) = 0$$

$$b_n \sim B_{\mathrm{I}} n^{3/5} \qquad B_{\mathrm{I}} = 2.0921467\underline{4}$$

$$y(0) = c_n \qquad y'(0) = 0$$

$$c_n \sim C_{\mathrm{I}} n^{2/5} \qquad C_{\mathrm{I}} = -1.030484\underline{4}$$

$$y_n(0) \sim -1.030\,484\,4 \left(n - \frac{1}{2}\right)^{\frac{2}{5}} \left[1 - \frac{0.009\,651\,8}{\left(n - \frac{1}{2}\right)^2} + \frac{0.0240}{\left(n - \frac{1}{2}\right)^4}\right]$$

Analytical asymptotic calculation of eigenvalues

$$B_{\mathrm{I}} = 2\left[\frac{5\sqrt{\pi}\,\Gamma(5/6)}{2\sqrt{3}\,\Gamma(1/3)}\right]^{3/5}$$

$$C_{\mathrm{I}} = -\left[\frac{5\sqrt{\pi}\,\Gamma(5/6)}{2\sqrt{3}\,\Gamma(1/3)}\right]^{2/5}$$

Obtained by using WKB to calculate the large eigenvalues of the

cubic PT-symmetric Hamiltonian

$$H = \frac{1}{2}p^2 + 2ix^3$$

Painlevé I corresponds to ε = 1

(Do you remember the cubic *PT*-symmetric Hamiltonian?!)

Analytical asymptotic calculation of eigenvalues

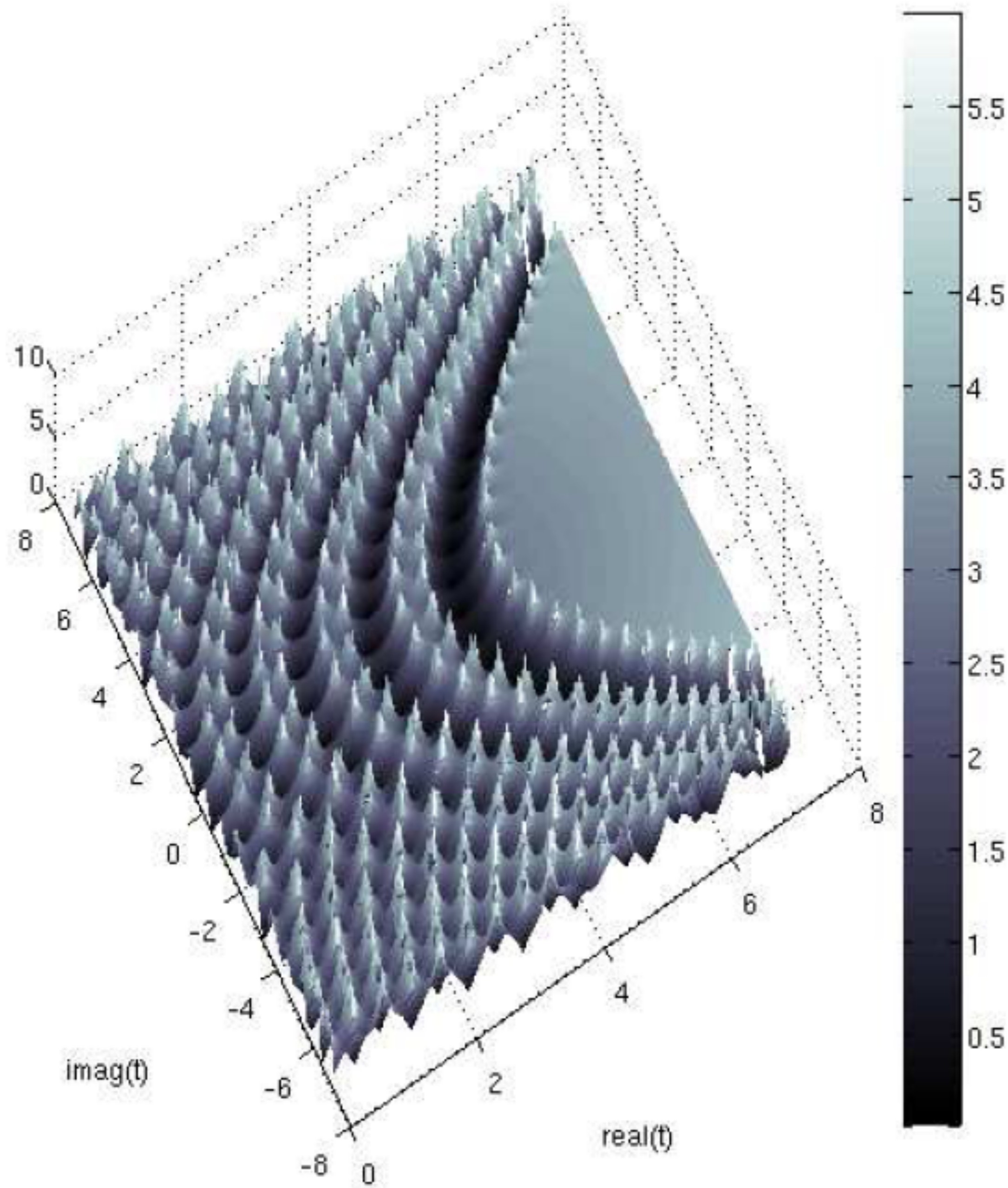
Multiply Painlevé I equation by *y'(t)*;
Then integrate from $t = 0$ to $t = x$:

$$H \equiv \frac{1}{2}[y'(x)]^2 - 2[y(x)]^3 = \frac{1}{2}[y'(0)]^2 - 2[y(0)]^3 + I(x),$$

where $I(x) = \int_0^x dt\, t y'(t).$

Take $|x|$ large at an angle of $\pi/4$, $I(x) \rightarrow 0$, and we get the *PT*-symmetric Hamiltonian for $\varepsilon = 1$.

D. Masoero noted connections between Painlevé I and $H = p^2 + ix^3$



10
5
0
8
6
4
2
0
-2
-4
-6
-8
imag(t)
0
2
4
6
8
real(t)
5.5
5
4.5
4
3.5
3
2.5
2
1.5
1
0.5

(2) Second Painlevé transcendent

$$y''(t) = 2[y(t)]^3 + ty(t), \qquad y(0) = b,\ y'(0) = c$$

Now, both solutions

$$+\sqrt{-t/2} \text{ or } -\sqrt{-t/2}$$

are unstable and 0 is stable.

Two types of solutions (not eigenfunctions):

Unstable

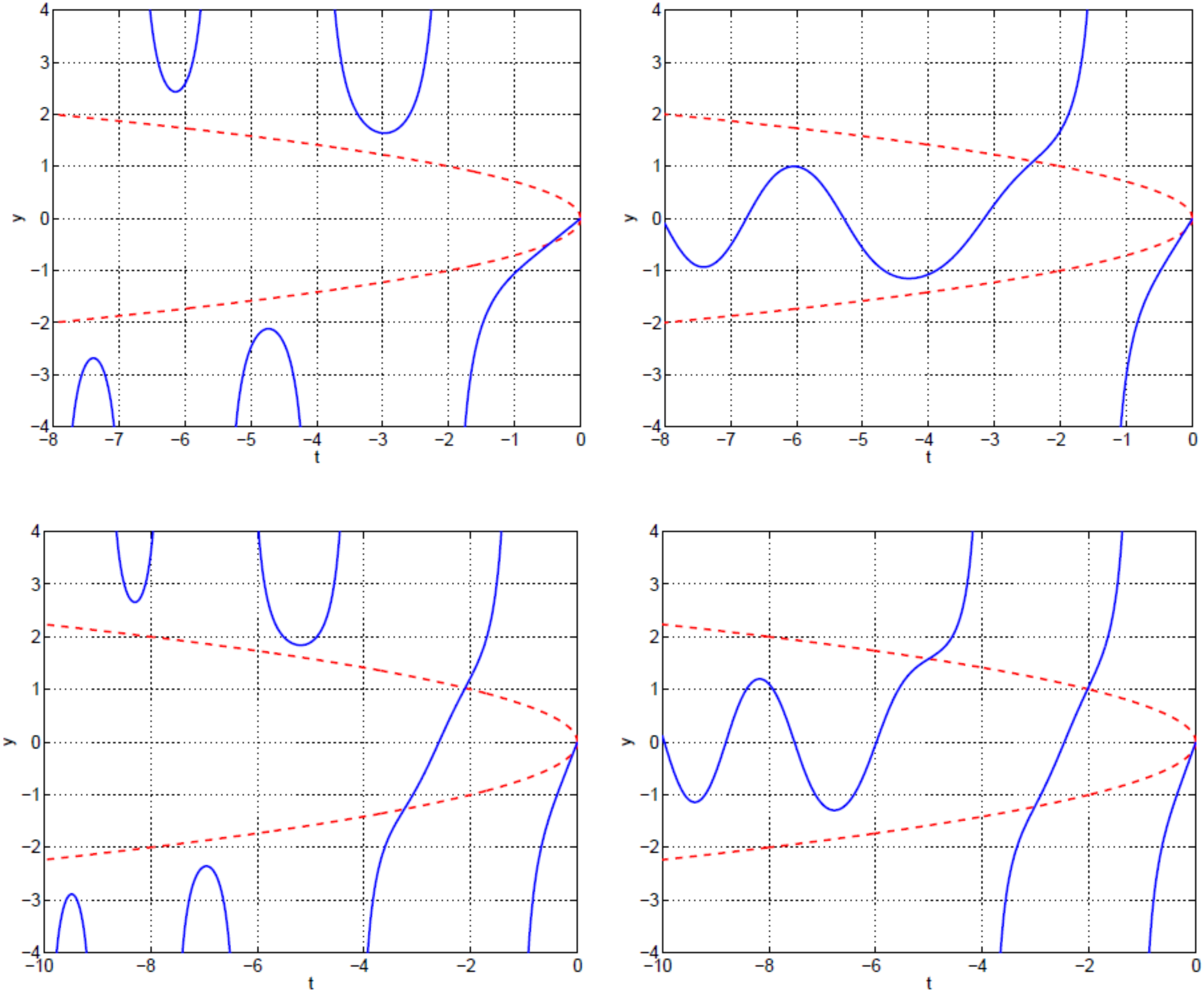
Stable

Unstable

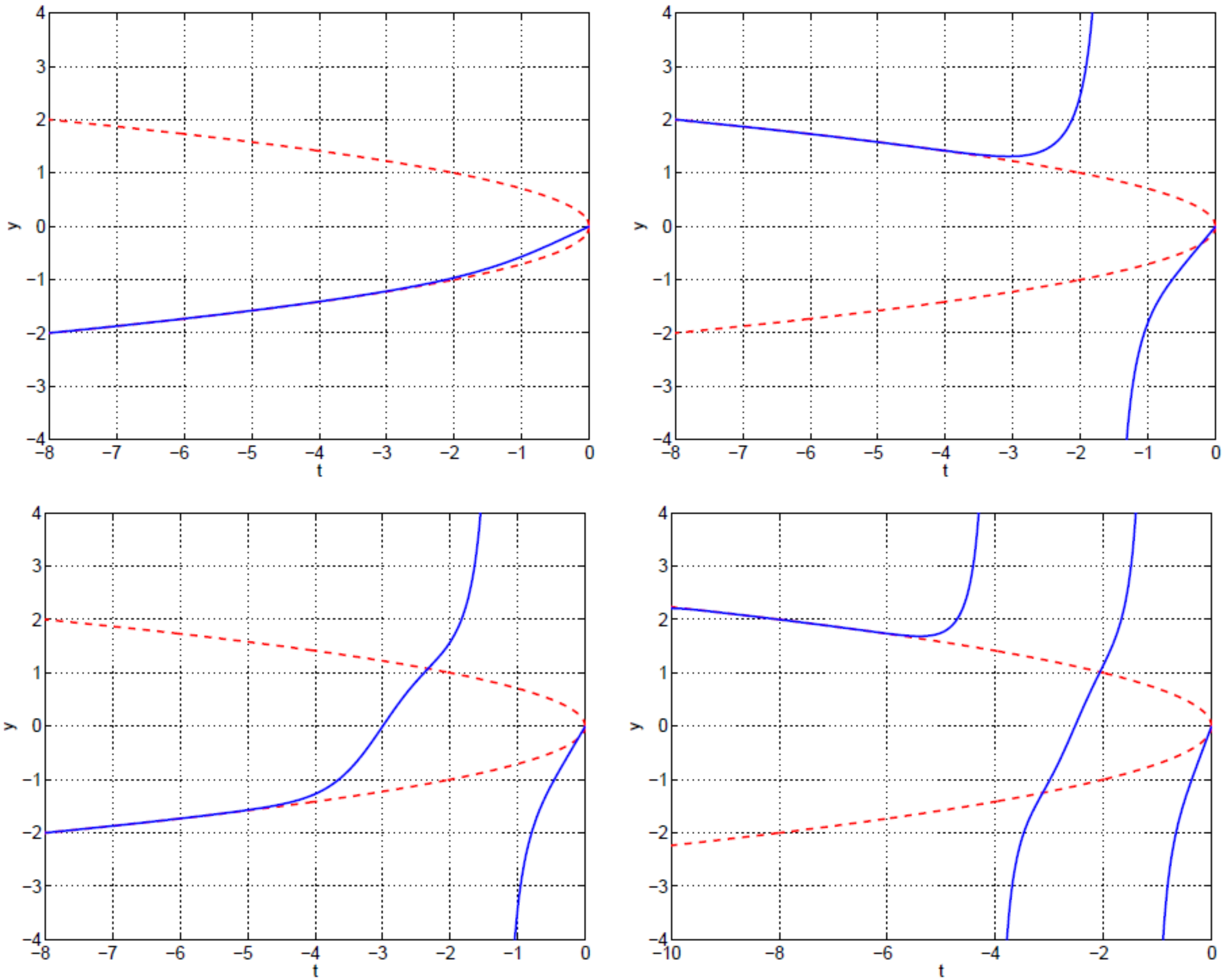
Unstable

Stable

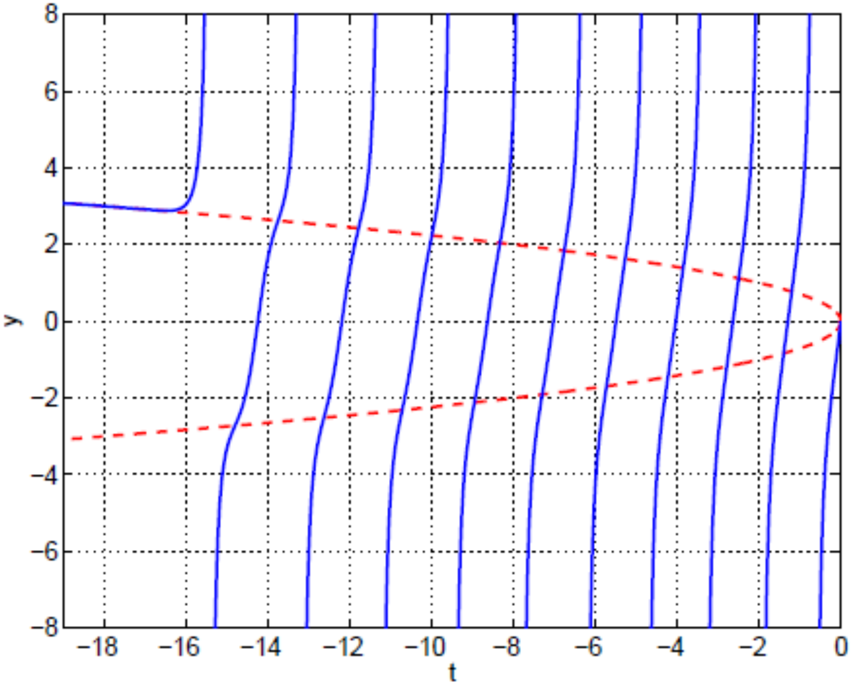
Unstable

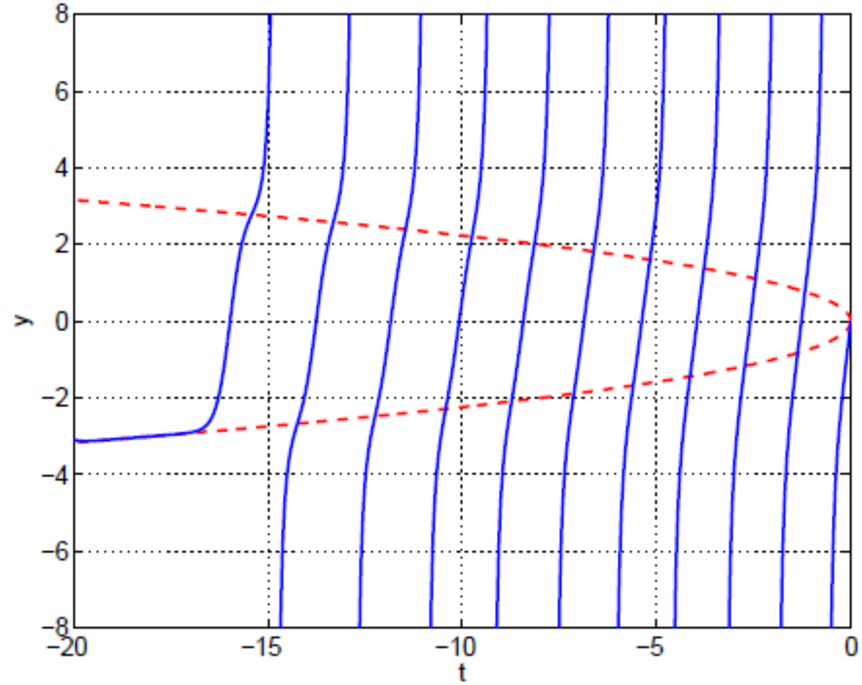


First four separatrix solutions with y(0)=0:

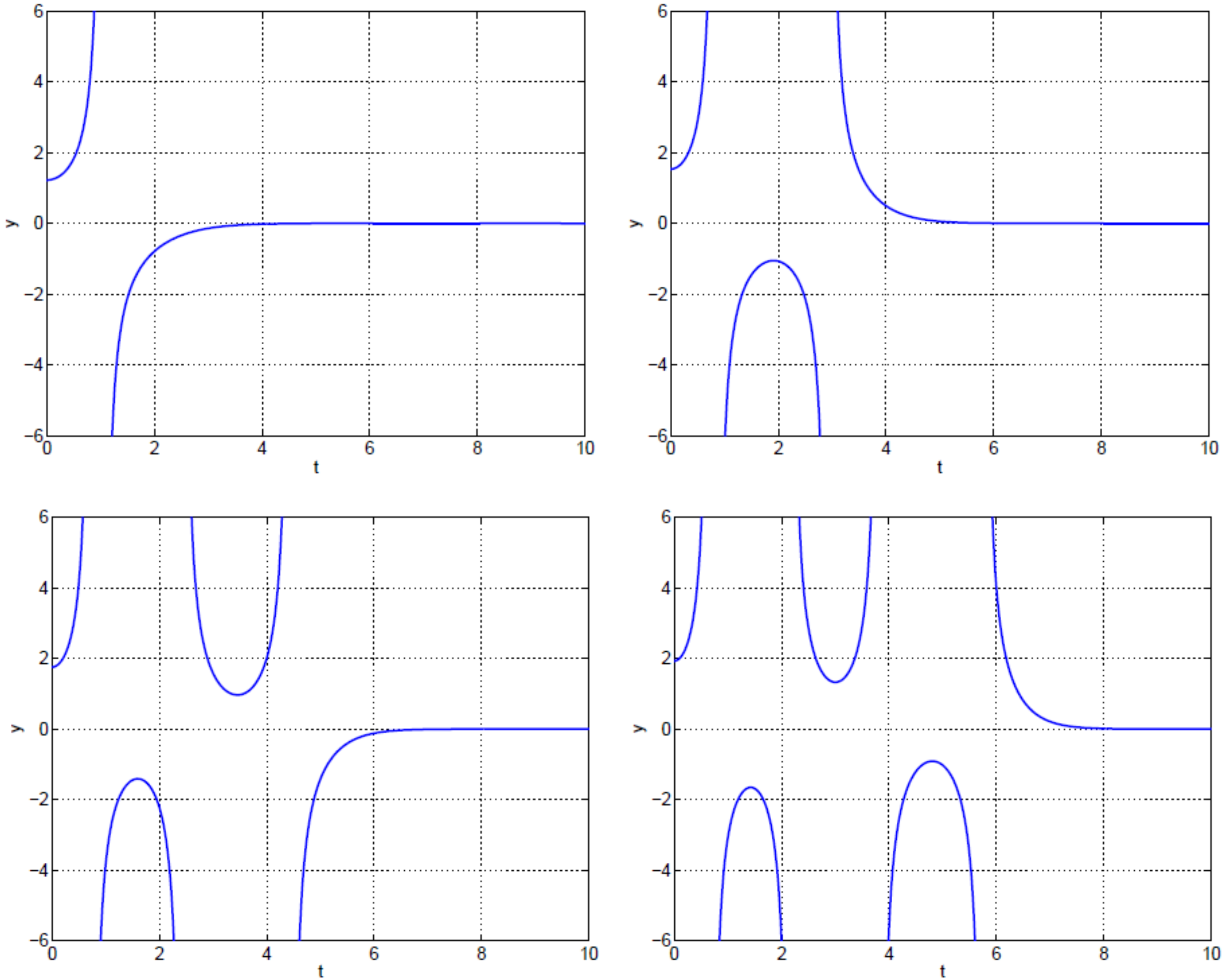


20^{th} and 21^{st} separatrix solutions:





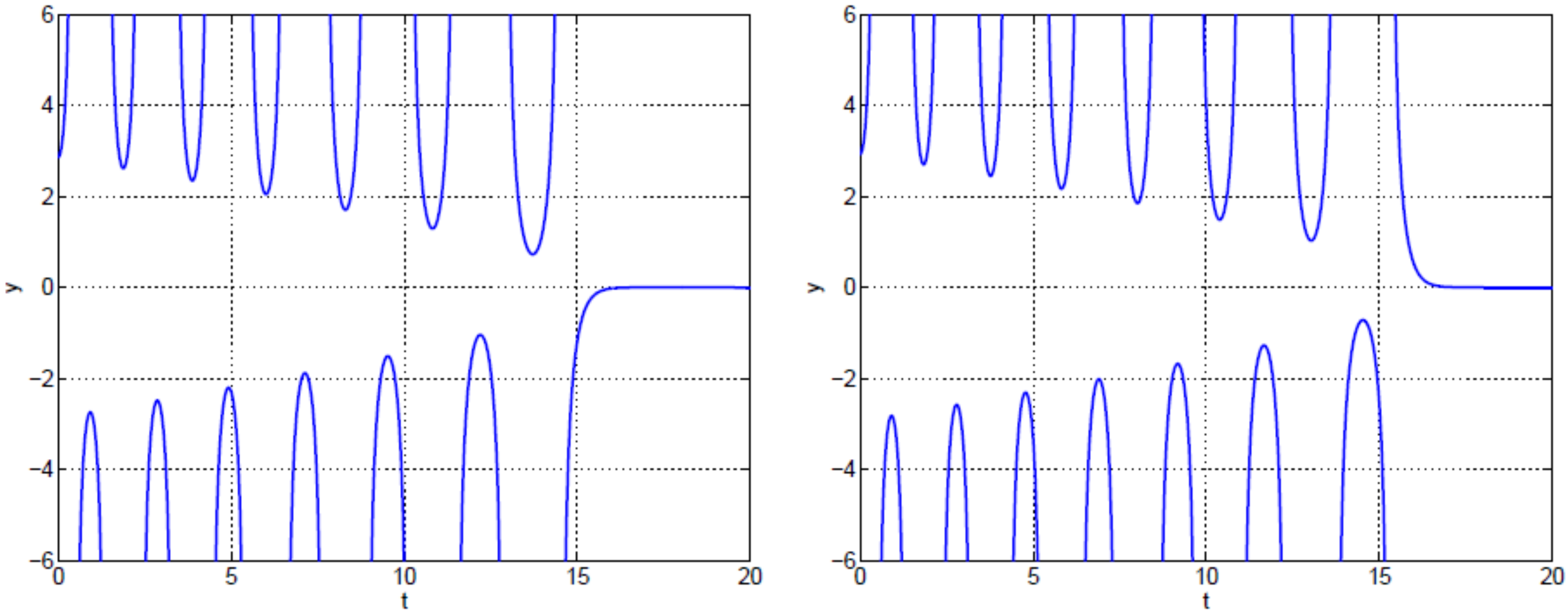
First four separatrices with vanishing initial slope $y'(0)=0$:



Unstable

Unstable

13th and 14th separatrices:



Numerical calculation of eigenvalues

$$y(0) = 0,\ b_n = y'(0)$$

$$c_n = y(0),\ y'(0) = 0$$

$$b_n \sim B_{\rm II} n^{2/3} \quad \text{and} \quad c_n \sim C_{\rm II} n^{1/3}$$

$$B_{\rm II} = 1.862412\underline{8} \qquad C_{\rm II} = 1.2158116\underline{5}$$

$$y_n(0) \sim 1.215\,811\,7\, n^{\frac{1}{3}} \left[1 + \frac{0.005\,254\,3}{n^2} + \frac{0.077}{n^4}\right]$$

CMB and J. Komijani
J. Physics A: Math. Theor. **48**, 475202 (2015)

Analytical calculation of eigenvalues

$$B_{\rm II} = \left[3\sqrt{2\pi}\,\Gamma\left(\tfrac{3}{4}\right)/\Gamma\left(\tfrac{1}{4}\right)\right]^{2/3}$$

$$C_{II} = \left[3\sqrt{\pi}\,\Gamma\left(\tfrac{3}{4}\right)/\Gamma\left(\tfrac{1}{4}\right)\right]^{1/3}$$

Obtained by using WKB to calculate the large eigenvalues of the

quartic PT-symmetric Hamiltonian

$$H = \frac{1}{2}p^2 - \frac{1}{2}x^4$$

Painlevé II corresponds to ε = 2

(Do you remember the quartic upside-down *PT*-symmetric Hamiltonian?!)

(3) Fourth Painlevé transcendent

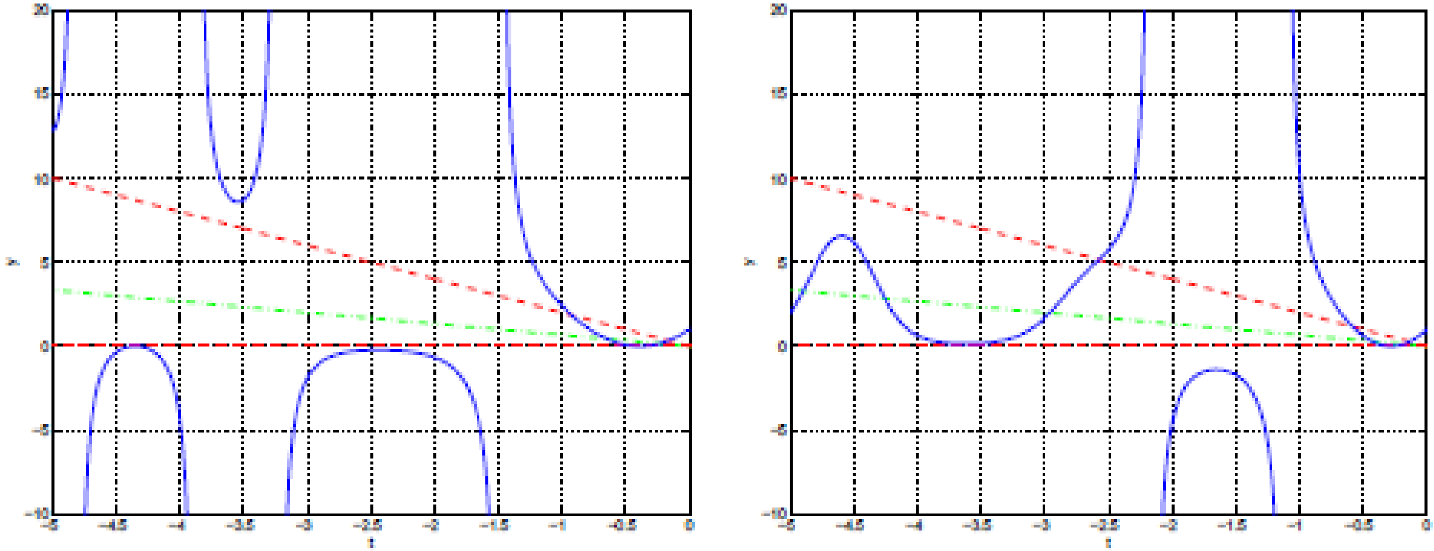
$$y(t)y''(t) = \frac{1}{2}[y'(t)]^2 + 2t^2[y(t)]^2 + 4t[y(t)]^3 + \frac{3}{2}[y(t)]^4$$

with $y(0) = c$ and $y'(0) = b$.

Two possible kinds of solutions (NOT eigenfunctions):

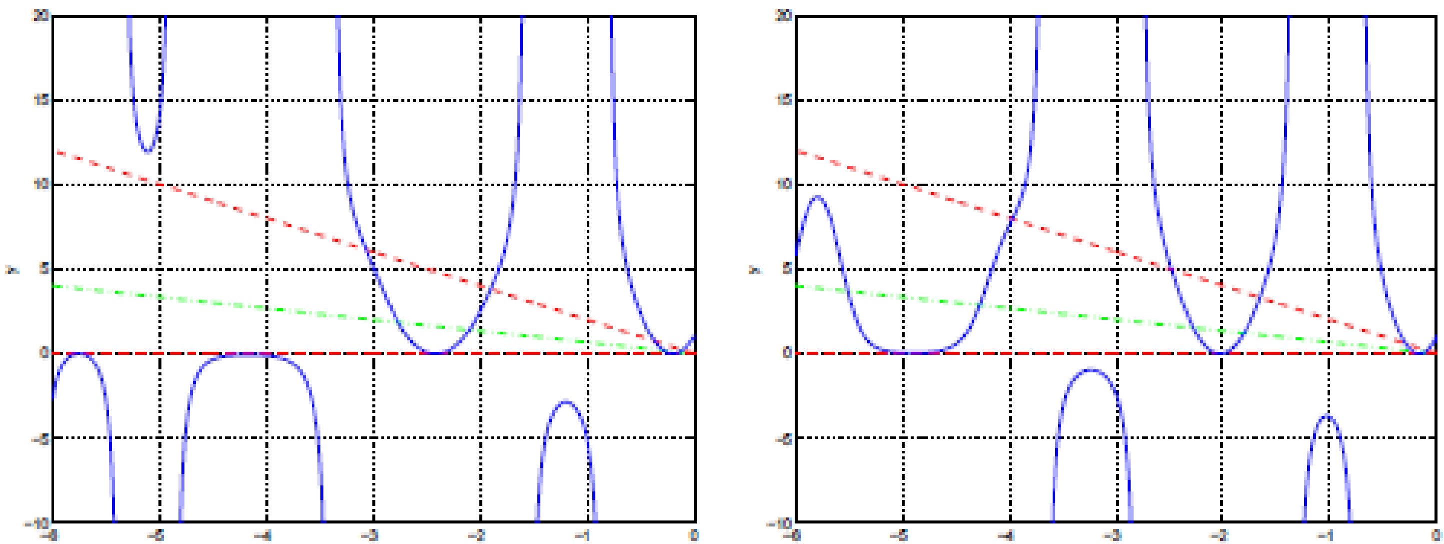
Unstable branch

Stable branch



Unstable branch

Stable branch



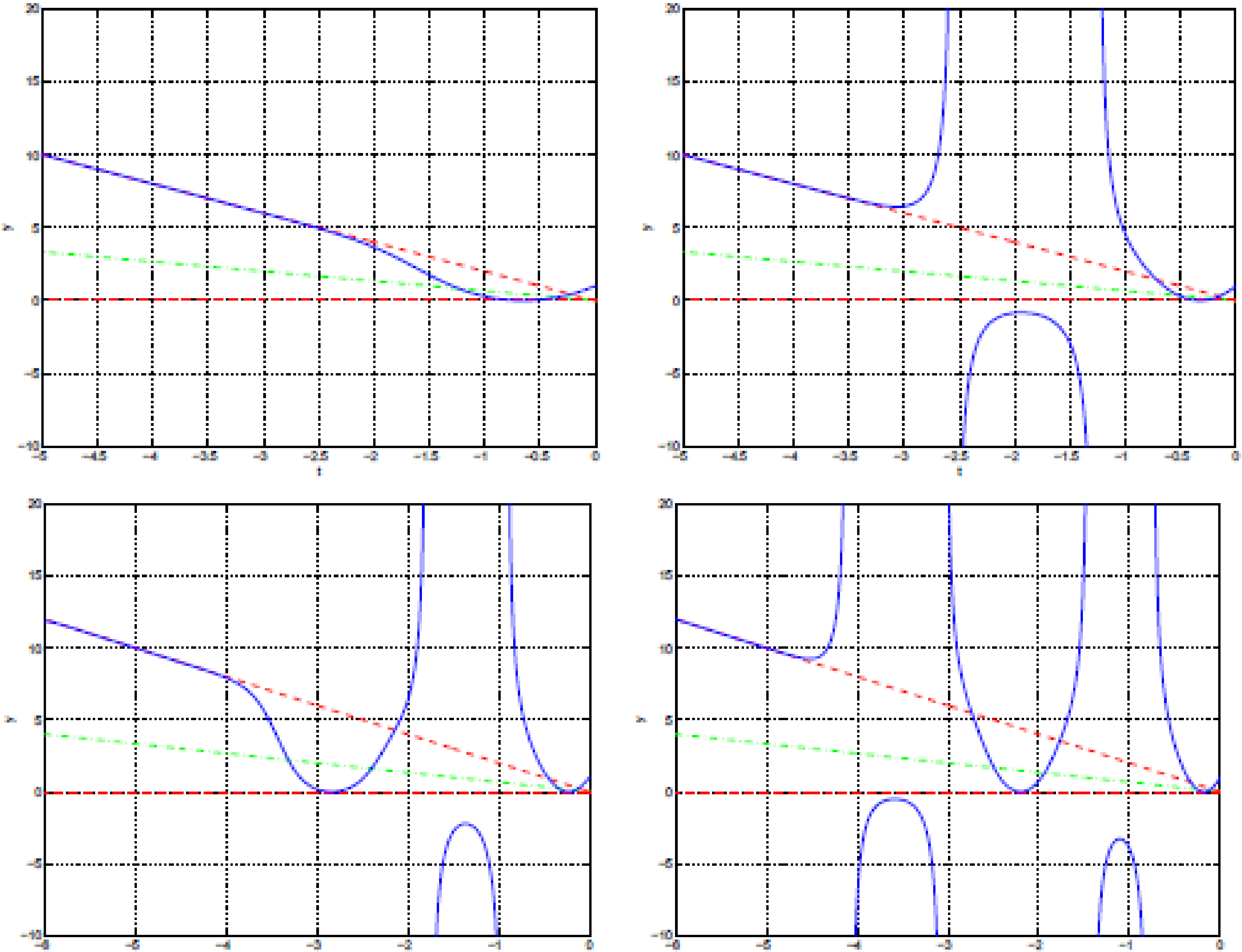
First four separatrix (eigenfunction) solutions [$y(0)=1$]:

Unstable

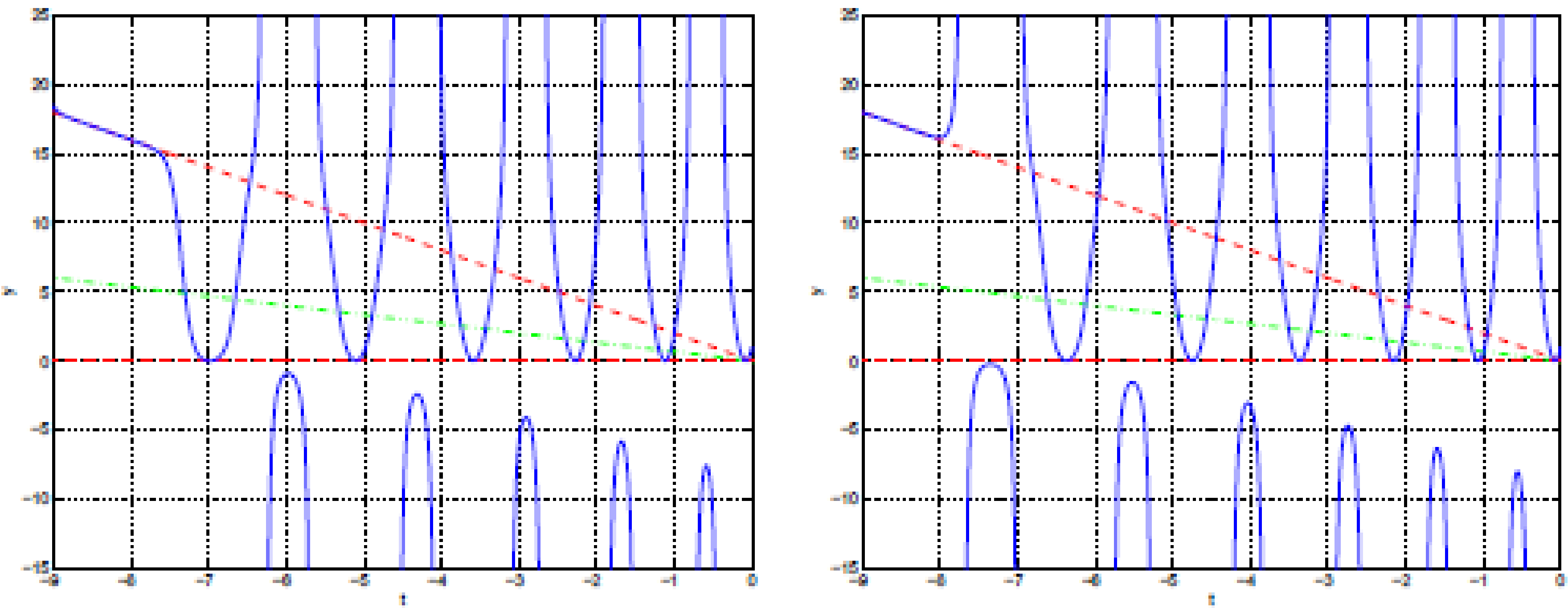
Stable

Unstable

Stable



Tenth and eleventh separatrix (eigenfunction) solutions:



***Slope* is the eigenvalue, initial value *y*(0) = 1**

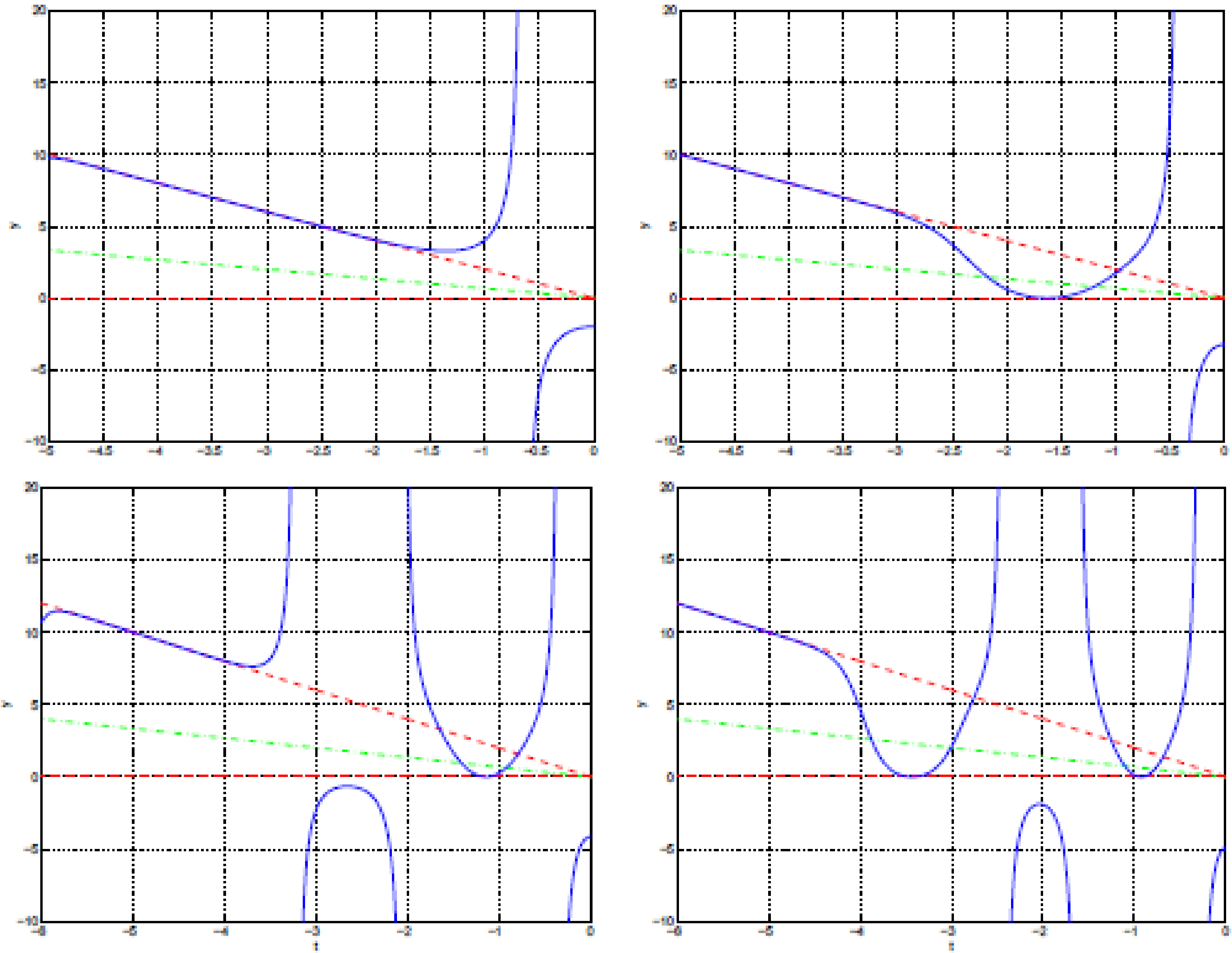
First four separatrix (eigenfunction) solutions [$y'(0)=0$]:

Unstable

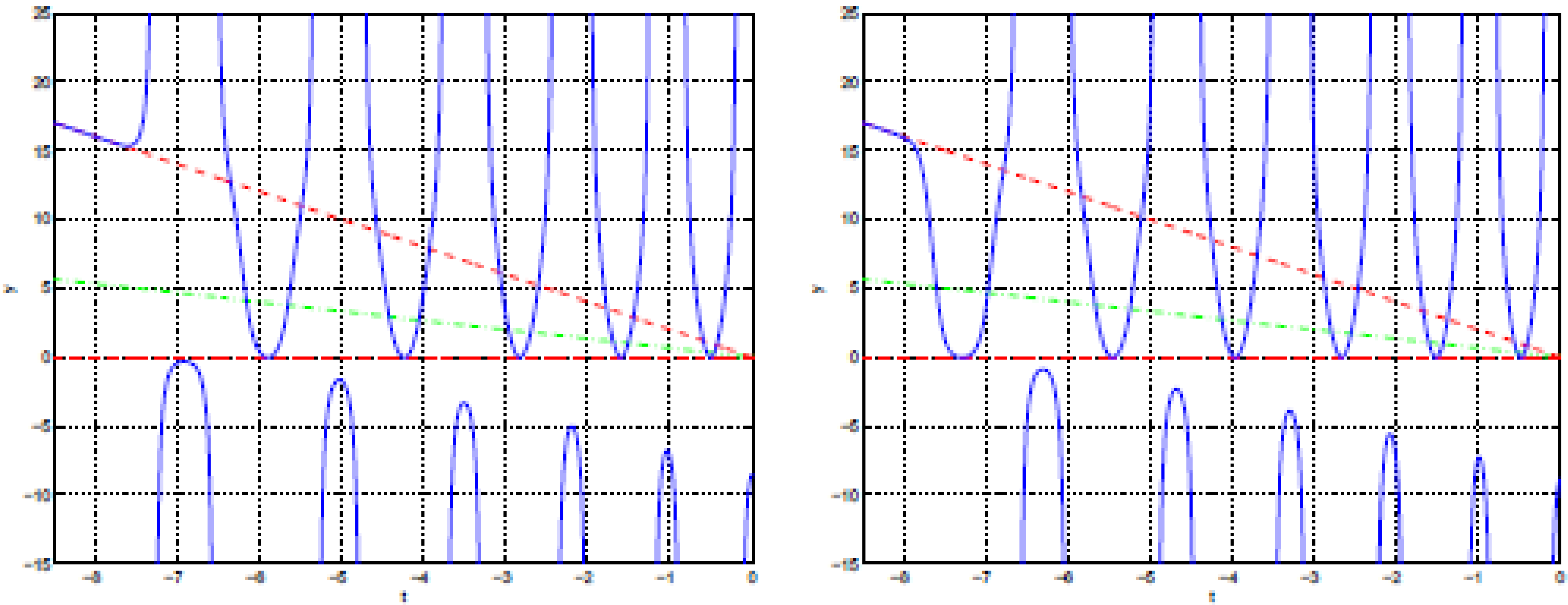
Stable

Unstable

Stable



Tenth and eleventh separatrix (eigenfunction) solutions:



$y(0)$ is the eigenvalue, initial slope is 0

Large n behaviour of eigenvalues: $b_n \sim B_{\mathrm{IV}} n^{3/4}$ and $c_n \sim C_{\mathrm{IV}} n^{1/2}$.

Numerical results using Richardson extrapolation:

$$B_{\mathrm{IV}} = 4.25684\underline{3}.$$

$$C_{\mathrm{IV}} = -2.62658\underline{7}$$

Analytic results using $\hat{H} = \frac{1}{2}\hat{p}^2 + \frac{1}{8}\hat{x}^6.$

$$B_{\mathrm{IV}} = 2^{3/2}\left[\sqrt{\pi}\,\Gamma\left(\tfrac{5}{3}\right)/\Gamma\left(\tfrac{7}{6}\right)\right]^{3/4},$$

$$C_{\mathrm{IV}} = -2\left[\sqrt{\pi}\,\Gamma\left(\tfrac{5}{3}\right)/\Gamma\left(\tfrac{7}{6}\right)\right]^{1/2}$$

Obtained by using WKB to calculate the large eigenvalues of the

sextic PT-symmetric Hamiltonian

Note:

Painlevé I, II, and IV correspond to $\varepsilon = 1, 2,$ and 4

This analysis extends to huge classes of equations beyond Painlevé. For example:

Super Painlevé:

$$y''(x) = \frac{2M+2}{(M-1)^2}[y(x)]^M + x[y(x)]^N$$

$$y'' = \frac{10}{9}y^4 + x$$

$$y(0) = 0$$

$$y_n'(0) \sim 1.1102\, n^{\frac{5}{11}} \left(1 + \frac{26.235}{11n}\right)$$

$$y_n'(0) \sim -1.109\, n^{\frac{5}{11}} \left(1 - \frac{11}{11n}\right)$$

$$y'(0) = 0$$

$$y_n(0) \sim 1.80547\, n^{\frac{2}{11}} \left(1 - \frac{1.998}{11n}\right)$$

$$y_n(0) \sim -1.226\, n^{\frac{2}{11}} \left(1 + \frac{3.03}{11n}\right)$$

$$y'' = \frac{10}{9}y^4 + xy$$

$$y(0) = 0$$

$$y_n'(0) \sim 2.1336 n^{\frac{5}{9}} \begin{cases} \left(1 - \frac{3.55}{9n}\right), & \text{odd } n \\ \left(1 - \frac{2.13}{9n}\right), & \text{even } n \end{cases}$$

$$y_n'(0) \sim -2.1336 n^{\frac{5}{9}} \begin{cases} \left(1 - \frac{3.24}{9n}\right), & \text{odd } n \\ \left(1 - \frac{2.08}{9n}\right), & \text{even } n \end{cases}$$

$$y'(0) = 0$$

$$y_n(0) \sim -1.59255 n^{\frac{2}{9}} \begin{cases} \left(1 - \frac{1.24}{9n}\right), & \text{odd } n \\ \left(1 - \frac{0.77}{9n}\right), & \text{even } n \end{cases}$$

$$y'' = \frac{10}{9}y^4 + xy^2$$

$$y(0) = 0$$

$$y_n'(0) \sim 2.9996 n^{\frac{5}{7}} \left(1 - \frac{0.959}{7n}\right)$$

$$y_n'(0) \sim -2.9996 n^{\frac{5}{7}} \left(1 - \frac{3.19}{7n}\right)$$

$$y'(0) = 0$$

$$y_n(0) \sim 1.098102 n^{\frac{2}{7}} \left(1 - \frac{2.00207}{7n}\right)$$

$$y_n(0) \sim -1.82502 n^{\frac{2}{7}} \left(1 - \frac{0.84}{7n}\right)$$

$$y'' = \frac{14}{25}y^6 + x$$

$$y(0) = 0$$

$$y_n'(0) \sim 2.3219 n^{\frac{7}{17}} \left(1 + \frac{2.52}{17n}\right)$$
$$y_n'(0) \sim -2.322 n^{\frac{7}{17}} \left(1 - \frac{7.3}{17n}\right)$$

$$y'(0) = 0$$

$$y_n(0) \sim 1.500998 n^{\frac{2}{17}} \left(1 - \frac{1.9992}{17n}\right)$$
$$y_n(0) \sim -1.652812 n^{\frac{2}{17}} \left(1 - \frac{0.6732}{17n}\right)$$

$$y'' = \frac{14}{25}y^6 + xy$$

$$y(0) = 0$$

$$y_n'(0) \sim 2.3569 n^{\frac{7}{15}} \begin{cases} \left(1 + \frac{4.137}{15n}\right), & \text{odd } n \\ \left(1 + \frac{0.4425}{15n}\right), & \text{even } n \end{cases}$$

$$y_n'(0) \sim -2.357 n^{\frac{7}{15}} \begin{cases} \left(1 - \frac{4.125}{15n}\right), & \text{odd } n \\ \left(1 - \frac{0.4275}{15n}\right), & \text{even } n \end{cases}$$

$$y'' = \frac{14}{25}y^6 + xy^2$$

$$y(0) = 0$$

$$y_n'(0) \sim 2.5979 n^{\frac{7}{13}} \begin{cases} \left(1 - \frac{1.2675}{13n}\right), & \text{odd } n \\ \left(1 - \frac{1.8473}{13n}\right), & \text{even } n \end{cases}$$

$$y_n'(0) \sim -2.598 n^{\frac{7}{13}} \begin{cases} \left(1 - \frac{5.148}{13n}\right), & \text{odd } n \\ \left(1 - \frac{5.727}{13n}\right), & \text{even } n \end{cases}$$

$$y'(0) = 0$$

$$y_n(0) \sim 1.73085 n^{\frac{2}{13}} \left(1 + \frac{0.195}{13n}\right)$$

$$y_n(0) \sim -1.7065 n^{\frac{2}{13}} \begin{cases} \left(1 - \frac{0.913}{13n}\right), & \text{odd } n \\ \left(1 - \frac{1.08}{13n}\right), & \text{even } n \end{cases}$$

$$y'' = \frac{14}{25}y^6 + xy^3$$

$$y(0) = 0$$

$$y_n'(0) \sim 1.7408 n^{\frac{7}{11}} \begin{cases} \left(1 + \frac{0.33}{11n}\right), & \text{odd } n \\ \left(1 - \frac{4.58}{11n}\right), & \text{even } n \end{cases}$$

$$y_n'(0) \sim -1.7408 n^{\frac{7}{11}} \begin{cases} \left(1 - \frac{2.44}{11n}\right), & \text{odd } n \\ \left(1 - \frac{7.34}{11n}\right), & \text{even } n \end{cases}$$

$$y'(0) = 0$$

$$y_n(0) \sim -1.52224 n^{\frac{2}{11}} \begin{cases} \left(1 + \frac{0.305}{11n}\right), & \text{odd } n \\ \left(1 + \frac{1.705}{11n}\right), & \text{even } n \end{cases}$$

$$y'' = \frac{14}{25}y^6 + xy^4$$

$$y(0) = 0, \ x < 0$$

$$y_n'(0) \sim 3.06787 n^{\frac{7}{9}} \begin{cases} \left(1 + \frac{1.5}{9n}\right), & \text{odd } n \\ \left(1 - \frac{1.87}{9n}\right), & \text{even } n \end{cases}$$

$$y_n'(0) \sim -3.06786 n^{\frac{7}{9}} \begin{cases} \left(1 - \frac{5.13}{9n}\right), & \text{odd } n \\ \left(1 - \frac{8.51}{9n}\right), & \text{even } n \end{cases}$$

$$y(0) = 0, \ x > 0$$

$$y_n'(0) \sim 2.9010 n^{\frac{7}{9}} \left(1 - \frac{1.67}{9n}\right)$$

$$y_n'(0) \sim -2.9010 n^{\frac{7}{9}} \left(1 + \frac{1.65}{9n}\right)$$

$$y'' = \frac{4}{9}y^7 + xy^4$$

$$y(0) = 0,\ x < 0$$

$$y_n'(0) \sim -1.38115 n^{\frac{2}{3}} \left(1 - \frac{0.4635}{3n}\right)$$

$$y(0) = 0,\ x > 0$$

$$y_n'(0) \sim -1.38114 n^{\frac{2}{3}} \left(1 - \frac{0.462}{3n}\right)$$

$$y'' = \frac{4}{9}y^7 + xy^5$$

$$y(0) = 0,\ x < 0$$

$$y_n'(0) \sim -1.86695 n^{\frac{4}{5}} \begin{cases} \left(1 - \frac{0.7375}{5n}\right), & \text{odd } n \\ \left(1 - \frac{3.2575}{5n}\right), & \text{even } n \end{cases}$$

$$y(0) = 0,\ x > 0$$

$$y_n'(0) \sim -2.29535 n^{\frac{4}{5}} \left(1 - \frac{0.225}{5n}\right)$$

Hyperfine splitting

$$y'' = \frac{1}{a^2}y^4 + xy^2, \quad \text{with} \quad a \equiv \frac{3}{\sqrt{10}}.$$

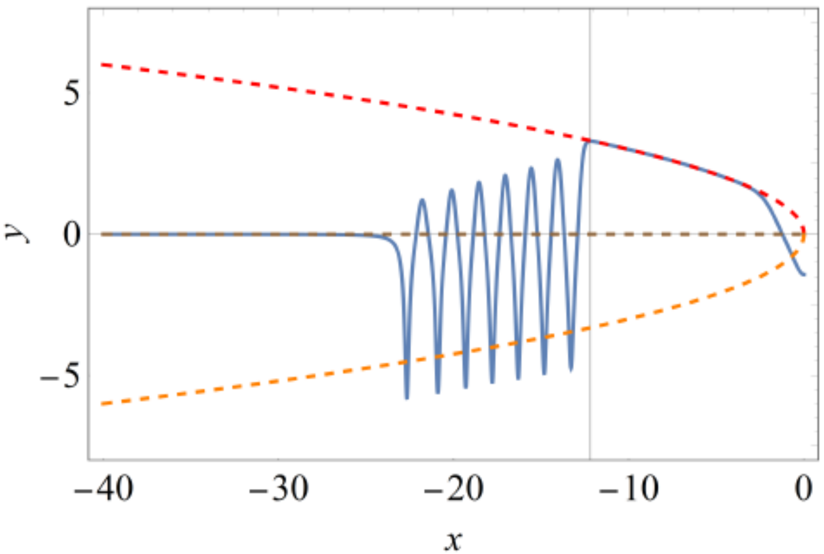
Let $y_{nm}(x) = Y_n(x) + \phi_m(x)$, where $Y_n(x)$ is a separatrix solution with

$$Y_n(x) \sim a\sqrt{-x}, \qquad x \to -\infty.$$

The new hyperfine solutions initially follow $Y_n(x)$.

Then they deviate from $Y_n(x)$ and oscillate *m* times about the curve $-a\sqrt{-x}$.

Finally, they level off for large x as $y_{nm}(x) \sim \frac{12}{x^3}, \qquad x \to -\infty.$



The initial values of ϕ are the hyperfine eigenvalues.

For example, for the lowest eigenfunction Y_0

$$\phi_m(0) \sim 4.1789\,\mathrm{e}^{-9.26201m}, \qquad m \to \infty.$$

The hyperfine oscillation separates at the negative values

$$T_m \sim \left(\frac{7}{4\sqrt{2a}} 9.26201m\right)^{\frac{4}{7}}, \qquad m \to \infty.$$

We hope we have opened a window to a new area of *nonlinear semiclassical* asymptotic analysis

Thanks for listening!