

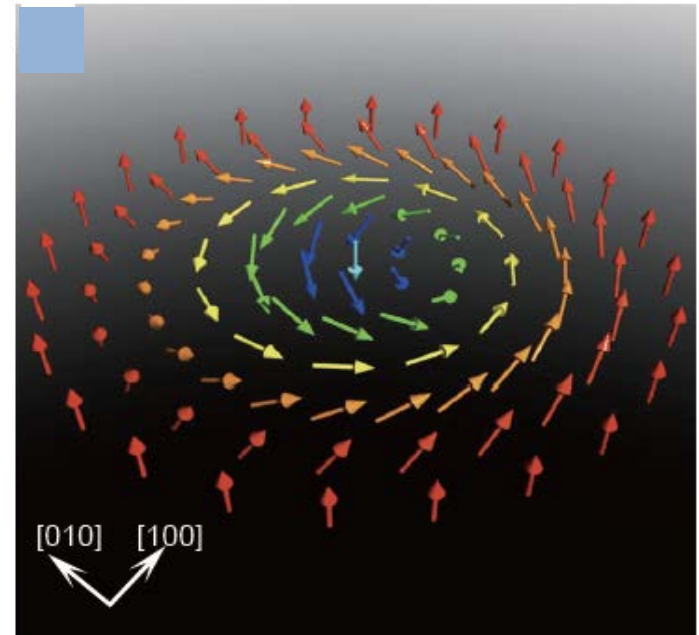
Skyrmion Dynamics and Topological Transport Phenomena

Yoshi Tokura

*RIKEN Center for Emergent Matter Science (CEMS)
Department of Applied Physics, University of Tokyo*



“**skyrmion**”, the concept originally introduced by Tony Skyrme (1922-87) to describe the state of nucleon: to model a **particle** as a **topological soliton**



Skyrmions and topological transport phenomena

Skyrmions in multiferroics toward *E*-control and light control

Collaborators

- **RIKEN (Japan)**

X. Z. Yu, S. Seki N. Nagaosa, Y. Tokunaga,
D. Okuyama, Y. Taguchi, M. Kawasaki,
T. Arima, N. Ogawa, M. Kubota

- **Univ. of Tokyo (Japan)**

Y. Kanazawa, Y. Onose,
S. Ishiwata, A. Tsukazaki,
M. Ichikawa, Y. Shiomi, K. Shibata,
F. Kagawa, Y. Okamura, M. Mochizuki

- **NIMS (Japan)**

Y. Matsui, K. Kimoto, W. Z. Zhang

- **Sung Kyun Kwan Univ. (Korea)**

J. H. Han, J. H. Park

- **MPI (Germany)**

D. S. Inosov, J. H. Kim, B. Keimer

- **PSI (Switzerland)**

J. White, N. Egetenmeyer, J. Gavilano

- **Groningen Univ. (Holland)**

M. Mostovoy



Xiuazhen Yu



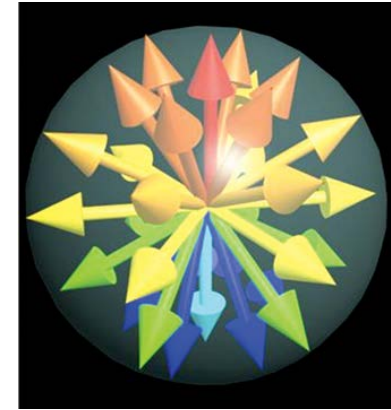
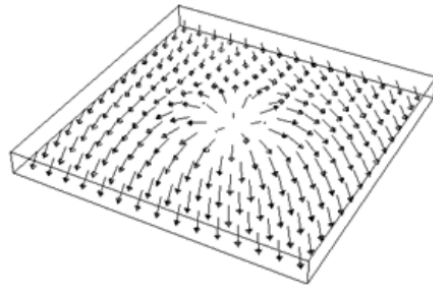
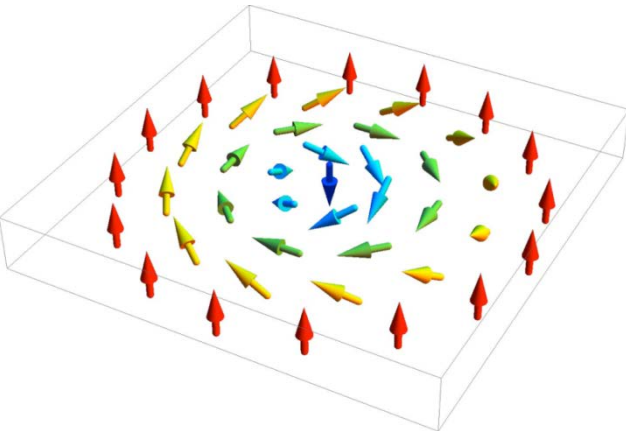
Naoya Kanzawa



Shinichiro Seki

What is magnetic skyrmion?

5 ~ 100 nm

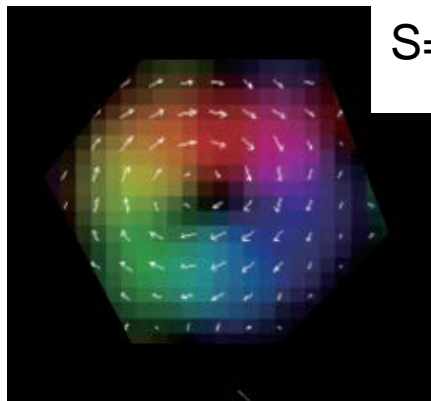


Topologically-stable spin vortex
with particle-like nature

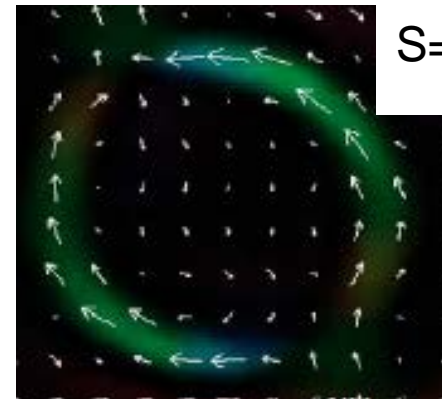
“skyrmion number”

Lateral component of
M of some bubbles

$$S = \frac{1}{4\pi} \int \vec{n} \cdot \frac{\partial \vec{n}}{\partial x} \times \frac{\partial \vec{n}}{\partial y} d\vec{r} = -1$$



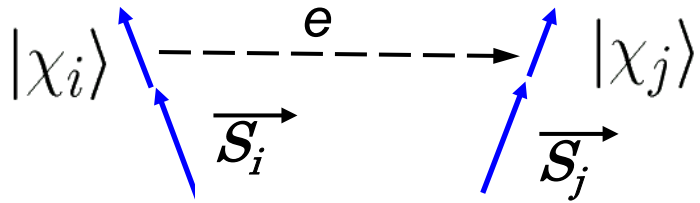
S=-1



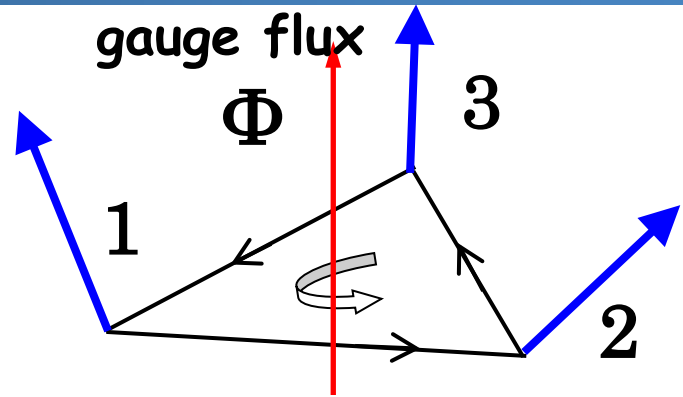
S=0

a pair of
Bloch lines

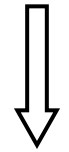
Spin Chirality



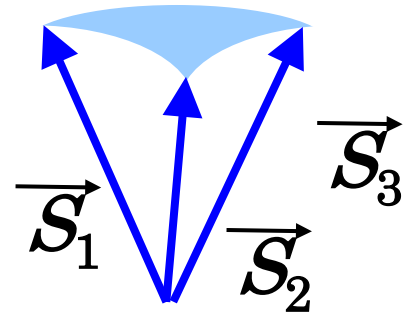
$$\begin{aligned}
 t_{ij} &= t \langle \chi_i | \chi_j \rangle \\
 &= t \left[\cos\left(\frac{\theta_i}{2}\right) \cos\left(\frac{\theta_j}{2}\right) + \sin\left(\frac{\theta_i}{2}\right) \sin\left(\frac{\theta_j}{2}\right) e^{i(\phi_i - \phi_j)} \right] \\
 &= t \cos\left(\frac{\theta_{ij}}{2}\right) \exp(i a_{ij})
 \end{aligned}$$



non-coplanar spin structure



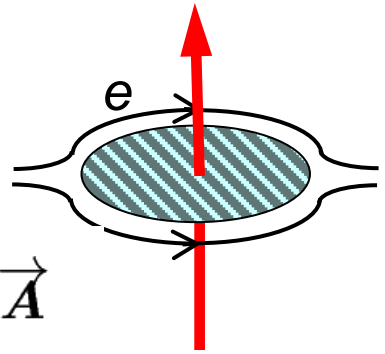
solid angle Ω



Fictitious "magnetic field" $\Phi = \frac{1}{2} \Omega$

$$= 2 \vec{S}_1 \cdot \vec{S}_2 \times \vec{S}_3$$

Aharonov - Bohm effect

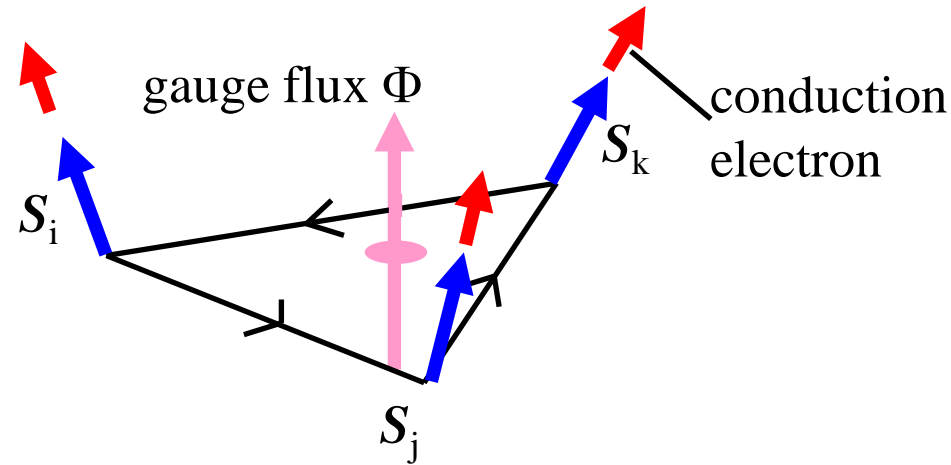


$$\vec{B} = \text{rot } \vec{A}$$

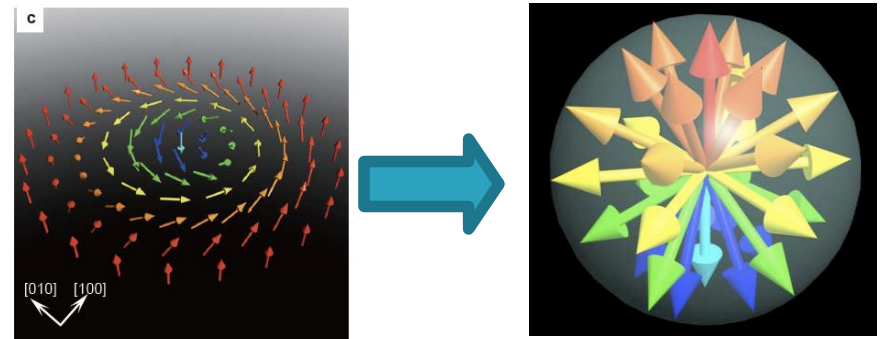
$$-\frac{ie}{\hbar} \oint \vec{A} d\vec{r} = -2\pi i \cdot \frac{\Phi}{\Phi_0}$$

$$\Phi_0 = \frac{h}{e} : \text{flux quantum}$$

Skyrmion



Mapping to a sphere

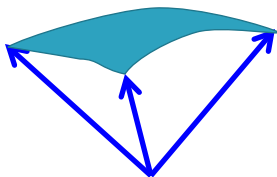


Solid angle $\Omega = 4\pi$

Cf. Spin chirality

$$\vec{S}_i \cdot (\vec{S}_j \times \vec{S}_k)$$

$= 1/2 \Omega$ Solid angle



Continuum approximation

Total spin Chirality

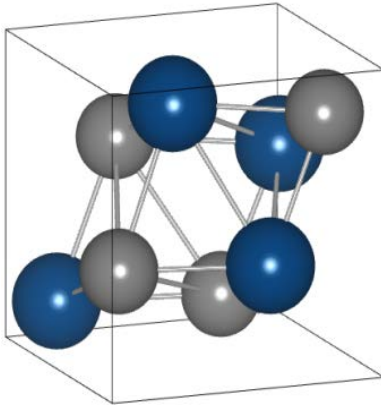
$$= \frac{1}{4\pi S^3} \int d^2\mathbf{r} \mathbf{S} \cdot (\nabla_x \mathbf{S} \times \nabla_y \mathbf{S})$$

$$= N_s \quad \text{Skyrmion number}$$

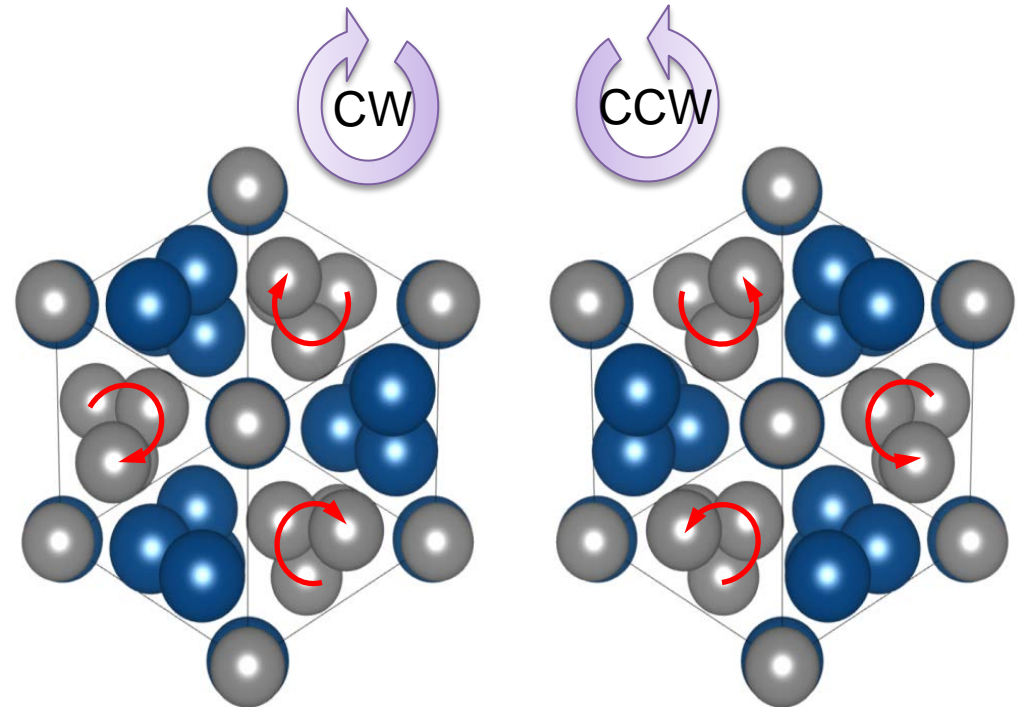
Skyrmion carries emergent magnetic field.

Helical spin order in B20-type crystals

Crystal structure



- : Transition-metal element
- : Group 14 element
- Cubic ($P2_13$)
- Noncentrosymmetric



Chiral lattice structure

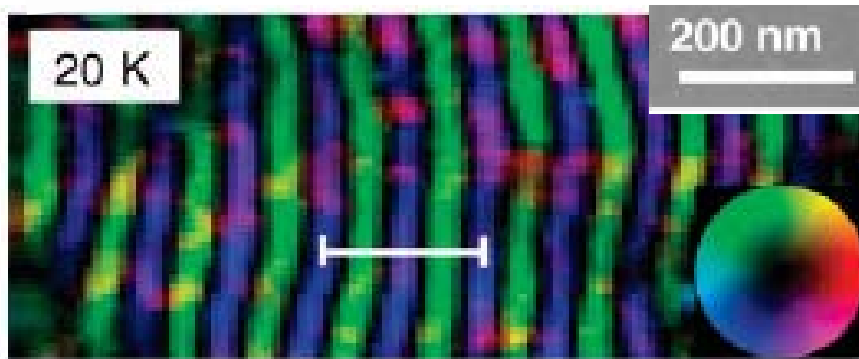
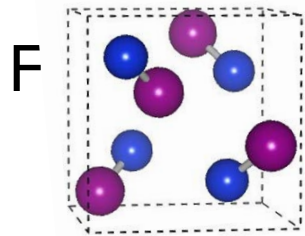
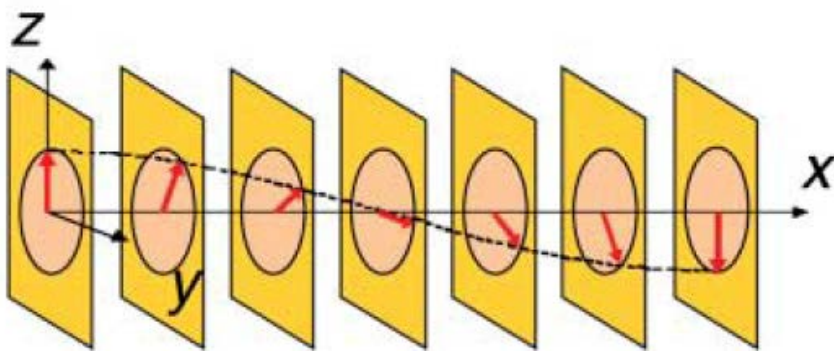
Magnetic structure

Three well-separated energy scales

ferromagnetic interaction ($\mathbf{S}_i \cdot \mathbf{S}_j$) > Dzyaloshinsky-Moriya interaction ($\mathbf{S}_i \times \mathbf{S}_j$) > magnetic anisotropy
→ one-handed helical spin structure
(a long wavelength 17.5 – 230 nm, weakly locked helix direction $\langle 111 \rangle$ or $\langle 100 \rangle$)

Toward real space observation of Skyrmion structure

structure



M. Uchida, Y. Onose, Y. Matsui, Y. Tokura,
Science (2006)

$$H = \sum \left(\underbrace{-J\vec{S}_i \cdot \vec{S}_j}_{\text{Ferro}} + \underbrace{\vec{D}_{ij} \cdot (\vec{S}_i \times \vec{S}_j)}_{\text{DM}} \right)$$

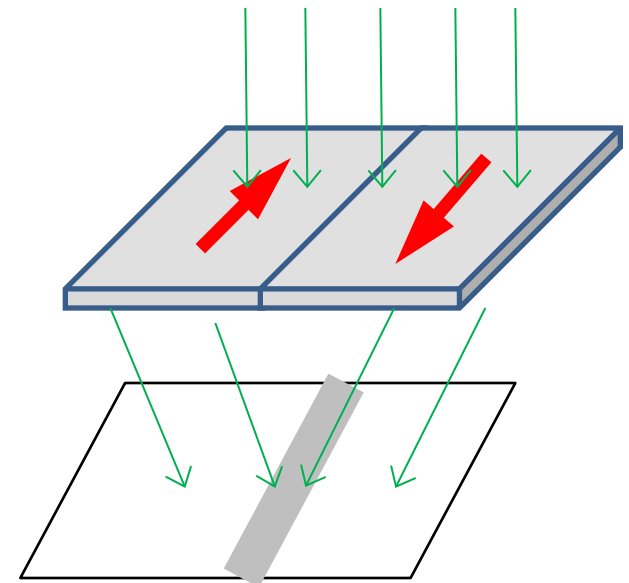
Ferro + DM



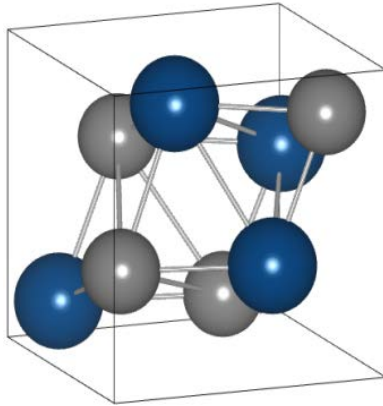
Helical spin structure

Long period $\sim aJ/D \sim 10\text{nm}-300\text{nm}$

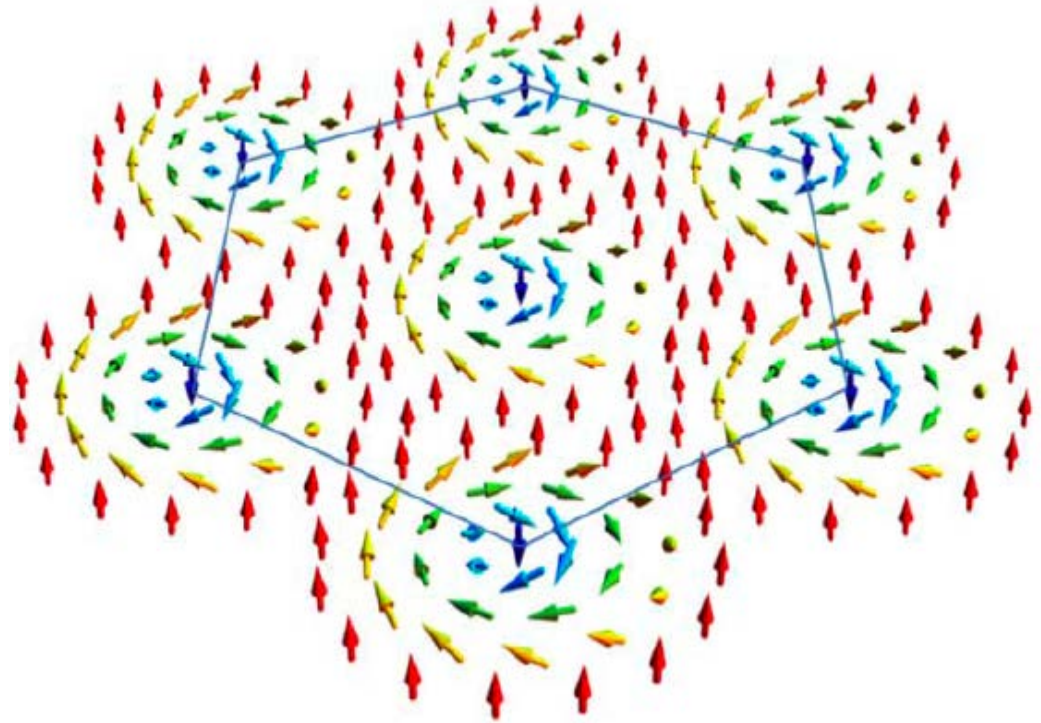
Lorentz microscope
electrons



Crystal structure



- : Transition-metal element
- : Group 14 element
- Cubic ($P2_13$)
- Noncentrosymmetric



Magnetic structure

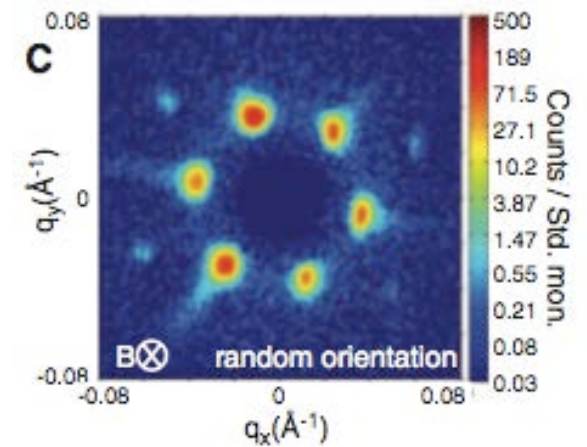
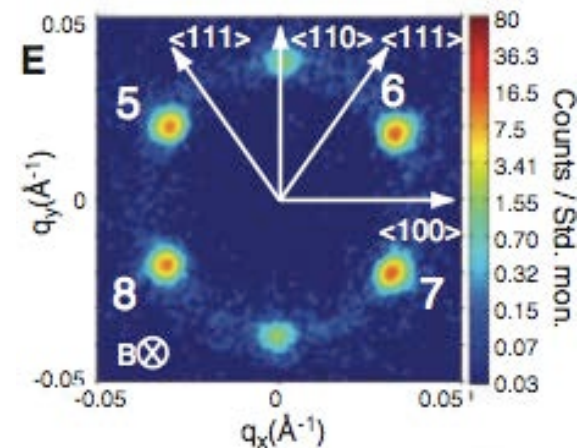
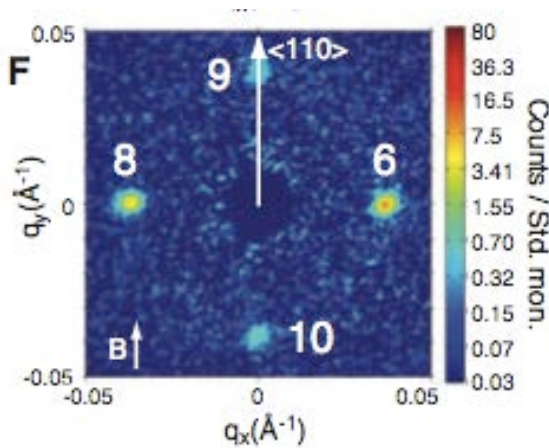
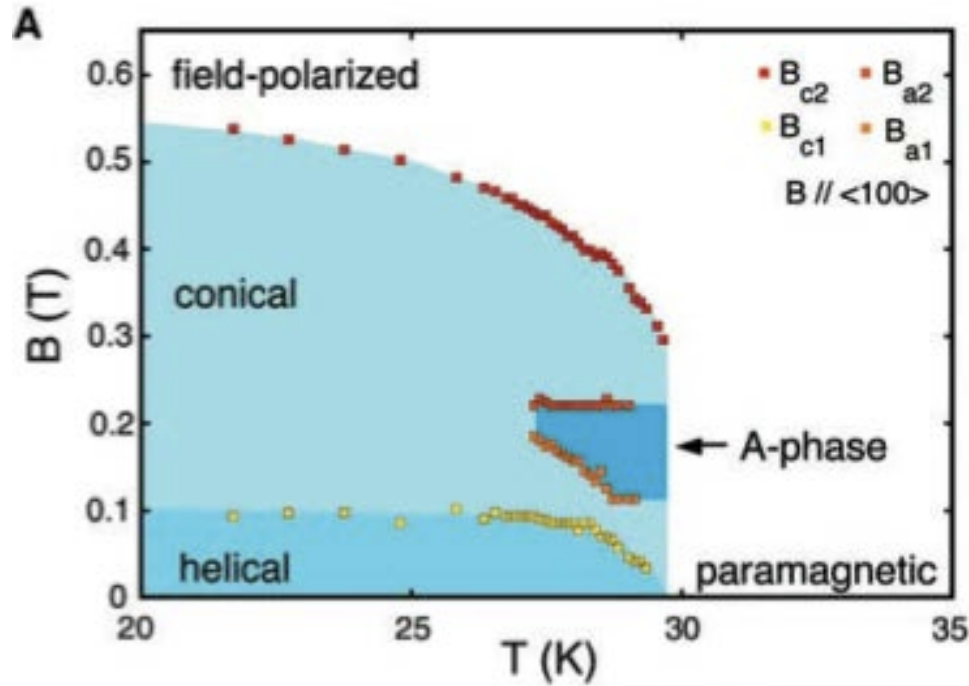
Three well-separated energy scales

ferromagnetic interaction ($\mathbf{S}_i \cdot \mathbf{S}_j$) > Dzyaloshinsky-Moriya interaction ($\mathbf{S}_i \times \mathbf{S}_j$) > magnetic anisotropy
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(a long wavelength 17.5 – 230 nm, weakly locked helix direction $\langle 111 \rangle$ or $\langle 100 \rangle$)

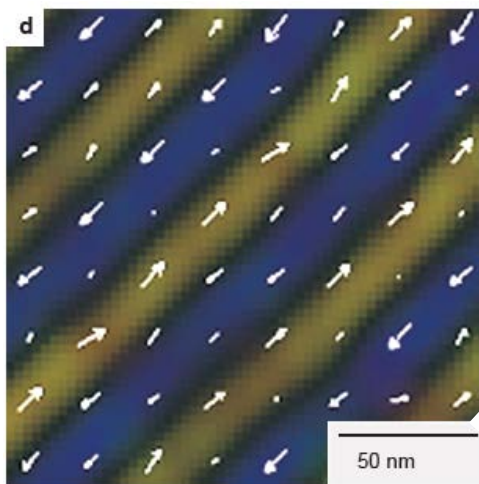
Small angle neutron scattering for Skyrmion Xtal

MnSi

S. Mühlbauer *et al.*,
Science 323 915
(2009)



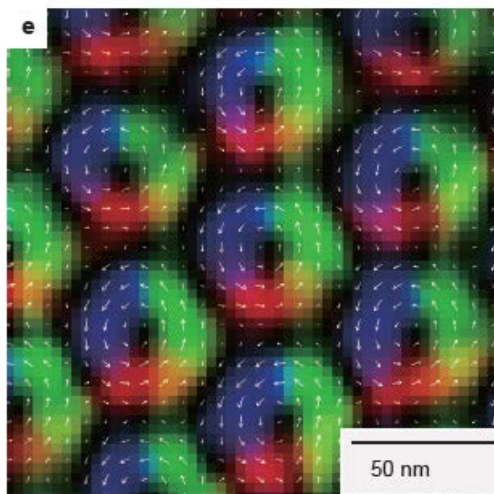
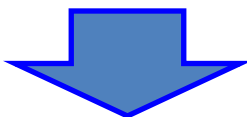
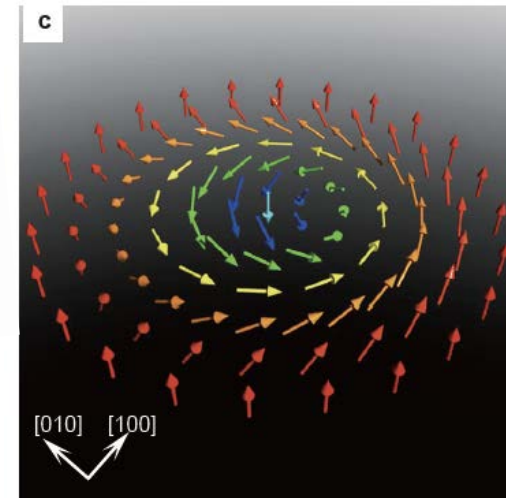
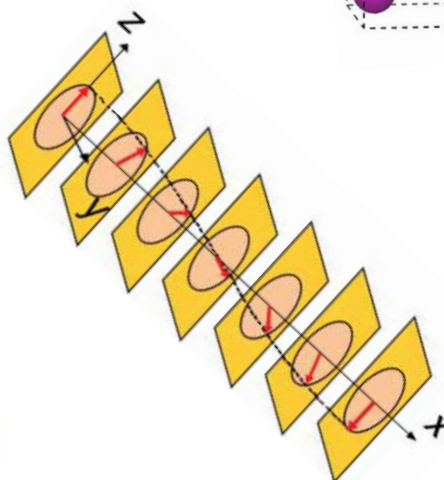
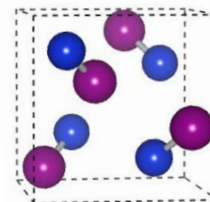
Real Space Observation of Skyrmion crystal



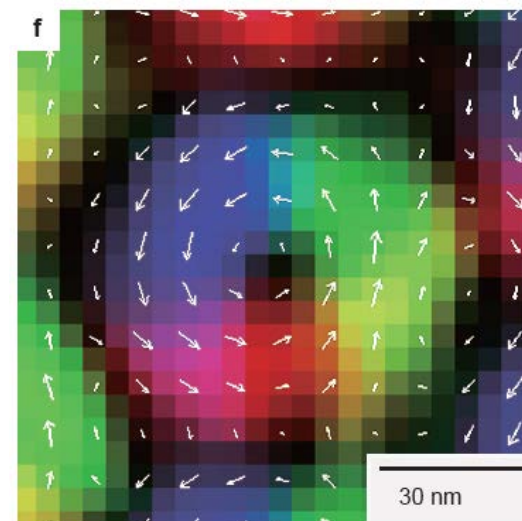
$\text{Fe}_{0.5}\text{Co}_{0.5}\text{Si}$

$T=25\text{K}$

$H=0$



$H=50\text{mT}$

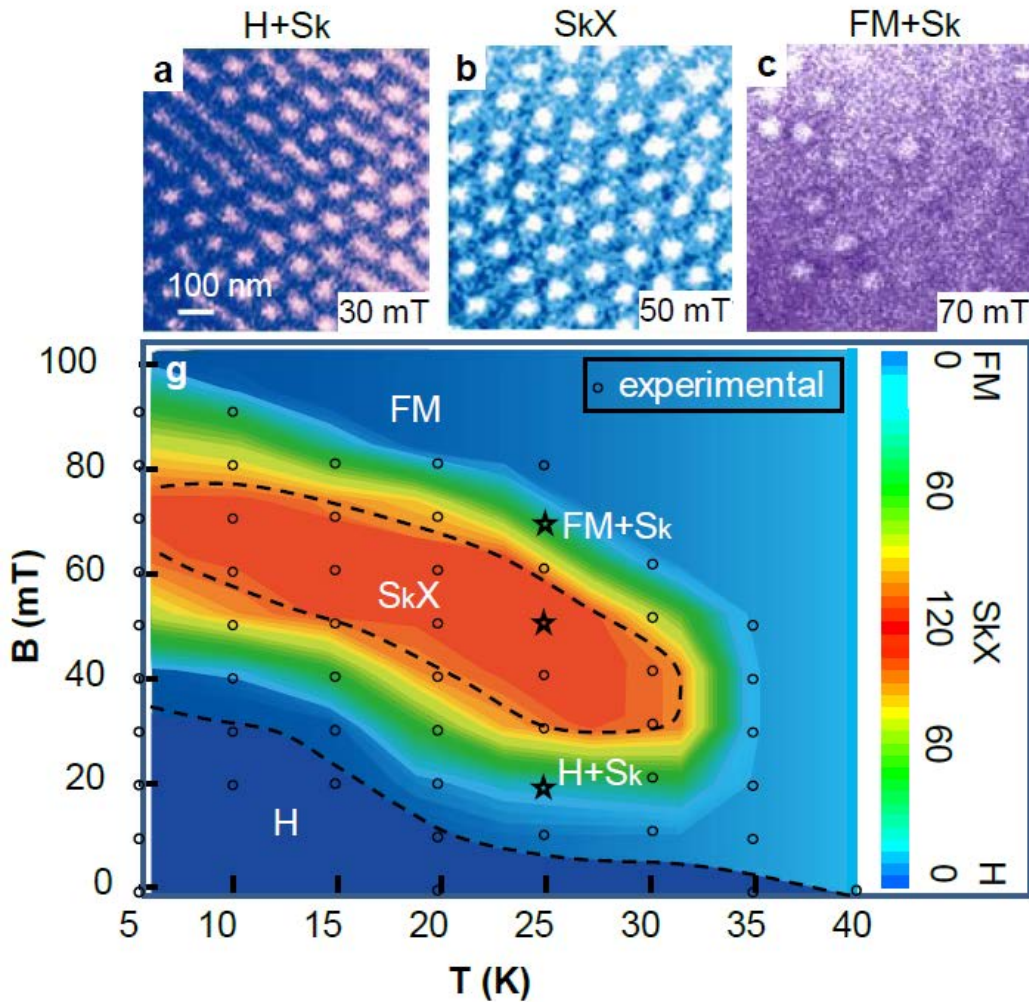


X.Z. Yu, Y.T. *et al.* Nature (2010).

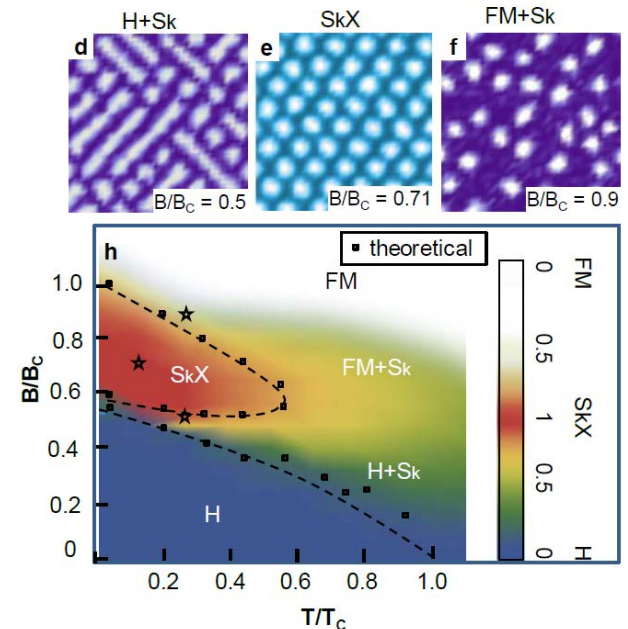
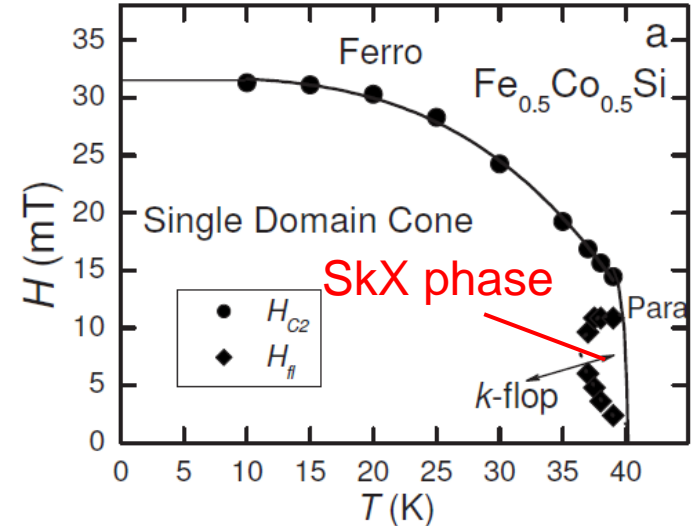
H-T Phase diagram

Bulk sample

20nm-thick film (Lorentz TEM)



Skx: Skyrmion Crystal



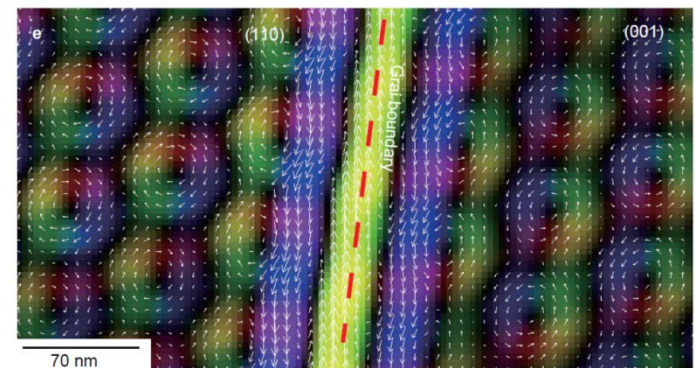
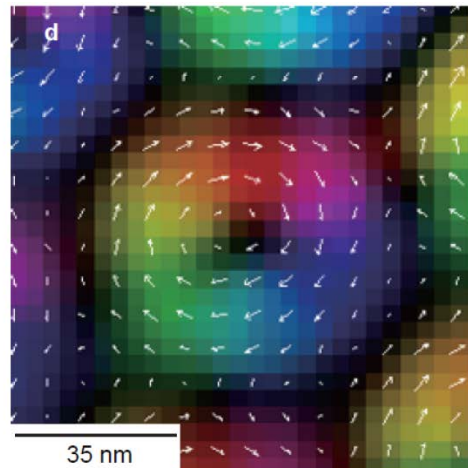
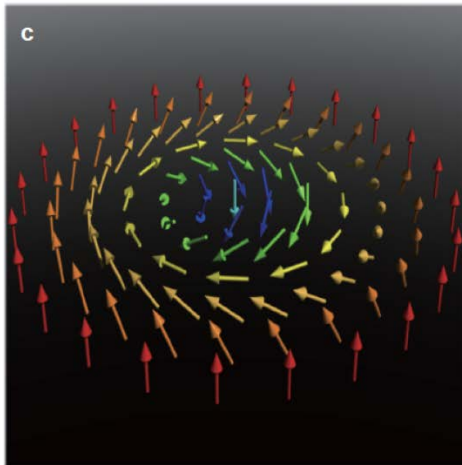
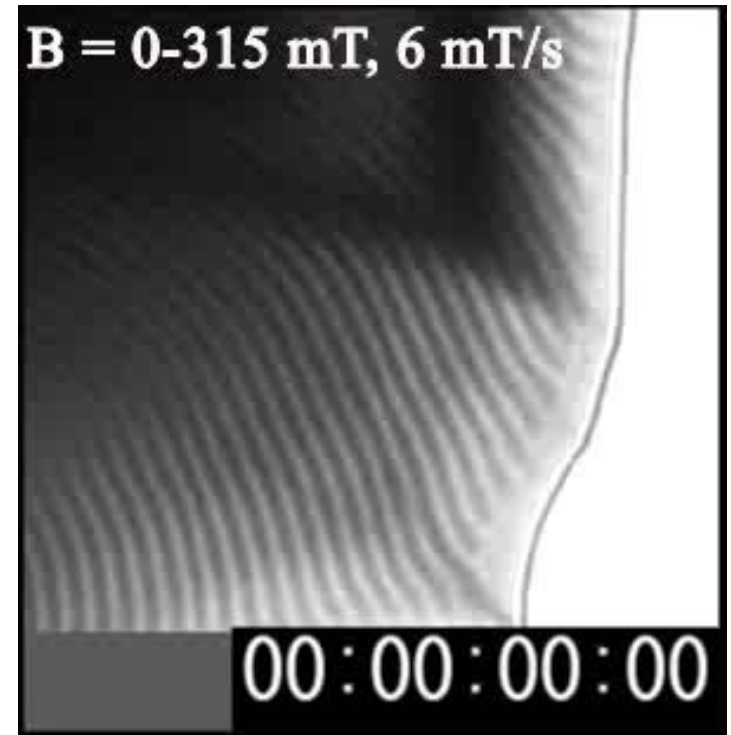
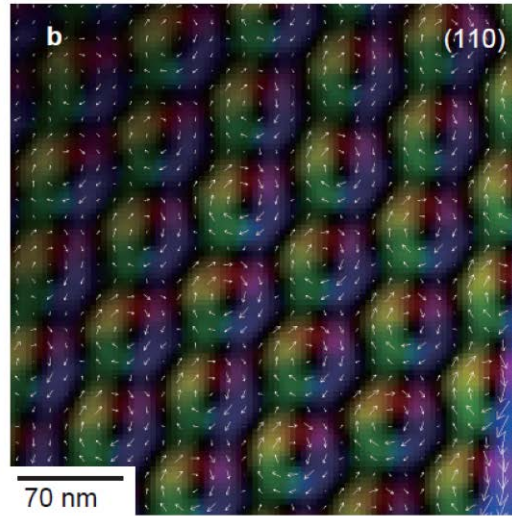
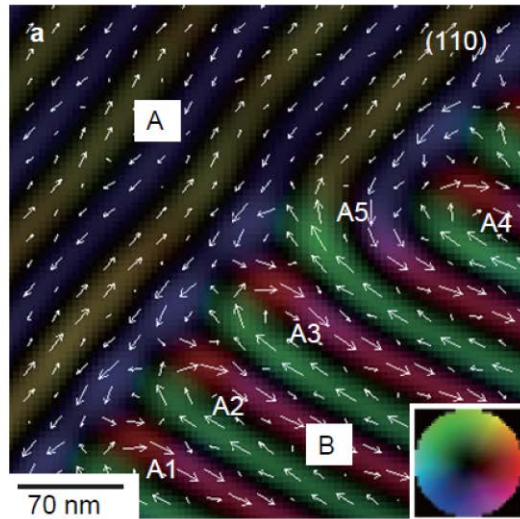
2D simulation

FeGe: from helical to skyrmion crystal at 260K

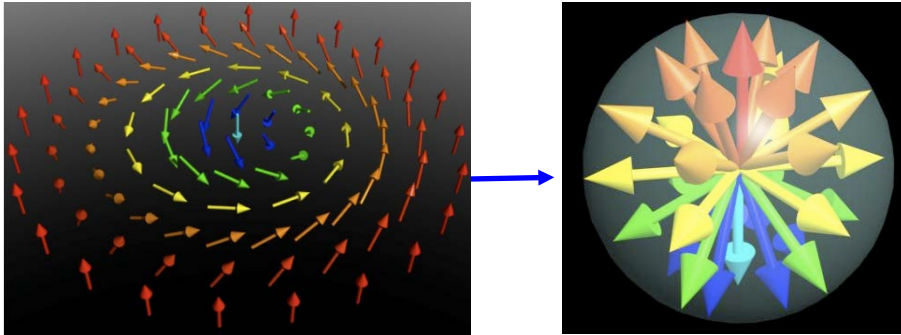
X.Z. Yu et al. Nat. Mater.(2010)

H=0

\odot H=0.1T



Real-space fictitious magnetic field in a skyrmion spin texture



Solid angle $\Omega = 4\pi$

In strong coupling case

One skyrmion

One magnetic flux ϕ_0

$$\phi_0 = h/e$$

Emergent magnetic field

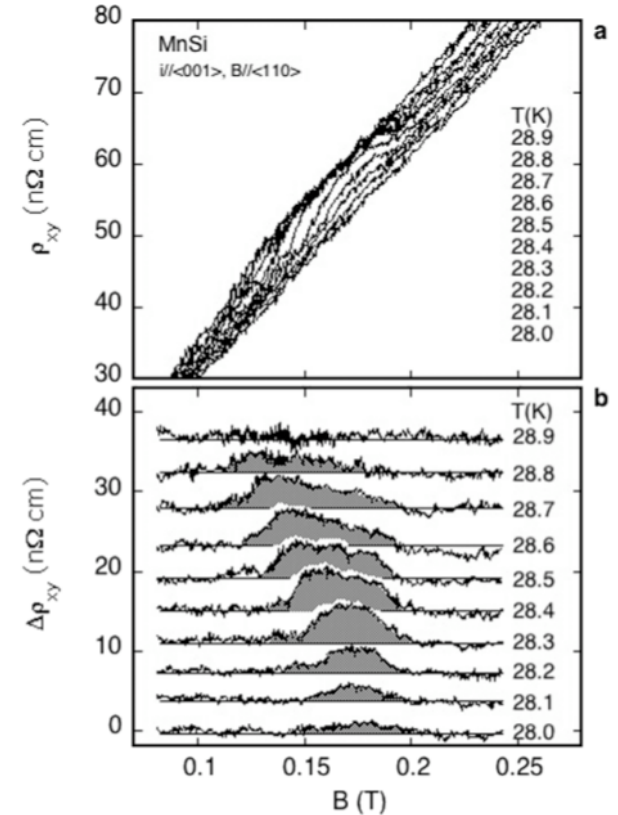
one skyrmion/nm²

$$\mathbf{B}_{\text{eff}}^z = -\phi_0/A$$

A: skyrmion size

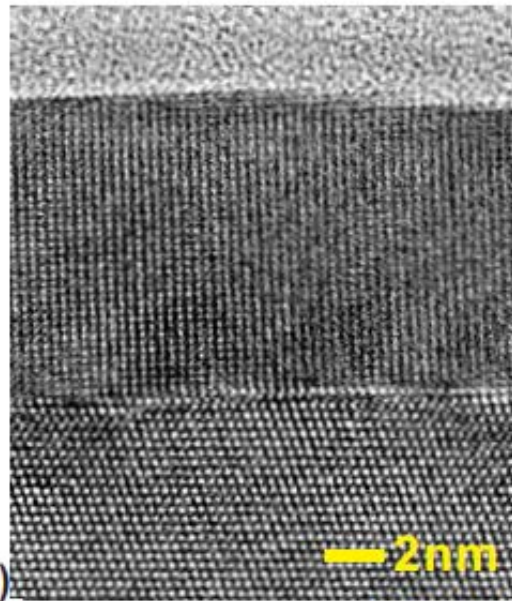
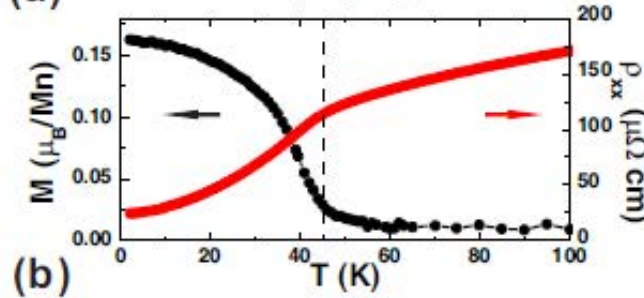
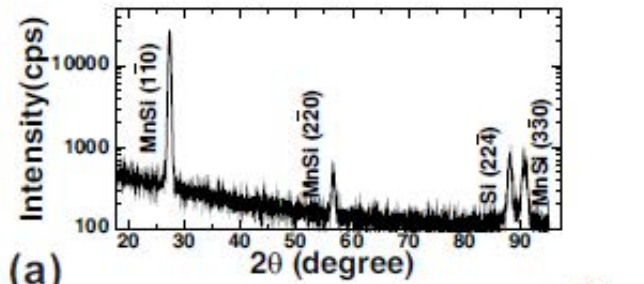
$B_{\text{eff}} \sim 4000\text{T}$

A. Neubauer *et al*, PRL 102 186602 (2009)



High skyrmion density \Leftrightarrow Large topological Hall Effect

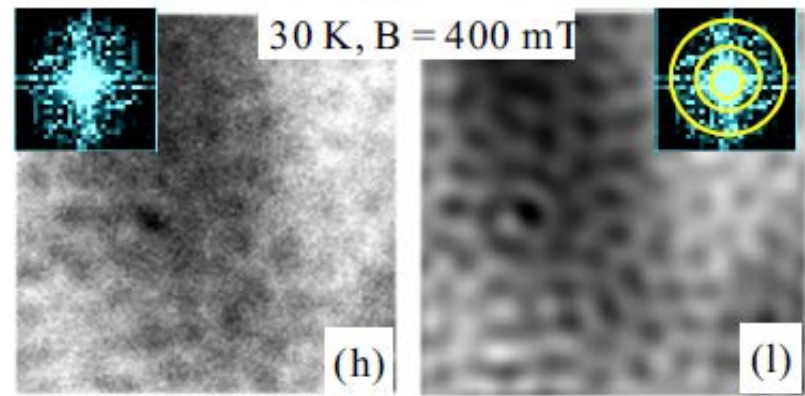
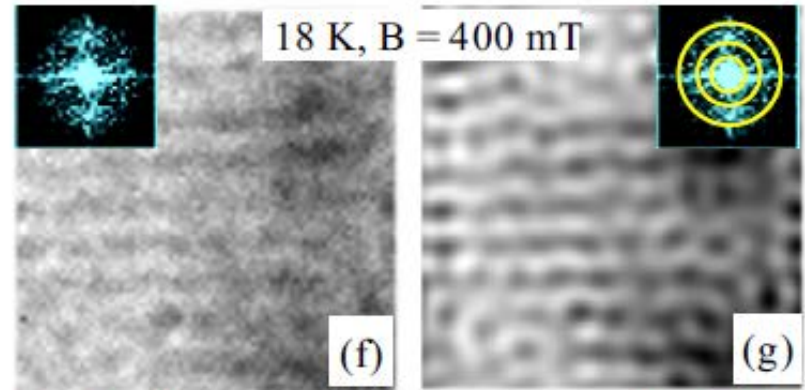
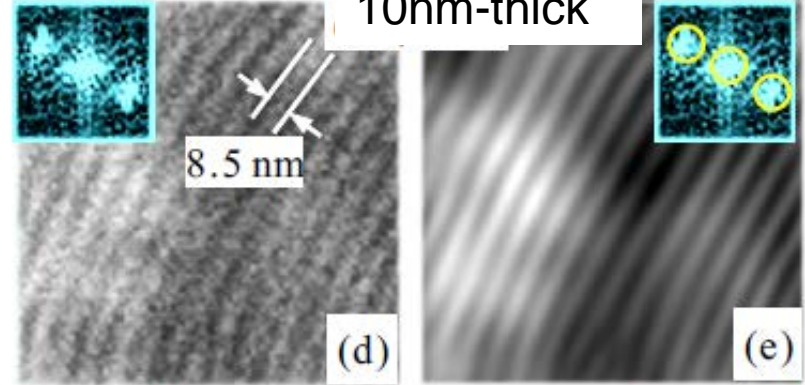
Ultrathin epitaxial thin films of MnSi



MnSi

Si substrate

10nm-thick



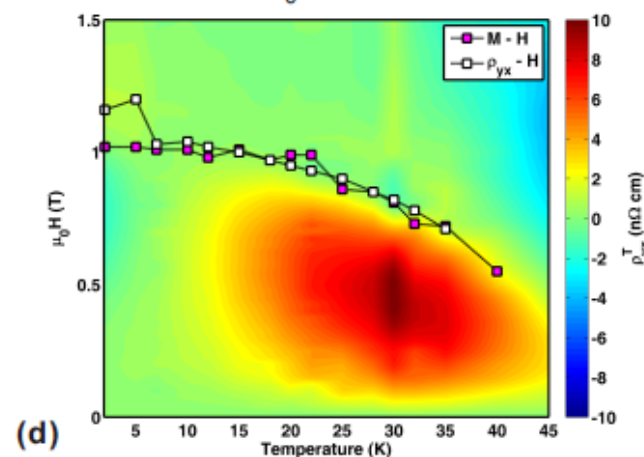
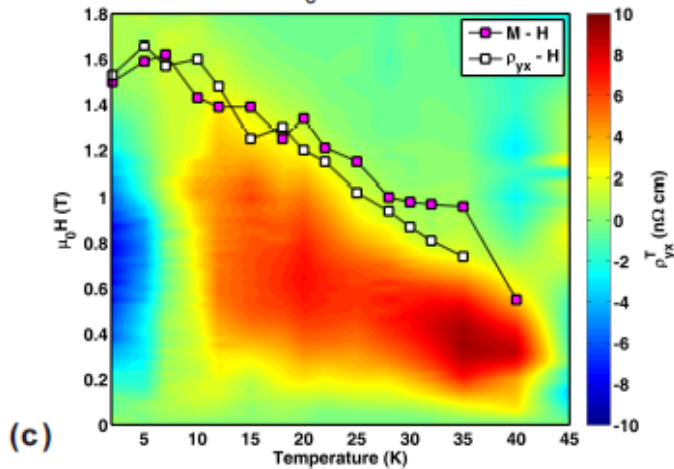
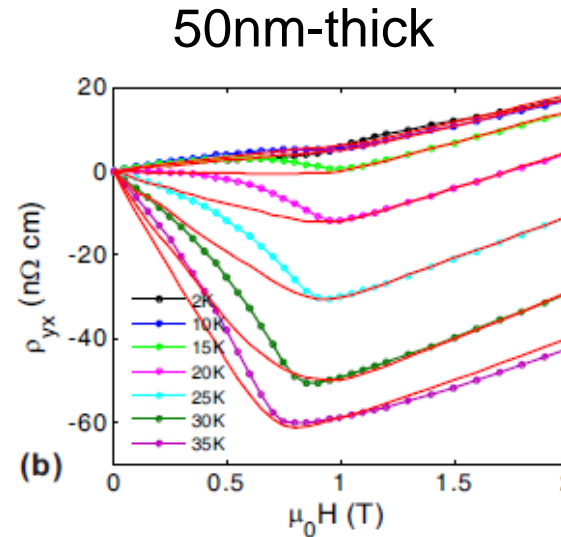
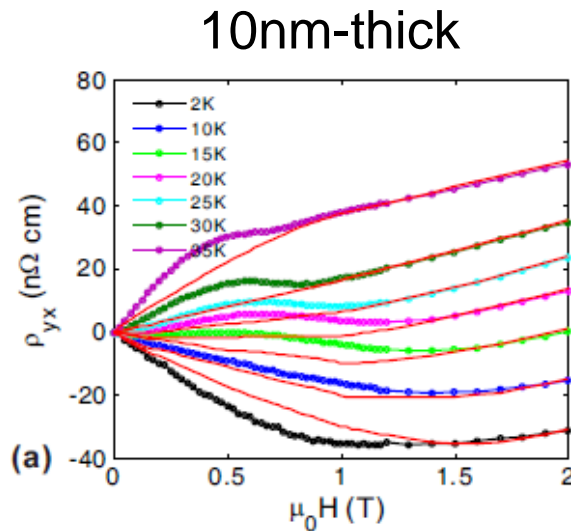
Skyrmion phase mapping by topological Hall resistivity

Yufan Li, Kanazawa, Kagawa et al. PRL (2013)

Conventional anomalous + normal Hall effects

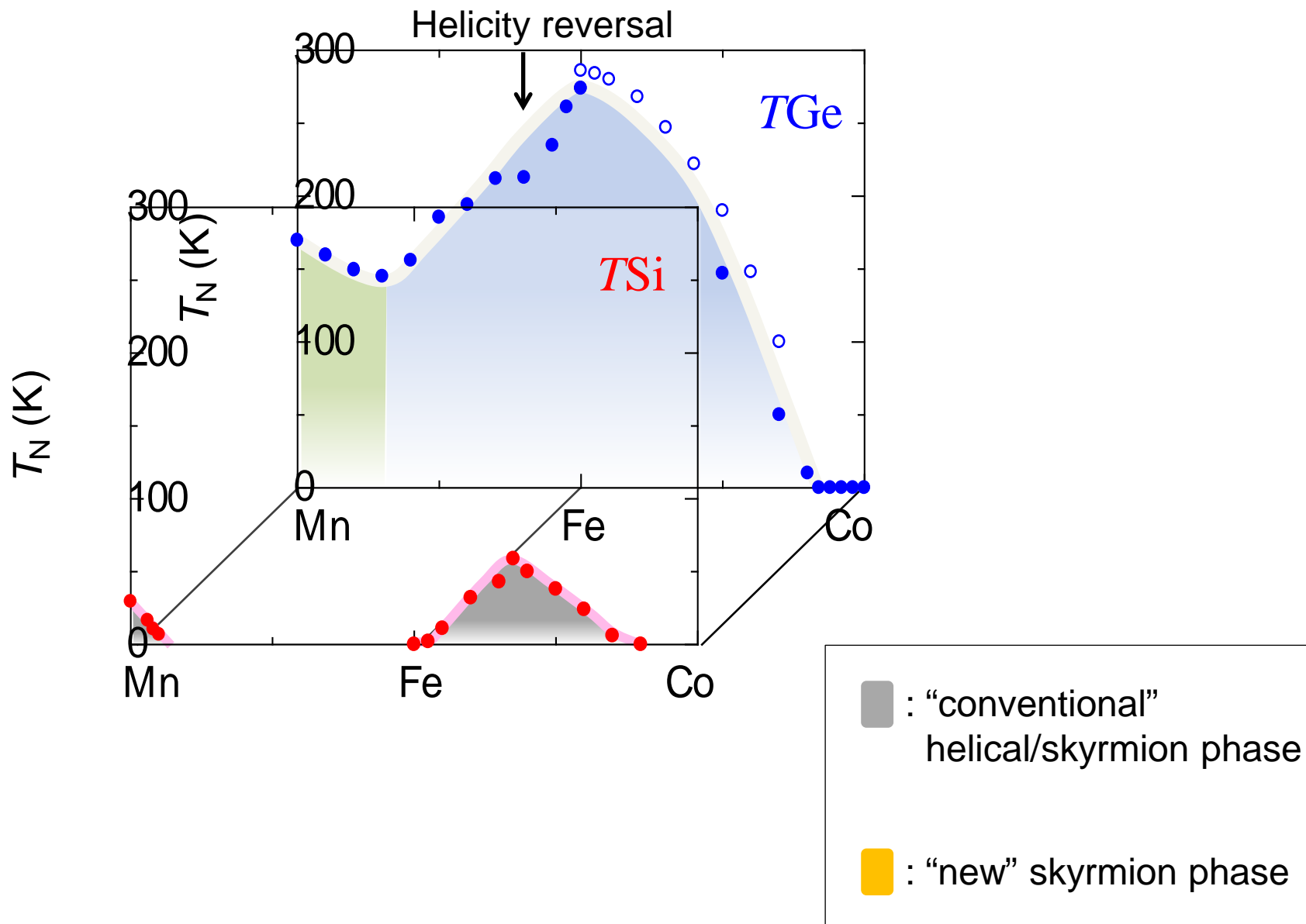
$$\rho_{yx}^A = R_0 B_z + \mu_0 R_S M_z$$

$$\mu_0 R_S = S_A \rho_{xx}^2$$



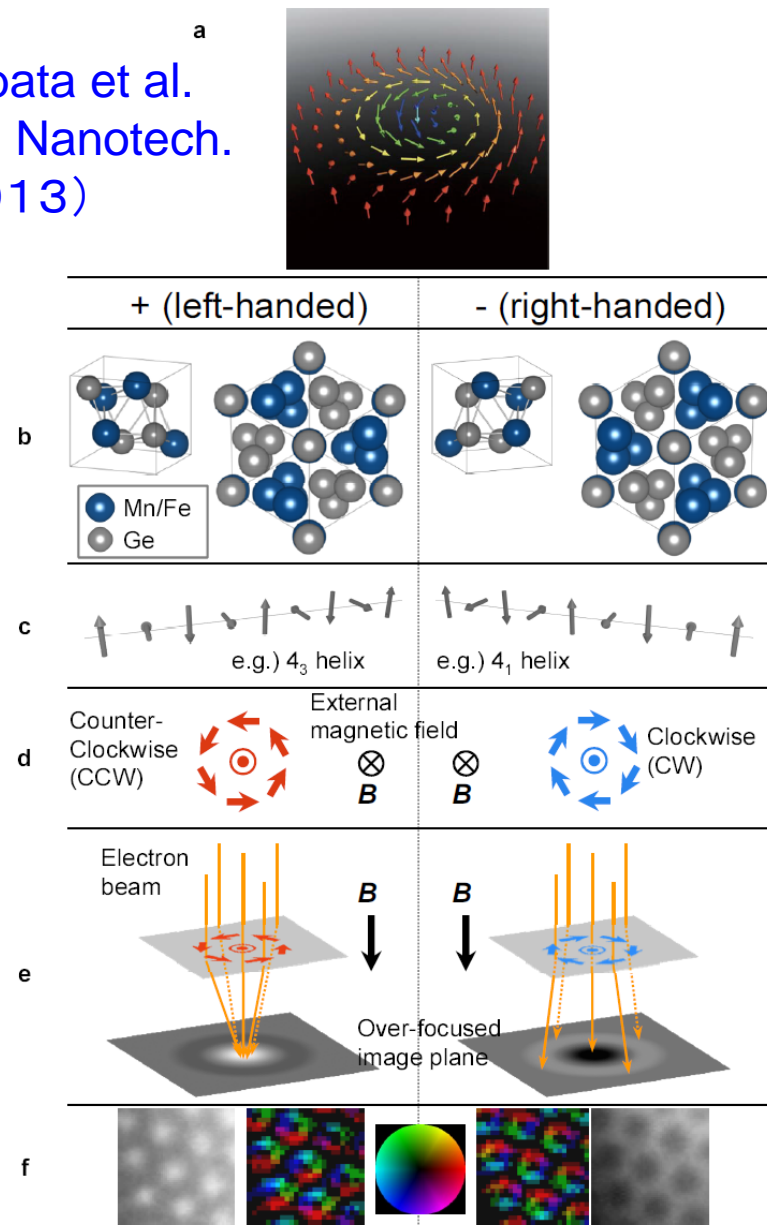
See also the late paper on FeGe thin film;
S. X. Huang and C. L. Chien, Phys. Rev. Lett. **108**, 267201 (2012)

Magnetic phase diagram in *B20* compounds

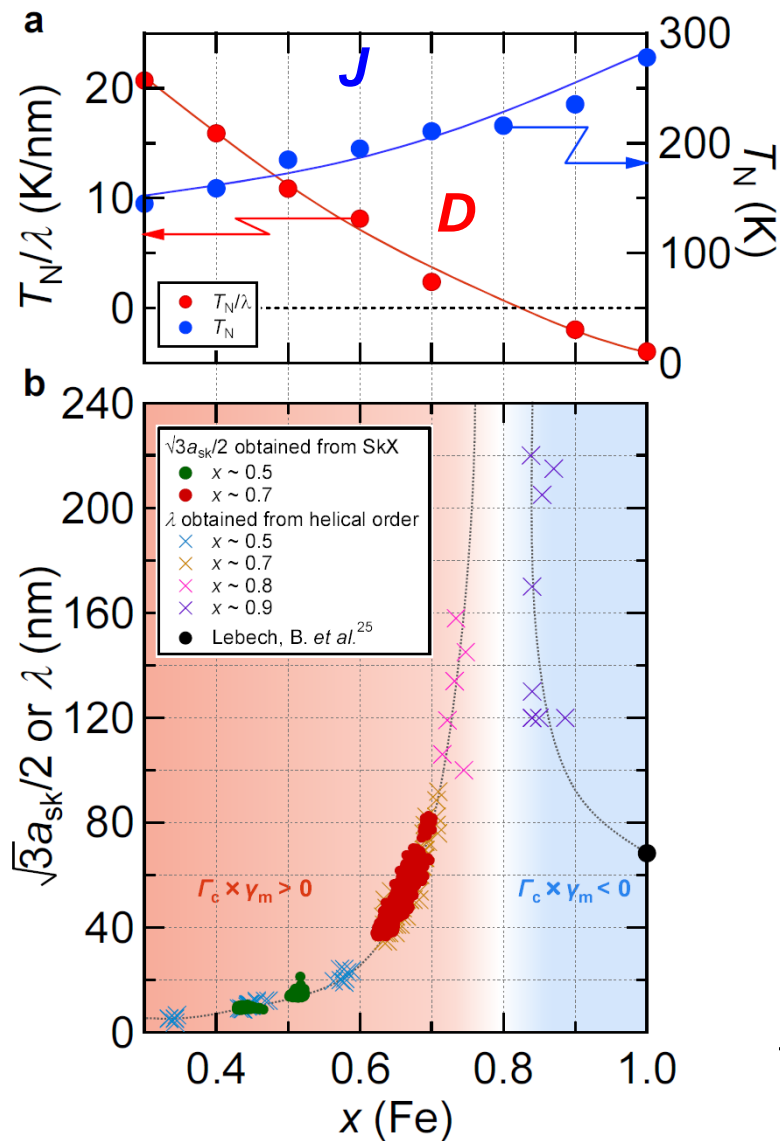


Mn_{1-x}Fe_xGe (Control of DM interaction)

Shibata et al.
Nat. Nanotech.
(2013)



$$H = \sum \left(-J\vec{S}_i \cdot \vec{S}_j + \vec{D}_{ij} \cdot (\vec{S}_i \times \vec{S}_j) \right)$$



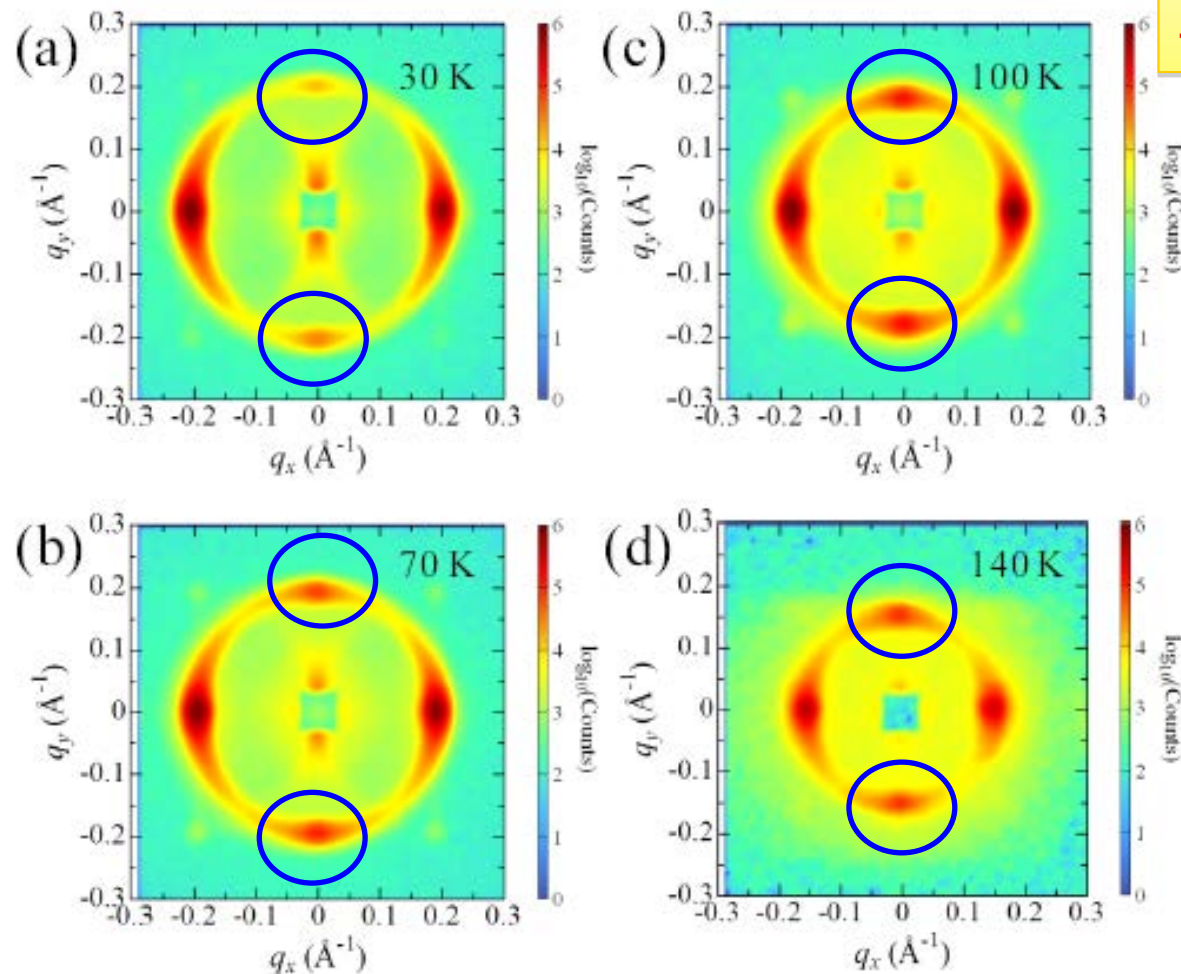
Small angle neutron scattering on MnGe (polyXtal)

B (10T) then B=0

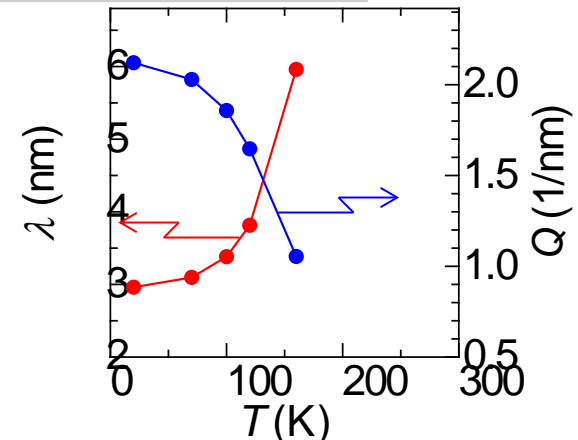


After application of high magnetic field

- $T_N = 170$ K
- Helical structure
- modulation vector : $Q \parallel \langle 100 \rangle$
- Helical period : $\lambda = 3$ nm – 6 nm



Helical period

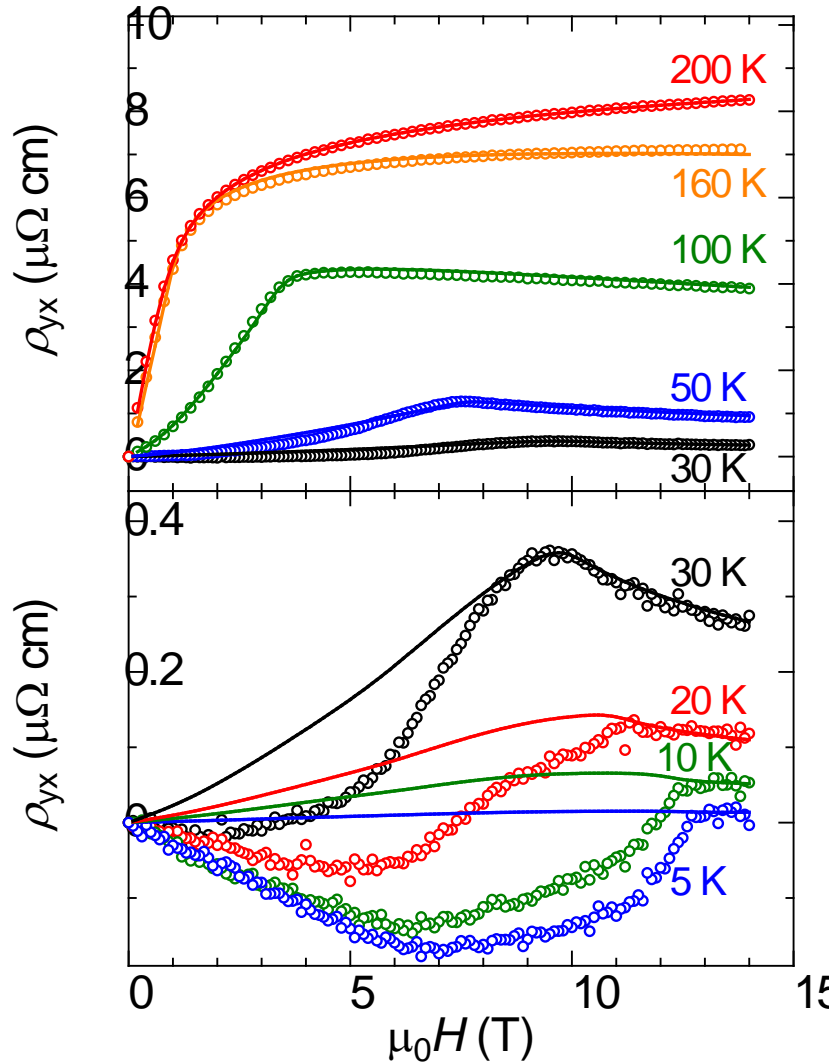


Huge topological Hall effect?

Evidence for multiple-q $\langle 100 \rangle$ structure even at B=0

in collaboration with Keimer group

Topological Hall effect in MnGe



$H > H_C$

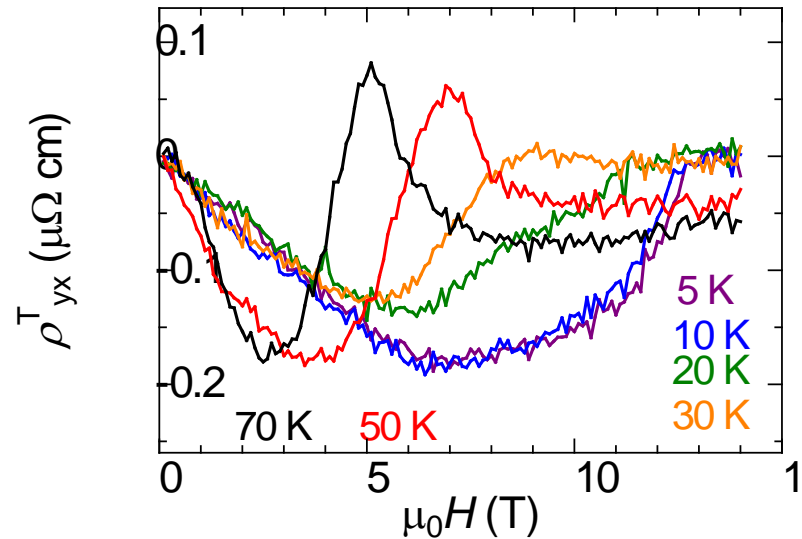
Induced ferromagnetic state
 → “Conventional” anomalous Hall effect

Solid lines: estimate of

$$\rho_{yx}^A = R_0 B_z + \mu_0 R_S M_z$$

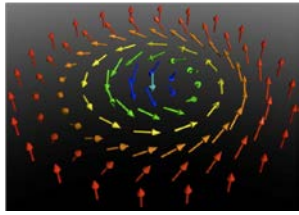
$$\mu_0 R_S = S_A \rho_{xx}^2$$

Components of THE

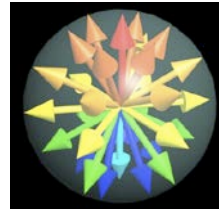


Nearly temperature independent

Real-space fictitious magnetic field in a skyrmion spin texture



Solid angle $\Omega = 4\pi$



In strong coupling case

One skyrmion



One magnetic flux ϕ_0

Emergent magnetic field

$$\mathbf{B}_{\text{eff}}^z = -\phi_0/A$$

A: skyrmion size

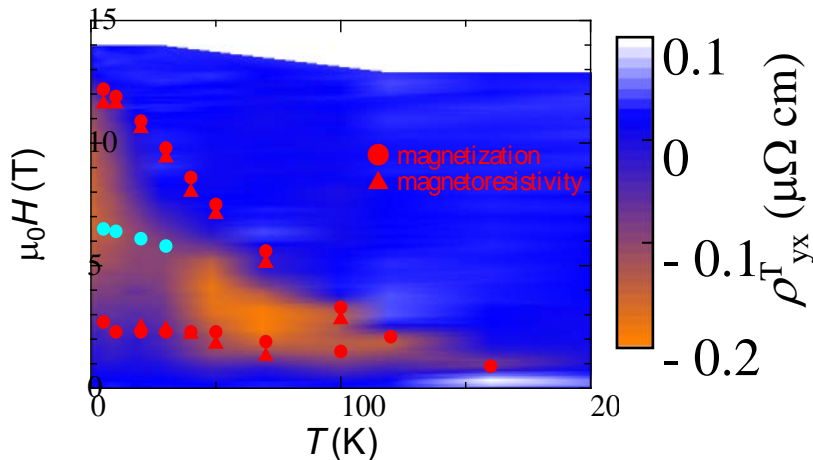
$$\phi_0 = h/e$$

one skyrmion/nm²



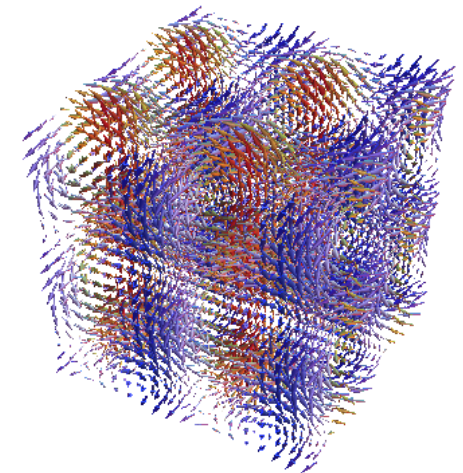
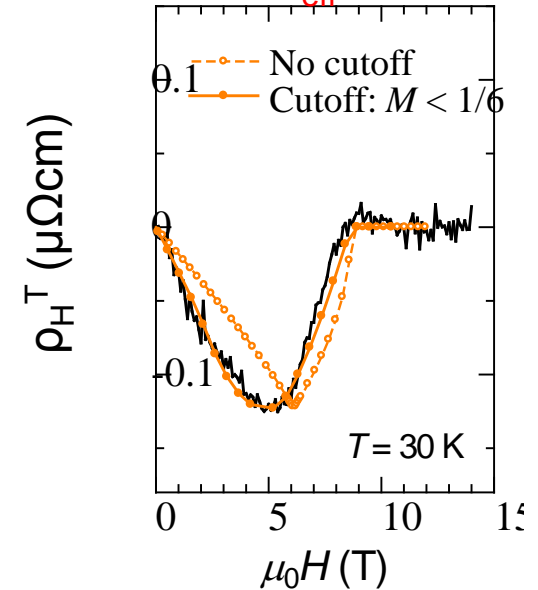
$B_{\text{eff}} \sim 4000\text{T}$

High skyrmion density \Leftrightarrow Large topological Hall Effect



MnGe topological Hall effect

$B_{\text{eff}} \sim 100-1000\text{T}!$



3D skyrmion X

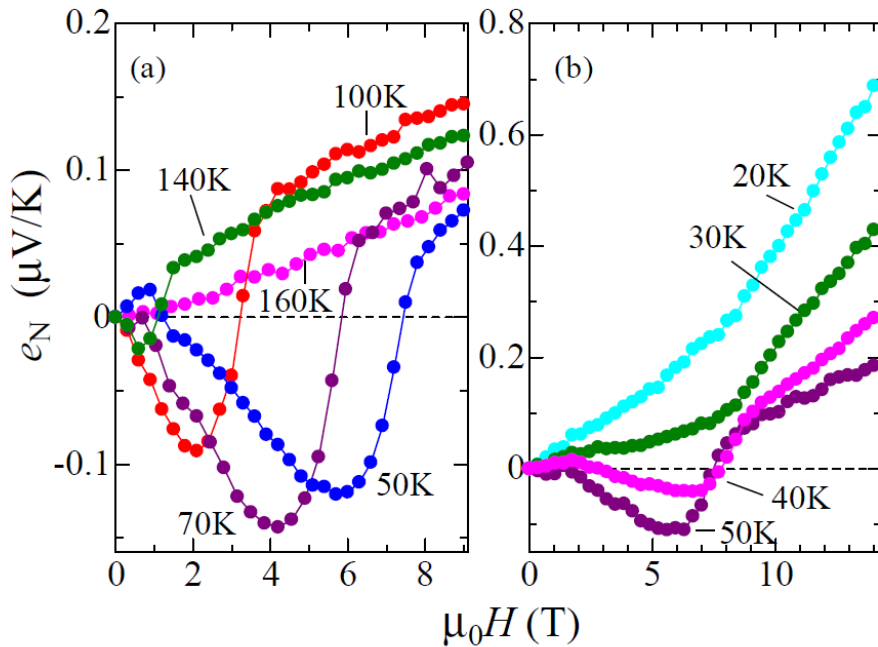
Topological Nernst Effect

Nernst effect

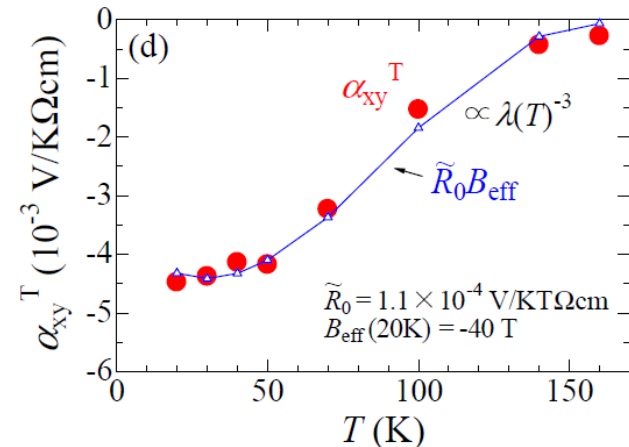
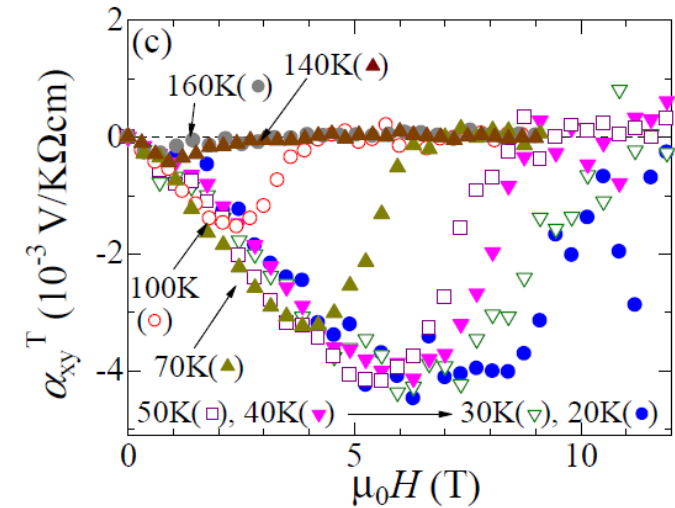
$$e_N = E_y / |\nabla_x T|$$

$$\alpha_{xy} = \sigma_{xx} [e_N + S_{xx} \{(\sigma_{xy}/\sigma_{xx}) + (\kappa_{xy}/\kappa_{xx})\}]$$

$$\text{Mott relation: } \alpha = eL_0T(\partial\sigma/\partial\varepsilon)_{\varepsilon=\mu}$$

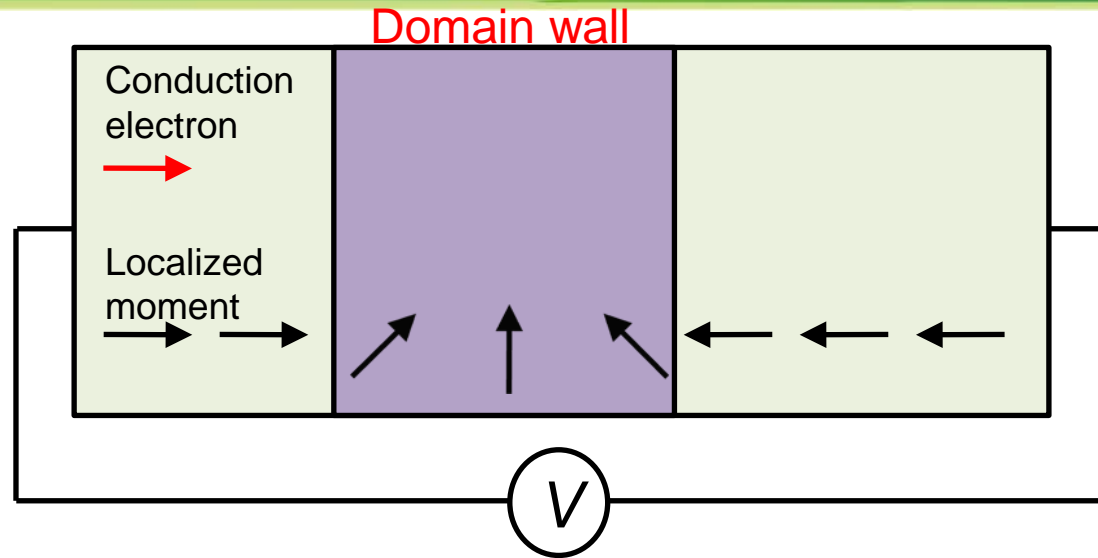


$$\alpha_{xy}^T = \frac{\pi^2 k_B^2 T}{3 e} \left(\frac{\partial \sigma_{xy}^T}{\partial \varepsilon} \right)_{\varepsilon=\mu} \approx \tilde{R}_0 B_{eff}$$

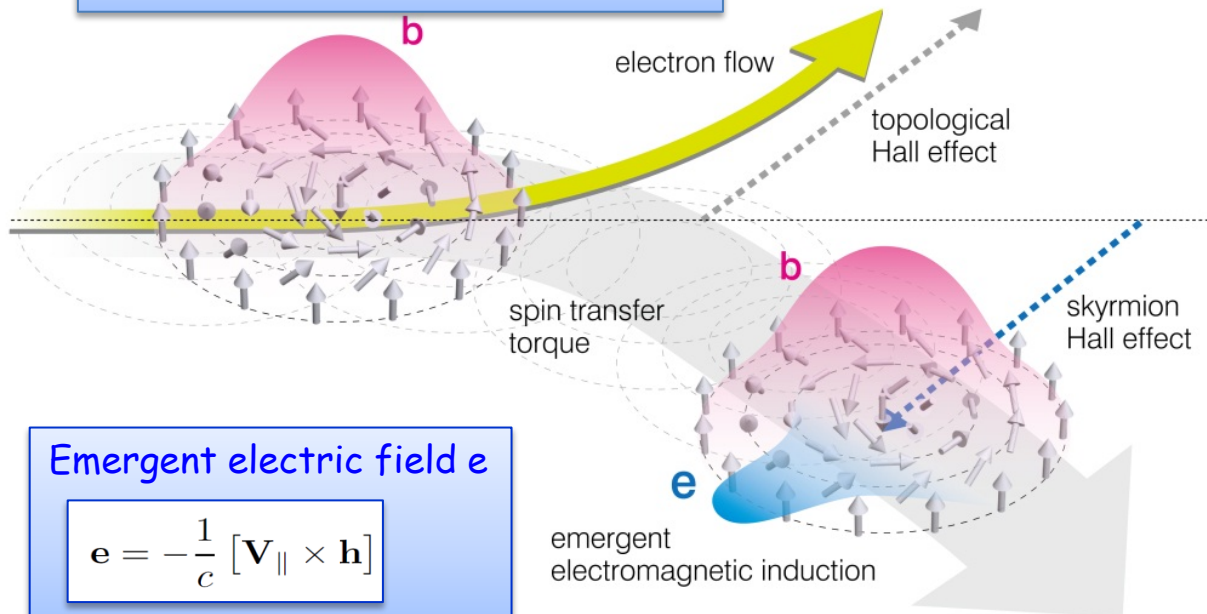


Current drive of skyrmions and emergent EM field

Domain wall motion
by spin transfer torque



Topological Hall effect (THE)
Emergent magnetic field b



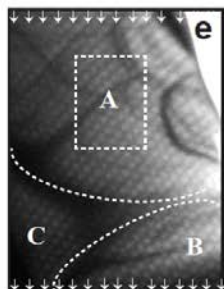
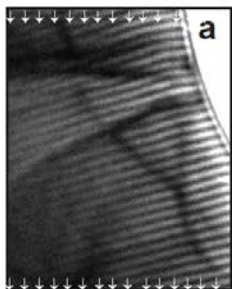
-counteraction of THE
→ skyrmion Hall effect

Emergent electric field e

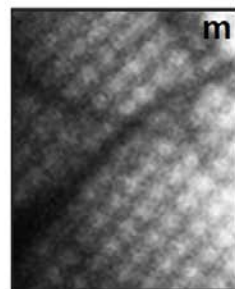
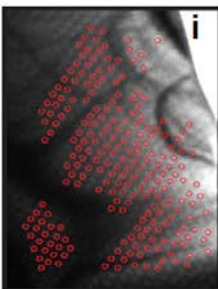
$$e = -\frac{1}{c} [\mathbf{V}_{\parallel} \times \mathbf{h}]$$

$T = 250 \text{ K}$ $B = 0$

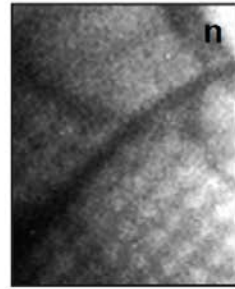
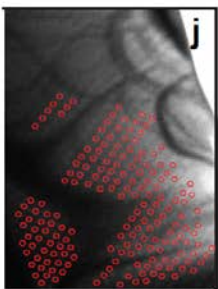
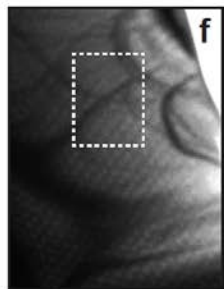
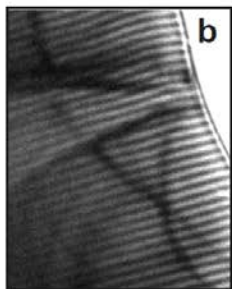
$t = 0$
 $I = 0$



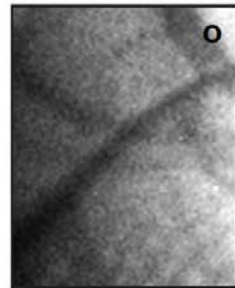
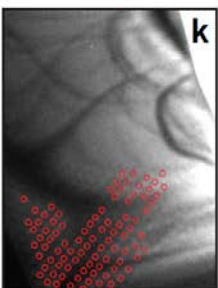
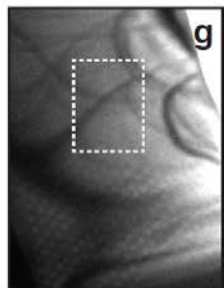
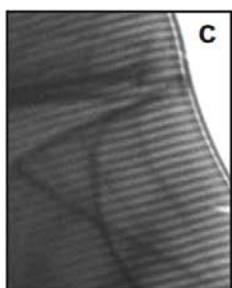
$B = 150 \text{ mT}$



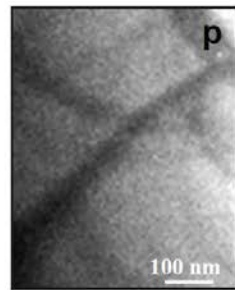
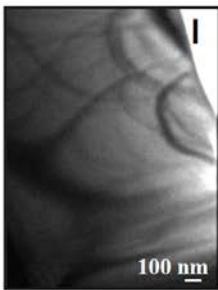
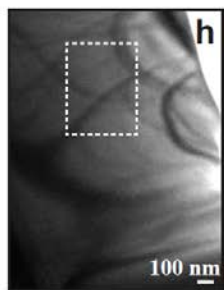
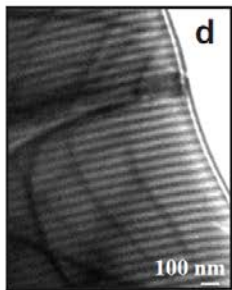
$t = 48 \text{ s}$
 0.41 mA



$t = 52 \text{ s}$
 0.50 mA



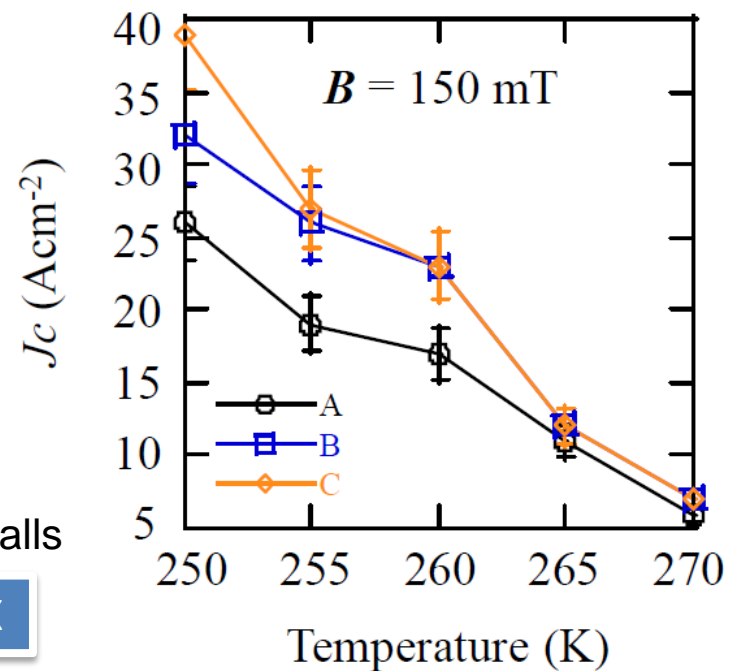
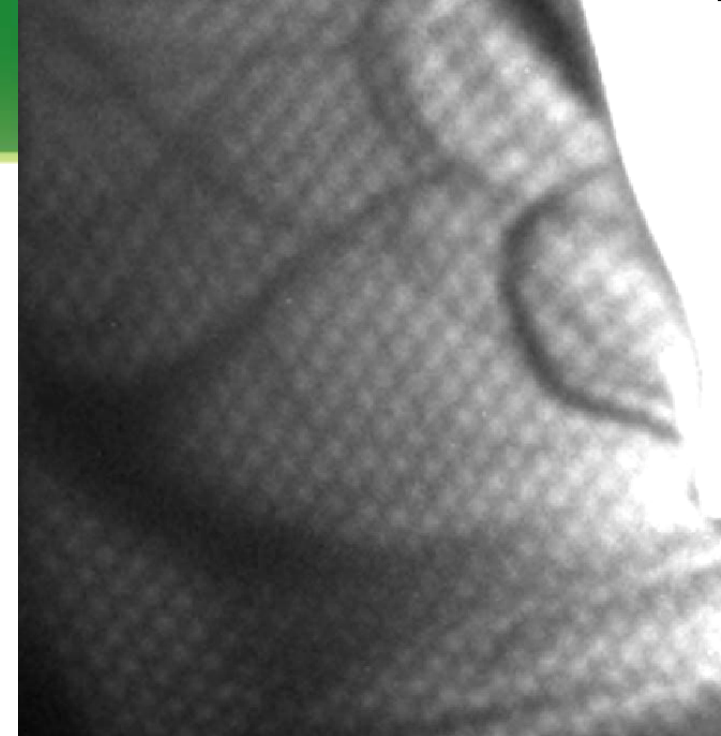
$t = 55 \text{ s}$
 0.61 mA



$J_c < 100 \text{ A/cm}^2$

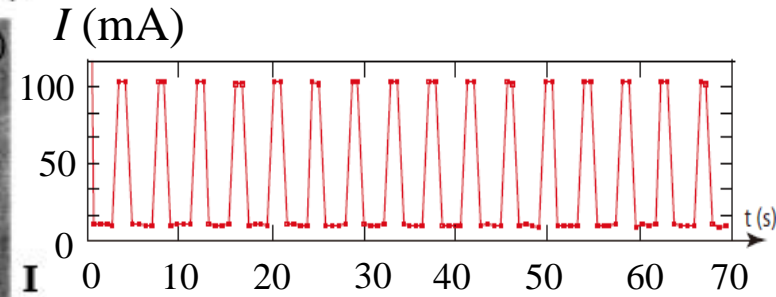
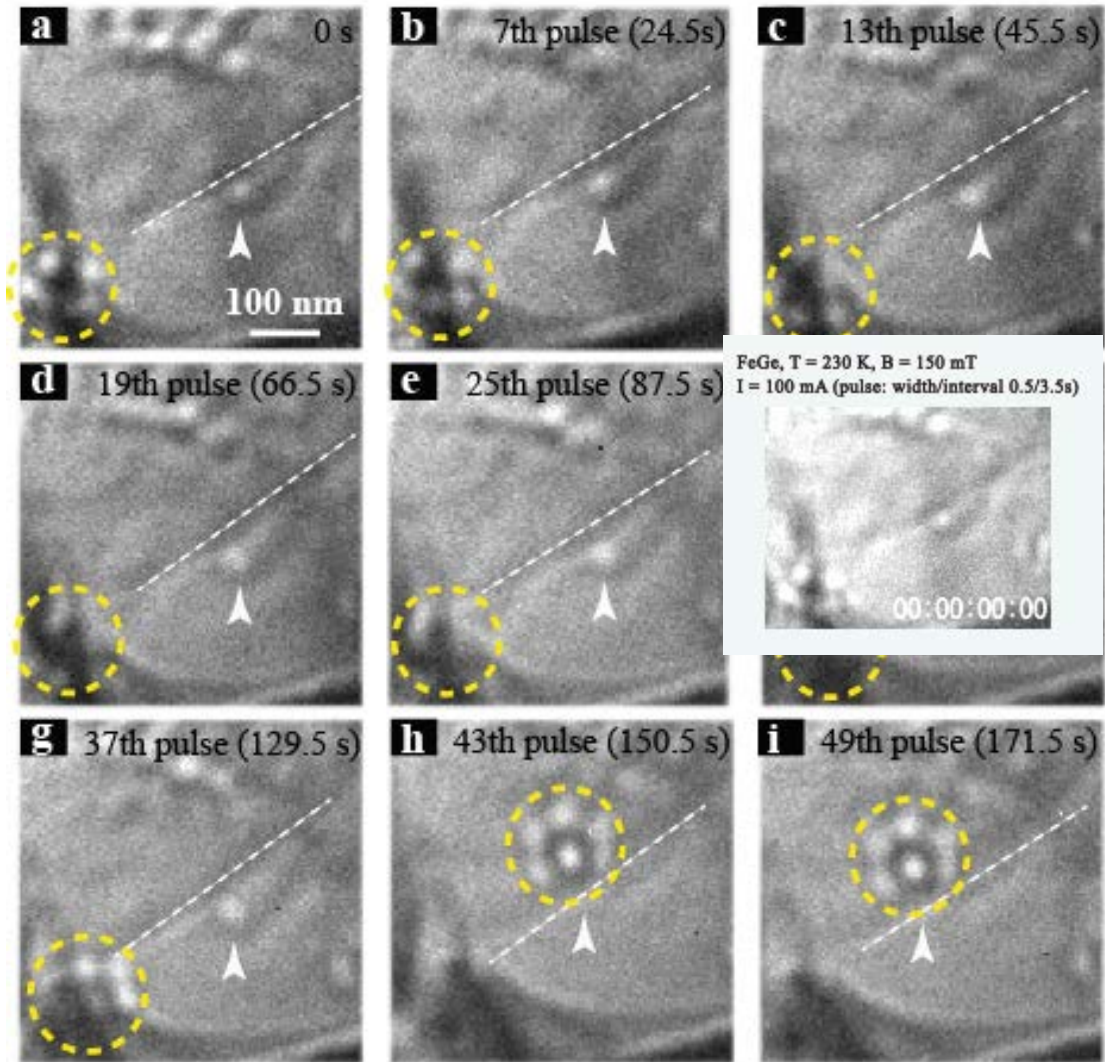
$J_c \sim 10^7 \text{ A/cm}^2$ for ordinary domain walls

no intrinsic / minimal extrinsic pinning effect on SkX



Manipulation of single skyrmion in FeGe by pulse currents

$T = 230 \text{ K}$, $B = 150 \text{ mT}$, $I = 100 \text{ mA}$ (pulse width/interval = 0.5/3.5 s)



FeGe, $T = 230 \text{ K}$, $B = 150 \text{ mT}$
 $I = 100 \text{ mA}$ (pulse: width/interval 0.5/3.5s)



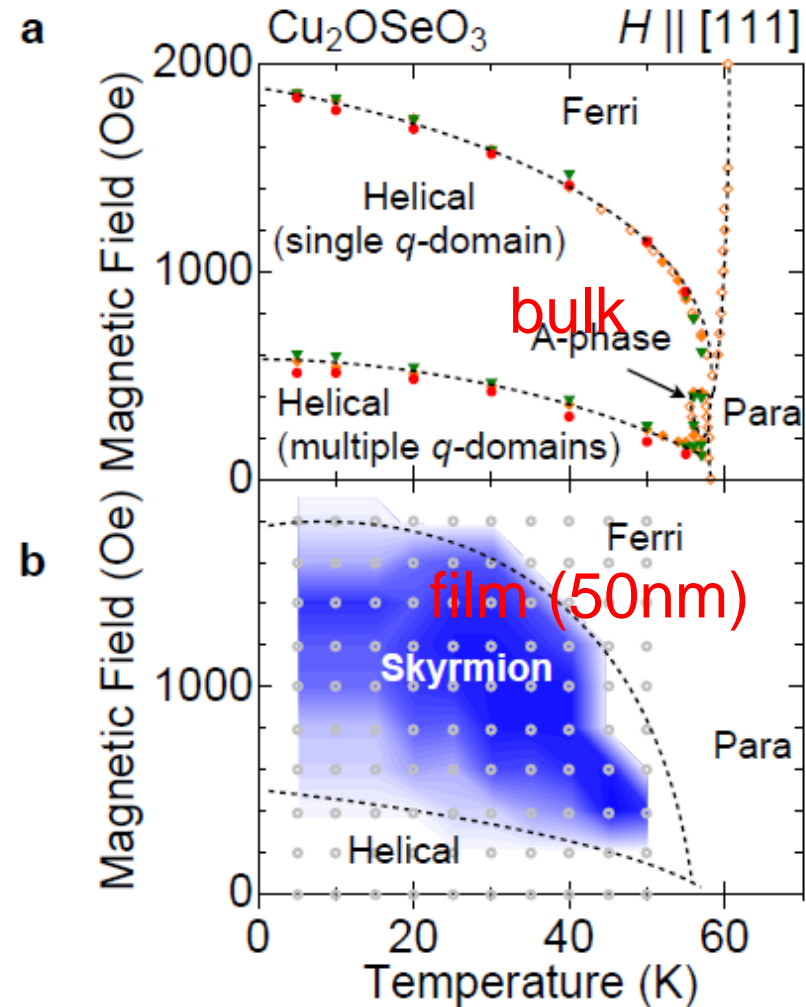
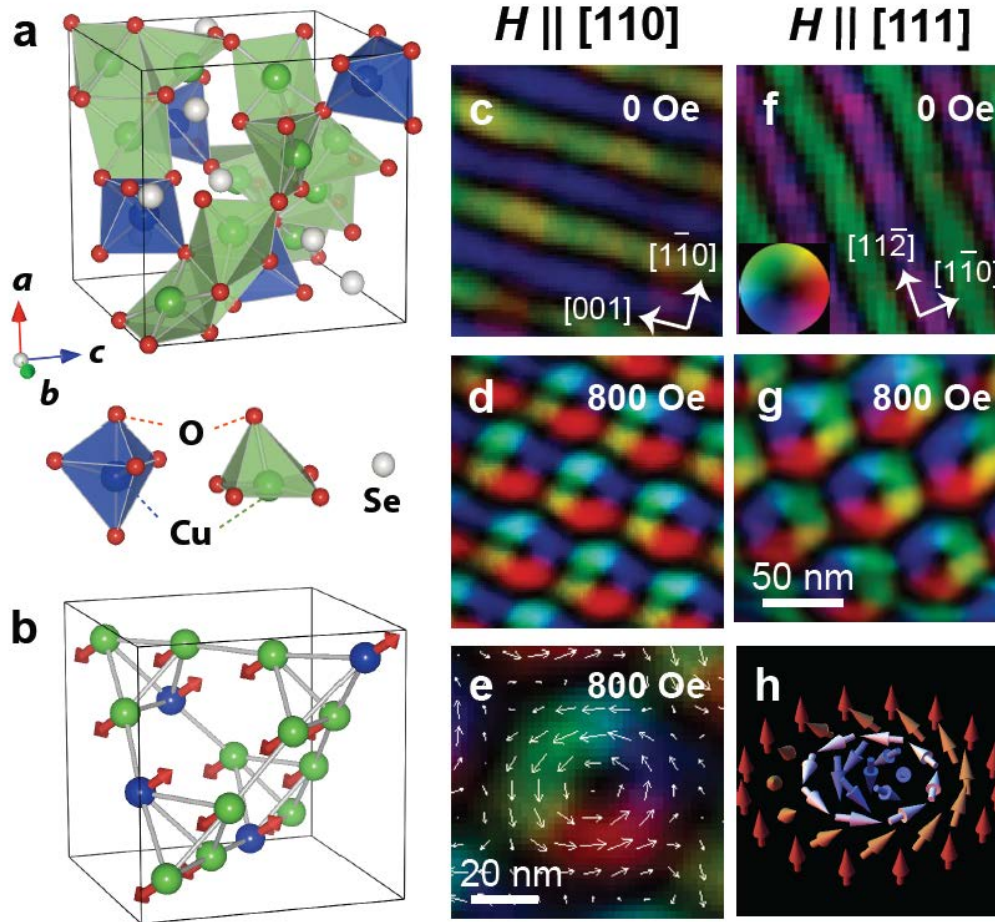
Lorentz TEM observation of thin flake of Cu_2OSeO_3

space group

$P2_13$

the same as
B 20 (MnSi)

Mott insulator with quantum spins



Cu₂OSeO₃ : *P* distribution in skyrmion

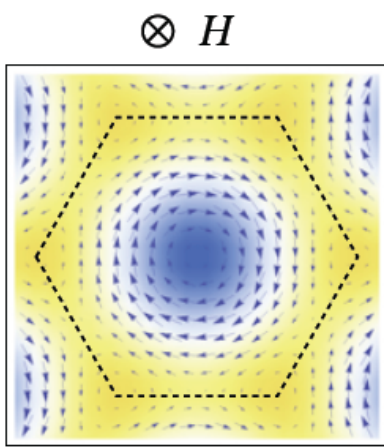
S. Seki *et al.*, PRB(2012)

d-p hybridization model

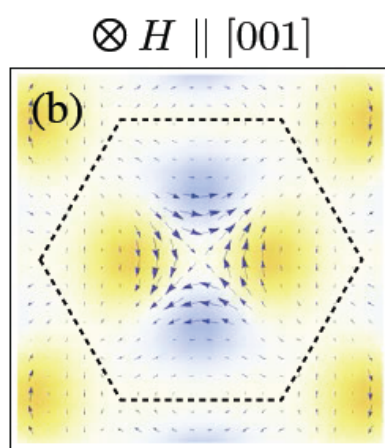
$$\vec{p}_{ij} \propto (\vec{e}_{ij} \cdot \vec{m}_i)^2 \vec{e}_{ij}$$

Local M

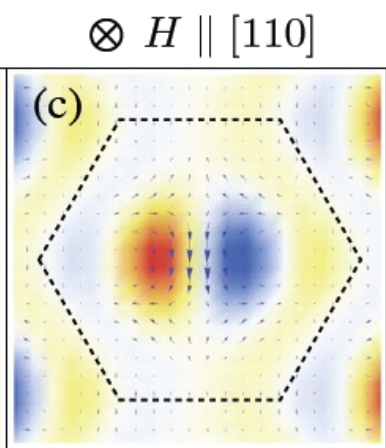
Local P



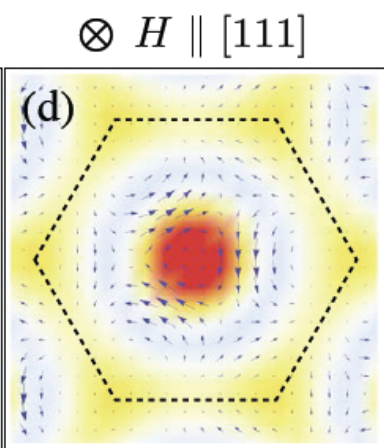
-1 *m^z, ρ* +1



$\begin{matrix} \rightarrow [\bar{1}10] \\ \downarrow [\bar{1}\bar{1}2] \end{matrix}$ ***P* = 0**



$\begin{matrix} \rightarrow [\bar{1}10] \\ \downarrow [001] \end{matrix}$ **↓ *P***

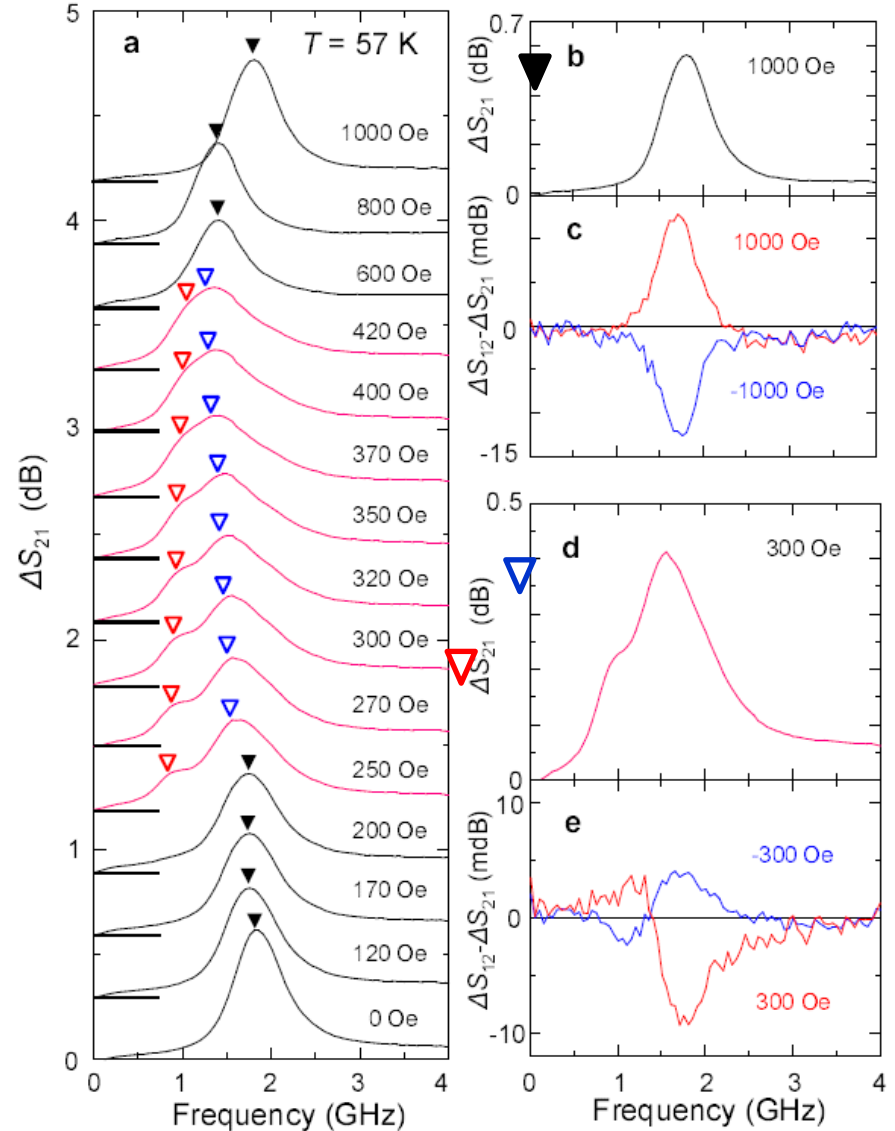
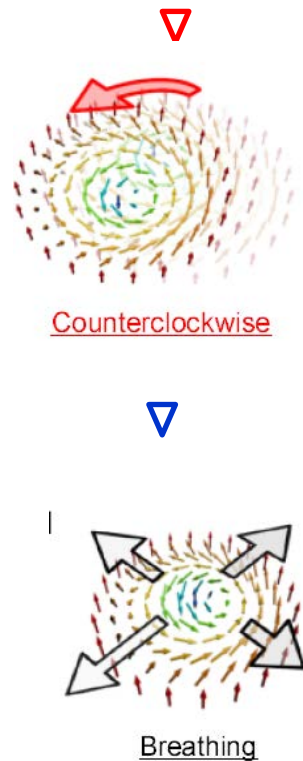
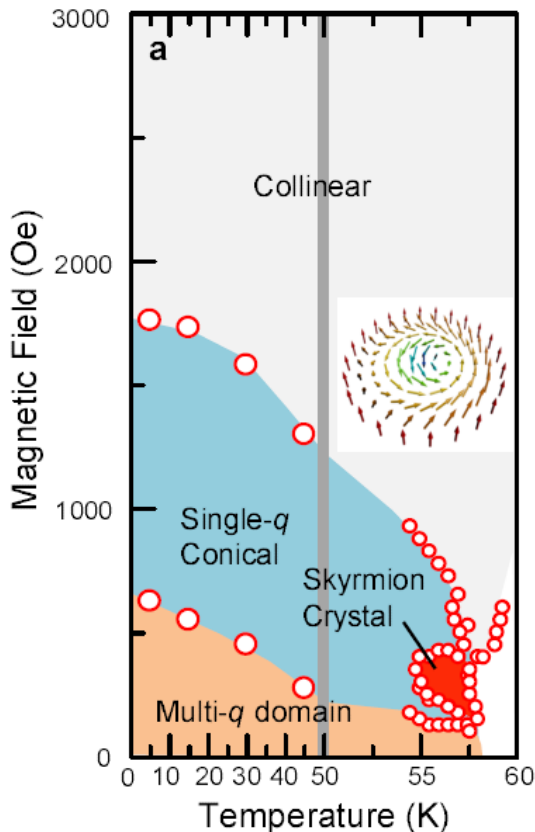
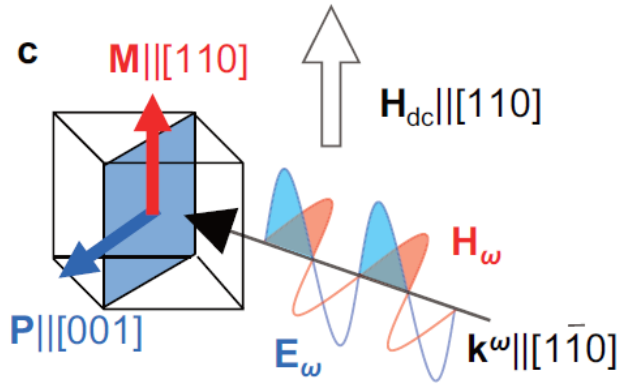


$\begin{matrix} \rightarrow [\bar{1}10] \\ \downarrow [\bar{1}\bar{1}0] \end{matrix}$ **⊗ *P***

Skyrmion particle can locally carry **electric dipole** or **quadrupole**

Skyrmion excitations as electromagnons showing directional dichroism

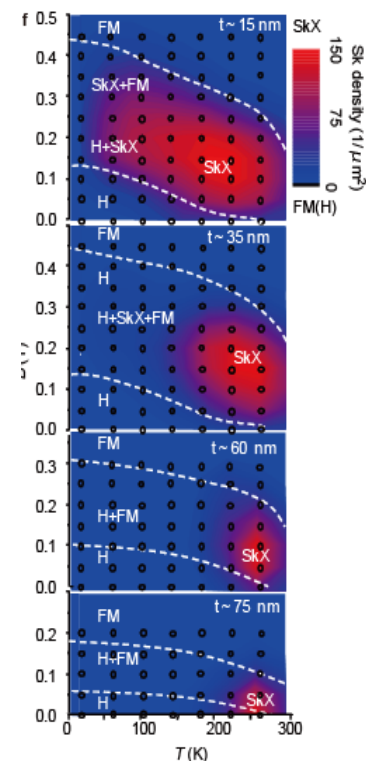
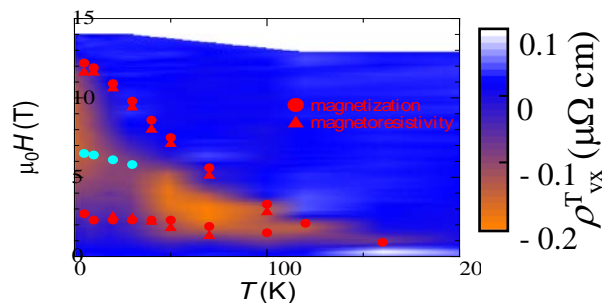
Y. Okamura et al. Nature Commun. (2013).



Stabilization Skyrmions in form of thin films, not only DM but also uniaxial magnets

Topological Hall effect as probe/detection for SkX

emergent EM fields



G. 3 Yu et al.

Skyrmion transport phenomena

- low-current drive of Skyrmions ($< 100 \text{ A/cm}^2$)
processing speed $\propto I^*$ (Sk density); energy-cost per bit $\propto I$
- optical, e-beam (spin-current) control,; E-drive (multiferroics)