

Instability of a falling jet of concentrated suspensions

P. Hébraud (IPCMS, Strasbourg)

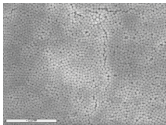
M. Liard (IPCMS, Strasbourg, Sika,Zürich), J. Sautel (IPCMS, Strasbourg, ENS, Lyon), A. Sato (IPCMS, Strasbourg), D. Lootens (Sika,Zürich)

Outline

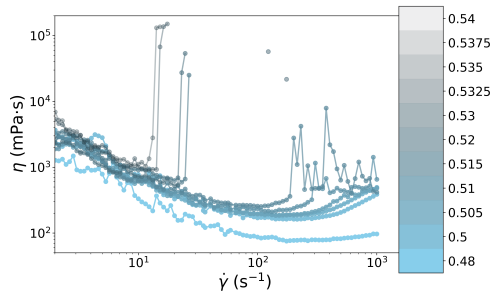
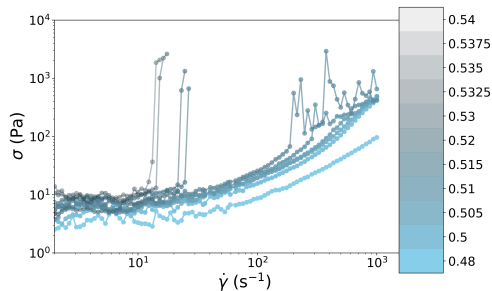
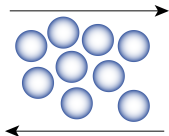
- 1 Introduction : shear thickening
 - Shear thickening in extension
 - Arrested state life time
- 2 Falling jet of a suspension : solid-like regime
 - Jet shape
 - Weight loss
 - Transverse waves propagation
- 3 Conclusion

Shear thickening of concentrated suspensions

- $d=2 \mu\text{m}$ silica particles,
- water pH=7

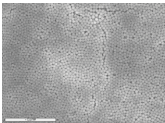


- Cone-plate
- Controlled shear rate

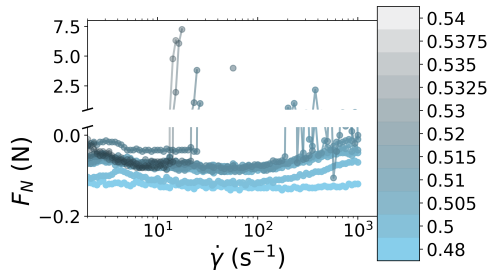
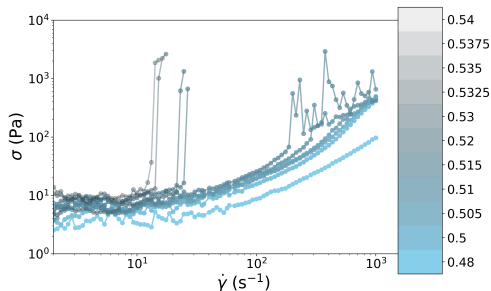
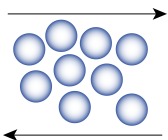


Shear thickening of concentrated suspensions

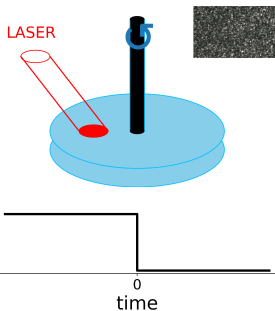
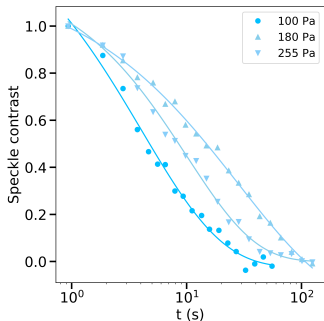
- $d=2 \mu\text{m}$ silica particles,
- water $\text{pH}=7$



- Cone-plate
- Controlled shear rate



Life time of the arrested state



- rough silica particles :

- diameter : 650 nm
- roughness : $\sqrt{\langle \delta r^2 \rangle} = 6.2 \text{ nm}$

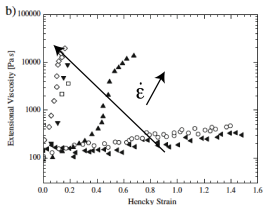
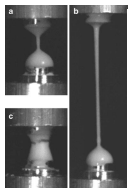
F.Ianni et al P.R.E., 2006

- Mustispeckle Diffusing Wave Spectroscopy
- Relaxation of the contrast of the speckles after a stress step

Shear thickening under extension

Suspensions of corn starch particles

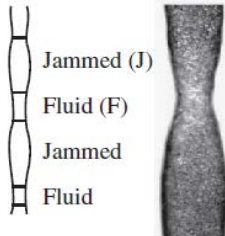
- $\dot{\epsilon}$ imposed
- Force and filament widths measured



$\phi = 55 \%$

E. E. Bischoff White, M. Chellamuthu, J.P. Rothstein, *Rheol. Acta*, 2010

Observation of the extended filament



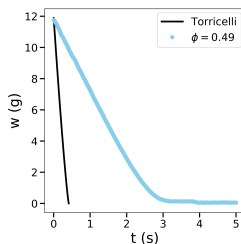
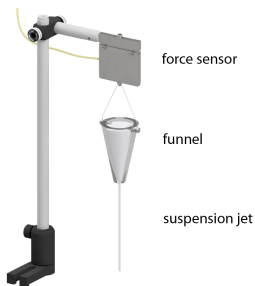
M. Roché, H. Kellay, H.A. Stone, *P.R.L.*, 2011

- pulling a column of sample
- heterogeneities : jammed regions linked with liquid bridges

Falling jet of a suspension

How does the suspension flow under its own weight ?
If g is sufficient :

- Flow under g \rightarrow shear-thicken
- Flow stops \rightarrow Relaxes back to the liquid state



Acquisition rate 500 im/s

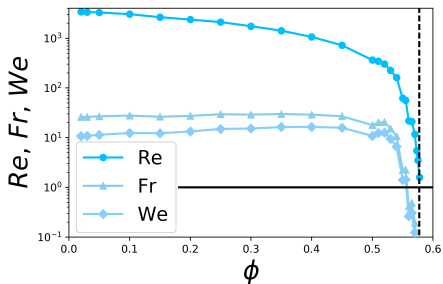
Non dimensional numbers

$$Re = \frac{\text{Inertia}}{\text{Viscosity}} = \frac{\rho U a}{\eta}$$

$$Fr = \frac{\text{Inertia}}{\text{Gravity}} = \frac{U^2}{ga}$$

$$We = \frac{\text{Inertia}}{\text{Surface}} = \frac{\rho U^2 a}{\gamma}$$

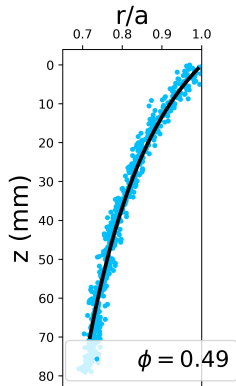
Flow regimes



- small ϕ : $Re, Fr, We \gg 1$: inertial regime,
- $\phi \rightarrow \phi_m$: surface effects

Inertial regime

Shape of the jet



Assuming a perfect fluid :

- Bernoulli's theorem
- Laplace pressure

$$\begin{aligned}\frac{r(z)}{r_0} &= \left[1 + \frac{2}{Fr} \frac{z}{a} + \frac{2}{We} \left(1 - \frac{a}{r(z)} \right) \right]^{-1/4} \\ &= \left[1 + \frac{2}{Fr} \frac{z}{a} \right]^{-1/4}\end{aligned}$$

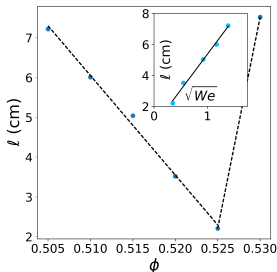
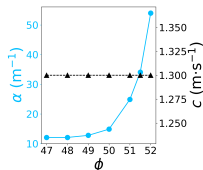
Capillary regime

Growth of the capillary wave amplitudes :

$$r(z) = r_0 + \epsilon e^{\alpha z}$$

- Constant celerity $c = v_{jet}$
- Decreasing α^{-1}

$$L \propto \frac{1}{\alpha} = \frac{U}{\sqrt{\Omega_c}} \propto \sqrt{We}$$



At high ϕ , the jet ruptures at longer length.

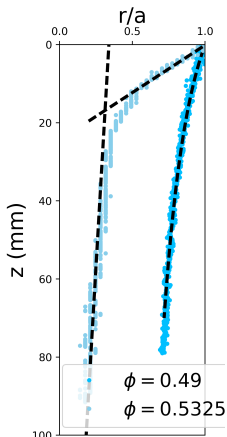
Jet shape in the "solid-like" regime

Jet instability

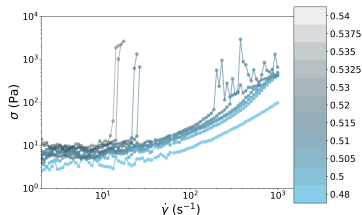


speed $\approx 1/10$
 $\phi = 0.535$

Average jet shape



- Rapid decrease of the jet radius at the output of the funnel,
- Thinner jet, smaller radius decrease

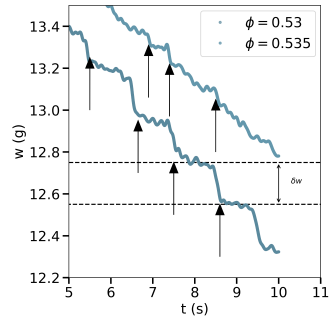
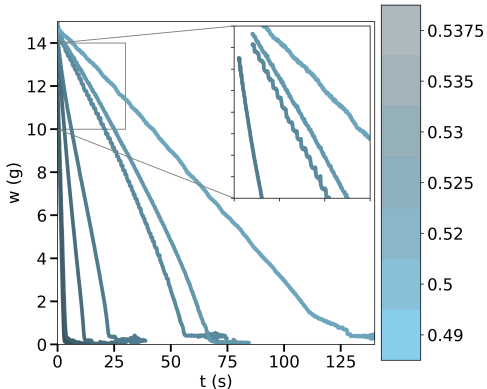


At the output of the funnel :

$$\dot{\epsilon} \approx 10^2 \text{ s}^{-1}$$

$1 \gg We, Fr$: inertia irrelevant

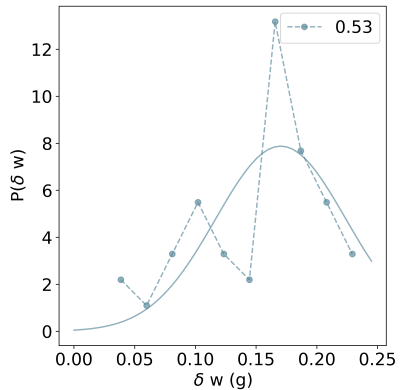
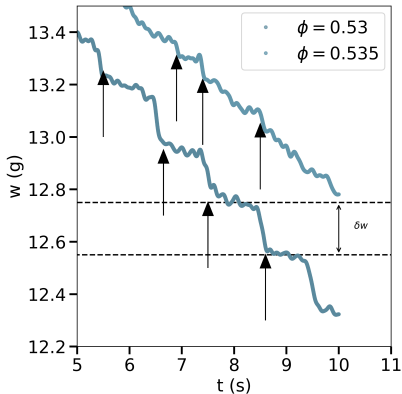
Weight loss in the "solid like" regime



At $\phi > 53\%$:

- the weight of the entire column is measured : tensile stress is transmitted
- well defined weight scales

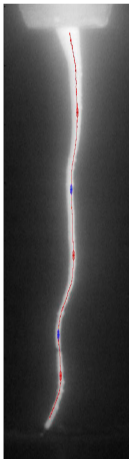
Weight loss in the "solid like" regime



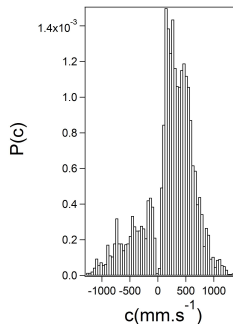
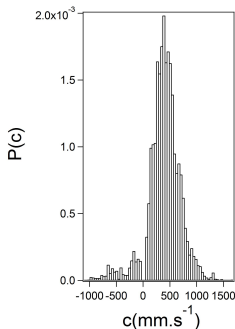
$$\langle \delta w \rangle = 0.17 \text{ g lineic mass} : 4.7 \cdot \text{g} \cdot \text{m}^{-1} \longrightarrow \delta \ell = 3.4 \text{ cm}$$

Wave propagation : solid rope model

Identify and follow maxima



Distribution of velocities



$\phi = 53.25 \%$, $\phi = 53.5 \%$

Superposition of up and down travelling waves

$c_{max} \approx 0.5 \text{ m} \cdot \text{s}^{-1}$ whereas $v_{jet} \approx 0.1 \text{ m} \cdot \text{s}^{-1}$

Wave propagation along the rope

Wave propagation along a solid rope

$$\frac{\partial^2 x}{\partial t^2} = \frac{T(x)}{\mu} \frac{\partial^2 x}{\partial z^2}$$

$$c = \sqrt{\frac{T}{\mu}} = \sqrt{Lg} \approx 1 \text{ m} \cdot \text{s}^{-1}$$

We observe :

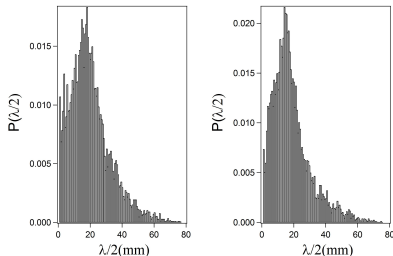
$$A(x, t) = \sum A_i \cos(\omega t \pm kx)$$

Assuming one wave travelling up and

one down :

Reflexion coeff : $R \approx 0.5$

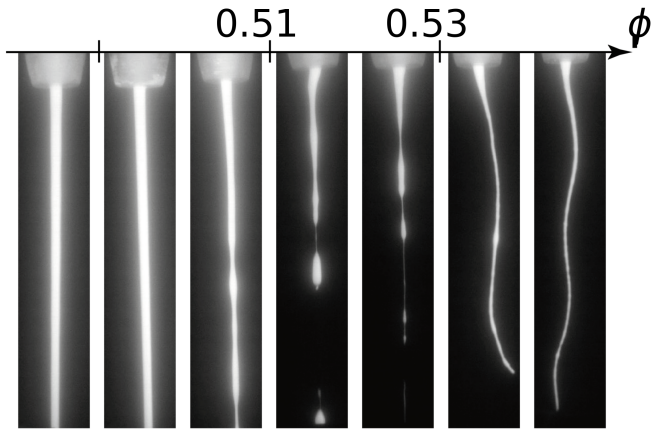
Distribution of wavelengths



$$\phi = 53.25 \%, \phi = 53.5 \%$$

$$\lambda_{max} = 20 \text{ mm}$$

Conclusion



Conclusion

At the shear thickening transition **under gravity** :

- the jet is heterogeneous
- it sustains tensile stress
- it lets solid chunks fall
- becomes radially unstable