

The Mixing Length Theory

Past, Present & Future

Probes of Transport in Stars
KITP

29 November, 2021

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- please argue with me

Past



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- 1968:** Cox & Giuli present an MLT formalism valid for optically thick material
(**2021**-- Adam Jermyn writes TDC in MESA, which reduces to this, “mostly”)

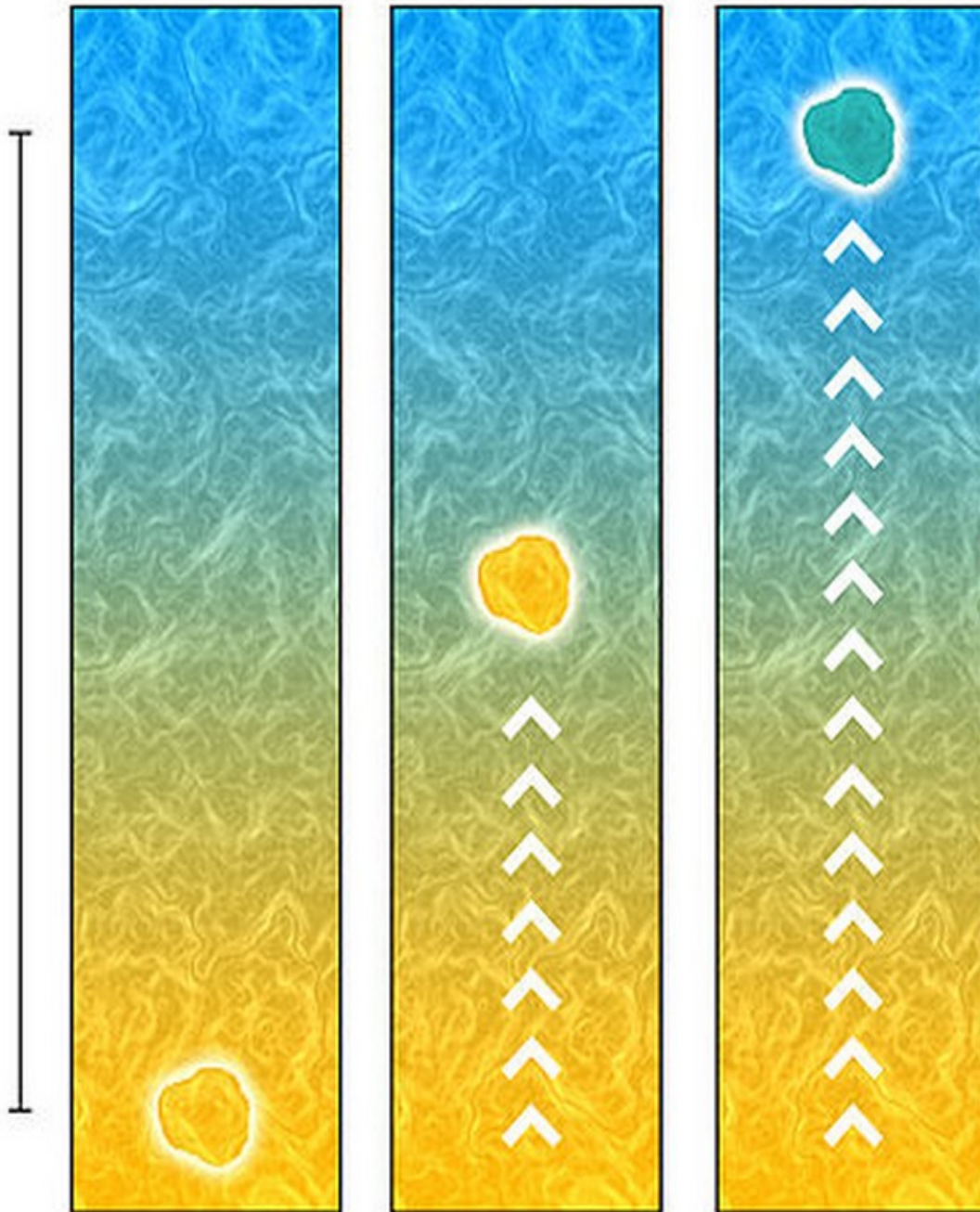
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- 1979:** Mihalas 1978, Kurucz 1979 – study of radiative transfer models of stellar atmospheres in connection to sub-surface convection zones characterized by MLT

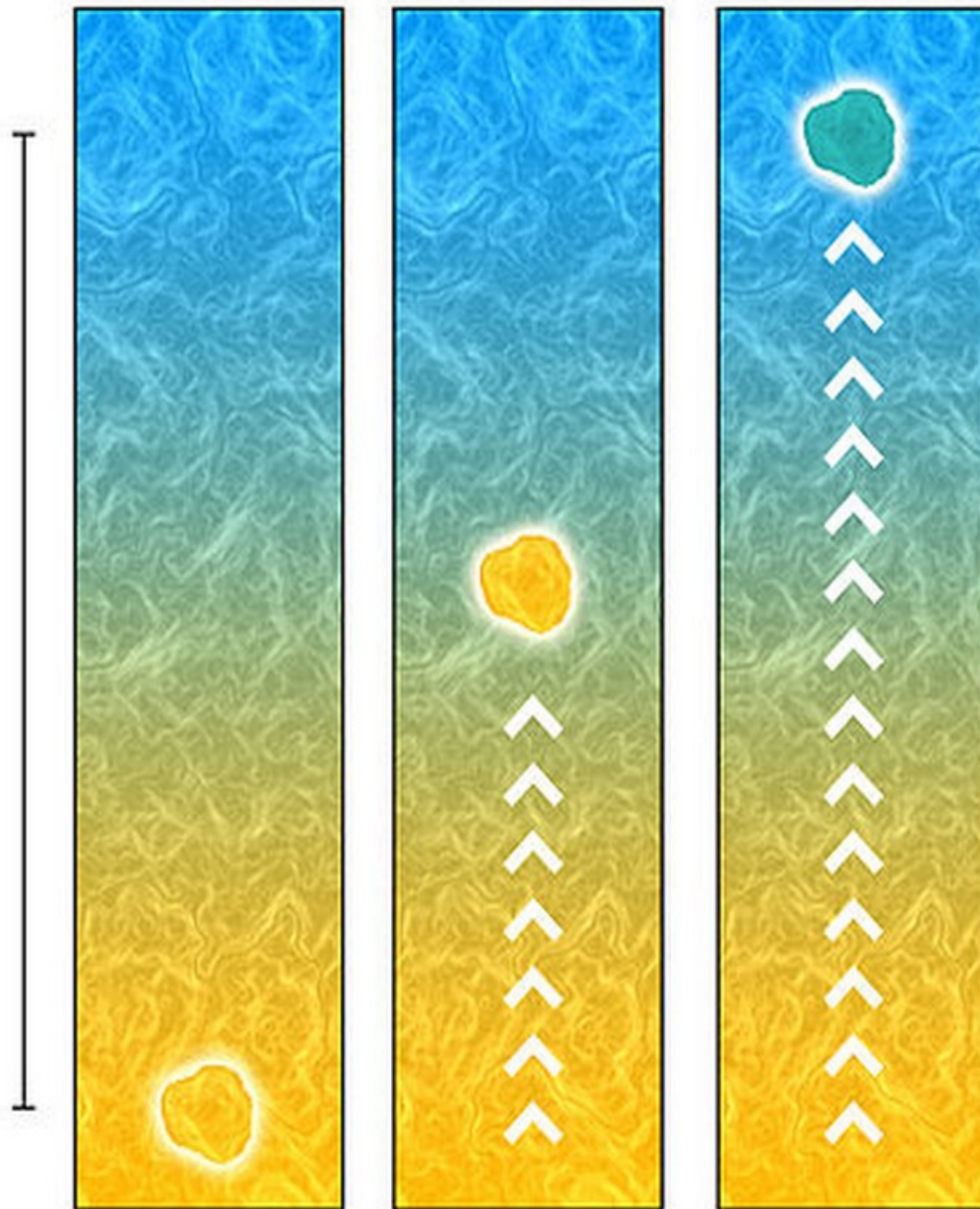
Mixing Length Theory (MLT) Formalism



-discrete parcels consist of fluid which are in pressure, but not thermal, equilibrium

Fig: Wikipedia commons public domain

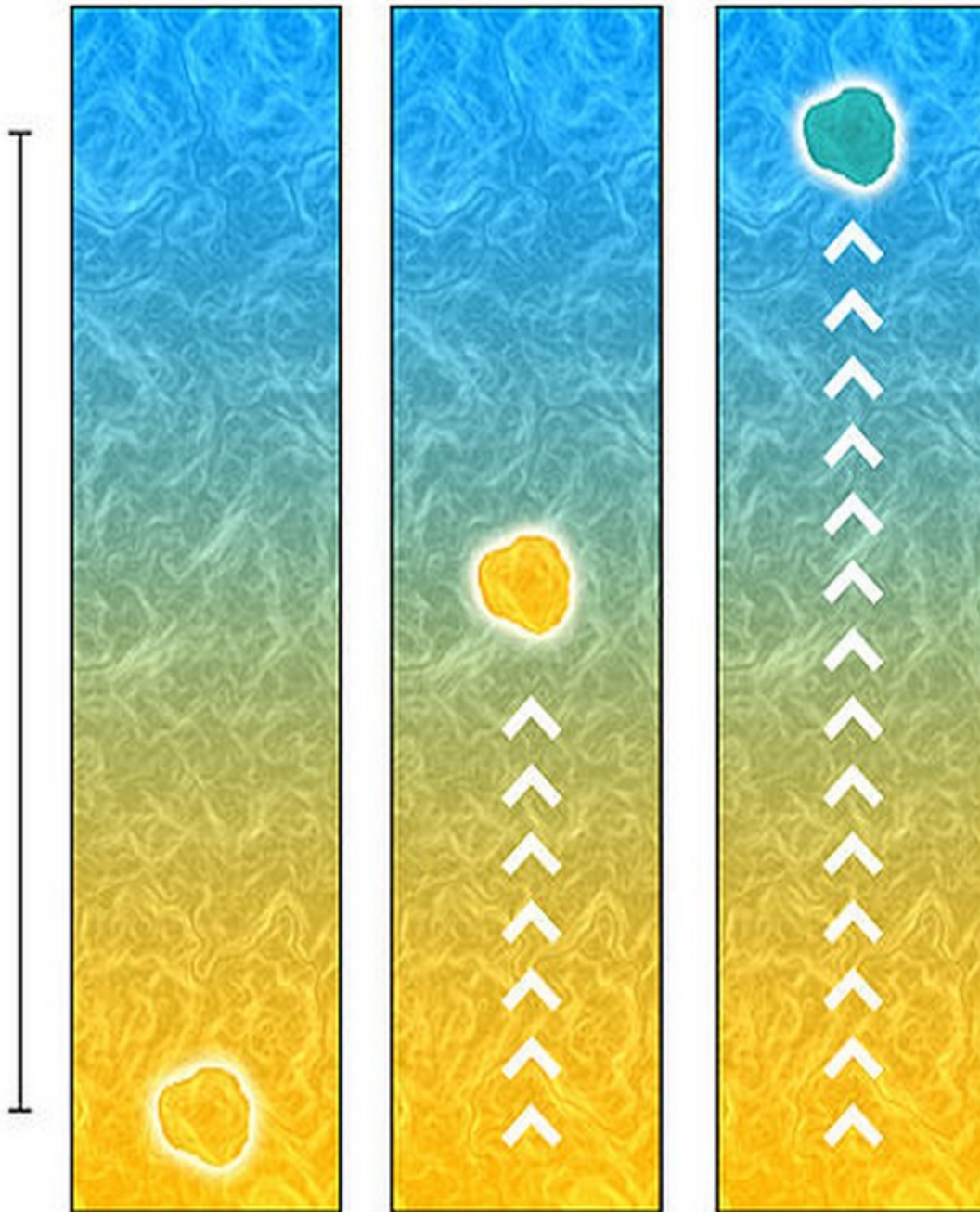
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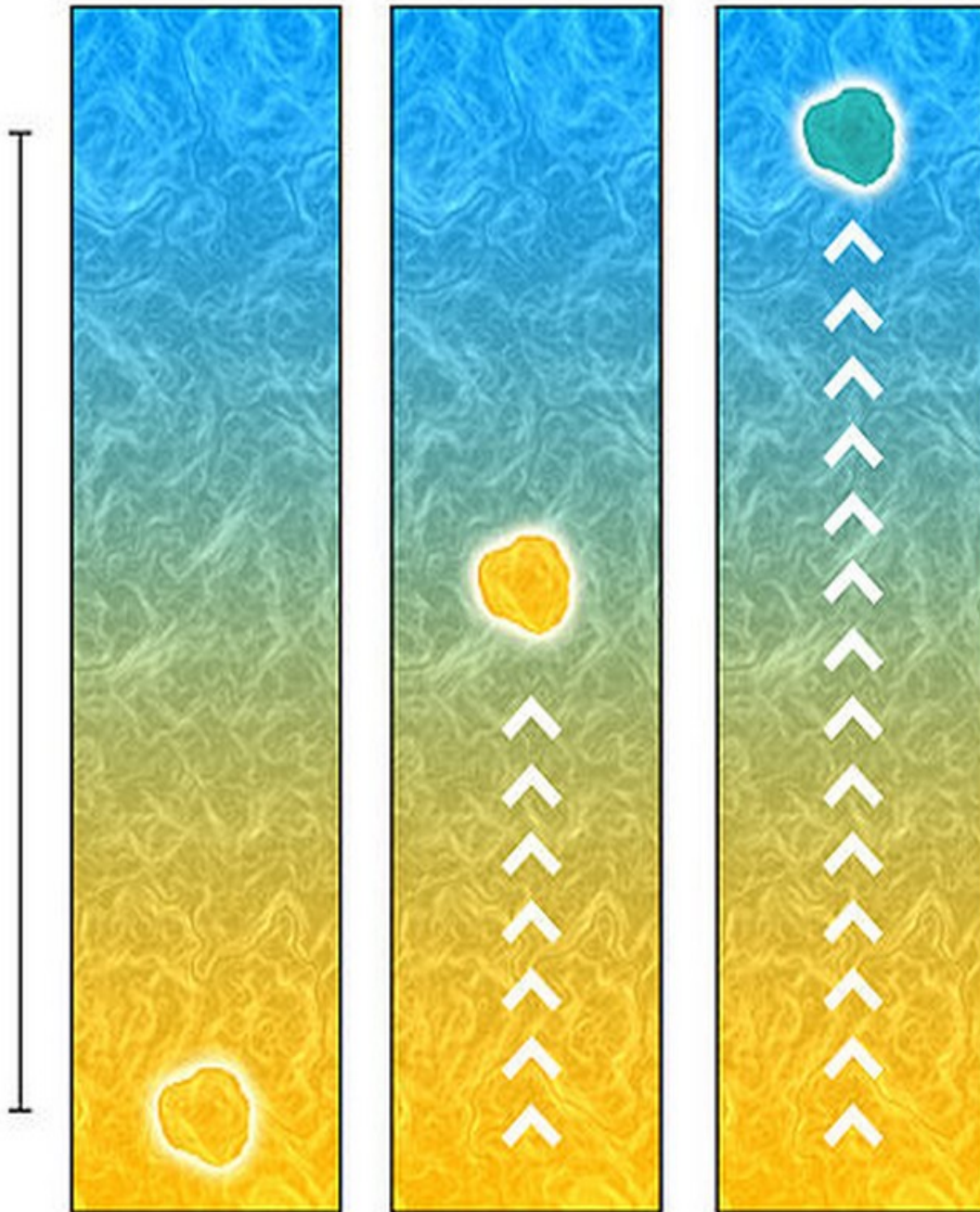
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 $H_p = d \ln(P) / d \ln(T)$

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$$F_{\text{conv}} = \frac{1}{2} \rho v c_p T \frac{\lambda}{H_P} (\nabla_T - \nabla_{\text{ad}}).$$

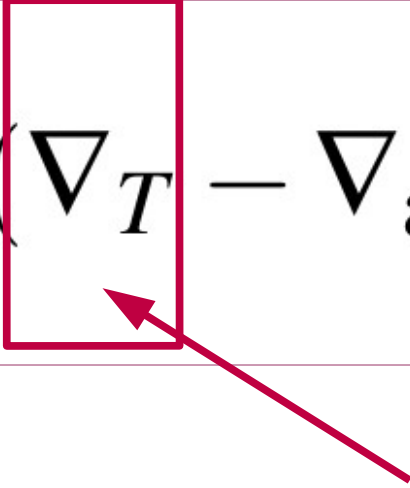
Let's break this down...

$$F_{\text{conv}} = \frac{1}{2} \rho v c_p T \frac{\lambda}{H_P} (\nabla_T - \nabla_{\text{ad}}).$$

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Reminder: Stellar evolution definitions:

$$\nabla_T = \left(\frac{d \ln T}{d \ln P} \right)$$

General temperature gradient

$$\nabla_{\text{rad}} = \left(\frac{d \ln T}{d \ln P} \right)_{\text{rad}} = \frac{3}{16\pi a c G} \frac{\kappa l P}{m T^4}$$

Radiative temperature gradient

$$\nabla_{\text{ad}} = \left(\frac{d \ln T}{d \ln P} \right)_{\text{ad}} = \frac{P \delta}{T \rho c_p}$$

Adiabatic temperature gradient

Let's break this down...

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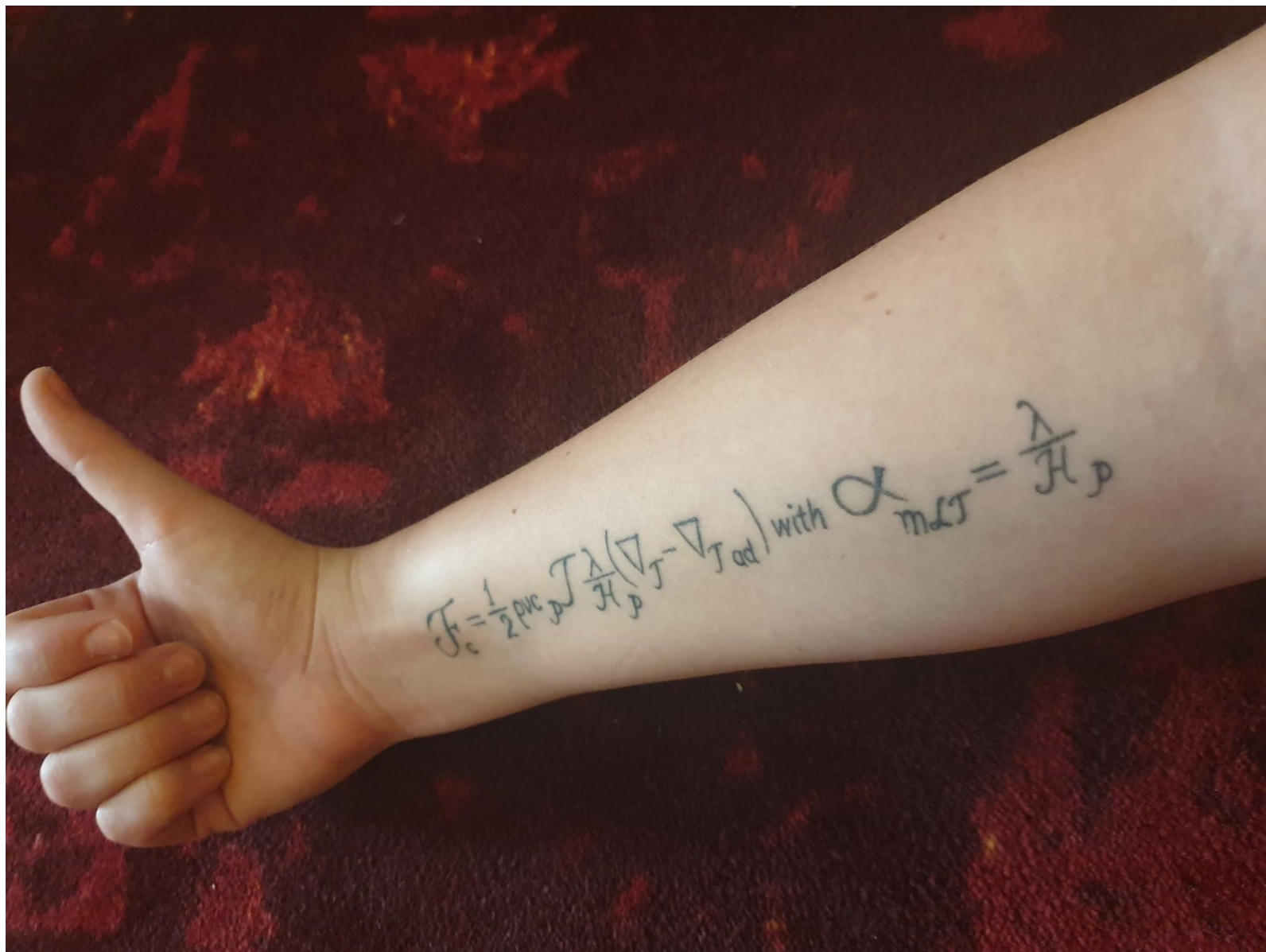
~ **Local flux of convective energy** (local pressure, specific heat, temperature, velocity)

Let's break this down...

$$F_{\text{conv}} = \frac{1}{2} \rho v c_p T \frac{\lambda}{H_P} (\nabla_T - \nabla_{\text{ad}}).$$

$$\alpha_{\text{MLT}} = \frac{\lambda}{H_P}$$

Mean free path λ in terms of the pressure scale height



$$F_c = \frac{1}{2} \rho v c_p \int \frac{\lambda}{H_p} (\nabla_T - \nabla_{T_{ad}}) \text{ with } \alpha_{mL_T} = \frac{\lambda}{H_p}$$

Quick review: convective stability criteria

- The Ledoux criterion for dynamical stability is given by

$$\nabla_{\text{rad}} < \nabla_{\text{ad}} + [\varphi/\delta] \nabla_{\mu}$$

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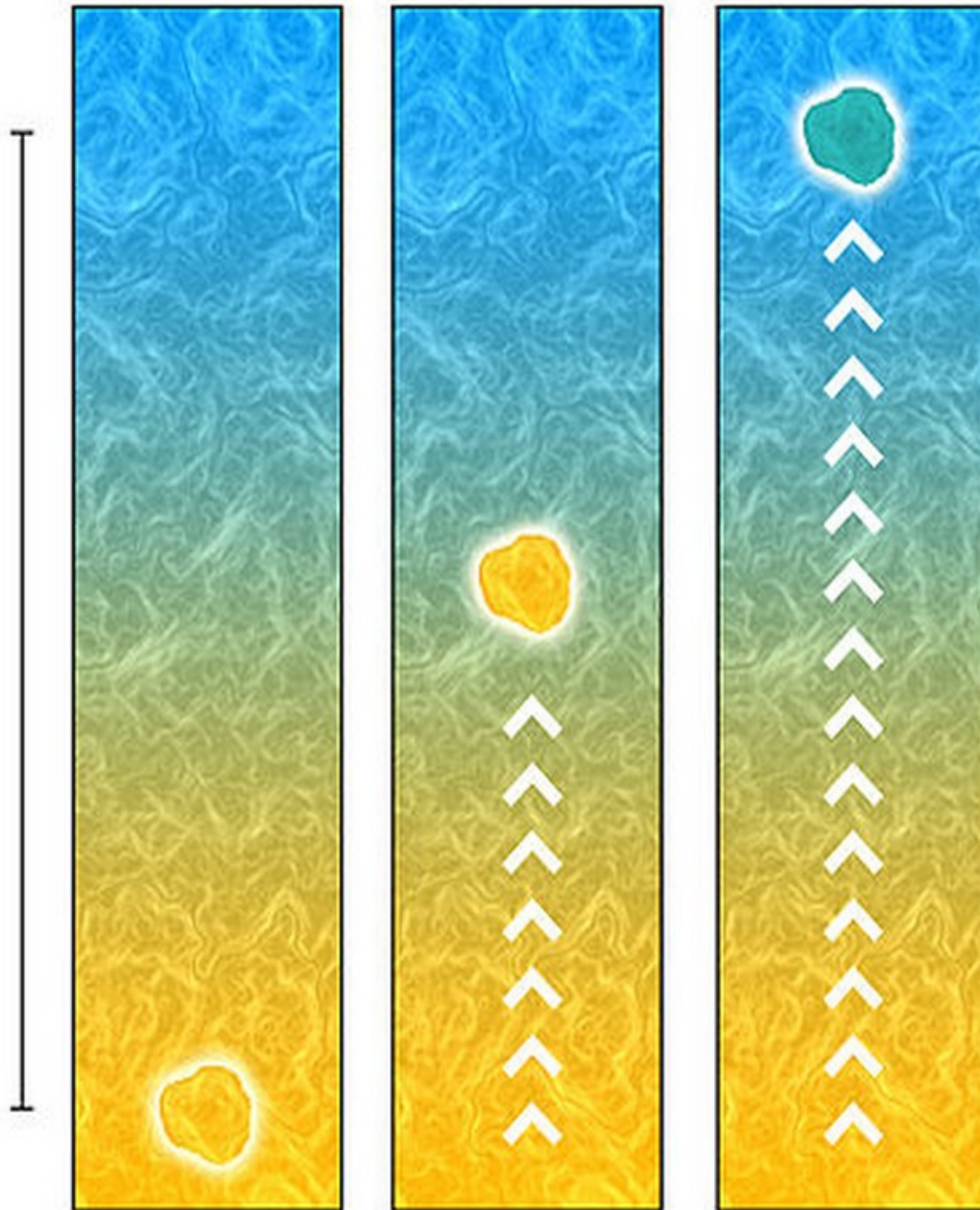
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- Dynamically stable regions do not produce convective motions, and the energy flux is carried out exclusively by radiation (or conduction).

- If the relevant condition is violated, convection will activate and share in the transport of flux

Summary: Mixing Length Theory



$$F_{\text{conv}} = \frac{1}{2} \rho v c_p T \frac{\lambda}{H_P} (\nabla_T - \nabla_{\text{ad}}).$$

$$\alpha_{\text{MLT}} = \frac{\lambda}{H_P} \quad \nabla_T = \left(\frac{d \ln T}{d \ln P} \right)$$

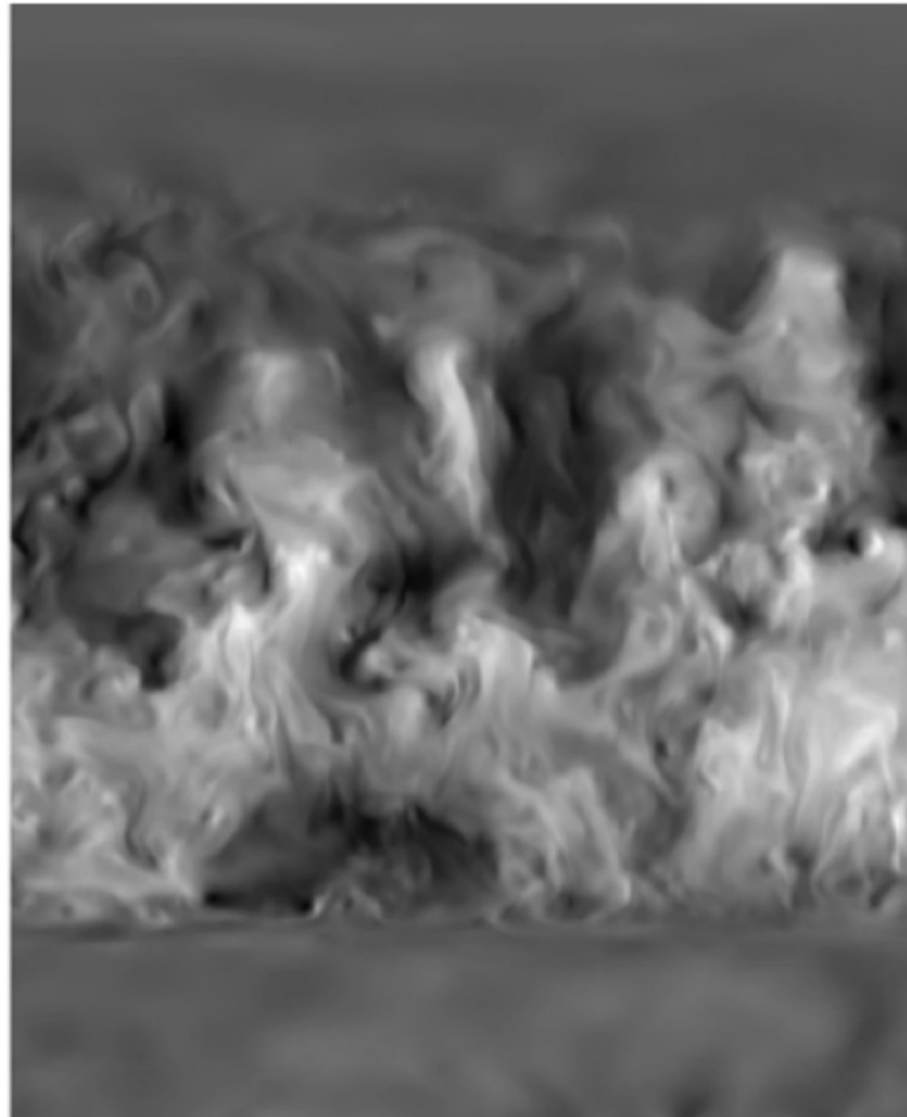
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How is MLT inadequate?

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Let us count the ways...

Nonconvective →



← Boundary layer

Body of
Convective
Region

← Boundary layer

Nonconvective →

Figure: Maxime Viallet. MUSIC

(1) Convective boundaries are not rigid; local mixing occurs. MLT cannot account for this

(2) There is no symmetry between upflows and downflows, which creates **negative kinetic flux**.
MLT cannot account for this.

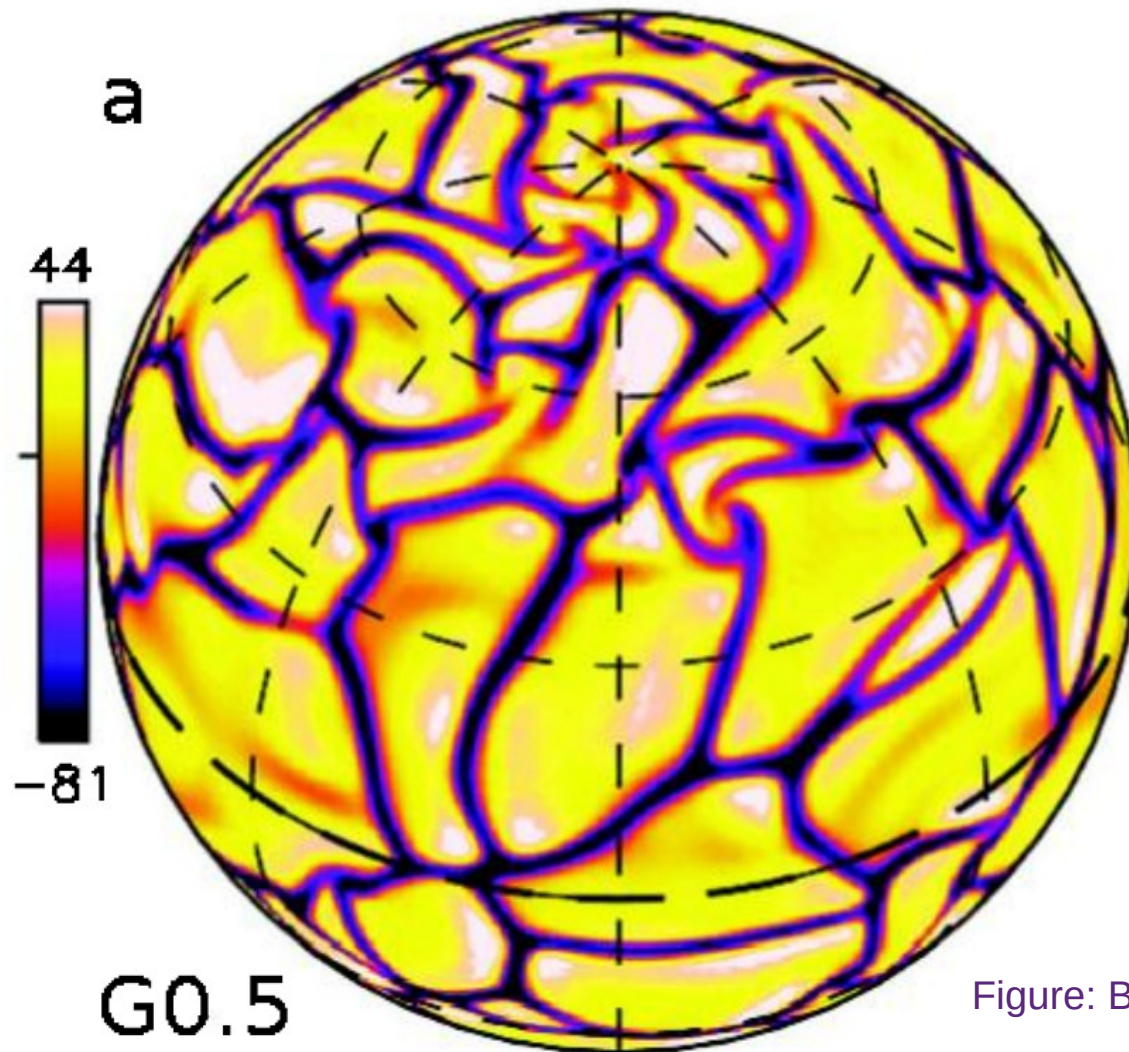
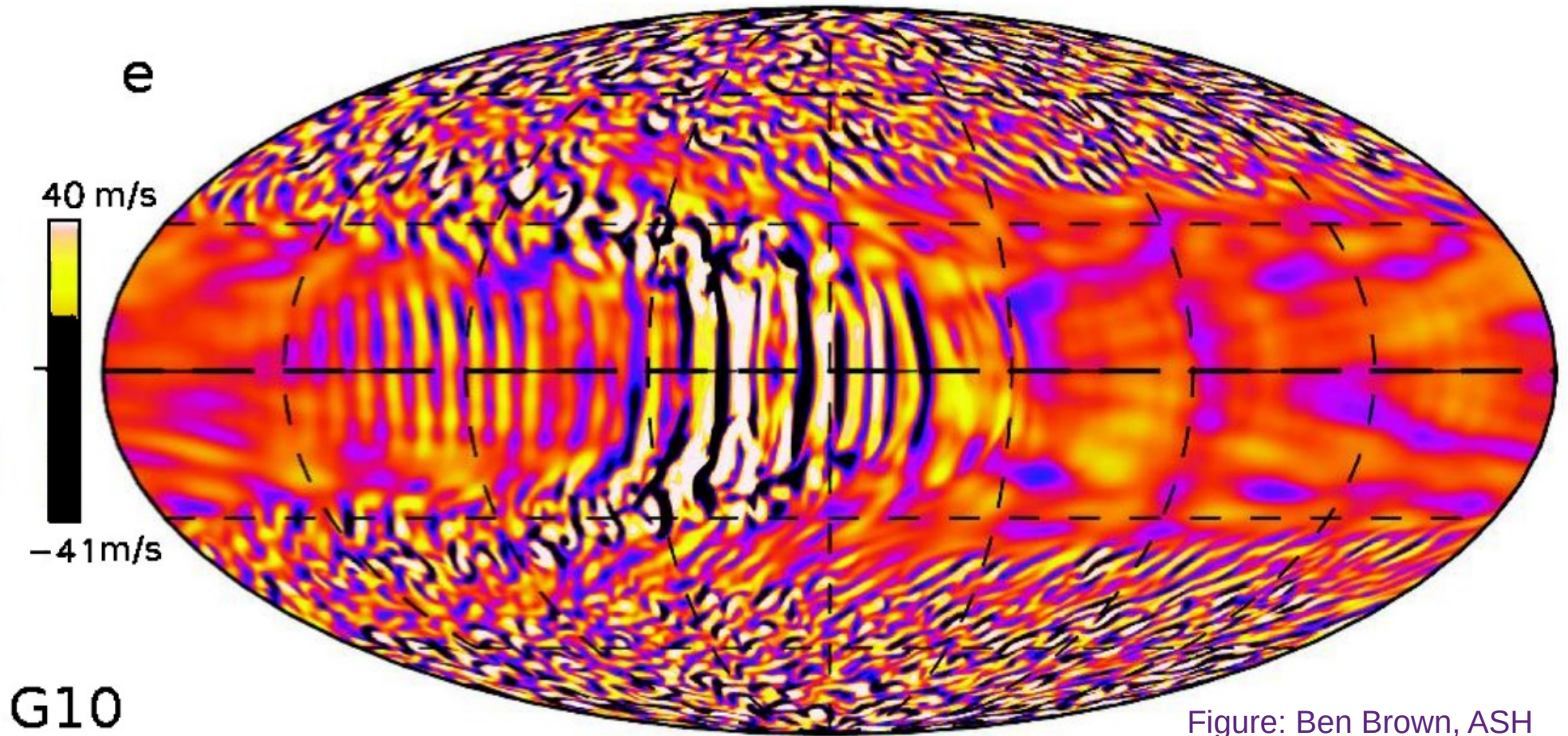


Figure: Ben Brown, ASH

(3) Real convective cells undergo shearing and fragmentation. MLT cannot account for this.

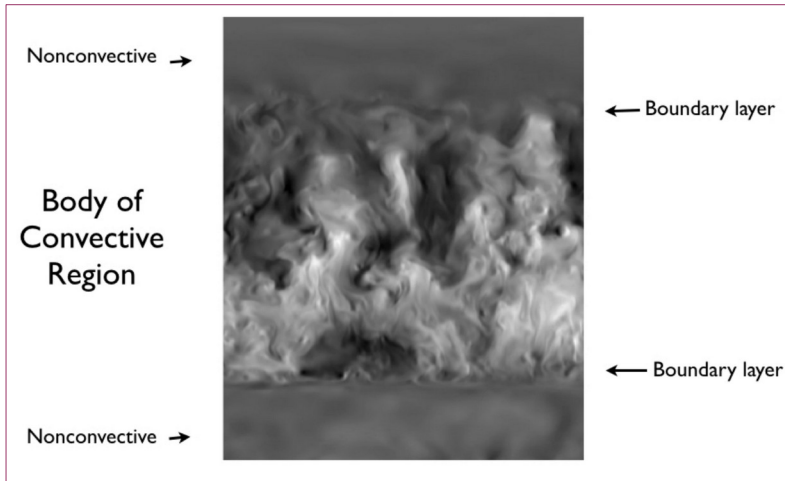


Recap: Shortcomings of MLT*

(*which means you cannot complain about these later in the talk :))

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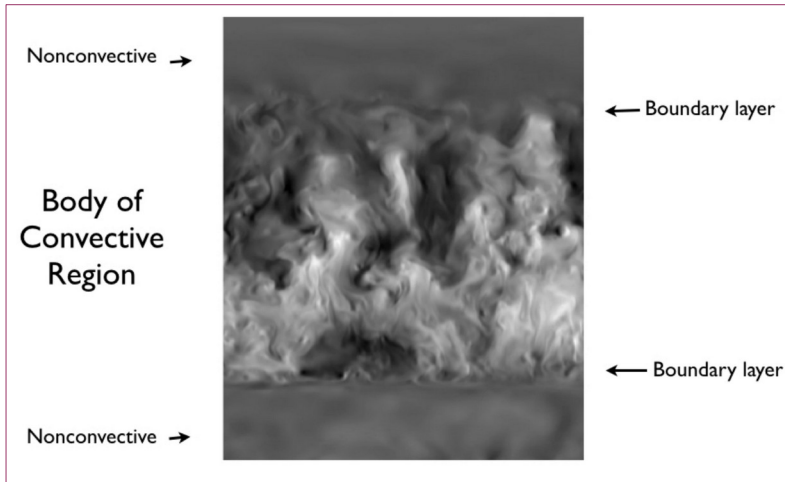
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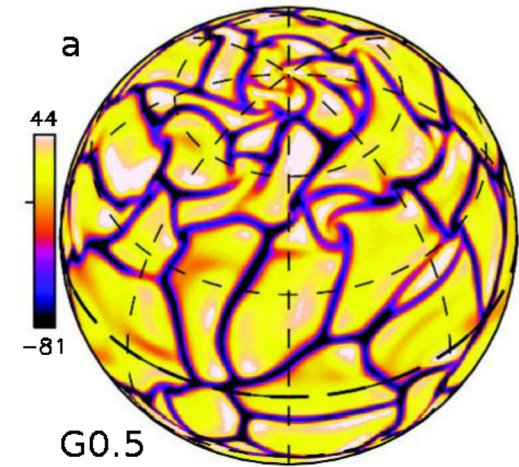
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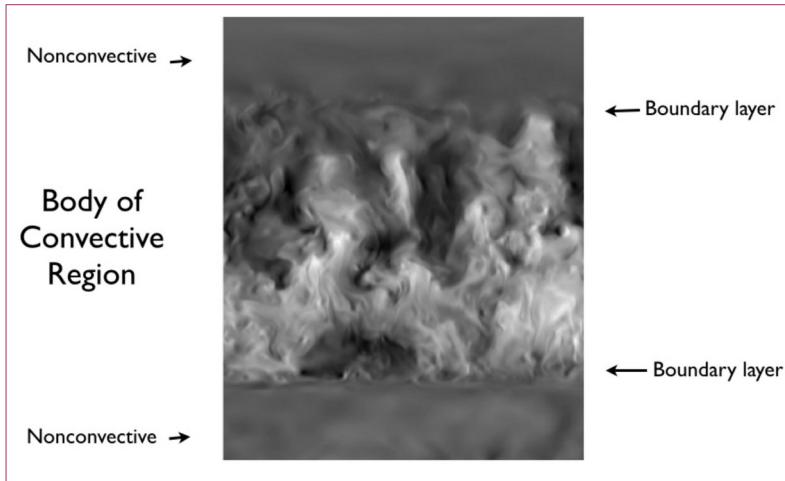
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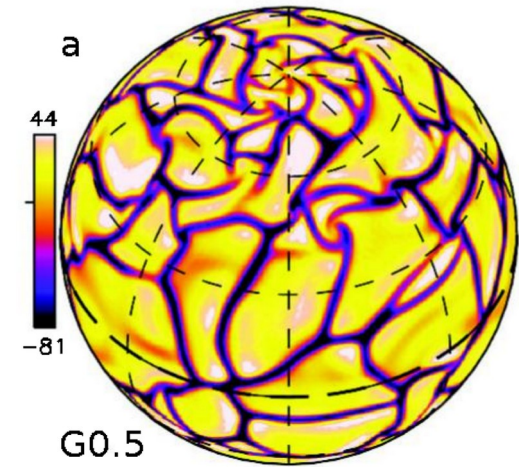
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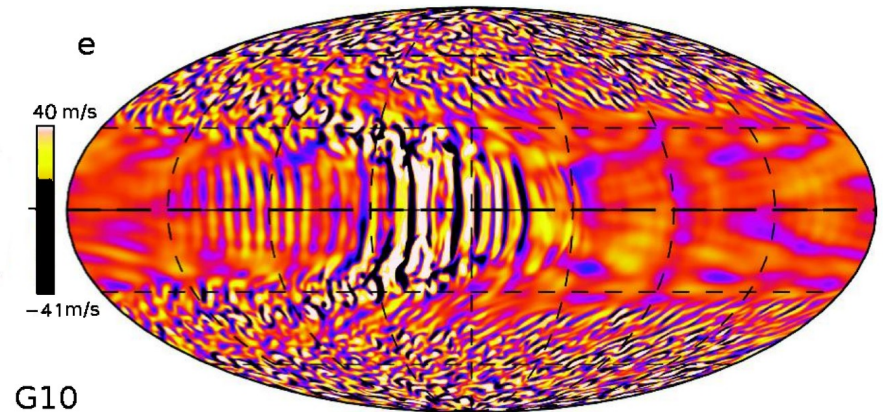
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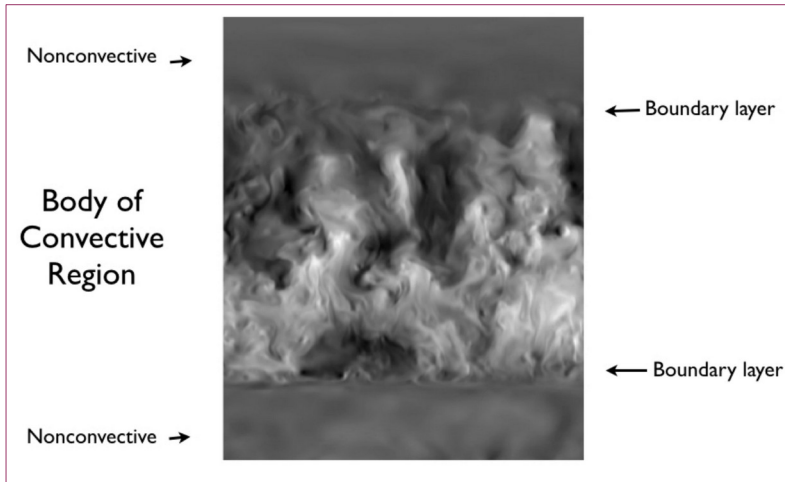
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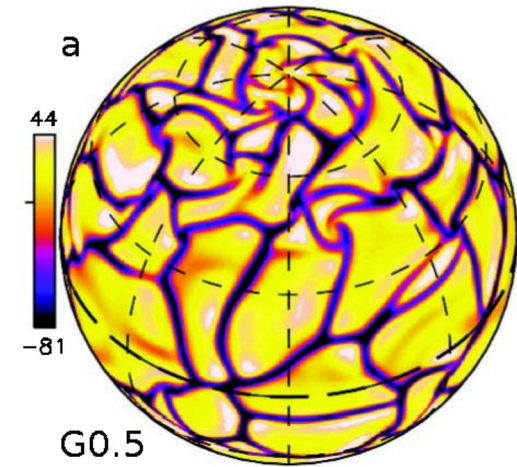
(3) No concept of shearing, fragmentation, nor any 2D or 3D effects

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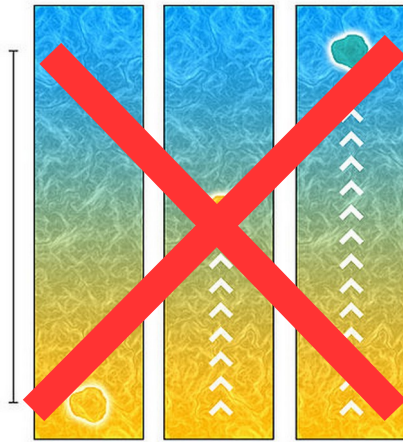
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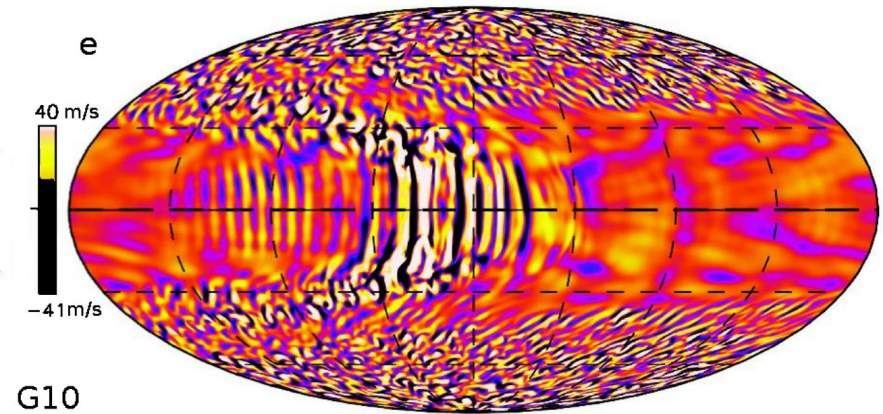
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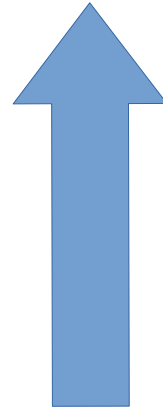
(4) BONUS: there is no such thing as a quantized unit of convection, so what does a "parcel" even mean?



(3) No concept of shearing, fragmentation, nor any 2D or 3D effects

Nevertheless, we persist...

Present

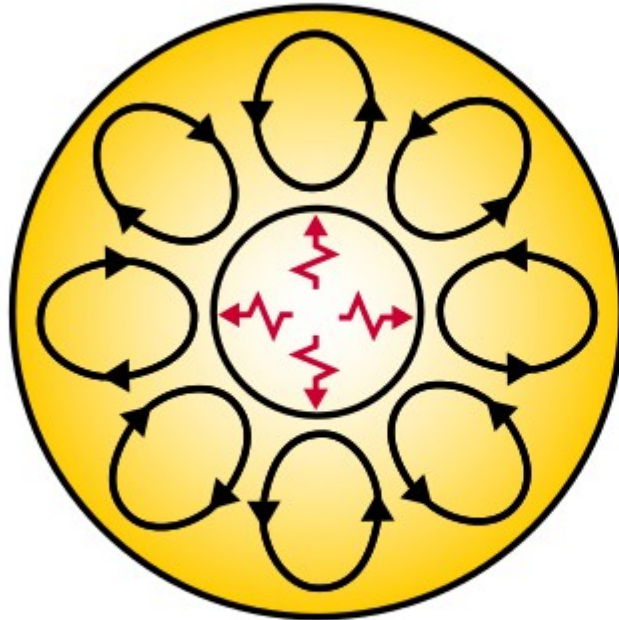


Let's talk about 1D “stars”

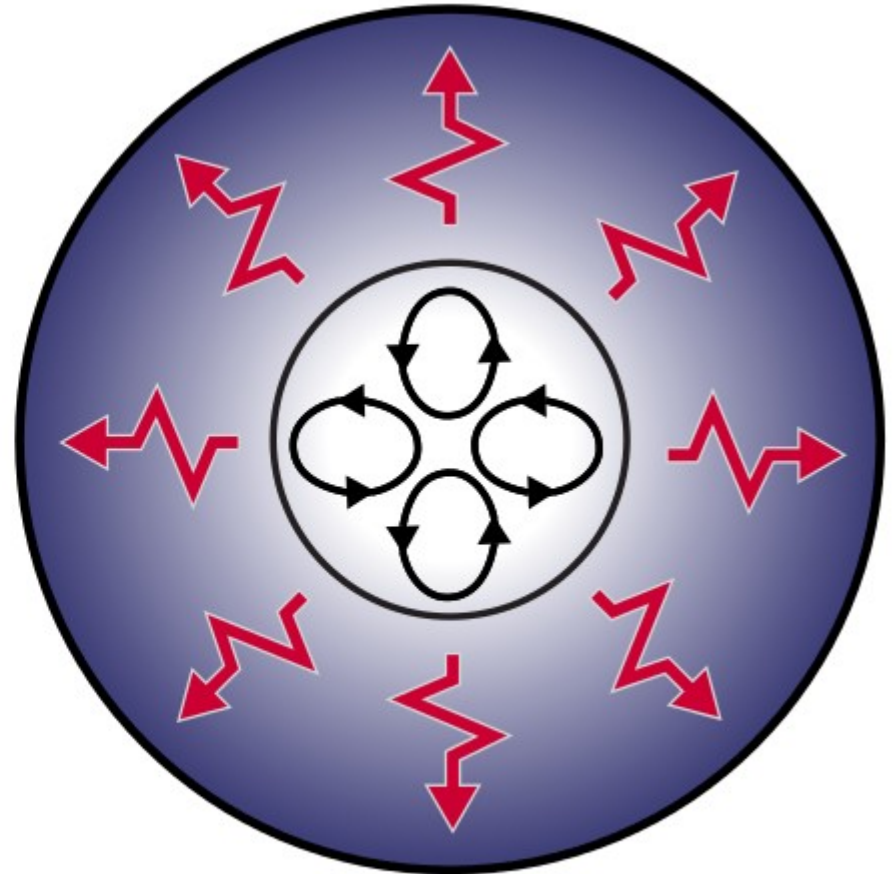
Where does MLT matter?



$M < 0.5$



$0.5 - 1.5$



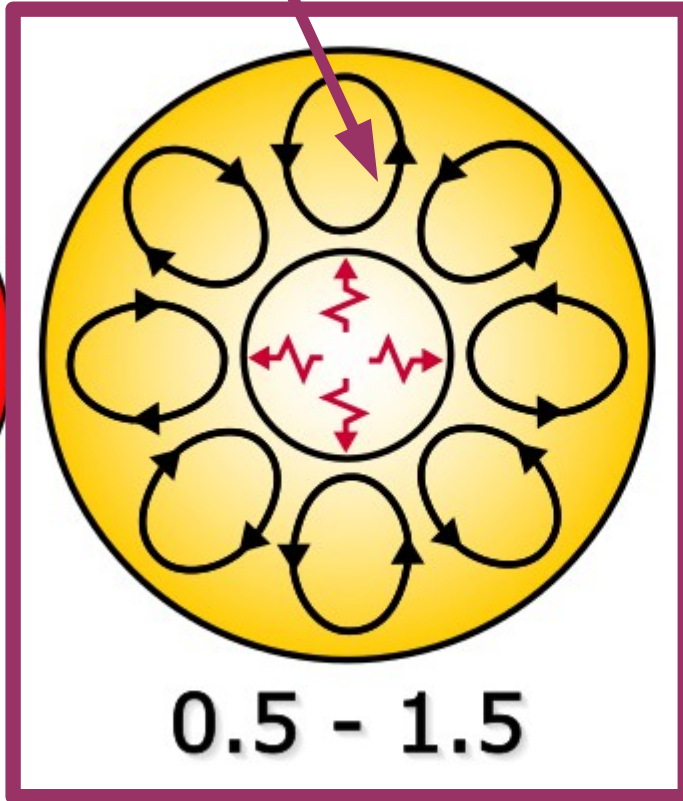
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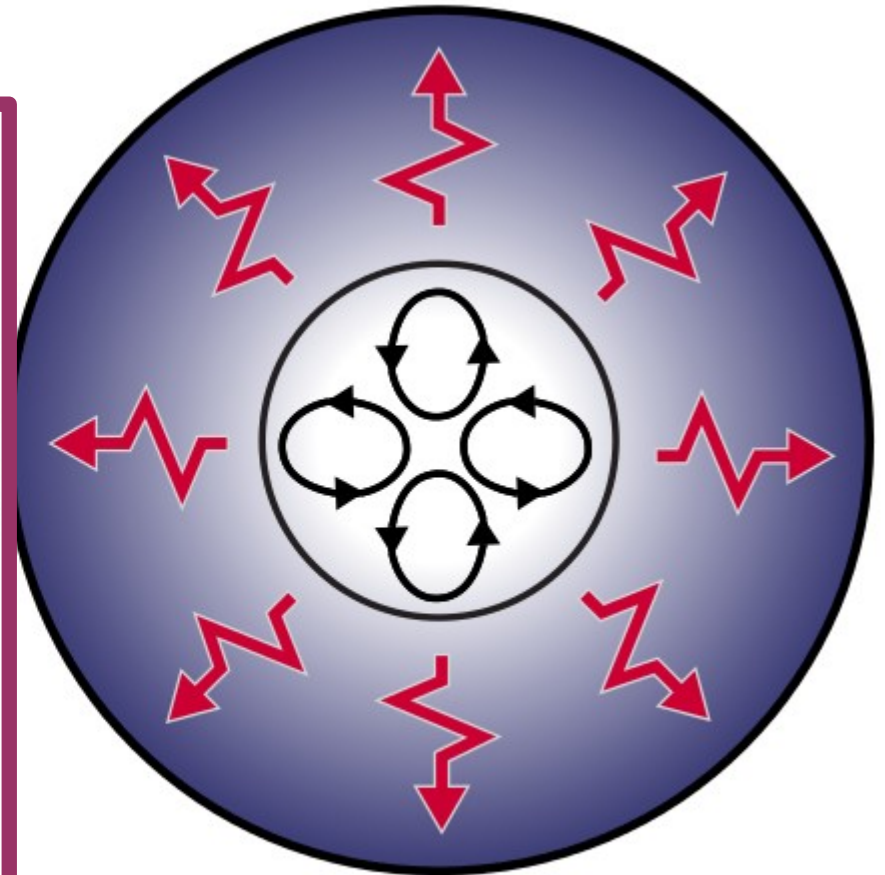
here



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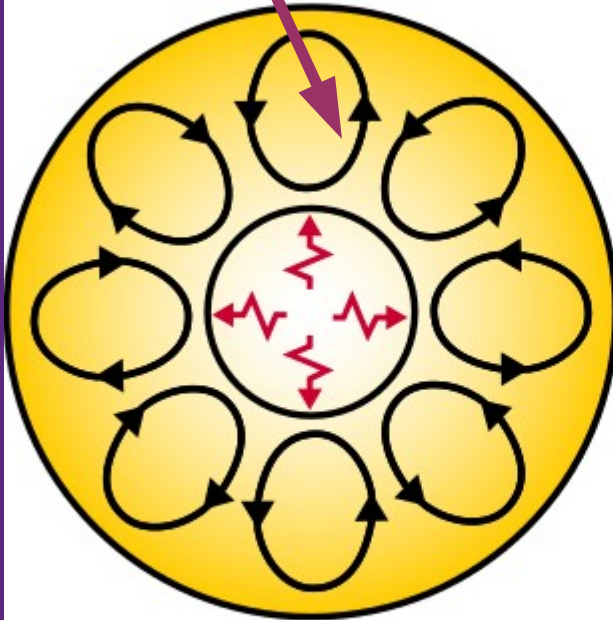
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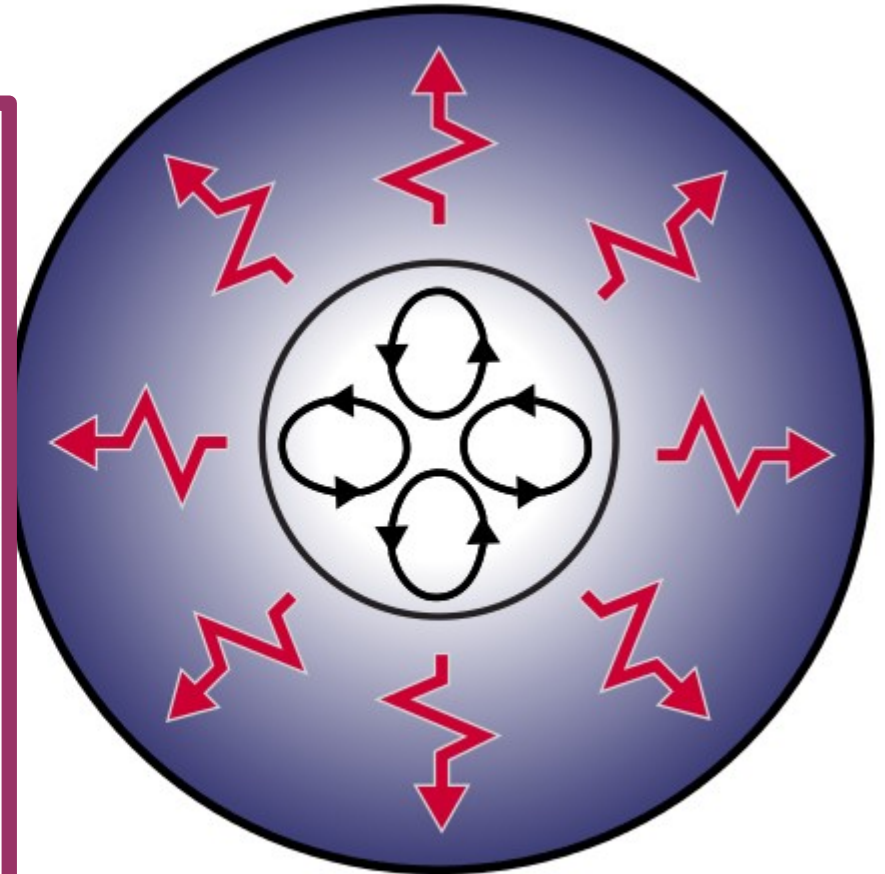
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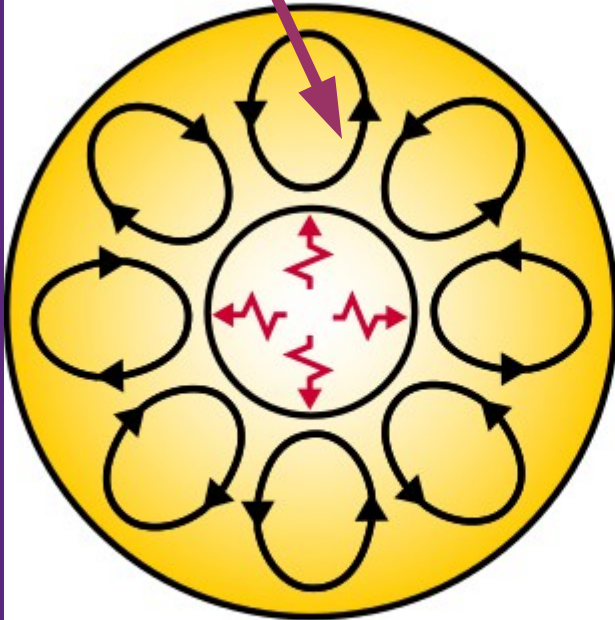
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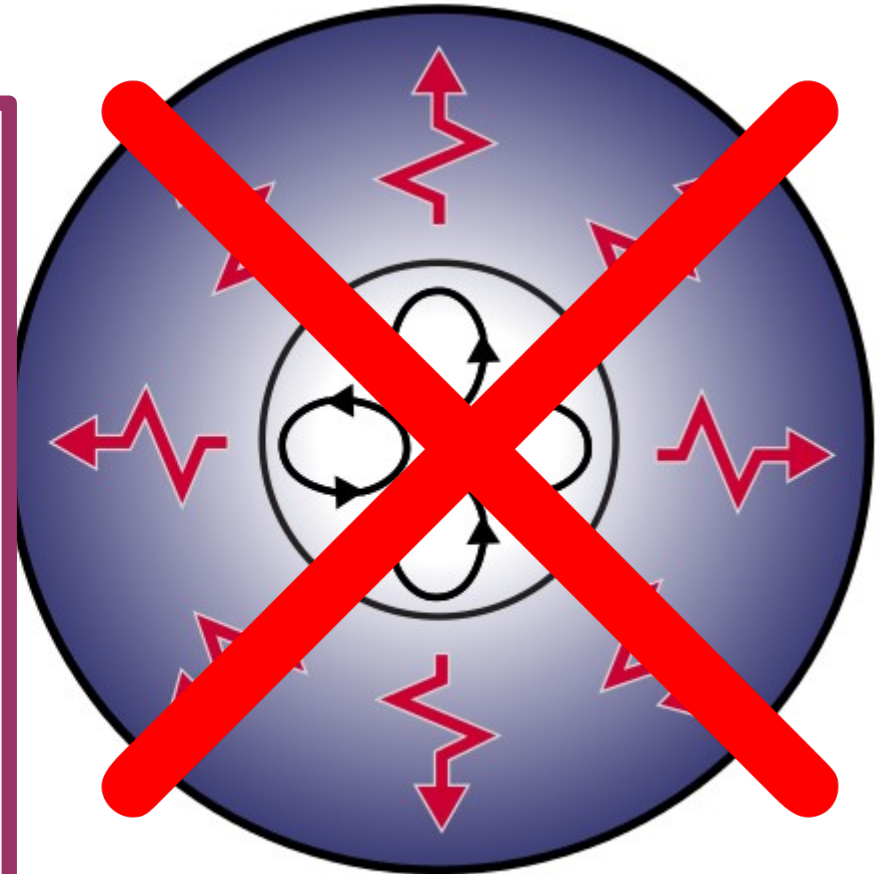
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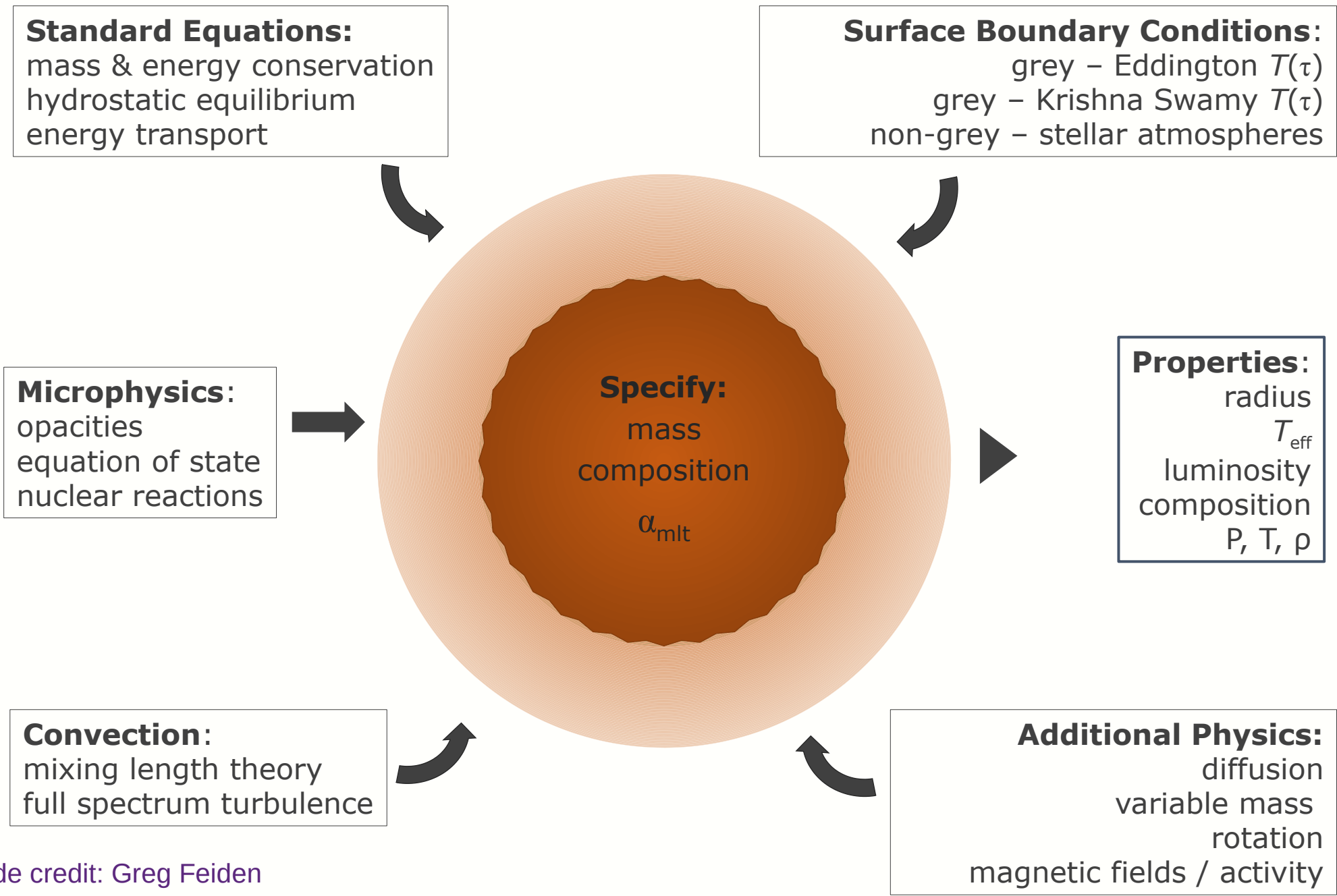
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- the entropy jump implies the depth of the convective envelope

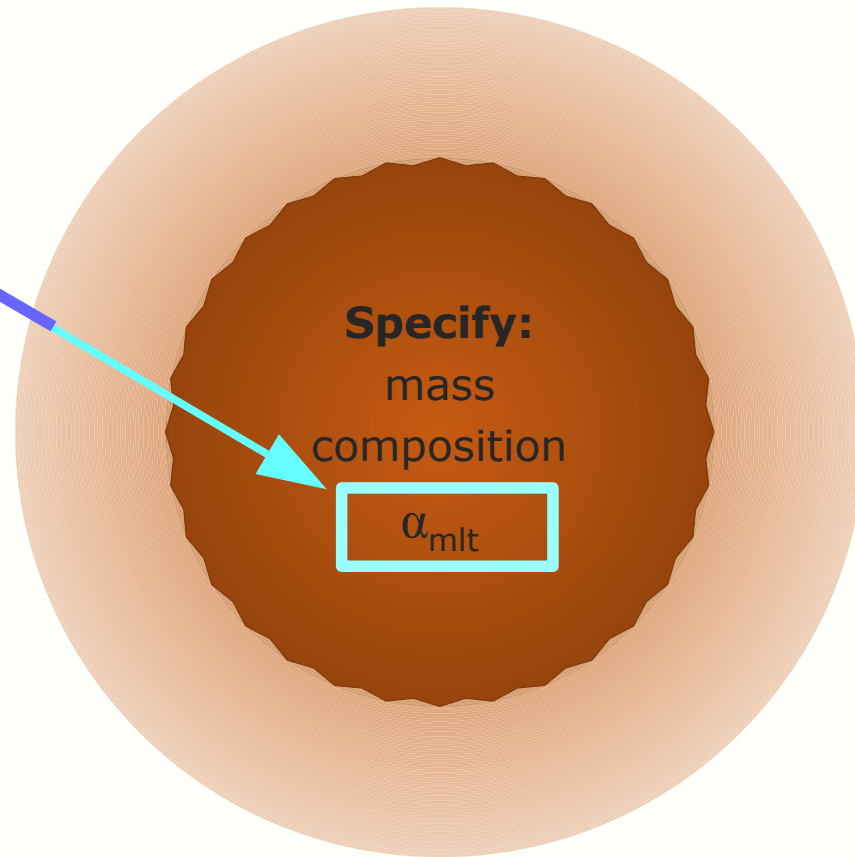
Components of a Stellar Structure & Evolution Program



Slide credit: Greg Feiden

Components of a Stellar Structure & Evolution Program

Need this to
be correct



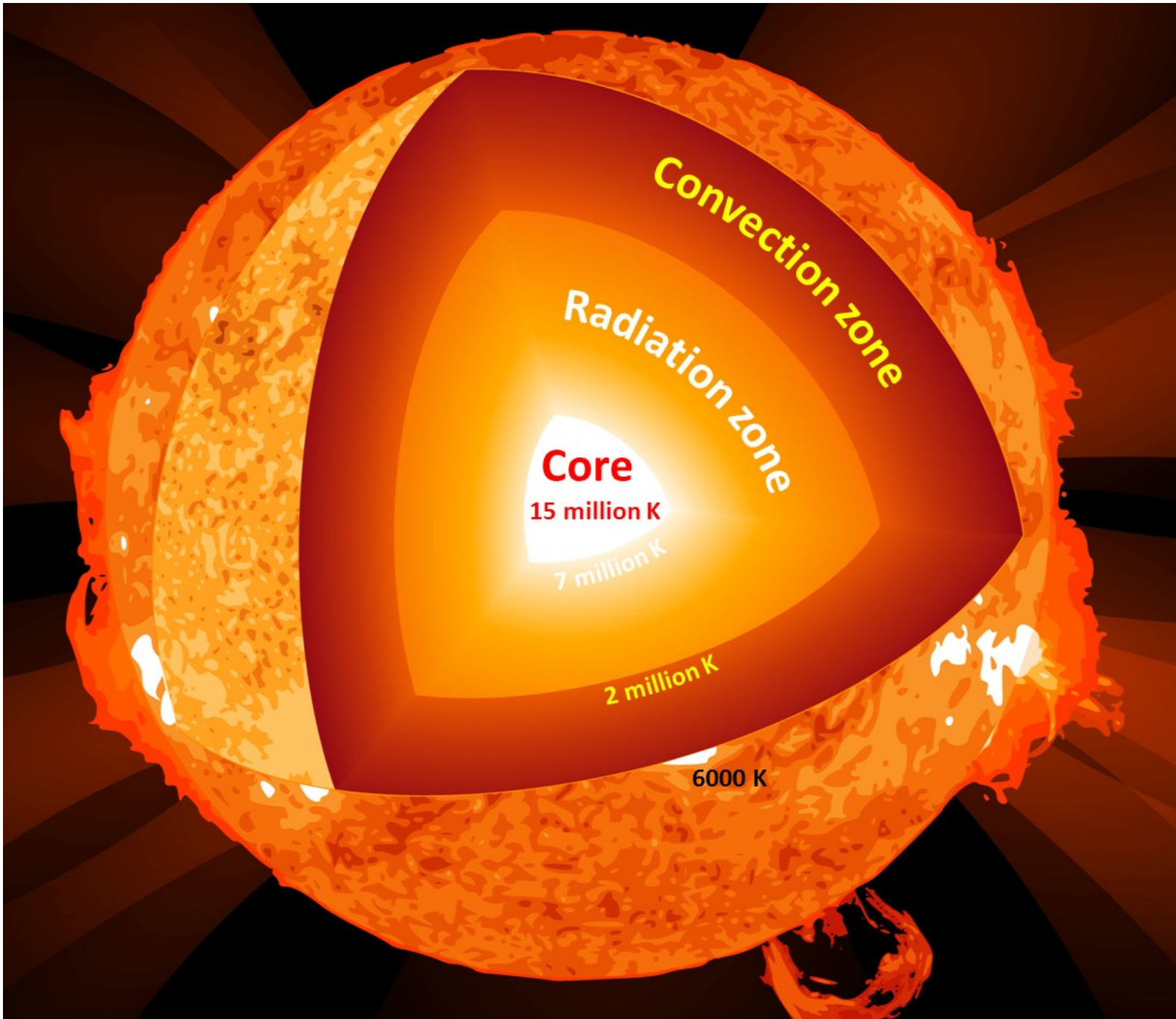
Specify:
mass
composition

α_{mlt}

Properties:
radius
 T_{eff}
luminosity
composition
 P, T, ρ

...for these to be correct

MLT in practice: The solar calibration



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—so yes, of course it is wrong...but is it still good enough?

A path forward: calibrating the mixing length to stars other than the Sun

Two separate science questions:

- (1) How does α_{MLT} vary among stars with different global properties?
- (2) How does α_{MLT} change within a single star's evolution?

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- > Tayar et al., 2017
- > Joyce & Chaboyer, 2018a, 2018b
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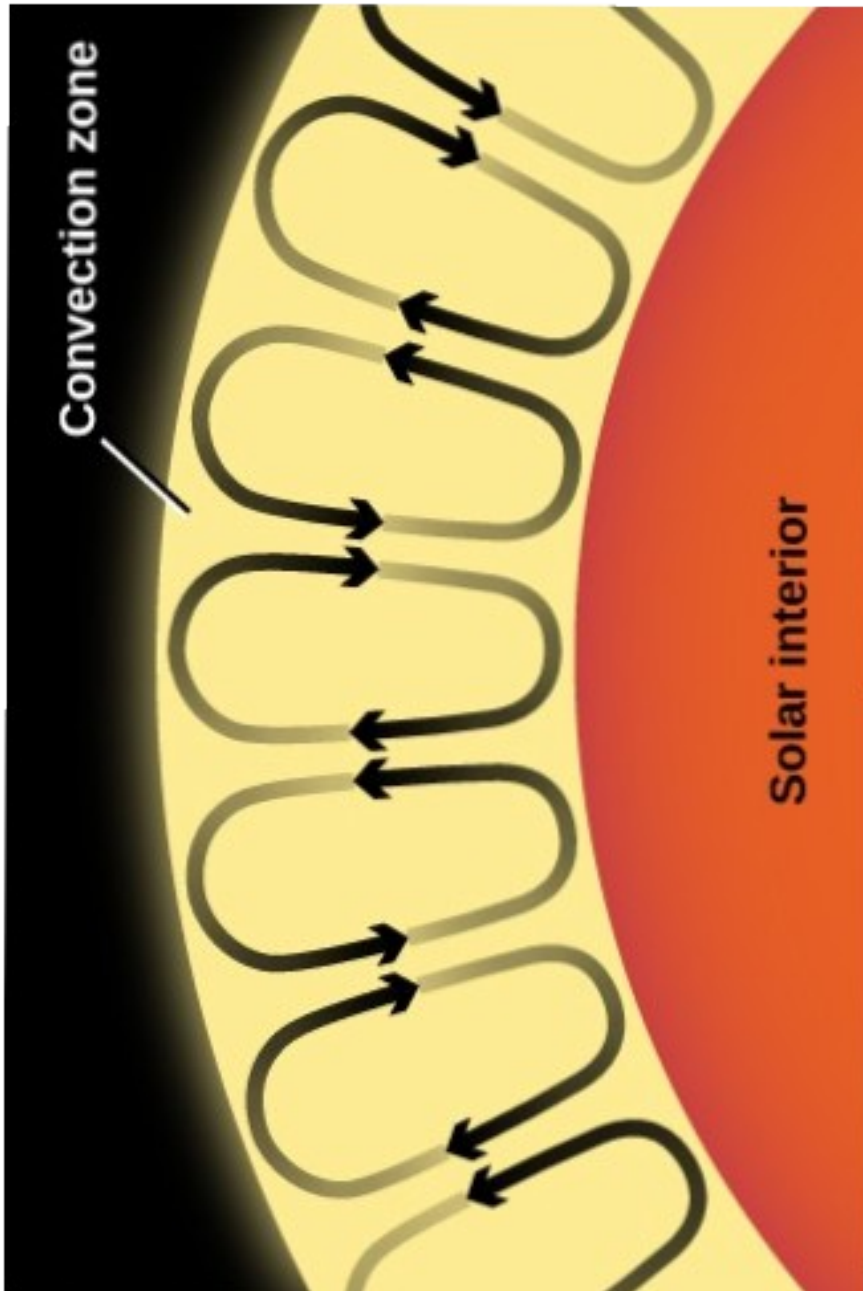
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(2) How should α_{MLT} change within a single star's evolution?

- > Trampedach et al., 2014
- > Spada et al., 2021

What we need:



Calibrate here!

- ♦ low mass stars (0.5 – 1.2 M_{\odot})
- ♦ Stars with sub-surface convective envelopes (preferably thick)
- ♦ main sequence, subgiant, or early RGB stars
- ♦ Star should ideally have at most one critical deviation from the Sun at a time
- ♦ As many independent measurements of the candidate's fundamental properties as possible

alpha Centauri A & B- the perfect lab

The high number of independent observational measurements removes degrees of freedom and **allows the mixing length parameter to be isolated** in the stellar evolution models

Mass – kinematics

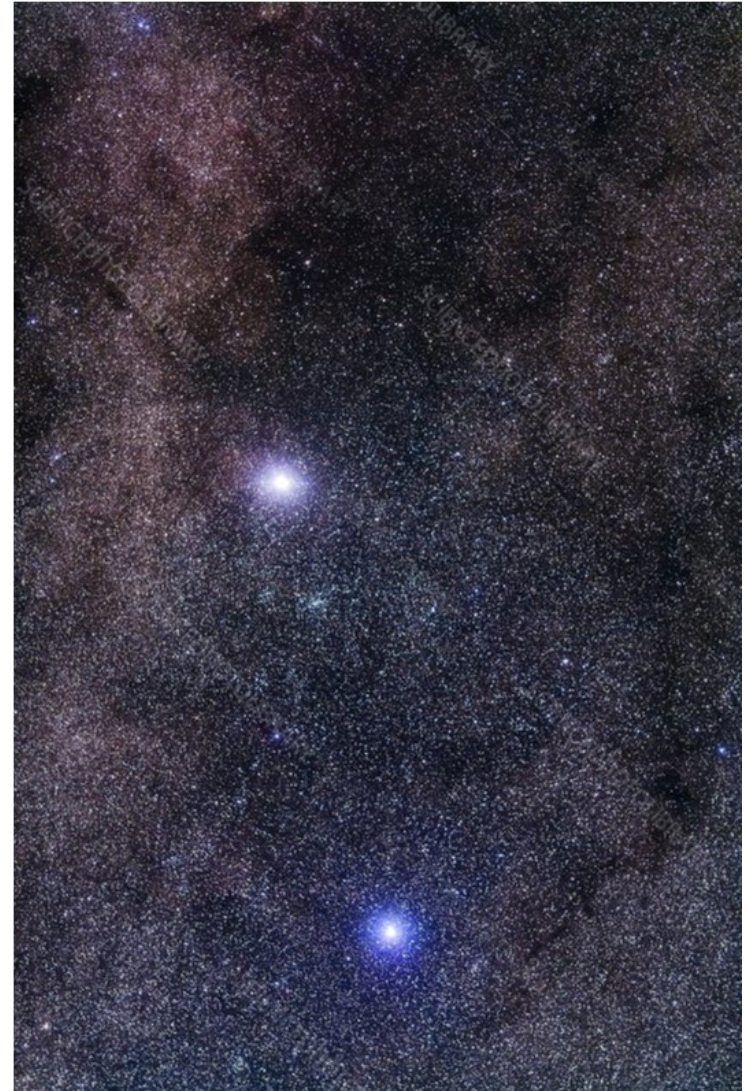
Radius – interferometry

Luminosity – photometry

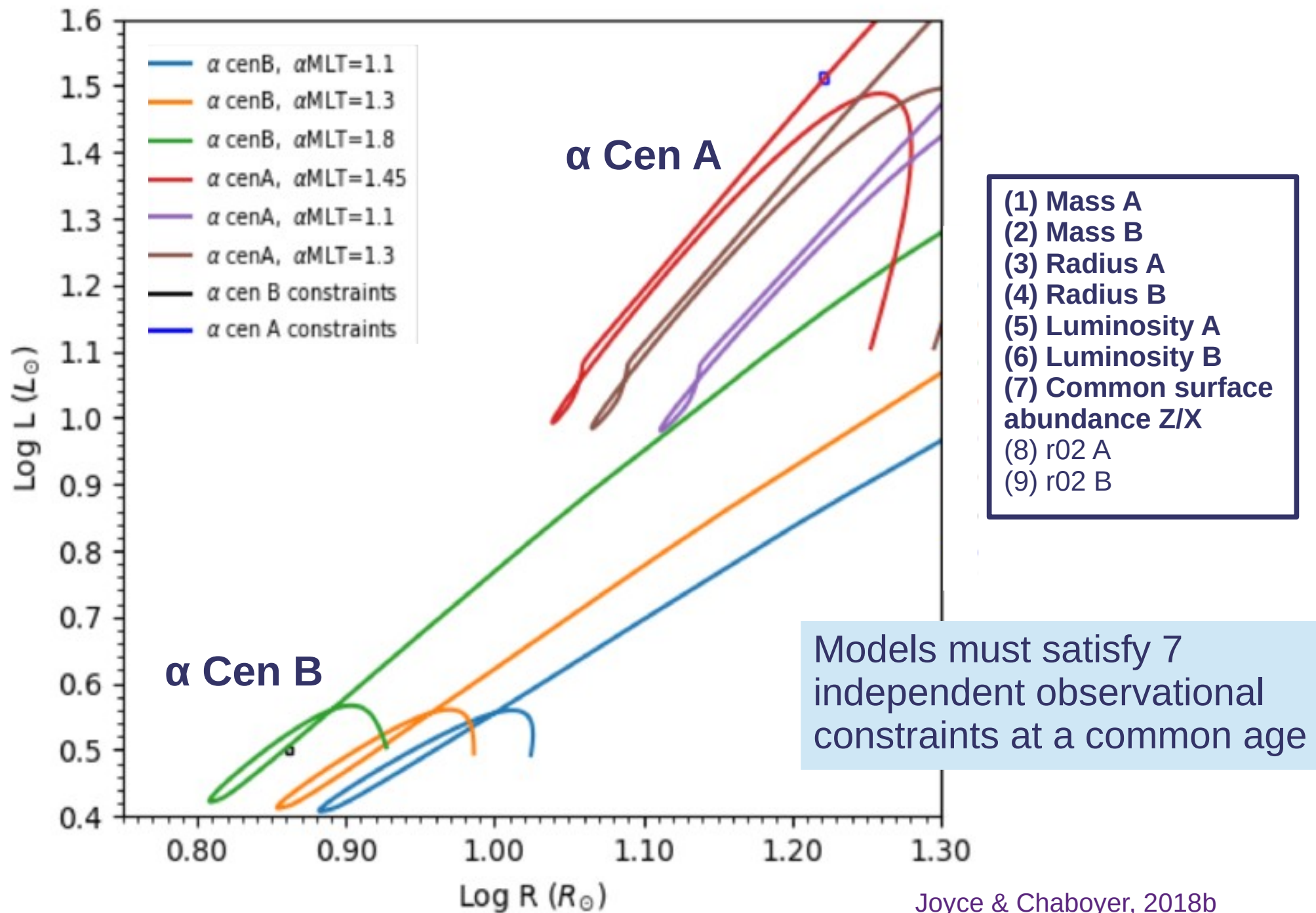
Surface abundance – high resolution spectroscopy

Stellar interior constraints **from which surface effects can be removed** – seismology

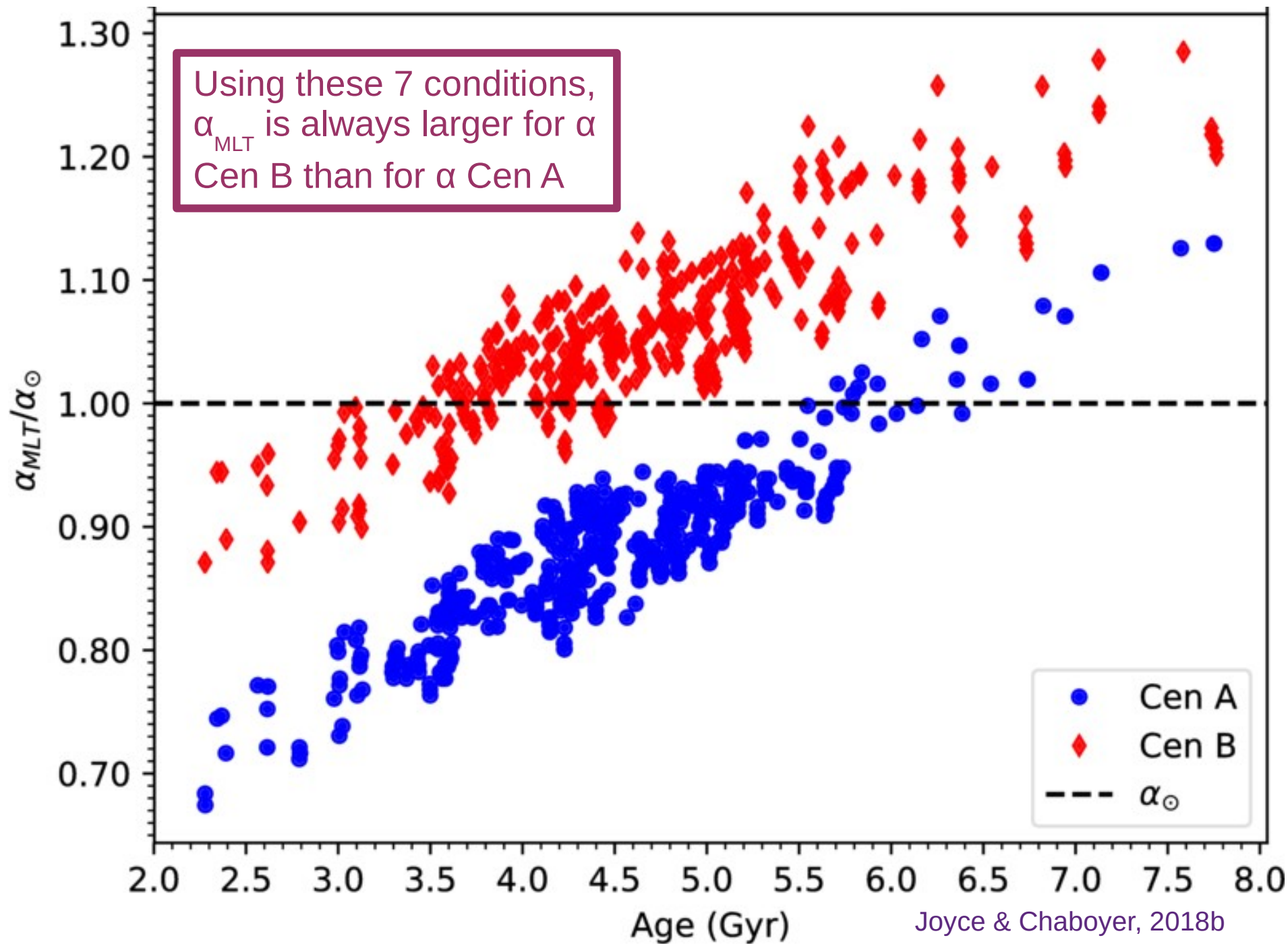
Because the candidate is a binary with all classical measurements satisfied in both components, we get a prior-independent age constraint for free



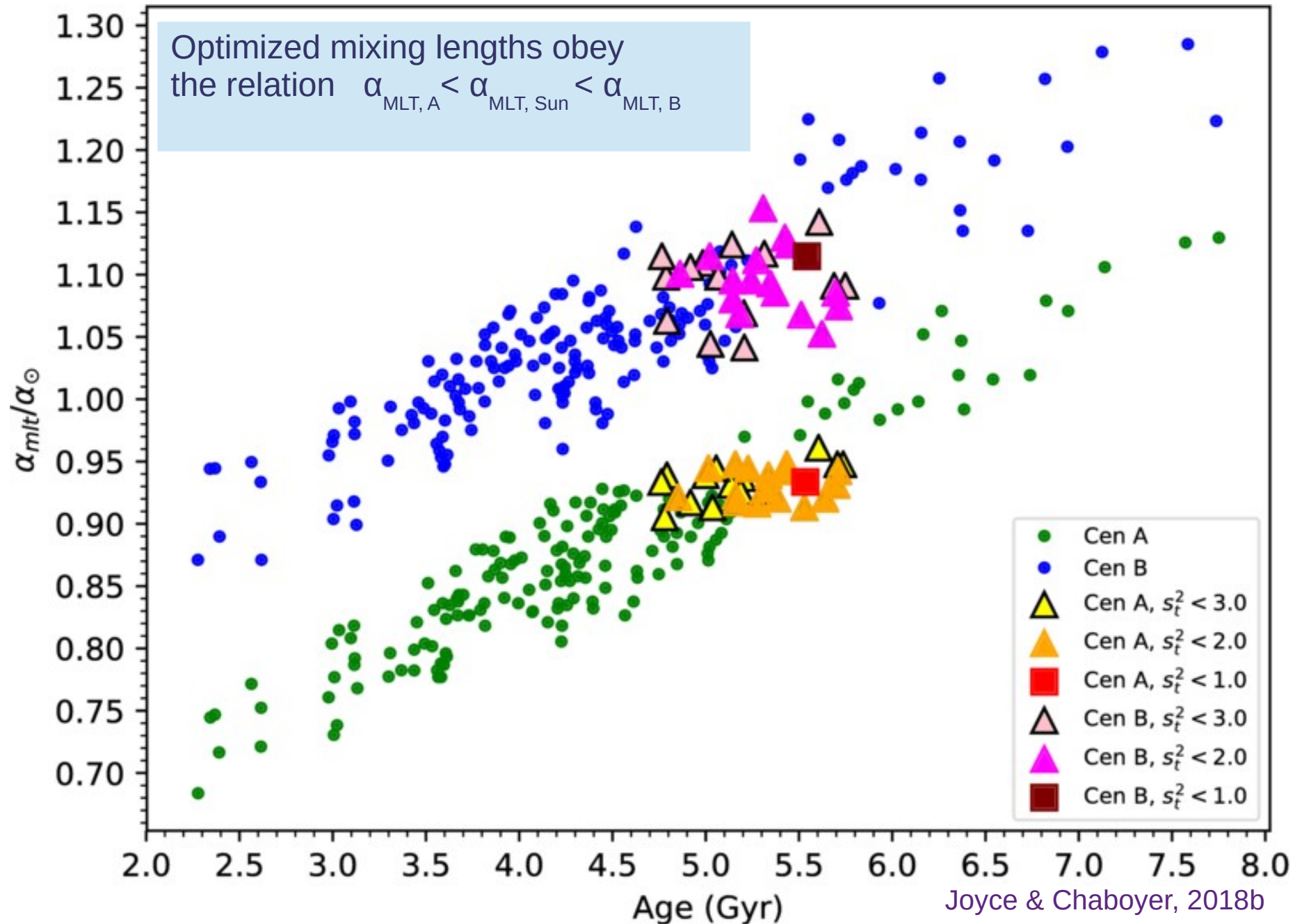
Mixing length calibrations to alpha Centauri A & B



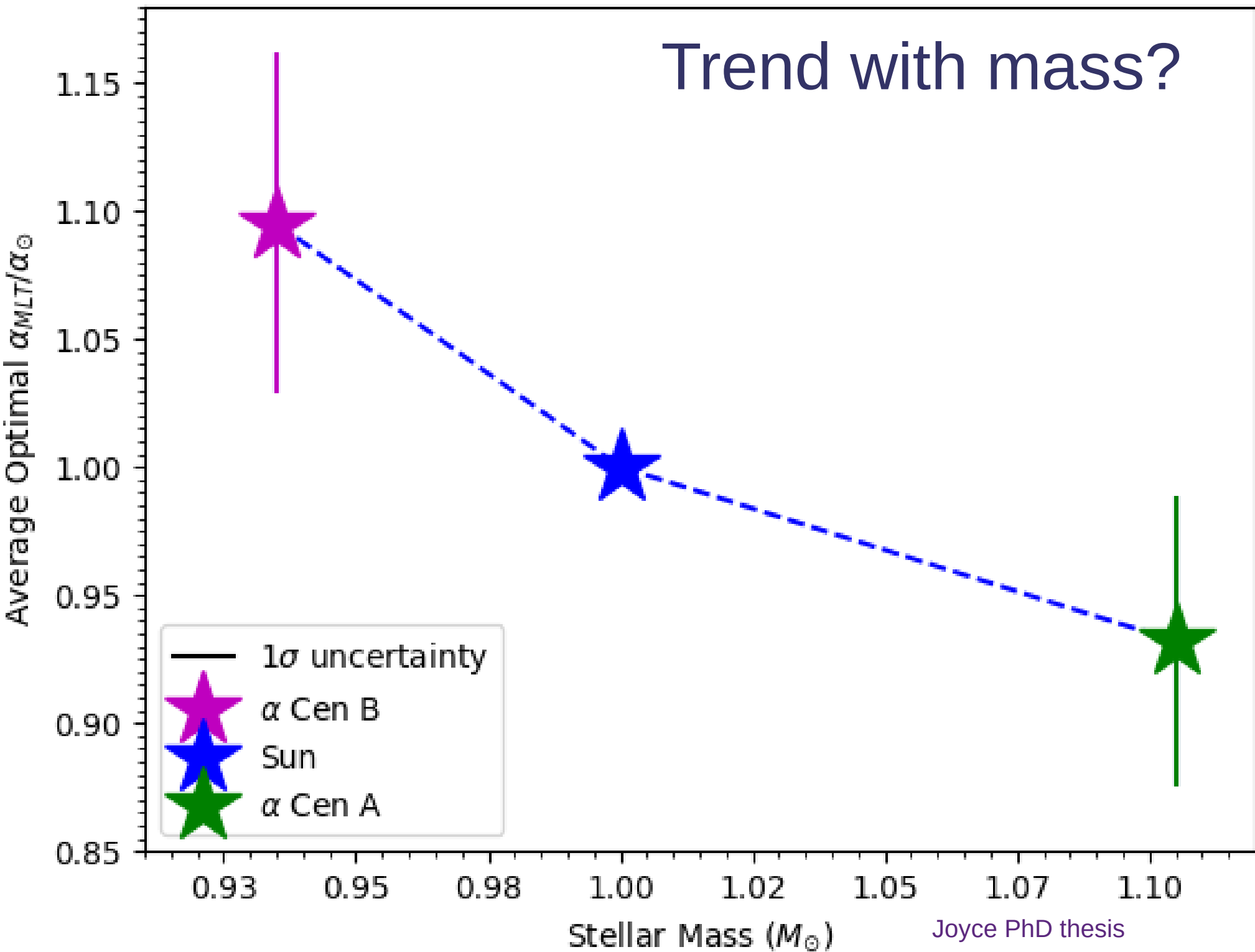
Mixing length calibrations to alpha Centauri A & B



When seismic constraints are added to the original 7 constraints, the models converge to precise values

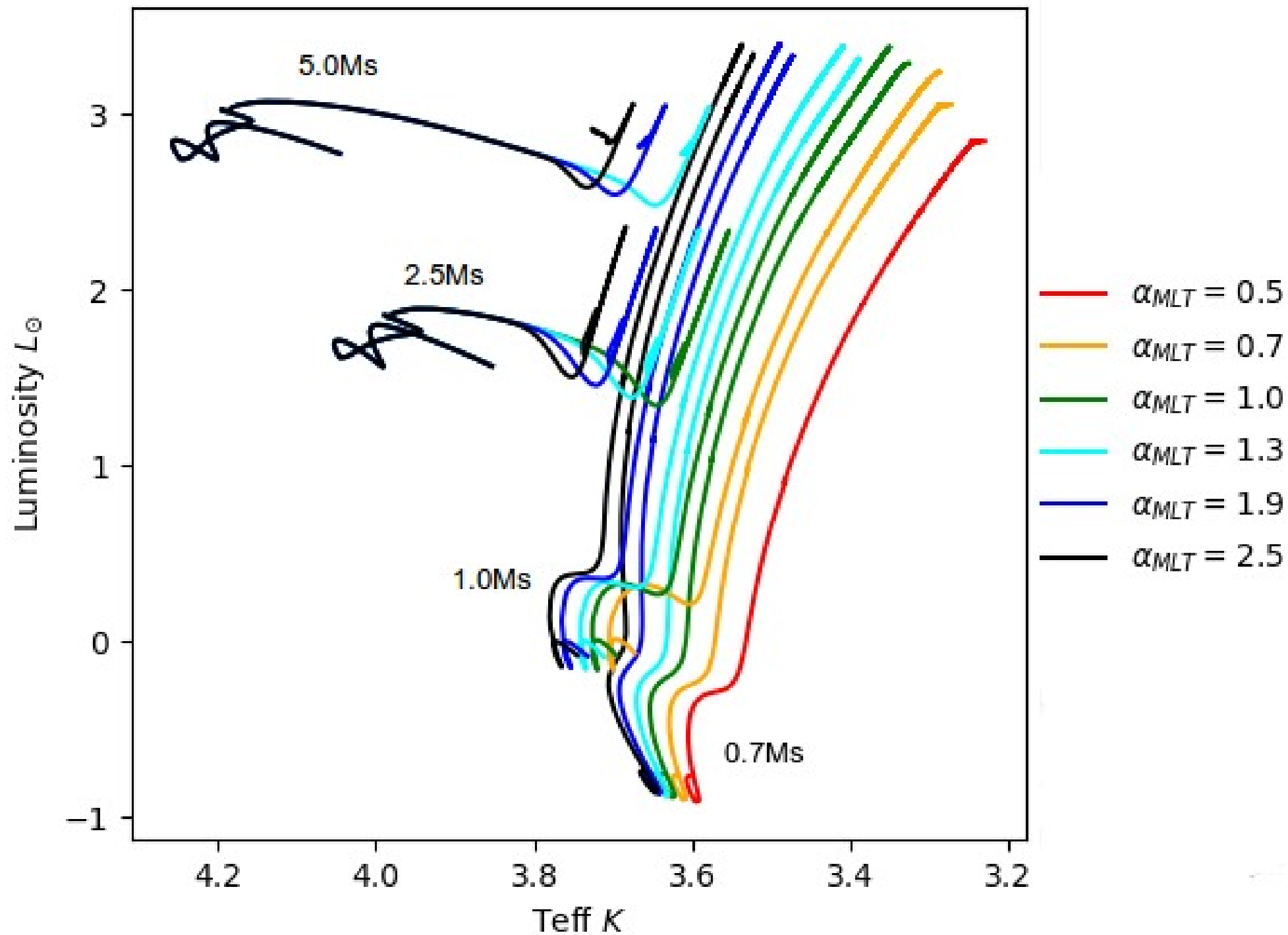


Trend with mass?



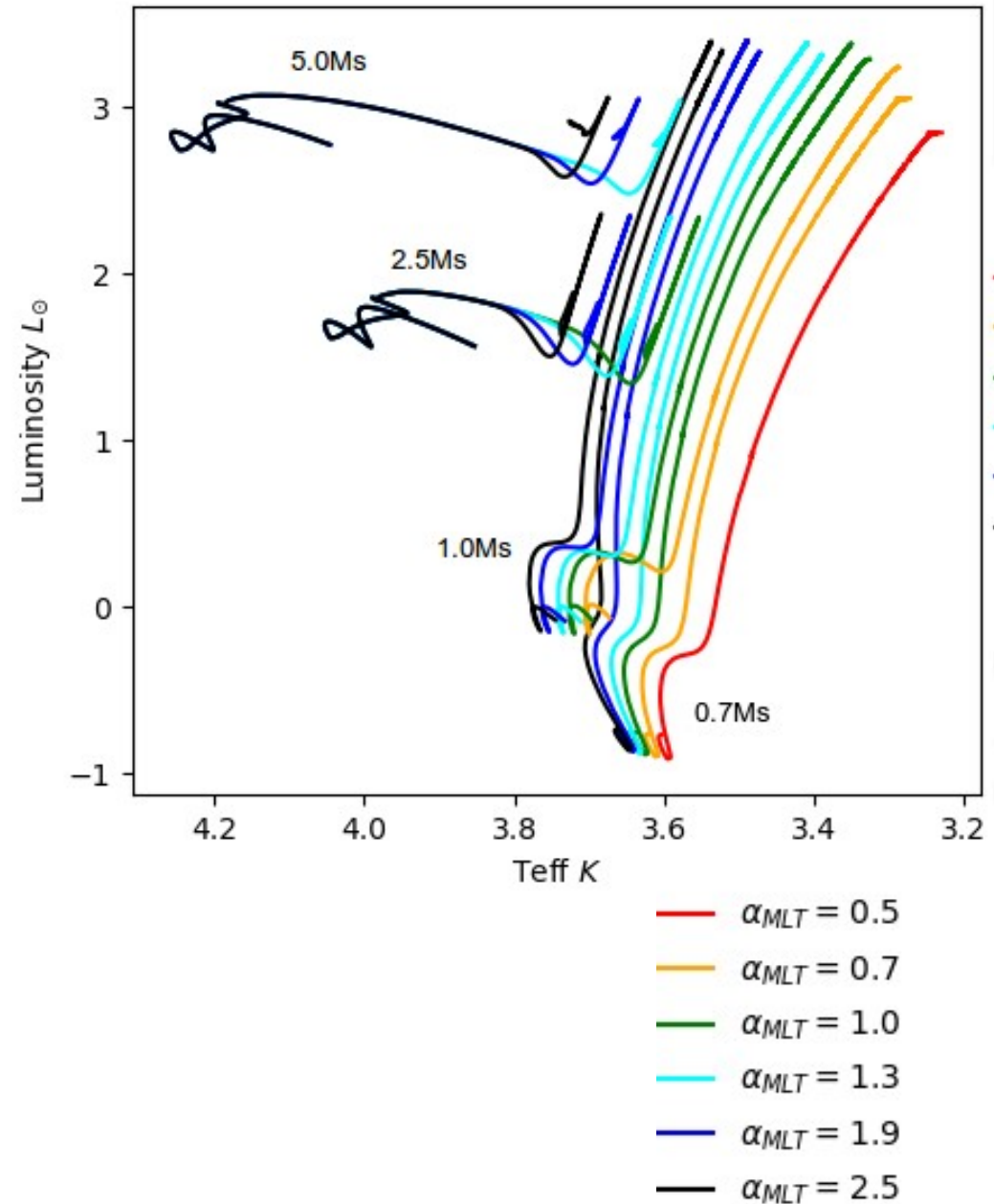
To appreciate the importance
of calibrations, we must
understand how changes to
 α_{MLT} impact stellar models,
and how much

DSEP (Dartmouth Stellar Evolution Program) tracks



Building intuition for general trends:

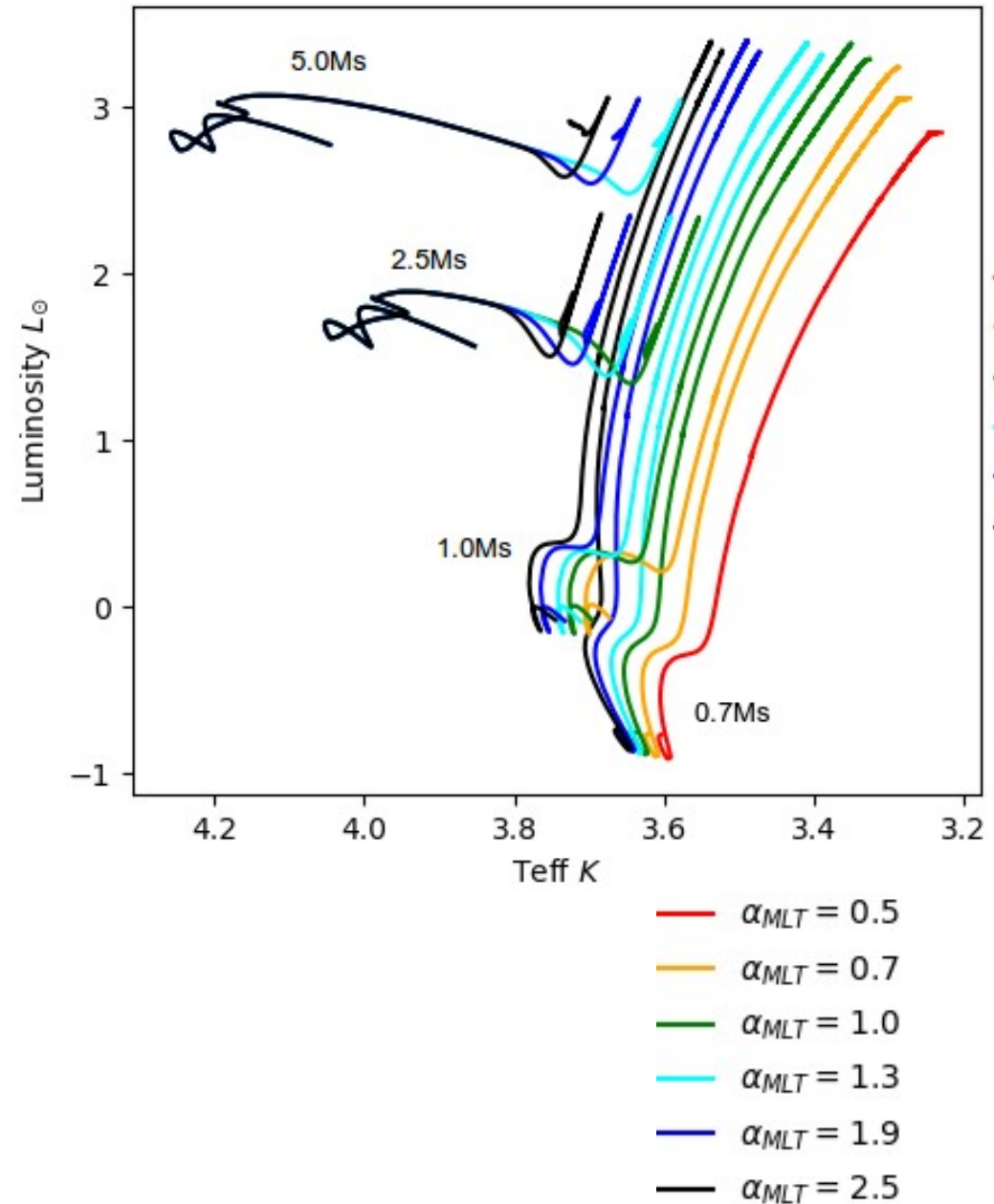
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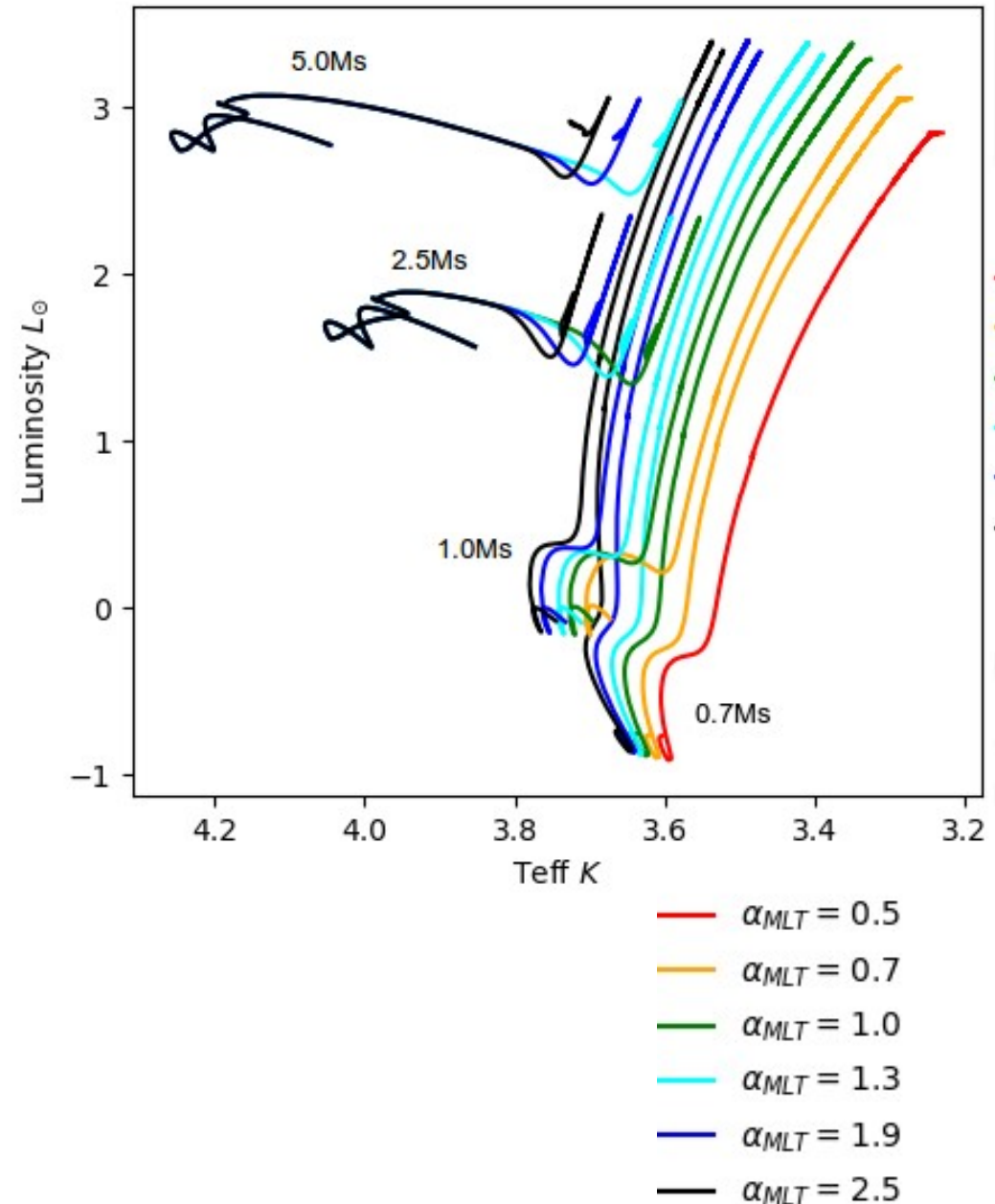


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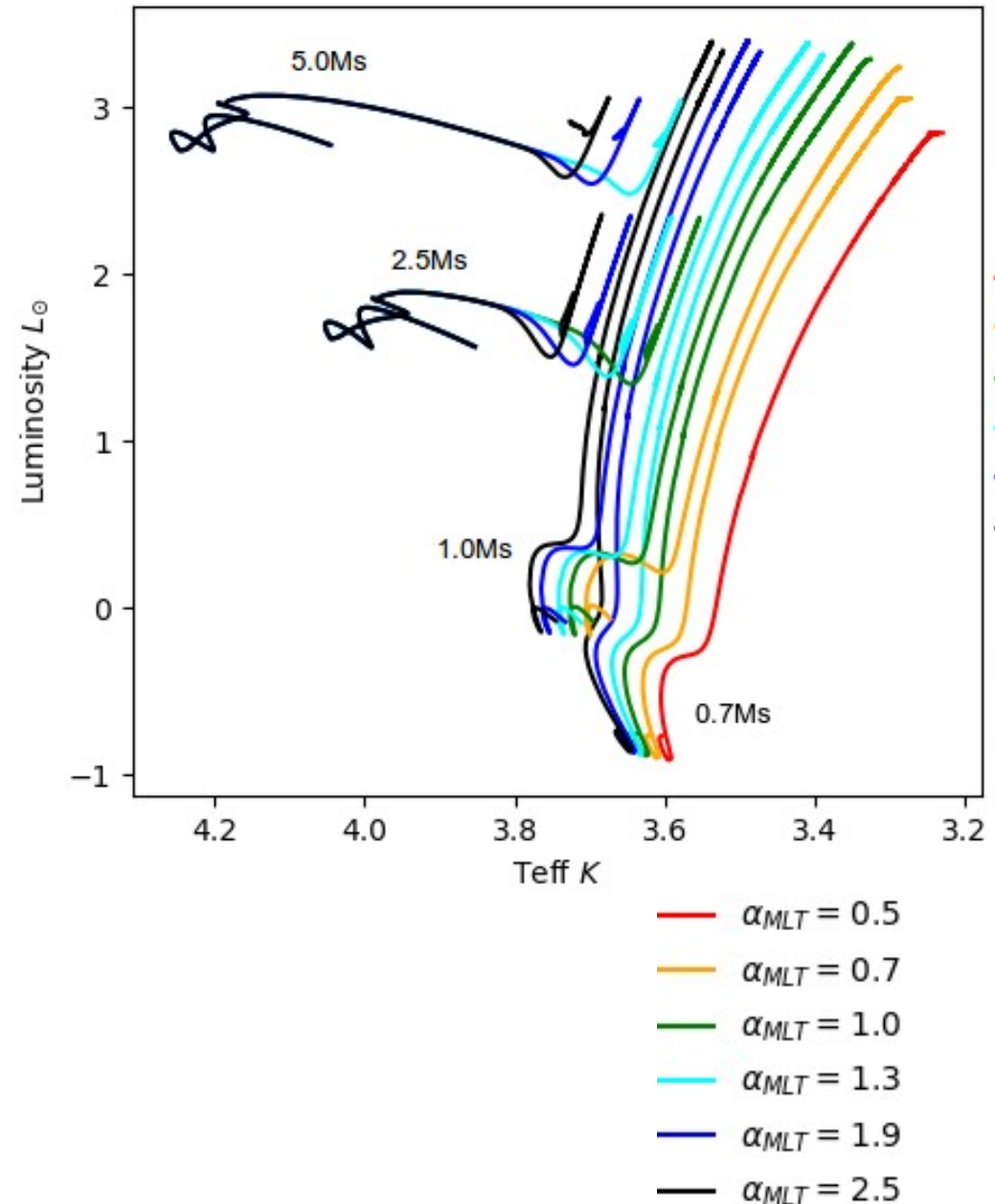
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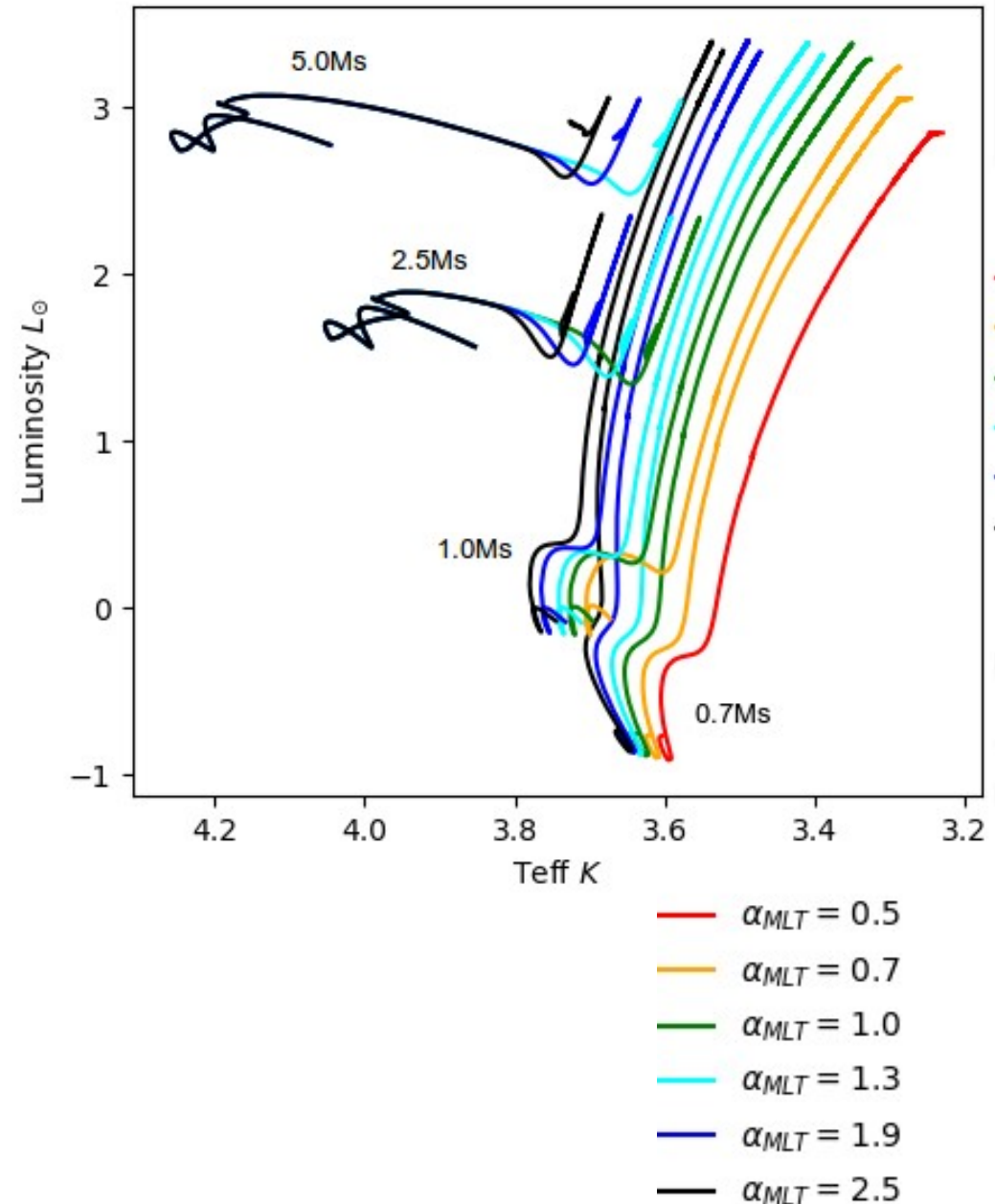
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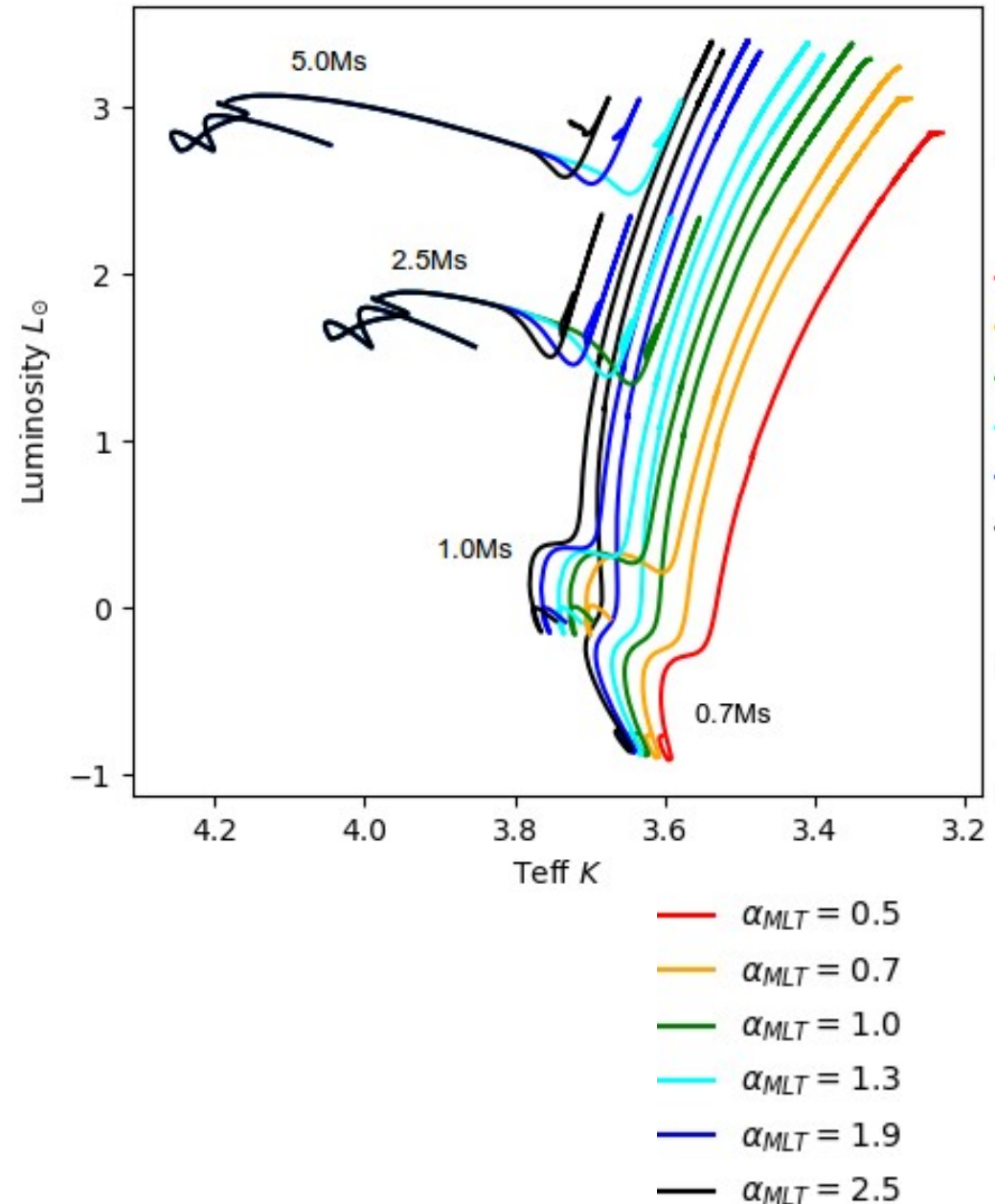
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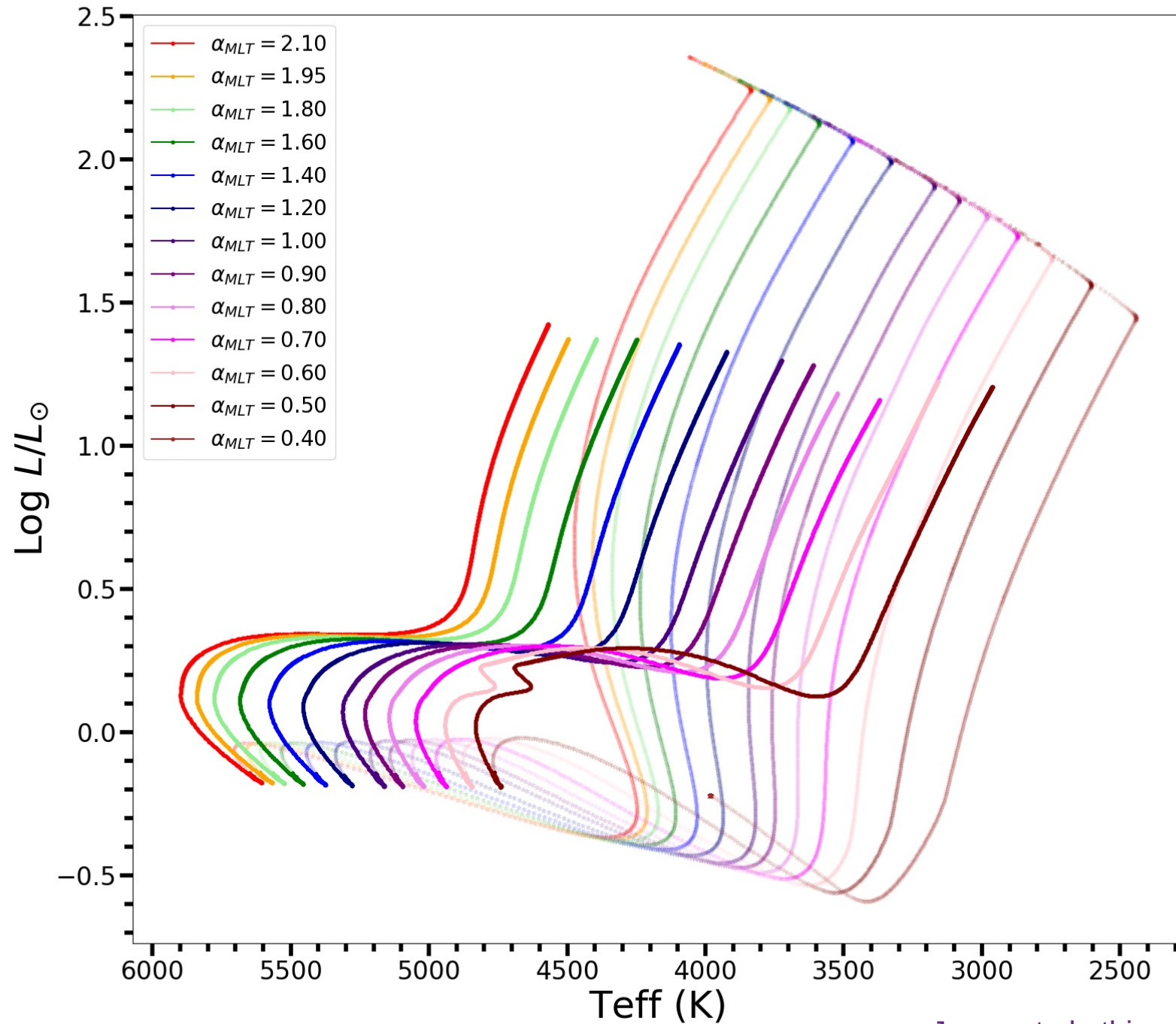
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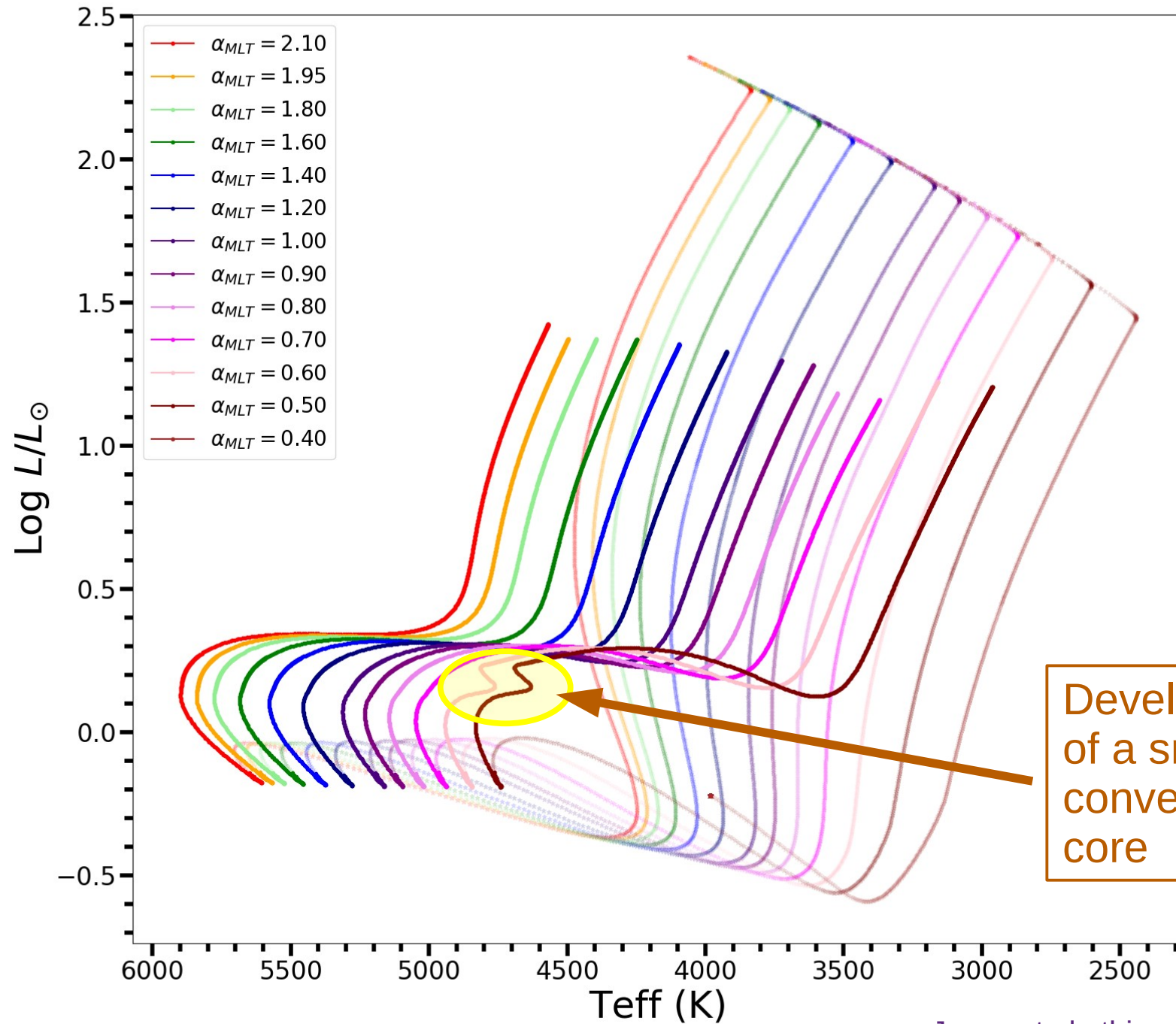
- Efficient transporters of energy should be hotter



Solar-like tracks, MESA v15140+



Solar-like tracks, MESA v15140+



Custom DSEP isochrones with non-solar α_{MLT} and fits to M92

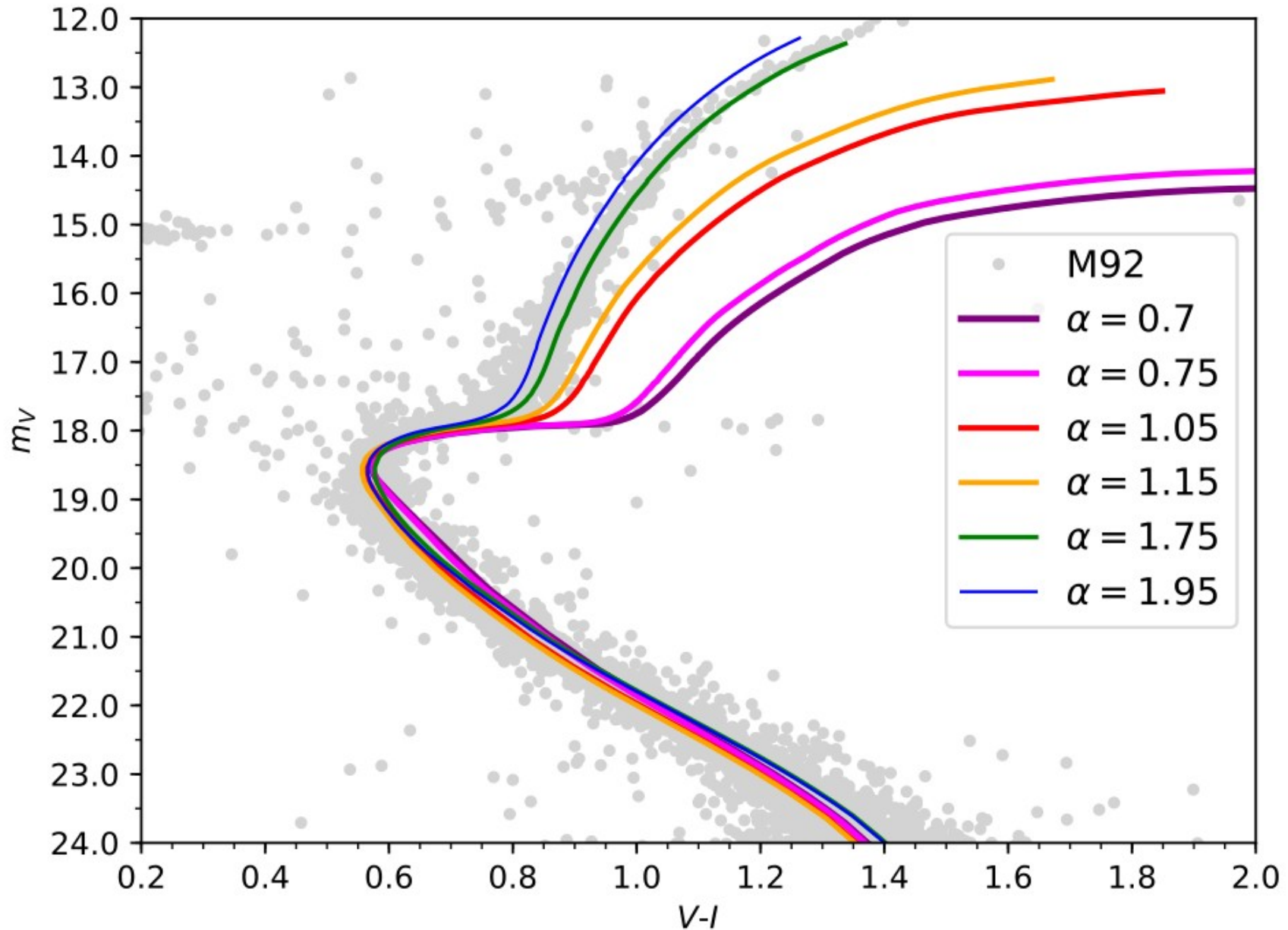
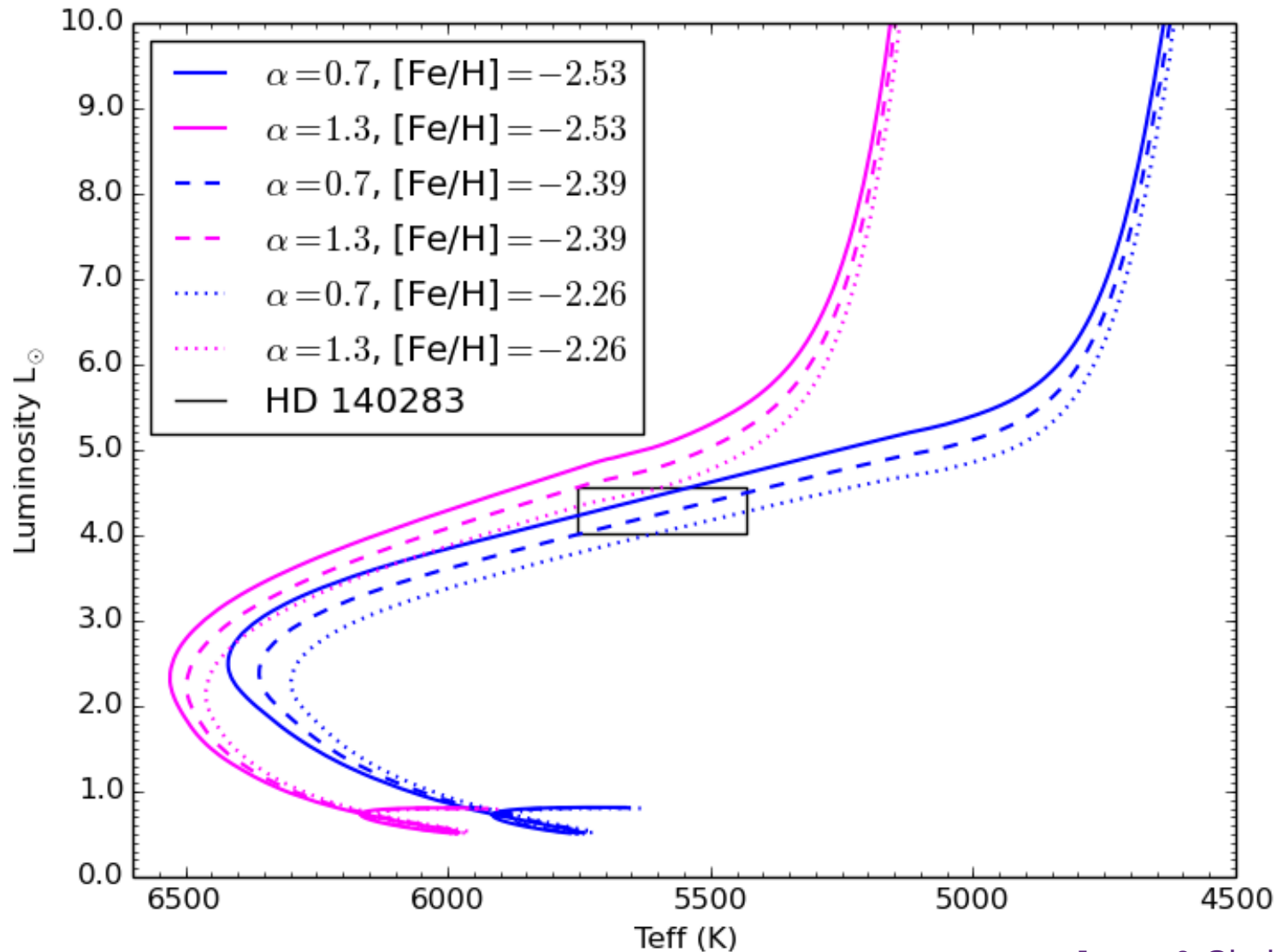
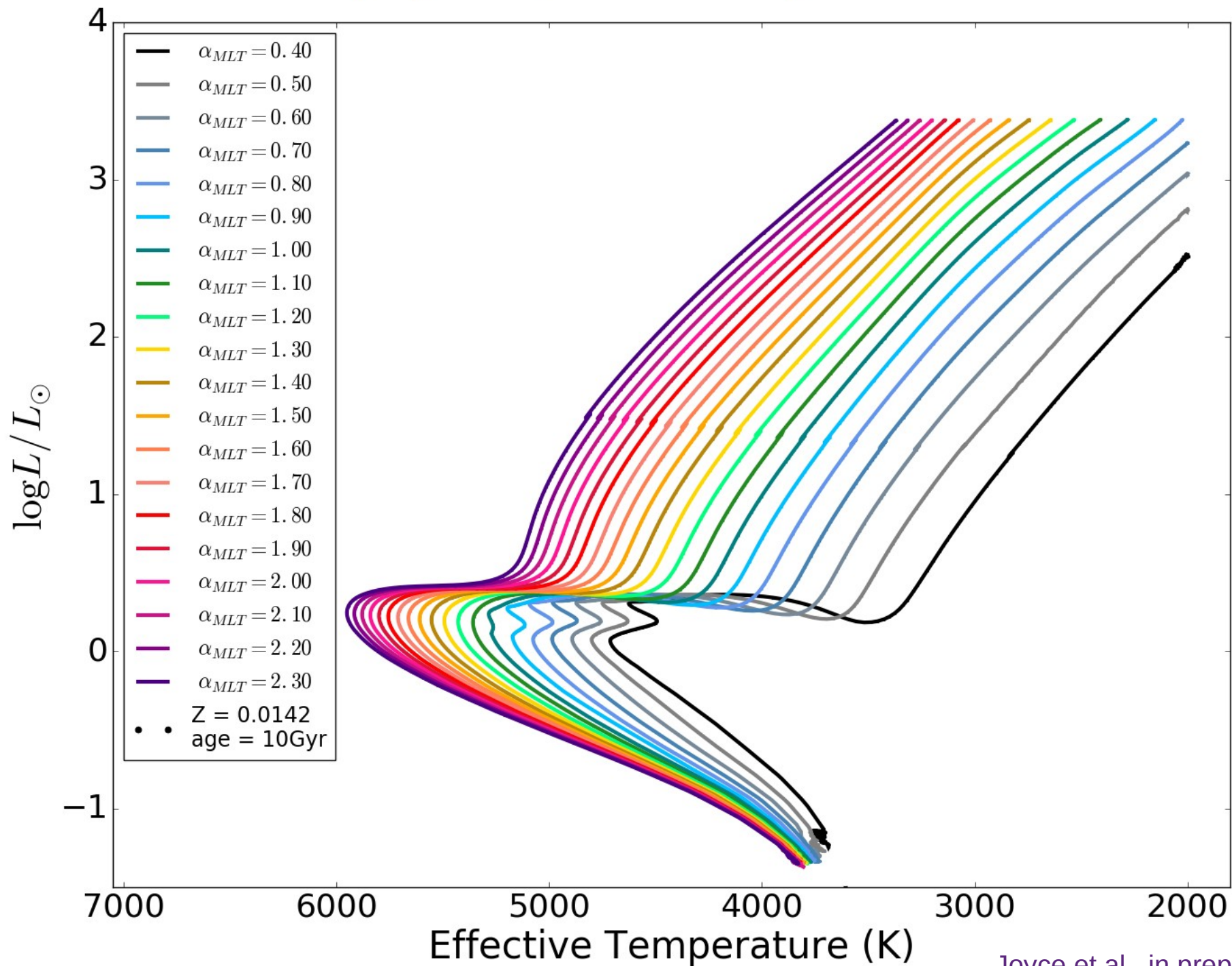


Figure 5. Six isochrones, each of age 13 Gyr, generated with different mixing lengths and shown against M92 for reference.

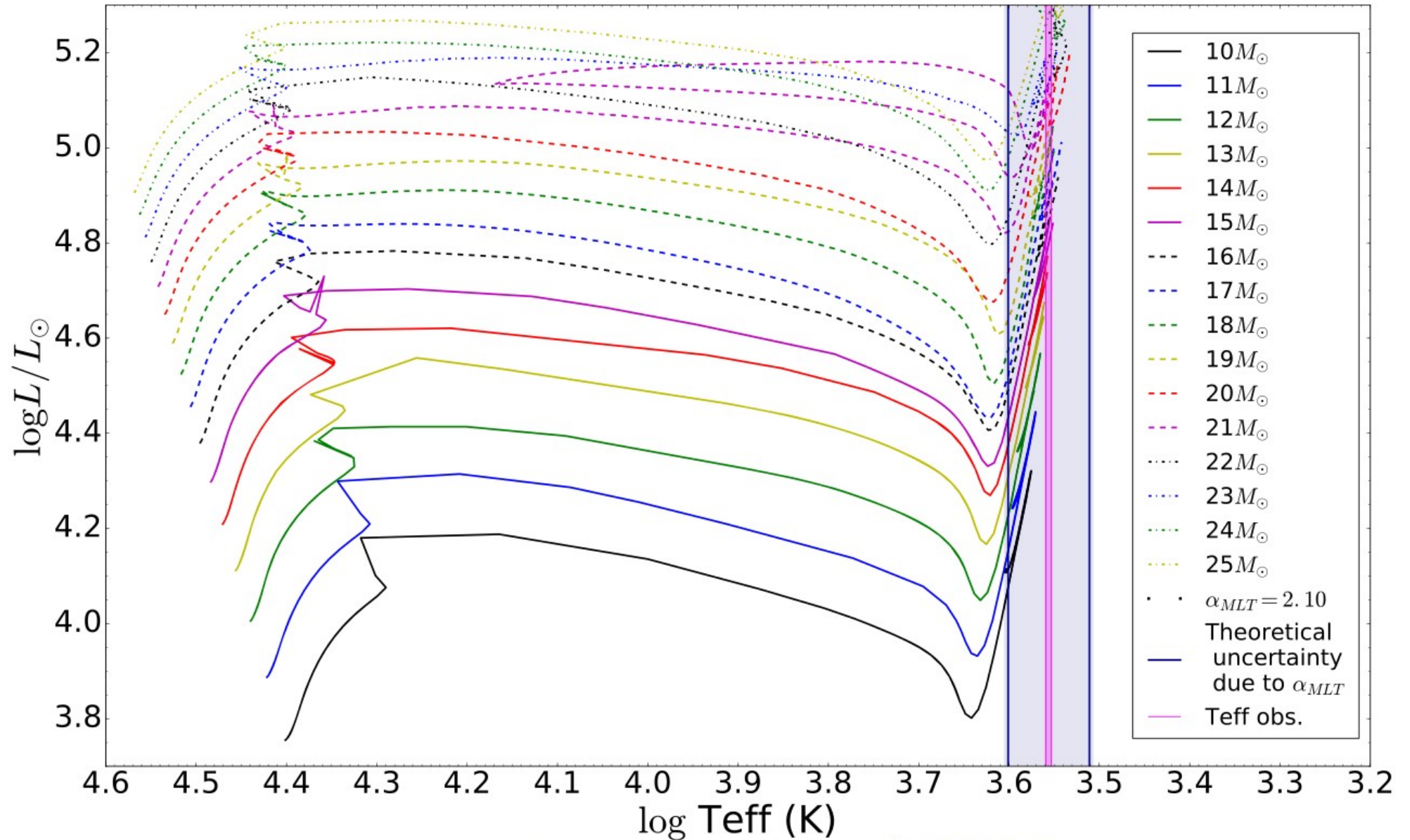
In stellar tracks, changes to the mixing length are indistinguishable from changes in chemical abundance (an observed quantity) below a certain precision



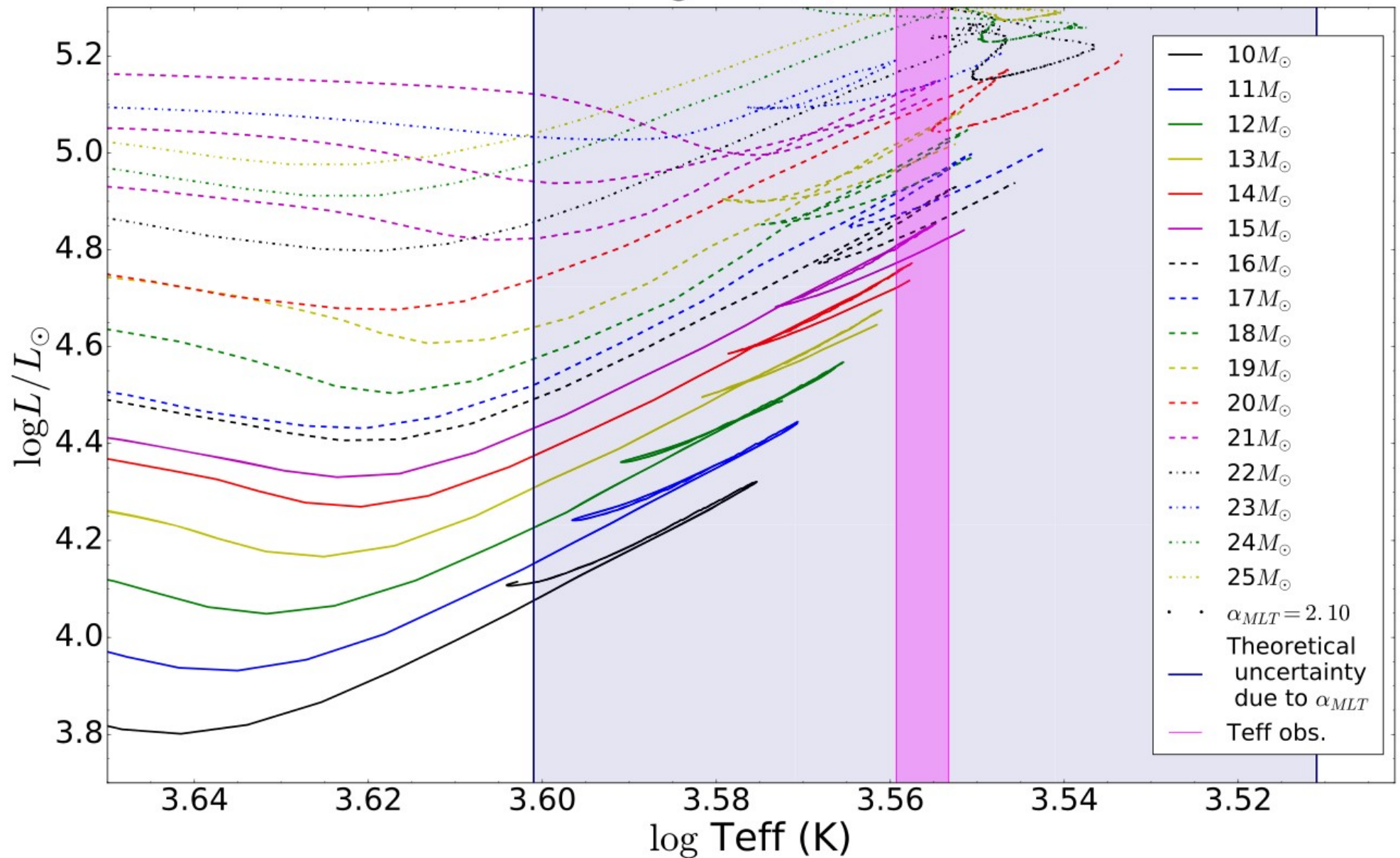
Custom isochrones based on MESA tracks



Observational vs modeling uncertainties: Betelgeuse



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Future



The future begins with the past: Summary of modern attempts to improve MLT

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- (5) direct calibrations of the mixing length parameter against hydrodynamic simulations (ex. Trampedach et al., 2014)

A deeper dive into three
of these modern
improvements

(2) Full spectrum turbulence models (Canuto et al., 1991)

Statement of problem:

1. *The MLT treats the energy spectrum of the turbulent eddies as if it consisted of only one large eddy.*

(2) Full spectrum turbulence models (Canuto et al., 1991)

Solution (without math):

In conclusion:

1. The new model describes turbulence in a more realistic way than the MLT, for it includes the large spectrum of eddies that characterize a nearly inviscid medium such as a stellar interiors.

2. The new model can easily be implemented in existing stellar codes.

3. The new model requires $\alpha < 1$, which is a welcome feature on its own right.

4. The new model gives rise to several differences with respect to the MLT; on the basis of items 1–3 above, they are expected to be in the right direction.

When all the above considerations are taken together, they suggest that the new expression for the convective flux, equation (32), even with the $\Lambda = \alpha H_p$ expression, is to be preferred to the standard MLT.

(2) Full spectrum turbulence models (Canuto et al., 1991)

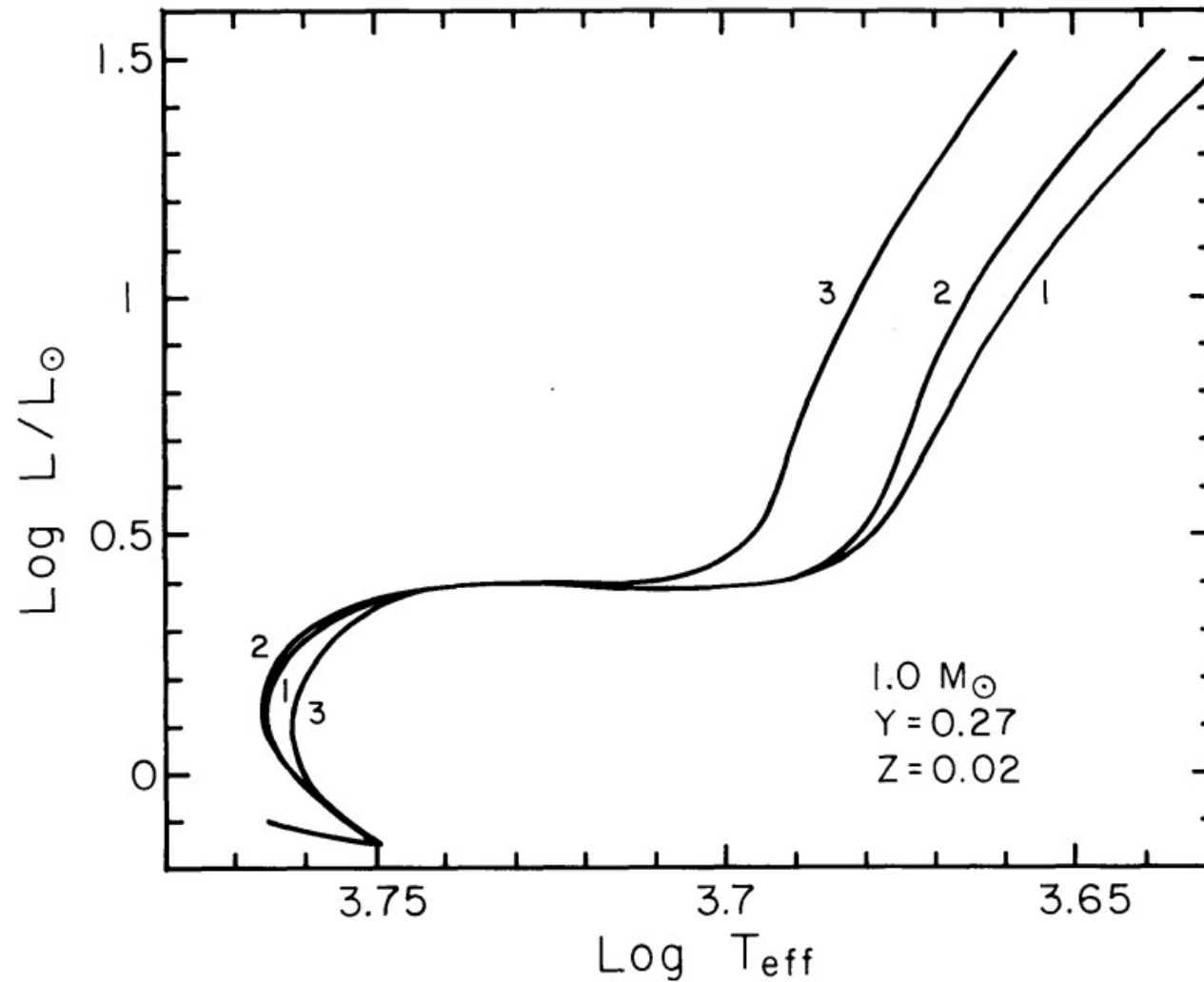


FIG. 6.—Evolutionary tracks in the H-R diagram for a $1 M_{\odot}$ star, with $Y = 0.27$, $Z = 0.02$, for the three cases (1) MLT, $\Lambda = 1.4H_p$; (2) new theory, $\Lambda = 0.7H_p$, no turbulent pressure; (3) new theory, $\Lambda = z$, with turbulent pressure.

(2) Full spectrum turbulence models (Canuto et al., 1991)

1. *The MLT treats the energy spectrum of the turbulent eddies as if it consisted of only one large eddy.*

We thus conclude that the one-eddy model adopted by the MLT may be considered a reasonable approximation for viscous flows, but is a poor approximation for nearly inviscid stellar interiors which are characterized by a wide spectrum of eddies of all sizes.

(4) Using 3D radiative hydrodynamics to inform the choice of α_{MLT} (Ludwig et al., 1999)

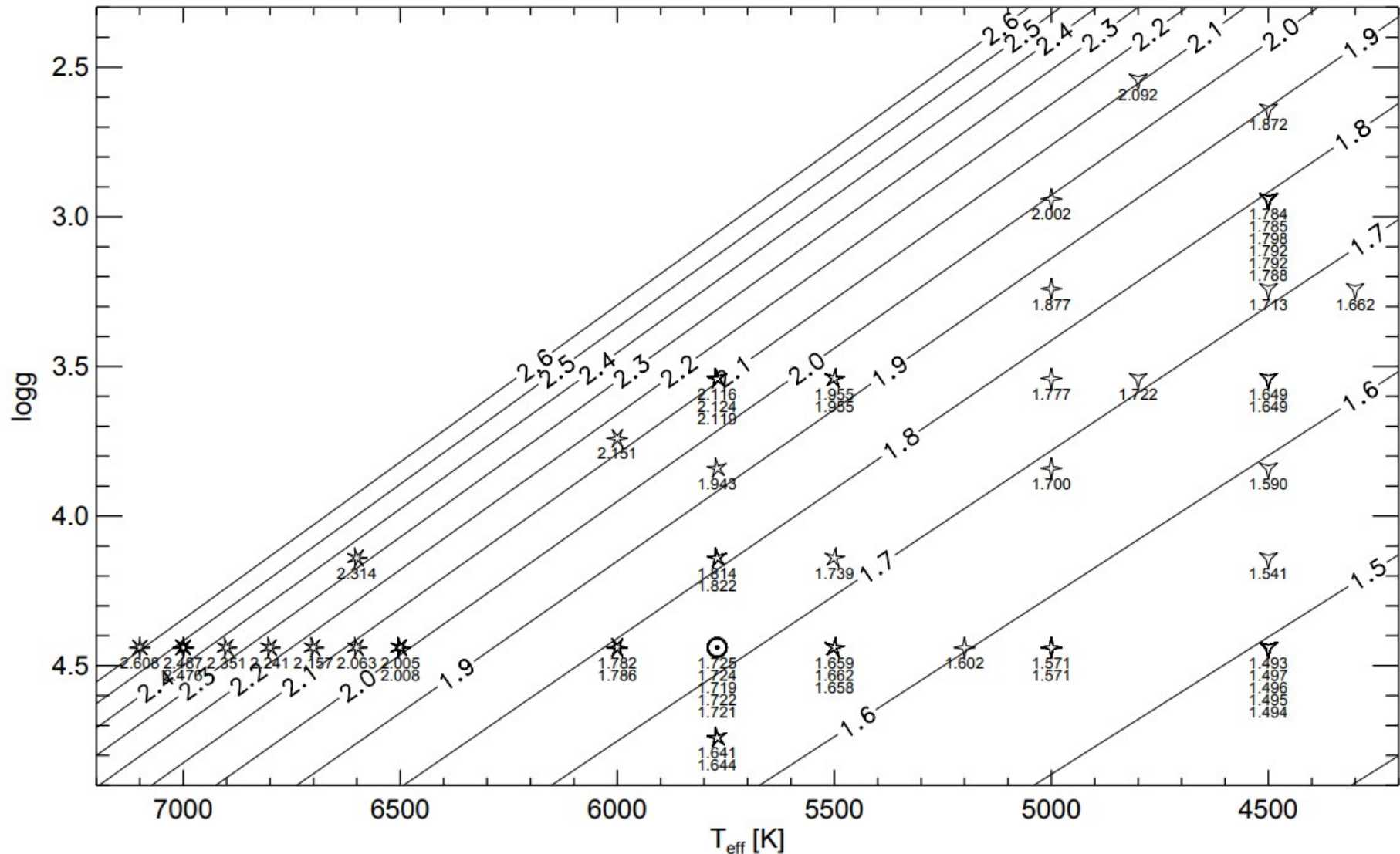


Fig. 3. s_{env} in units of $s_0 = 10^9 \text{ erg/g/K}$ from the grid of RHD models for $Y = 0.28$ and solar metallicity. Symbols indicate RHD models. Attached to the symbols the actual s_{env} values are given; the contour lines present a fit to them. The fitting function is given in Appendix C.

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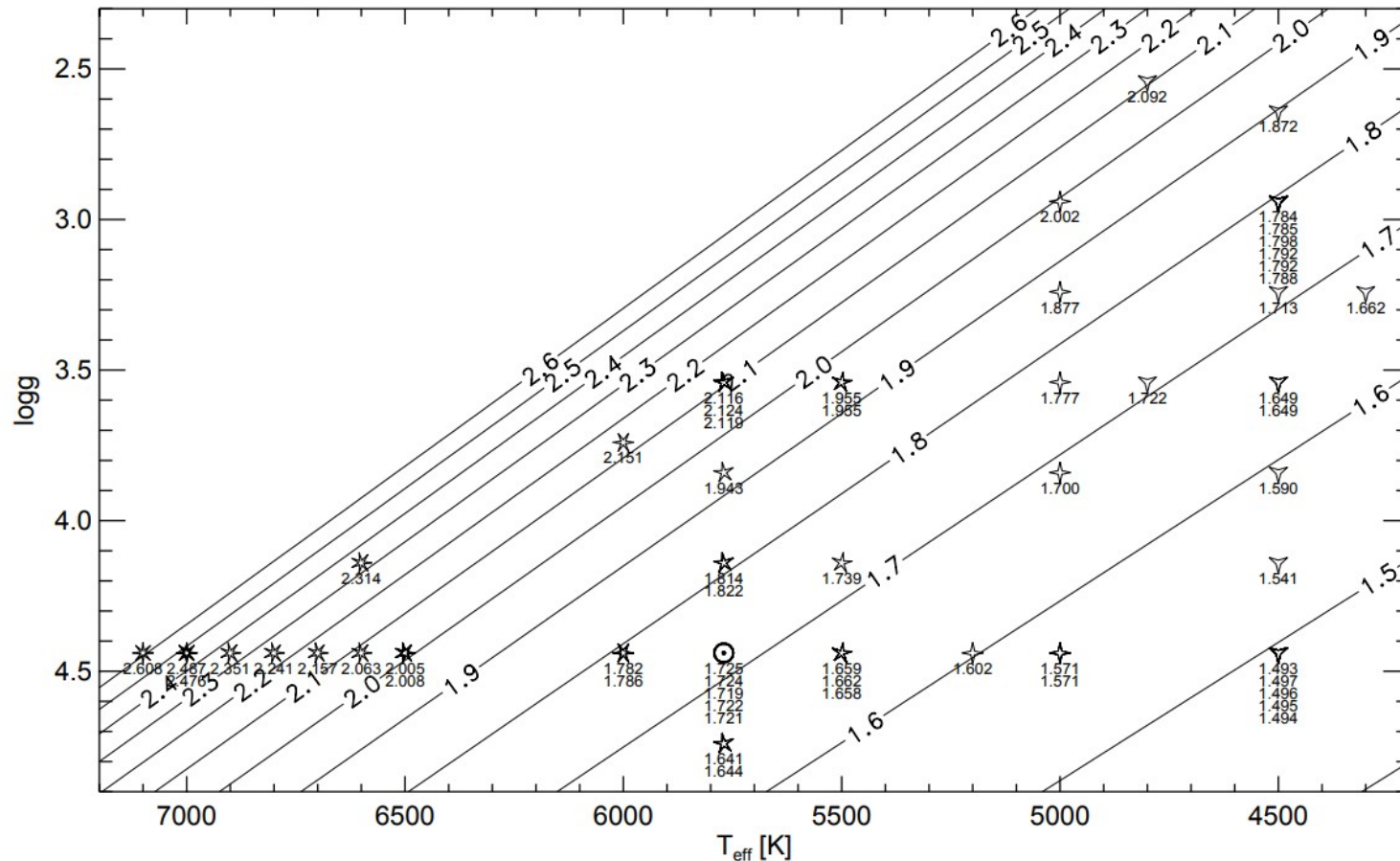


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- trends are qualitatively the same as those shown by 1D sims
- in their own words: “do not extrapolate outside of this graph!”

(5) Direct calibration against 3D hydrodynamics
(Trampedach et al., 2014)

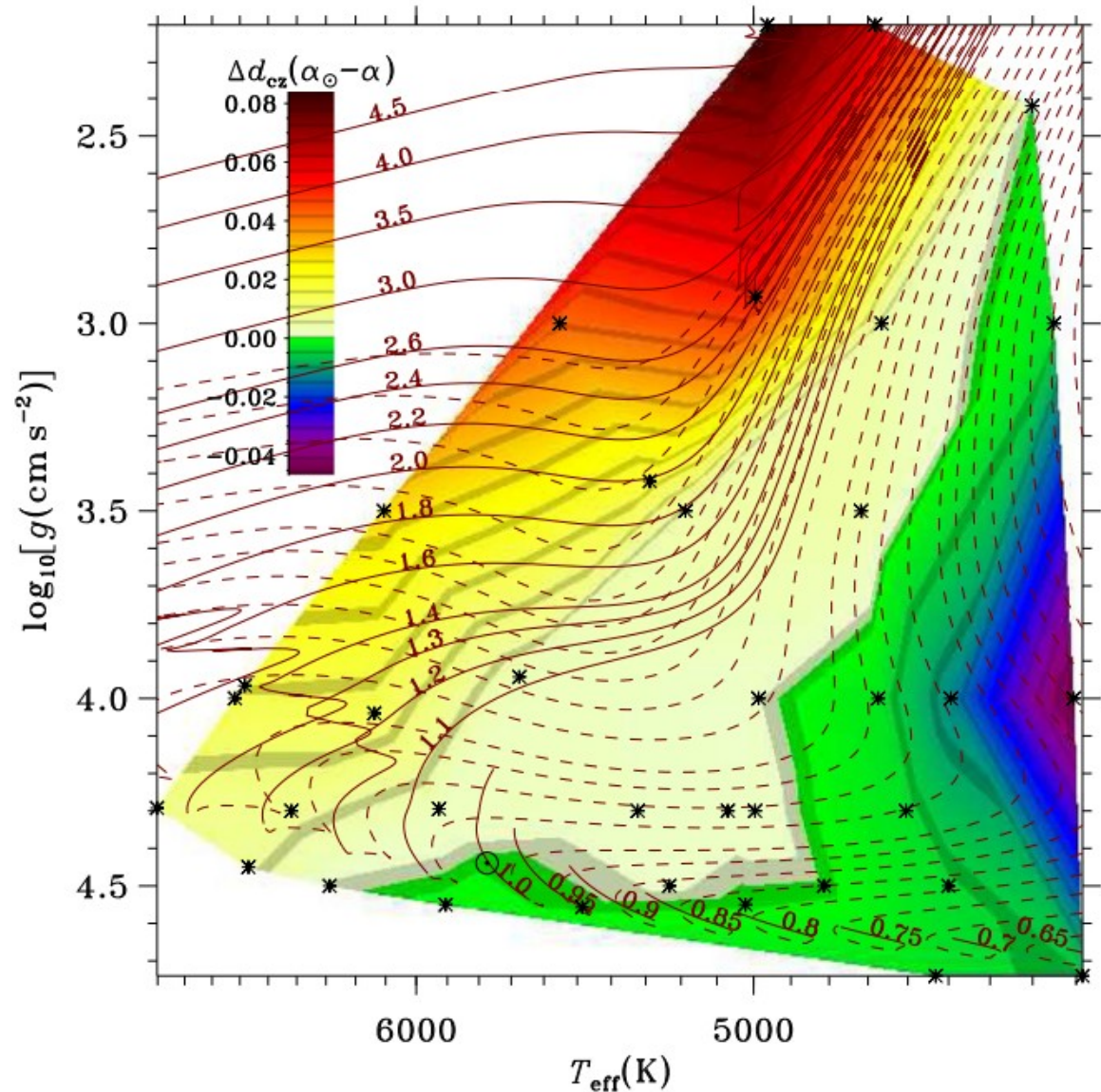


Figure 11. As Fig. 1, but showing differences in the depth of the convective envelope between models using solar-calibrated and individually calibrated values of α , in the sense that warm giants with α_{\odot} have deeper convective envelopes.

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Trend: lower values of alpha
(less efficient convection) are
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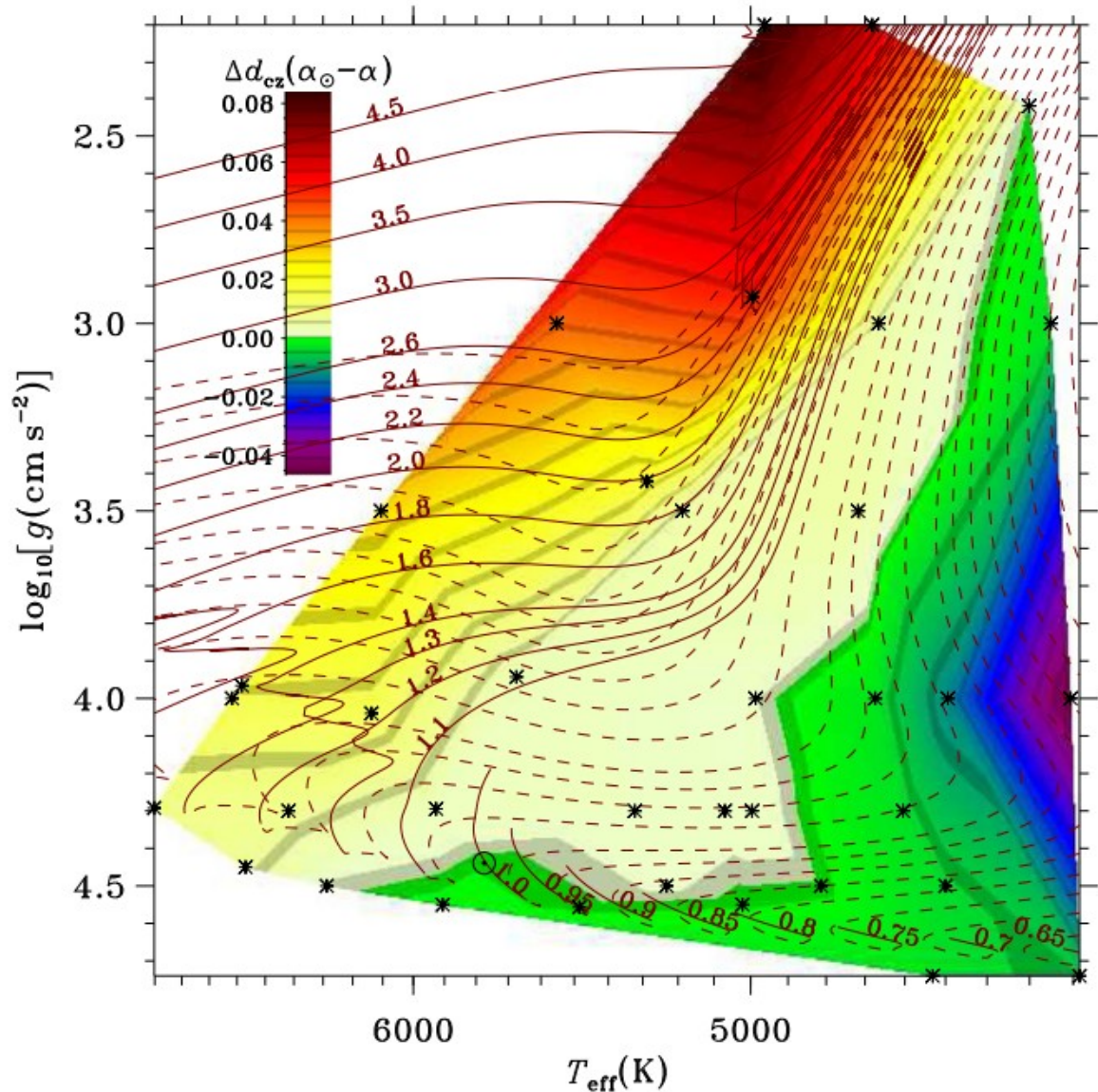


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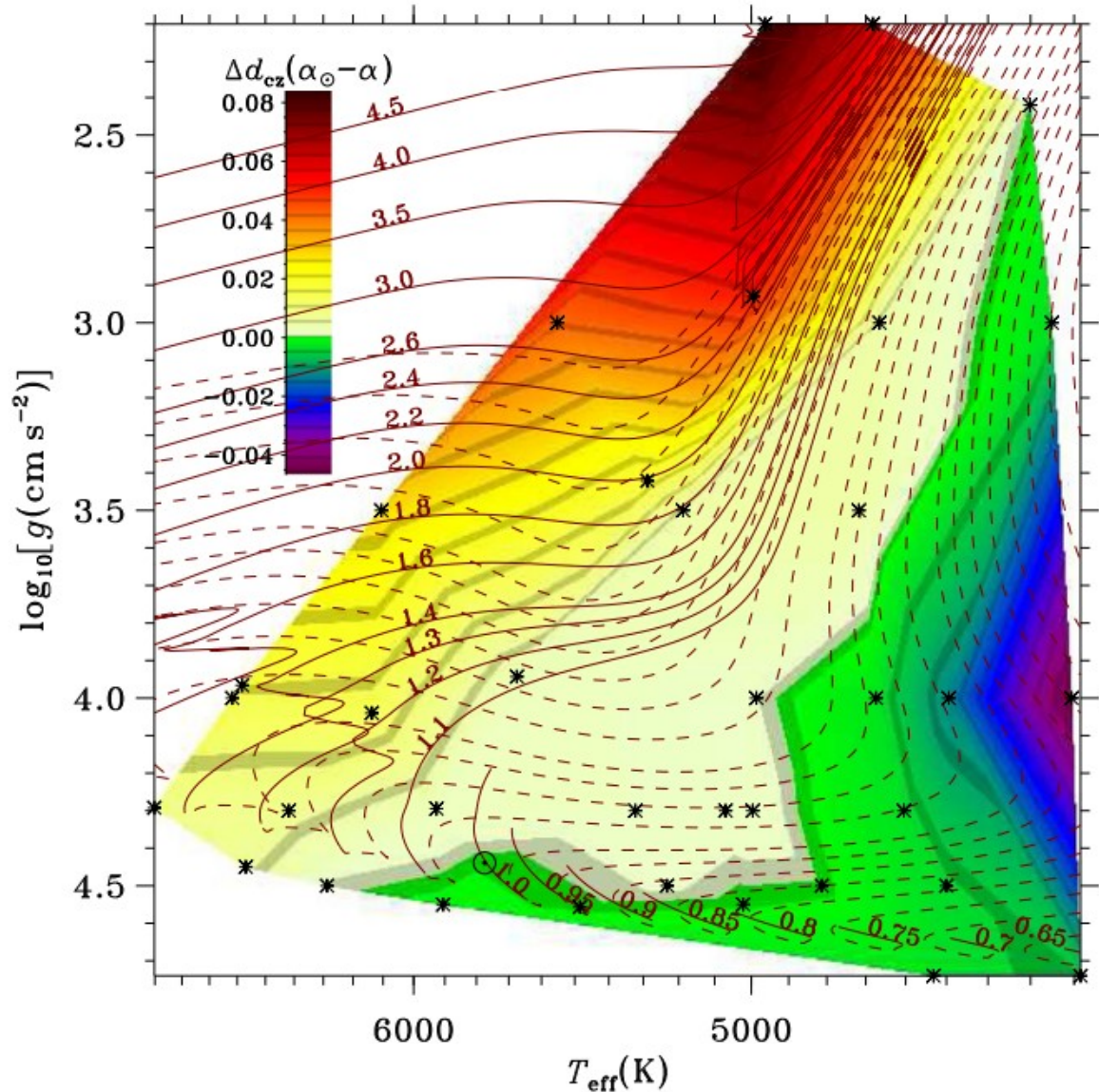


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Question: how is this compatible with other work suggesting fully convective M-dwarfs require ultra-low alpha_MLT values?
(See Mann et al., 2015, and references therein)

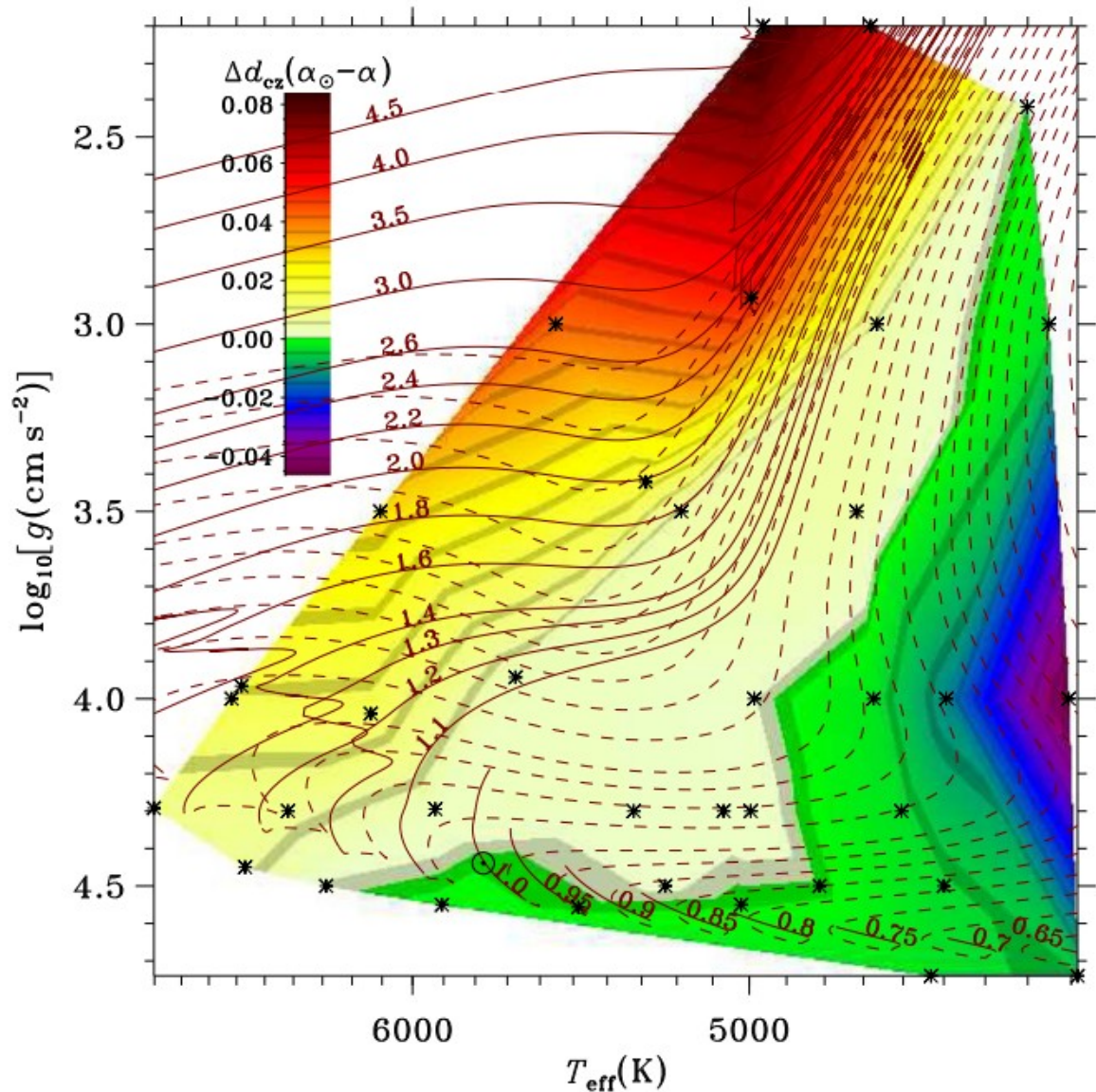


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Where do we go from here?

Q: Is mixing length theory going away?

Q: Is this formalism going to get better?

Q: What role does 3D hydro play?

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*and it most likely is not, unless you are me or one of a handful of other people intentionally probing this regime. Your choices for, e.g., atmospheric boundary conditions or color corrections should be examined first

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Why- because MLT is a calibration-based formalism inherently, and stellar evolution codes are designed around this treatment

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...but these are both in tension with MLT-vs-fundamental stellar parameter trends estimated for large populations of stars (by machine learning techniques and other methods)

As more targets become available for precision 1D calibrations, we must continue to validate these results against 3D simulations with the goal of building a calibration library that is **informed in equal parts by observations and the truth of fluid dynamics**

Fin

$$F_c = \frac{1}{2} \rho v c_D J \frac{\lambda}{H_p} (\nabla_{T^*} - \nabla_{T_{ad}}) \text{ with } \alpha_{mLJ} = \frac{\lambda}{H_p}$$