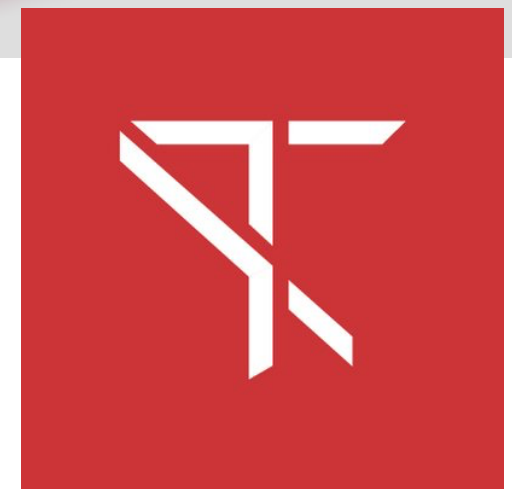
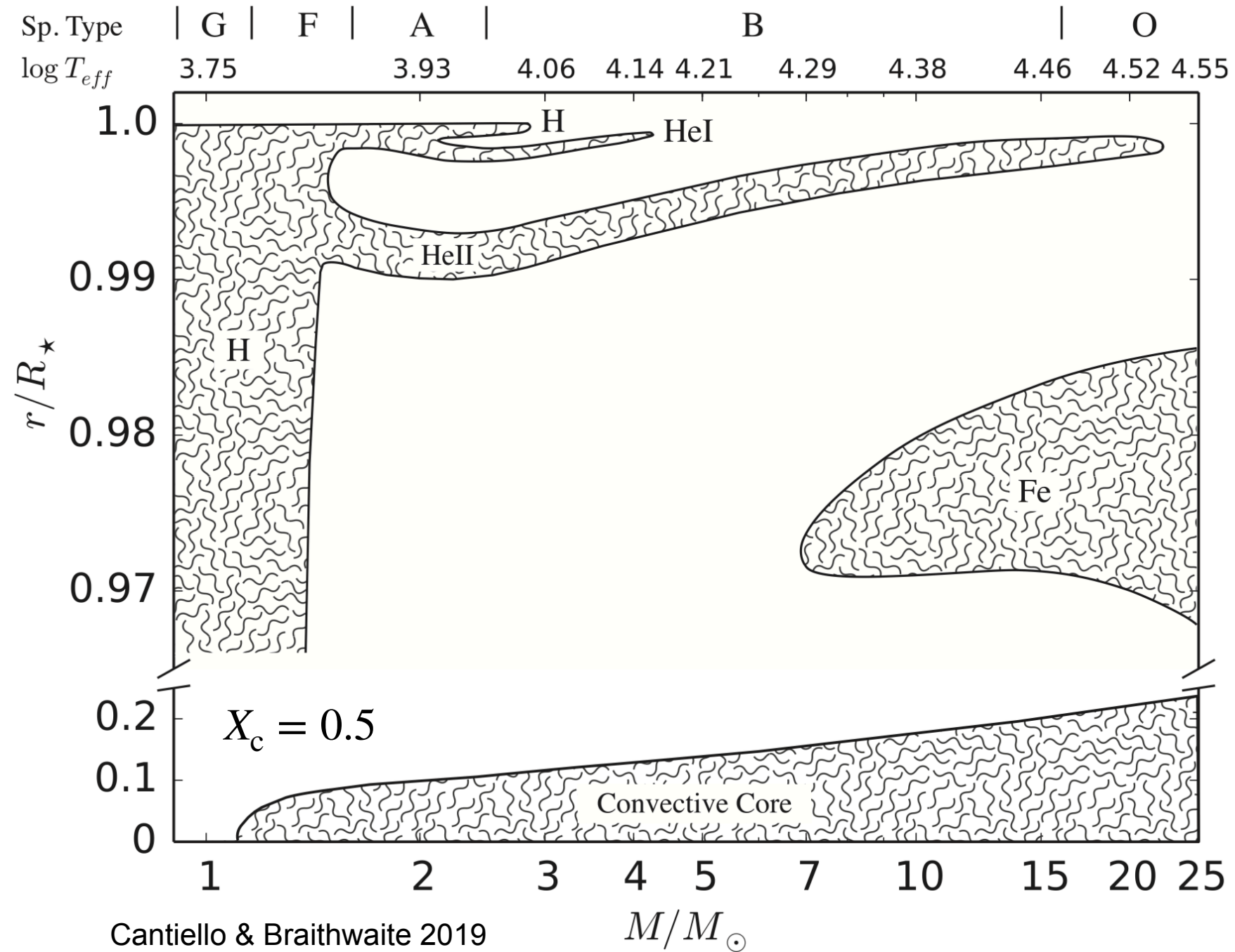


In collaboration with Lars Bildsten and Yan-Fei Jiang
TranStar21 Program @ KITP
December 16th 2021



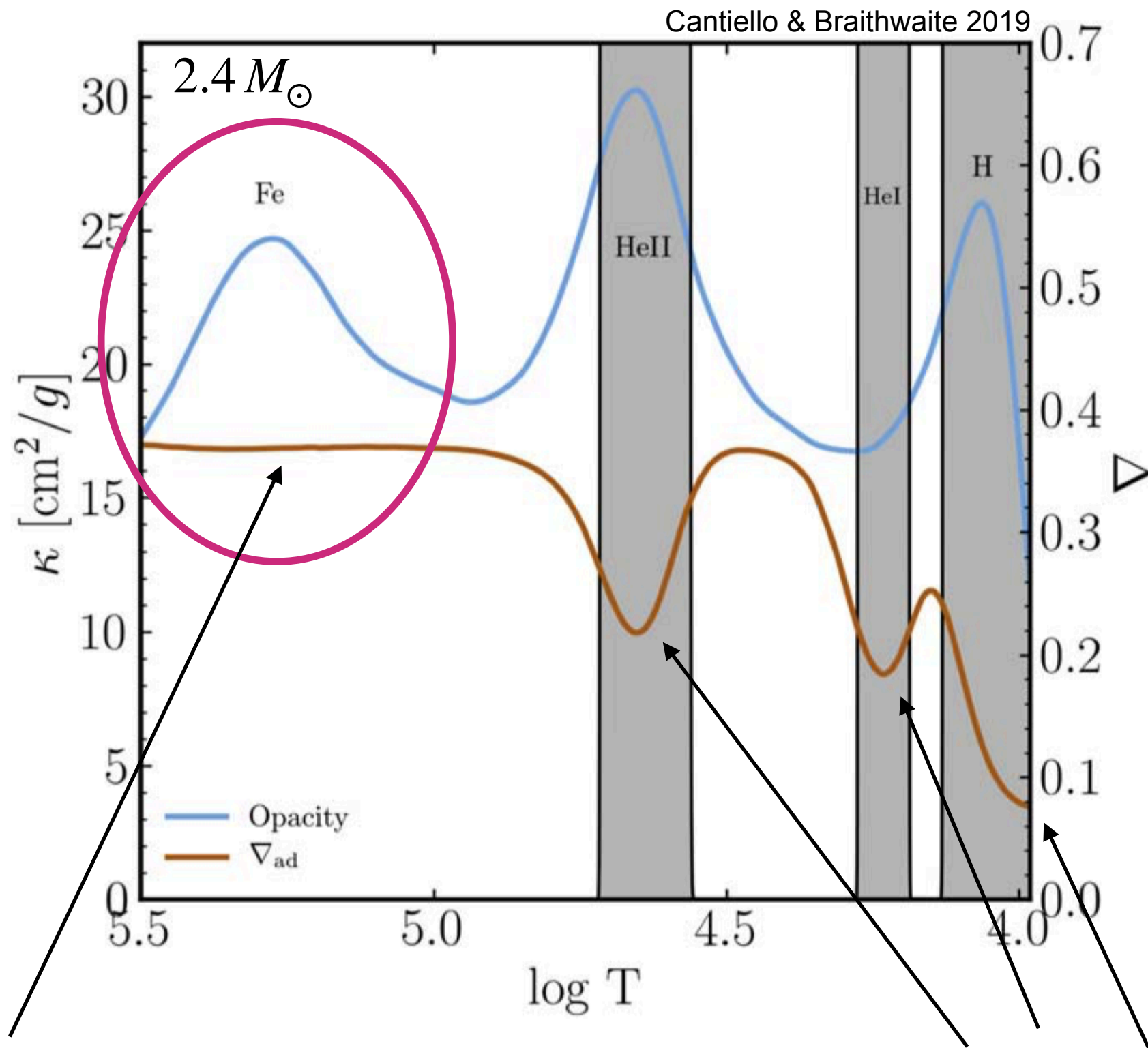
Sub-Surface Convection Zones in HRD



Sub-surface convection zones ubiquitous in massive stars

“Sub-surface” is very close to the surface. Maybe too close ...

Causes of Sub-Surface Convection Zones

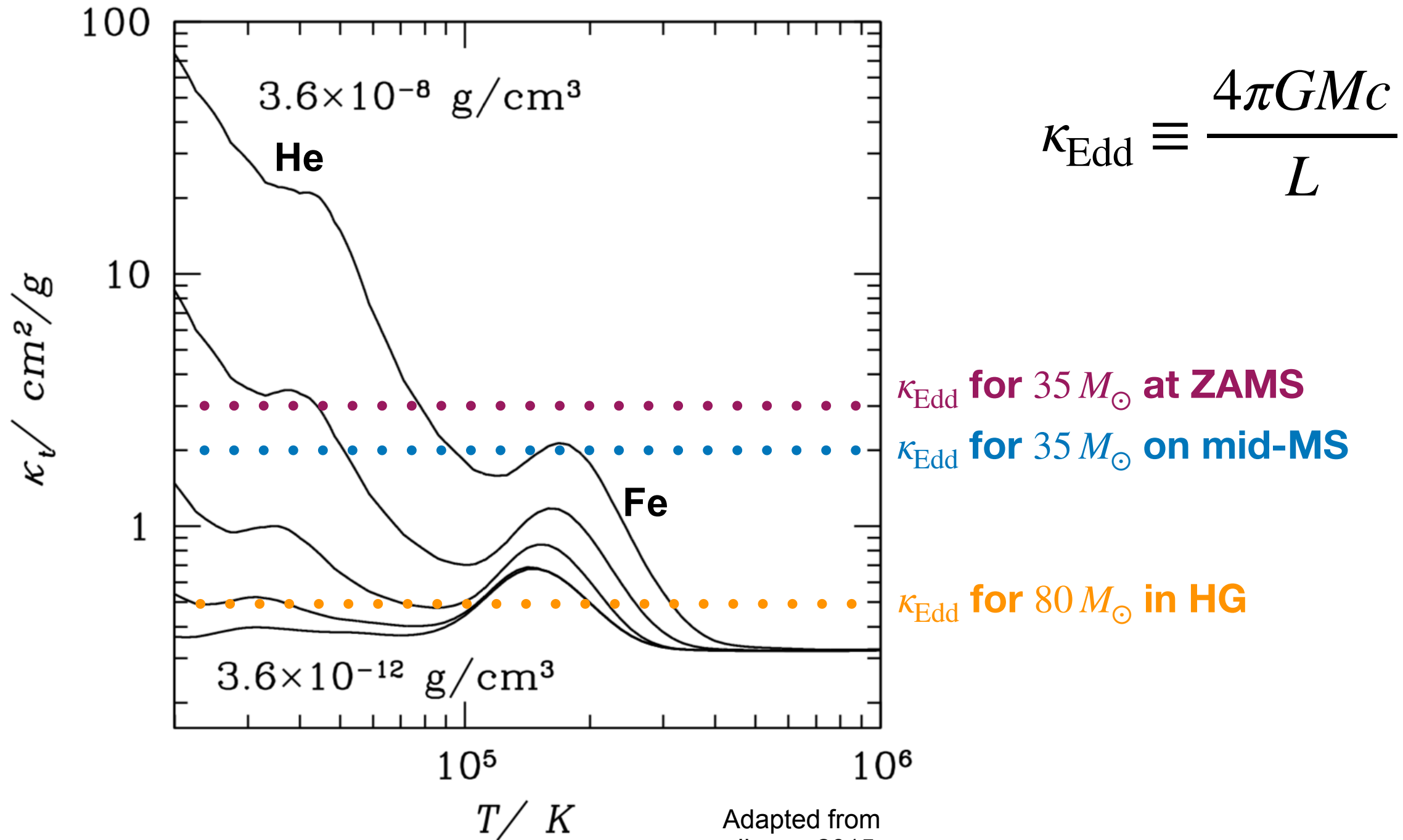


Need $L \sim L_{\text{Edd}}$ to
excite convection

($M \gtrsim 10 M_{\odot}$)

Partial ionization reduces ∇_{ad}

Convection Excited by Opacity Peaks



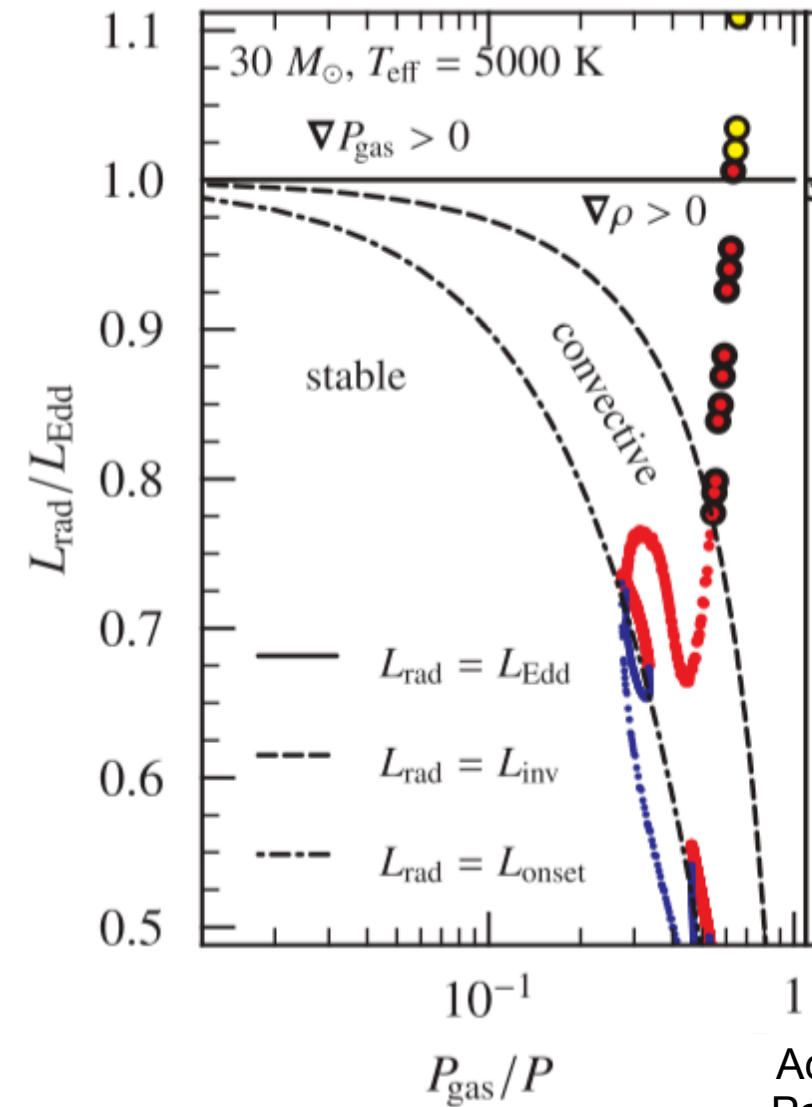
Adapted from
Jiang+ 2015

Density Inversions in 1D Models

With sufficient Eddington ratio

- Convection excited
- Density inversions form
- Gas pressure inversions can form

MLT++ was designed to reduce/
eliminate density inversions



Adapted from Paxton+ 2013

Started 3D exploration using Athena++

Jiang+ 2015
Plane parallel
RHD

Jiang+ 2017
Plane parallel
M+RHD

Jiang+ 2018
1/4 star LBV
RHD

Schultz+ 2020
Exploring
LBV models

Schultz+ in press
Narrow Box MS
RHD

Convective Efficiency γ

From ∇

$$\gamma \equiv \frac{\nabla - \nabla_{\text{eddy}}}{\nabla_{\text{eddy}} - \nabla_{\text{ad}}} \qquad \nabla \equiv \frac{d \ln T}{d \ln P}$$

$\gamma \gg 1$ convection is efficient

$\gamma \ll 1$ convection is inefficient

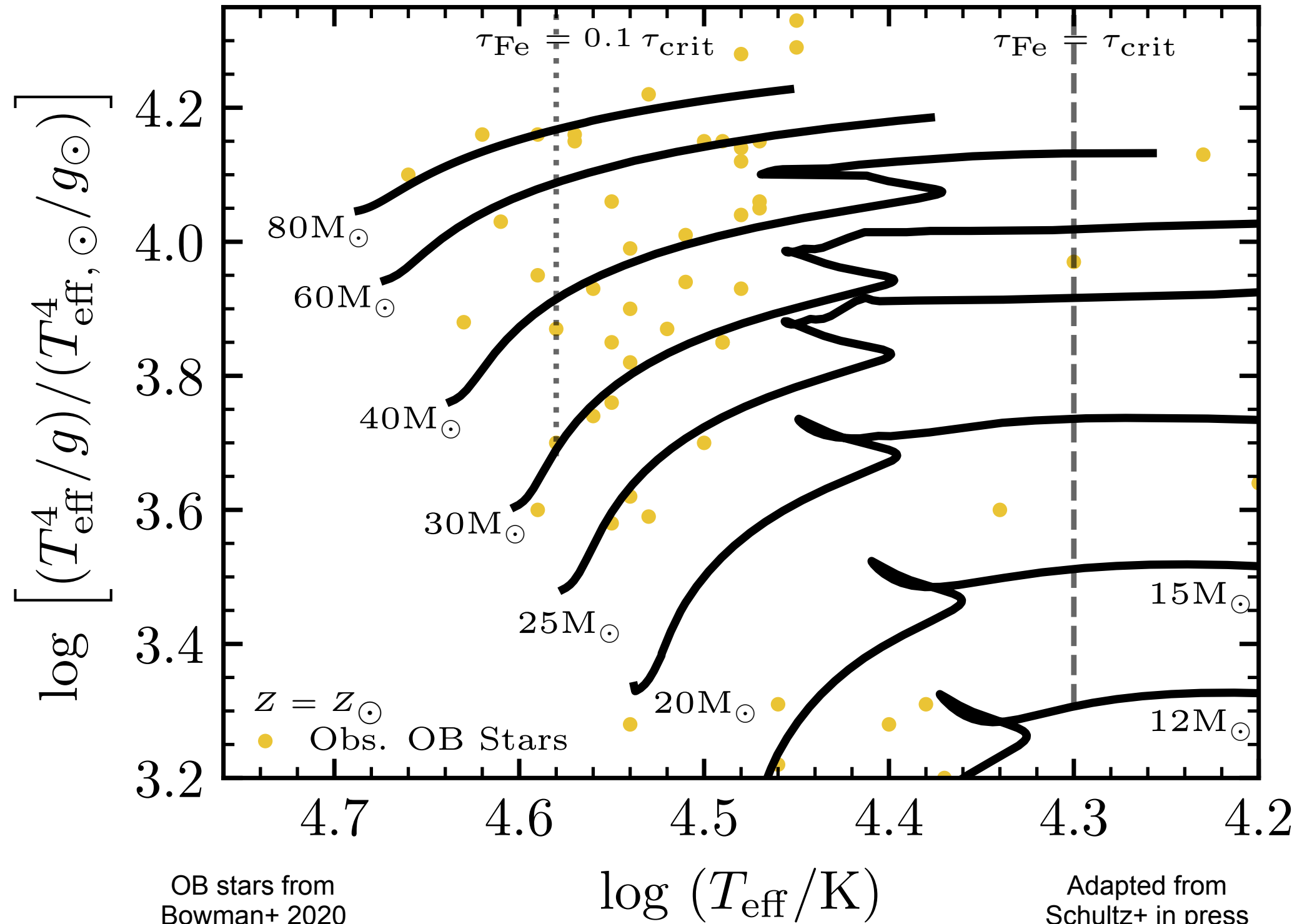
In optically thick limit

$$\gamma \sim \frac{F_{\text{conv}}}{F_{\text{rad}}} \sim \frac{(P_{\text{rad}} + P_{\text{gas}})v_c}{P_{\text{rad}} \left(\frac{c}{\tau} \right)} \qquad \tau_{\text{crit}} \equiv \frac{c}{v_c} \frac{P_{\text{rad}}}{(P_{\text{rad}} + P_{\text{gas}})}$$

If convection occurs at $\tau < \tau_{\text{crit}}$, expect inefficiencies due to radiative losses

For massive stars, $\tau_{\text{crit}} \sim 10^3 - 10^4$

τ_{crit} Across the High L sHRD



OB stars from
Bowman+ 2020

MESA models from
Cantiello+ 2021

Adapted from
Schultz+ in press

3D Rad-Hydro Modeling using Athena++

Modeling Method

- Initializing Models

- ▶ Use MESA to get $T, \rho, L, M_{\text{in}}$
- ▶ Take T, ρ at iron opacity peak, integrate to find radiative only model
- ▶ Use radiative solution as initial condition and integrate 3D model with L, M_{in} from MESA

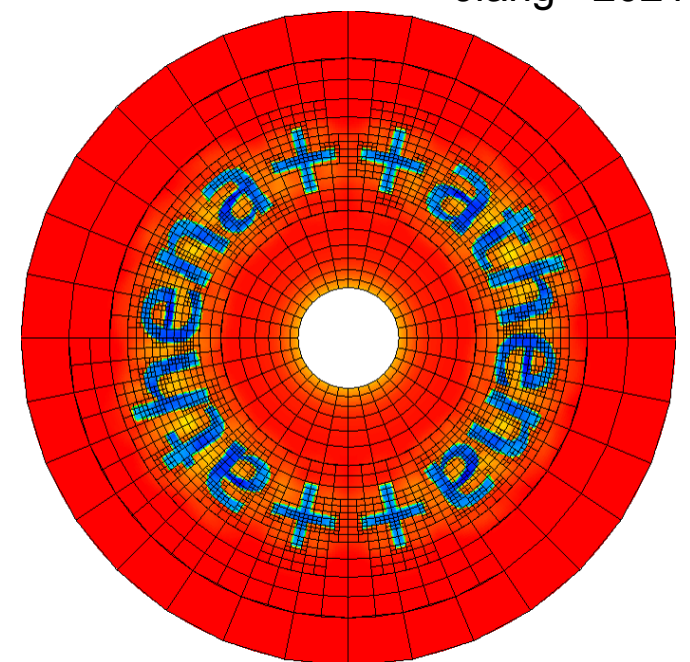
- Solve full time-dependent radiation transport equation implicitly

See Jiang+ 2021

- ▶ Using ~ 100 angles per cell
- ▶ Calculates specific intensity $\rightarrow T_{\text{rad}}, E_{\text{rad}}, F_{\text{rad}}$
 - ▶ Our models are mostly in LTE

- Dimensionless fluid numbers (e.g. Re, Pe) not explicitly specified

Jiang+ 2012
Davis+ 2012
Jiang+ 2021



Checking Assumption/Limitations

- Constant μ
 - Good for FeCZ, less so at surface/HeCZ
- Gravitational acceleration $\propto r^2$
 - $M_{\text{env}}/M_{\text{star}} < 0.2\%$
- Solar Metallicity
 - Also running lower metallicity models
- Gray OPAL opacity
 - Good in optically thick
 - LDI is important in optically thin region
- No magnetic field
 - FeCZ damped in high B field
 - Requires $>10\text{kG}$ field in CZ
- No rotation
 - Rotation period \gg eddy turnover time in FeCZ
 - (~ 3 days \gg ~ 1 hour for $35 M_{\odot}$ mid-MS)



Two Classes of 3D Models

Jiang+ 2018
Schultz+ 2020

Global

$$\phi : 0 \rightarrow \pi$$

$$\theta : -\pi/4 \rightarrow \pi/4$$

$$\Delta r_{\text{box}}/R_{\text{star}} \sim 1$$

Resolution:

$$(512, 256, 512)$$

$$(r, \theta, \phi)$$

$\sim 10^8$ CPU hours

Box

Schultz+ in press

$$\phi : -\pi/20 \rightarrow \pi/20$$

$$\theta : -\pi/20 \rightarrow \pi/20$$

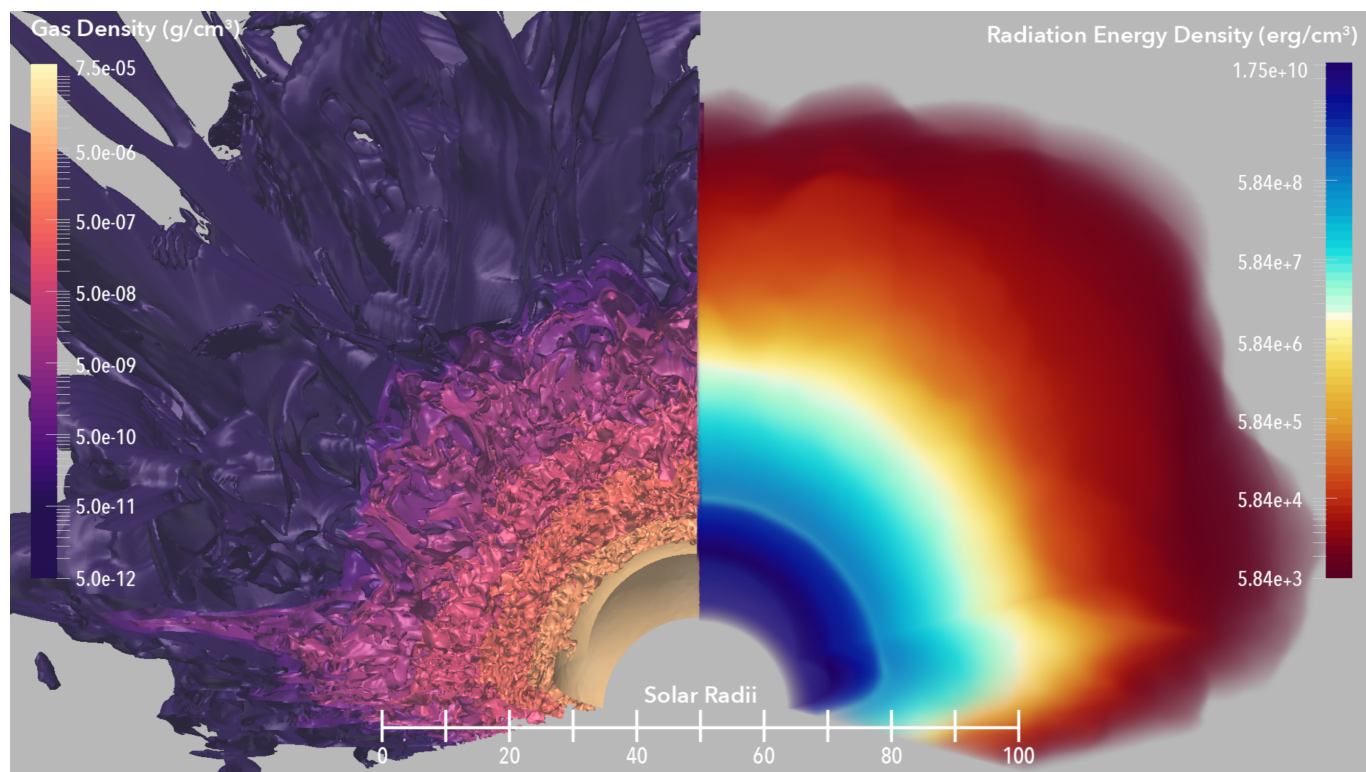
$$\Delta r_{\text{box}}/R_{\text{star}} \sim 0.3$$

Resolution:

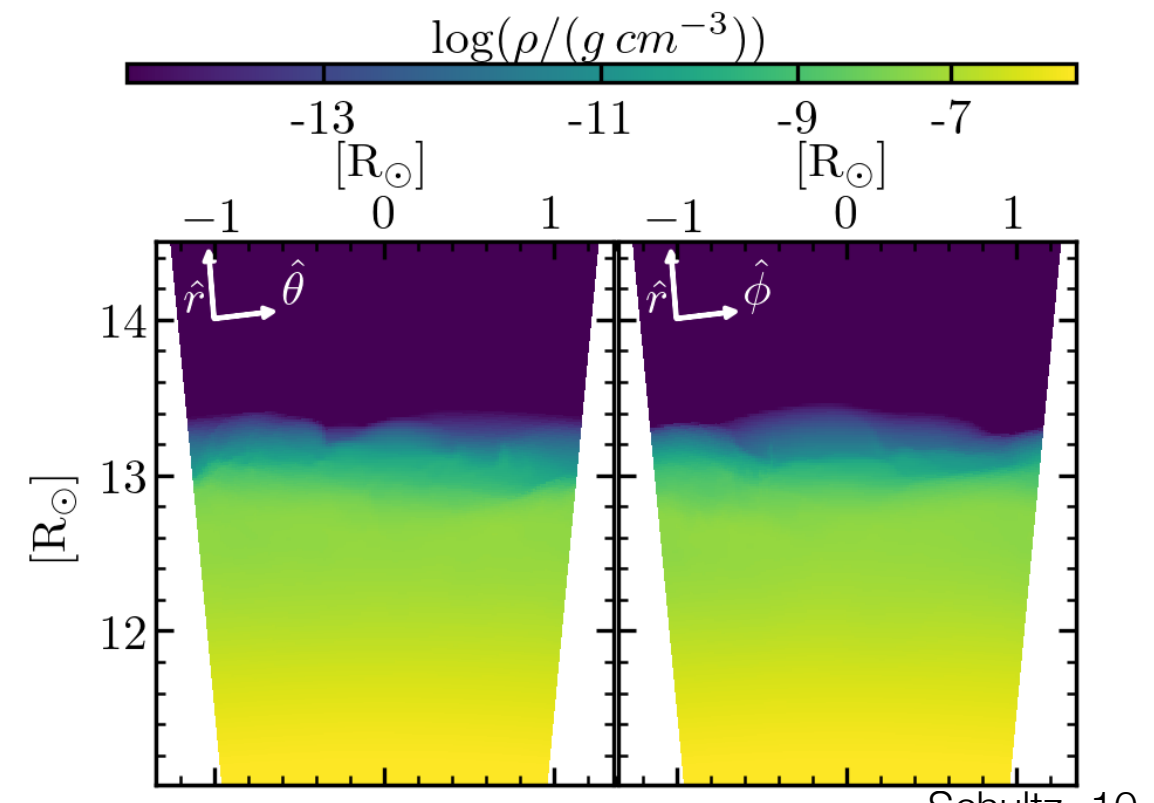
$$(265, 256, 256)$$

$$(r, \theta, \phi)$$

$\sim 10^6$ CPU hours

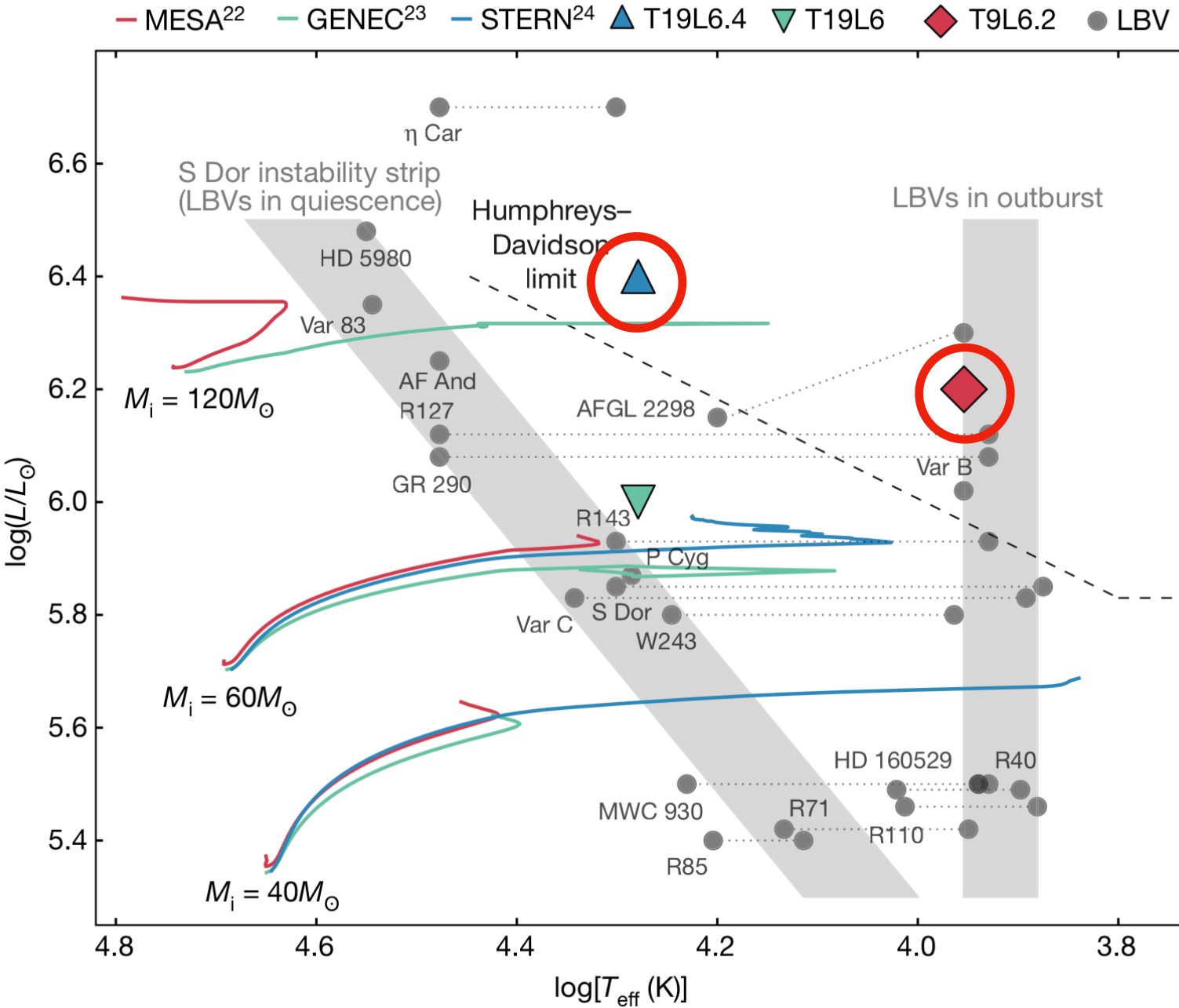


Jiang+ 2018



Schultz 10

Global Models



T19L6.4
80 solar masses
~100 solar radii

$$\tau_{\text{Fe}} \approx 3\tau_{\text{crit}}$$

T9L6.2
56 solar masses
~300 solar radii

$$\tau_{\text{Fe}} \approx 10\tau_{\text{crit}}$$

Jiang+ 2018

Table 1: Properties of the 3D Stellar Models

Name	M_{core}	M_{env}	T_{eff}	L	L_{Edd}^1	r_{base}	r_{max}	r_{ph}^2	r_{Fe}	τ_{Fe}	Z
	(M_{\odot})	(M_{\odot})	(10^3 K)	($\log(L/L_{\odot})$)	($\log(L/L_{\odot})$)	(R_{\odot})	(R_{\odot})	(R_{\odot})	(R_{\odot})		(Z_{\odot})
T9L6.2	56	0.13	9	6.20	6.26	35.0	809.8	353.3	80.3	28,000	1
T19L6.4	80	0.011	19	6.40	6.42	16.3	335.5	99.0	44.0	5,400	1

¹ For an assumed electron scattering opacity.

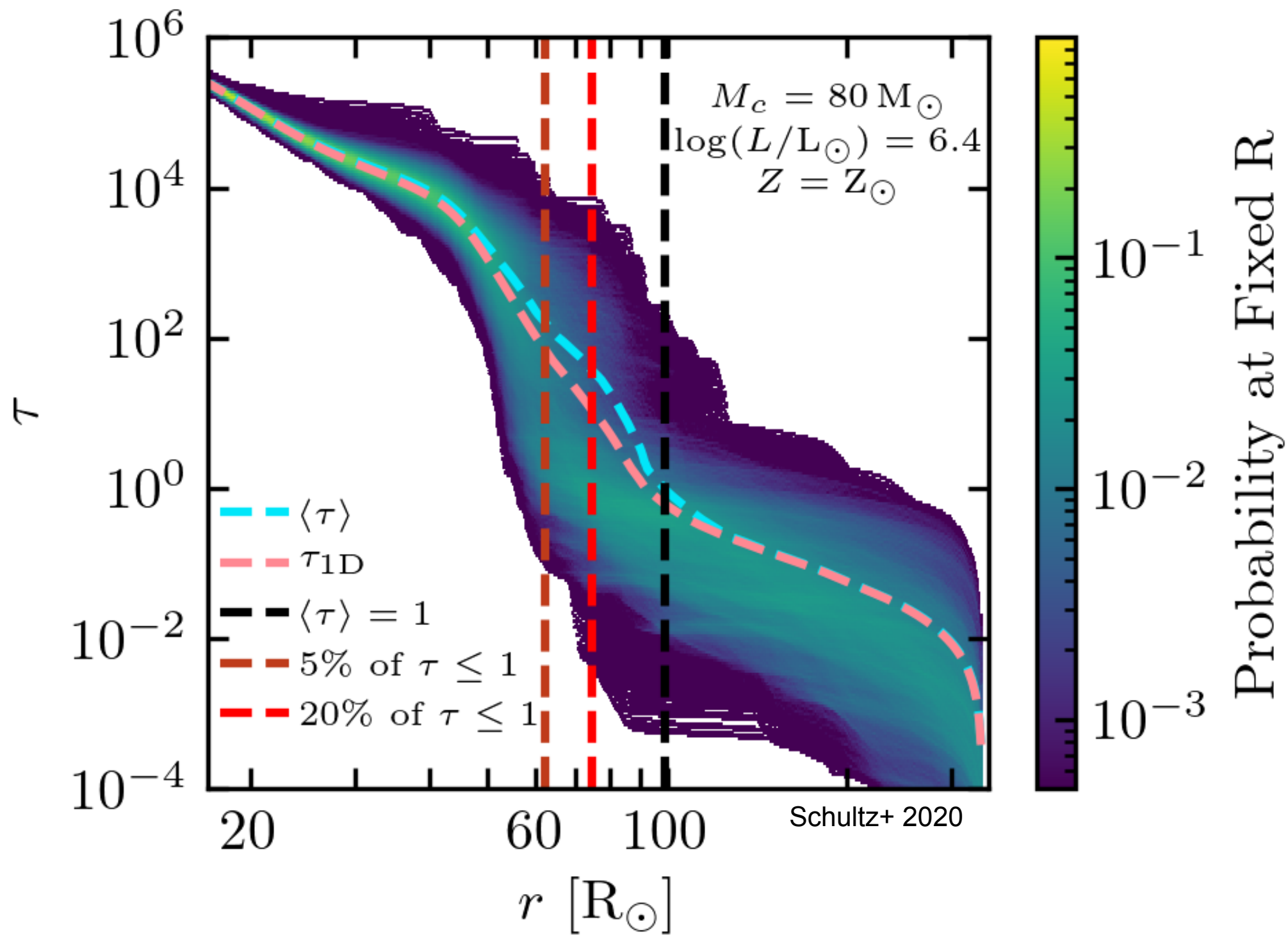
² The photosphere radii specified are where $\langle \tau(r_{\text{ph}}) \rangle = 1$.

Modified from
 Schultz+ 2020

T19L6.4



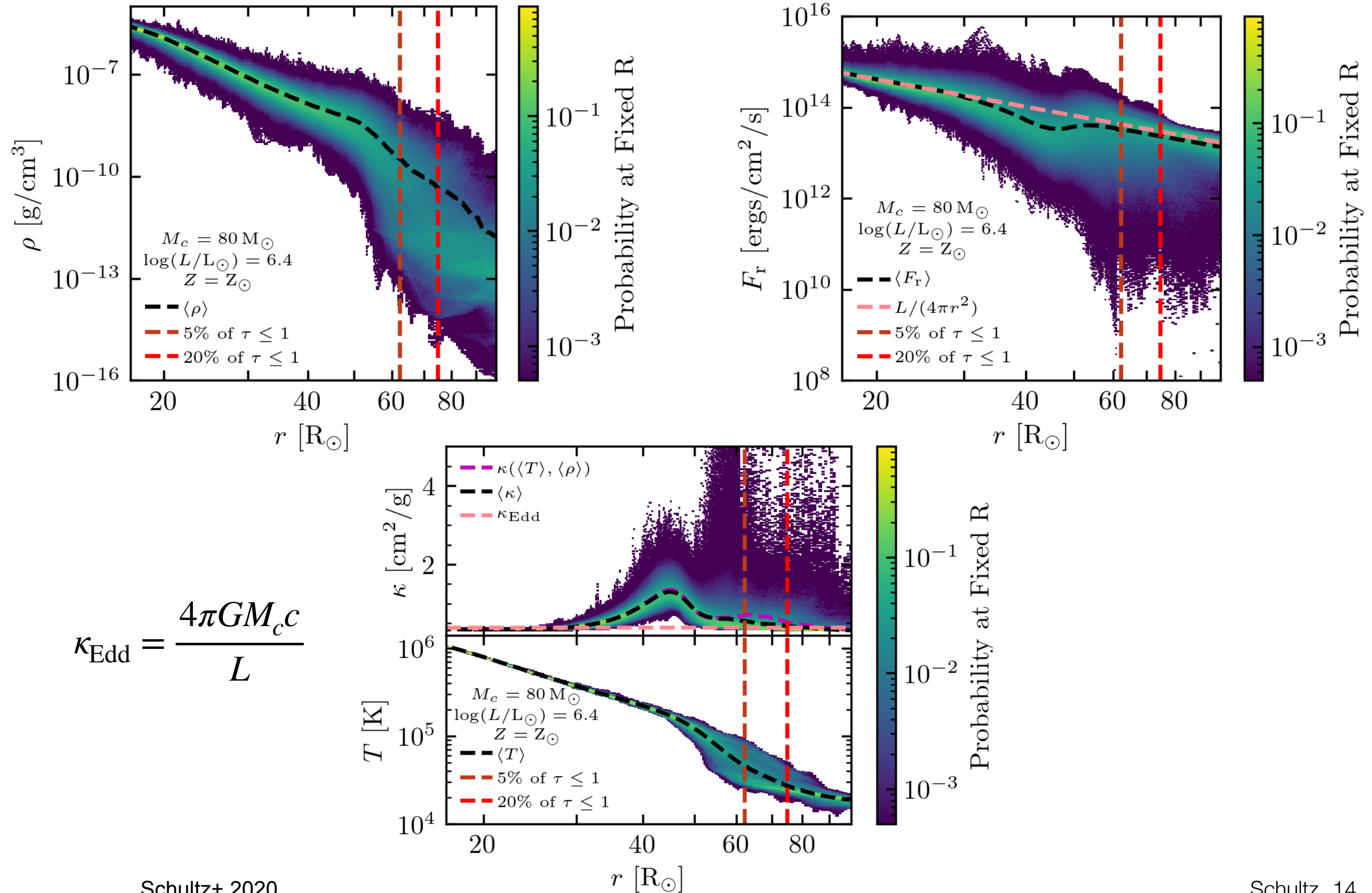
T19L6.4: Optical Depth Profile



$$\tau_{1D}(r) = \int_r^{r_{max}} \langle \kappa(r') \rangle \langle \rho(r') \rangle dr'$$

* τ is calculated along radial lines of sight only

T19L6.4: Spatial Variability



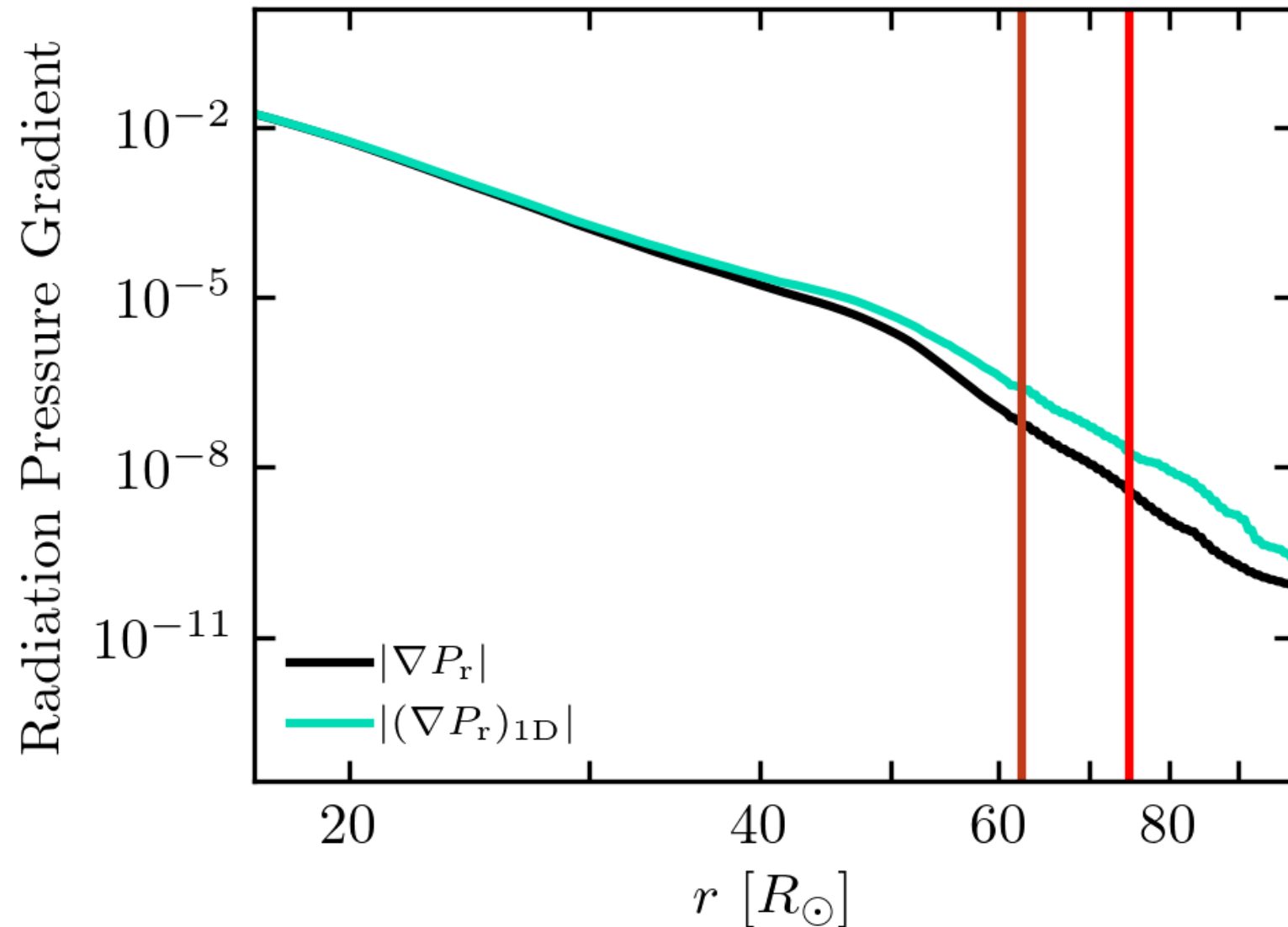
Correlations Impact ∇P_r

$$\nabla P_r = \frac{-1}{c} \langle \kappa F_r \rho \rangle$$

$$(\nabla P_r)_{1D} = \frac{-1}{c} \langle \kappa \rangle \langle F_r \rangle \langle \rho \rangle$$

From spatial variability,

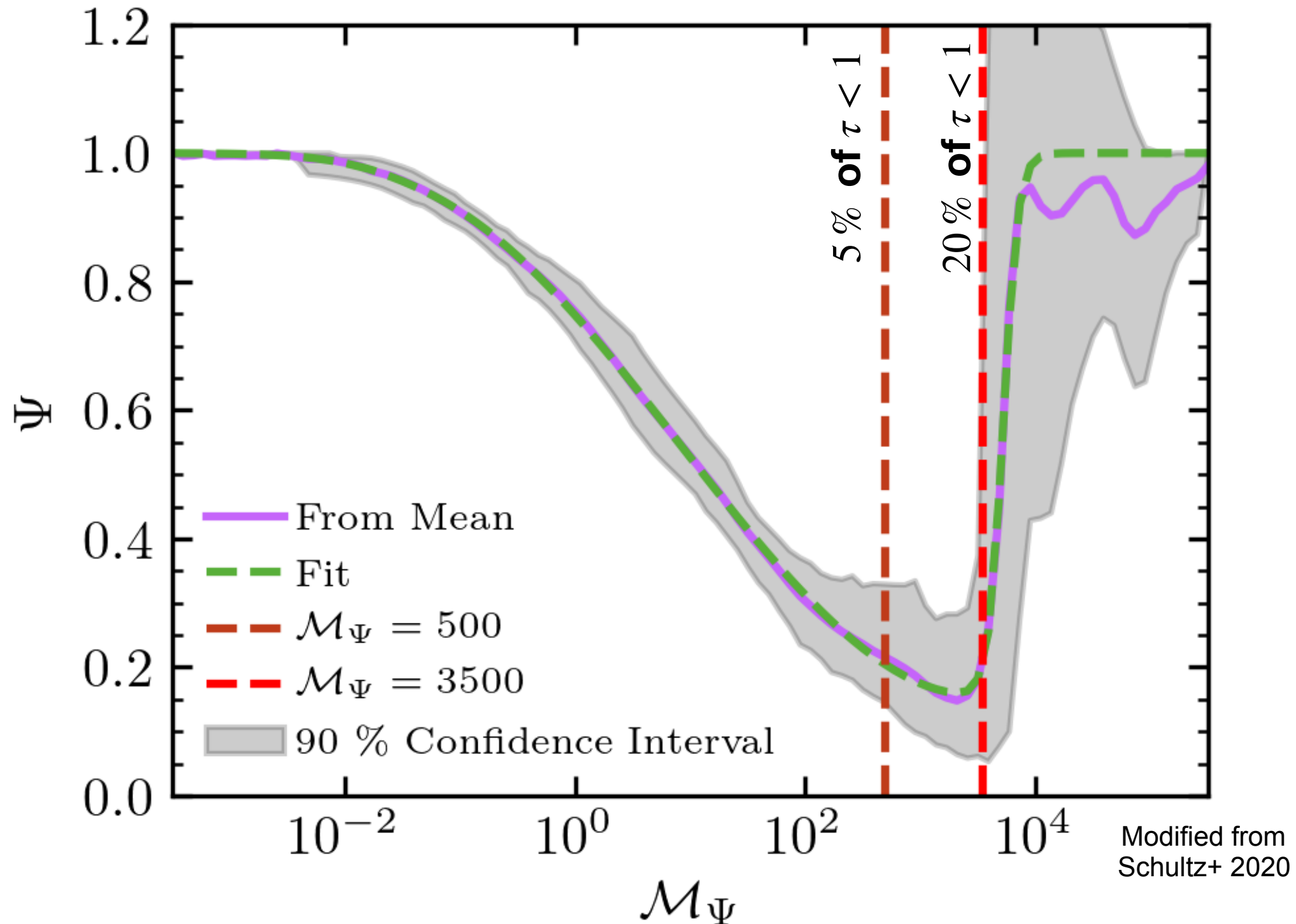
$$\nabla P_r \approx \frac{-1}{c} \langle \kappa \rangle \langle F_r \rho \rangle$$



Define a quantity Ψ such that $\nabla P_r \approx (\nabla P_r)_{1D} \Psi$

$$\Psi \equiv \frac{\langle F_r \rho \rangle}{\langle F_r \rangle \langle \rho \rangle}$$

Ψ in T19L6.4 and T9L6.2



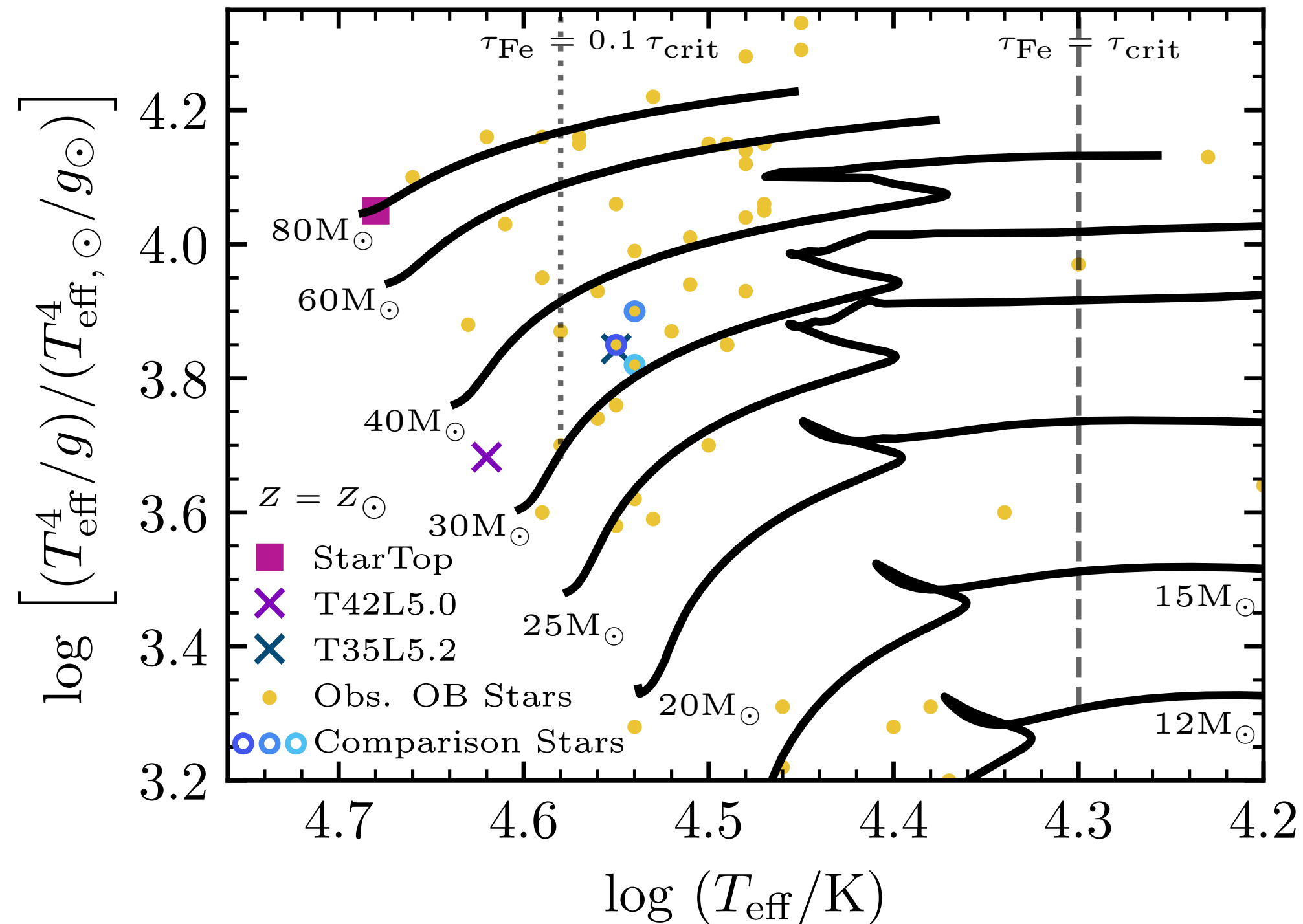
$\mathcal{M}_\Psi = f(L, T, r)$ is analogous to radius

See Owocki & Sundqvist 2018
and references there-in
for more on porosity

- Porosity is important near the surface of massive stars
- Decreased pressure support counteracted by turbulent pressure

Box Models

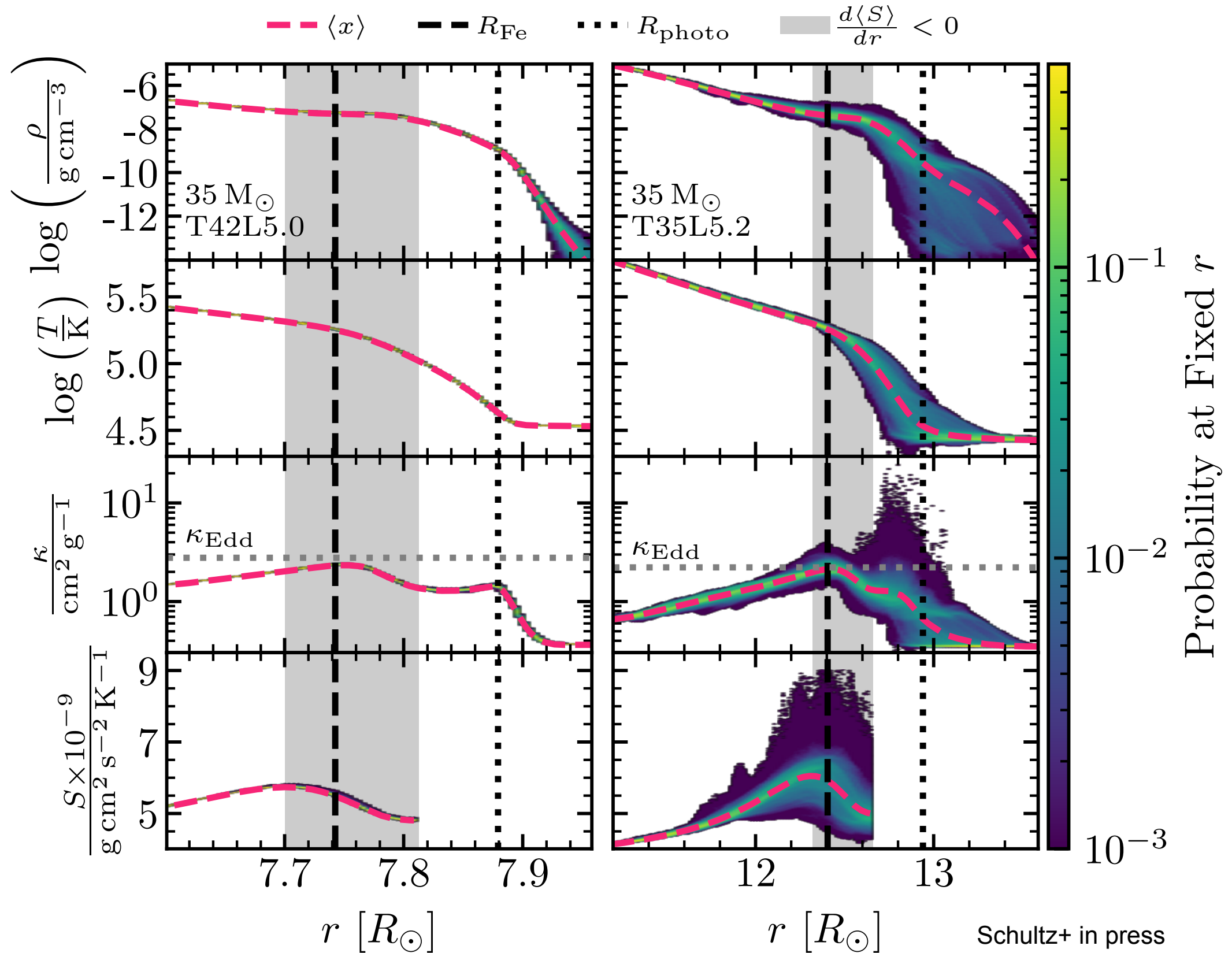
Schultz+ in press



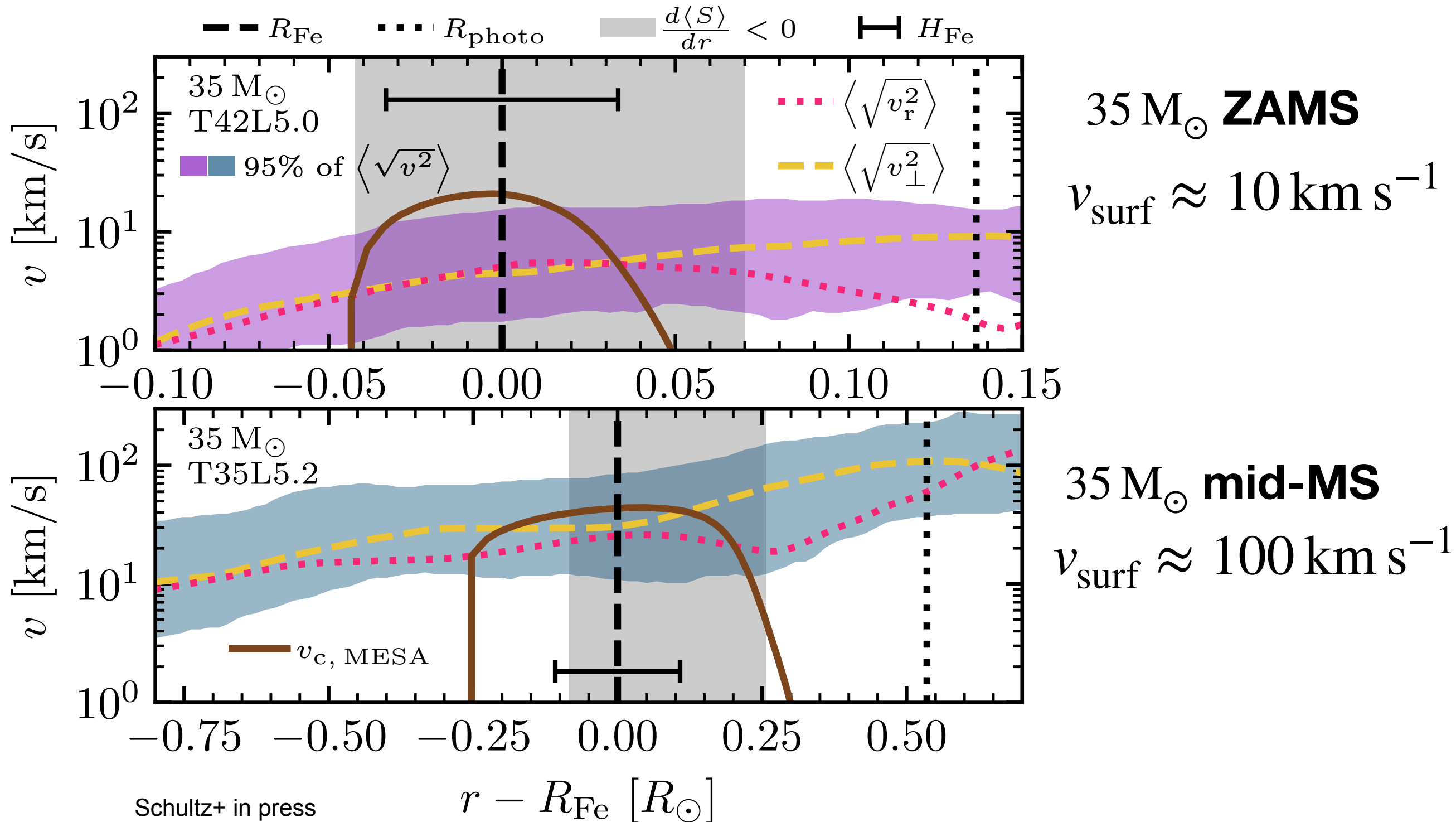
T42L5.0 (ZAMS)
35 solar masses
~8 solar radii
 $\tau_{\text{Fe}} \approx 0.01 \tau_{\text{crit}}$

T35L5.2 (mid-MS)
35 solar masses
~13 solar radii
 $\tau_{\text{Fe}} \approx 0.2 \tau_{\text{crit}}$

Top Model Spatial Variability

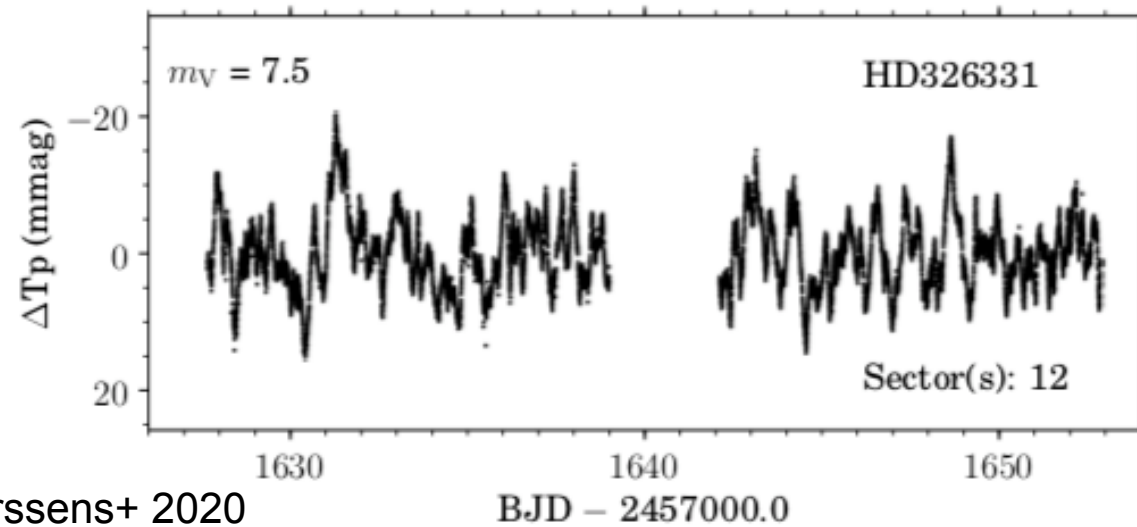


O Star Turbulent Surface Velocities

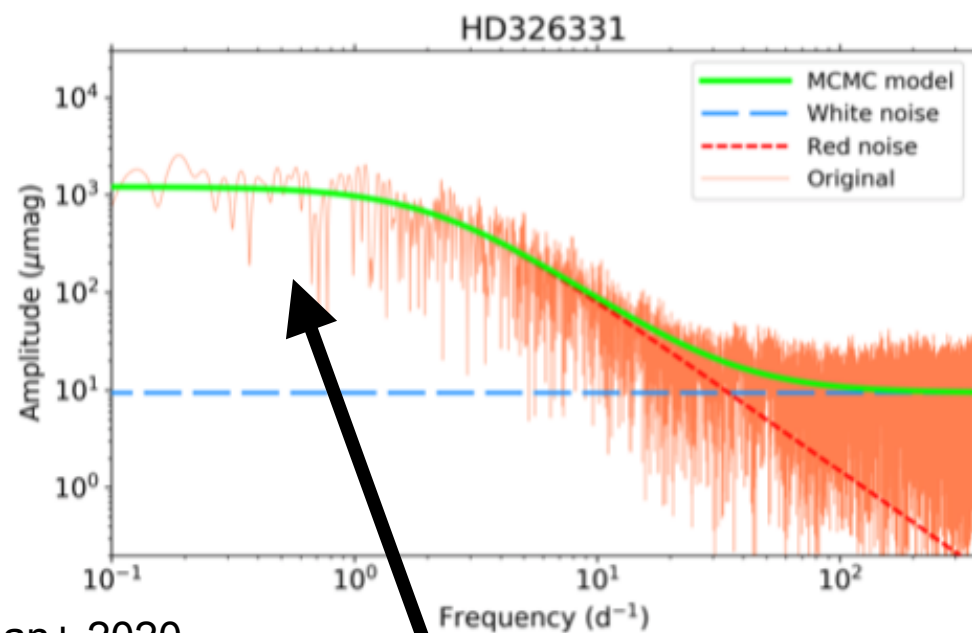


- Velocities not confined sub-surface convection zone in 3D
- Convective fluxes carry small amount of energy

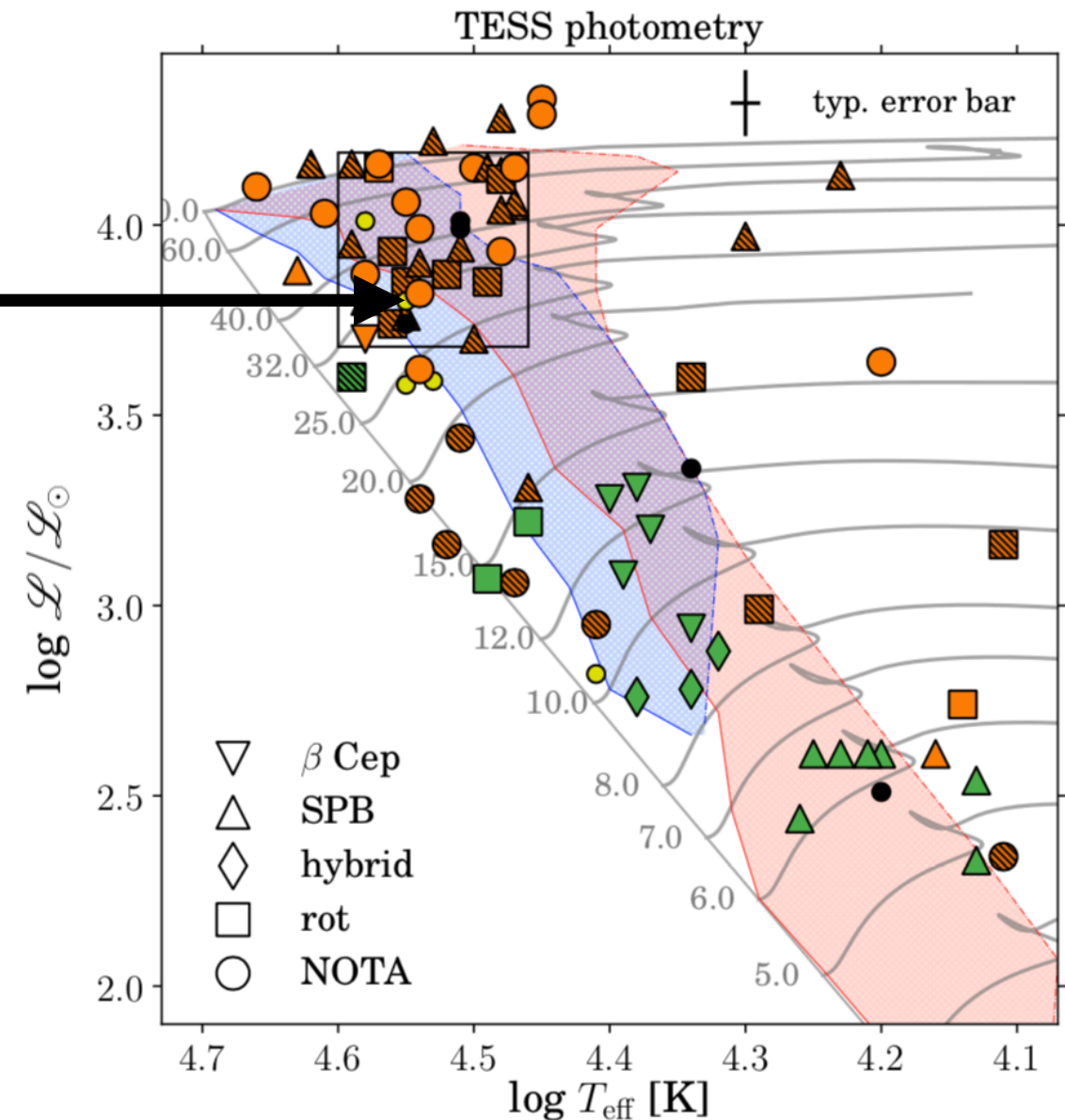
TESS Characterizations of Massive Stars



Burssens+ 2020



Bowman+ 2020



Burssens+ 2020

Stochastic Low-Frequency Variability: 3 possible causes

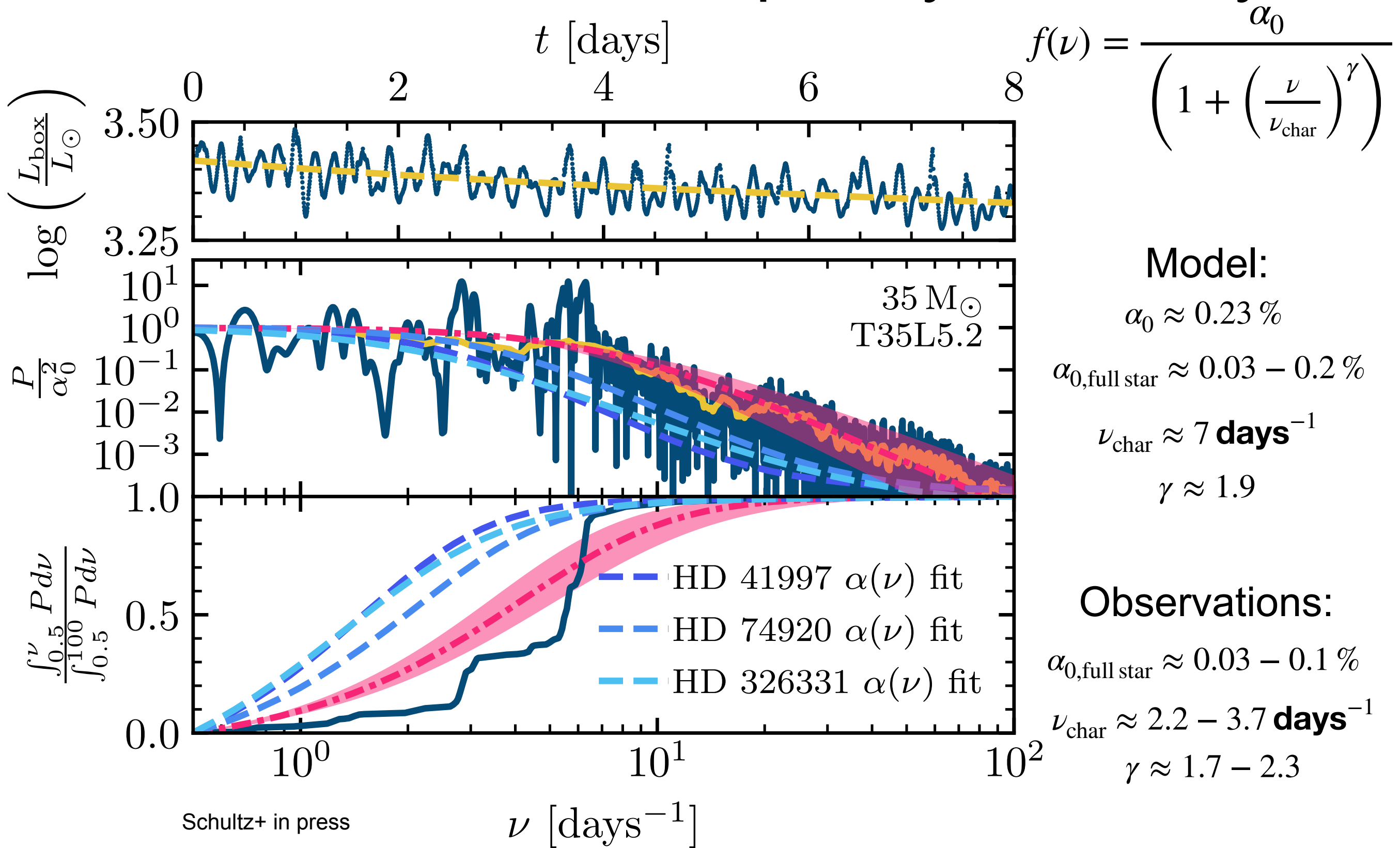
1. Internal gravity waves
2. Near surface convective zones
3. Wind interactions

e.g. Aerts+ 2009, Simón-Días & Herrero 2010, 2014, 2017, Edelmann+ 2019, Bowman+ 2020

e.g. Lecoanet & Quataert 2013, Grassitelli+ 2015, Cantiello+ 2021, Schultz + in press

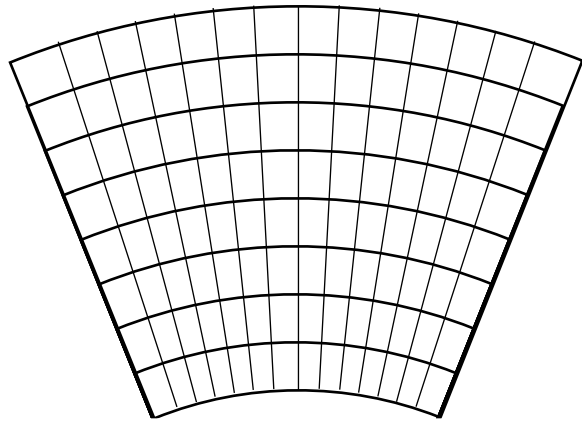
e.g. Moffat+ 2008, David-Uraz+ 2017, Aerts+ 2018, Simón-Días+ 2018, Krlicka & Feldmeier 2018,

Stochastic Low-Frequency Variability



Brief Resolution Study

T35L5.2



ν_{surf} **characteristics**

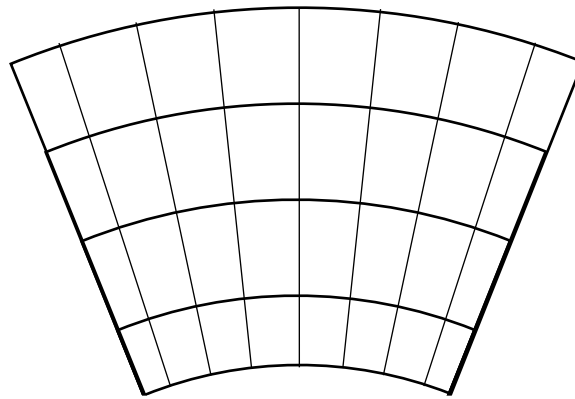
QPO frequencies

$$\alpha_0 \approx 0.23 \%$$

$$\nu_{\text{char}} \approx 7 \text{ days}^{-1}$$

$$\gamma \approx 1.9$$

Lower Resolution

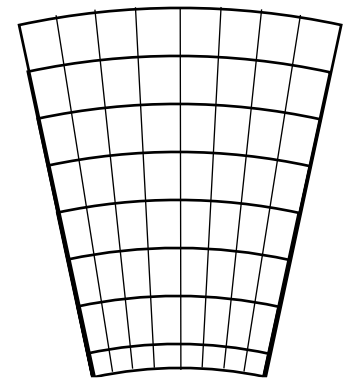


$$\alpha_0 \approx 0.25 \%$$

$$\nu_{\text{char}} \approx 7 \text{ days}^{-1}$$

$$\gamma \approx 1.9$$

Smaller Box



$$\alpha_0 \approx 0.16 \%$$

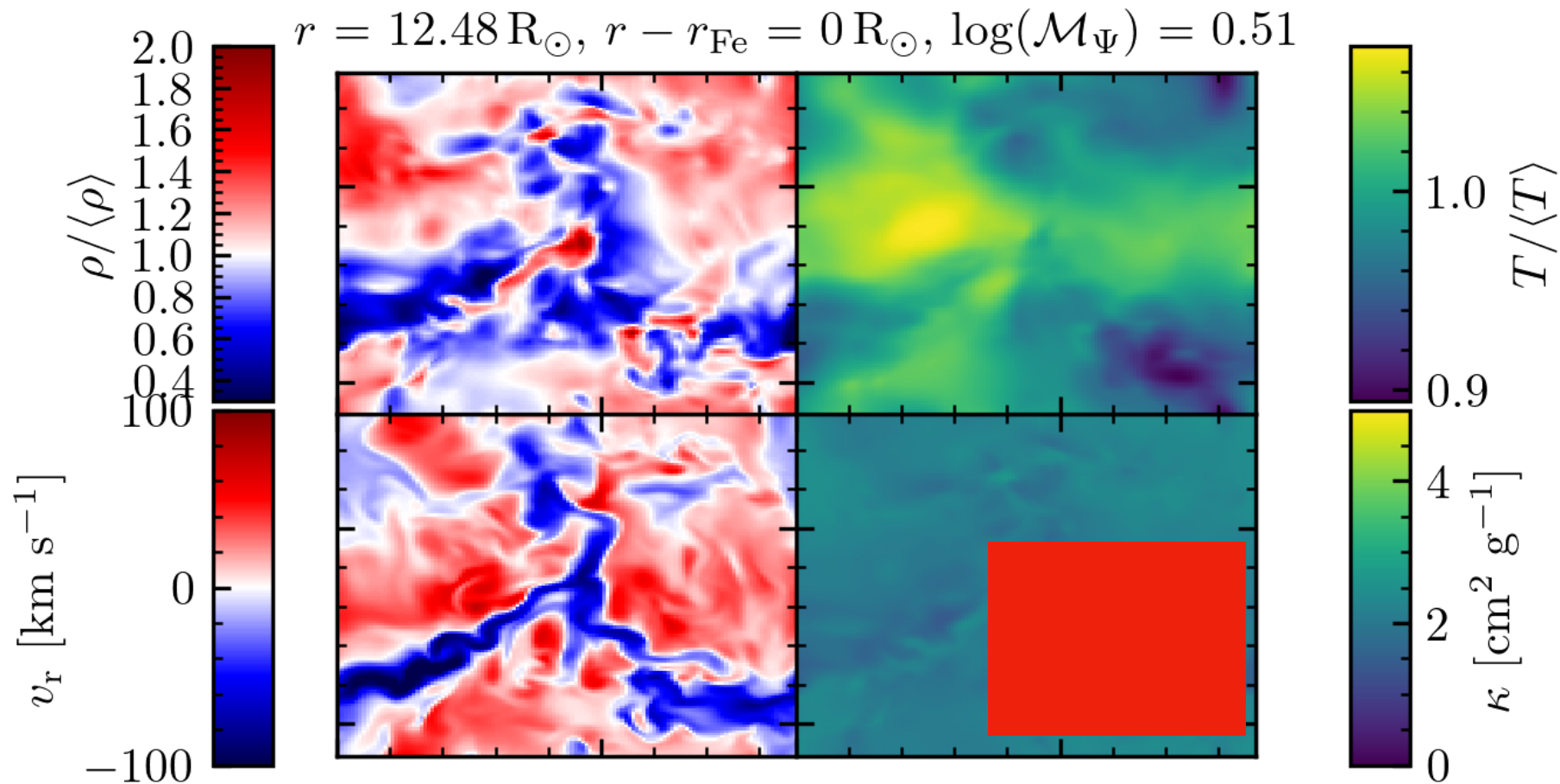
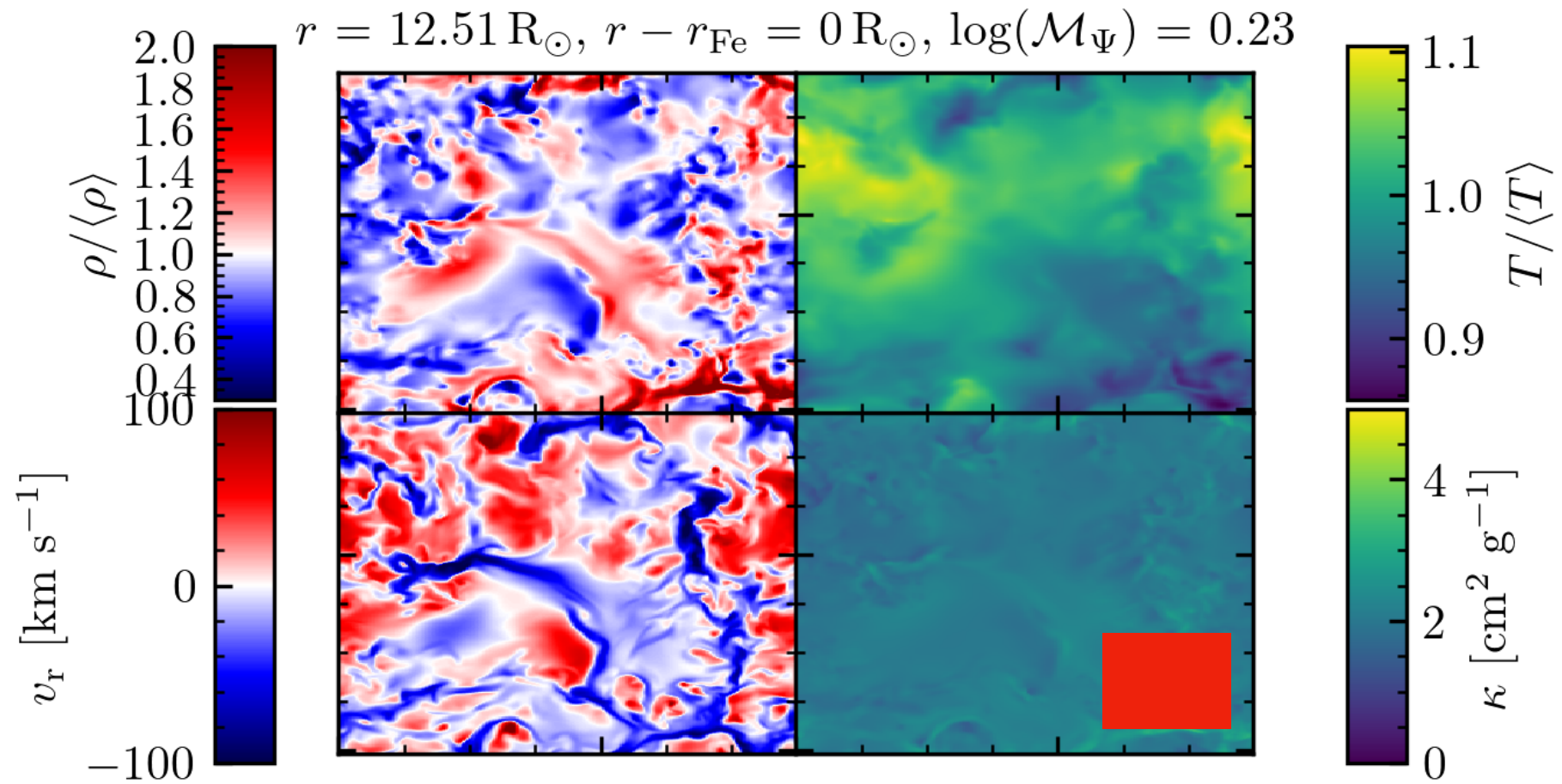
$$\nu_{\text{char}} \approx 10 \text{ days}^{-1}$$

$$\gamma \approx 1.9$$

Model Cross-Sections

T35L5.2

Many plumes
Can be extrapolated to
full star



Smaller Box
Few plumes
Plumes are correlated so
cannot extrapolate

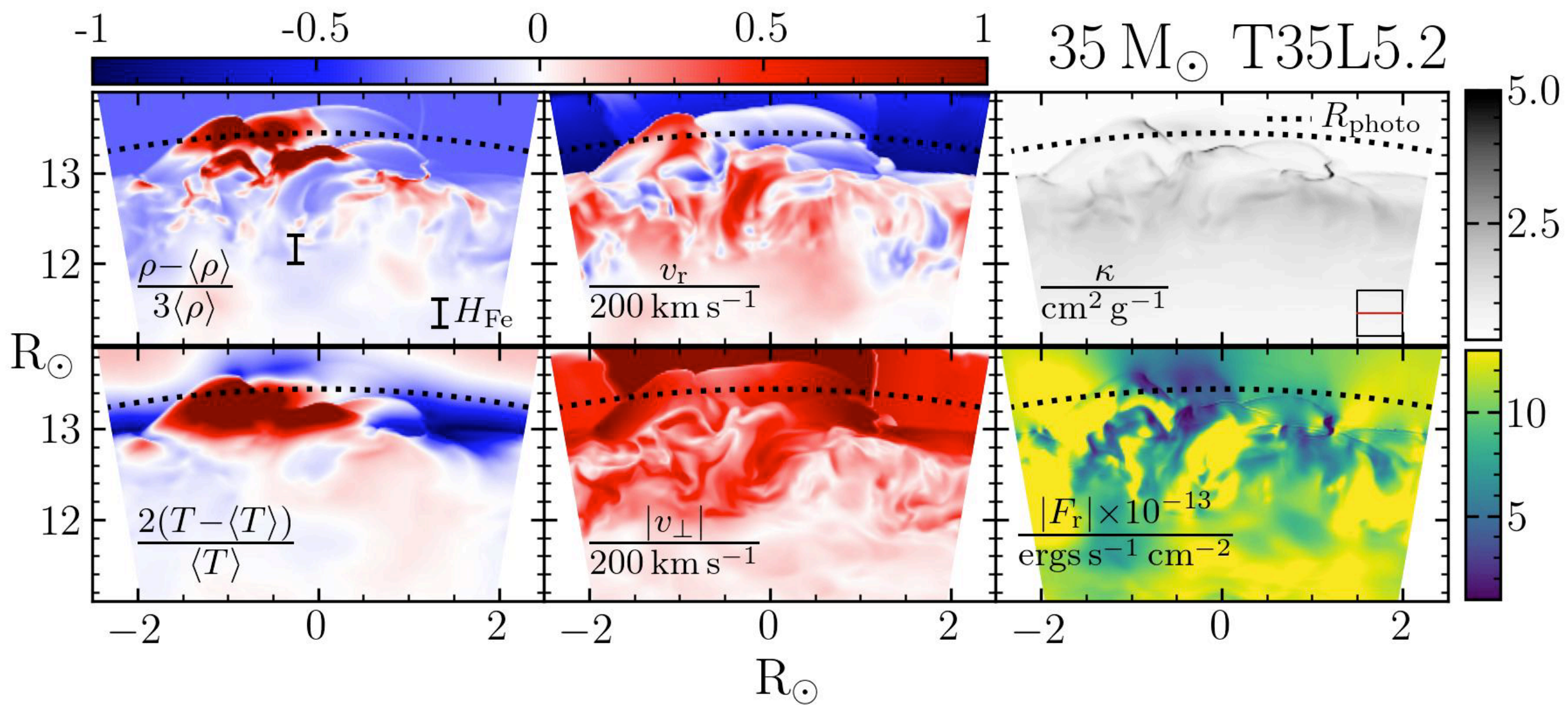
Main Takeaways

- Porosity/clumping is important at the surface of massive stars
- Turbulent motion from near-surface CZs persists to the surface of massive stars at $|v_{\text{surf}}| \sim 100 \text{ km s}^{-1}$
- Near-surface CZs can induce SLFV resembling recent TESS observations

Future Outlook

- More models running
 - ▶ Mid-MS $35 M_{\odot}$ with lower metallicity
 - ▶ $13 M_{\odot}$ TAMS model
- Investigate line effects in opacity
- How would turbulent pressure affect 1D model evolution?

Thank you!



Backup

Choosing Model Independent Variable

For radiation pressure dominated system, we chose

$$F_{c, \text{Rad dominant}} = (E_r) v_c$$

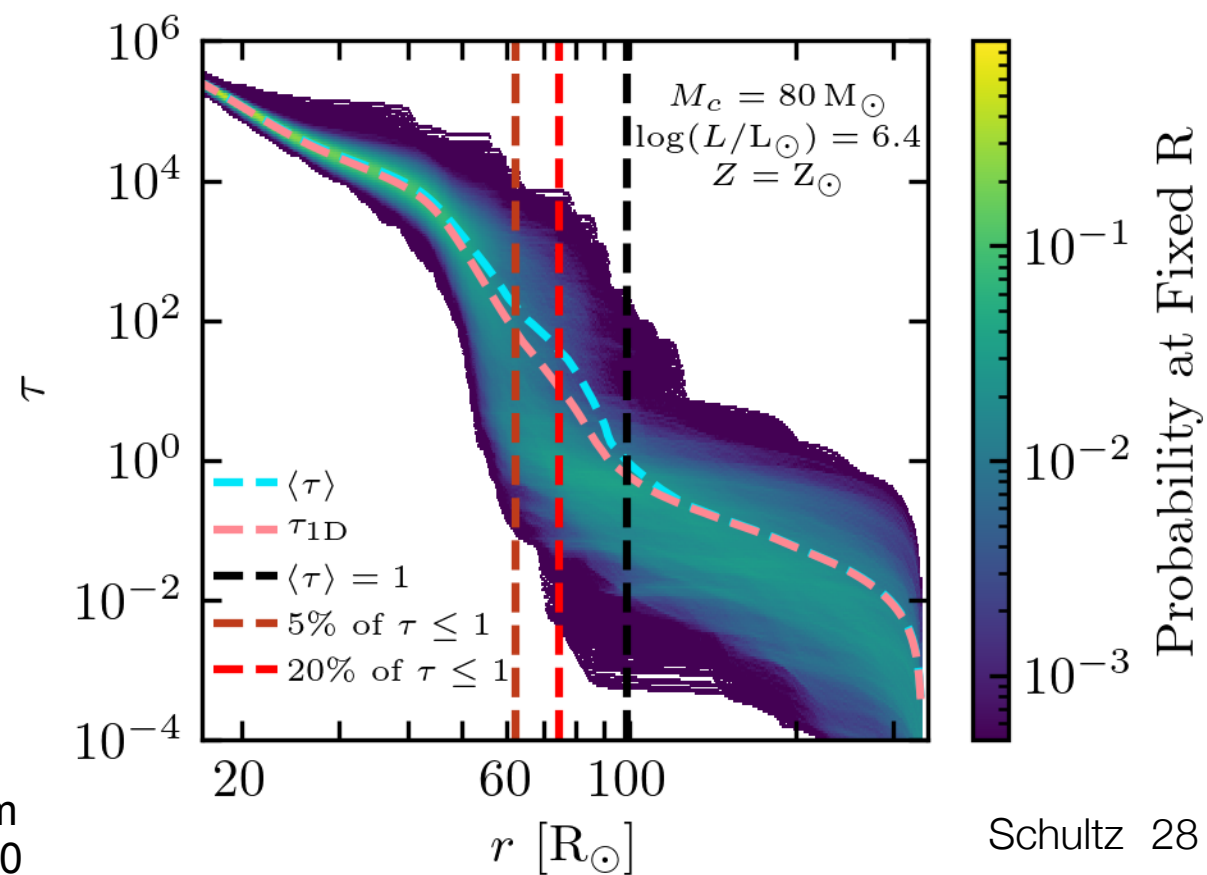
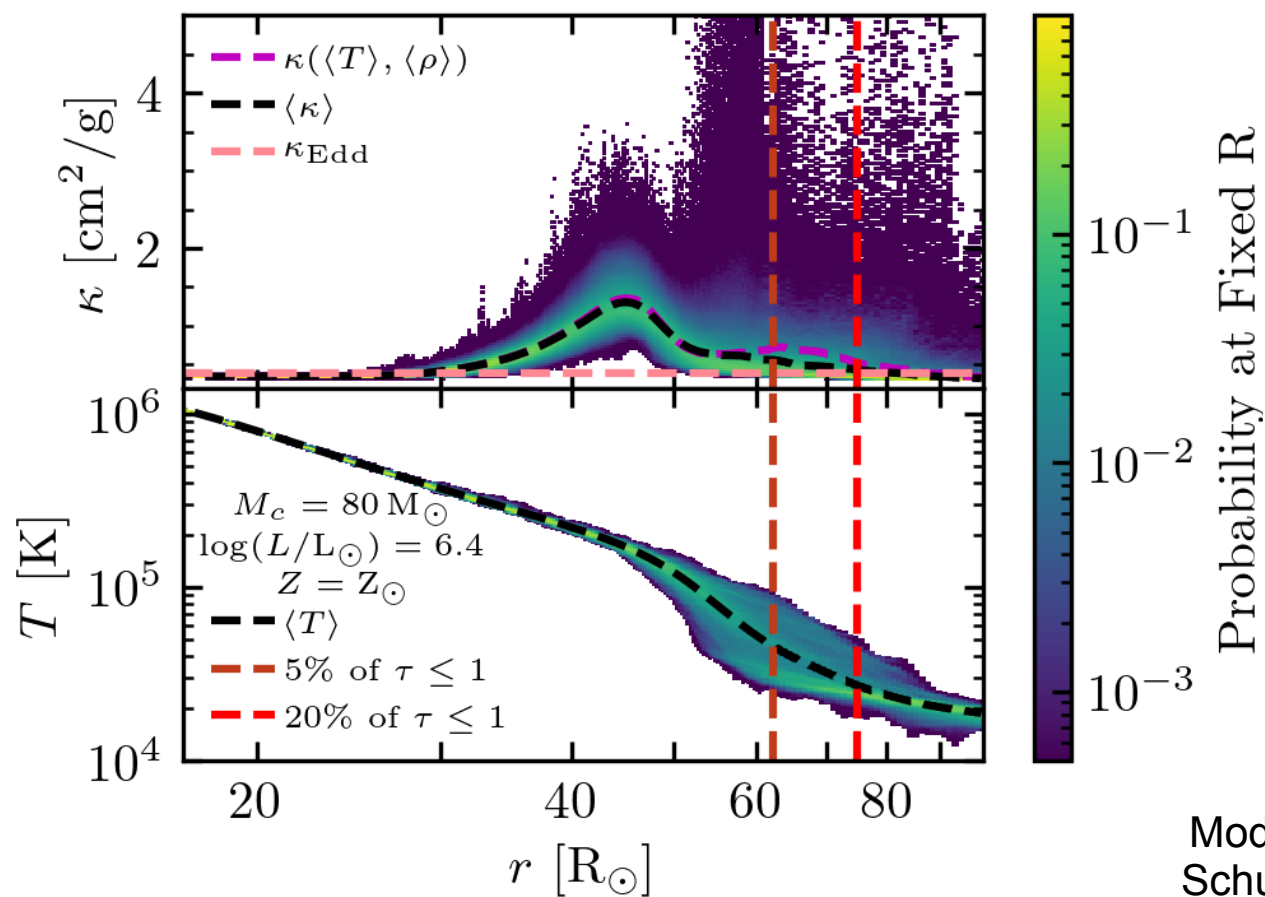
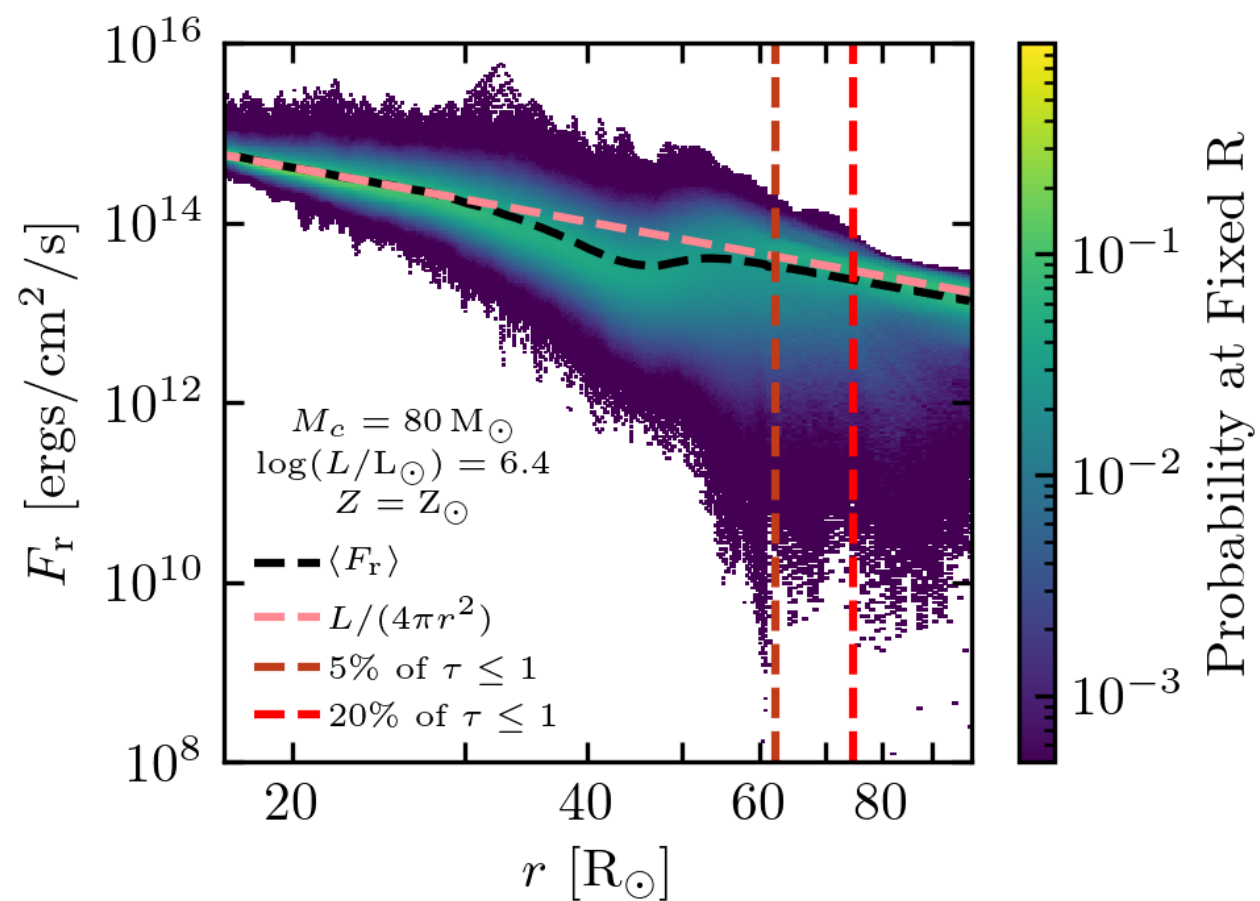
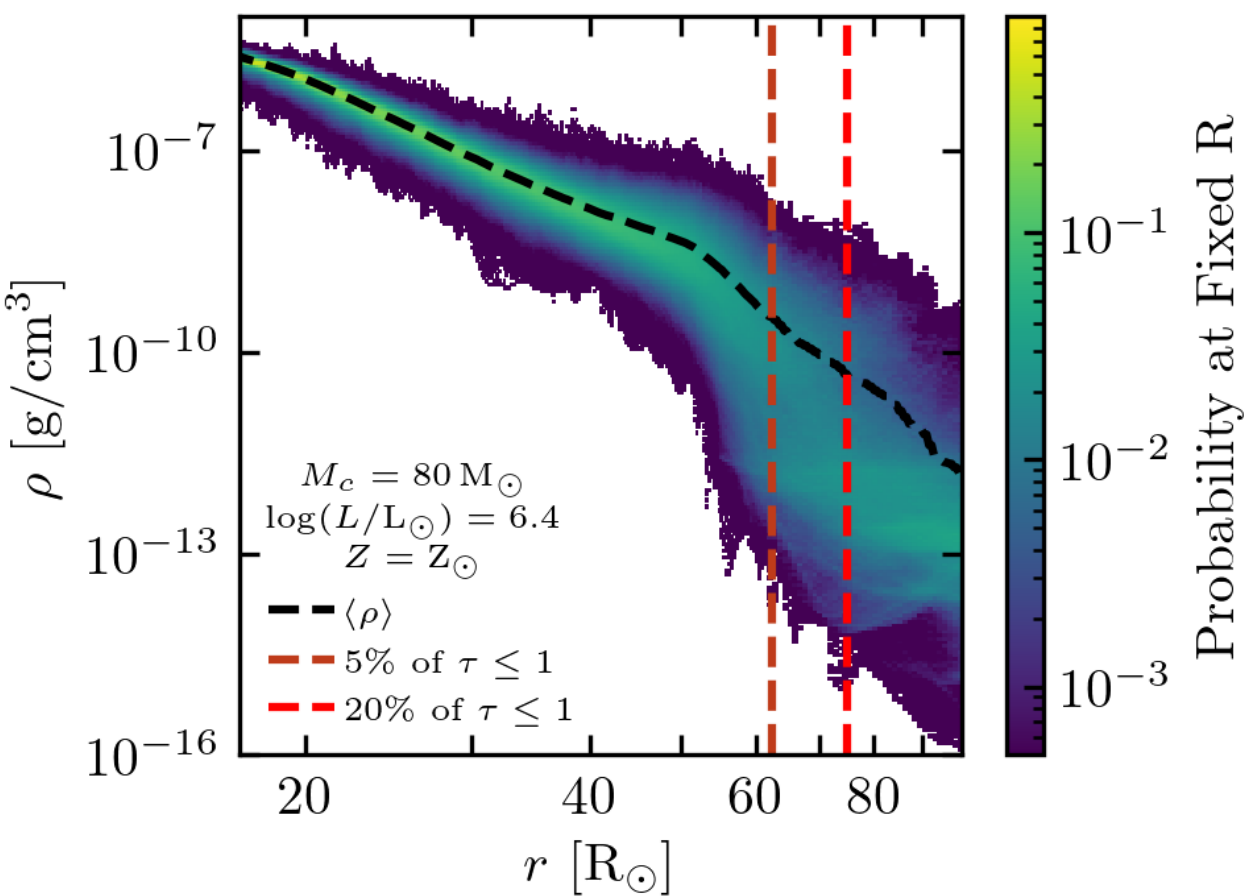
We then choose the velocity needed for this flux to carry the entire luminosity

$$\frac{L}{4\pi r^2} \equiv E_r v_L \quad \text{or} \quad v_L \equiv \frac{L}{4\pi r^2 E_r} = \frac{L}{4\pi r^2 a T^4}$$

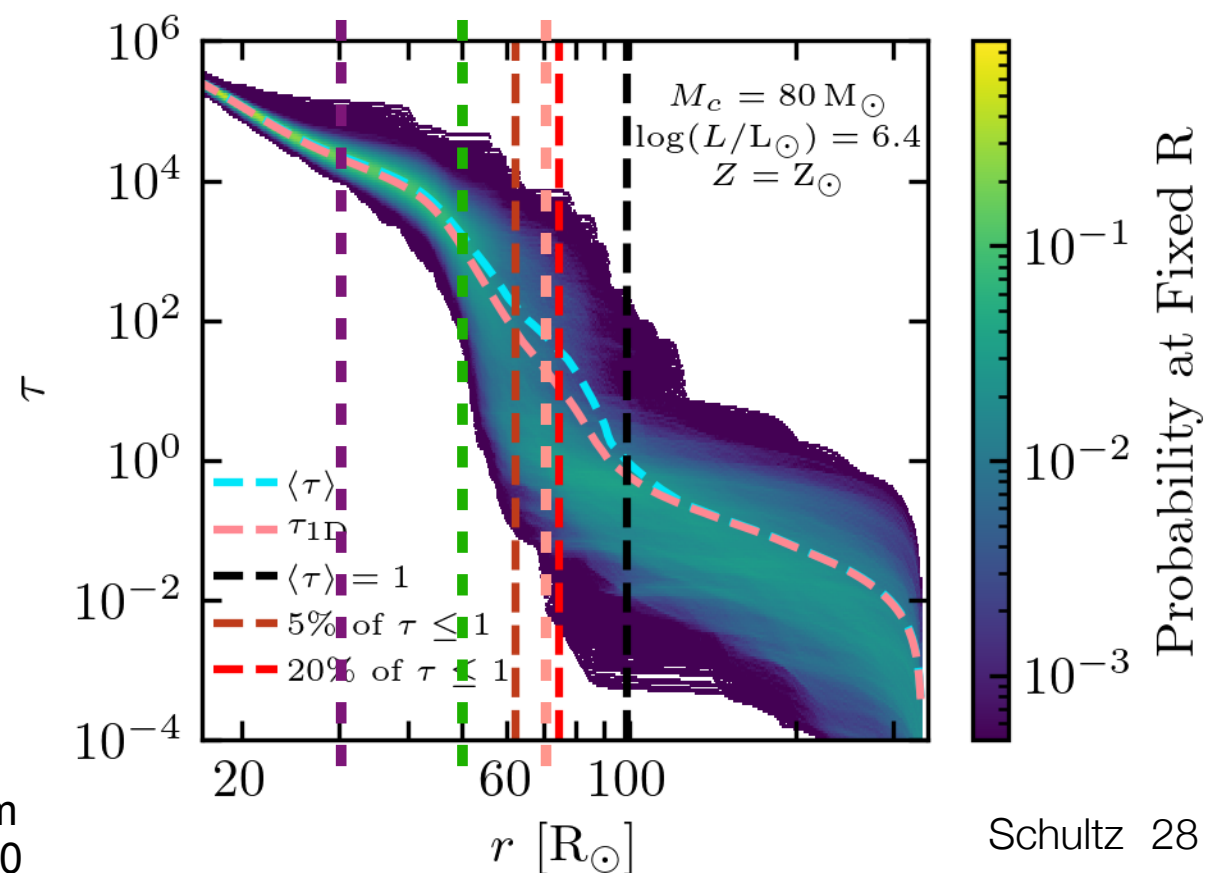
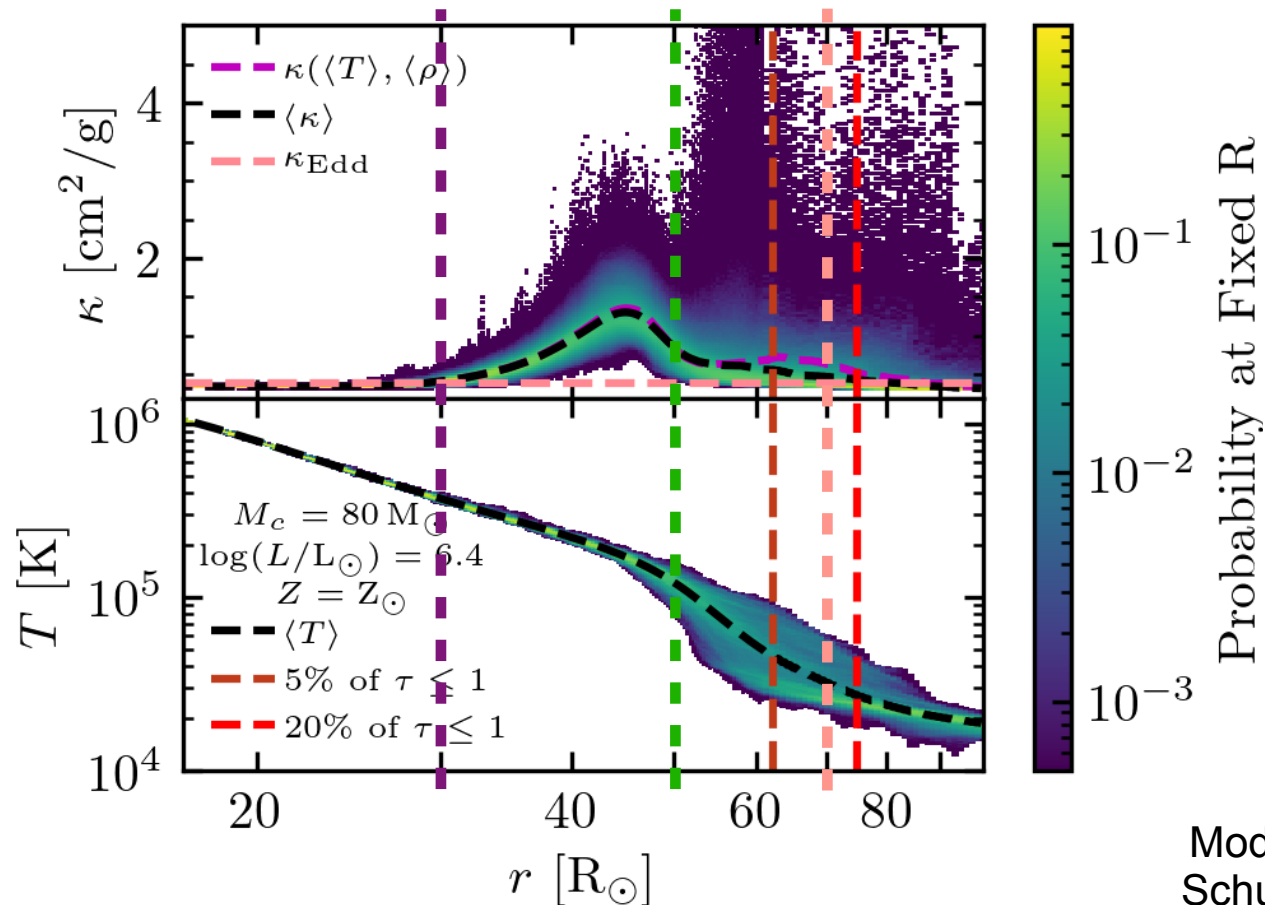
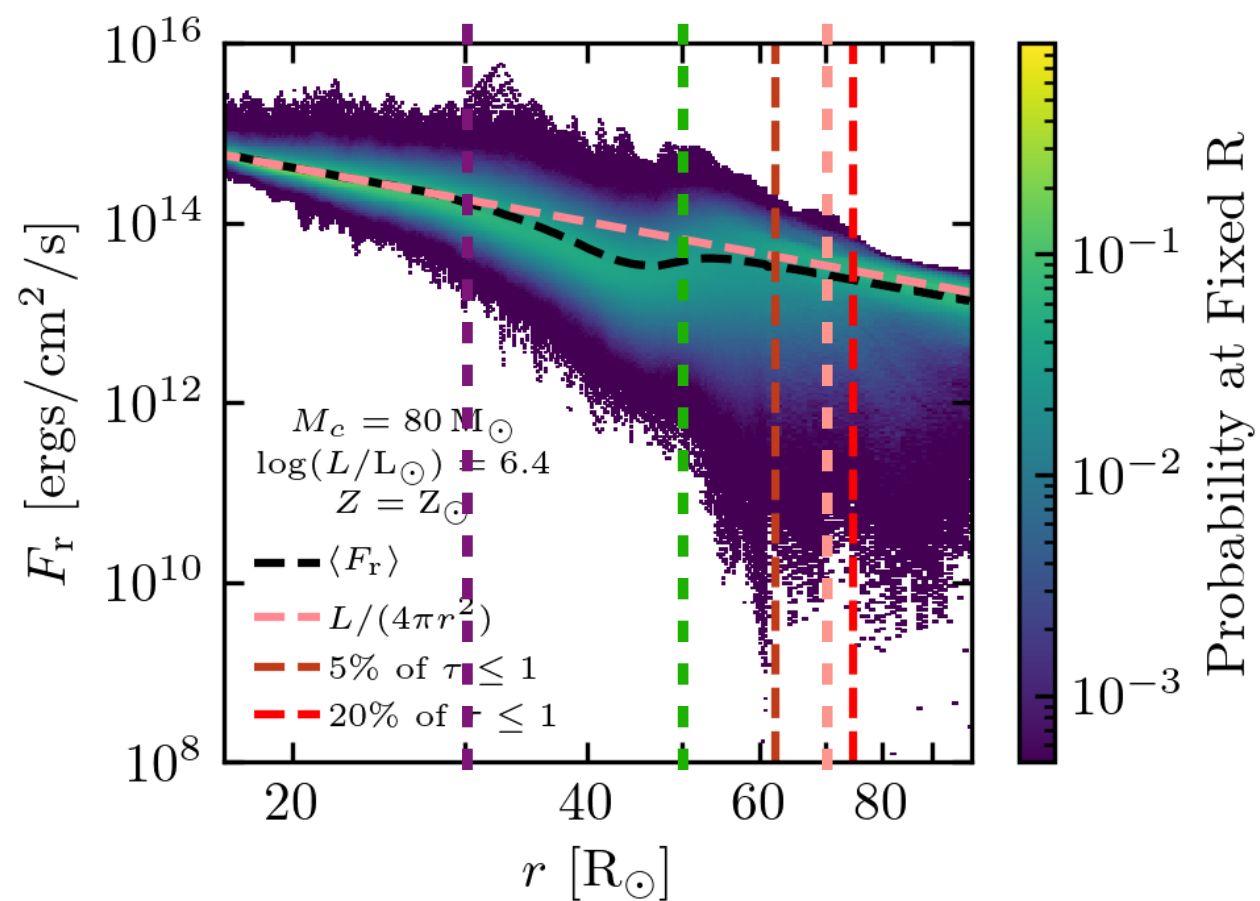
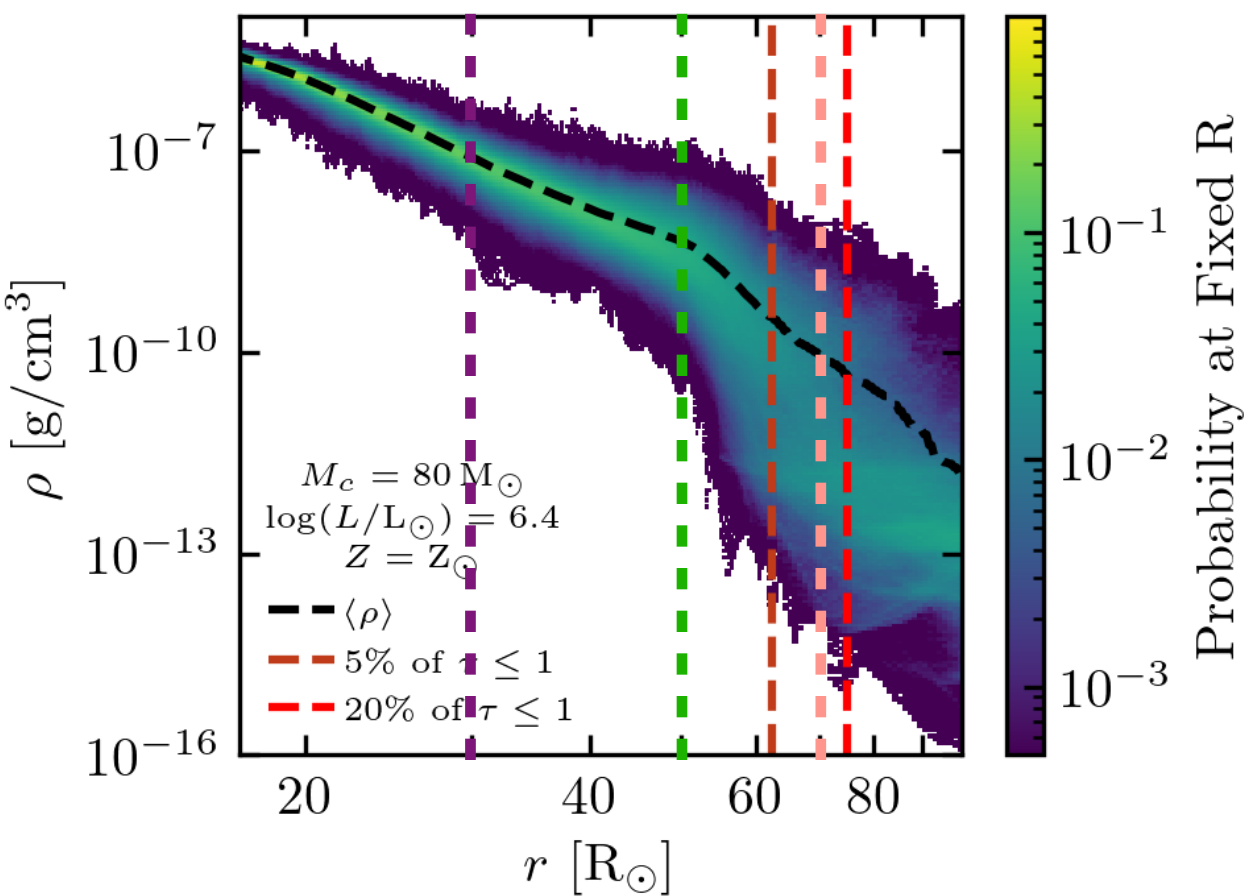
Dividing by the isothermal gas sound speed, we arrive at a pseudo-Mach number

$$\mathcal{M}_\Psi(r) \equiv \frac{v_L(r)}{\sqrt{\frac{k_B T(r)}{\mu m_p}}} = \frac{L}{4\pi r^2 a T^{4.5}(r)} \left(\frac{\mu m_p}{k_B} \right)^{1/2}$$

Are There Correlations?

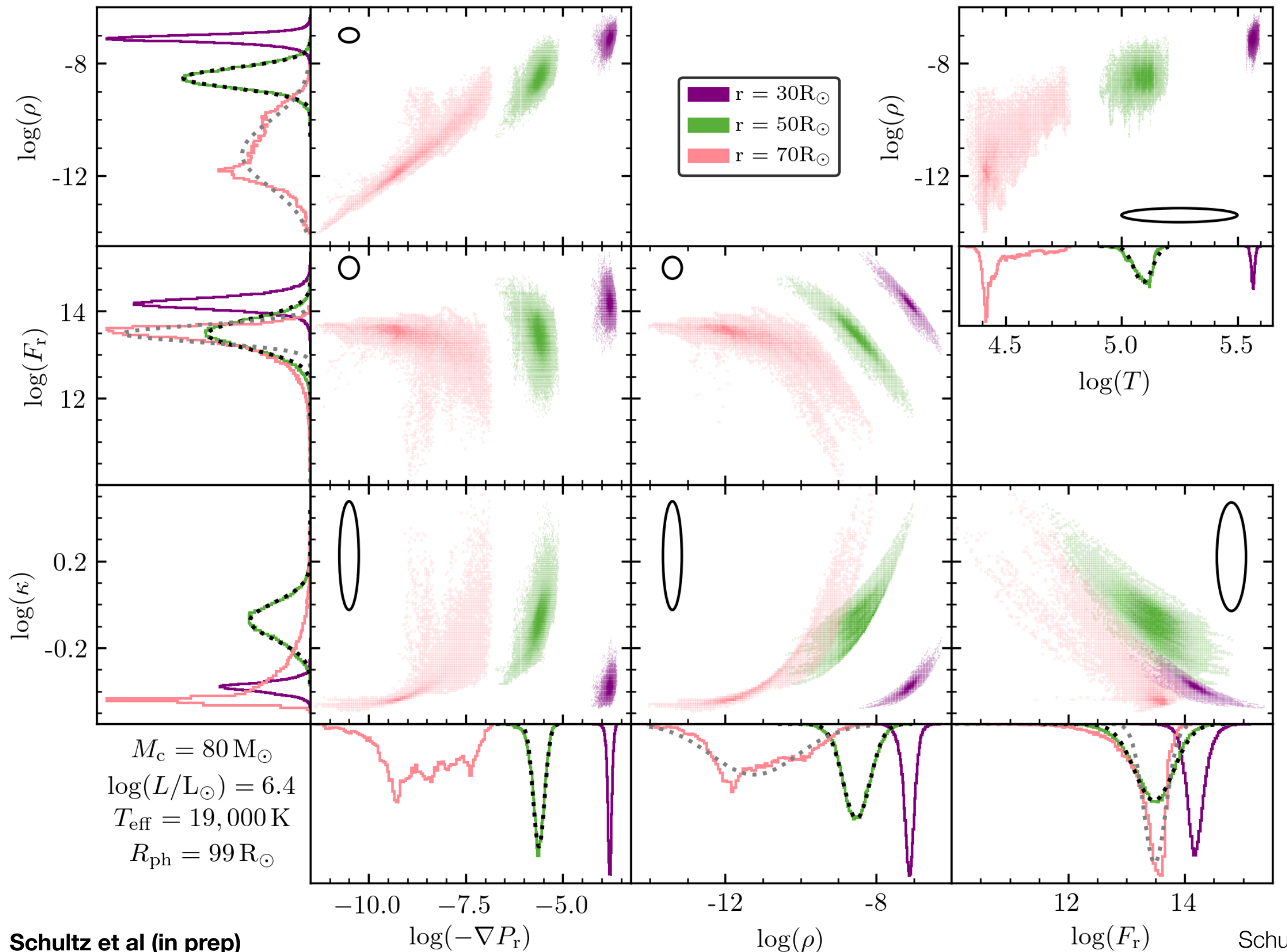


Are There Correlations?

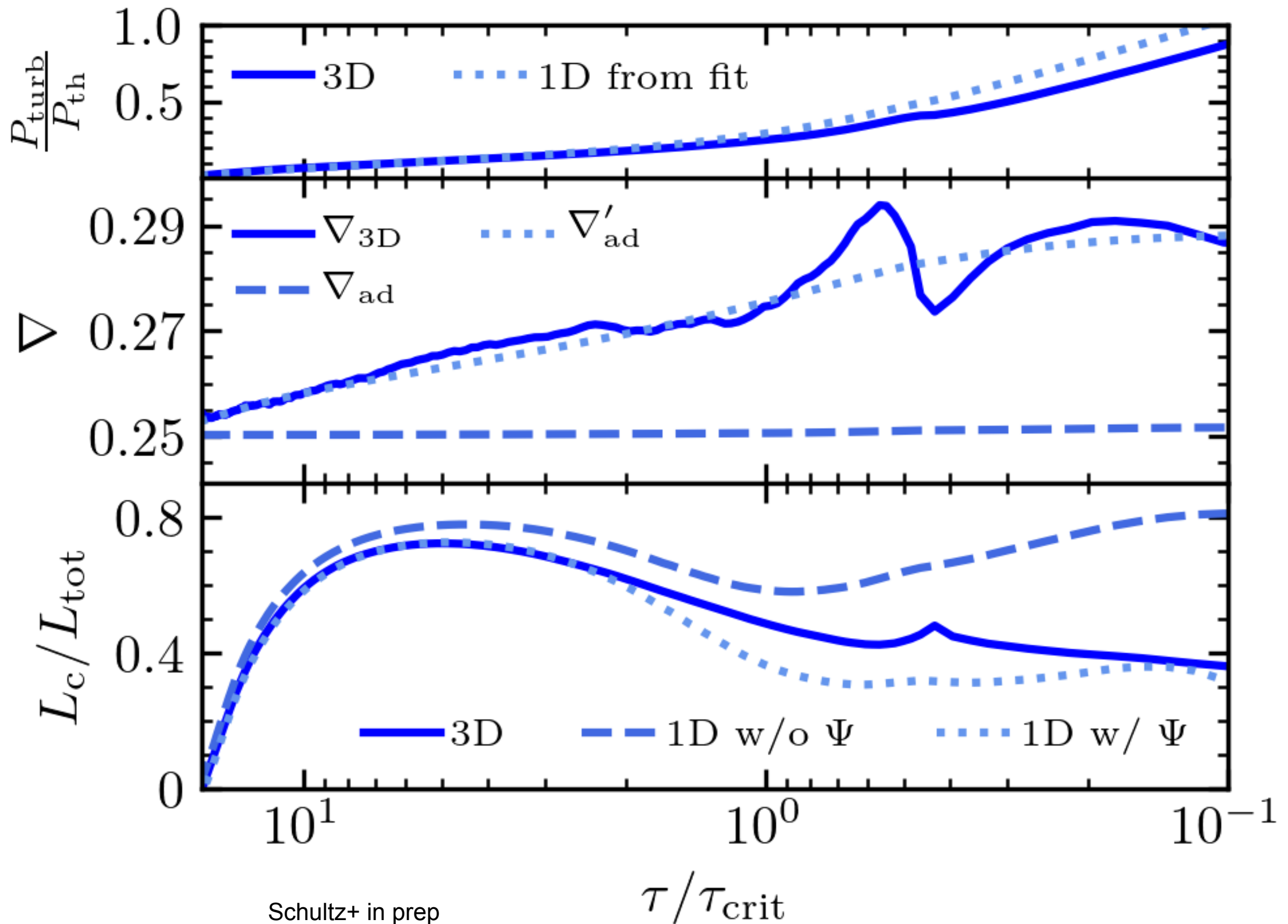


Correlations

$$f(x) = \frac{1}{\sqrt{2\pi}\sigma_{\ln(x)}} \exp\left(-\frac{(\ln(x) - \langle \ln(x) \rangle)^2}{2\sigma_{\ln(x)}^2}\right)$$



Turbulent pressure in grad

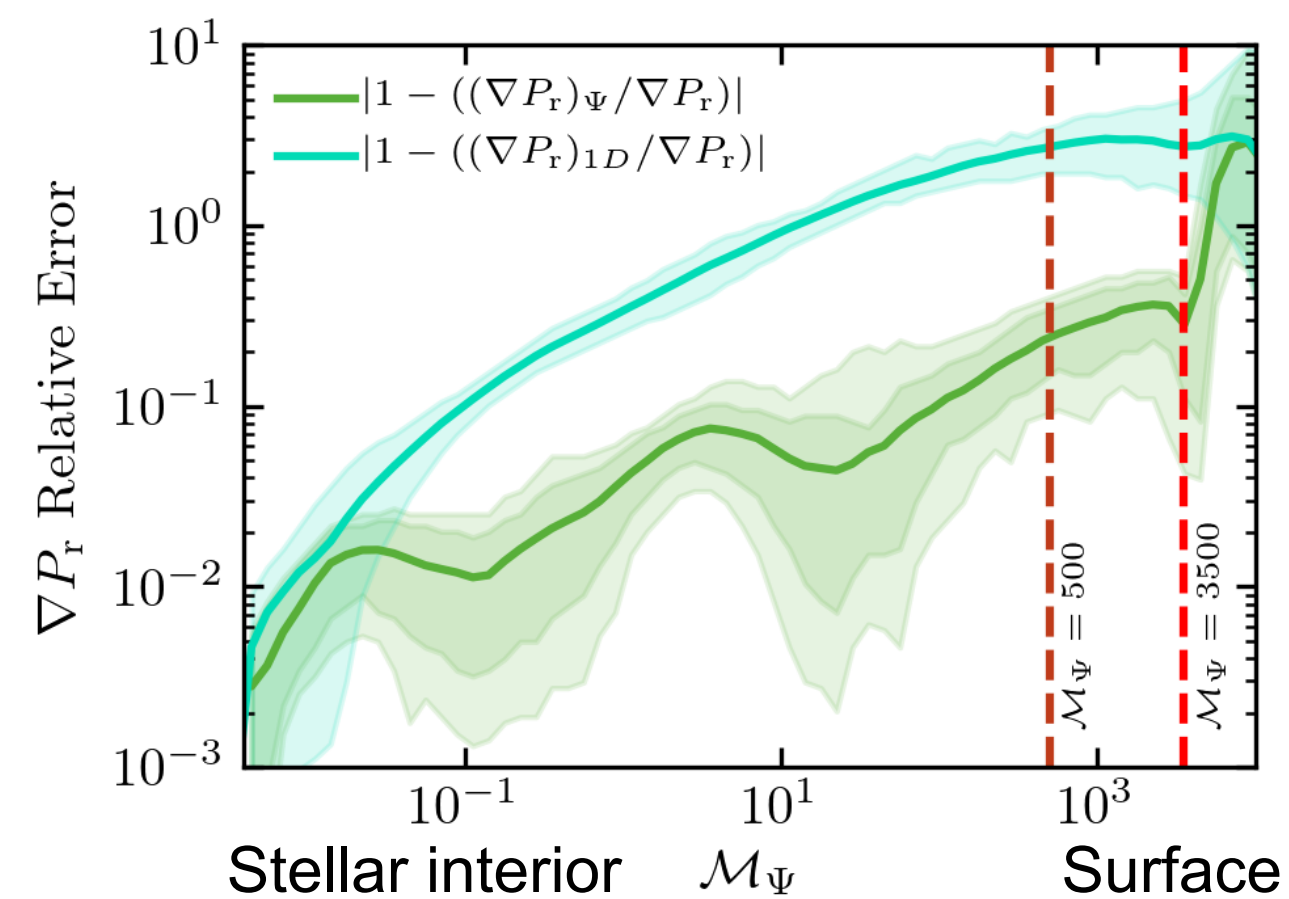


Schultz+ in prep

$\tau / \tau_{\text{crit}}$

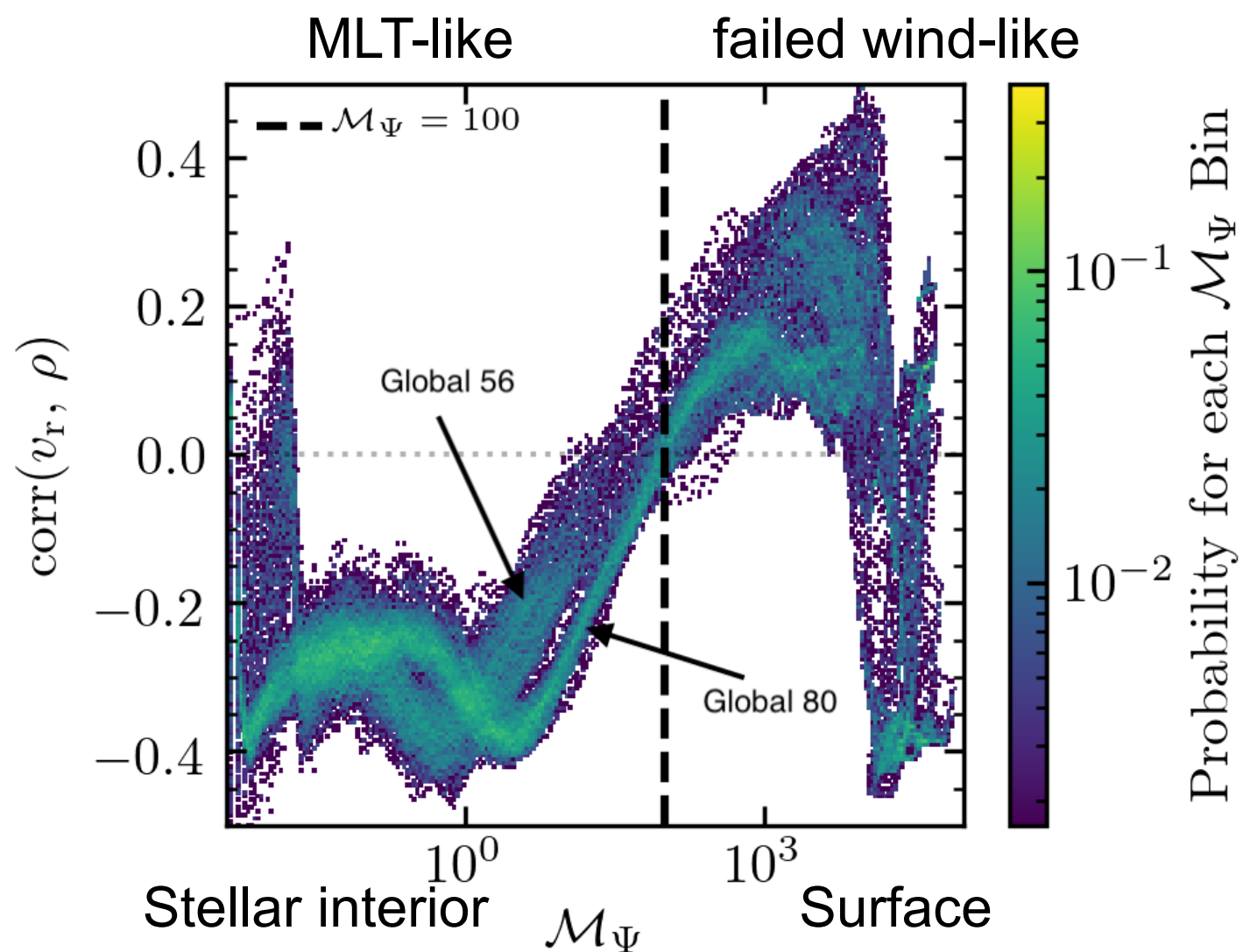
Global Models Shed Light on Turbulent Convection

Turbulence reduces radiation pressure gradient



Schultz, Bildsten, Jiang (2020)

Two types of turbulence



Schultz et al (in prep)