

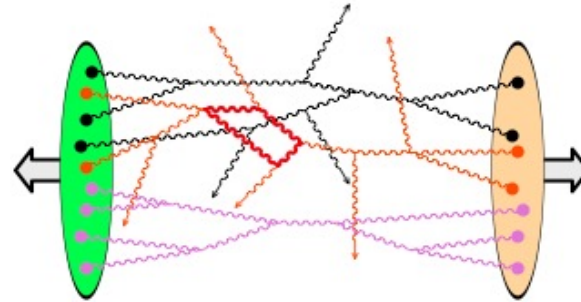
# Thermalization in quantum chromodynamics: Ab initio approaches and interdisciplinary connections

Raju Venugopalan  
Brookhaven National Laboratory

KITP Conference on non-equilibrium universality in many-body systems, September 28, 2021

# Boiling the QCD vacuum in ultrarelativistic heavy-ion collisions

Collisions of heavy-ions at **200 GeV/nucleon** at BNL's Relativistic Heavy Ion Collider (RHIC) and **5.5 TeV/nucleon** at CERN's Large Hadron Collider (LHC) create a deconfined non-Abelian fluid, the **Quark-Gluon Plasma (QGP)**

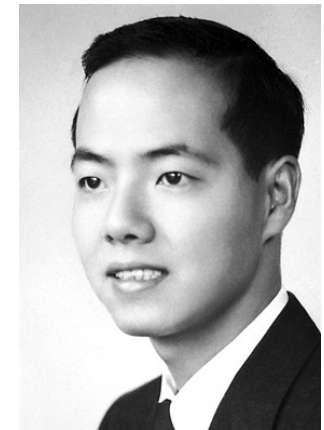


Hottest matter on earth:  
 $T \sim 4 * 10^{12}$  Kelvin



At lower energy, physics sensitive to critical phenomena in the  $T-\mu_B$  phase diagram

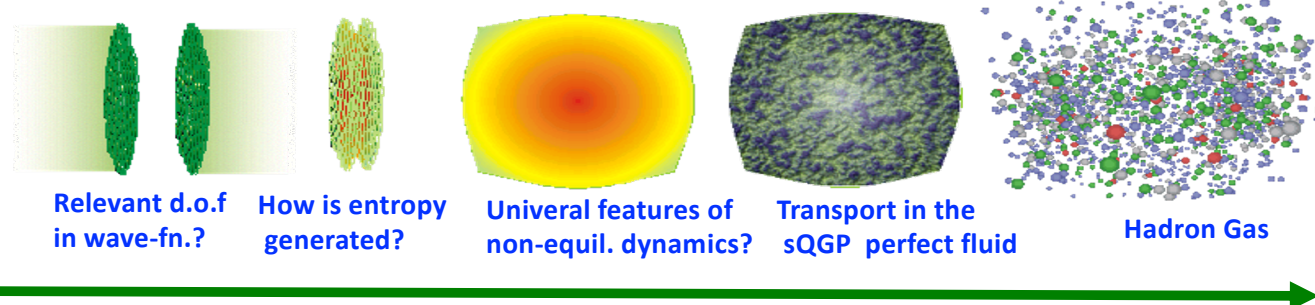
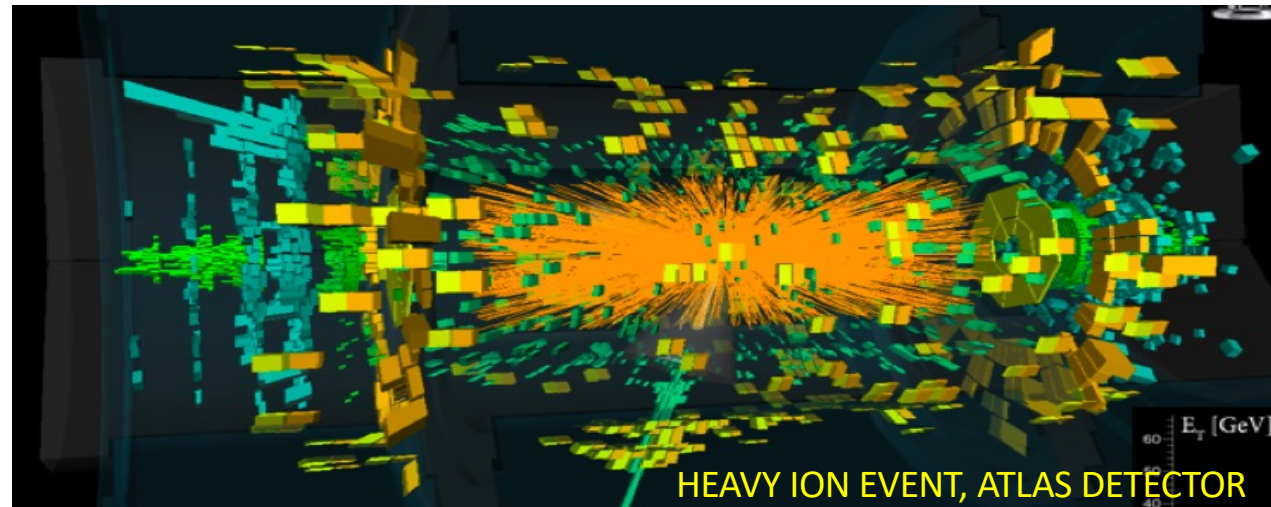
Beam energy scan at RHIC



TD Lee  
Nobel Laureate (1957)

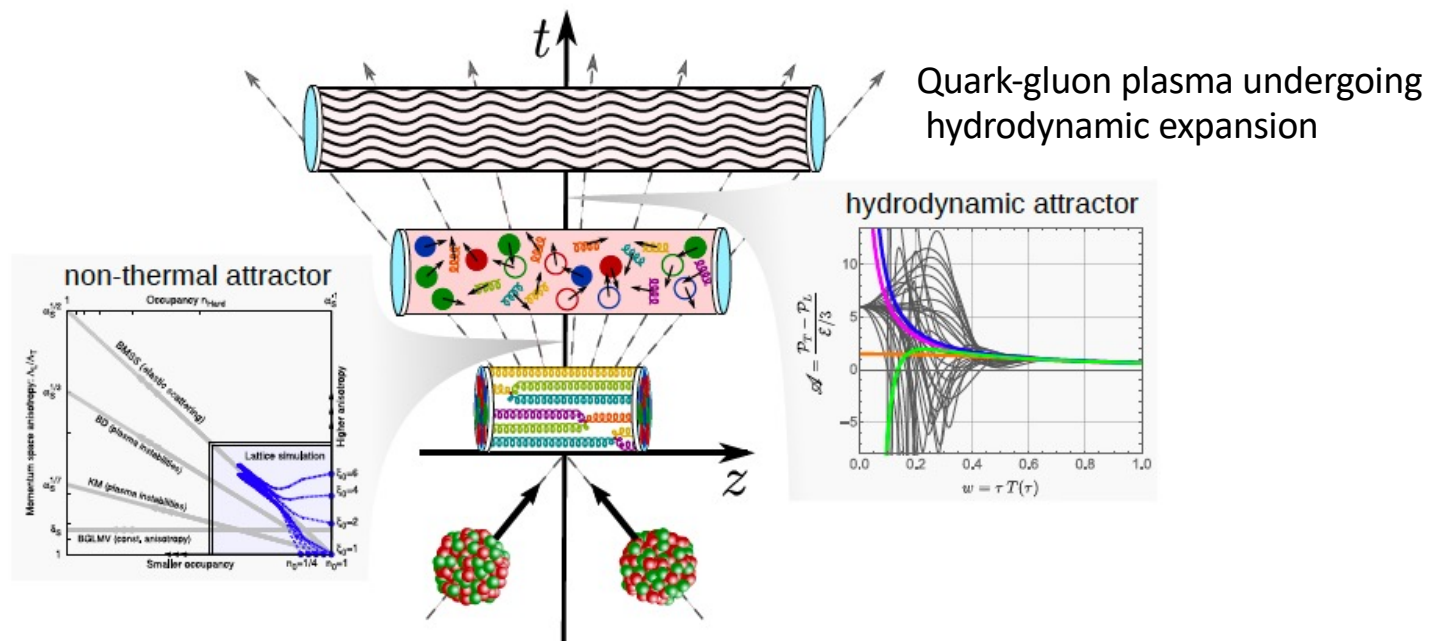
# “Standard model” of a heavy-ion collision

Initial energy density  $> 10^3$  that at cross-over to hadron gas



Thermalization a very difficult problem in quantum field theory even for simple systems: Ab initio theory, phenomenology and insight from other sub-fields, all essential for progress

# Spacetime evolution of a heavy-ion collision

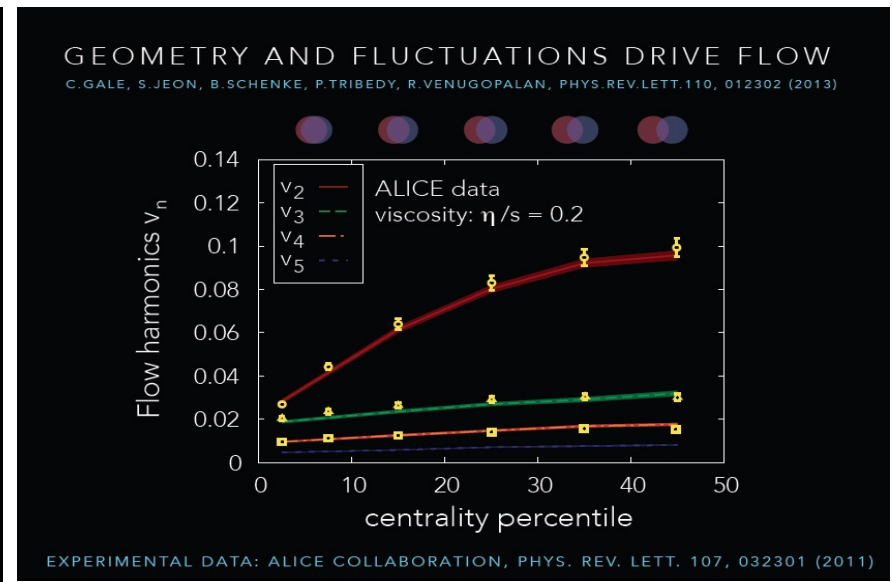
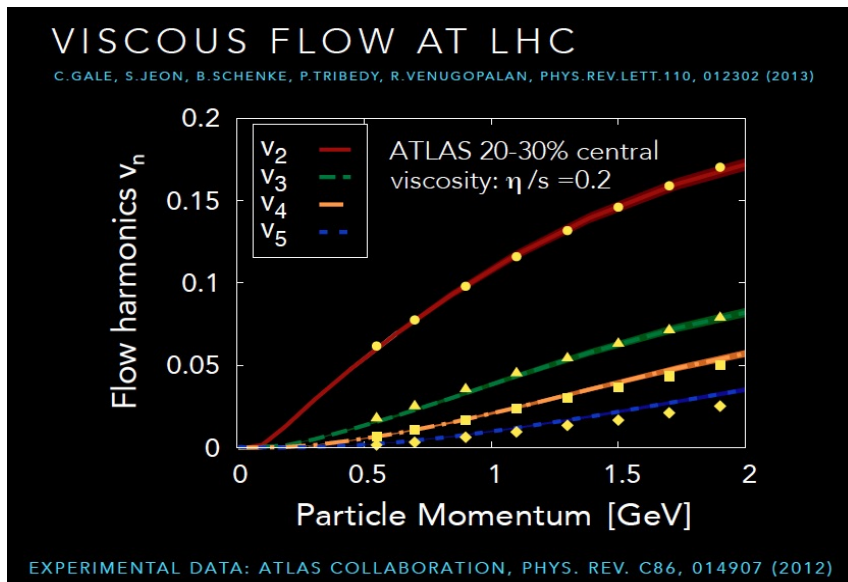
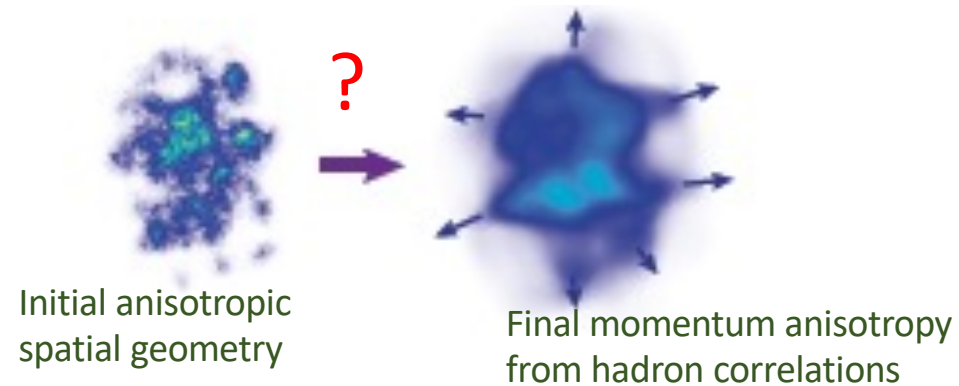


*QCD thermalization: Ab initio approaches and interdisciplinary connections*

Jürgen Berges, Michal P. Heller, Aleksas Mazeliauskas, and Raju Venugopalan

Rev. Mod. Phys. **93**, 035003 (2021)

# The unreasonable effectiveness of hydrodynamics



# Approaches to thermalization in heavy-ion collisions

Two “clean” *ab initio* theoretical approaches:

- Holographic thermalization (paradigmatic example: AdS/CFT duality of strongly coupled  $(g^2 N_c \rightarrow \infty, N_c \rightarrow \infty)$  **N=4 SUSY YM** to classical gravity in  $\text{AdS}_5 \times \text{S}_5$ )  
Not the right theory but valuable insight into universal features of strongly correlated non-equilibrium dynamics. Example: transport coefficients, hydrodynamics far-from-equilibrium
- Highly occupied (occupancy  $f \gg 1$ ) QCD at weak coupling ( $g^2 \rightarrow 0, g^2 f \sim 1$ )

Much of this talk will be describe the second of these two approaches

Applications to heavy-ion collisions require *phenomenological extrapolations* of ab initio theory

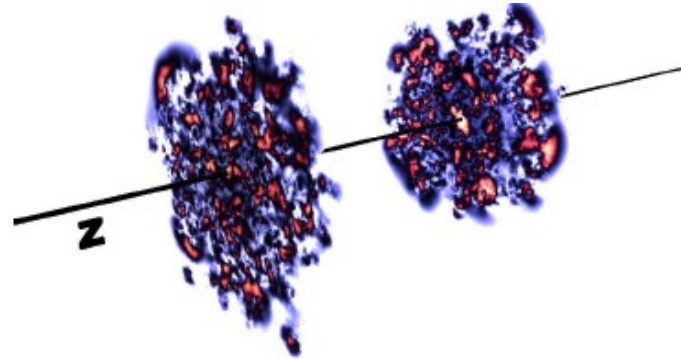
Universal features of the dynamics (non-thermal fixed points, hydrodynamization), and *interdisciplinary connections* thereof, provide powerful guidance

# Highly occupied glue far off-equilibrium: the Glasma

In QCD's high energy (Regge) limit,  
 nuclei can be described as highly occupied  $\sim \frac{1}{\alpha_S}$   
 gluon shock waves with emergent hard scale  $Q_S$

EFT description: Color Glass Condensate

Review: Gelis, Jalilian-Marian, Iancu, RV, arXiv:1002.0333

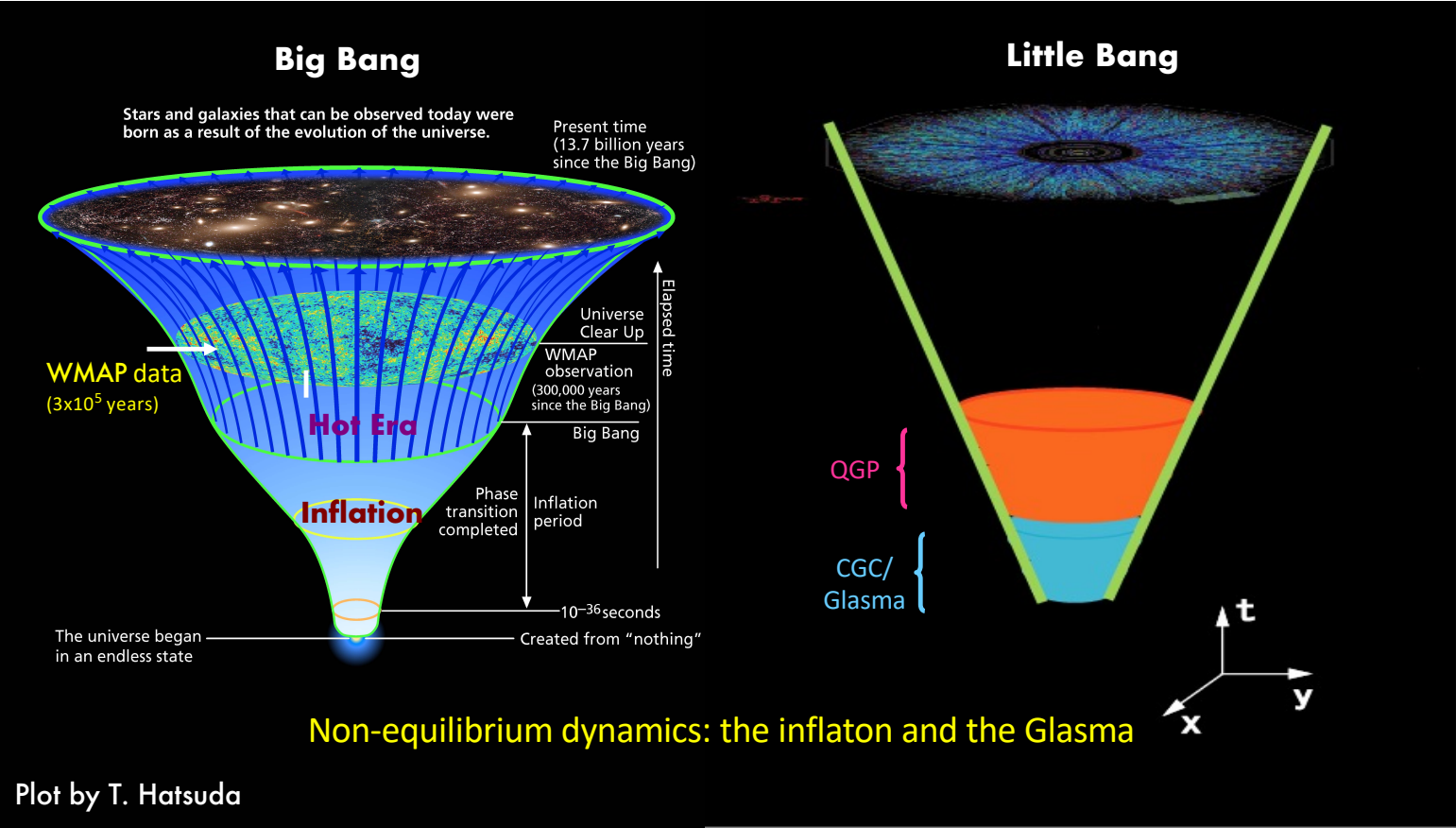


Collisions of "lumpy" gluon shock waves

Leading order solution shock wave solution: QCD Yang-Mills eqns in presence of light-cone sources

$$D_\mu F^{\mu\nu,a} = \delta^{\nu+} \rho_A^a(x_\perp) \delta(x^-) + \delta^{\nu-} \rho_B^a(x_\perp) \delta(x^+)$$

$$\langle \rho_{A(B)}^a(x_\perp) \rho_{A(B)}^a(y_\perp) \rangle = Q_{S,A(B)}^2 \delta^{(2)}(x_\perp - y_\perp)$$





## Big Bang vs. Little Bang

Decaying Inflaton  
with occupation #  $1/g^2$



Decaying Glasma  
with occupation #  $1/g^2$

Explosive amplification of low  
momentum small fluctuations  
(parametric resonance in preheating)



Explosive amplification of low  
momentum small fluctuations  
(Weibel instabilities)

Interaction of fluctuations/inflaton  
-> thermalization?

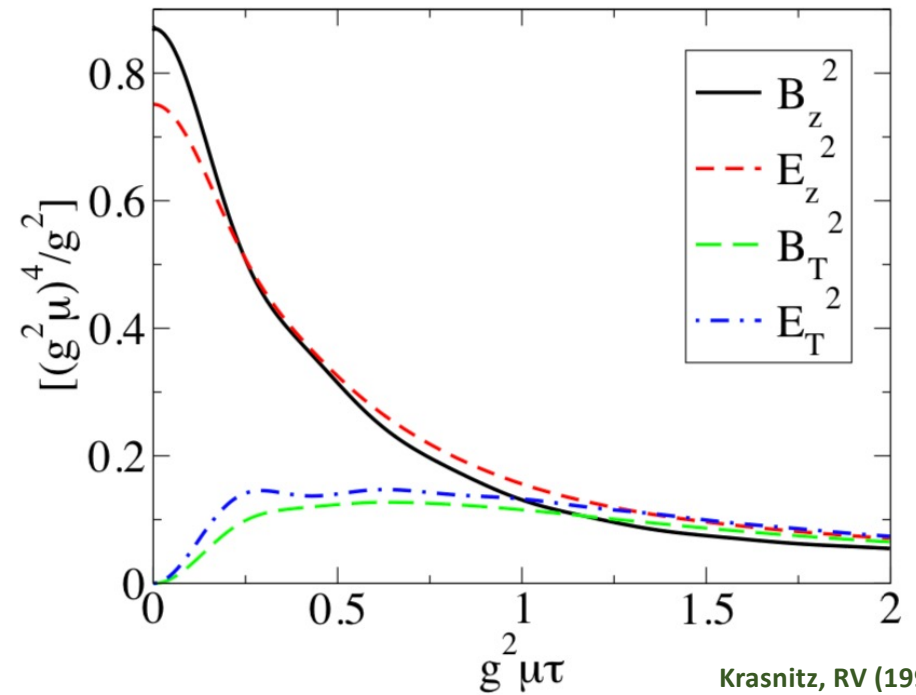
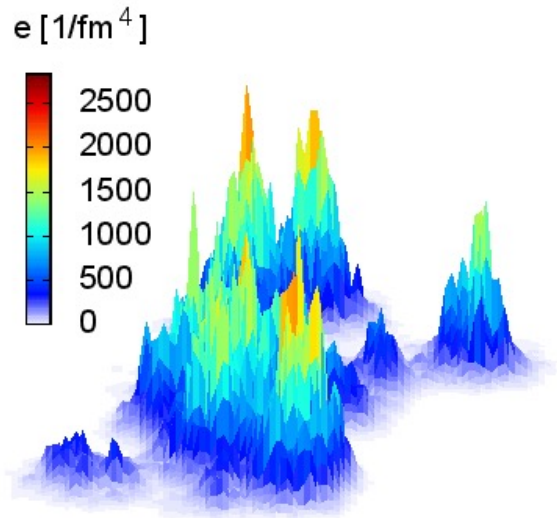


Interaction of fluctuations/Glasma  
-> thermalization?

Other common features: turbulence, topological defects,...

*Turbulent thermalization,*  
Micha and Tkachev, arXiv:hep-ph/0403101

## Boost invariant Yang-Mills: numerical solutions



Krasnitz, RV (1998)  
Lappi, hep-ph/0606207

Glasma energy density and pressure

$$T_{\mu\nu}(\tau = 0) = \frac{1}{2}(B_z^2 + E_z^2) \times \text{diag}(1, 1, 1, -1)$$

Initial longitudinal pressure is negative: Goes to  $P_L = 0$  from below with time evolution

# Decoherence and prethermalization from quantum fluctuations

“Toy” example: scalar  $\Phi^4$  theory

$$\phi(\tau, \eta, x_{\perp}) = \phi_{\text{cl.}}(\tau, x_{\perp}) + \frac{1}{2} \int \frac{d\nu}{2\pi} d\mu_k c_{\nu k} e^{i\nu\eta} \chi_k(x_{\perp}) H_{i\nu}(\lambda_k \tau) + c.c$$

Gaussian random variable

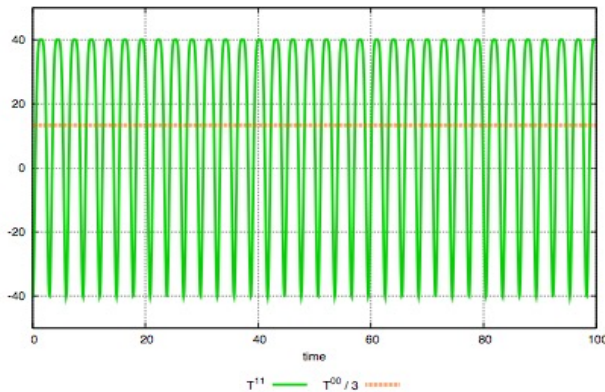
$$\langle c_{\nu k} c_{\mu l} \rangle = 0$$

$$\langle c_{\nu k} c_{\mu l}^* \rangle = 2\pi \delta(\nu - \mu) \delta_{kl}$$

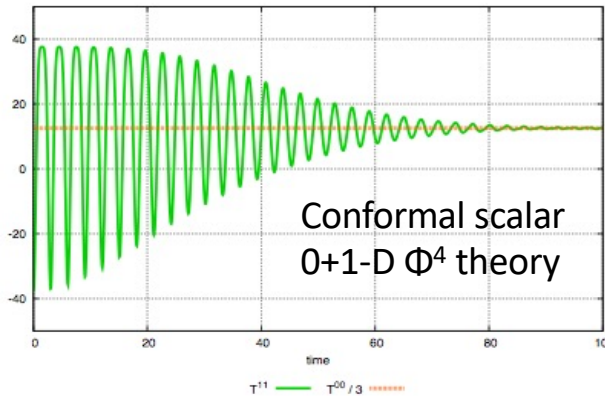
satisfies “small fluctuation” equation

$$[-\partial_{\perp}^2 + V''(\phi_0)]\chi_k = \lambda_k^2 \chi_k$$

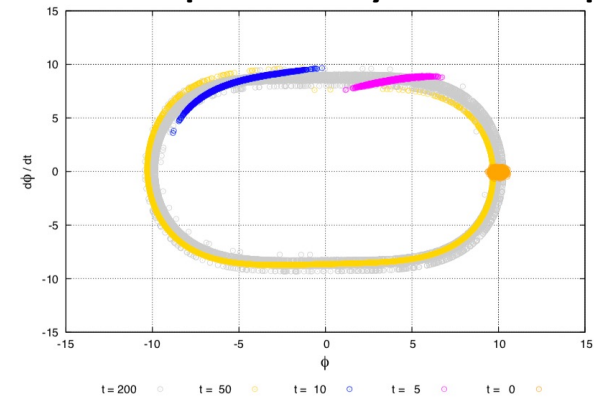
Energy density and pressure without averaging over fluctuations



Energy density and pressure after averaging over fluctuations



Phase space density in Poincaré plane



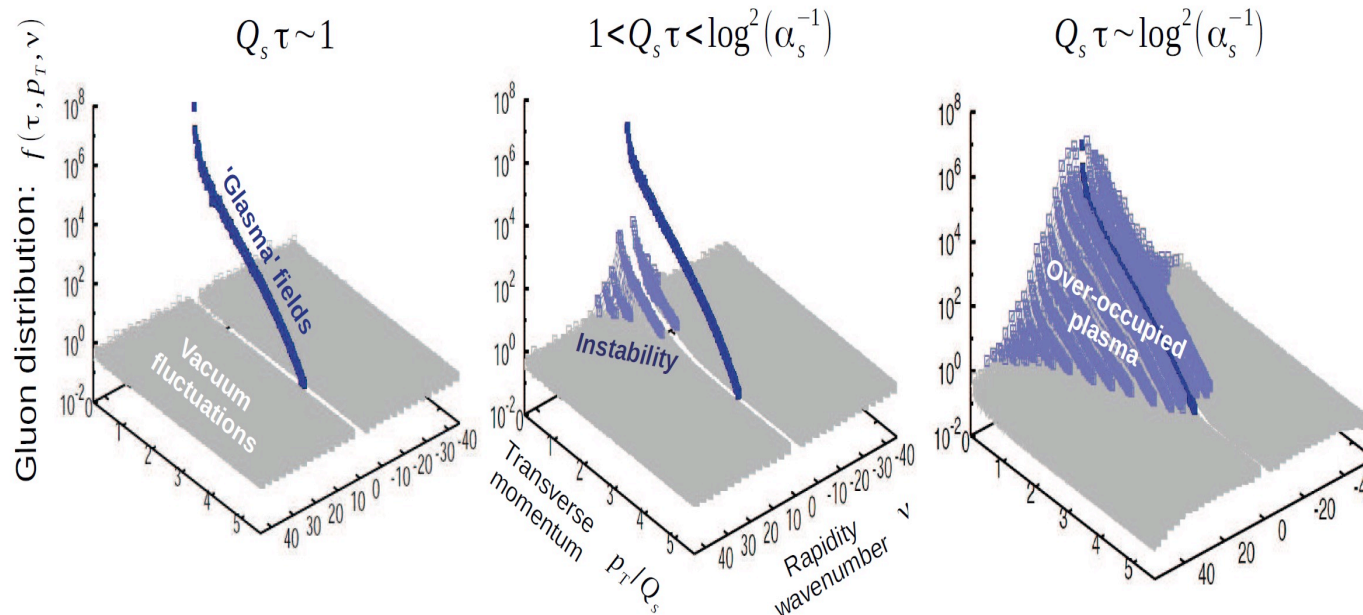
These quantum modes satisfy an “eigenstate thermalization” criteria conjectured by Berry (and significantly developed by Srednicki and others) as essential for thermalization of a quantum fluid

Berges, Borsanyi, Wetterich (2004)  
Dusling, Epelbaum, Gelis, RV (2011)

# Explosive amplification of quantum fluctuations

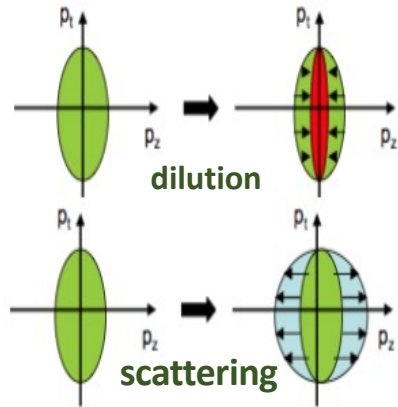
Longitudinally expanding Glasma fields are unstable to quantum fluctuations... leading to an explosive “Weibel”-like instability.

Rapid decoherence and overpopulation of all momentum modes



Classical-statistical lattice simulations of 3+1-D gluon fields exploding into the vacuum

# Classical-statistical simulations of 3+1-D Yang-Mills

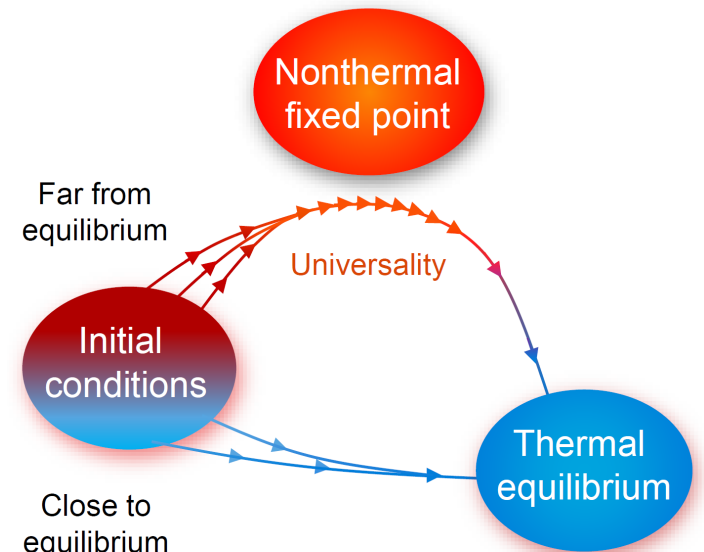
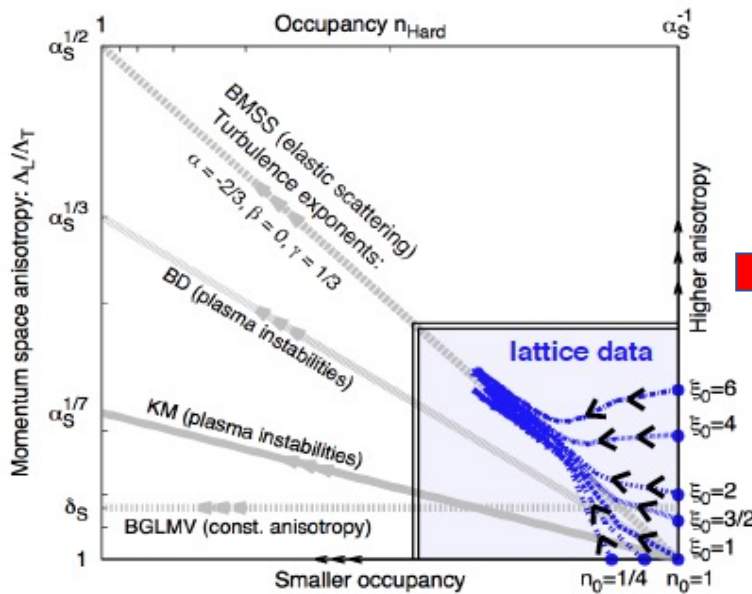


After rapid scrambling of information by quantum fluctuations, competition between dilution due to expansion and isotropization due to scattering

Single particle distributions become self-similar in time characterized by universal exponents – helps identify “right” kinetic theory

$$f(p_{\perp}, p_z, t) = t^{\alpha} f_S(t^{\beta} p_T, t^{\gamma} p_z)$$

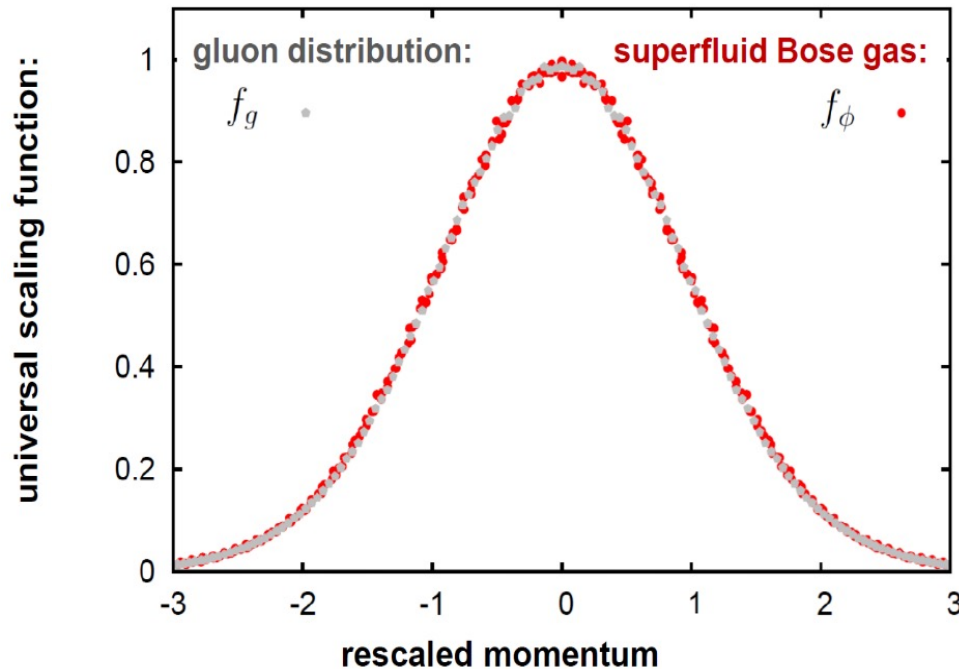
Increasing pressure anisotropy



Berges, Boguslavski, Schlichting, RV (2014)

# The Glasma and overoccupied ultracold quantum gases

Simulations of self-interacting scalar fields with identical initial conditions demonstrates remarkable *universality of longitudinally expanding world's hottest and coolest fluids*



In a wide inertial range, scalar & gauge fields have identical scaling exponents & functions

$$f(p_T, p_z, \tau) = \tau^\alpha f_S(\tau^\beta p_T, \tau^\gamma p_z)$$

$$\tau = \sqrt{t^2 - z^2}$$

$$\alpha = -\frac{2}{3}, \beta = 0, \gamma = 1/3$$

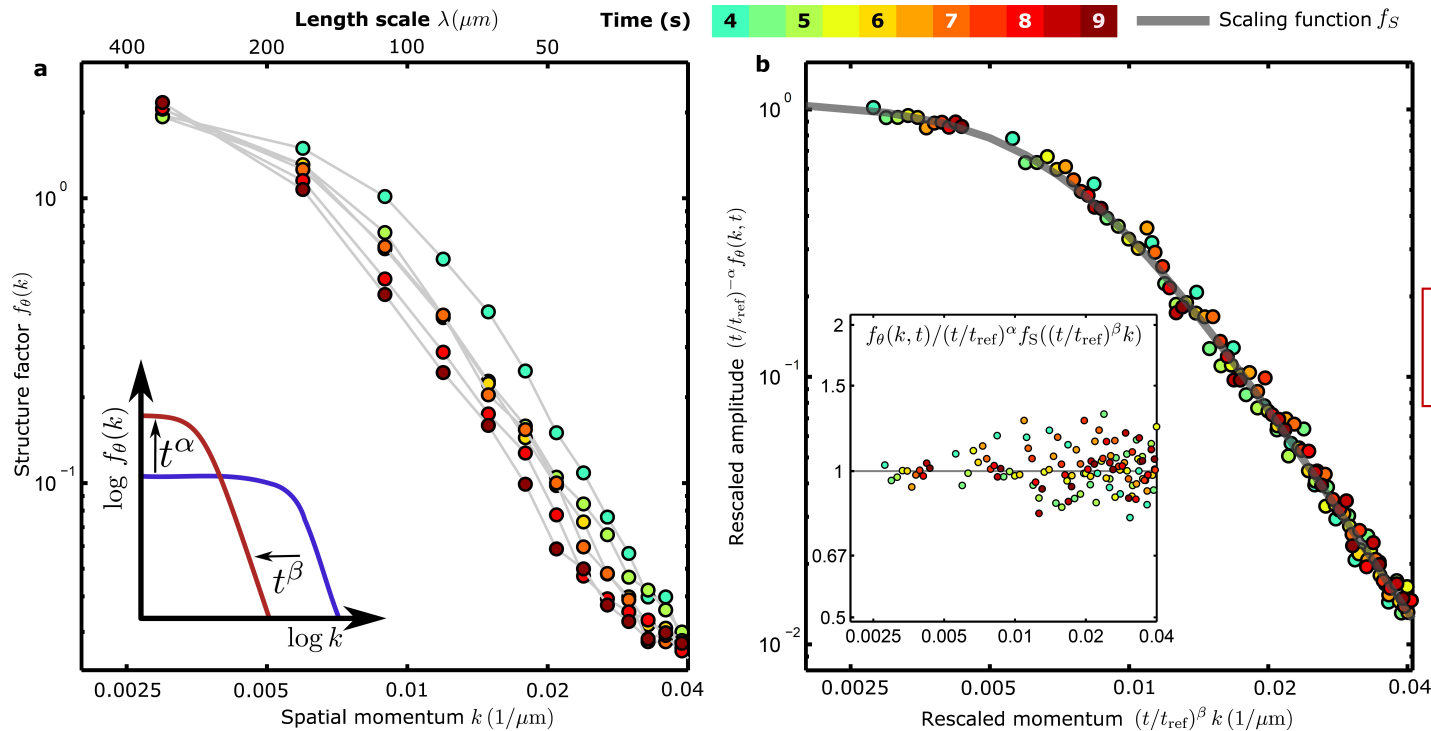
Berges, Boguslavski, Schlichting, RV, PRL (2015)  
Editor's suggestion

# The Glasma and over-occupied quantum gases

Similar non-thermal fixed points discovered in cold atom experiments

- albeit only for static geometry so far

## $^{87}\text{Rb}$ BEC in a quasi 1D optical trap

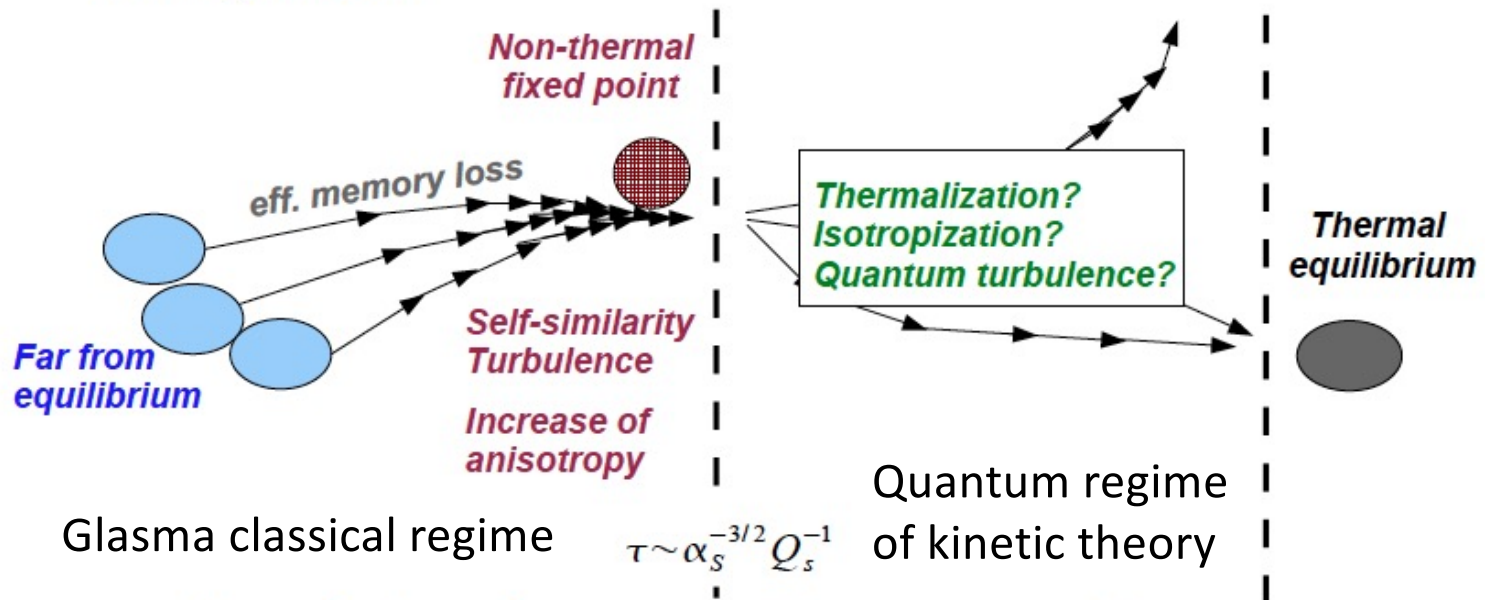


$$f_\theta(k, t) = t^\alpha f_S(t^\beta k)$$

$$\alpha = 0.33 \pm 0.08 \quad \beta = 0.54 \pm 0.06$$

Oberthaler BEC Labs  
Prüfer et al, arXiv:1805.11881, *Nature* (2018)

# From nuts to soup: bottom-up thermalization



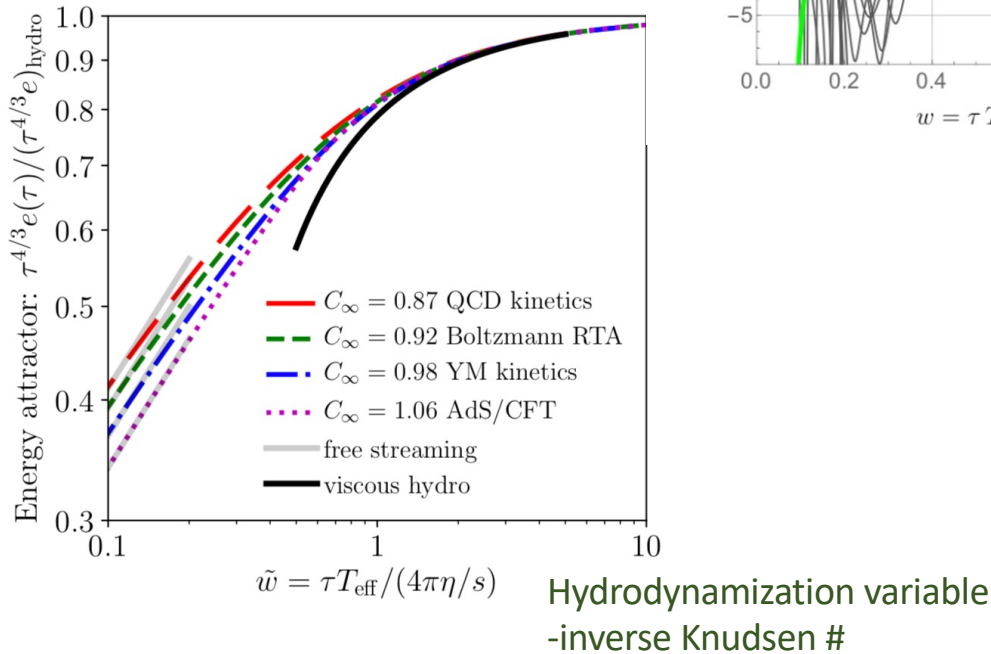
Thermalized soft bath of gluons for  $\tau > \frac{1}{\alpha_S^{5/2}} \frac{1}{Q_S}$

Thermalization temperature of  $T_i = \alpha_S^{2/5} Q_S$

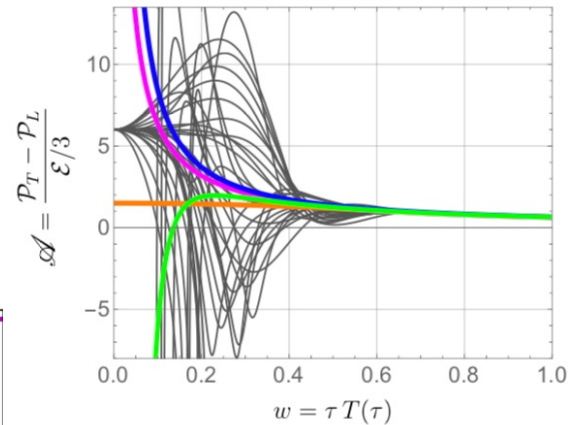


# Hydrodynamics far-from-equilibrium

Rapid convergence to viscous hydrodynamics



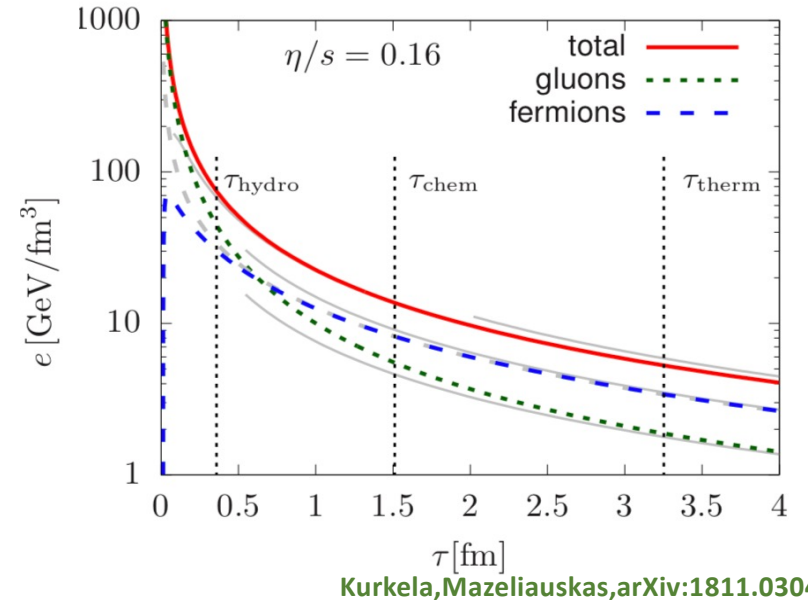
Giacalone, Schlichting, Mazeliauskas, arXiv:1908.02866



Magenta, Blue, Green, 1<sup>st</sup>, 2<sup>nd</sup>, 3<sup>rd</sup>, hydro constitutive relations. Orange curve is hydro attractor; different init. conditions (grey)

Heller, Janik, Witaszczyk, arXiv:1203.0755  
Romatschke, arXiv:1704.08699

Hydrodynamization < chemical < thermal equilibration



Kurkela, Mazeliauskas, arXiv:1811.03040

# From nuts to soup: bottom-up thermalization

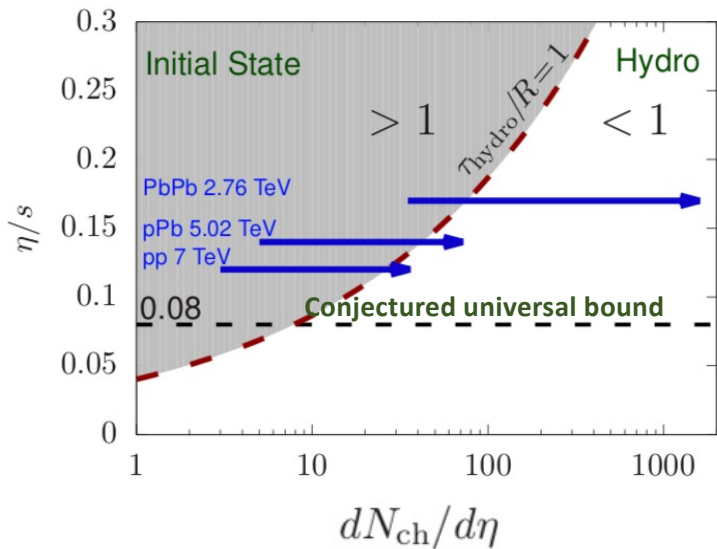
Thermalized soft bath of gluons for  $\tau > \frac{1}{\alpha_S^{5/2}} \frac{1}{Q_S}$

Thermalization temperature of  $T_i = \alpha_S^{2/5} Q_S$

Since  $\alpha_S \propto \frac{1}{\log Q_S}$

then  $\tau_{\text{therm}} \propto \frac{(\log Q_S)^{5/2}}{Q_S} \rightarrow 0$  as  $Q_S \rightarrow \infty$

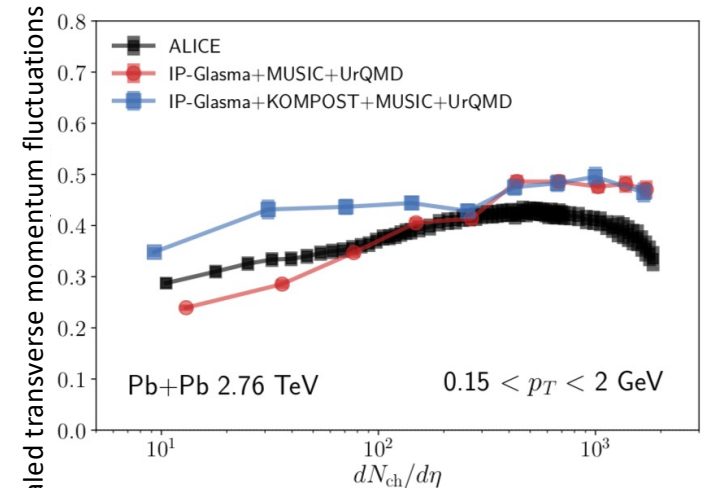
Glasma/bottom-up prediction: *In the Regge limit, matter thermalizes almost instantaneously*



Mazeliauskas, arXiv:1807.05586  
See also Kurkela, Wiedemann, Wu, arXiv:1905.05139

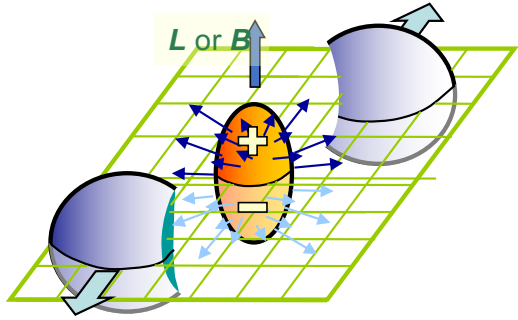
Event engineering:

Promising tool to quantify thermalization times in small systems

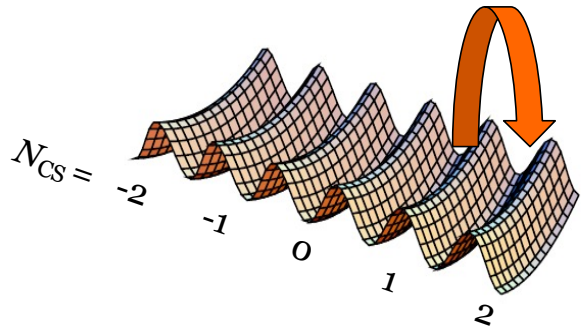


Schenke, Shen, Teaney, arXiv:2004.00690  
**KOMPOST implementation of bottom-up:**  
Kurkela et al., arXiv:1805.00961

# Early-time dynamics and topology: The Chiral Magnetic Effect



External (QED) magnetic fields  
 10<sup>18</sup> Gauss, of Magnetar strength!  
 -dies very rapidly (~1 fm) after collision

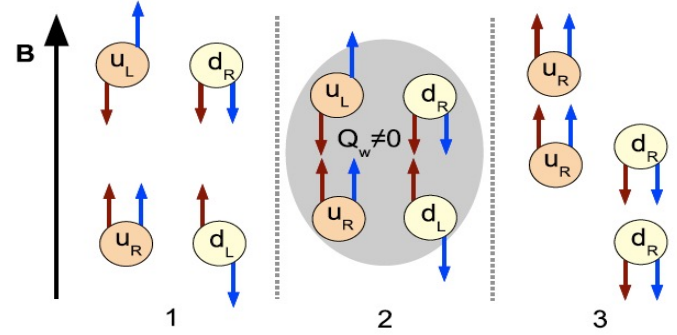


Over barrier topological (sphaleron) transitions ... analogous to proposed mechanism for electroweak baryogenesis

Kharzeev, McLerran, Warringa (2007)  
 Kharzeev, Fukushima, Warringa (2008)

Considerable experimental activity in HI collisions and condensed matter systems

$$\vec{J}_{CME} = \frac{e^2}{2\pi^2} \mu_5 \vec{B}$$



Chiral magnetic effect

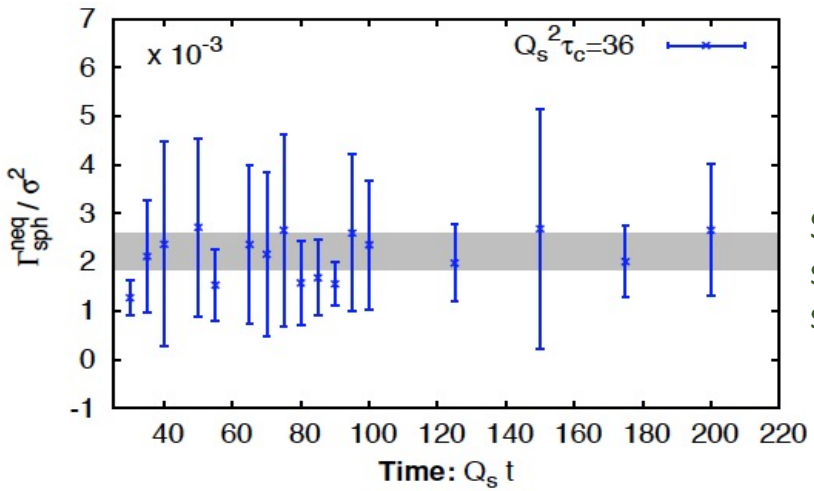
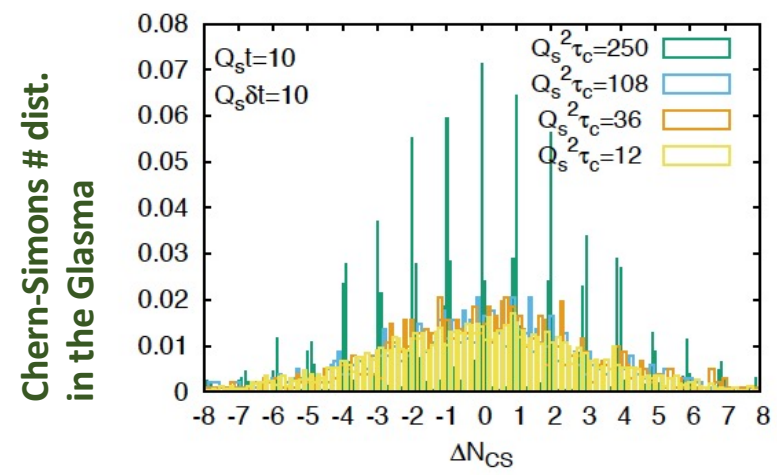
Ruled out definitively in isobar RHIC collisions at the highest energy.

STAR collaboration,  
 arXiv:2109.00131

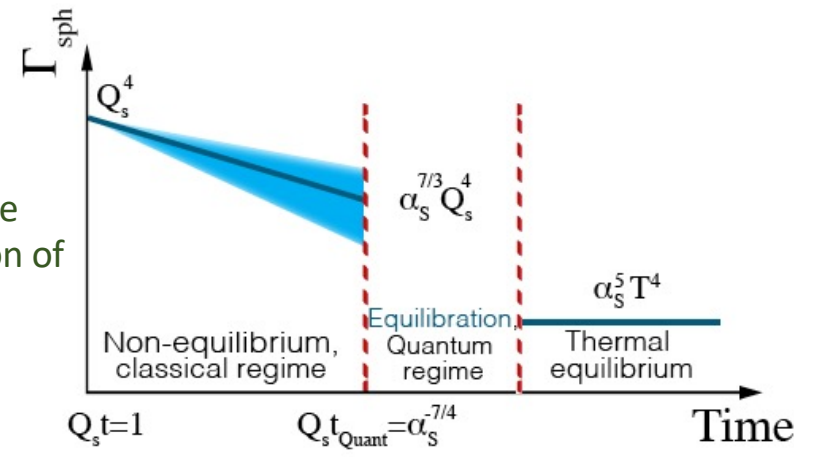
Prospects may still exist at lower energies

# Early-time dynamics and topology: The Chiral Magnetic Effect

Mace, Schlichting, RV: PRD (2016)



Sphaleron transition rate scales with string tension of spatial Wilson loops



# My perspective on some outstanding open questions

What is the role of entanglement in this complex many-body system ?

Example: Entanglement “induced” thermalization of a 1+1-D QCD string

**Berges, Floerchinger, RV,  
arXiv: 1705.05338, 1712.09362**

Are there other systems in nature which have dynamical features universal to the CGC/Glasma ?

Example: a CGC-Black Hole correspondence

**Dvali, RV, arXiv:2106.11989**

Can some of its real-time dynamical features be simulated by quantum computers ?

Example: Since particle digitization strategy motivated by Feynman’s parton model

**Barata, Mueller, Tarasov, RV  
arXiv:2012.00020**