

# **The Energy Interpretation of Interacting Static Black Holes**

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# The Energy Interpretation of Interacting Static Black Holes

## Outline:

- BHs in extra dimensions
- BHs with electric charge

# Black Holes in Extra Dimensions

**Based on:**

Scott Fraser, Doug Eardley

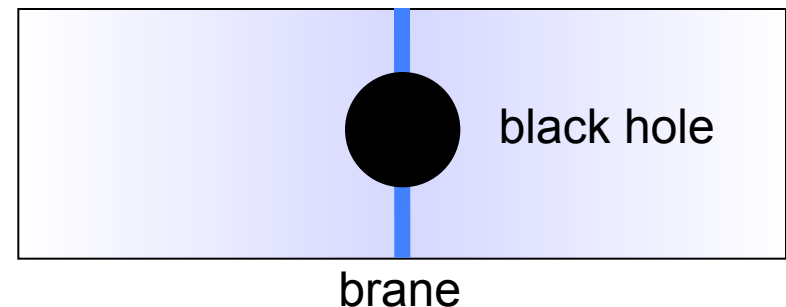
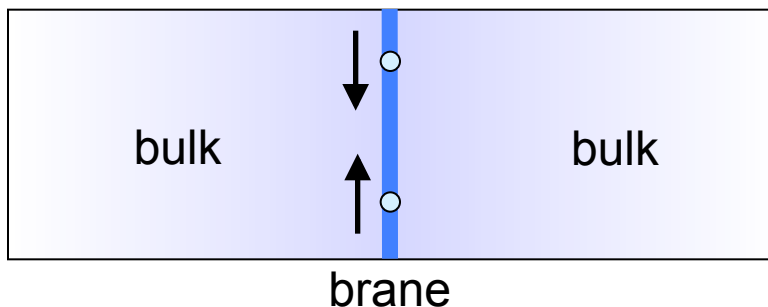
arXiv:1408.4425 [gr-qc]

arXiv:1409.0884 [gr-qc]

**(to appear in Phys. Rev. D)**

# Introduction: Extra Dimensions

- Our observed universe has 4 dimensions (3 spatial, 1 time).
- Use general relativity with extra dimensions (total  $D > 4$ ).
- Model: observed universe = 4-dimensional surface (brane).
- Only gravity (including BHs) extends into the bulk ( $D > 4$ ).
- Gravity acts 4-dim. (large distance),  $D$ -dim. (small distance)



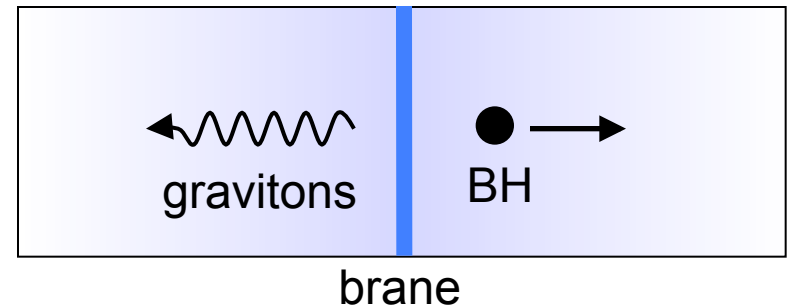
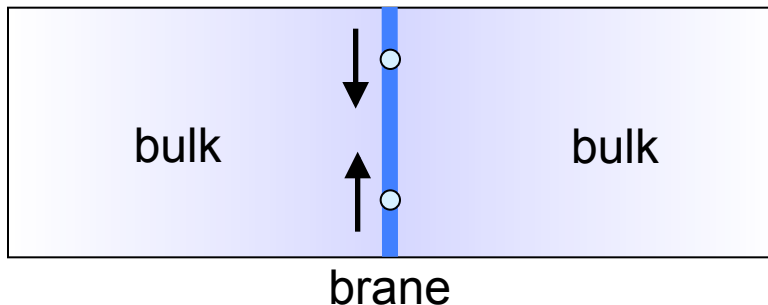
# Introduction: Extra Dimensions

- Geometry of extra dimensions could explain the hierarchy:

$$M_{4, \text{Planck}} / (\text{TeV}/c^2) \sim 10^{16} \quad \text{if} \quad M_{5, \text{Planck}} \sim \text{TeV}/c^2$$

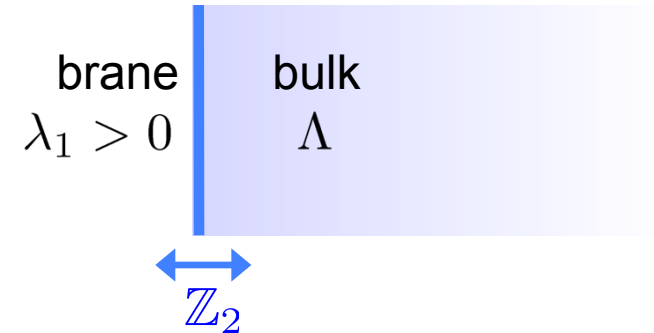
(Henceforth: use units with  $c = 1$ , use mass = energy)

- If we live on a brane (explaining hierarchy): high-energy collisions on the brane would produce small black holes!
- Could the black hole fall into the bulk?

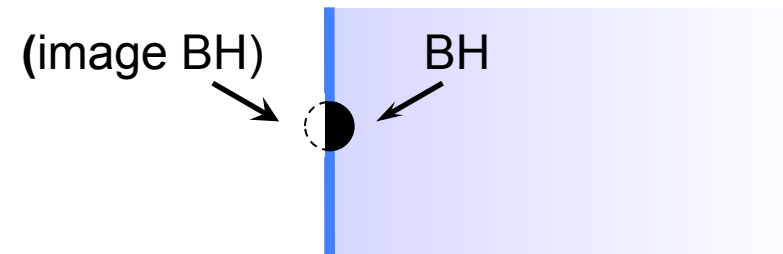


# Randall-Sundrum Model (RS2)

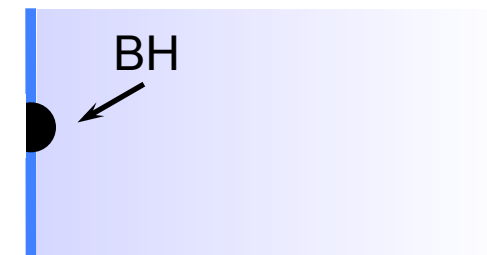
- **Brane: positive tension  $\lambda_1$**   
(gravitationally repulsive)
- **Bulk: cosmological constant  $\Lambda$**
- **Orbifold ( $\mathbb{Z}_2$  mirror) symmetry:**



- **Covering space:**  
**Identify** symmetric points across brane. All physics must be mirror-symmetric.



- **Physical space:**  
on **one** side of the brane.



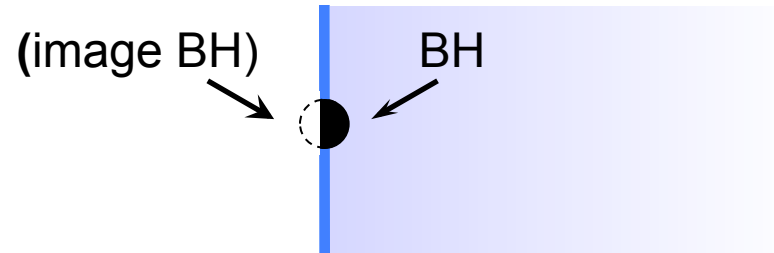
# Small Static Black Holes in RS2

- Well-known (numerically): a static small BH on the brane.

[Kudoh et al, 2003]

[Figueras and Wiseman, 2011]

[Abdolrahimi et al, 2013]

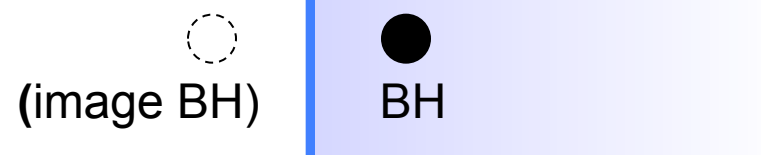


- Is this BH stable?

- The BH experiences:

- **repulsion** from brane (due to positive brane tension)
- **attraction** to brane (due to orbifold image)

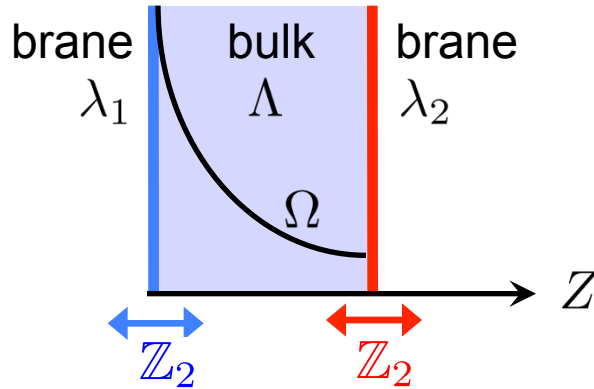
- Where these effects balance:  
expect a static small bulk BH



- Is this BH stable?

# Randall-Sundrum Asymptotics

- **RS1** (two branes; can explain the hierarchy if  $\Omega_1/\Omega_2 \sim 10^{16}$ )



$$ds^2 = \Omega^2 (\eta_{ab} dx^a dx^b + dZ^2)$$

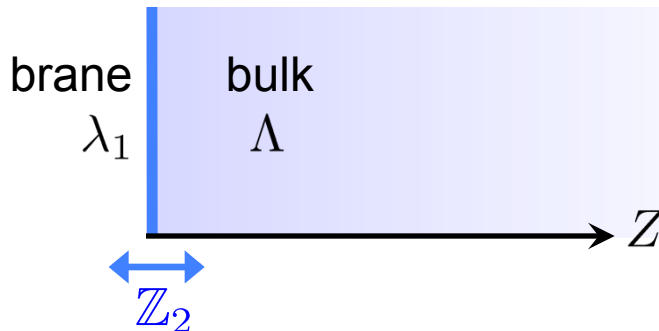
$$\Omega = \ell/Z$$

$$\lambda_1 = -\lambda_2 = \frac{2(D-2)}{8\pi G_D \ell}$$

$$P = -\frac{\Lambda}{8\pi G_D} = \frac{(D-1)(D-2)}{16\pi G_D \ell^2}$$

- Interbrane length:  $L_b = \ell \ln(\Omega_1/\Omega_2)$

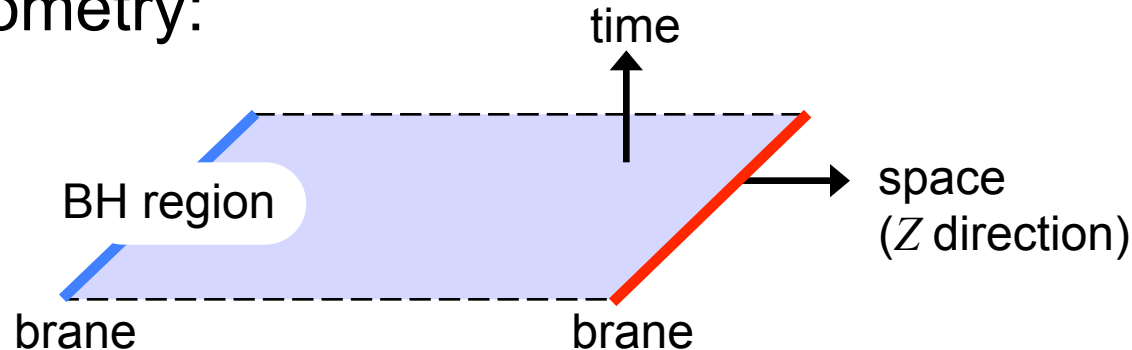
- **RS2** (single-brane limit; can't solve hierarchy problem)





# A Static BH Extremizes Mass

- Spatial geometry:

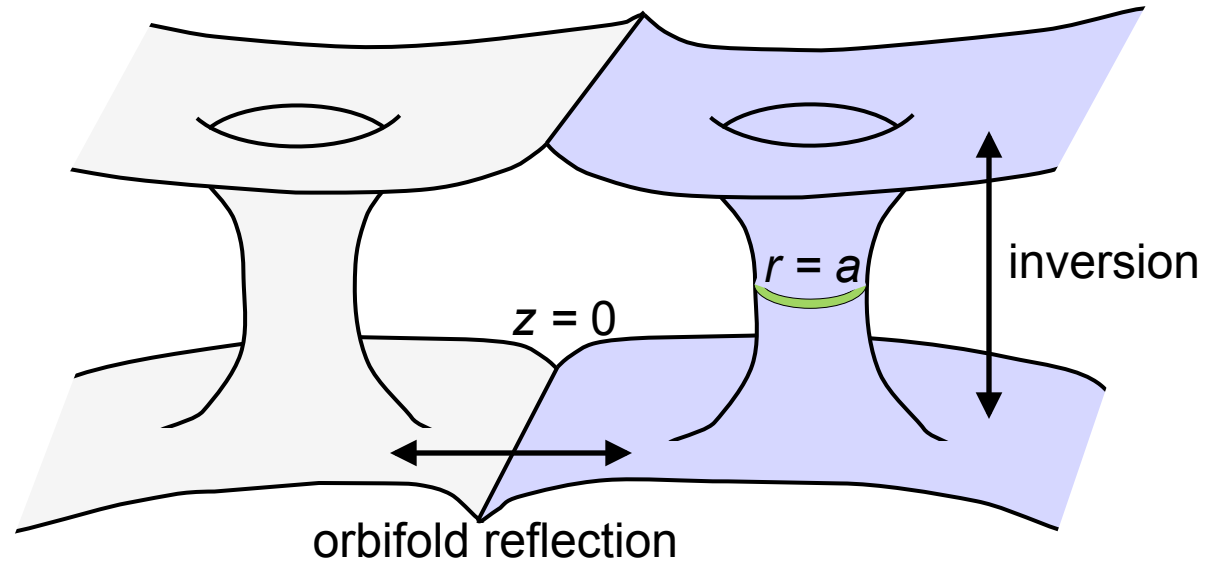


- Keep most parameters fixed:  $(L_b, \ell, \Omega_1, \Omega_2)$
- A static BH obeys a First Law:  $\delta M = \frac{\kappa}{8\pi G_5} \delta A$
- So: if a BH is static, it **extremizes** mass  $M$ , for fixed area  $A$ .
- Apparent horizon (AH): best approximation to event horizon.
- **We prove a Variational Principle** (for BH initially at rest):  
If mass  $M$  is **extremized** at fixed area  $A$ , the BH is **static**.

# Geometry for BH (initially at rest)

Visualize spatial geometry as:

(if BH is far from brane)



■ Small black hole:

$$a \ll \ell$$

■ Trial geometry:

$$ds^2 = \psi^{4/(D-3)} d\mathbf{x}^2 \quad , \quad \mathbf{x} = (\vec{\rho}, z) \quad , \quad Z = \ell + z$$

■ Equation for geometry:

$$\nabla_{\mathbf{f}}^2 \psi = \frac{(D-1)(D-3)}{4\ell^2} \psi^{(D+1)/(D-3)}$$

■ At throat ( $r = a$ ):

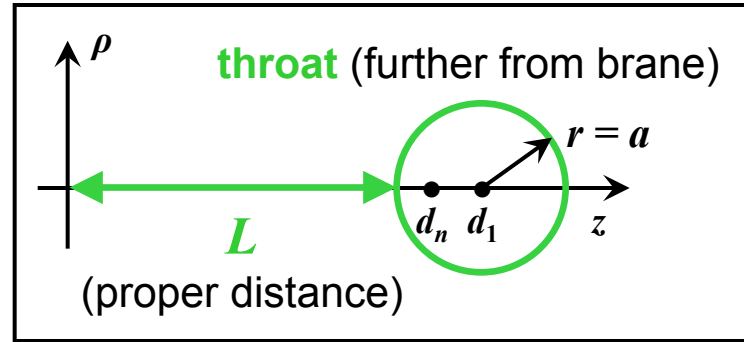
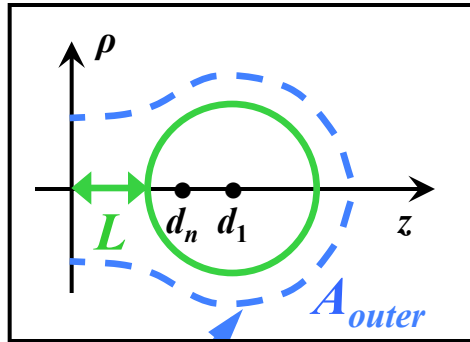
$$0 = r \partial_r \psi + \left( \frac{D-3}{2} \right) \psi$$

■ At brane ( $z = 0$ ):

$$0 = \partial_z \psi + \left( \frac{D-3}{2\ell} \right) \psi^{(D-1)/(D-3)}$$

# Black Holes (near the brane)

- Approximate  $\ell \rightarrow \infty$ :  $\nabla_f^2 \psi = 0$  ,  $\partial_z \psi \Big|_{z=0} = 0$
- Method of images [Misner 1963]:  $\psi = 1 + \sum_{n=1}^{\infty} \left( \frac{q_n}{|\mathbf{x} + \mathbf{d}_n|^{D-3}} + \frac{q_n}{|\mathbf{x} - \mathbf{d}_n|^{D-3}} \right)$



$$a = c \operatorname{csch} \mu_0$$

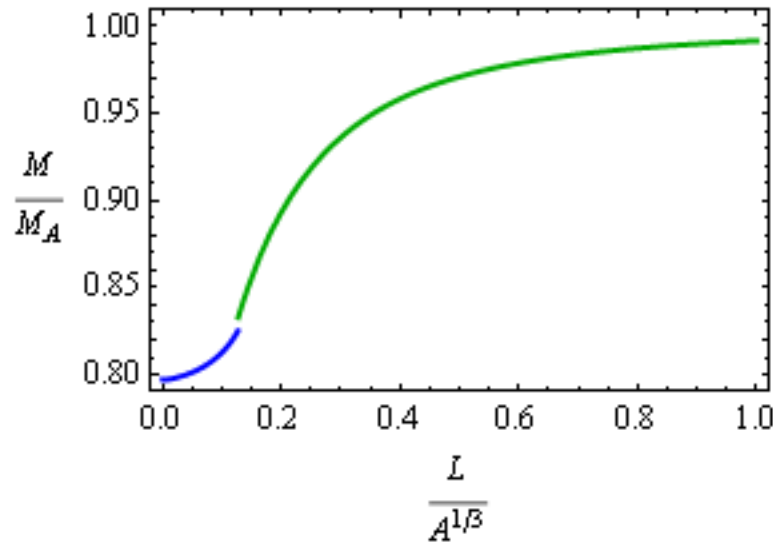
$$d_n = c \operatorname{coth} n\mu_0$$

$$q_n = (c \operatorname{csch} n\mu_0)^{D-3}$$

- **Find numerically:**  $A_{outer}$  ,  $\bar{\mu}_0$  if throat very near brane ( $\mu_0 < \bar{\mu}_0$ )
- Find analytically:  $M = \frac{(D-2)\omega_{D-2}}{4\pi G_D} \sum_n q_n$  ,  $A_{throat}$  ,  $L_D$  ,  $L_5 = c$
- Constant area  $A = A_{AH} = (A_{throat} \text{ or } A_{outer})$ :  $c(\mu_0) = \left[ \frac{A}{f(\mu_0)} \right]^{1/(D-2)}$

# Static BH, Stability, Binding Energy

Mass  $M$  (constant  $A$ )



$L \rightarrow 0$ :  
 $M \rightarrow M_0$

large  $L$ :

$$M \simeq M_A - \frac{2}{3\pi} \frac{G_5 (M_A)^2}{(2L)^2}$$

rest energy:

$$M_A = \frac{3\pi}{8G_5} \left( \frac{A}{2\pi^2} \right)^{2/3}$$

- Each point: an initially static BH (**on brane**, or **off brane**).
- **The BH on the brane at  $L \rightarrow 0$ :**
- Variational Principle: this mass extremum is a **static** BH.
- It's **stable** against translations (mass is a local minimum).
- **High** binding energy: for brane with  $\mathbb{Z}_n$  orbifold symmetry,

$$E_B = M_A - M_0 = \left[ n^{1/(D-2)} - 1 \right] M_0$$

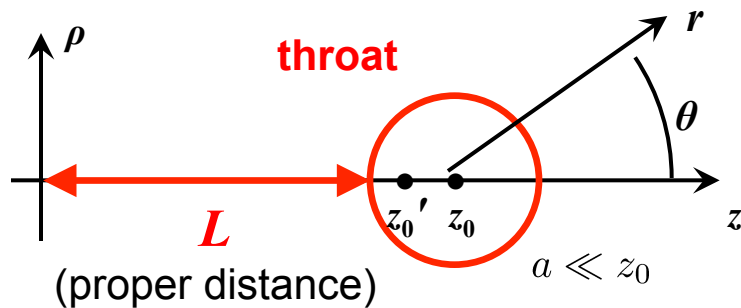
# Black Holes (far from brane)

- Farther from brane, the approximation  $\ell \rightarrow \infty$  breaks down.
- Solve the nonlinear problem with a perturbation series:

$$\psi = \psi_0 + (\psi_0)^2 \phi_1 + (\psi_0)^2 \sum_{i \geq 2} \phi_i \quad , \quad \psi_0 = \frac{\ell}{\ell + z}$$

- All main results from:  $\phi_1$       Perturbations:  $\phi_i$  ( $i \geq 2$ )

- For  $\rho \gg \ell, z, z_0$ :  $\phi_1(\rho, z) \simeq \frac{2G_5 m}{\ell \rho}$  (determines mass  $m$ )



$$L = \int_0^{z_0 - a} dz \psi \simeq -\ell \ln \psi_0(z_0)$$

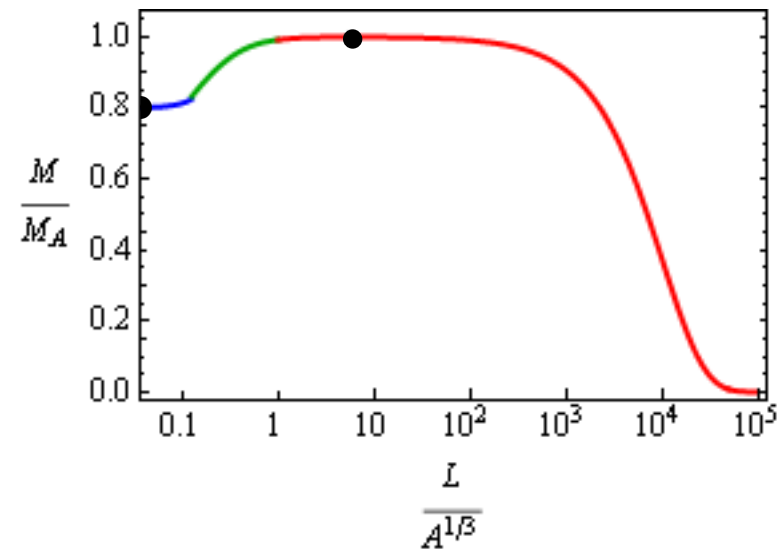
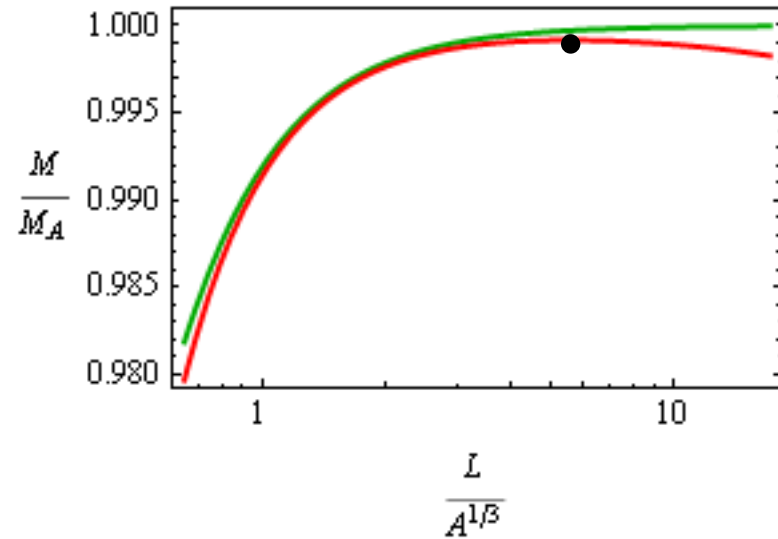
- Mass in terms of area  $A$ :  $M \simeq \psi_0(z_0) M_A - \frac{2}{3\pi} \frac{G_5 (M_A)^2}{(2z_0)^2 \psi_0(z_0)}$

# Static BH, Stability, Binding Energy

Mass  $M$  (constant  $A$ )

$$A^{1/3} = 10^{-4} \ell$$

Mass  $M$  (constant  $A$ )



- Mass extremum at:  $L_{\text{ext}} \simeq (z_0)_{\text{ext}} \simeq \left( \frac{G_5 M_A \ell}{3\pi} \right)^{1/3}$
- Variational Principle: this extremum is a **static** black hole.
- It is **unstable** to translations (mass is a local maximum).
- **Small** contribution to binding energy from brane repulsion:

$$E_B = M_{\text{ext}} - M_0 \simeq \left[ 2^{1/3} - 1 - \frac{3}{2} \left( \frac{2G_5 M_A}{3\pi \ell^2} \right)^{1/3} \right] M_0$$

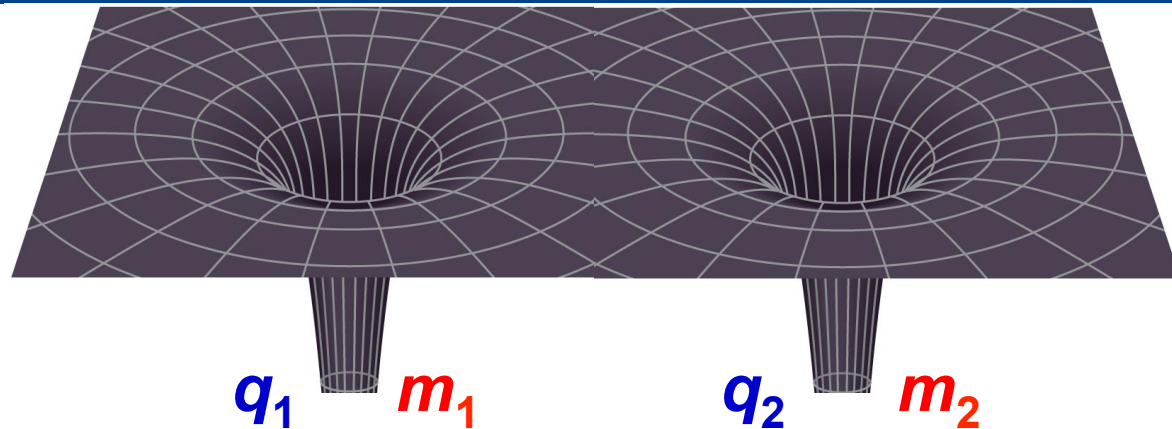
# Black Holes with Electric Charge

**To appear on arXiv**

Scott Fraser, Shaker Funkhouser \*

\* First prize, 2015 CSU Student Research Competition

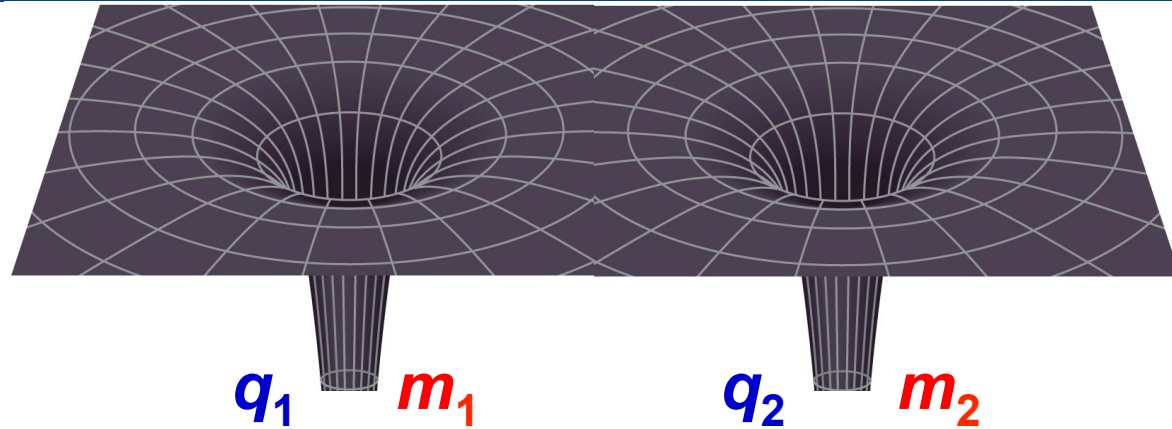
# Introduction: A Set of Static BHs



- Intuition: the black holes are static due to balanced forces.
- **Masses  $m$ : gravitationally attract.**
- **Electric charges  $q$  (same sign): repel.**
  
- For nearly 50 years: this interpretation has prevailed.
- But in general relativity, gravity is due to spacetime geometry, not a force.



# Introduction: A Set of Static BHs



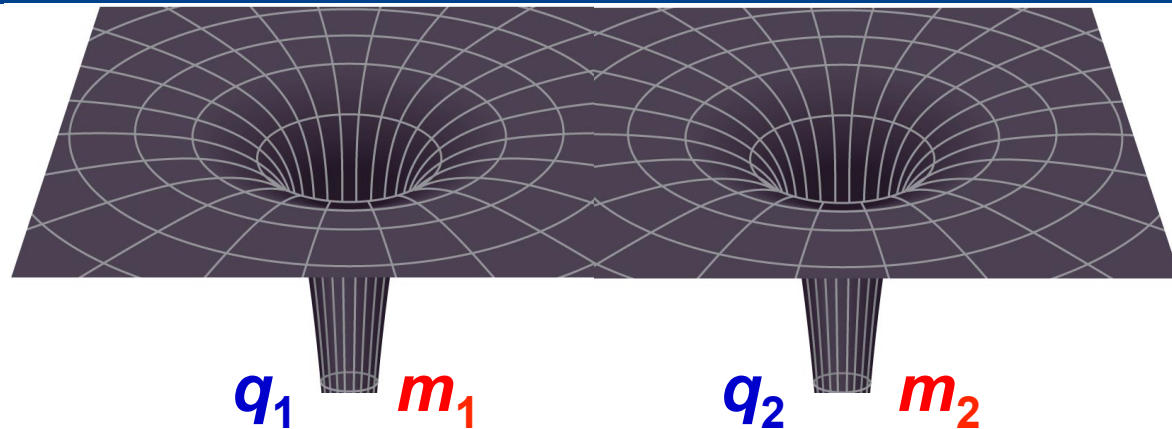
- Known static condition, in units with  $G = 1$  and  $1/(4\pi\epsilon_0) = 1$ :

$$|q_1| = m_1 \quad \text{and} \quad |q_2| = m_2$$

Concisely:  $|q_i| = m_i \quad (i = 1, 2)$  [Majumdar and Papapetrou, 1947]  
[Hartle and Hawking, 1972]

- We rediscover this condition using energy.

# Static BHs Extremize Energy



- Our result:

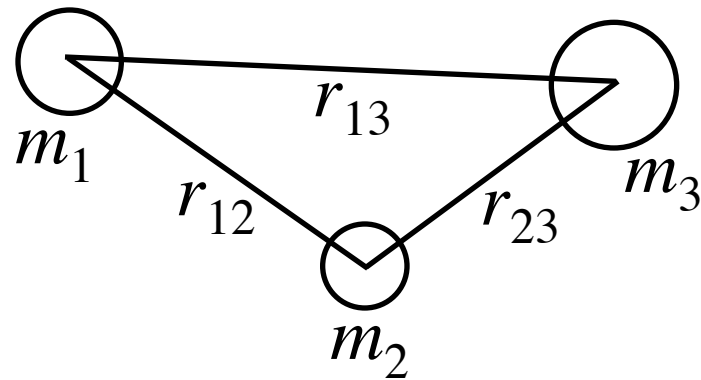
The first energy interpretation of the static condition:

$$|q_i| = m_i \quad (i = 1, 2, \dots, N)$$

- Our method:

Prove that  $|q_i| = m_i$  extremizes the energy of a known geometry with arbitrary same-sign charges:  $|q_i| \leq m_i$

# A Set of $N$ Initially Static BHs

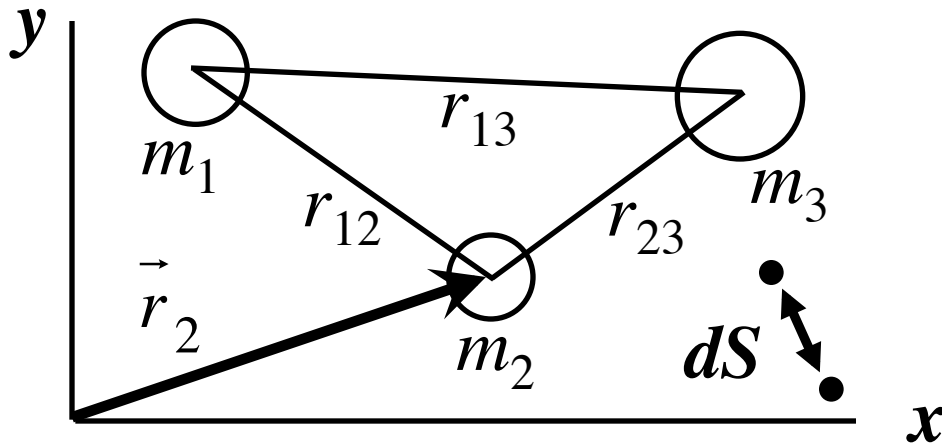


example shown:  $N = 3$

## Procedure:

- Review known geometry for  $N$  initially static BHs.
- Calculate each black hole's area.
- Convenient to use an expansion in large distances  $r_{ij}$ .
- Evaluate total energy and extremize it.
- Show: the extremum yields the static condition  $|\mathbf{q}_i| = m_i$ .

# The known geometry



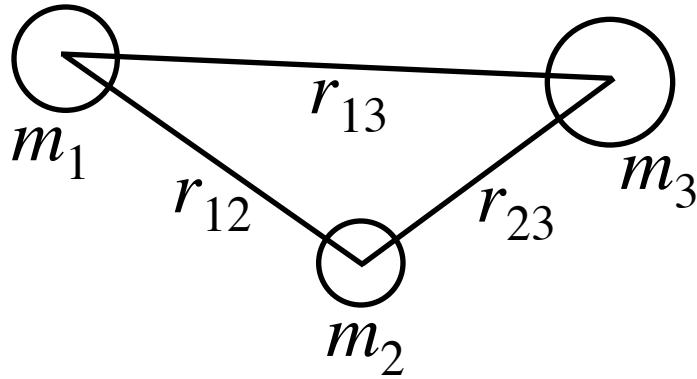
- Known geometry with  $|\mathbf{q}_i| \leq m_i$ :  $dS^2 = f^2 (dx^2 + dy^2 + dz^2)$

[Brill and Lindquist, 1963]

$$f = \left( 1 + \sum_{i=1}^N \frac{\alpha_i}{|\vec{r} - \vec{r}_i|} \right) \left( 1 + \sum_{i=1}^N \frac{\beta_i}{|\vec{r} - \vec{r}_i|} \right)$$

- Constants  $\alpha_i$  and  $\beta_i$  are related to  $m_i$ ,  $\mathbf{q}_i$ , total energy  $E$ .

# The known geometry



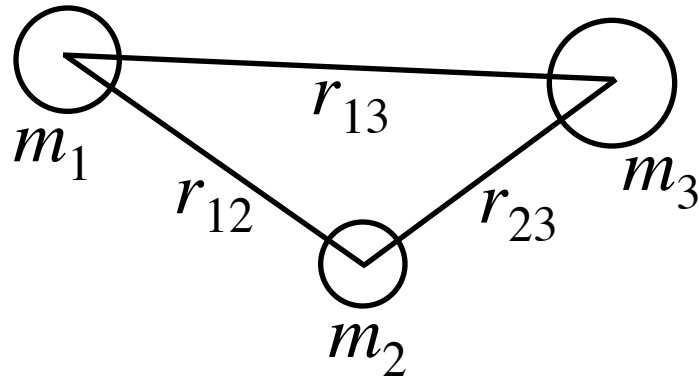
- Constants  $\alpha_i$  and  $\beta_i$  are related to  $m_i$ ,  $q_i$ , total energy  $E$ .

$$m_i = \alpha_i + \beta_i + \sum_{j \neq i} \frac{(\alpha_i \beta_j + \alpha_j \beta_i)}{r_{ij}}$$

$$q_i = \beta_i - \alpha_i + \sum_{j \neq i} \frac{(\beta_i \alpha_j - \beta_j \alpha_i)}{r_{ij}}$$

$$E = \sum_{i=1}^N (\alpha_i + \beta_i)$$

# The known geometry



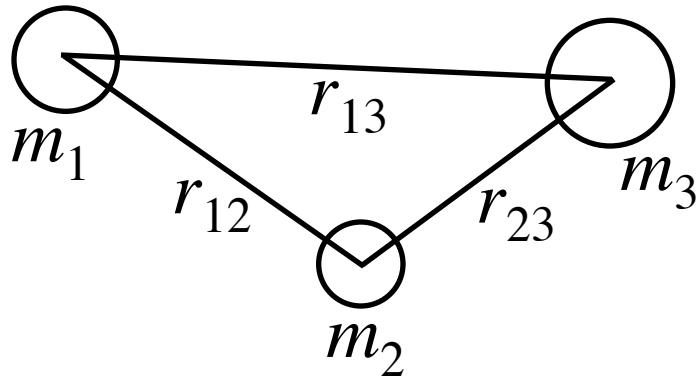
- Constants  $\alpha_i$  and  $\beta_i$  are related to  $m_i$ ,  $q_i$ , total energy  $E$ .

$$\alpha_i = \frac{(m_i - q_i)}{2} \left[ 1 - \frac{1}{2} \sum_{j \neq i} \frac{(m_j + q_j)}{r_{ij}} \right] + \left( \text{terms with } \frac{1}{(r_{ij})^n}, n \geq 2 \right)$$

$$\beta_i = \frac{(m_i + q_i)}{2} \left[ 1 - \frac{1}{2} \sum_{j \neq i} \frac{(m_j - q_j)}{r_{ij}} \right] + \left( \text{terms with } \frac{1}{(r_{ij})^n}, n \geq 2 \right)$$

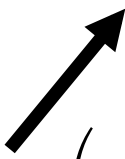
$$E = \sum_{i=1}^N (\alpha_i + \beta_i)$$

# Black hole area



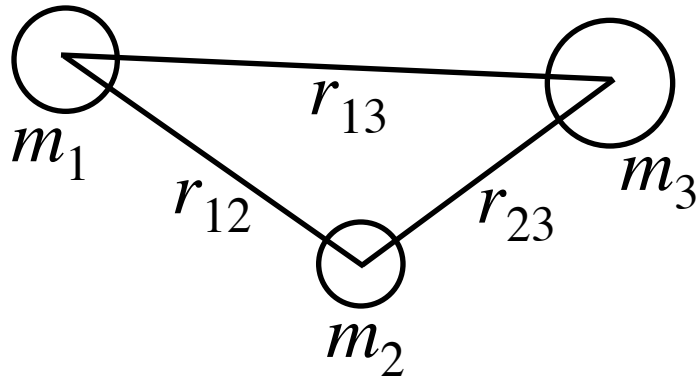
- Calculate the area  $A_i$  of black hole  $i$ :

$$A_i = 4\pi R_i^2 f^2 \qquad R_i^2 = \alpha_i \beta_i \left(1 + \sum_{j \neq i} \frac{\alpha_j}{r_{ij}}\right)^{-1} \left(1 + \sum_{j \neq i} \frac{\beta_j}{r_{ij}}\right)^{-1}$$



$$f = \left(1 + \frac{\alpha_i}{R_i} + \sum_{j \neq i} \frac{\alpha_j}{r_{ij}}\right) \left(1 + \frac{\beta_i}{R_i} + \sum_{j \neq i} \frac{\beta_j}{r_{ij}}\right) + \left(\text{terms with } \frac{1}{(r_{ij})^n}, n \geq 2\right)$$

# Black hole area



- Evaluate area  $A_i$  in terms of mass and charge:

$$\sqrt{\frac{A_i}{4\pi}} = m_i + \sqrt{m_i^2 - q_i^2} + \mathbf{0} + \left( \text{terms with } \frac{1}{(r_{ij})^n}, n \geq 2 \right)$$

*(terms involving  $1/r_{ij}$  cancel)*

- Solve for mass:  $m_i = \sqrt{\frac{\pi}{A_i} \left( \frac{A_i}{4\pi} + q_i^2 \right)} + \mathbf{0}$



# Evaluate Energy

- Recall:  $E = \sum_{i=1}^N (\alpha_i + \beta_i)$  and  $m_i = \sqrt{\frac{\pi}{A_i}} \left( \frac{A_i}{4\pi} + q_i^2 \right) + \mathbf{0}$

- Find:

$$E = \sum_{i=1}^N m_i + \sum_{i=1}^N \sum_{j>i}^N \frac{(q_i q_j - m_i m_j)}{r_{ij}} + \left( \text{terms with } \frac{1}{(r_{ij})^n}, n \geq 2 \right)$$

- Extremize:  $E(\mathbf{A}_i, \mathbf{q}_i, \mathbf{r}_{ij})$  while holding  $\mathbf{q}_i$  and  $\mathbf{r}_{ij}$  constant.

- One extremum is  $N$  conditions:  $\frac{\partial E}{\partial A_i} = 0$

- Find:  $\frac{\partial E}{\partial A_i} = 0$  yields the static condition  $|\mathbf{q}_i| = \mathbf{m}_i$ .

# Conclusions and Outlook

- **Black holes in RS models:**

- Static BHs extremize energy (variational principle).
- Small BHs on the brane are strongly bound:  
an important result for LHC experiments.

- **Charged black holes:**

- By extremizing energy, we found a new interpretation of the long known static condition,  $|\mathbf{q}_i| = m_i$ .
- Proof at higher orders is in progress.